

# *Metallic films: conductivity properties vs. work function. New opportunities for accelerators and other technological applications*



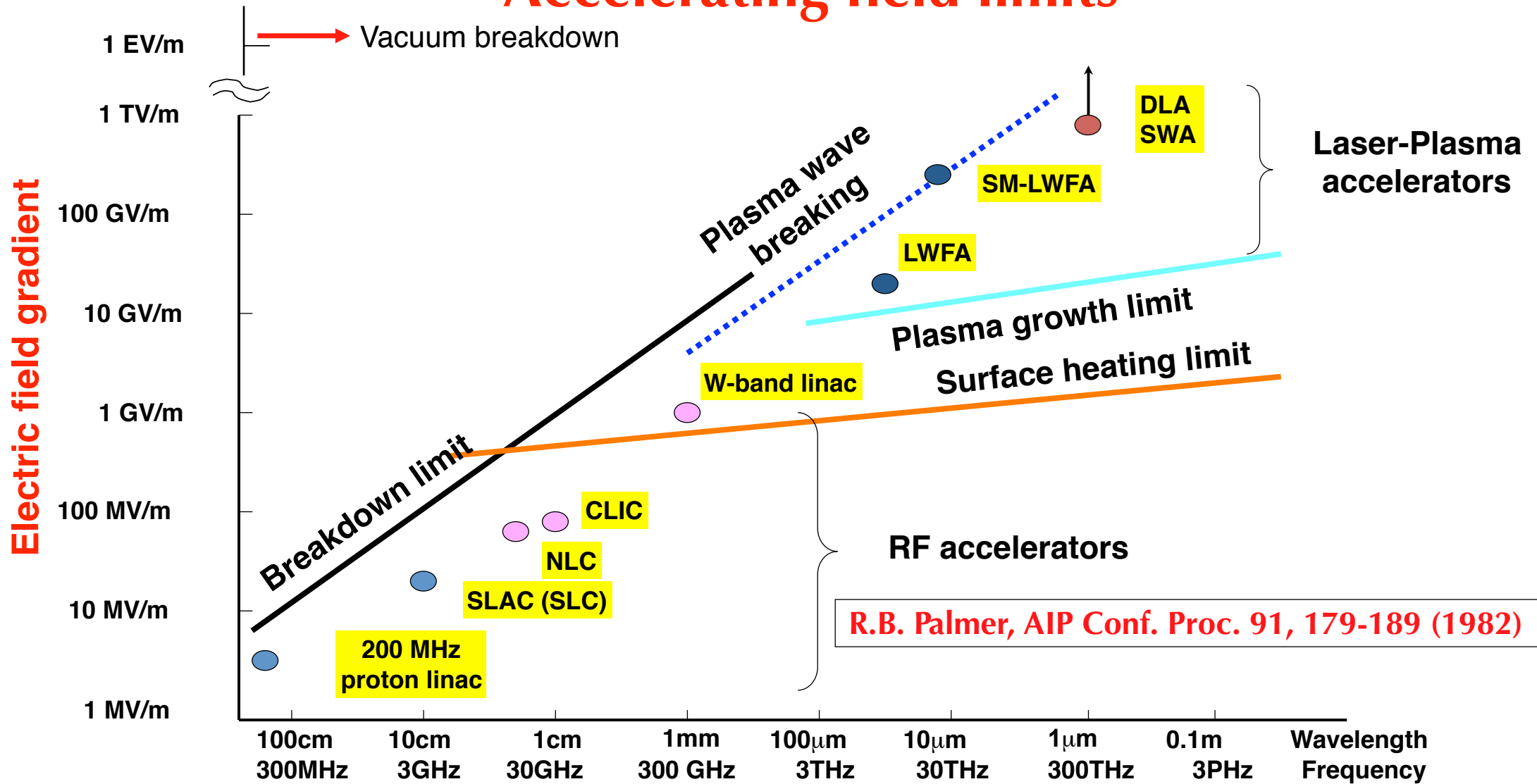
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*Frascati, October 3, 2017*

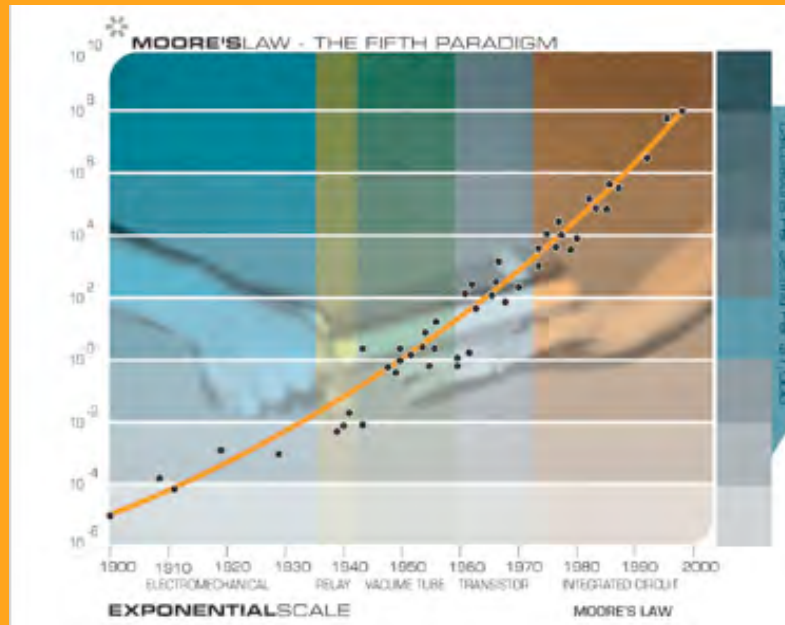
# *Layout*

- Introduction: materials, breakdown rates, work function, etc.
- Mesoscale science, metals, alloys and TM oxides
- The case of Mo and MoO<sub>3</sub> coatings
- Perspectives & conclusions

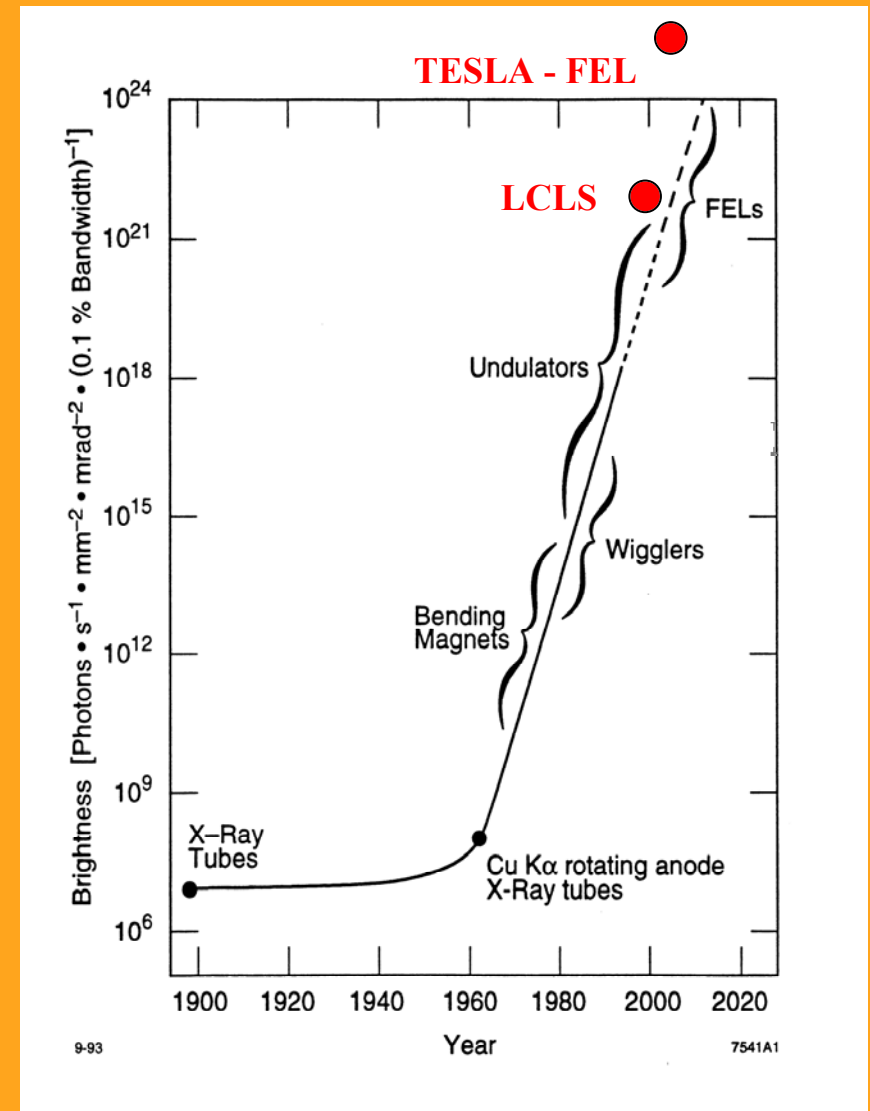
# Accelerating field limits



# EXPONENTIAL GROWTH



## MOORE LAW

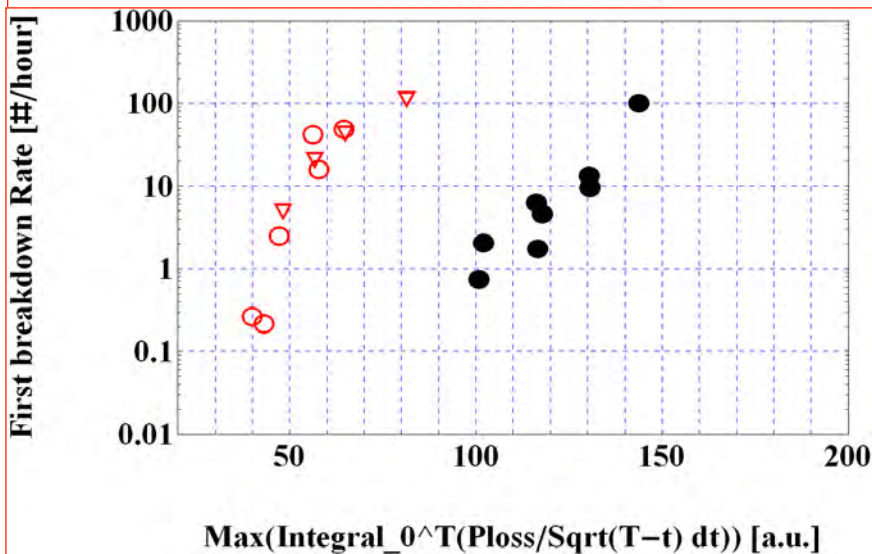
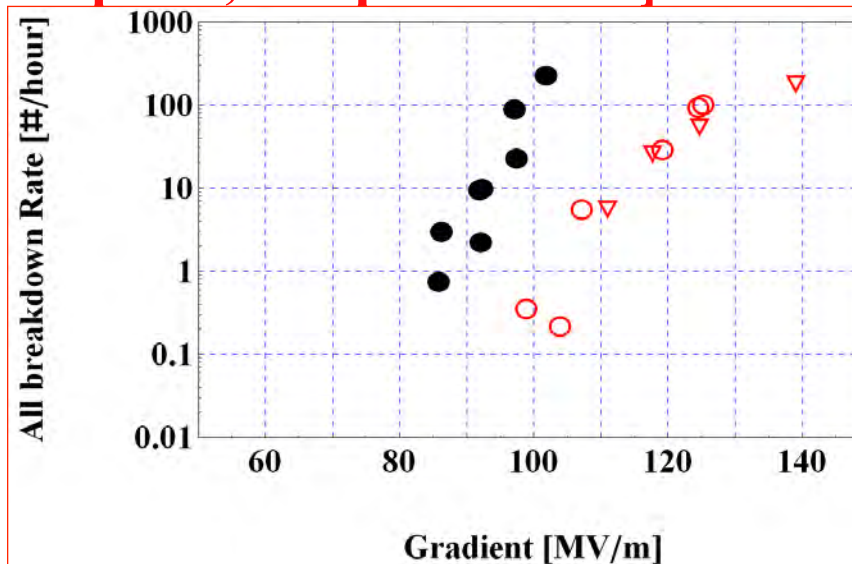


Frascati, October 3, 2017

# Comparison of peak pulse heating for two Cu and one Mo structure [shaped pulse, flat part 150 ns]



Spataro et al. NIM-A-657 (2011) 114-121



Frascati, October 3, 2017



# On the materials basis of modern society

T. E. Graedel<sup>1</sup>, E. M. Harper, N. T. Nassar, and Barbara K. Reck

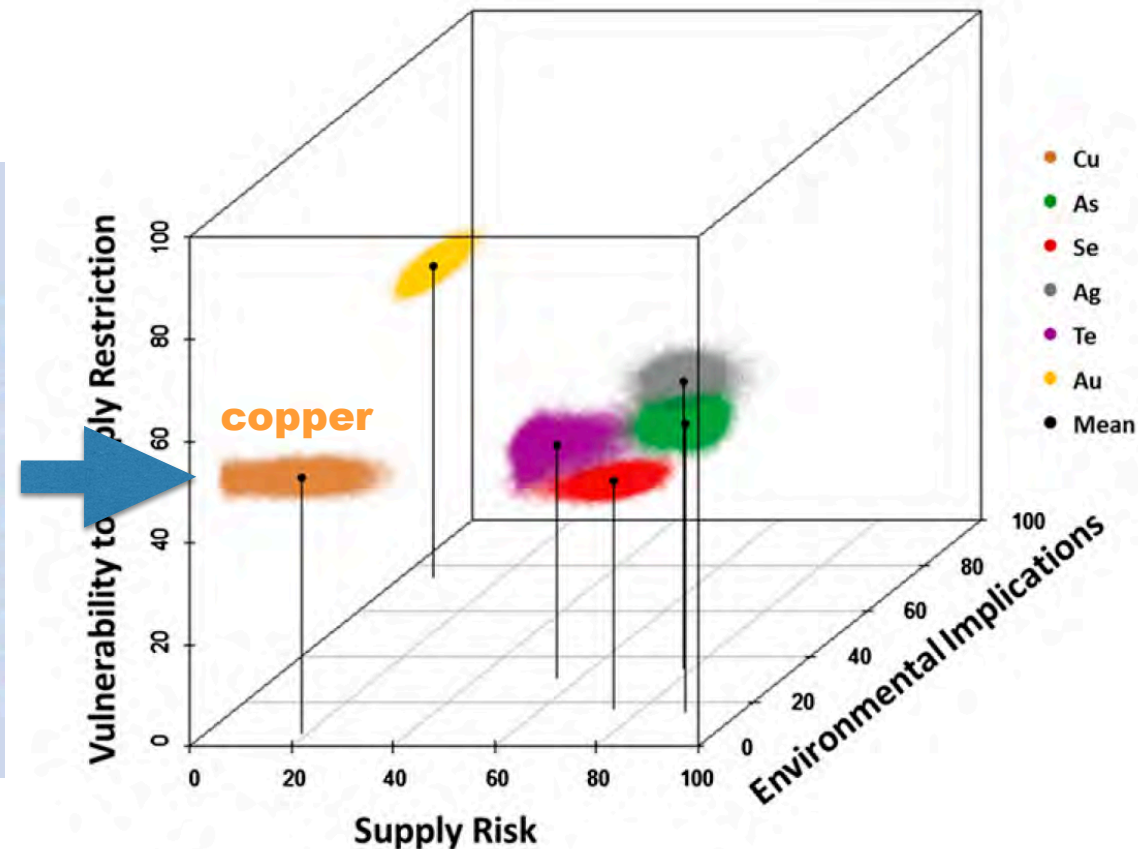
School of Forestry and Environmental Studies, Center for Industrial Ecology, Yale University, New Haven, CT 06511

Edited by William C. Clark, Harvard University, Cambridge, MA, and approved October 11, 2013 (received for review July 29, 2013)

[www.pnas.org/cgi/doi/10.1073/pnas.1312752110](http://www.pnas.org/cgi/doi/10.1073/pnas.1312752110)

## Significance

Modern life is enabled by the use of materials in its technologies. Over time, these technologies have used a larger and more diverse array of materials. Elemental life cycle analyses yield an understanding of these materials, and a definite concern that arises is that of possible scarcity of some of the elements as their use increases. We studied substitution potential for 62 different metals in their major uses. For a dozen different metals, the potential substitutes for their major uses are either inadequate or appear not to exist at all. Further, for not 1 of the 62 metals are exemplary substitutes available for all major uses.



## Observation of Field-Emission Dependence on Stored Energy

Jiahang Shao,<sup>1,2,\*</sup> Sergey P. Antipov,<sup>2,3</sup> Sergey V. Baryshev,<sup>2,3</sup> Huaibi Chen,<sup>1</sup> Manoel Conde,<sup>2</sup> Darrell S. Doran,<sup>2</sup> Wei Gai,<sup>2</sup> Chunguang Jing,<sup>2,3</sup> Wanming Liu,<sup>2</sup> John Power,<sup>2</sup> Jiaqi Qiu,<sup>2,3</sup> Jiaru Shi,<sup>1</sup> Dan Wang,<sup>1,2</sup> Faya Wang,<sup>4,†</sup> Charles E. Whiteford,<sup>2</sup> Eric Wisniewski,<sup>2</sup> and Liling Xiao<sup>4</sup>

<sup>1</sup>*Department of Engineering Physics, Tsinghua University, Beijing 100084, People's Republic of China*

<sup>2</sup>*Argonne National Laboratory, Lemont, Illinois 60439, USA*

<sup>3</sup>*Euclid Techlabs LLC, Solon, Ohio 44139, USA*

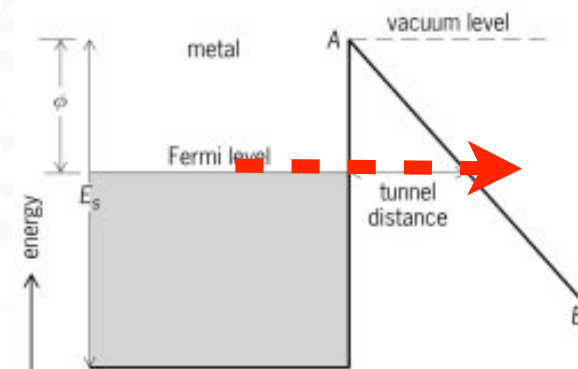
<sup>4</sup>*SLAC National Accelerator Laboratory, Menlo Park, California 94025, USA*

(Received 24 August 2015; published 23 December 2015)

Field emission from a solid metal surface has been continuously studied for a century over macroscopic to atomic scales. It is general knowledge that, other than the surface properties, the emitted current is governed solely by the applied electric field. A pin cathode has been used to study the dependence of field emission on stored energy in an  $L$ -band rf gun. The stored energy was changed by adjusting the axial position (distance between the cathode base and the gun back surface) of the cathode while the applied electric field on the cathode tip is kept constant. A very strong correlation of the field-emission current with the stored energy has been observed. While eliminating all possible interfering sources, an enhancement of the current by a factor of 5 was obtained as the stored energy was increased by a factor of 3. It implies that under certain circumstances a localized field emission may be significantly altered by the global parameters in a system.

DOI: 10.1103/PhysRevLett.115.264802

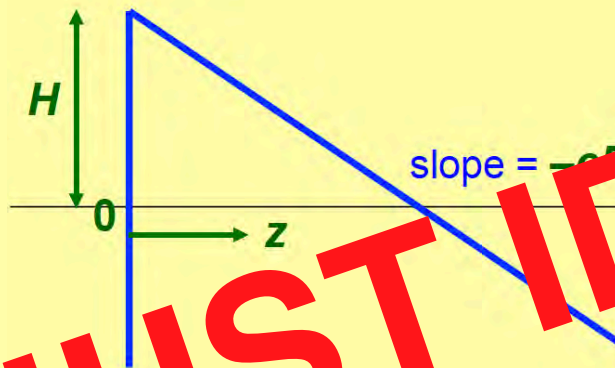
PACS numbers: 29.20.Ej, 52.80.Pi, 79.70.+q





Two well-known special **barrier forms** exist:

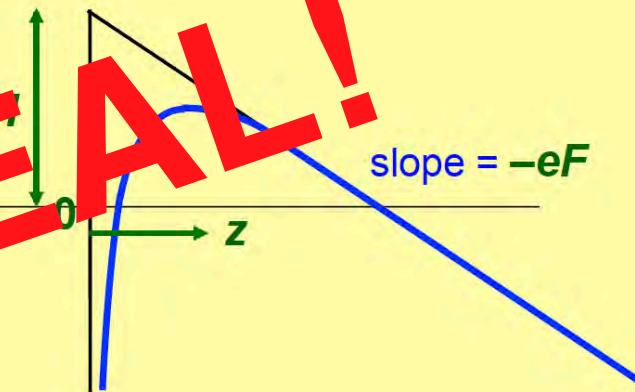
Exact triangular (ET) barrier



$$M(z) = H - eFz$$

used by **Fowler & Nordheim**

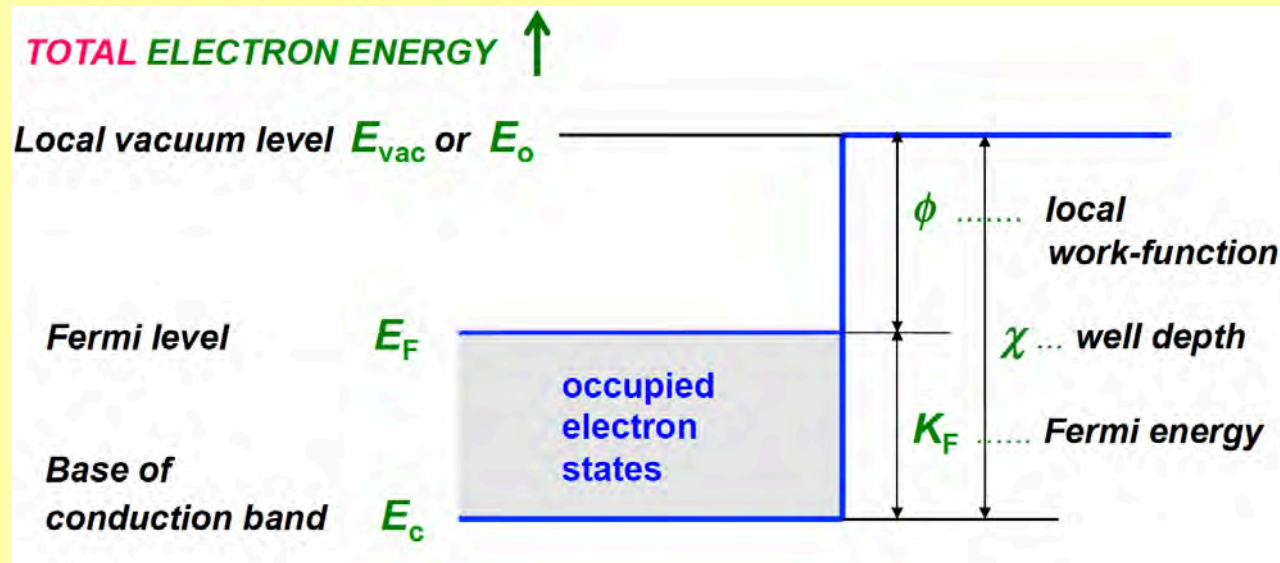
Schottky-Nordheim (SN) barrier



$$M(z) = H - eFz - 1/16\pi\epsilon_0 z$$

used by **Murphy & Good**





The well depth and the local work-function have TWO components

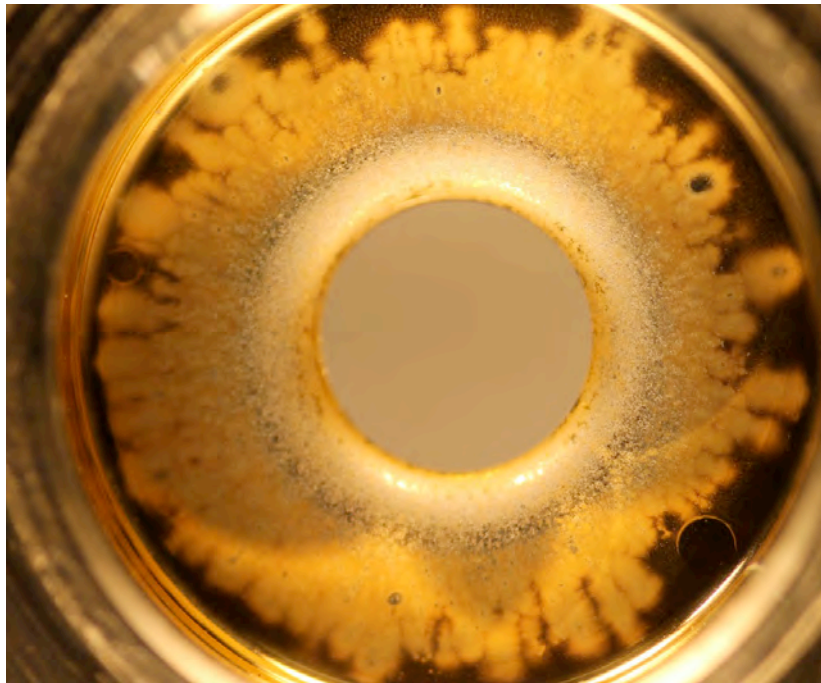
- a bulk (or "chemical") component
- a surface (or "electric dipole") component.



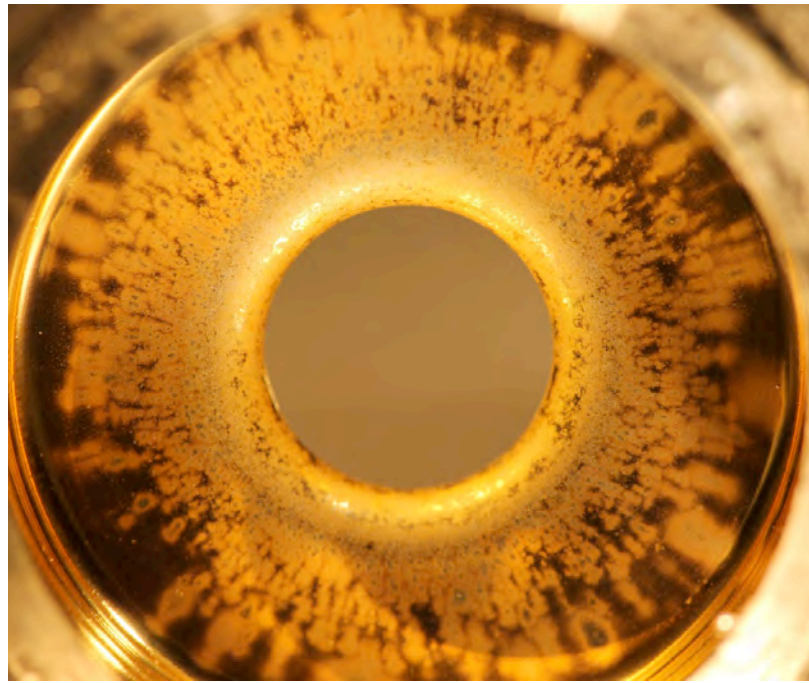
The bulk component explains why different materials have different work functions; the surface component explains why different faces of a single material have different work-functions.

In summary, a strong correlation between the field-emission current and the stored energy has been observed on pin cathodes. This study has excluded mechanisms that may affect the conclusion, such as multipacting in the cavity, background emission from other surfaces, secondary electron emission from the Faraday cup, the beam loading effect, and the space charge limited emission effect. We conclude that the observation is fundamental and inconsistent with the Fowler-Nordheim equation. This indicates that macroscopic parameters like stored energy are affecting the microscopic emission. The findings suggest a new territory to be explored while developing field emission electron sources and high gradient devices.

# Autopsy of an electroformed Au-Ni structure



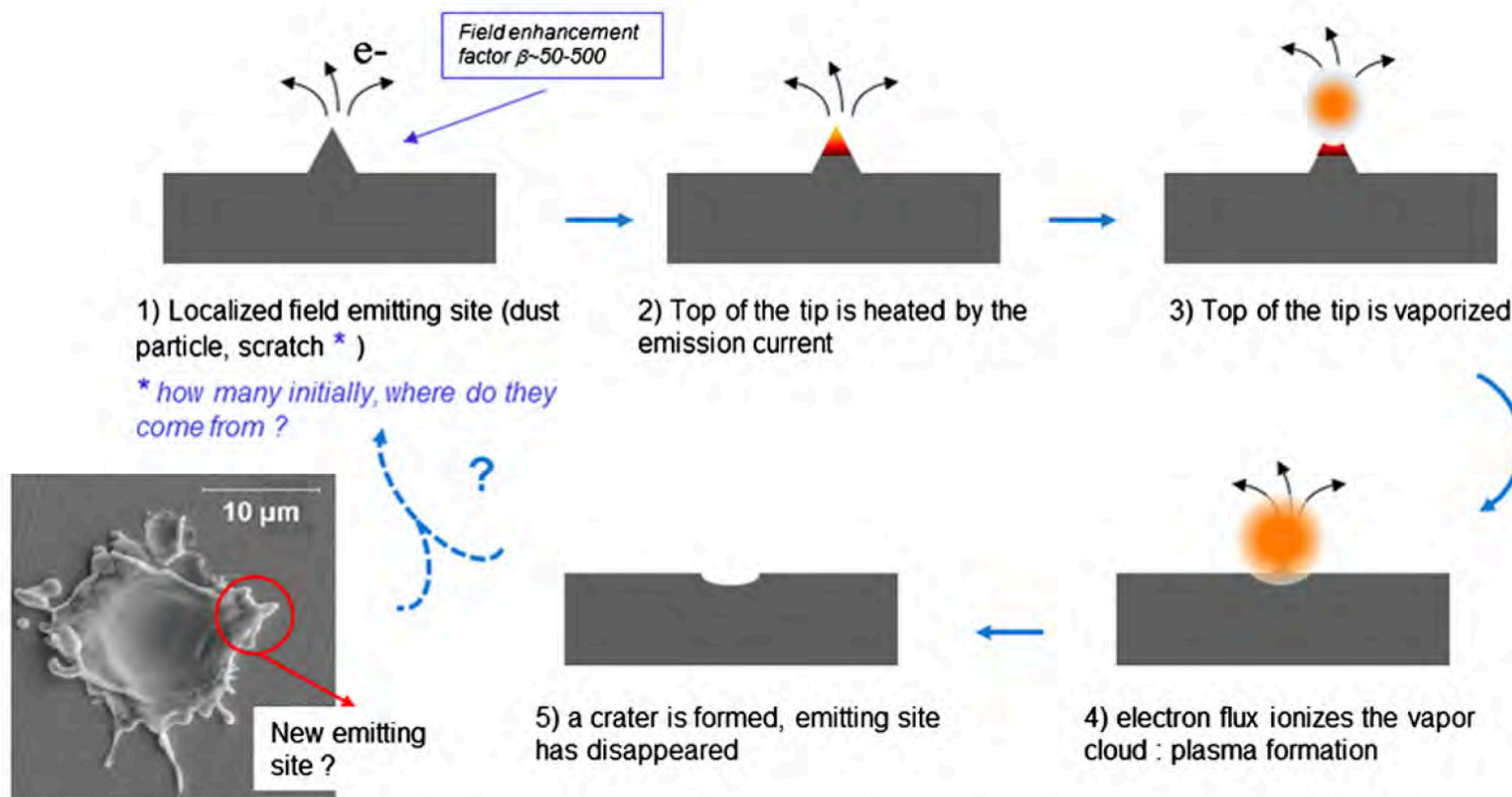
High gradient side of  
*Center cell*



High gradient side of  
*End cell*

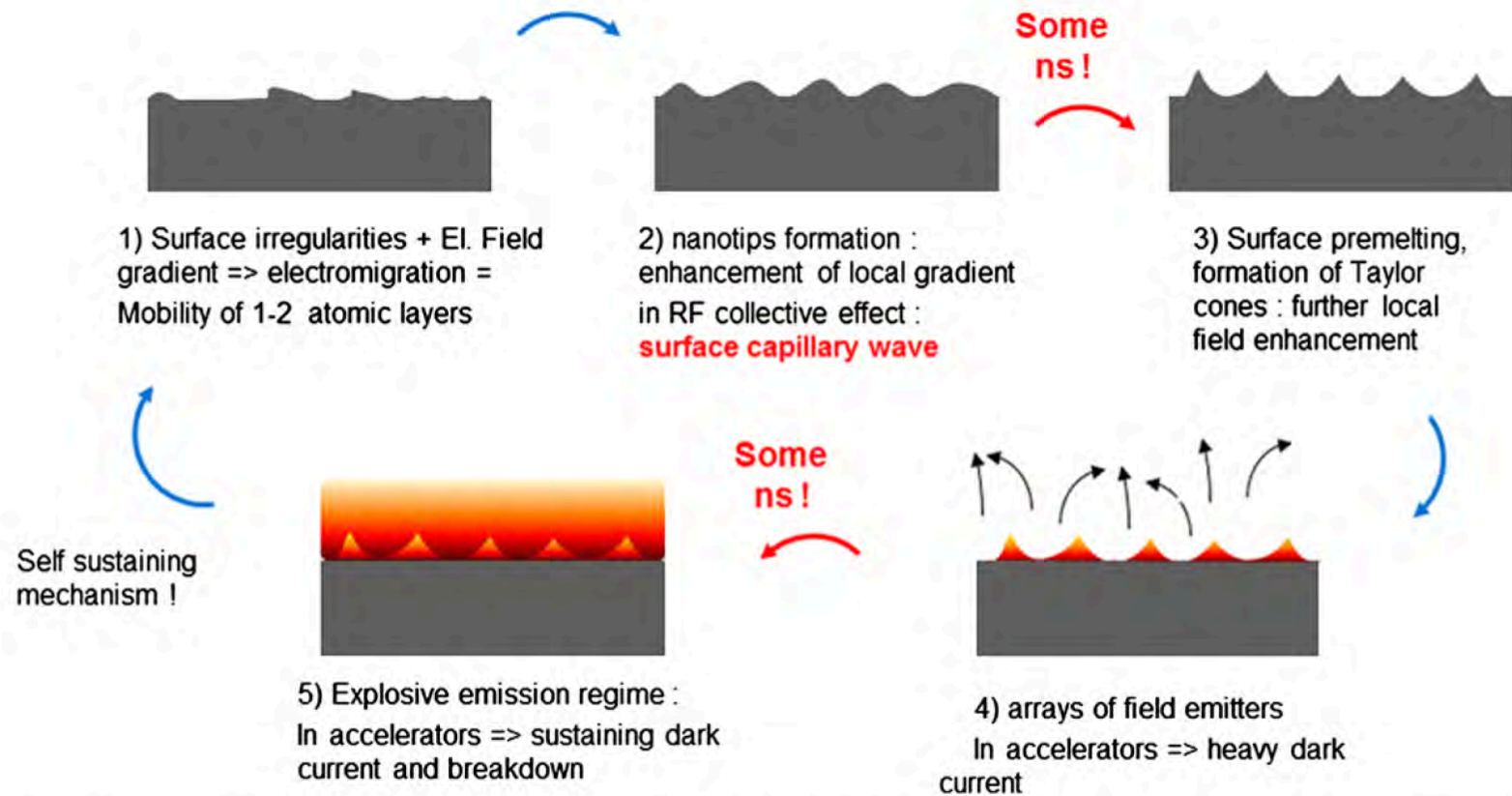
**Au coating is removed from the high-electric field area  
by multiple breakdowns**





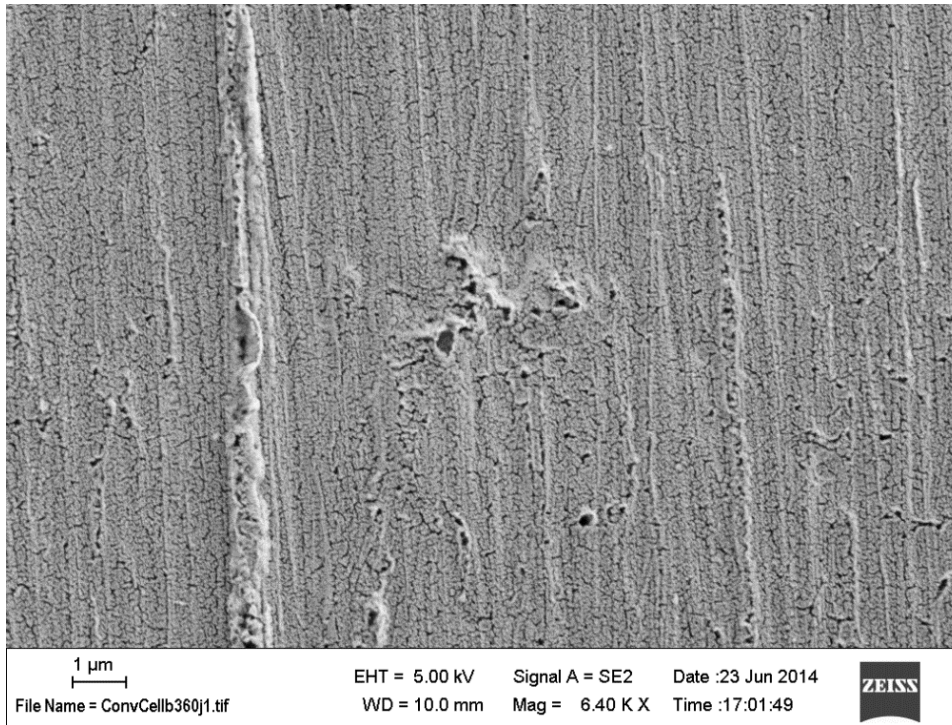
**Fig. 3.** Classical breakdown scenario: BD is initiated by a local defect (inclusions, dust particle, cutting flake, etc.) which provide sufficient field enhancement to induce field emission. The current at the apex of the emitting site becomes sufficient to heat up and vaporize the material, where the emitting current can trigger plasma and generate the explosion of the emitting site. This scenario does not consider the origin of emitting sites.



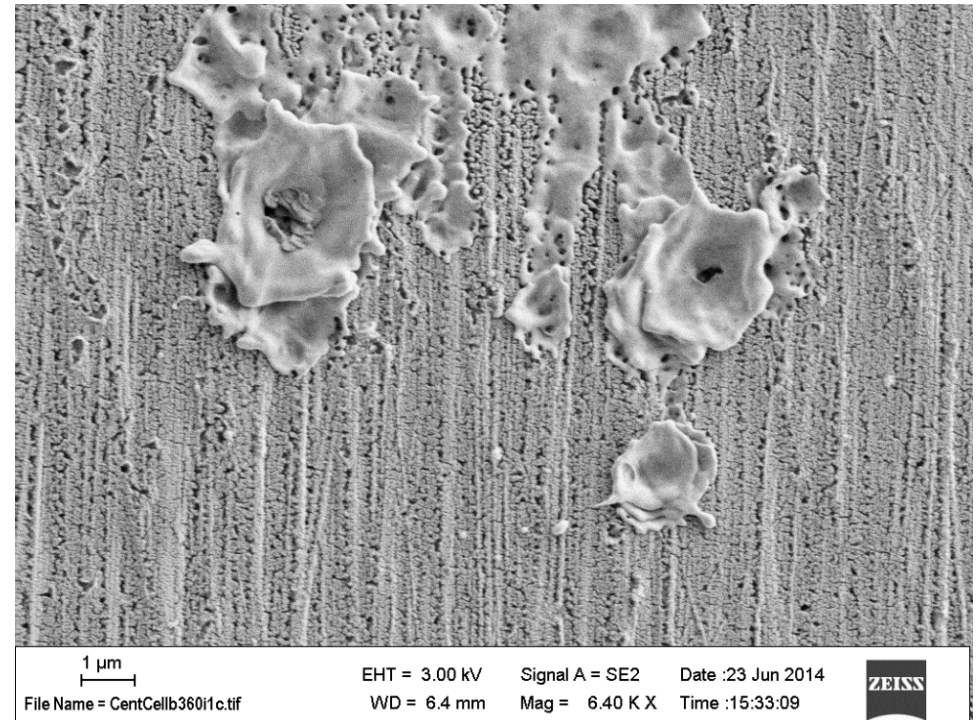


**Fig. 4.** Alternative scenario for breakdown: electromigration tends to enhance the surface irregularities because atoms tend to migrate along field gradient toward the top of the “hill” and form nanotips. In certain conditions, such as the presence of RF field, this effect becomes collective and will help form an array of nanotips. Once this phenomena starts, it can only run away because the electric field gradient become stronger, which in turn enhances electromigration, etc. until the surface becomes liquid like. The electromigration scenario shows that in some conditions, emitters can be formed simply due to the presence of electric field and can trigger breakdown continuously (see text for details).

# Autopsy of an electroformed Au-Ni structure

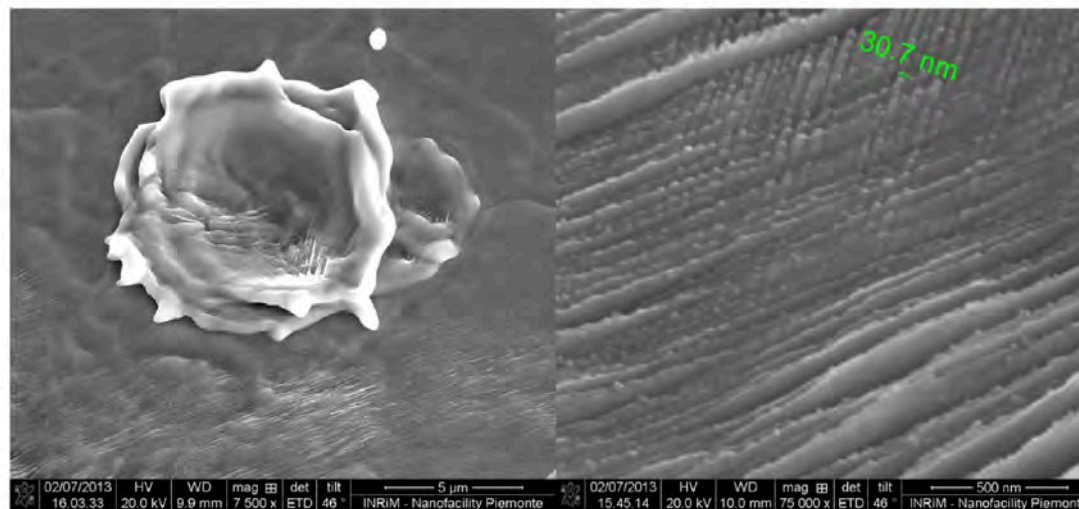


**Coupler cell, near 2b diameter,  
low pulse heating area**

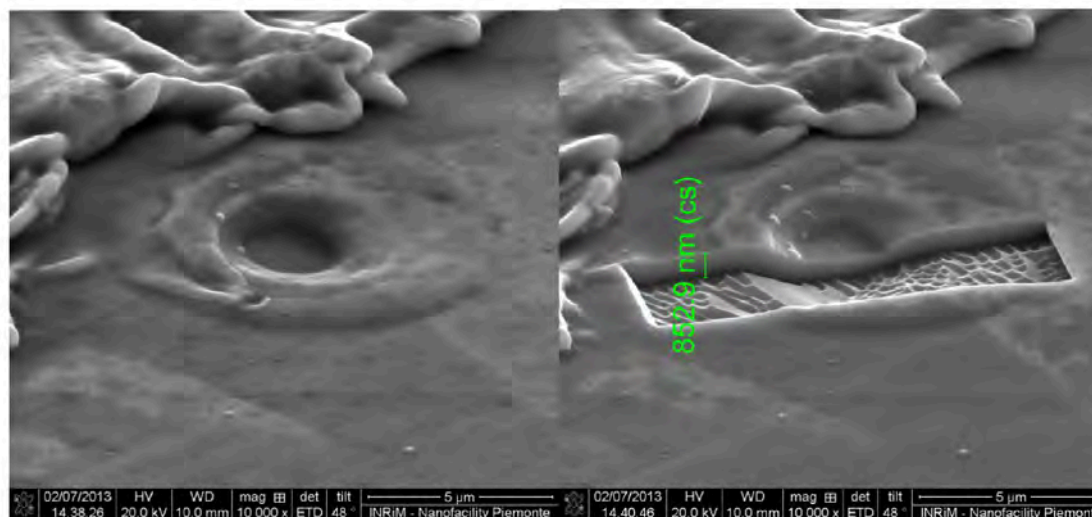


**Center cell,  
high pulse heating area**





**Fig. 4.** Details of the surface of a cavity exposed at SLAC (courtesy of V. Dolgashev). left) an isolated defect consisting in a molten Cu crater showing the texture ascribed to thermal fatigue; right) a magnified view of texture at the nanometer scale induced by the cyclic pulsed heating.



**Fig. 5.** Left: damaged copper surface at the boundary of fully melted regions near the iris of a cell (cavity exposed at SLAC-courtesy of V. Dolgashev). Right: the typical depth of the isolated corona damage is of the order of few micrometers as shown by the FIB cross section.

## Next Generation High Brightness Electron Beams From Ultra-High Field Cryogenic Radiofrequency Photocathode Sources

J.B. Rosenzweig<sup>1</sup>, A. Cahill<sup>1</sup>, V. Dolgashev<sup>2</sup>, C. Emma<sup>1</sup>, A. Fukusawa<sup>1</sup>, R. Li<sup>2</sup>, C. Limborg<sup>2</sup>, J. Maxson<sup>1</sup>, P. Musumeci<sup>1</sup>, A. Nause<sup>1</sup>, R. Pakter<sup>1</sup>, R. Pompili<sup>3</sup>, R. Roussel<sup>1</sup>, B. Spataro<sup>3</sup>, and S. Tantawi<sup>2</sup>

<sup>1</sup>*Department of Physics and Astronomy, University of California, Los Angeles*

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<sup>3</sup>*Istituto Nazionale di Fisica Nucleare, Laboratori Nazionali di Frascati, Frascati, Italy*

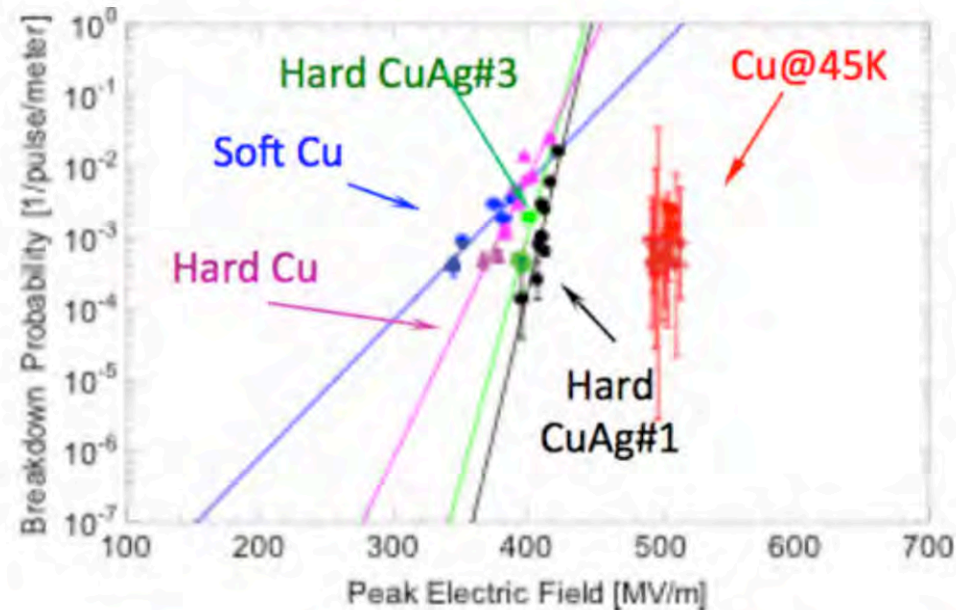


Figure 2. Breakdown probability as a function of peak electric field in single cell X-band accelerating structure tests. The introduction of a harder alloy improves the breakdown as predicted; the effect of operation at 45 deg K is more dramatic, permitting surface electric fields of 500 MV/m.



TABLE II. Fundamental constants for copper.

Thermal conductivity [W/m K]	400
Volumetric heat capacity [MJ/m <sup>3</sup> K]	3.45
Resistivity@300 K [nΩ m]	17
Melting temperature [K]	1358

**The advantages conferred by cryogenic operation of RF cavities for high field operation arise from the hardness of the material, the lowered surface dissipation and concomitant pulsed heating.**

**The improved yield strength of the copper material (new alloys?) is the effect that enables very high fields to be reached, while the changes in dissipation dictate important design features of the RF cavity system in its practical manifestation.**

# Mesoscale science

Mesoscale science, where atomic granularity, quantization of energy, and simplicity of structure and function give way to continuous matter and energy, complex structures, and composite functionalities, is a broad and rich horizon for innovative materials and chemistry.

We assisted to a continuous increase of investigations at ever smaller length and time scales that reveal the atomic, molecular, and nanoscale origins of macroscopic behaviour. Now, we have to begin to reverse the approach, using the knowledge of nanoscale phenomena to understand and control mesoscale architectures to promote the emergence of new behaviours.

Actually the mesoscale is an inherently dynamic regime, where energy and information captured at the nanoscale are processed and transformed at the mesoscale. Experiments have to look at the dynamics, possibly *in situ*, and have to be multimodal being the ordinary “static” characterisation not more sufficient.

*Frascati, October 3, 2017*

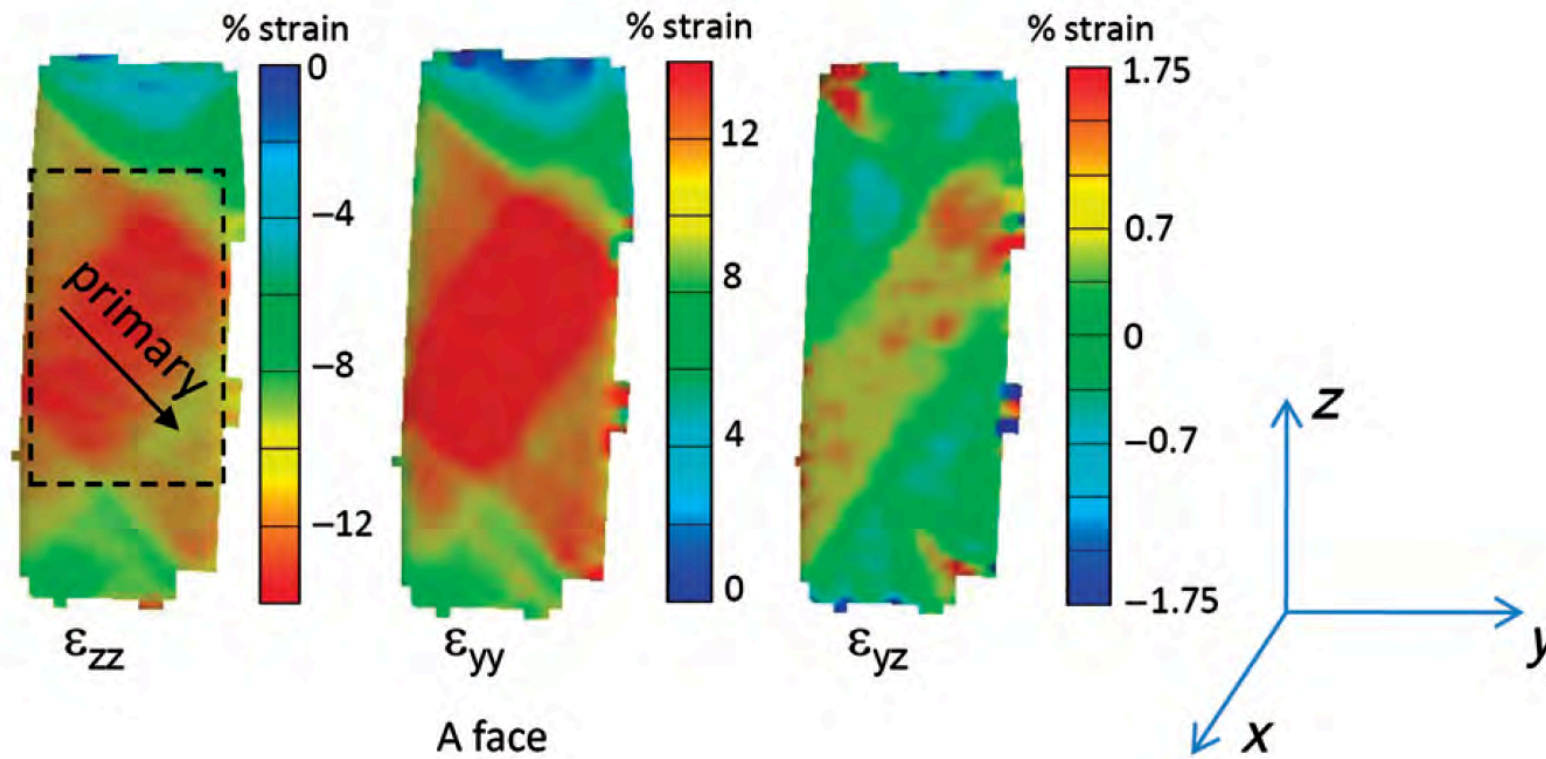
## **Mesoscale strain measurement in deformed crystals: A comparison of X-ray microdiffraction with electron backscatter diffraction**

D.P. Field<sup>a\*</sup>, K.R. Magid<sup>b</sup>, I.N. Mastorakos<sup>a</sup>,  
J.N. Florando<sup>c</sup>, D.H. Lassila<sup>c</sup> and J.W. Morris Jr.<sup>d</sup>

<sup>a</sup>*School of Mechanical and Materials Engineering, Washington State University, Pullman, WA 99164-2920, USA;* <sup>b</sup>*Laboratory for Nanometallurgy, Department of Materials, ETH-Zurich, CH-8093 Zurich, Switzerland;* <sup>c</sup>*Lawrence Livermore National Laboratory, Livermore, CA 94551, USA;* <sup>d</sup>*Department of Materials Science and Engineering, University of California, Berkeley, CA 94720, USA*

*(Received 29 June 2009; final version received 6 October 2009)*

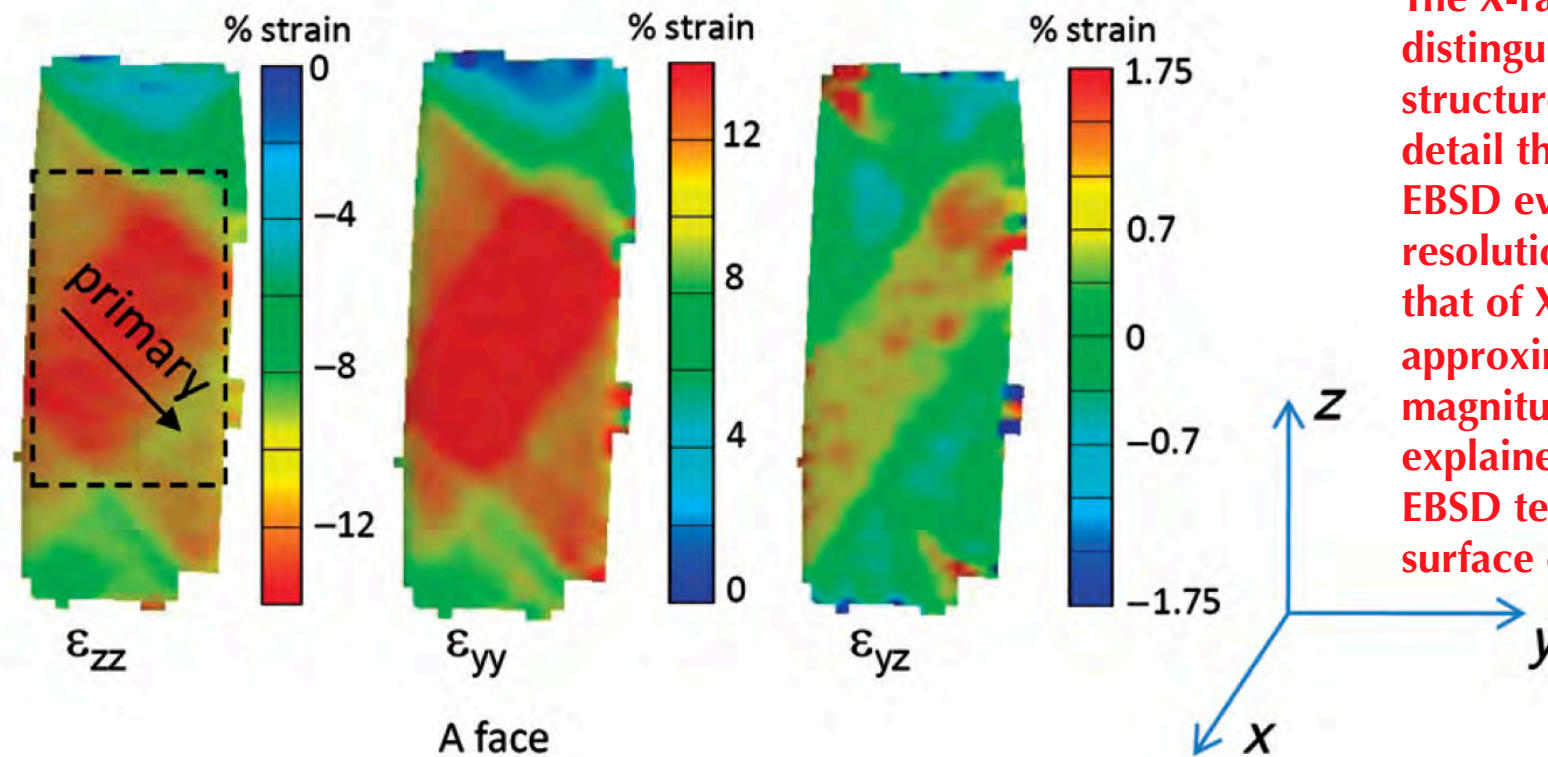
## Mesoscale strain measurement in deformed crystals: A comparison of X-ray microdiffraction with electron backscatter diffraction



October 3, 2017



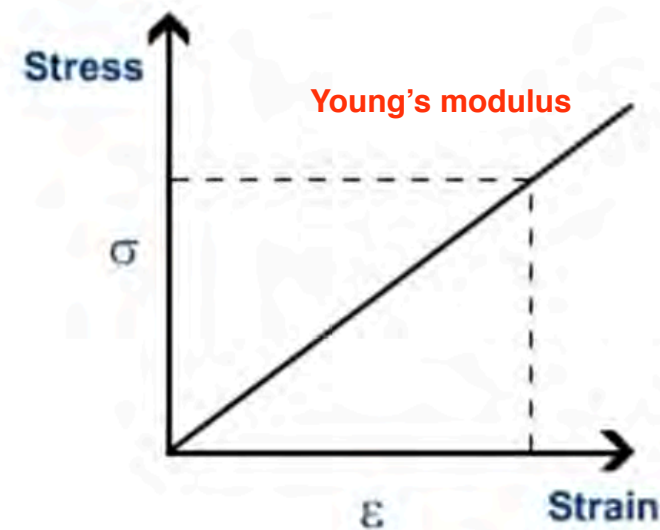
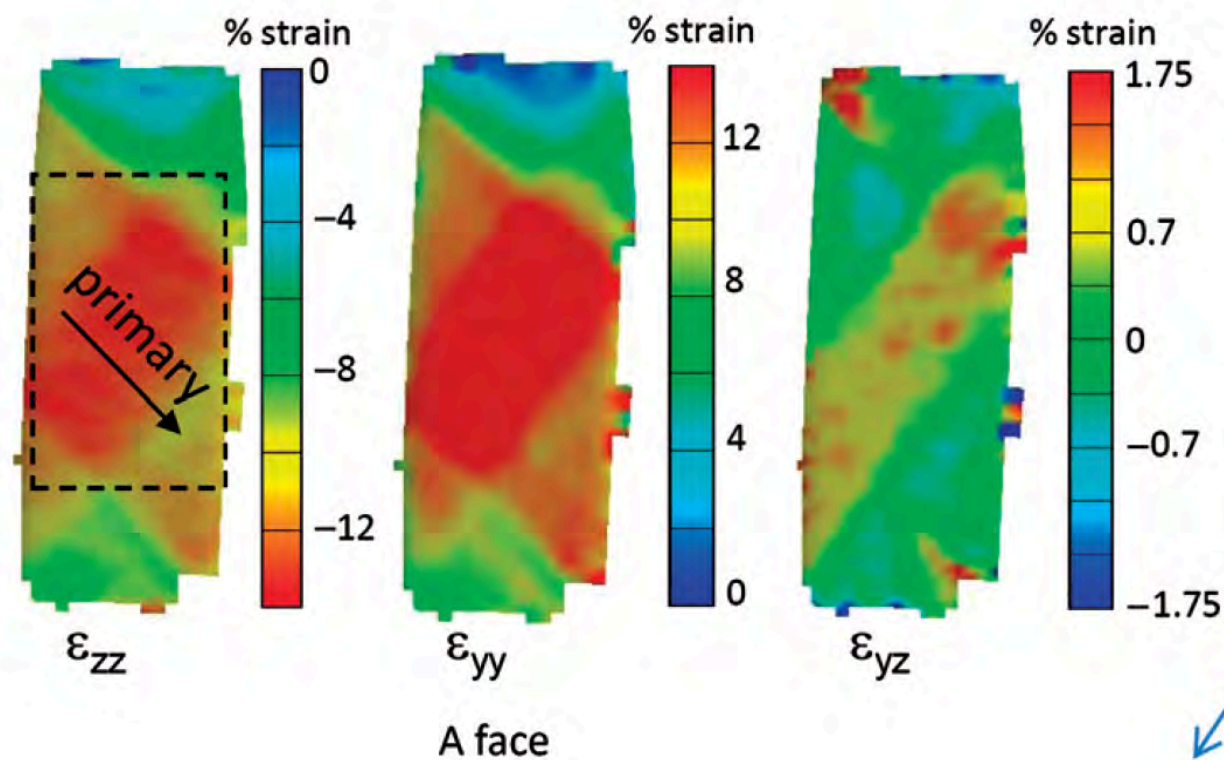
## Mesoscale strain measurement in deformed crystals: A comparison of X-ray microdiffraction with electron backscatter diffraction



The X-ray technique can distinguish features in the structure that are much finer in detail than those observed using EBSD even though the spatial resolution of EBSD is superior to that of X-ray diffraction by approximately two orders of magnitude. The results are explained by the sensitivity of the EBSD technique to the specimen surface condition.

October 3, 2017

## Mesoscale strain measurement in deformed crystals: A comparison of X-ray microdiffraction with electron backscatter diffraction



October 3, 2017

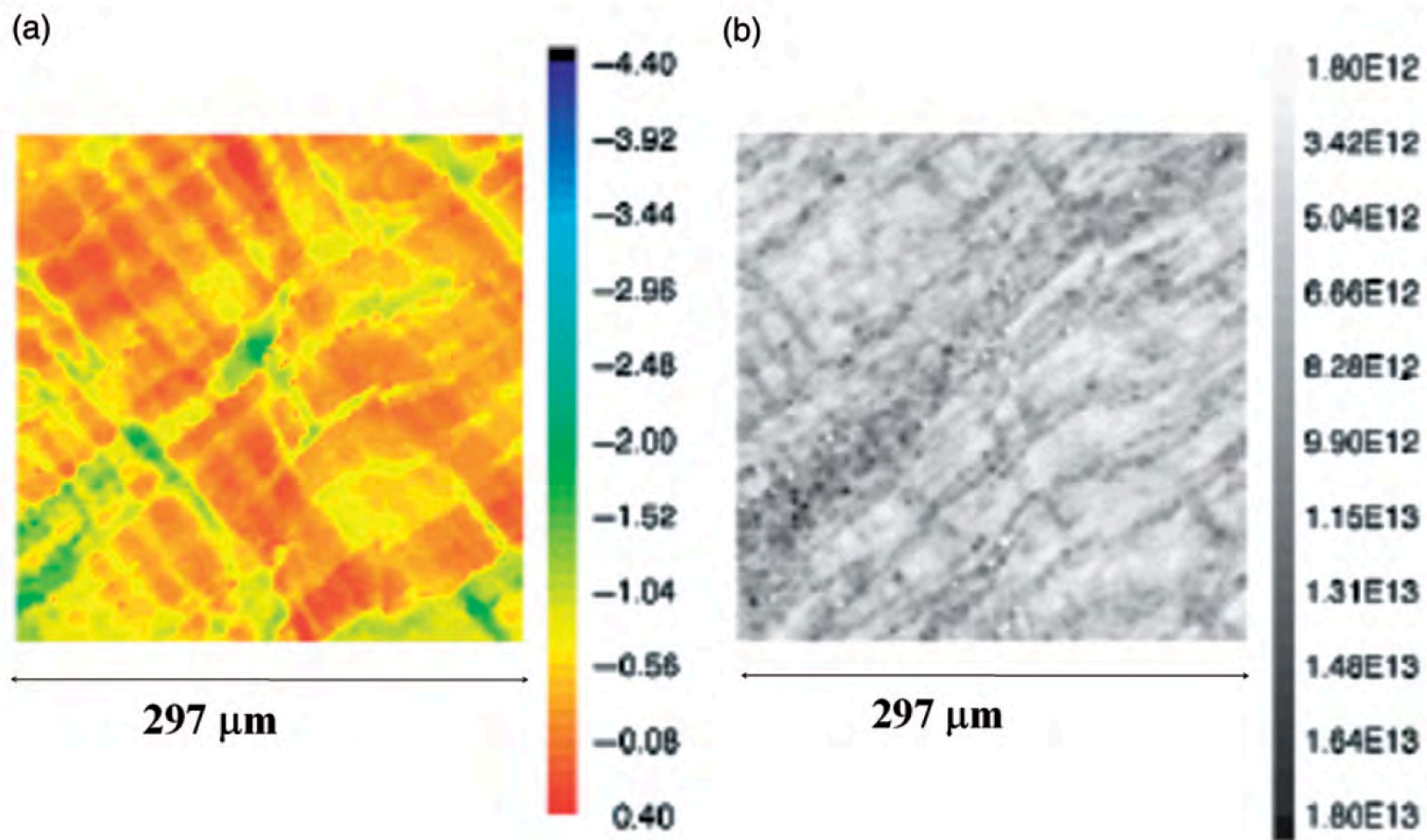
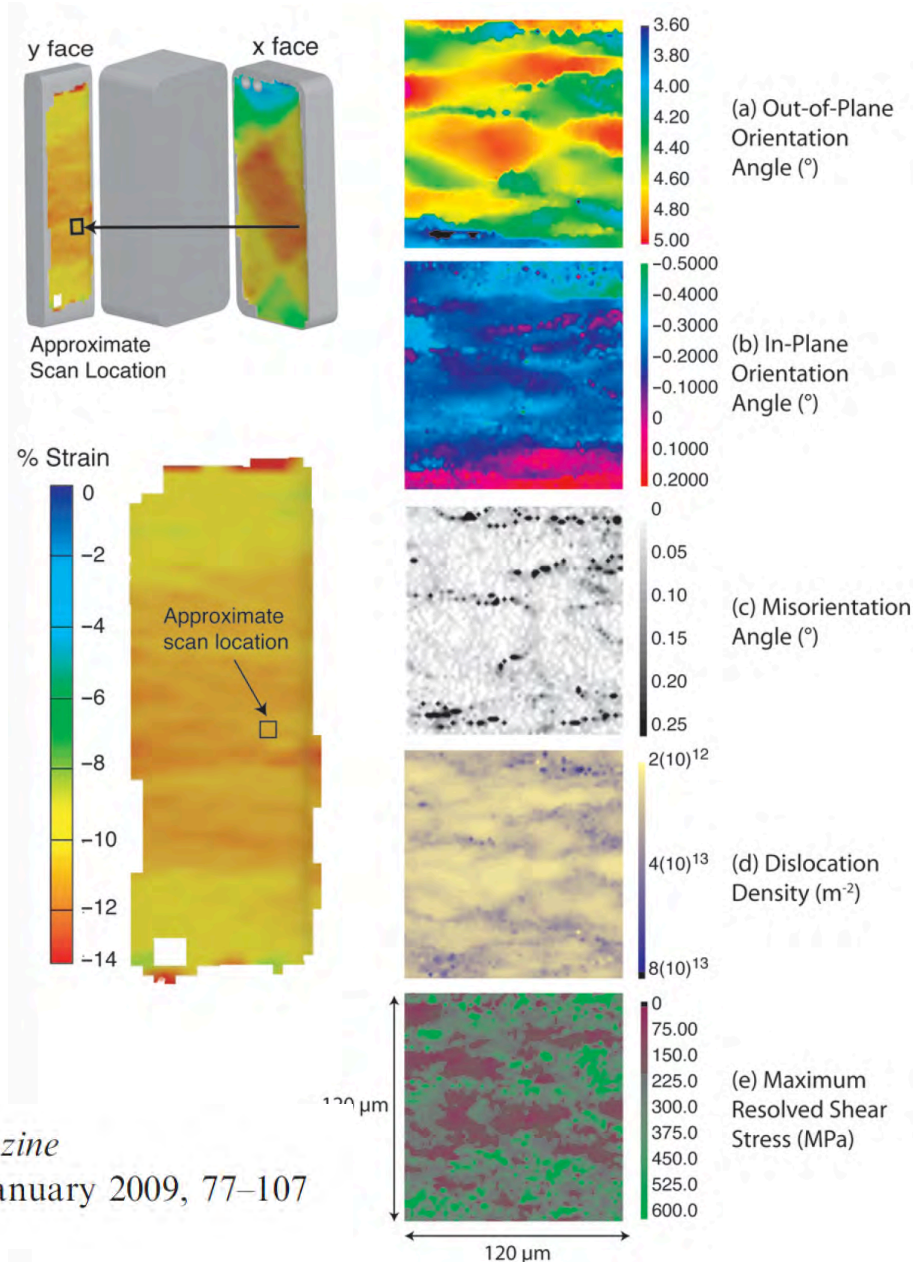


Figure 3. (Colour online). X-ray microdiffraction results showing (a) orientation image and (b) dislocation density map of the deformed structure.





**It is practical and useful to employ ‘multiscale characterization’ with modern tools to clarify the nature and interaction of the heterogeneous deformation patterns and mechanisms that operate at different length scales. In particular, the combination of image correlation techniques to map the macroscopic strain and synchrotron (or EBSD) techniques provides detailed information on mesoscopic defect patterns.**



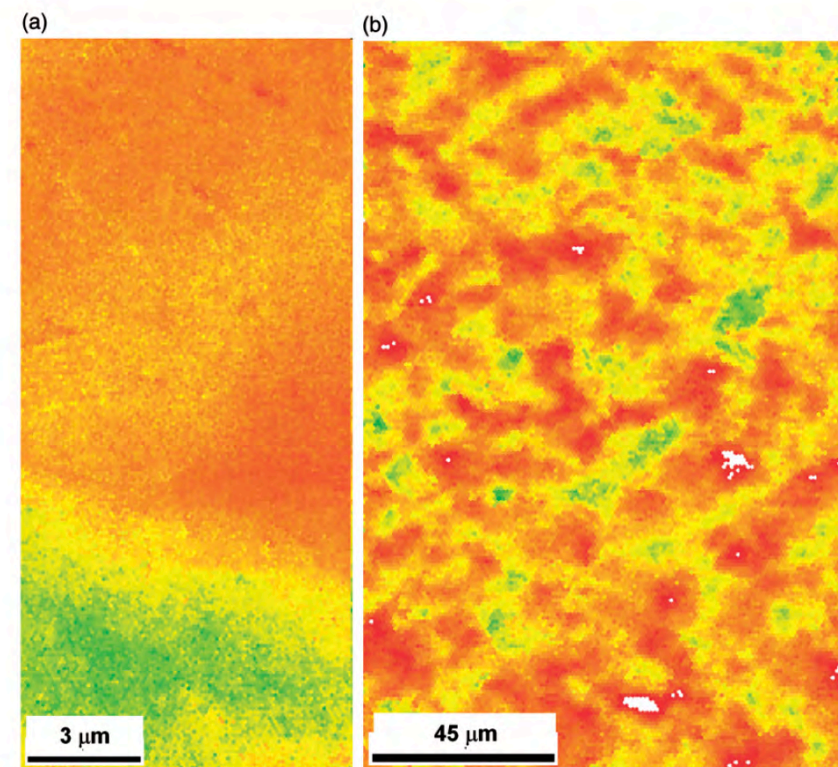
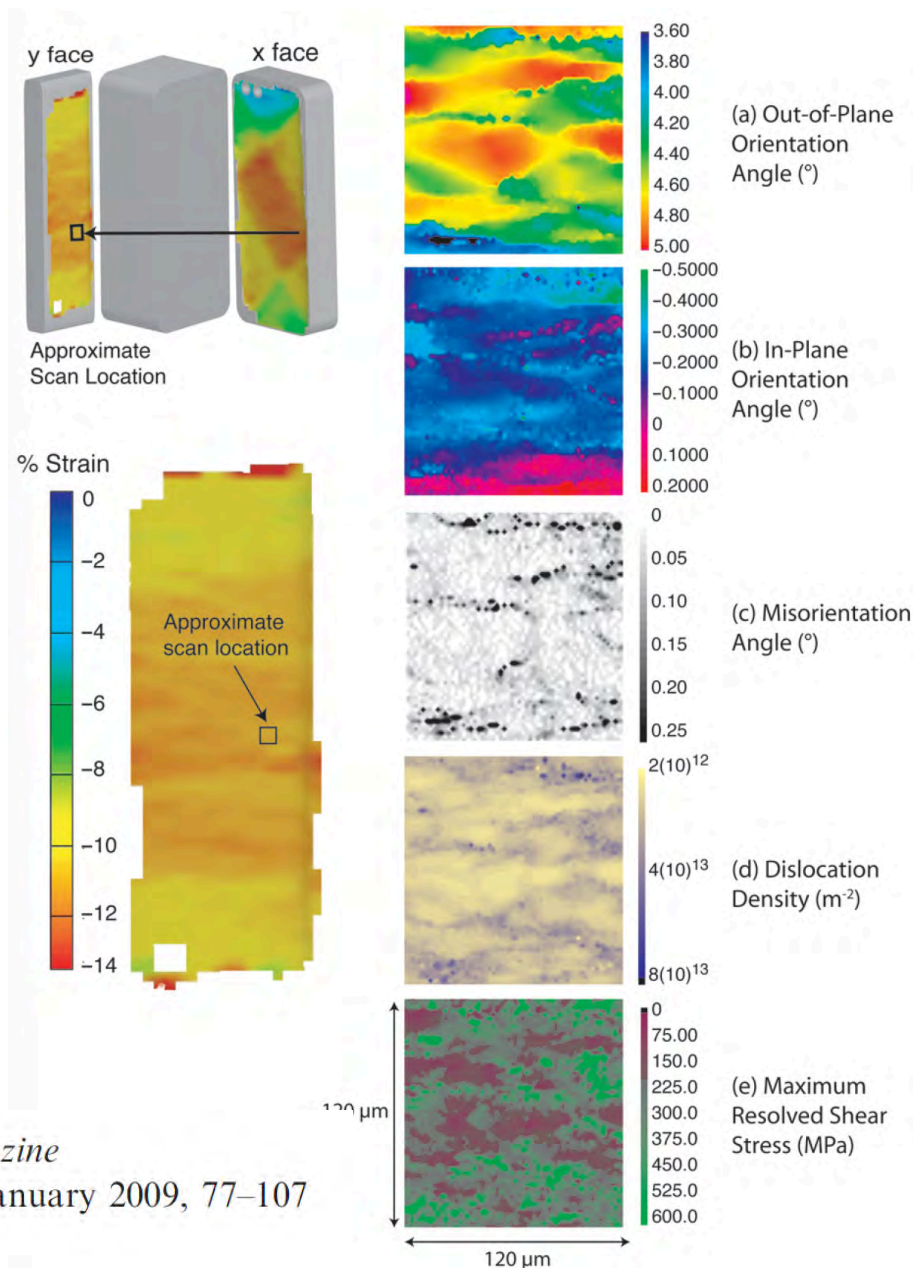
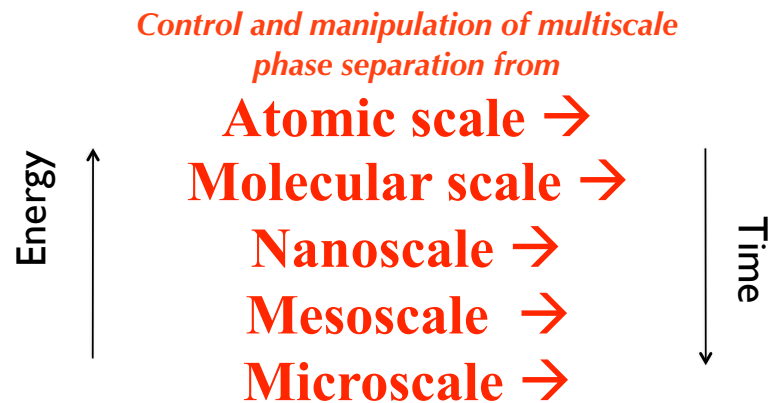


Figure 5. (Colour online). Orientation maps taken at EBSD step sizes of (a) 0.1  $\mu\text{m}$  and 1.0  $\mu\text{m}$ , with the colour scale being the same as that shown in Figure 3a.

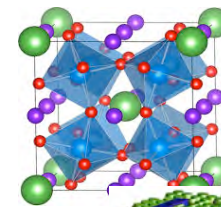
# MULTISCALE PHASE SEPARATION



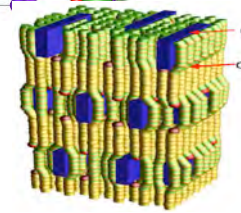
may determine new fundamental material functionality in many challenging open problems in

**Energy**  
**Biology & Biomedicine**  
**Technology**  
**Environmental science**  
**Fundamental science**

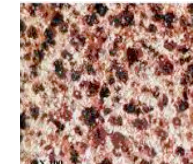
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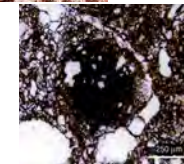
perovskite



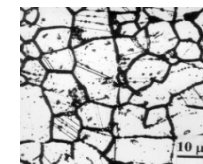
bone



chocolate



soil



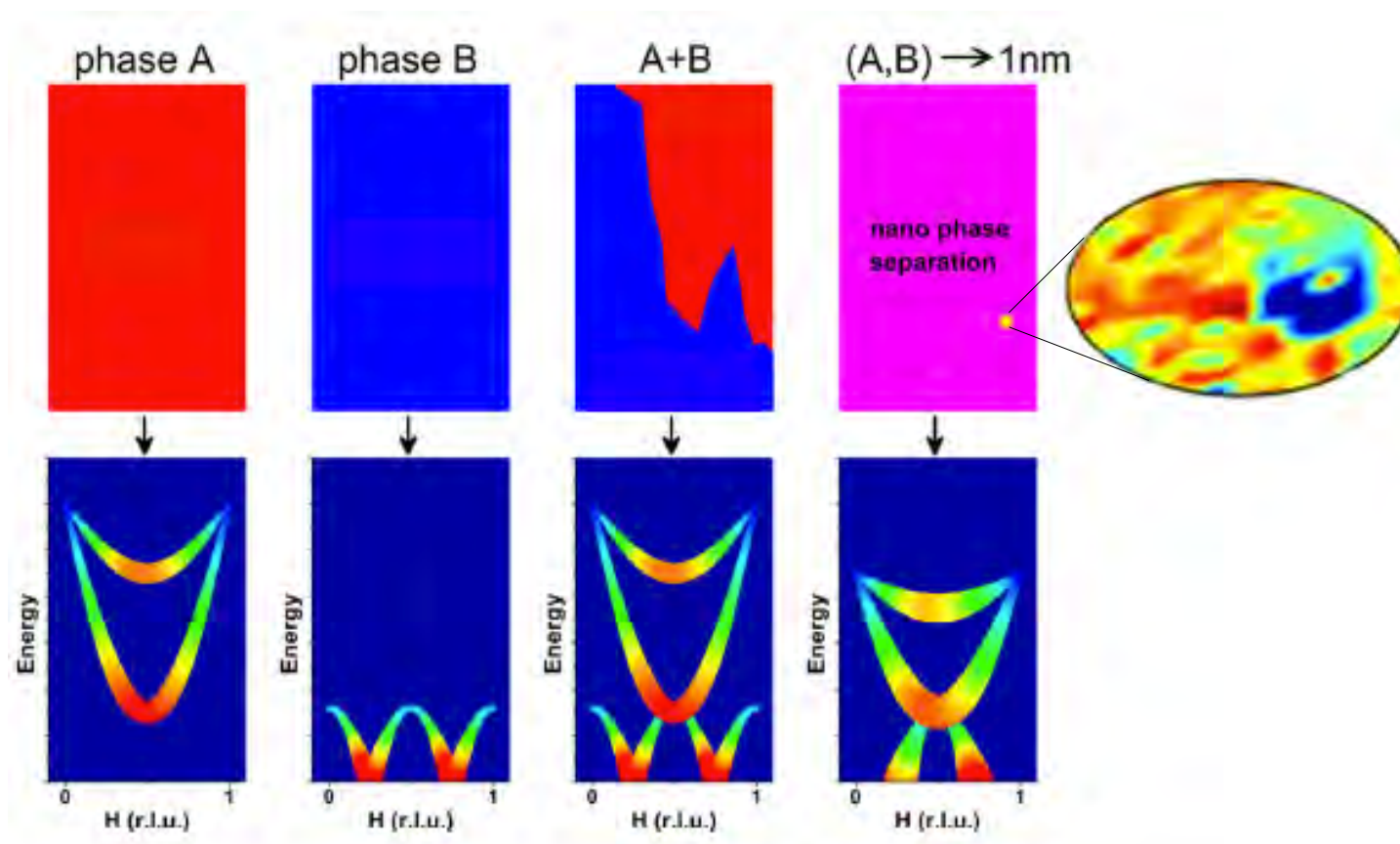
steel

**commentary**

# Multiscale materials modelling at the mesoscale

Sidney Yip and Michael P. Short

The challenge to link understanding and manipulation at the microscale to functional behaviour at the macroscale defines the frontiers of mesoscale science.





**The new physics of this century is focusing to understand novel intricate non trivial mesoscale inhomogeneity in heterogeneous materials where novel low energy Quantum Physics is emerging.**

**The mesoscopic world separates the atomic scale (0.1-10 nm) from the macroscopic world ( $d > 100$  microns).**

## New materials for accelerators

To increase the accelerating gradient of RF cavities working at high frequencies we investigated copper materials coated by Mo via sputtering under vacuum. Structural and electronic characterisation of Mo coatings obtained via the sputtering and annealed up to 600°C have been performed.

This method is a promising approach to obtain homogeneous coatings suitable to increase performances of RF cavities working at high frequencies. We combined FIB imaging to visualise at high spatial resolution the morphology of Mo films and to accurately measure their thickness, with transport experiments to measure the resistivity, while XRD and XAS were used to evaluate the degree of crystallinity, identify different ordered phases and probe the local structure and their electronic properties.

Dedicated RF devices with Mo have been already manufactured at LNF. Mo and Mo oxides with a full characterisation of the conductivity properties are under investigation to optimise manufacturing procedures capable to deposit high-performance coatings.

# Why study TMs and TM oxides?

These materials can have unusual and useful electronic and magnetic properties. Many of these properties strongly depend on defects such as vacancies, dislocations, stacking faults, grain boundaries, etc. These “defects” affect the local oxygen bonding.

We can grow films trying to improve functional properties, e.g., working with doping, strain, stress, different kinetics of growth (non-equilibrium processes), annealing, etc.

Can strain/stress effects be correlated with changes in electronic properties (e.g., electronic band structure)?

Can structural parameters (e.g., strain, doping, defects, ...) be correlated with changes in magnetic/electronic behaviour?

What type of defects are stable/meta-stable?

Can results be understood in terms of surface chemistry, activity, selectivity, .... etc.?

.....

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Comparison of various exotic conductors as well as a few conventional metals (Ag and Cu) based on the ratio of the experimental kinetic energy to that calculated within band theory.

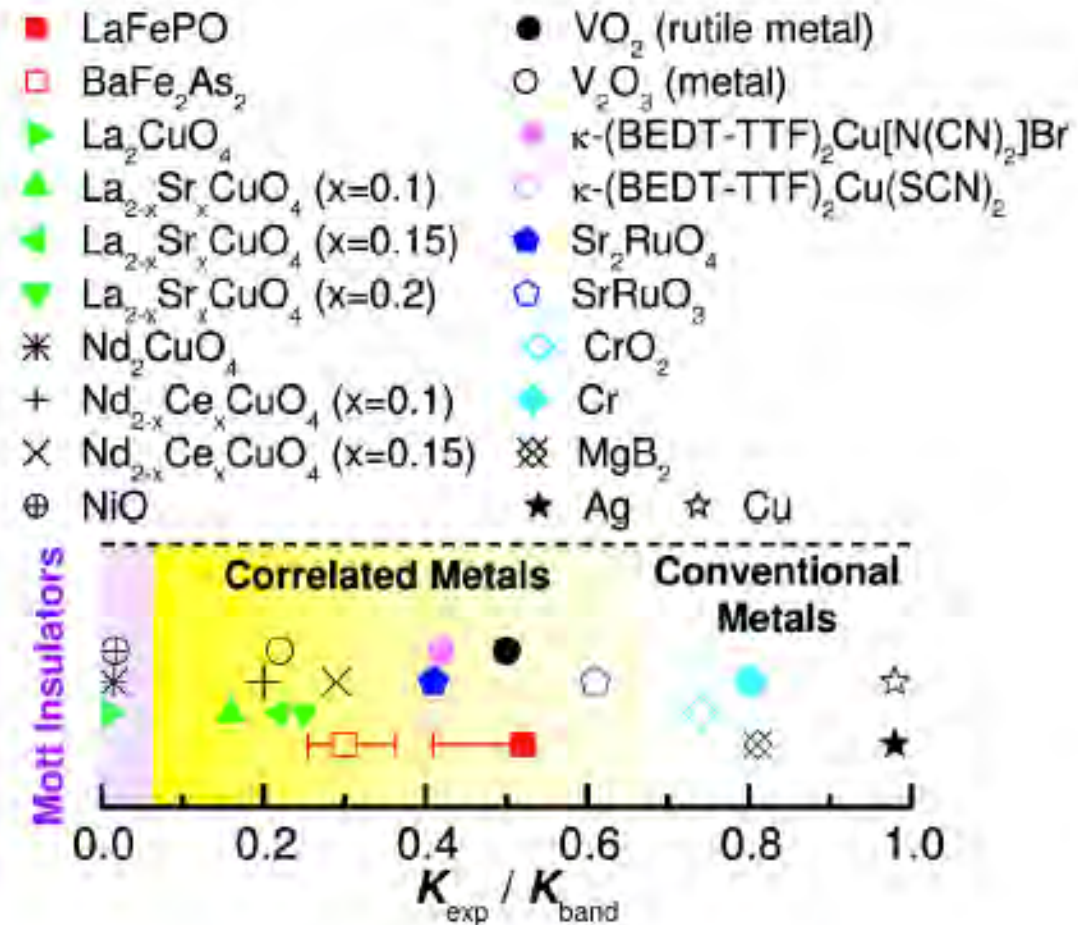
These exotic conductors share at least:

- a) high-T<sub>c</sub> superconductivity
- b) itinerant magnetism and/or
- c) electronic correlations.

The correlated metal regime is characterized by a substantially reduced empirical kinetic energy of charge carriers compared with theory.

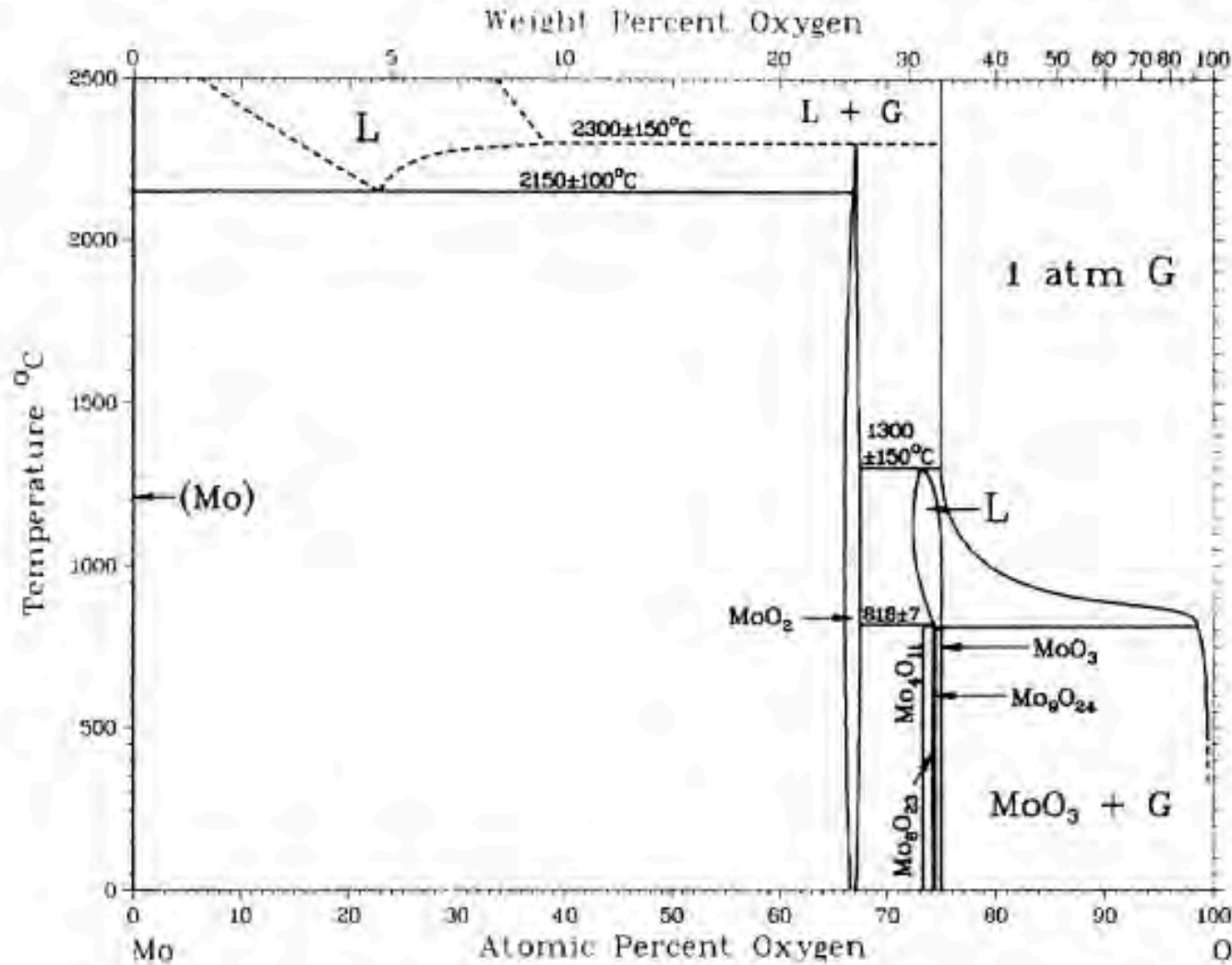
The electron-phonon coupling does not result in a significant reduction of the electron's kinetic energy (unless it is extremely strong as to lead to polaron formation). As an example, the electron-phonon coupling in the MgB<sub>2</sub> superconductor results in  $K_{\text{exp}} = K_{\text{band}} < 1$  in the metallic state.

**Qazilbash et al., Nature Phys. 2009**





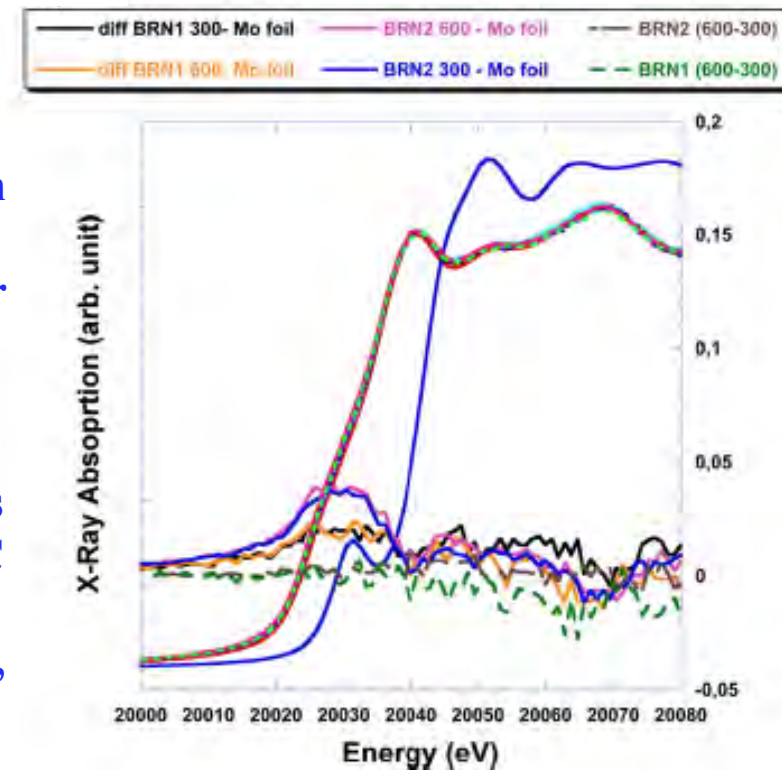
# Mo oxides



# Mo films

Although the solubility of oxygen in solid Mo is extremely low, the growth of Mo films is strongly affected by the presence of oxygen atoms. Because of the complexity of the phase diagram of the Mo–O system with many phases corresponding to oxide compounds with both transparent and insulating phases, e.g.,  $\text{MoO}_3$ , or metallic phases such as  $\text{MoO}_2$ , the different films are multiphase systems of high-interest for both fundamental and technological applications.

An extensive literature is available on Mo-based films and coatings and many researches aim to clarify the role of the internal mechanical stress, the impurity concentration, the texture or the surface roughness and their correlation with deposition parameters, such as pressure, DC power or temperature of the substrate. In spite of the data available, due to the presence of different oxide compositions, amorphous phases, microstructure defects and grain size the properties of Mo films is still not clear what is the lowest achievable resistivity of relatively thick films such as those necessary in high gradient accelerating structures ( $\sim \mu\text{m}$  range).



A. Marcelli et al., *Surface & Coatings Technology*  
261 (2015) 391-397

*Frascati, October 3, 2017*

# Experimental methods

**Rutherford Backscattering (RBS) data show the elemental profile of Mo coatings.**

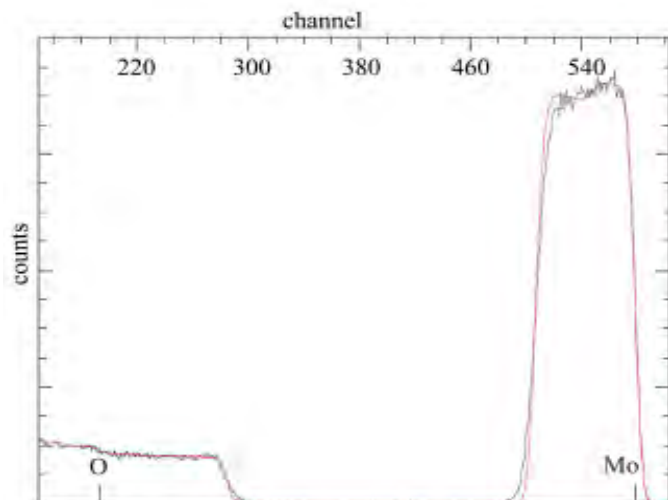


Fig. 12. 2 MeV He (normal incidence, scattering angle  $165^\circ$ ) RBS spectrum (black line) obtained on a Mo film deposited on a  $\text{SiO}_2$  substrate. The simulation (red line) indicates that the Mo film contains oxygen with a maximum concentration of 25%. The film contains  $8.5 \times 10^{17} \text{ Mo/cm}^2$  that corresponds to a Mo coating of about 130 nm thickness. Mo and O surface scattering contributions are reported as green labels.

**Mo 3d photoemission (XPS) spectra giving the oxide content in metallic Mo coatings.**

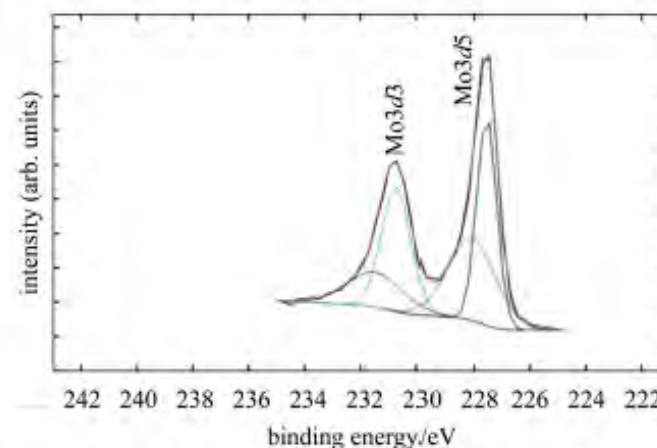


Fig. 14. Deconvolution of the 3d doublet of the molybdenum spectrum. The measured peak (150 nm deep inside the 300 nm film) is well represented by two doublets corresponding to electrons from the 3d shell, the first one at 227.55 in binding energy corresponds to the metallic Molybdenum, the second one at 228.05 shows the typical chemical shift corresponding to the Molybdenum from  $\text{MoO}_2$ .



# Experimental methods

The electrical properties of Mo sputtered coatings have been measured via sample surface impedance probed as a function of the frequency through a Corbino disk geometry with a Vector Network Analyzer Anritsu 37297D. The connection between the cable and the sample was made through a double spring method.

S. Bini, et al., Chinese Physics C 37 (2013) 097005-7

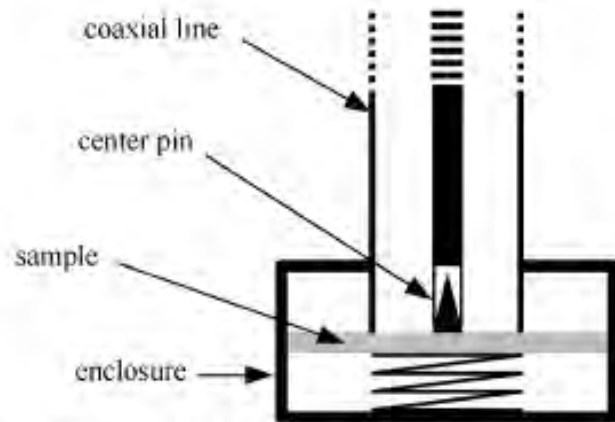


Fig. 5. Sketch of the contact between sample and coaxial line.

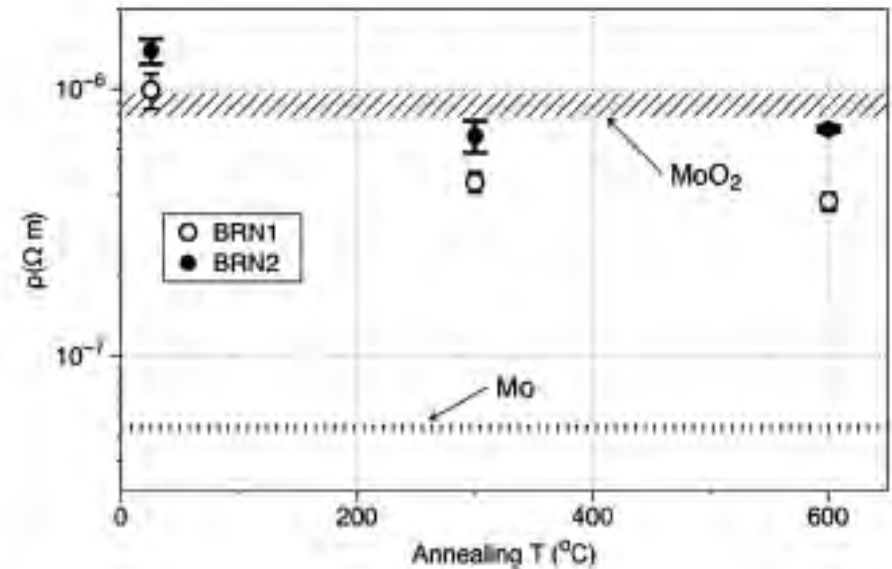


Fig. 2. Resistivity values of BRN1 and BRN2 samples for different annealing temperatures. Mo and MoO<sub>2</sub> values available from the literature are also shown.

In a logarithmic scale we compare the behaviour of the electrical resistivity of some Mo coatings. Although the resistivity values are lower in the thinner sample, after the annealing a clear decrease of the resistivity occurs in both samples, and the observed variation is larger in the thicker samples.

Frascati, October 3, 2017

# Experimental methods

To characterise Mo coatings we used a focused ion beam (FIB) microscope. FIB visualises the coating morphologies at high spatial resolution and allows to measure the thickness of the coating. A typical FIB image from a Mo film grown on  $\text{Al}_2\text{O}_3$  (sample B5) is shown on the right. It is characterised by a thickness of ~615 nm. On such coatings we collected XAS spectra at the Mo K edge.

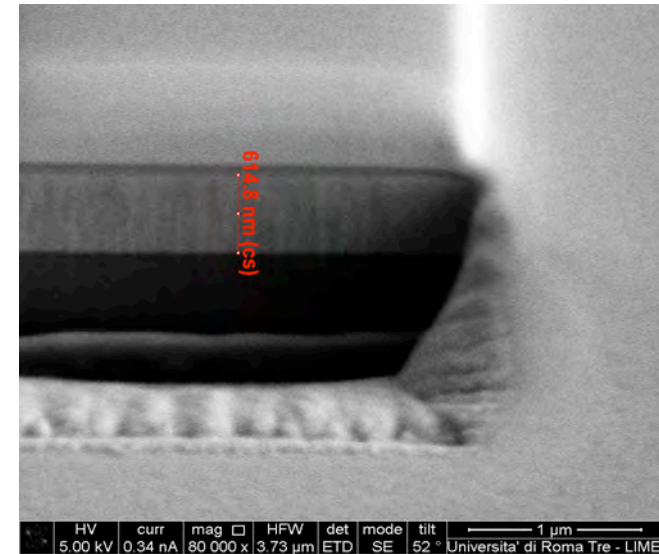


Table 1 Measured thickness of all investigated Mo coatings.

Sample	Mo coatings	Thickness /nm <sup>a</sup>
A1	Mo/ Cu	630
A2	Mo/Cu	1300
B1	Mo/ $\text{Al}_2\text{O}_3$	70-75
B2	Mo/ $\text{Al}_2\text{O}_3$	130-140
B3	Mo/ $\text{Al}_2\text{O}_3$	205-225
B4	Mo/ $\text{Al}_2\text{O}_3$	310
B5	Mo/ $\text{Al}_2\text{O}_3$	615
B6	Mo/ $\text{Al}_2\text{O}_3$	1030

<sup>a</sup>measured with a Focused Ion Beam (FIB)

# X-RAY ABSORPTION SPECTROSCOPY

## LOCAL AND PARTIAL EMPTY DENSITY OF STATES

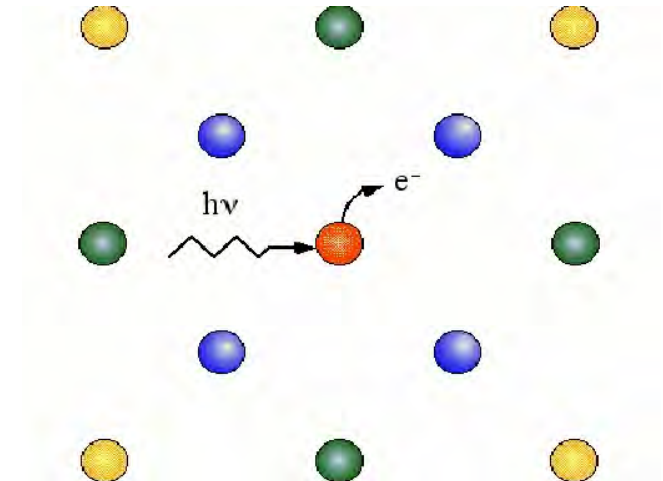
**CROSS SECTION**

$$\sigma(\omega) = 4\pi^2 \alpha \hbar \omega \sum_{if} \{ |\langle f | \hat{\epsilon} \cdot \mathbf{r} | i \rangle|^2 + (1/4) |\langle f | \hat{\epsilon} \cdot \mathbf{r} \mathbf{k} \cdot \mathbf{r} | i \rangle|^2 \} \delta(E_f - E_i - \hbar\omega).$$

ONLY DIPOLE AND ELECTRIC QUADRUPOLE TERMS ARE CONSIDERED

$$\sigma(\omega) = \sigma_\alpha(\omega) [1 + \chi(\omega)] = \sigma_\alpha(\omega) + \sigma_\alpha(\omega) \chi(\omega)$$

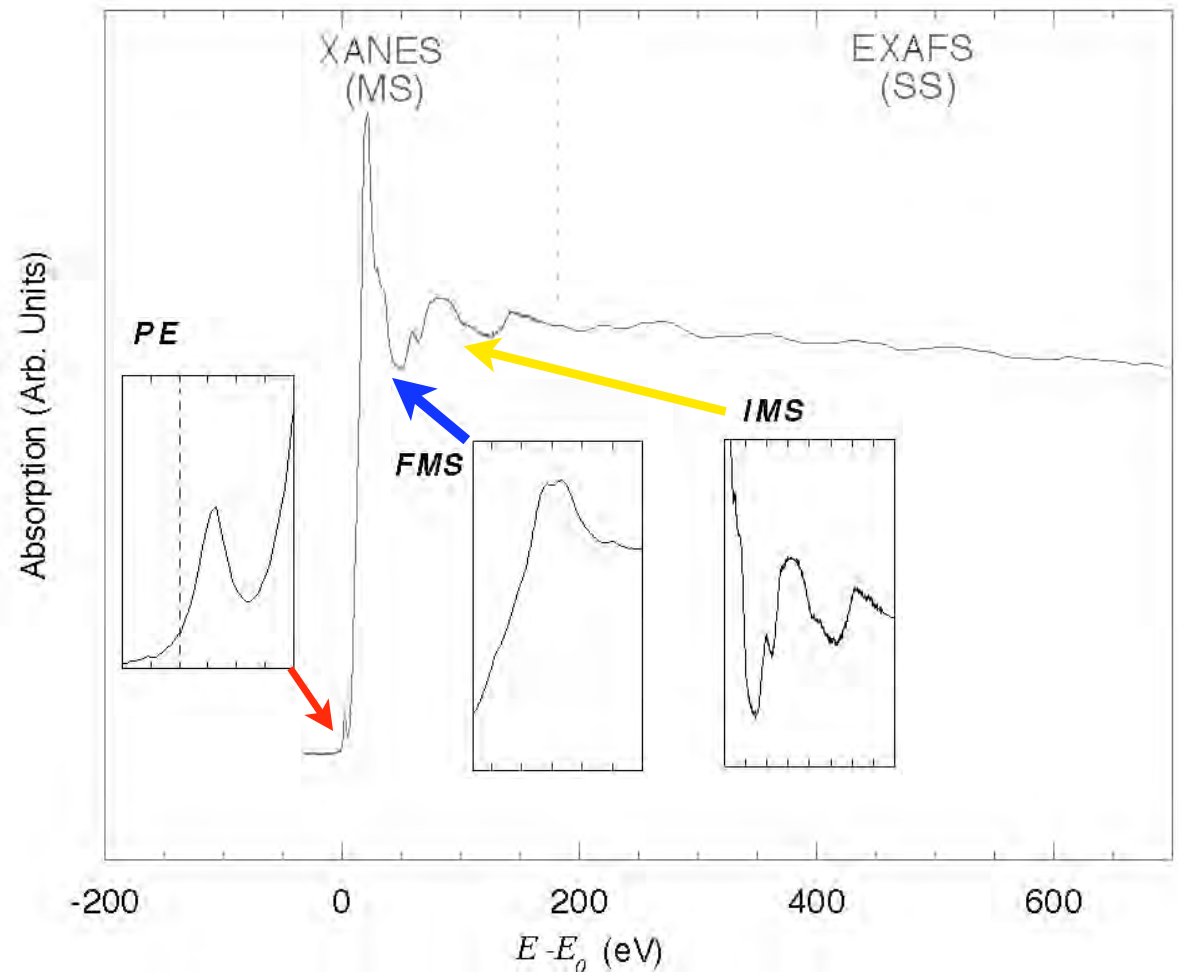
HOW IT WORKS?



Frascati, October 3, 2017

# XANES spectroscopy

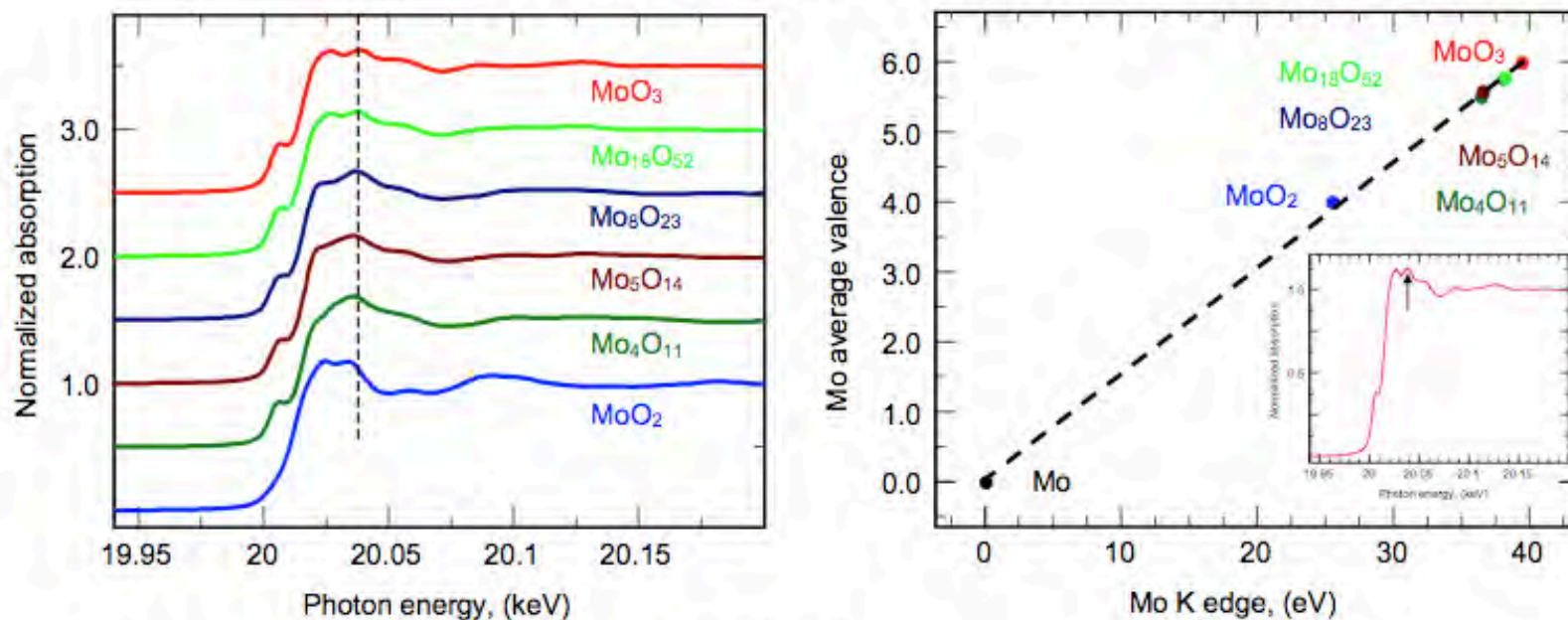
Being a fast ( $10^{-15}$  s) and local probe (few Å) of a selected atomic species, which also does not require a long range order, XANES spectroscopy represents an almost unique spectroscopic technique. Using extremely small spots, without space averaging, it is suited to investigate the local structure and the electronic properties of different polymorphs, systems with vacancies/impurities/defects, materials where complex magnetic, electronic and structural phase separations occur.





# XANES Analysis: Oxidation State

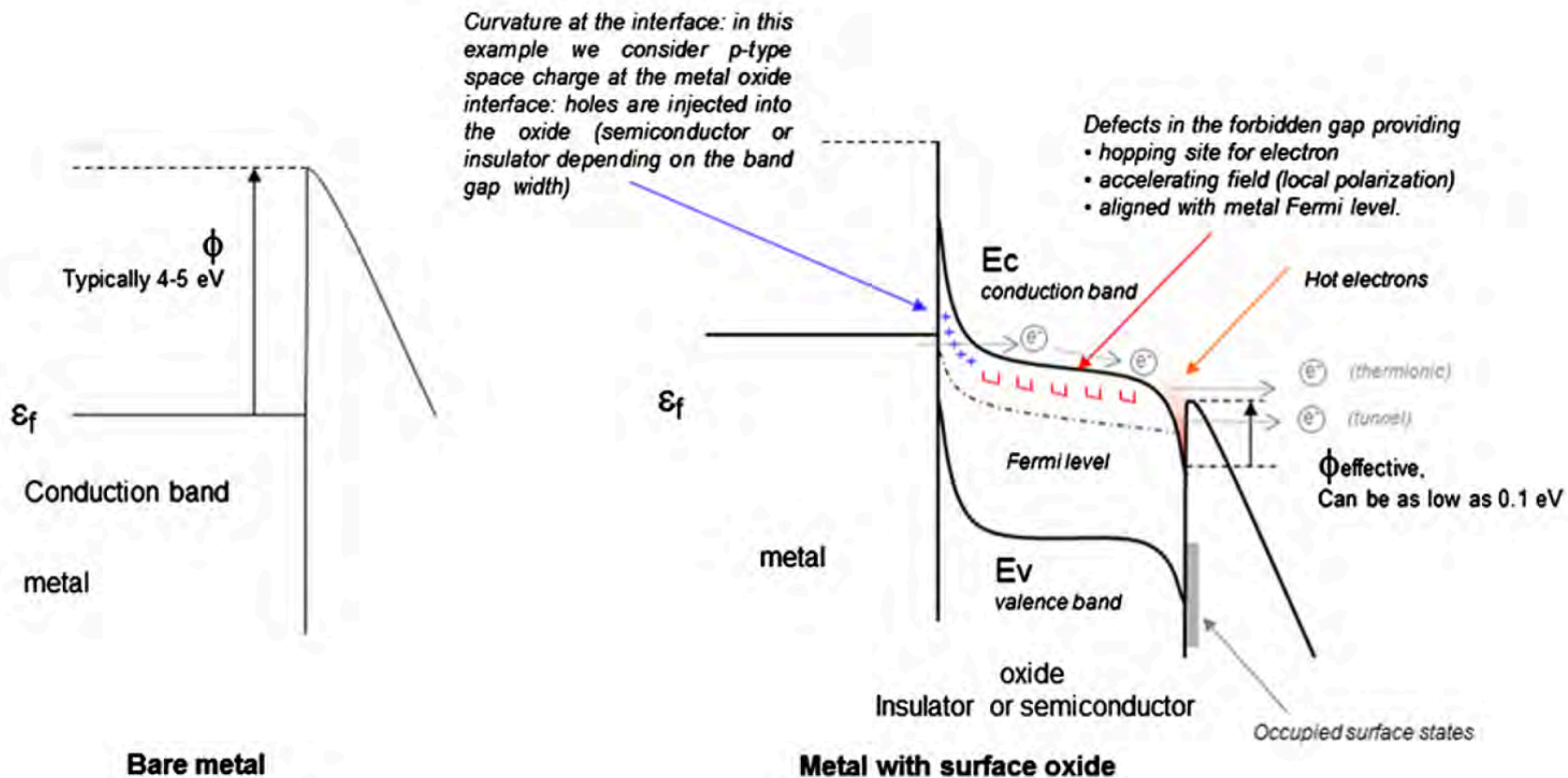
## Mo K-edge XANES of Mo oxides



- Linear fit of Mo valence with K-edge position only obtained using a feature above the absorption edge!

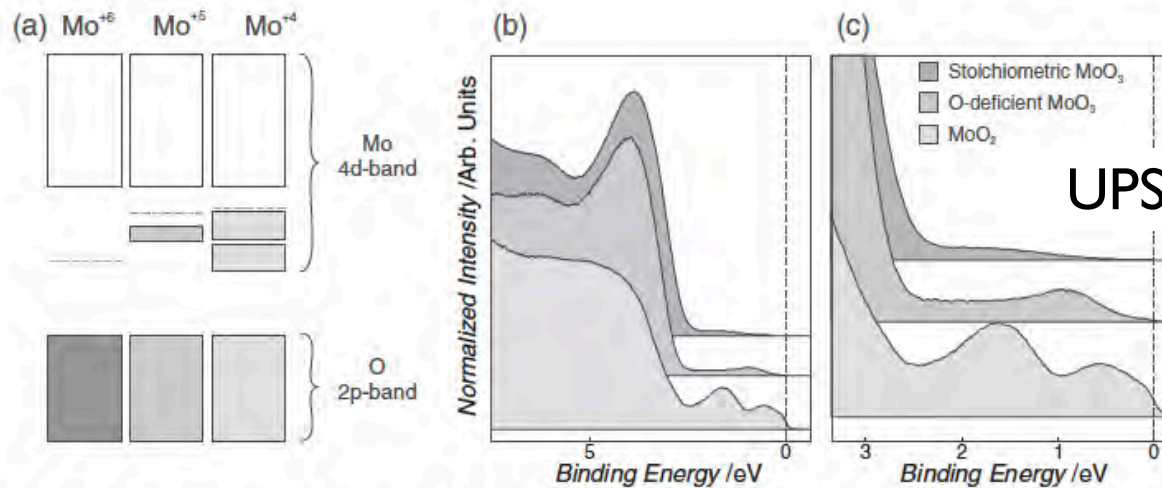
T. Ressler et al. J. Cat 210 (2002) 67

Frascati, October 3, 2017

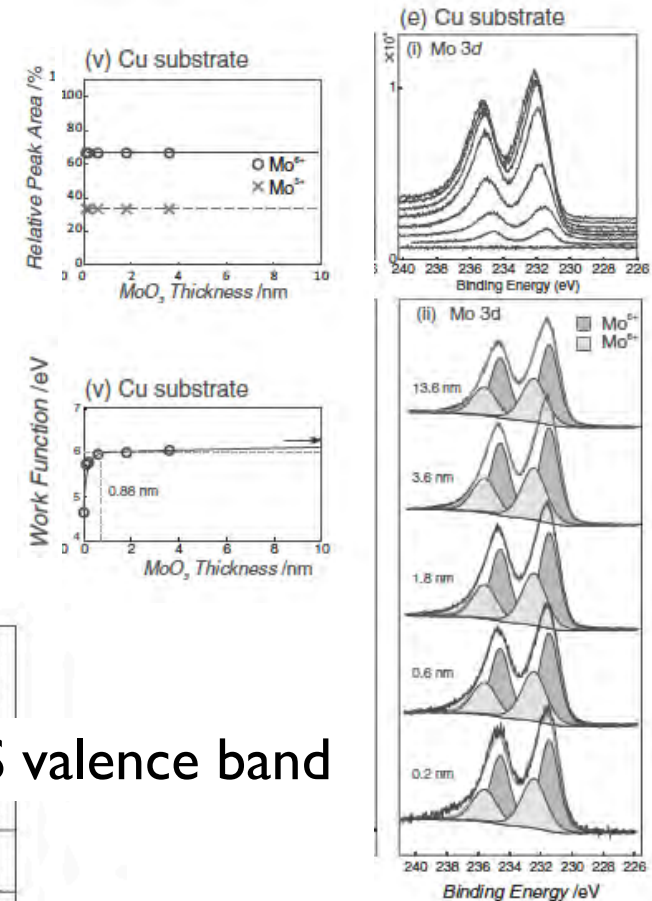


# TM oxides on TM surfaces: Mo on Cu

- calculations of transport properties of coatings of Cu/MoO<sub>3</sub> and alloys such as CuAg/MoO<sub>3</sub>, etc. with thickness from few nm to tens of nm;
- growth of oxides films on Cu, CuAg, CuV, etc. metallic surfaces;
- characterization of coatings and oxides materials;
- breakdown tests and of field emission on these surfaces.



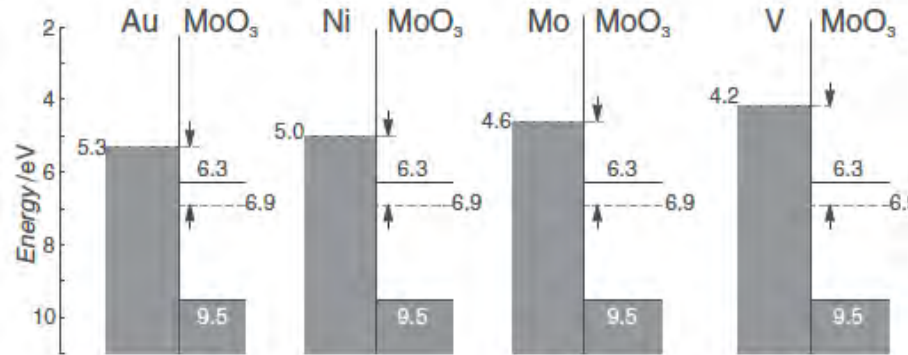
**Figure 2.** a) Schematic energy-level diagram of MoO<sub>3</sub>, O-deficient MoO<sub>3</sub>, and MoO<sub>2</sub>. b) UPS valence band spectra and c) expanded view of the "defect" feature of MoO<sub>3</sub>, O-deficient MoO<sub>3</sub> and MoO<sub>2</sub>.



Mark T. Greiner et al. Adv. Funct. Mater. 23, 215 (2013)

Frascati, October 3, 2017

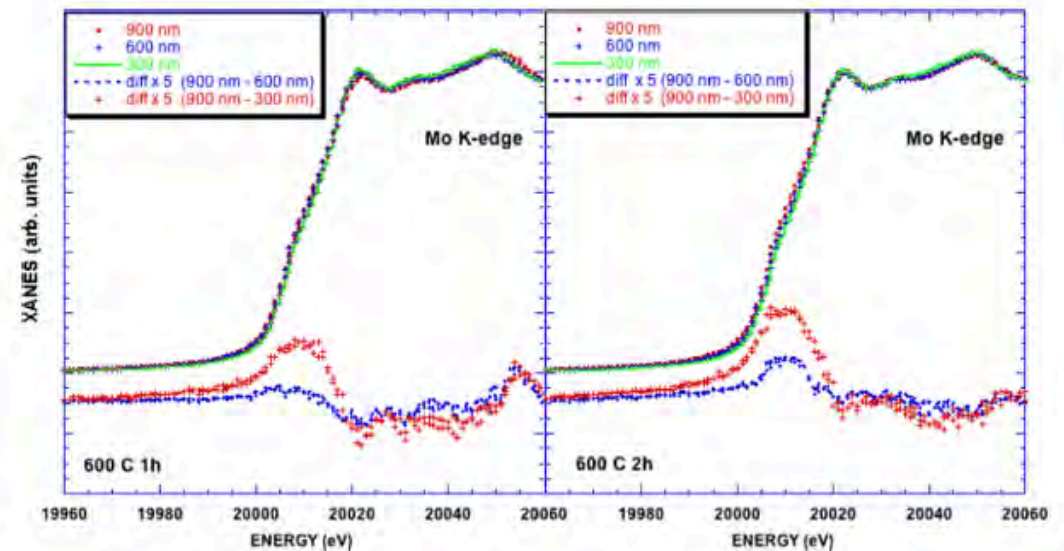
# New coating materials & methods



Mark T. Greiner et al. Adv. Funct. Mater. 23, 215 (2013)

**XAS experiment on Mo growth on different substrates.**

On the left different WFs for MoO<sub>3</sub> on different TMs substrates



An extensive and accurate characterisation of the conductivity and mechanical properties of Mo and Mo oxides coatings may identify reliable procedures capable to growth stable and hard semi-metallic films with a high work function to reduce the dark current.

Frascati, October 3, 2017



## Effects of Surface Nonuniformities on the Mean Transverse Energy from Photocathodes

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(Received 24 March 2015; revised manuscript received 29 July 2015; published 24 August 2015)

The performance of photoinjectors is limited by the lowest value of the mean transverse energy of the electrons obtained from photocathodes. The factors that influence the mean transverse energy are poorly understood. In this paper, we develop models to calculate the effect of spatial work-function variations and subnanometer-scale roughness and surface defects on the mean transverse energy. We show that these can limit the lowest value of mean transverse energy achieved and that atomically perfect surfaces will be required to further reduce the mean transverse energy obtained from photocathodes.

DOI: 10.1103/PhysRevApplied.4.024015

JOURNAL OF APPLIED PHYSICS 114, 243302 (2013)



### Local changes of work function near rough features on Cu surfaces operated under high external electric field

Flyura Djurabekova,<sup>1,a)</sup> Avaz Ruzibaev,<sup>1</sup> Eero Holmström,<sup>2,3</sup> Stefan Parviainen,<sup>1</sup> and Mikko Hakala<sup>2</sup>

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<sup>2</sup>Department of Physics, University of Helsinki, P.O. Box 64, FIN-00014 Helsinki, Finland

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(Received 30 October 2013; accepted 9 December 2013; published online 30 December 2013)

**“The change in the value of the work function indicates the strong influence of intrinsic defects on this parameter, which requires further investigation. Moreover, we were able to validate the approach which we developed previously for the dynamic simulation of the effect of electric field on extended metal surfaces.”**

*Frascati, October 3, 2017*

## Characterization of the transport properties of MoO<sub>3</sub> films on copper

C. Bonavolontà<sup>1)</sup>, M. Valentino<sup>1,2)</sup>, M. de Lucia<sup>1)</sup>, M. Ambrosio<sup>1)</sup>,  
C. Aramo<sup>1)</sup>, S. Macis<sup>3)</sup>, I. Davoli<sup>3)</sup>, G. Castorina<sup>4,5)</sup>, F. Monforte<sup>6)</sup>,  
B. Spataro<sup>4)</sup>, M. Scarselli<sup>3)</sup>, S. Lupi<sup>7)</sup> and A. Marcelli<sup>4,8)</sup>

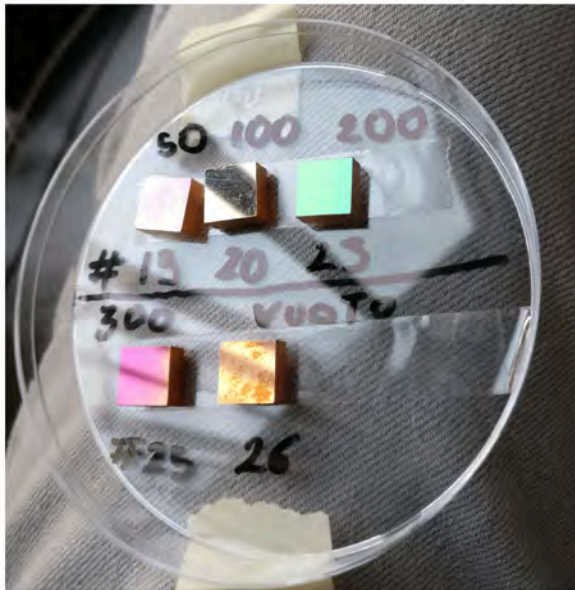


Figure 2: Photograph showing evaporated samples of MoO<sub>3</sub> on copper. Due to the different thickness, each film can be easily recognized by its different colour.

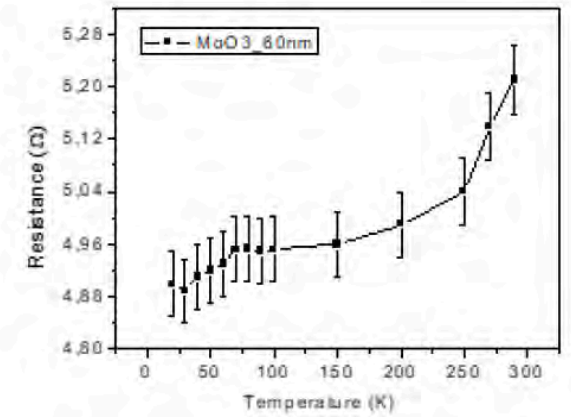
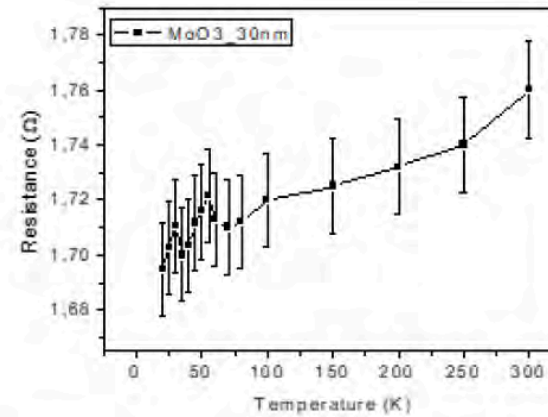
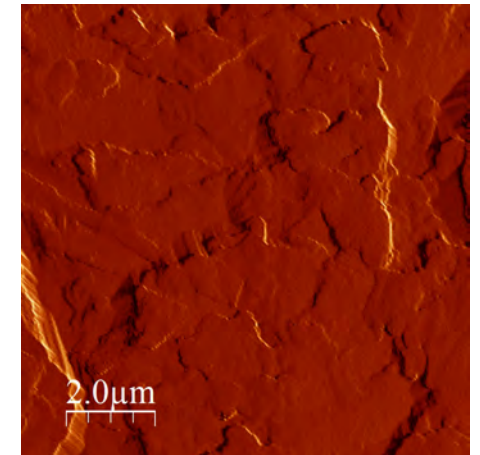
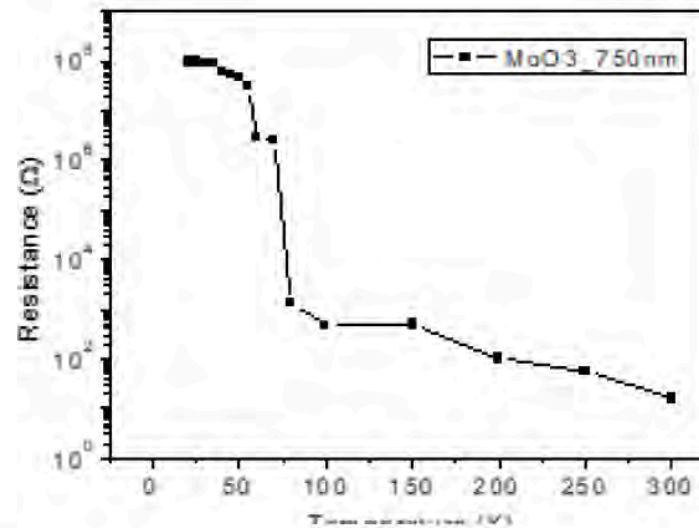


Figure 5: Electrical resistance vs. temperature for MoO<sub>3</sub> samples with thickness: 30 nm (left) and 60 nm (right).

Thin films are conductive and the resistance decreases at low temperature.



Frascati, October 3, 2017

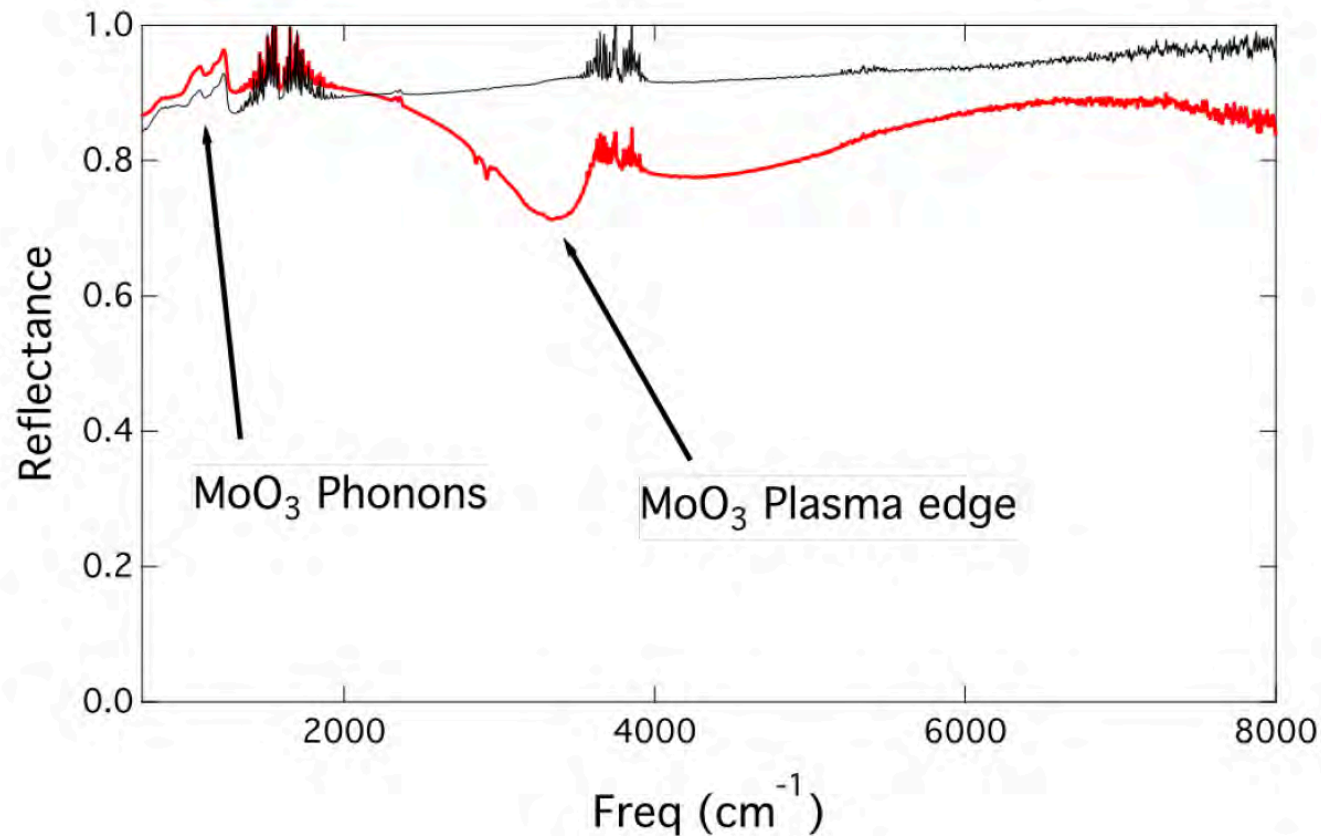


Figure 12: Comparison of the absolute reflectance of MoO<sub>3</sub> films of 30 and 300 nm thickness. While the thinner film shows, at room temperature, a metallic behaviour with a well evident plasma edge in the mid-infrared region, the 300 nm thick film shows an insulating behaviour with a flat ( $<1$ ) reflectance. The narrow peaks around 1600 and 3900 cm<sup>-1</sup> are due to the water vapour absorption in the spectrometer.



# XAS of Mo coatings

XAS (X-ray Absorption Spectroscopy) is a well-established technique for a direct quantitative measurement of local structure properties such as geometry and coordination numbers in ordered and disorder systems such as coatings or interface. To characterise the chemical status of Mo atoms in our coatings we performed XAS experiments at the Mo K edge at B18, Core XAS beamline at the Diamond Light Source (UK).

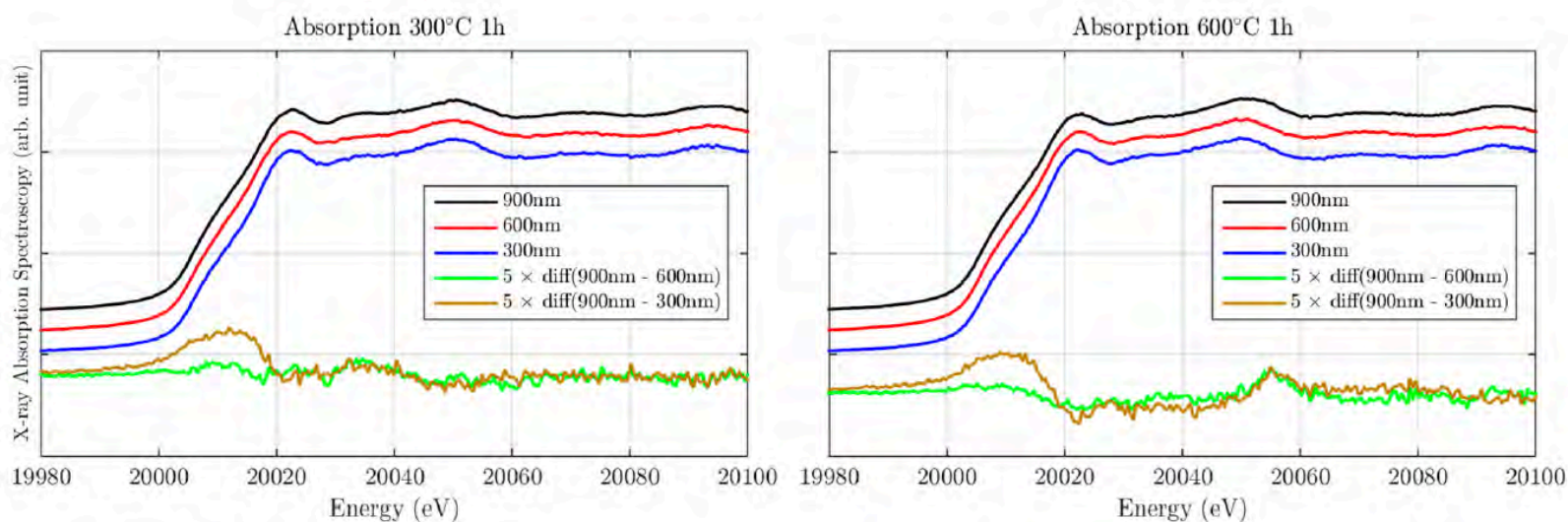
Data were collected with focusing radiation in a spot of  $\sim 200 \times 200 \text{ }\mu\text{m}^2$  in the fluorescence mode. Experiments have been performed both at normal incidence and at grazing incidence to enhance the signal associated to thin surface layers and in different areas of the Mo coatings.



Article

# Materials and Breakdown Phenomena: Heterogeneous Molybdenum Metallic Films

Augusto Marcelli <sup>1,2,3,\*</sup>, Bruno Spataro <sup>1</sup>, Giovanni Castorina <sup>1,4</sup>, Wei Xu <sup>2,5</sup>, Stefano Sarti <sup>6</sup>,  
Francesca Monforte <sup>7,8</sup> and Giannantonio Cibin <sup>9</sup>



**Figure 2.** (Left panel) comparison of Mo K-edge spectra of samples of 300 nm, 600 nm, and 900 nm annealed at 300 °C for 1 h. The bottom curves are the differences between spectra of different thickness (multiplied 5×); (Right panel) the same as in the left panel for the samples with the same thickness, but annealed at 600 °C for 1 h.

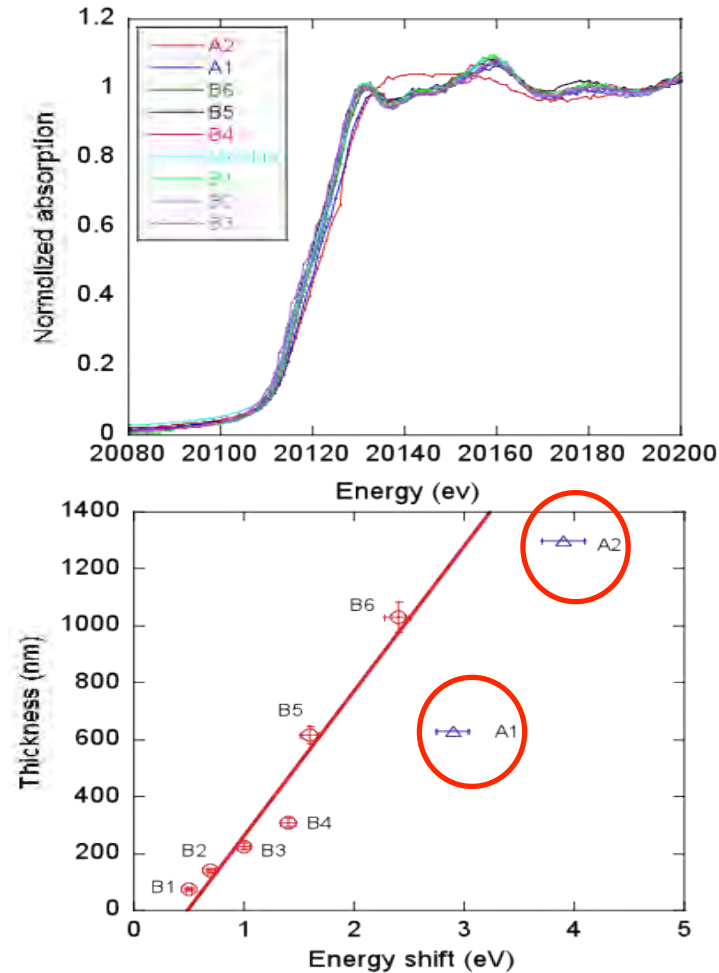
# Mo K edge XANES

Comparison of XANES spectra of Mo coatings (top right). Spectra measured in the transmission geometry point out that all coatings has a slightly disordered structure with a negligible Mo oxide contribution. Coatings are similar except for the thick A2 sample (i.e., Mo/Cu) which has a significant Mo oxide contribution. In the bottom right we plot the edge energy vs. thickness where the Mo K-edge has been determined from the normalized absorption at the height of 0.8. The edge of samples grown on  $\text{Al}_2\text{O}_3$  (B1-B6) shift vs. thickness due to the changes of the Fermi energy. For samples grown on Cu (A1-A2), XANES spectra are slightly different (in particular the thicker one). Mo coatings are all characterized by a linear expansion of the next-nearest distances around the Mo atoms that clearly depends by the thickness.

A. Marcelli, B. Spataro, S. Sarti, et al. Characterization of thick conducting molybdenum films: enhanced conductivity via thermal annealing, *Surface & Coatings Technology* 261, 391 (2015)

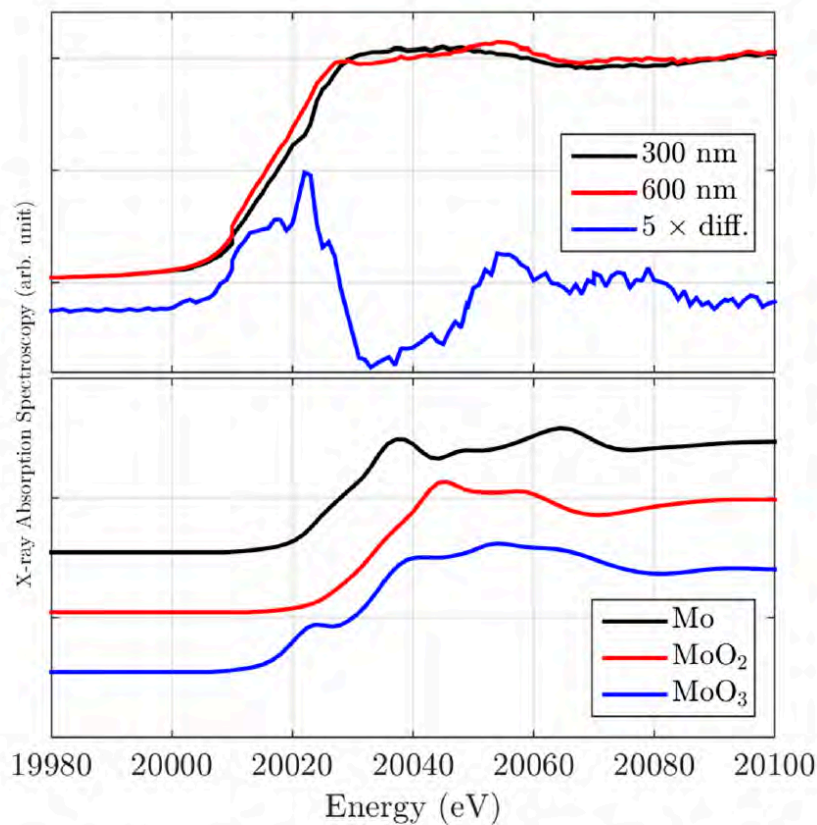
S. Bini, B. Spataro, A. Marcelli, S. Sarti, et al., Molybdenum sputtering film characterization for high gradient accelerating structures, *Chinese Physics C* 37, 097005 (2013)

Y. Xu, B. Spataro, S. Sarti, et al., Structural and morphological characterization of Mo coatings for high gradient accelerating structures, *Journal of Physics: Conference Series* 430, 012091 (2013)

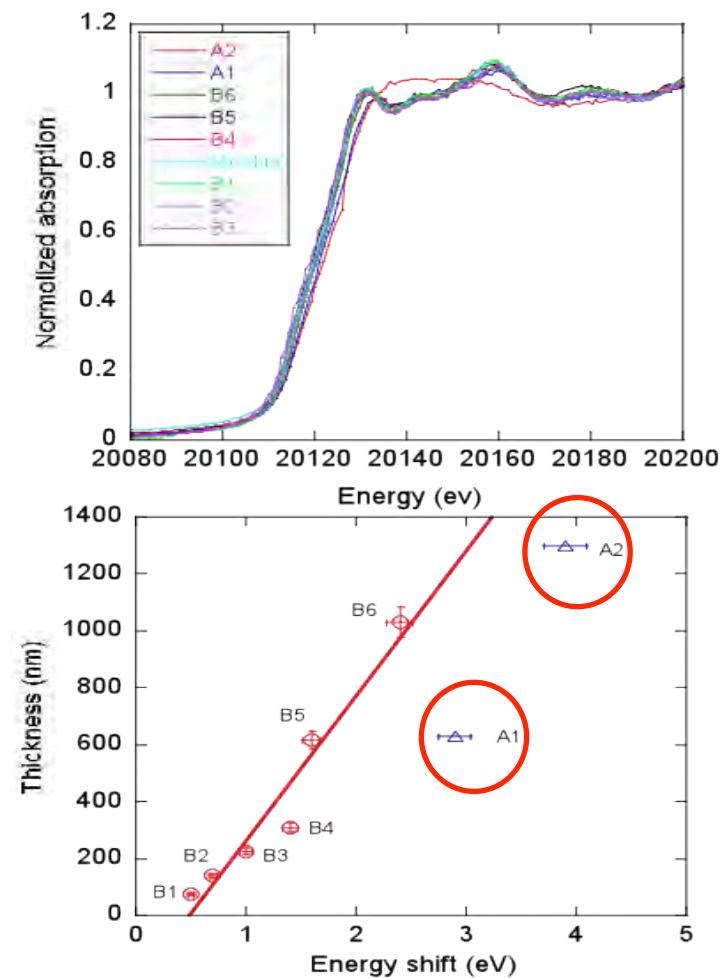


Frascati, October 3, 2017

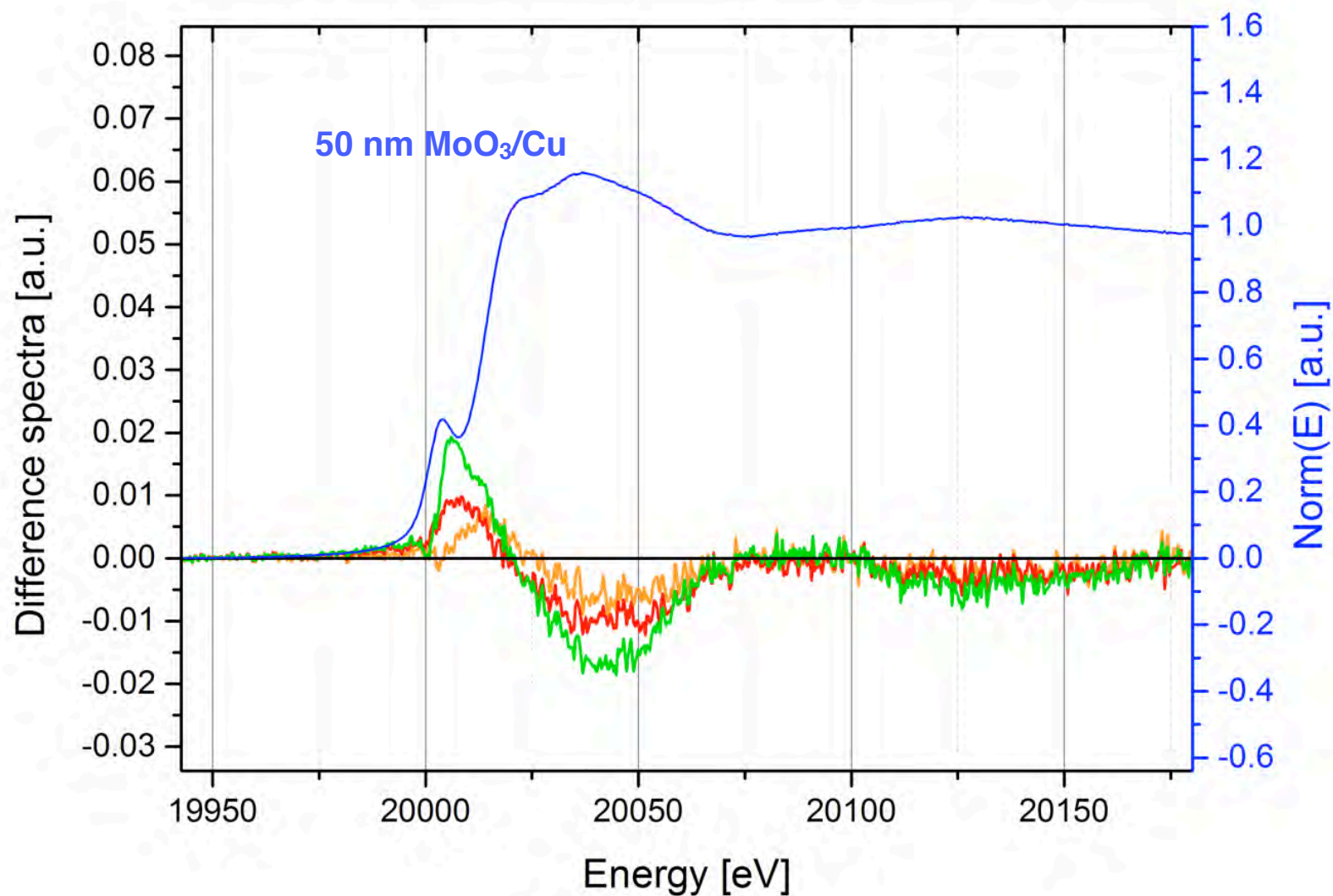
# Mo K edge XANES



**Figure 3.** (Top panel) comparison of Mo K-edge XANES spectra of two Mo films: 300 nm and 600 nm thick grown on OFHC copper substrates. The blue curve is the difference between the two experimental spectra multiplied by five to enhance the difference at the edge (see text); (Bottom panel) comparison among Mo K-edge XANES spectra of Mo metal (black), MoO<sub>2</sub> (red), and MoO<sub>3</sub> (blue) simulated using the full multiple scattering (FMS) theory (see text).



Frascati, October 3, 2017



XAS(100nm-50nm)

XAS(200nm-50nm)

XAS(300nm-50nm)



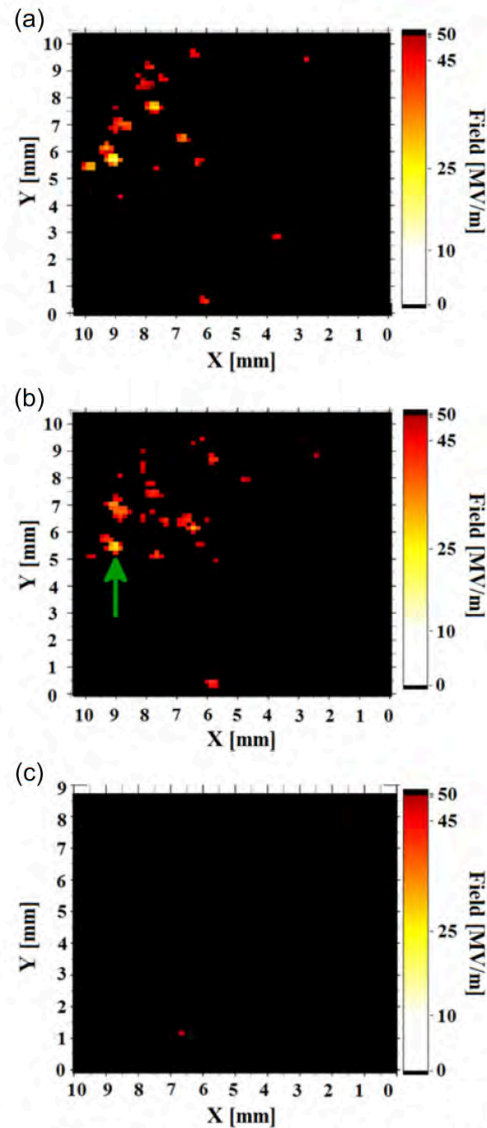


Fig. 4. FE maps up to  $E = 50$  MV/m of the single crystal Mo sample after HT at  $500\text{ }^{\circ}\text{C}$  (a) and  $600\text{ }^{\circ}\text{C}$  (b) and of the much better cleaned polycrystalline sample after HT at  $600\text{ }^{\circ}\text{C}$  (c).

Eur. Phys. J. Appl. Phys. (2015) 70: 21301  
 DOI: [10.1051/epjap/2015150167](https://doi.org/10.1051/epjap/2015150167)

THE EUROPEAN  
 PHYSICAL JOURNAL  
 APPLIED PHYSICS

Regular Article

## Prevention of electron field emission from molybdenum substrates for photocathodes by the native oxide layer

Stefan Lagotzky<sup>1,a</sup>, Roman Barday<sup>2</sup>, Andreas Jankowiak<sup>2</sup>, Thorsten Kamps<sup>2</sup>, Carola Klimm<sup>2</sup>, Jens Knobloch<sup>2</sup>,  
 Günter Müller<sup>1</sup>, Boris Senkovskiy<sup>3</sup>, and Frank Siewert<sup>2</sup>

<sup>1</sup> Bergische Universität Wuppertal, FB C Abteilung Physik, Gauss-Str. 20, 42097 Wuppertal, Germany

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Received: 20 March 2015 / Accepted: 21 April 2015

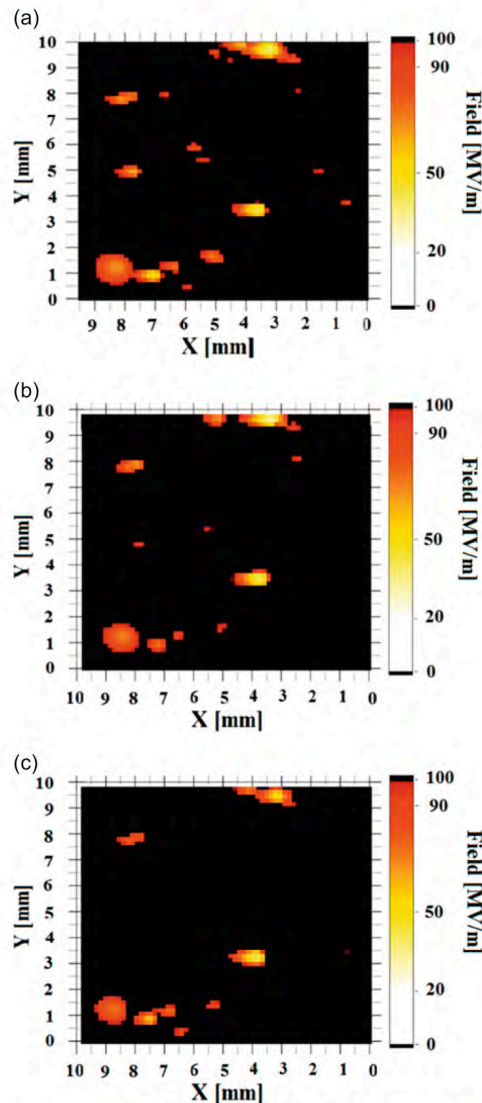
Published online: 25 May 2015 – © EDP Sciences 2015

FE maps up to  $E = 50$  MV/m on the single crystal Mo sample after HT at  $500\text{ }^{\circ}\text{C}$  (a) and  $600\text{ }^{\circ}\text{C}$  (b) and of the much better cleaned polycrystalline sample after HT at  $600\text{ }^{\circ}\text{C}$  (c).

Well-polished and dry-ice cleaned Mo samples with native oxide did not show parasitic FE up to a field level of 50 MV/m required for photoinjector cavities.

No FE was observed after HT up to  $400\text{ }^{\circ}\text{C}$ , but 19 (26) emitters occurred after HT at  $500$  ( $600$ ) $^{\circ}\text{C}$  as shown on the left.

*Frascati, October 3, 2017*



**Fig. 6.** FE maps of the polycrystalline Mo sample up to  $E = 100$  MV/m after HT at  $600\text{ }^{\circ}\text{C}$  (a) and after different reoxidation steps with  $\text{O}_2$  pressure:  $p = 10^5$  Pa at  $25\text{ }^{\circ}\text{C}$  (b) and  $p = 1$  Pa at  $400\text{ }^{\circ}\text{C}$  (c).

.... the same polycrystalline Mo sample was scanned up to  $100\text{ MV/m}$  to enhance its FE significantly. Beside both formerly present emitters, about 18 additional ones with  $E_{\text{on}}$  values of  $30\text{--}90\text{ MV/m}$  could be activated now as shown in Figure 6a. It is remarkable that some new emitter show rather low values for  $E_{\text{on}}$  now. This well-known activation effect for field emitters is usually explained by the creation of conducting channels in the insulating oxide layer .....

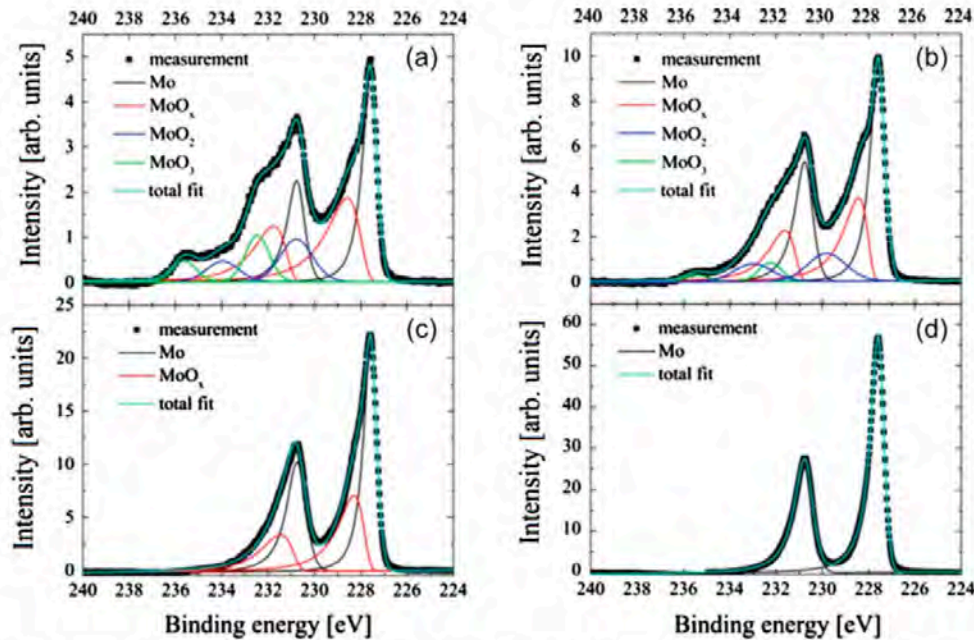


Fig. 11. Mo 3d photoemission spectra after HT at 100 °C (a), 300 °C (b), 400 °C (c) and 600 °C (d).

After the XPS measurement the sample was moved into a preparation chamber for a HT. Four treatment cycles were performed in situ with increasing temperature of 100, 300, 400 and 600°C for 1 h. After each treatment step the sample was transferred back into the analyzing chamber for the XPS measurement.

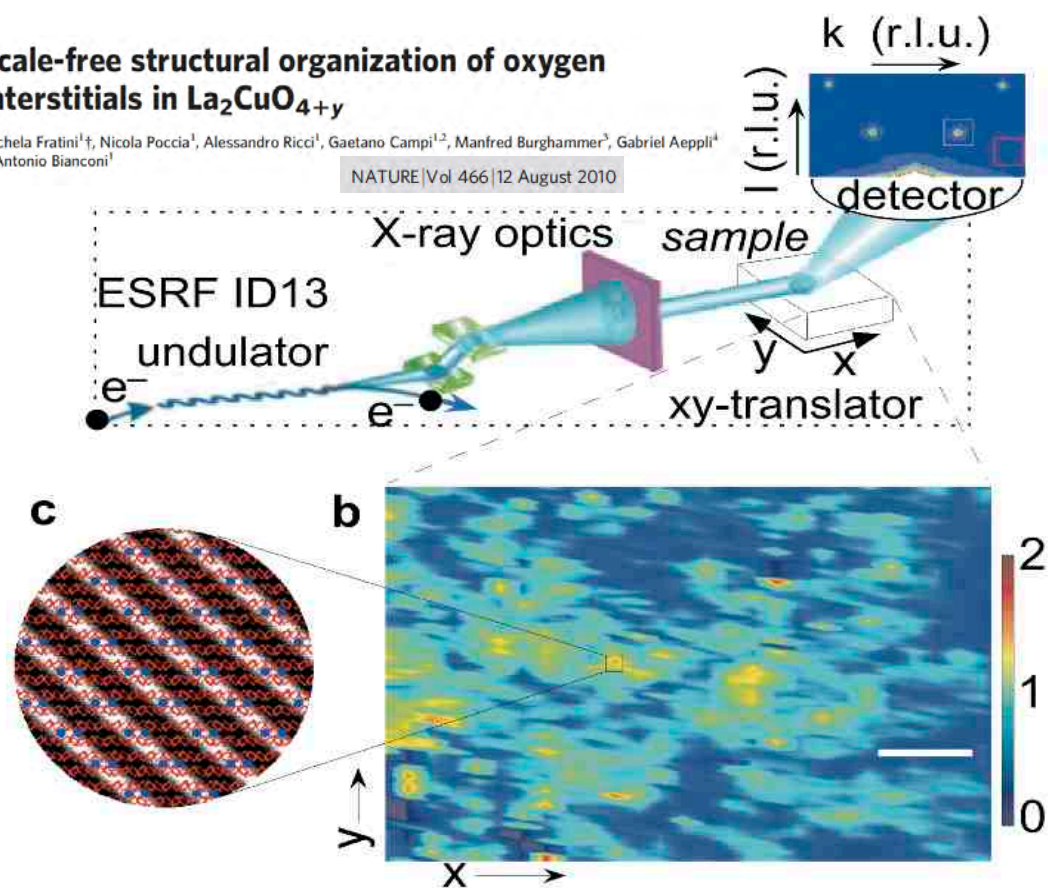
The position of the illuminated area remains unchanged. The weight  $f$  for MoO<sub>3</sub> reduces from ~0.33 for the unannealed sample to 0.05 after HT at 300–600°C.

After HT at 400°C the molybdenum trioxide layer is dissolved into the bulk (Fig. 11). After the last HT at 600°C all Mo oxide correlated peaks disappeared, indicating that the oxide film was removed completely.

# Scale-free structural organization of oxygen interstitials in $\text{La}_2\text{CuO}_{4+y}$

Michela Fratini<sup>1†</sup>, Nicola Poccia<sup>1</sup>, Alessandro Ricci<sup>1</sup>, Gaetano Campi<sup>1,2</sup>, Manfred Burghammer<sup>3</sup>, Gabriel Aeppli<sup>4</sup> & Antonio Bianconi<sup>1</sup>

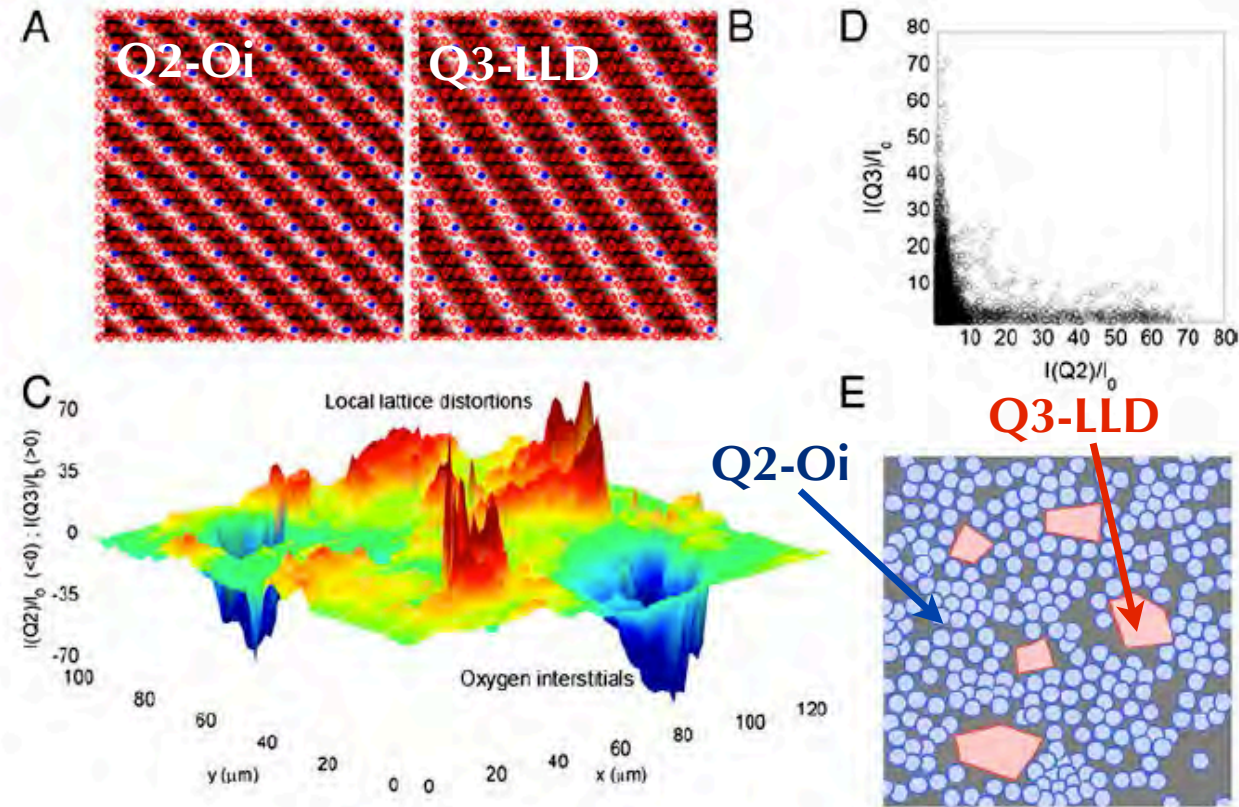
NATURE | Vol 466 | 12 August 2010





Even in optimal superconducting samples the oxygen defect order can be highly inhomogeneous. Actually, the glue regions contain incommensurate modulated local lattice distortions, whose spatial extent is most pronounced for the best superconducting samples. For an underdoped single crystal with mobile oxygen interstitials in the spacer  $\text{La}_2\text{O}_{2+y}$ , the incommensurate modulated local lattice distortions form droplets anti-correlated with the ordered Oi. In the simplest among high temperature superconductors, two networks of ordered defects coexist. They can be tuned to achieve optimal superconductivity.

Actually, for a given stoichiometry, the highest transition temperature is obtained when both the ordered oxygen (Oi) and the lattice defects (LLD) form fractal patterns.



## Optimum inhomogeneity of local lattice distortions in $\text{La}_2\text{CuO}_{4+y}$

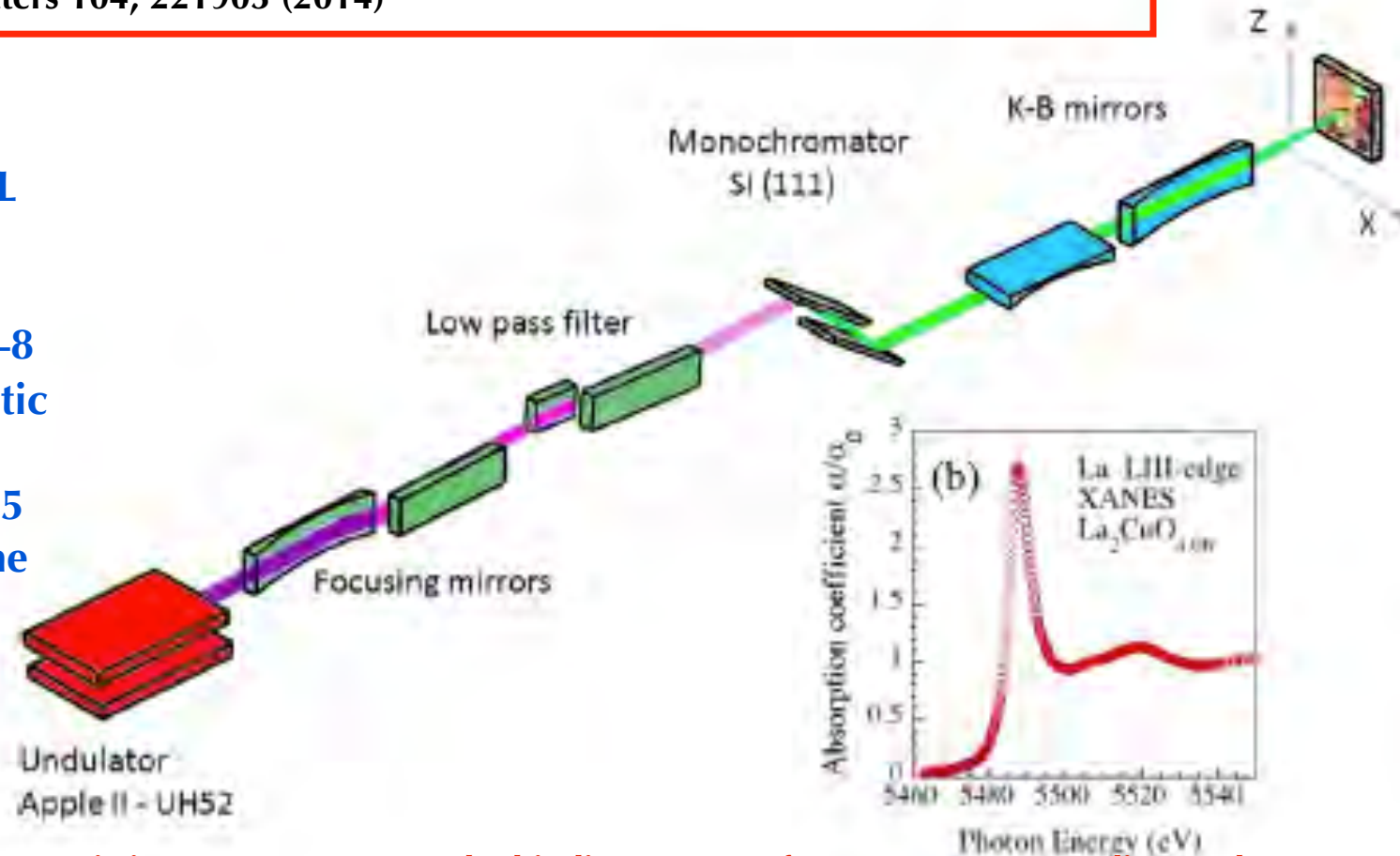
Nicola Poccia<sup>a,b</sup>, Alessandro Ricci<sup>a,c</sup>, Gaetano Campi<sup>d</sup>, Michela Fratini<sup>a,e</sup>, Alessandro Puri<sup>f</sup>, Daniele Di Gioacchino<sup>f</sup>, Augusto Marcelli<sup>f</sup>, Michael Reynolds<sup>b</sup>, Manfred Burghammer<sup>b</sup>, Naurang Lal Saini<sup>a</sup>, Gabriel Aeppli<sup>g</sup>, and Antonio Bianconi<sup>a,h,i,1</sup>

PNAS | September 25, 2012 | vol. 109 | no. 39 | 15685–15690

Frascati, October 3, 2017

N. Poccia, M. Chorro, A. Ricci, Wei Xu, A. Marcelli, G. Campi and A. Bianconi  
**Percolative superconductivity in  $\text{La}_2\text{CuO}_{4.06}$  by lattice granularity patterns with scanning micro x-ray absorption near edge structure**  
 Applied Physics Letters 104, 221903 (2014)

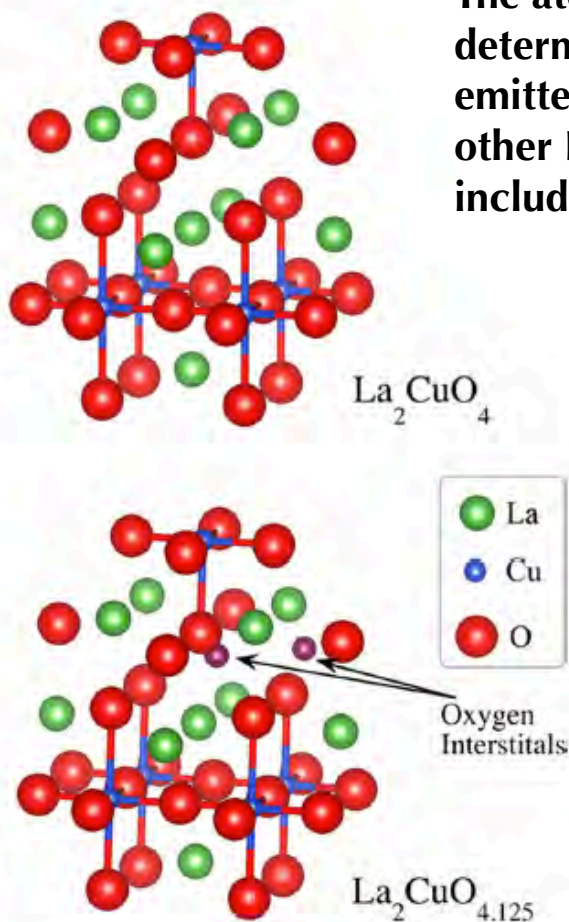
$\mu$ -XANES measurements were performed using LUCIA at SOLEIL with the high brilliance ( $1.6 \cdot 10^{11}$  ph/s/400 mA) Apple II-UH52 undulator soft X-rays source (0.8–8 keV). The polarized monochromatic photon beam was focused with a couple of KB mirrors on a  $2.5 \times 2.5 \mu\text{m}^2$  spot. At each x-z position, the lanthanum  $L_3$ -edge XANES signal was recorded in the fluorescence mode with a 4 elements Si drift Bruker detector.



**Characteristic La  $L_3$  spectrum at the binding energy of 5383 eV corresponding to the electronic transition from the 2p level to 5d empty states.**

*Frascati, October 3, 2017*

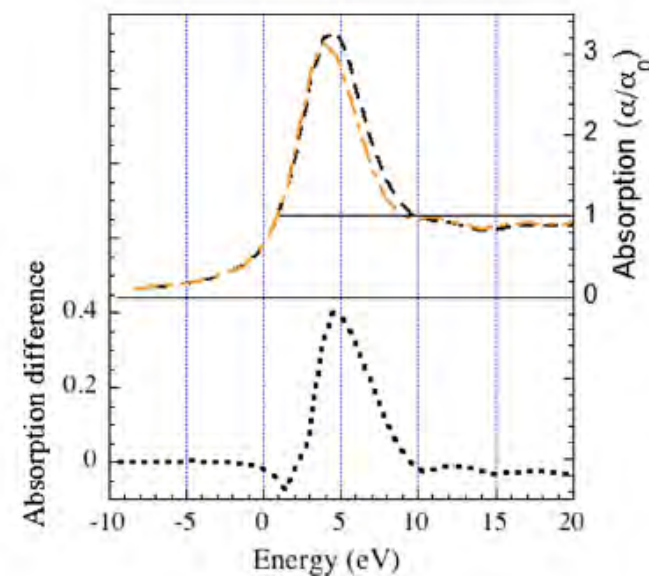
The atomic cluster ( $\sim 8 \text{ \AA}$ ) surrounding La (green) in a  $\text{La}_2\text{CuO}_{4+y}$  crystal that determines the shape resonance in the La  $L_3$ -XANES. The photoelectron wave emitted by the absorbing atoms are scattered by Cu (blue), oxygens (red) and other La ions. In oxygen rich puddles (as in the  $\text{La}_2\text{CuO}_{4.125}$  crystal) the cluster includes additional Oi (purple spheres indicated by the arrows).



FMS La  $L_3$  XANES calculation for the  $\text{La}_2\text{CuO}_4$  (brown) and the  $\text{La}_2\text{CuO}_{4.125}$  (black). In the bottom the difference between the two spectra. As shown in the lower panel, Oi induce an increase of the peak due to the first and strongest MSR.

The superimposition of two XANES spectra in the ORPs ( $\text{La}_2\text{CuO}_{4+y}$  with  $y \sim 0.125$  Oi) and in the OPP ( $\text{La}_2\text{CuO}_4$  without Oi with  $y \sim 0$ ) shows an increase by 0.4 in units of the La  $L_3$ -edge atomic absorption jump and a shift of the white-line in correspondence of the ORP.

## FMS theory



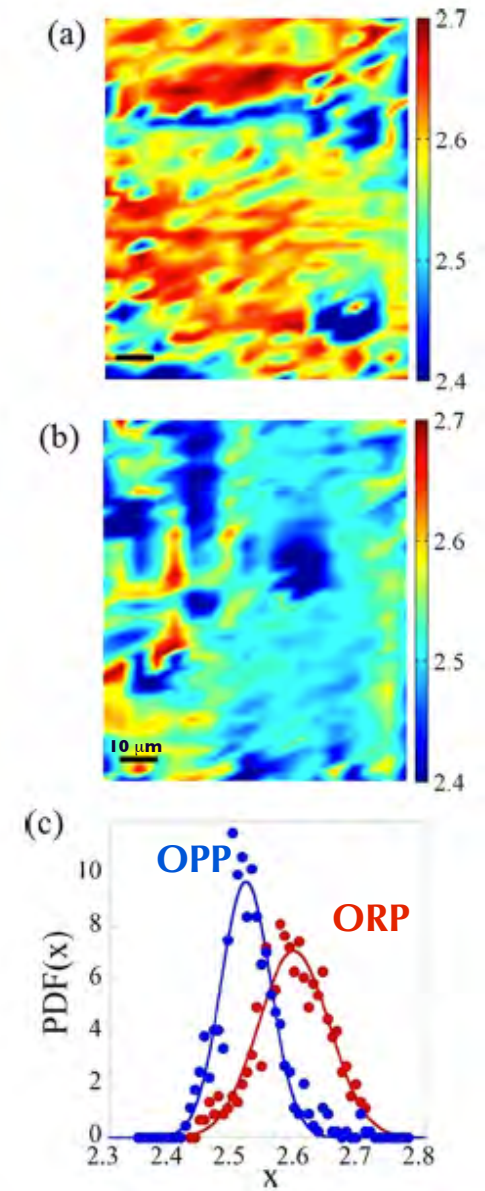


Color maps of the position dependent intensity of the MSR in the XANES spectra of the  $\text{La}_2\text{CuO}_{4.06}$  crystal. Maps are 36x18 pixels (4  $\mu\text{m}$  pixel size). The exposition time was 10 s/step. In (c) are compared the probability density function PDF(x) of the MSR of the **ORP** and **OPP** regions.

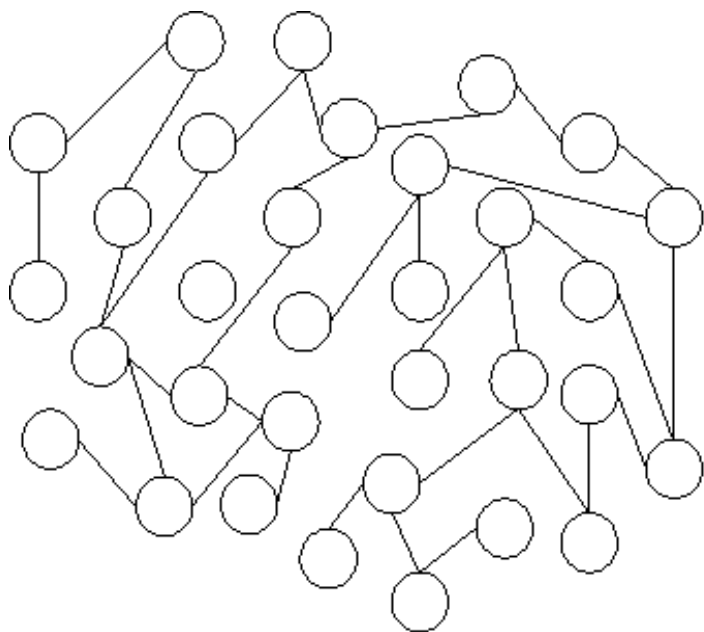
We reconstructed the intensity of the main MSR peak in the XANES spectra using a spot of 4  $\mu\text{m}$ . The spatial distribution of ORPs and OPPs is revealed looking at the MSR peak and comparing the images where different spatial distributions of ORP occur.

The statistical analysis, i.e., the distribution curves of the two regions are shown in panel (c). The x axis is the intensity of the MSR, which ranges from 2.4 to 2.7 in units of the absorption jump. The Gaussian distribution of the map (a) peaked at 2.6 while the map (b) at 2.5 [see panel (c)]. **The width of the oxygen poor distribution (OPP) ranges from 2.45 to 2.55**, while for the **oxygen rich region (ORP) it ranges from 2.52 to 2.65**, i.e., we have different distributions in different locations.

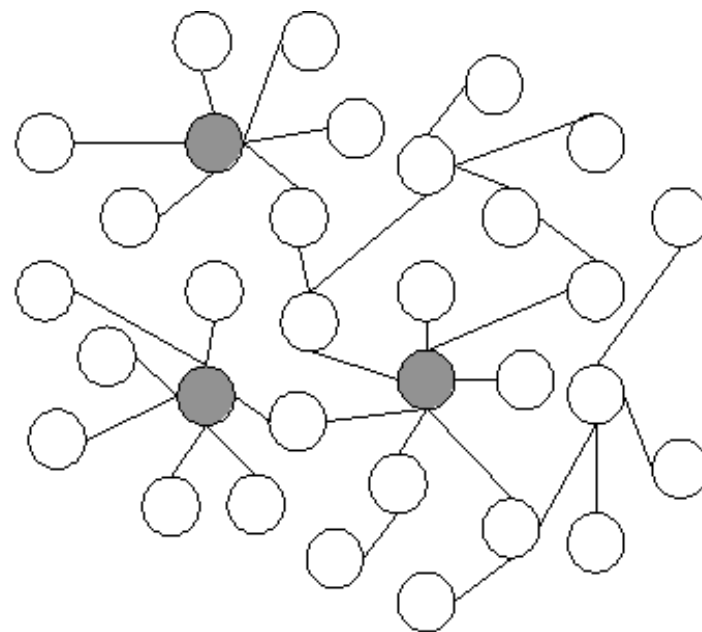
The variation of the MSR intensity is smaller respect to the calculation. Actually, as showed by scanning  $\mu\text{-XRD}$  experiments, **OPP** regions are <4  $\mu\text{m}$  and with our micro-beam it was impossible to probe regions free from **ORPs**.







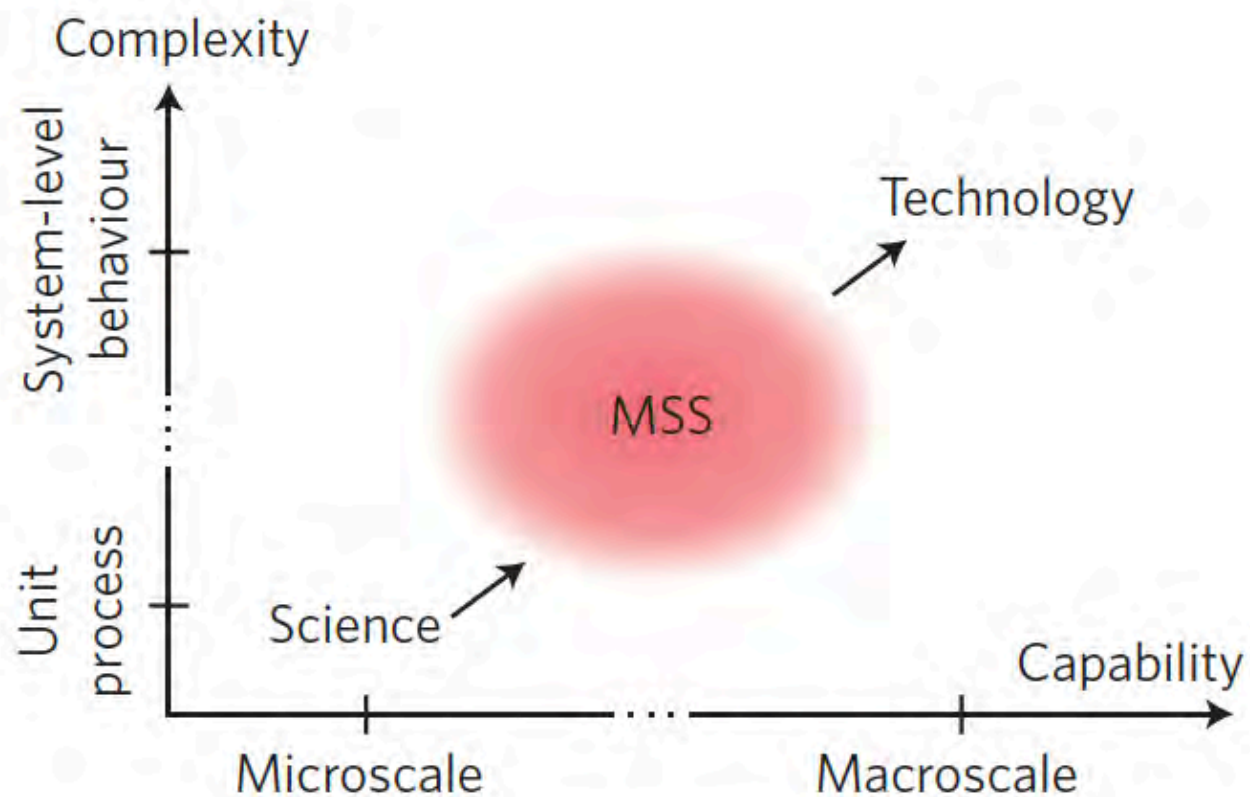
(a) Random network



(b) Scale-free network

# Conclusions

- RF applications are limited by operational conditions & materials properties, e.g., yield stress, Young's modulus, etc.;
- TM oxides show many exotic properties and tunable (percolative) conductivity mechanisms;
- Mo oxides intercalated by electron-donating cations may become good conductors, and even metallic. These mixed type conductors show both ionic and electronic conductivity;
- The advantages of cryogenic operation of metallic RF cavities at high field arise from the hardness of the material, the lowered surface dissipation and the concomitant pulsed heating;
- Spectroscopy, diffraction and hyperspectral-imaging techniques are necessary to visualise the nature of metals, alloys and multi-scale systems such as metals, alloys and TM oxides;
- These researches can be extremely useful not only to identify coatings and technology suitable to increase performances of RF cavities working at high frequencies (X-band, W-band, ...), but to many other applications where conductivity, field emission, dark current and other properties of technological interest have to be characterised and possibly tuned/optimized.





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