

The Hunt for Axions.

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Strong Case for Particles Beyond the Standard Model

- > Standard Model (SM) of particle physics describes properties of known matter and forces to a great precision

Drei Generationen
der Materie (Fermionen)

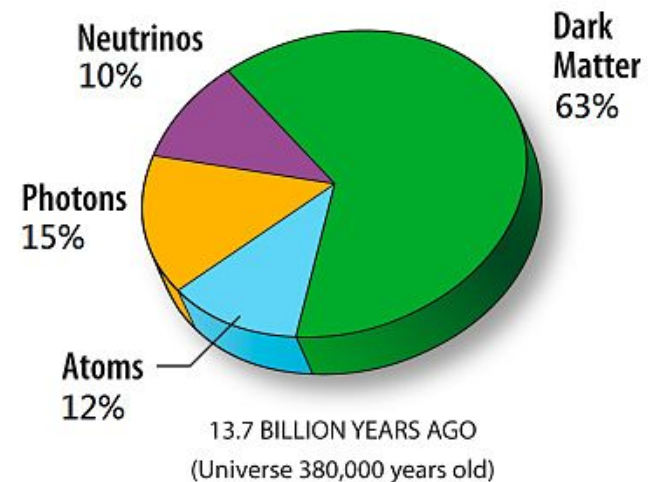
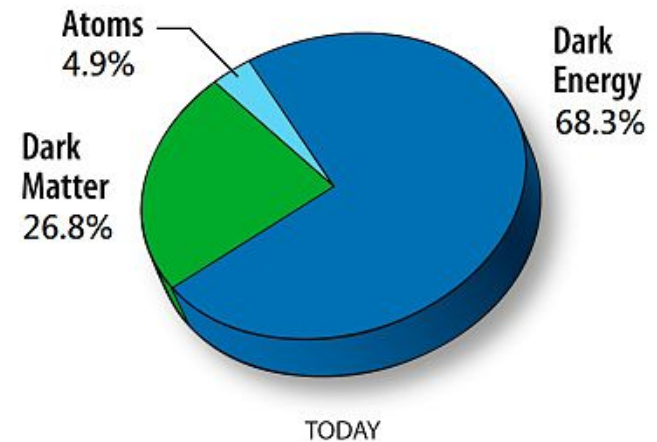
| | I | II | III | | |
|----------|--|--|---|--|-------------------------|
| Masse → | 2,3 MeV | 1,275 GeV | 173,07 GeV | 0 | 125,9 GeV |
| Ladung → | $\frac{2}{3}$ | $\frac{2}{3}$ | $\frac{2}{3}$ | 0 | 0 |
| Spin → | $\frac{1}{2}$ | $\frac{1}{2}$ | $\frac{1}{2}$ | 1 | 0 |
| Name → | u up | c charm | t top | γ Photon | H Higgs Boson |
| Quarks | 4,8 MeV $-\frac{1}{3}$ $\frac{1}{2}$ d down | 95 MeV $-\frac{1}{3}$ $\frac{1}{2}$ s strange | 4,18 GeV $-\frac{1}{3}$ $\frac{1}{2}$ b bottom | 0 0 1 g Gluon | |
| | <2 eV 0 $\frac{1}{2}$ ν_e Elektron-Neutrino | <0,19 MeV 0 $\frac{1}{2}$ ν_μ Myon-Neutrino | <18,2 MeV 0 $\frac{1}{2}$ ν_τ Tau-Neutrino | 91,2 GeV 0 1 Z⁰ Z Boson | |
| | 0,511 MeV -1 $\frac{1}{2}$ e Elektron | 105,7 MeV -1 $\frac{1}{2}$ μ Myon | 1,777 GeV -1 $\frac{1}{2}$ τ Tau | 80,4 GeV ±1 1 W[±] W Boson | Eichbosonen |
| Leptonen | | | | | |

[wikipedia]



Strong Case for Particles Beyond the Standard Model

- > Standard Model (SM) of particle physics describes properties of known matter and forces to a great precision
- > SM not a complete and fundamental theory:
 - No explanation of the origin of dark energy and dark matter (DM)

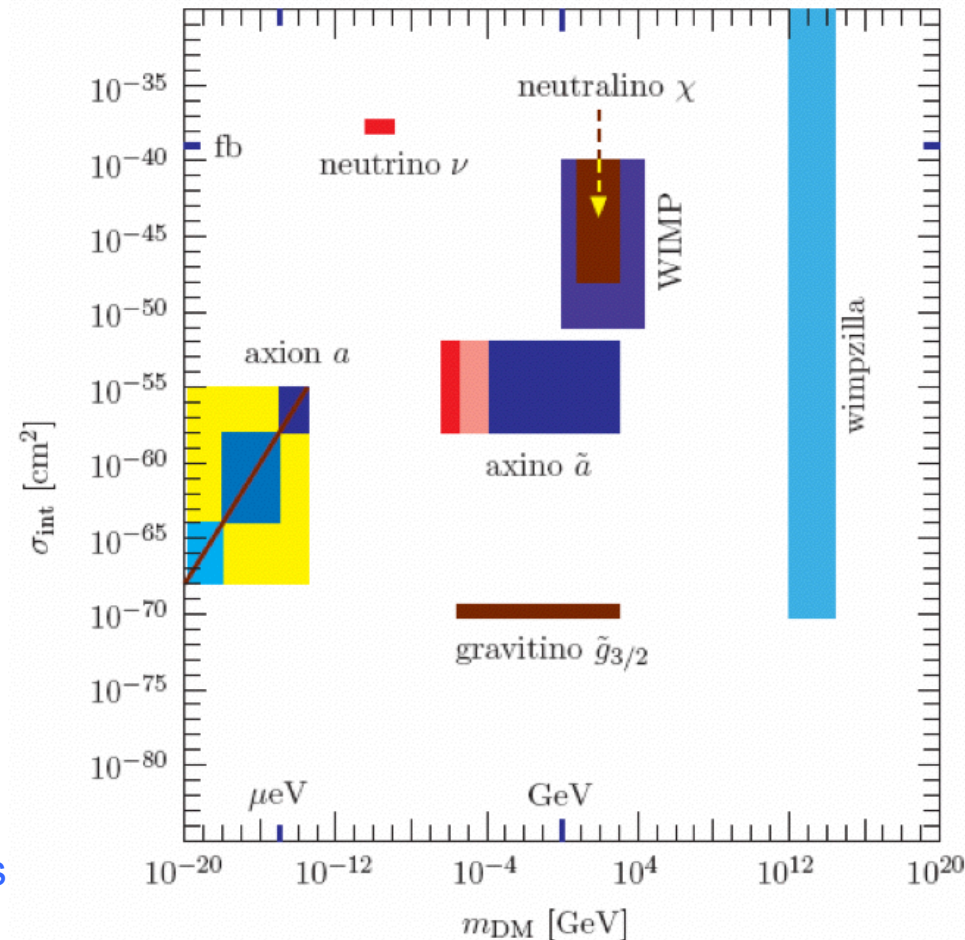


[wikipedia]



Strong Case for Particles Beyond the Standard Model

- > Standard Model (SM) of particle physics describes properties of known matter and forces to a great precision
- > SM not a complete and fundamental theory:
 - No explanation of the origin of dark energy and dark matter (DM)
- > Plenitude of DM candidates, notably:
 - Weakly Interacting Massive Particles (**WIMPs**), such as **neutralinos**
 - Very Weakly Interacting Slim (=ultra-light) Particles (**WISPs**), such as **axions**
- > Stand out because of their convincing physics case and the variety of experimental probes



[Kim, Carosi 10]



Natural Candidates for WISPs: Nambu-Goldstone Bosons

- > Nambu-Goldstone boson arising from breaking of global, e.g. U(1), symmetry
- > Hidden Higgs field:

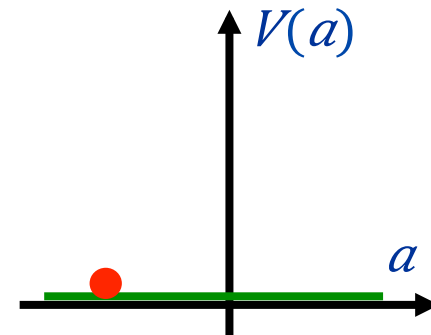
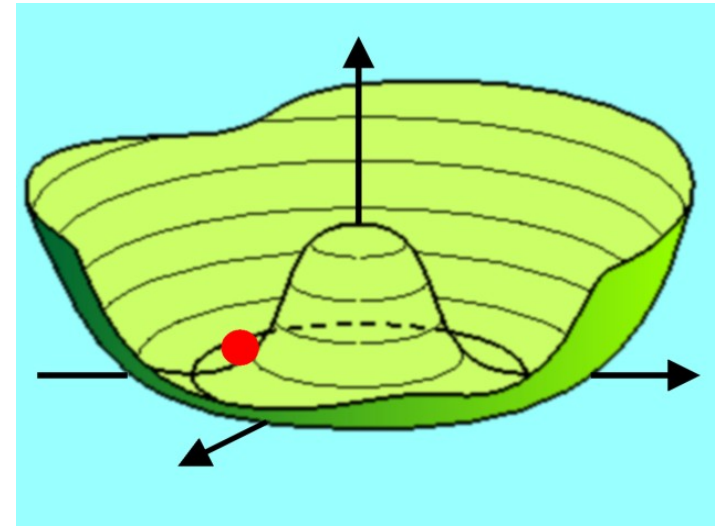
$$H_h(x) = \frac{1}{\sqrt{2}} [v_h + h_h(x)] e^{ia(x)/v_h}$$

Massive modulus, massless phase:

$$m_{h_h} \sim v_h \quad m_a = 0$$

- > Interactions with SM particles small, if scale of symmetry breaking much larger than SM Higgs vacuum expectation value,

$$v_h \gg v = 246 \text{ GeV}$$

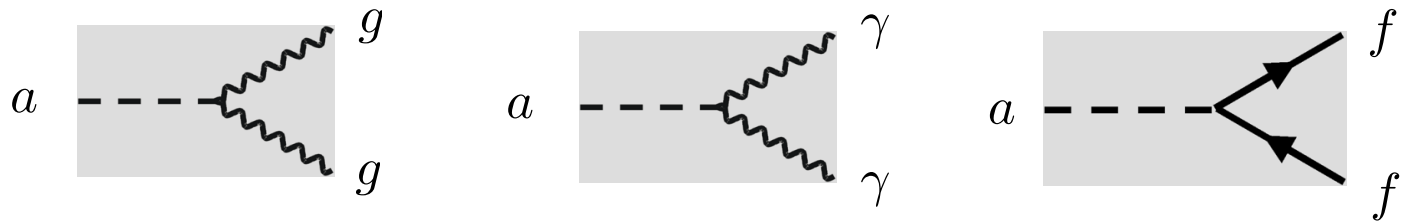


[Raffelt]

Natural Candidates for WISPs: Nambu-Goldstone Bosons

- > Couplings to SM suppressed by powers of $f_a \sim v_h \gg v = 246 \text{ GeV}$

$$\mathcal{L} \supset -\frac{\alpha_s}{8\pi} \frac{C_{ag}}{f_a} a G_{\mu\nu}^b \tilde{G}^{b,\mu\nu} - \frac{\alpha}{8\pi} \frac{C_{a\gamma}}{f_a} a F_{\mu\nu} \tilde{F}^{\mu\nu} + \frac{1}{2} \frac{C_{af}}{f_a} \partial_\mu a \bar{\psi}_f \gamma^\mu \gamma_5 \psi_f$$



- > Coefficients C_{ag} , $C_{a\gamma}$ determined by loops over particles charged under hidden U(1). C_{af} can arise at tree or loop level.
- > Global symmetry not necessarily exact: Nambu-Goldstone boson will acquire a small mass vanishing in the limit that the global hidden symmetry is exact
 - Example in SM: Pions pseudo Nambu-Goldstone bosons of chiral symmetry breaking in QCD ... mass vanishes for vanishing quark masses

Natural Candidates for WISPs: Nambu-Goldstone Bosons

- Often, there is more than one global symmetry and therefore more than one Nambu-Goldstone boson
 - Global lepton number symmetry: **Majoron** [Chikashige et al. 78; Gelmini, Roncadelli 80]
 - Global family symmetry: **Familon** [Wilczek 82; Berezhiani, Khlopov 90]

$$\mathcal{L} \supset -\frac{\alpha_s}{8\pi} \frac{C'_{ig}}{f_{a'_i}} a'_i G_{\mu\nu}^b \tilde{G}^{b,\mu\nu} - \frac{\alpha}{8\pi} \frac{C'_{i\gamma}}{f_{a'_i}} a'_i F_{\mu\nu} \tilde{F}^{\mu\nu} + \frac{1}{2} \frac{C'_{a'_i f}}{f_{a'_i}} \partial_\mu a'_i \bar{\psi}_f \gamma^\mu \gamma_5 \psi_f$$

- The particle corresponding to the linear combination

$$\frac{A(x)}{f_A} \equiv \frac{C'_{ig}}{f_{a'_i}} a'_i(x)$$



is called **Axion** (= laundry detergent): it cleans up the strong CP problem

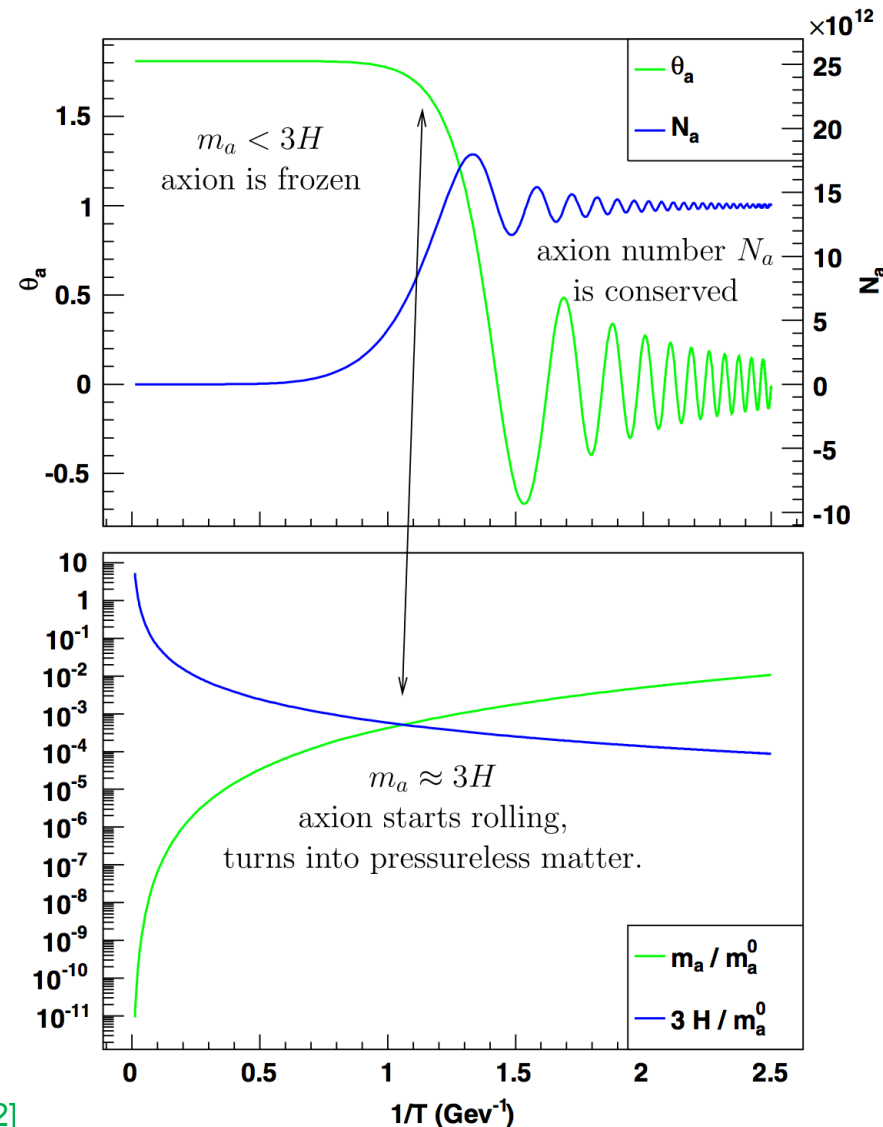
[Peccei, Quinn 77; Weinberg 78; Wilczek 78]

- Particle excitations of the fields orthogonal to the axion field are called **Axion-Like-Particles (ALPs)**
- String theory suggests a plenitude of ALPs [Witten; Arvanitaki et al., Cicoli, Goodsell, AR]



Axion/ALP Dark Matter?

- > In early universe, axion/ALP frozen at random initial value
- > Later, field feels pull of mass towards zero and oscillates around it
- > Spatially uniform oscillating classical field = coherent state of many, extremely non-relativistic particles = cold dark matter (CDM)



[Preskill et al 83; Abbott, Sikivie 83; Dine, Fischler 83,....; Arias et al. 12]

[Wantz, Shellard 09]



Axion/ALP Dark Matter?

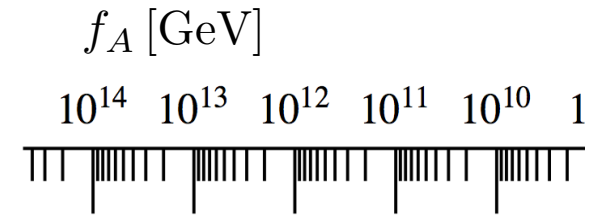
- > In early universe, WISP (axion/ALP) frozen at random initial value
- > Later, field feels pull of mass towards zero and oscillates around it
- > Spatially uniform oscillating classical field = coherent state of many, extremely non-relativistic particles = cold dark matter (CDM)
- > Energy density proportional to initial field amplitude squared,

$$\rho_a = n_a m_a \sim f_a^2$$

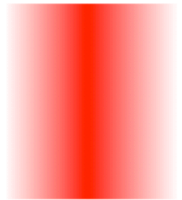
- > Axion/ALP CDM prefers:

$$f_a \gtrsim 10^{10 \div 12} \text{ GeV}$$

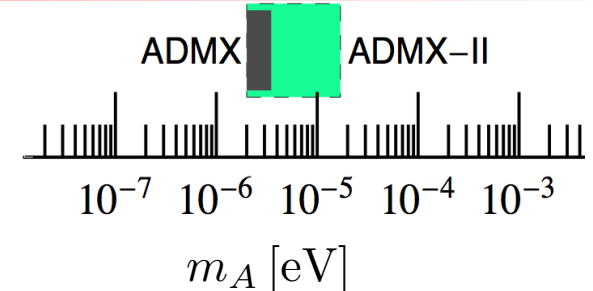
[Preskill et al 83; Abbott, Sikivie 83; Dine, Fischler 83,....; Arias et al. 12]



postinflation PQ
(realignment+cosmic strings+DWs)



preinflation PQ
(only realignment)



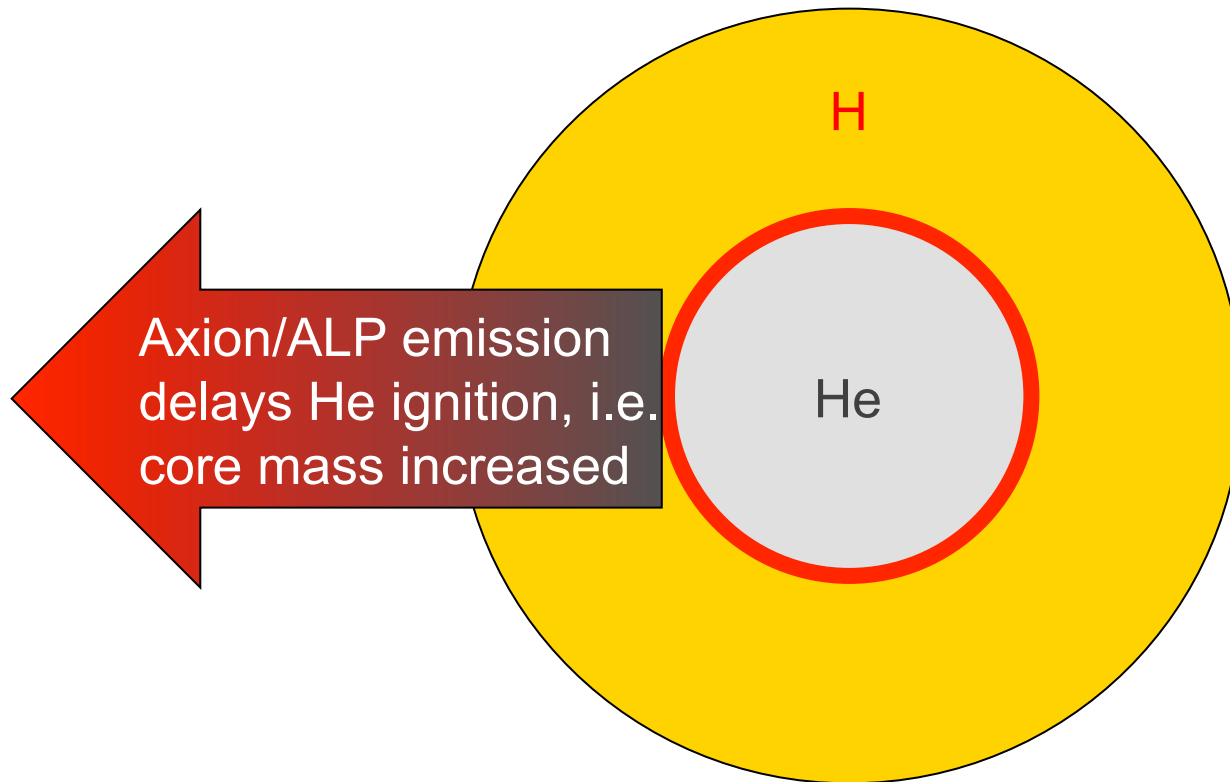
[adapted by from Essig et al. 1311.0029]



Axion/ALP Energy Losses of Stars in Globular Clusters?

- > **Red Giants** (RGs) in globular clusters mildly prefer additional energy losses due to axion/ALP emission via Bremsstrahlung $e + Ze \rightarrow Ze + e + a$

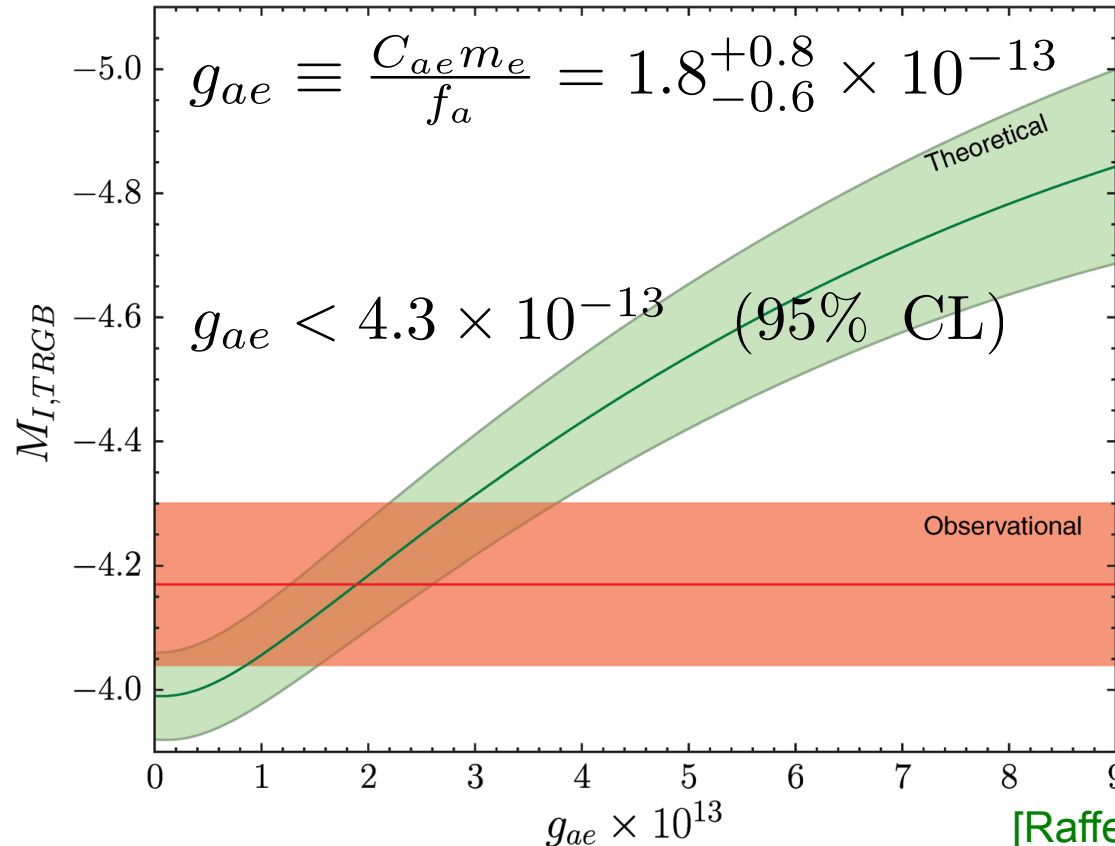
[Viaux et al. 13]



Red Giant

Axion/ALP Energy Losses of Stars in Globular Clusters?

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[Viaux et al. 13]

[Raffelt 14]

- Mild hints of anomalous energy loss of White Dwarfs (WDs) could also be explained by same parameter values

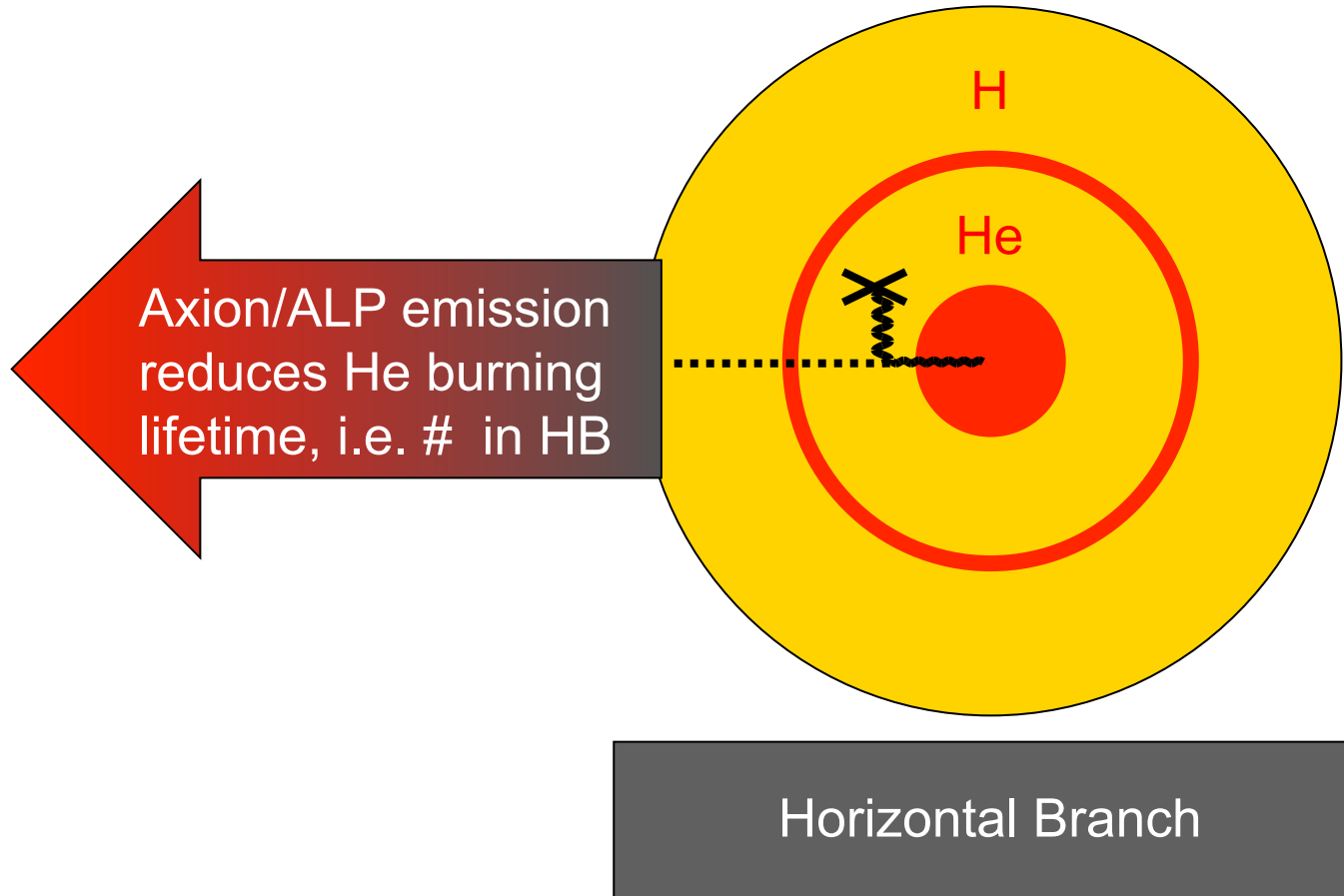
[Isern et al.]



Axion/ALP Energy Losses of Stars in Globular Clusters?

- > **Horizontal Branch** (HB) stars in globular clusters mildly prefer additional energy losses due to axion/ALP emission via Primakoff $\gamma + Ze \rightarrow Ze + a$

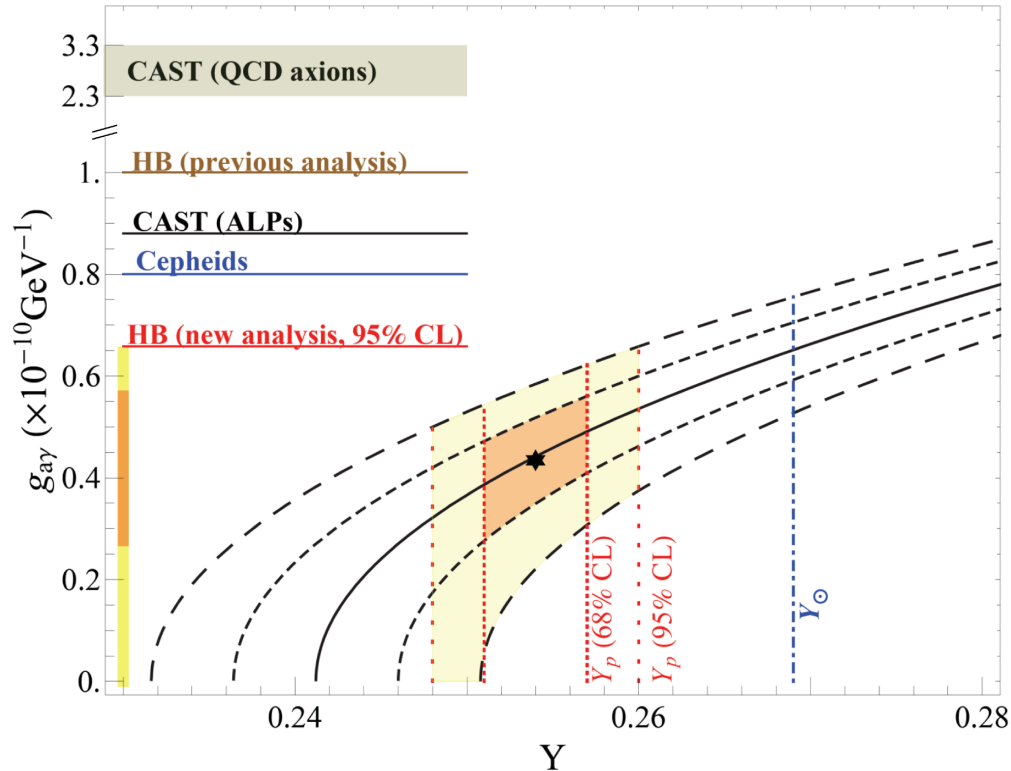
[Ayala et al. 14]



Axion/ALP Energy Losses of Stars in Globular Clusters?

- **Horizontal Branch** (HB) stars in globular clusters mildly prefer additional energy losses due to axion/ALP emission via Primakoff $\gamma + Ze \rightarrow Ze + a$

[Ayala et al. 14]



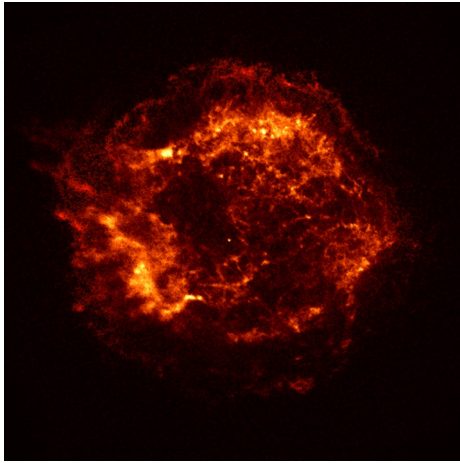
$$g_{a\gamma} \equiv \frac{\alpha C_{a\gamma}}{2\pi f_a} = 0.45^{+0.12}_{-0.16} \times 10^{-10} \text{ GeV}^{-1}$$

$$g_{a\gamma} < 0.66 \times 10^{-10} \text{ GeV}^{-1} \quad (95\% \text{ CL})$$



Axion/ALP Energy Losses of Neutron Star in Cas A?

> Neutron star in Cas A:

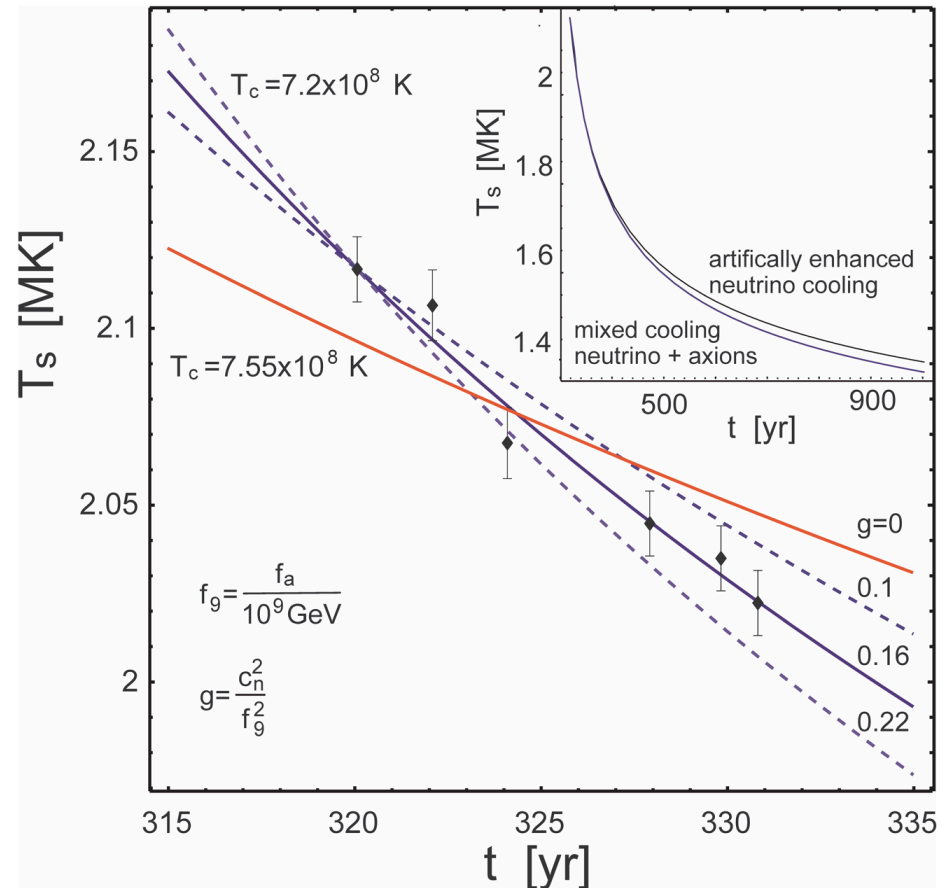


- Measured surface temperature over 10 years reveals unusually fast cooling rate
- Hint on extra cooling by axion/ALP due to nucleon bremsstrahlung



- Required coupling to neutron:

$$g_{an} \equiv \frac{C_{an} m_n}{f_a} \sim 4 \times 10^{-10}$$

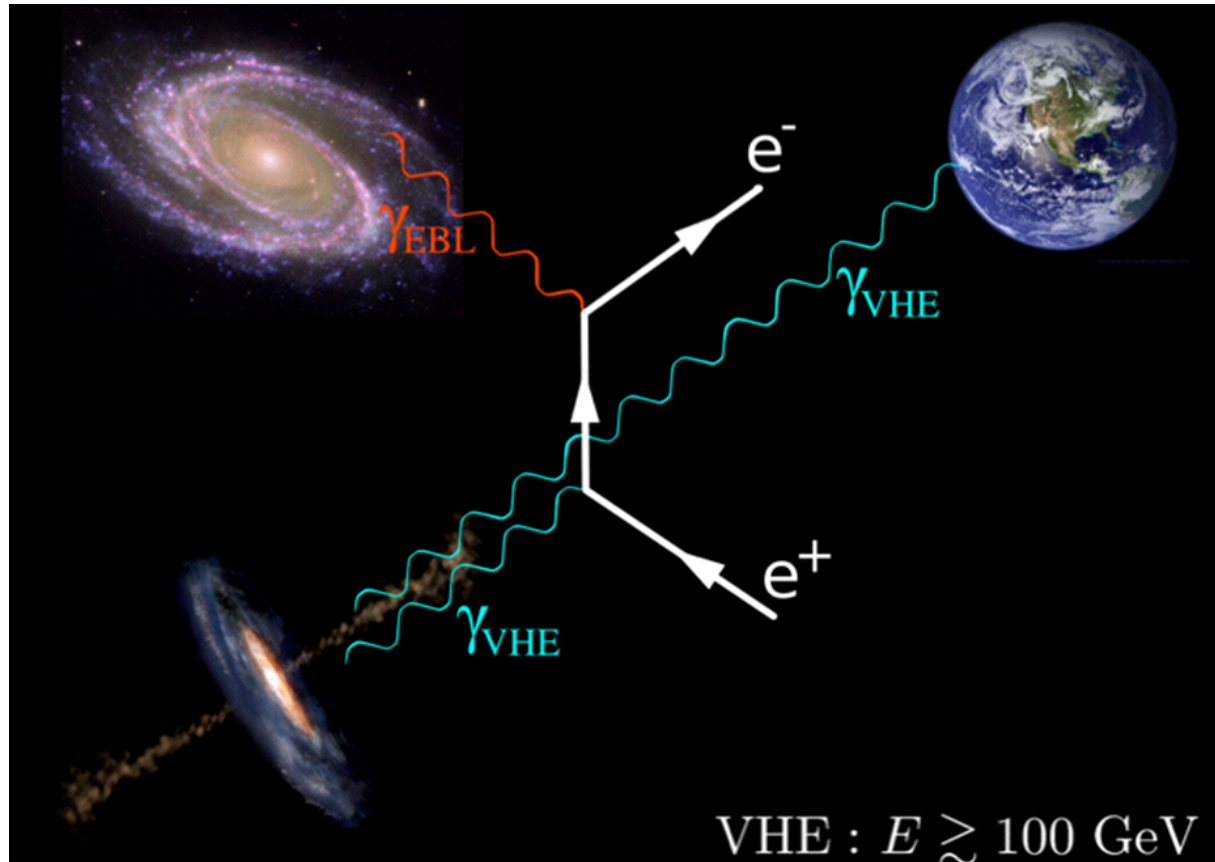


[Leinson 14]



Photon – ALP Conversion in Cosmic Magnetic Fields?

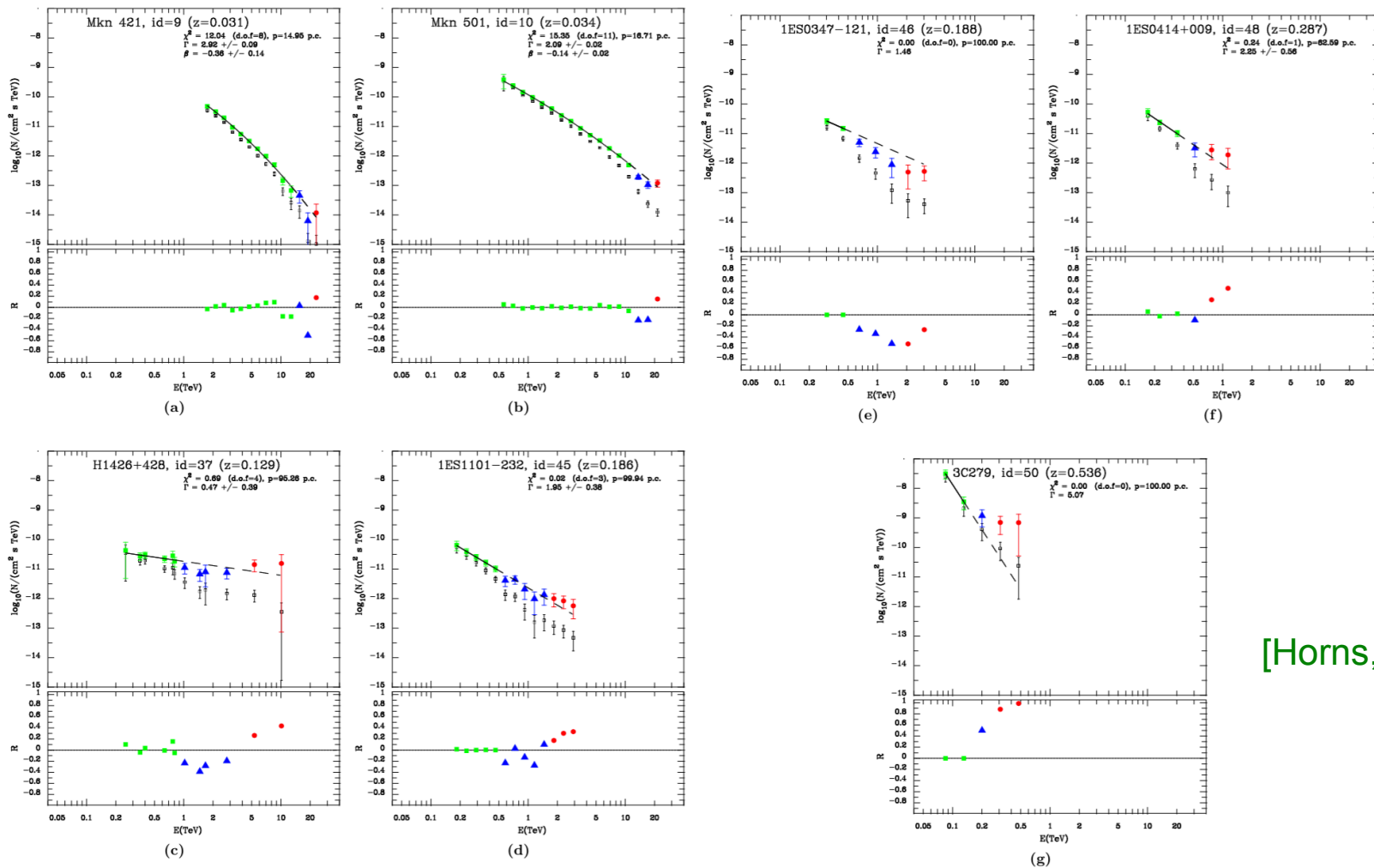
- Gamma ray spectra from distant AGNs should show an energy and red-shift dependent exponential attenuation, due to pair production at Extragalactic Background Light (EBL)



[Manuel Meyer 12]

Photon – ALP Conversion in Cosmic Magnetic Fields?

- Indication of anomalous gamma transparency: attenuation observed by **IACT** and Fermi-LAT too small [Aharonian et al. 07; de Angelis, Roncadelli et al. 07; ...; Horns, Meyer 12; ...; Rubtsov, Troitsky 14]

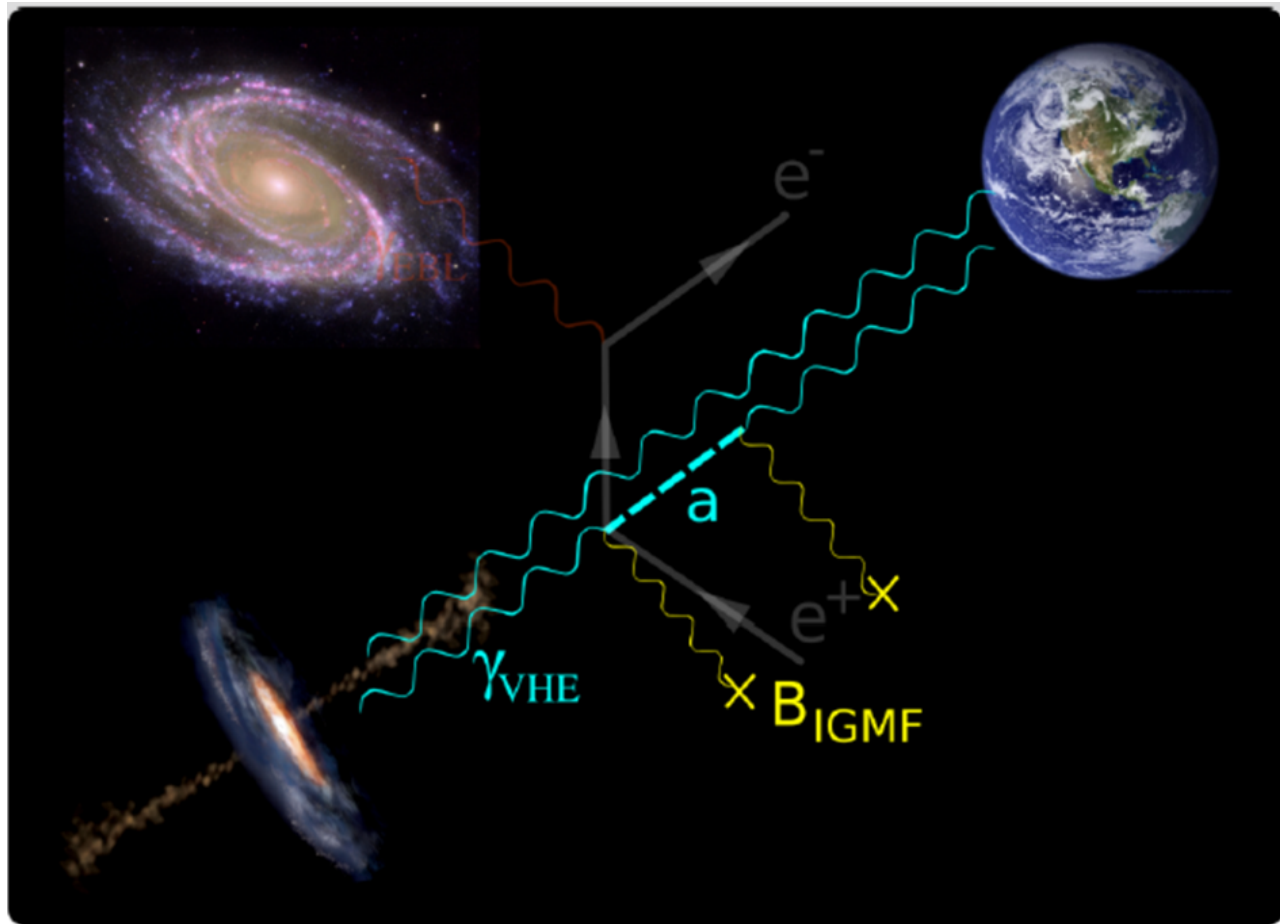


[Horns, Meyer 12]



Photon – ALP Conversion in Cosmic Magnetic Fields?

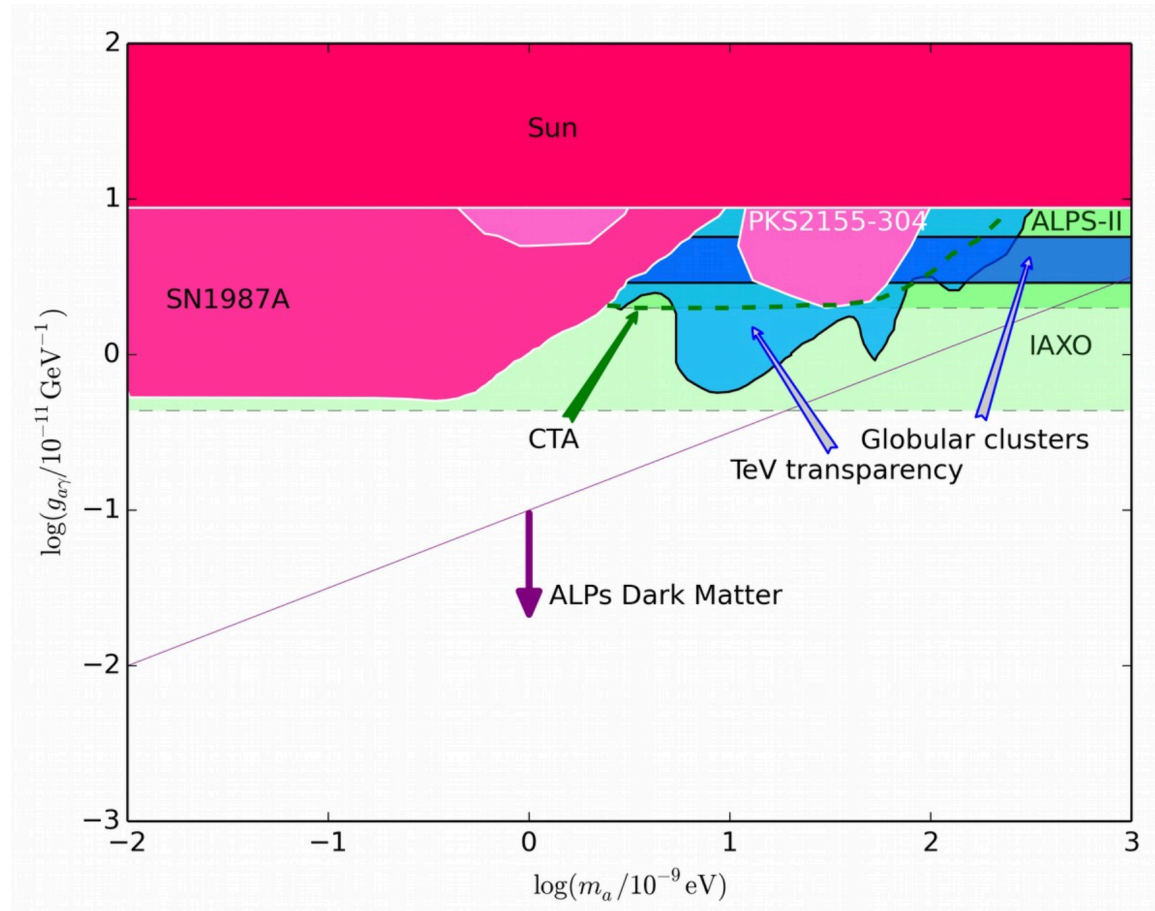
- Possible explanation: photon \leftrightarrow ALP conversions in magnetic fields
[De Angelis et al 07; Simet et al 08; Sanchez-Conde et al 09; Meyer,Horns,Raue 13]



[Manuel Meyer 12]

Photon – ALP Conversion in Cosmic Magnetic Fields?

- Possible explanation: photon \leftrightarrow ALP conversions in magnetic fields
[De Angelis et al 07; Simet et al 08; Sanchez-Conde et al 09; Meyer,Horns,Raue 13]

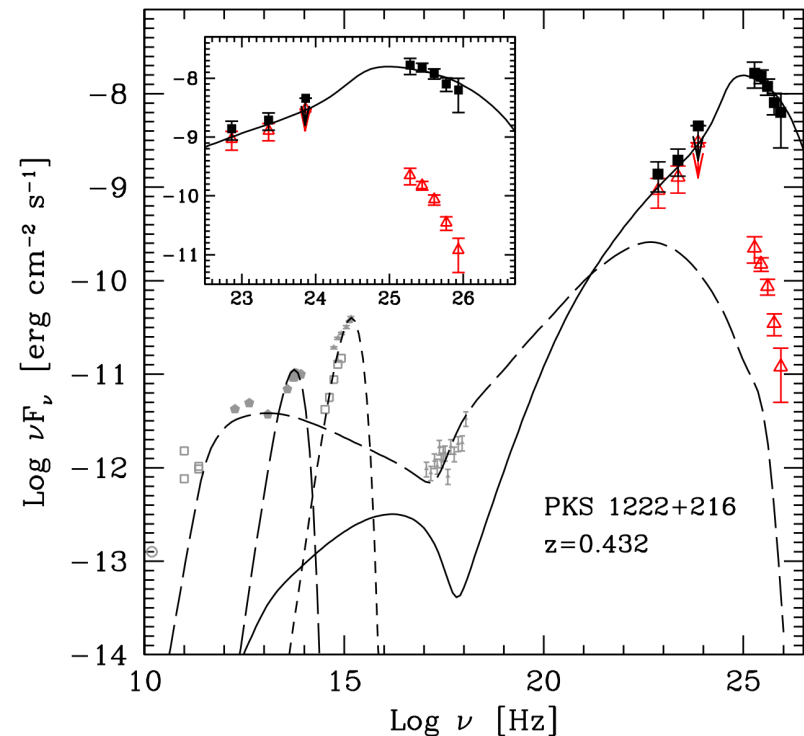
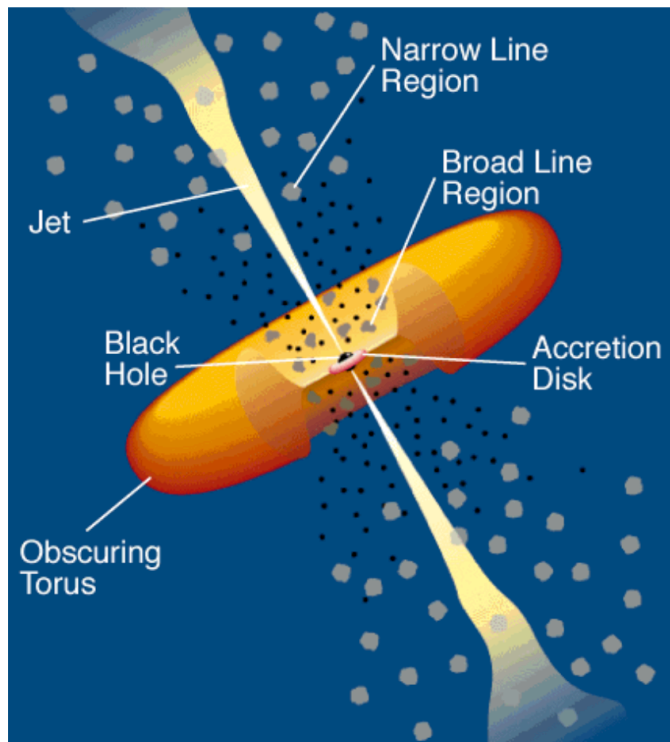


[Horns 15]

- Required photon coupling overlaps with preferred region from HBs in GCs

Photon – ALP Conversion in Cosmic Magnetic Fields?

- Photon-ALP conversion with a strength $g_{a\gamma} \sim \text{few} \times 10^{-11} \text{ GeV}^{-1}$ explains also puzzling observation of VHE photons from Flat Spectrum Radio Quasars (FSRQs) [Tavecchio, Roncadelli, Galanti, Bonnoli 12]
- Pair production on UV photons in broad line region should prevent photons produced by the central engine to escape



Summary of Astrophysical Hints for Axion/ALPs

> Symmetry breaking scale inferred from astrophysical hints:

1. RGs + WDs: $f_a = 3 \times 10^9 \text{ GeV } C_{ae} \left(\frac{2 \times 10^{-13}}{g_{ae}} \right)$
2. n star in Cas A: $f_a = 2 \times 10^9 \text{ GeV } C_{an} \left(\frac{4 \times 10^{-10}}{g_{an}} \right)$
3. HB stars + AGN spectra: $f_a = 2 \times 10^7 \text{ GeV } C_{a\gamma} \left(\frac{5 \times 10^{-11} \text{ GeV}^{-1}}{g_{a\gamma}} \right)$

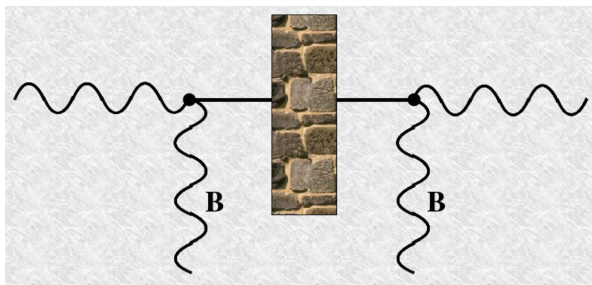
> Astrophysical hints can be explained by

- **ALP** with $f_a \sim 10^7 \text{ GeV}$, $m_a \lesssim 0.1 \text{ } \mu\text{eV}$, $C_{a\gamma} \sim 1$, $C_{ae} \sim C_{an} \sim 10^{-2}$
- **Axion** with $f_A \sim 10^9 \text{ GeV}$, $C_{An} \sim C_{A\gamma} \sim C_{Ae} \sim 1$

plus

ALP with $f_a \sim 10^7 \text{ GeV}$, $m_a \lesssim 0.1 \text{ } \mu\text{eV}$, $C_{a\gamma} \sim 1$, $C_{ae} \sim C_{an} \ll 10^{-2}$

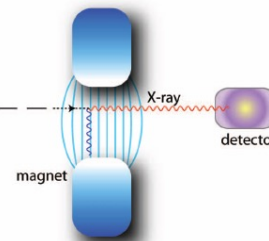
> In reach of upcoming generation of terrestrial experiments:



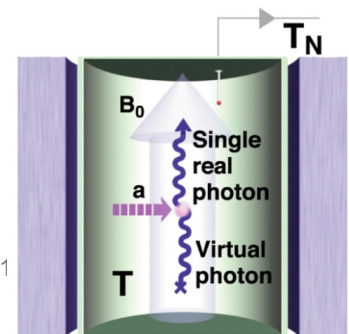
Palazzo

Sun

axion



I, 3 March 201



Axion/ALP Experiments Worldwide

An incomplete selection of (mostly) small-scale experiments:

| Experiment | Type | Location | Status |
|------------|------------------------------|----------|-------------|
| ALPS II | Light-shining-through-a-wall | DESY | preparation |
| CROWS | | CERN | finished |
| OSQAR | | CERN | running |
| REAPR | | FNAL | proposed |
| CAST | Helioscopes | CERN | running |
| IAXO | | ? | proposed |
| SUMICO | | Tokyo | running |
| ADMX | Haloscopes | Seattle | running |
| CASPEr | | Mainz | preparation |
| QUAX | | Legnaro | preparation |

[adapted from Axel Lindner `14]



Light-shining-through-a-wall Searches

- Any Light Particle Search (ALPS) at DESY (in coll. with LZH, AEI)

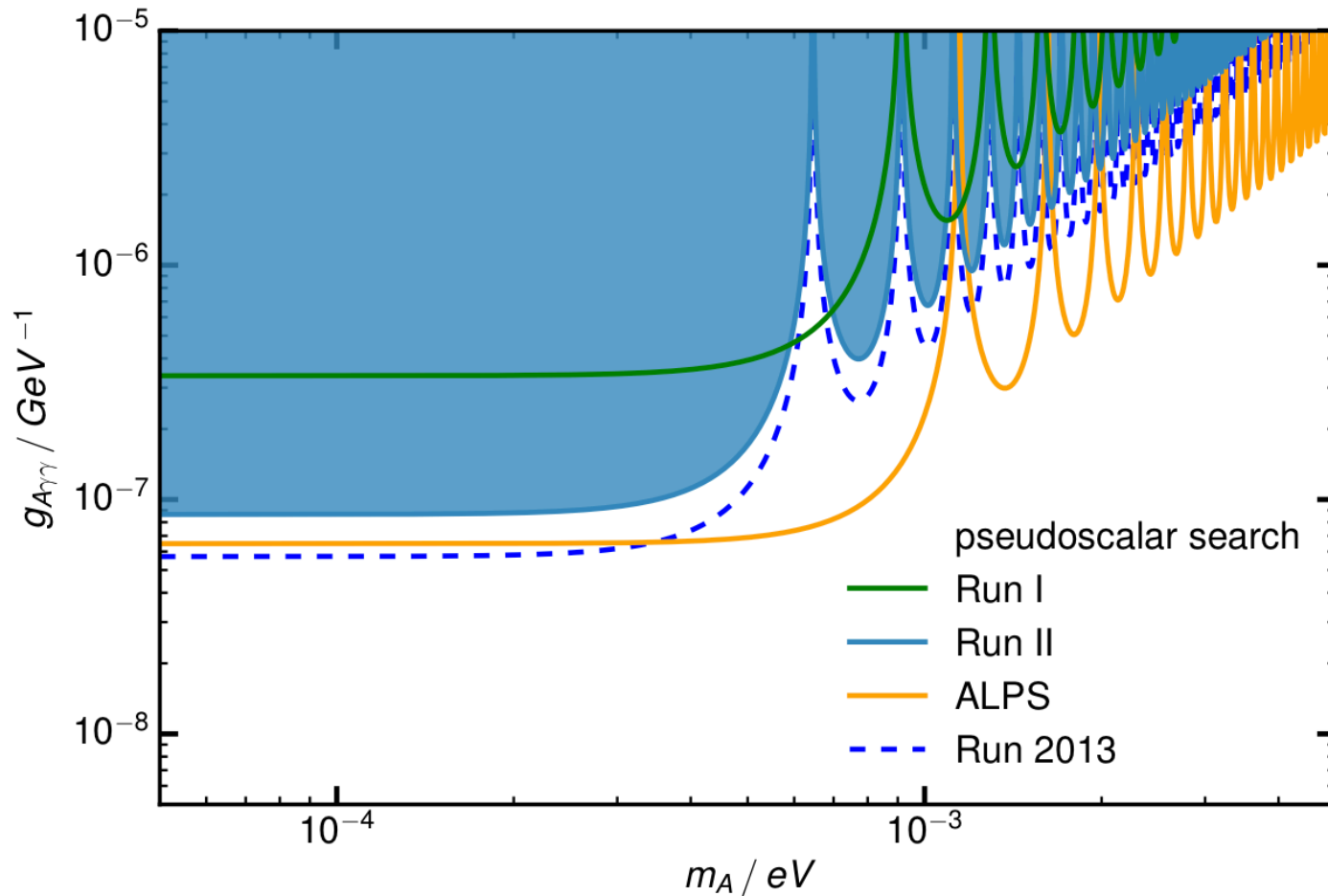


$$P(a \leftrightarrow \gamma) = 4 \frac{(g_{a\gamma} \omega B)^2}{m_a^4} \sin^2 \left(\frac{m_a^2}{4\omega} L_B \right)$$

[Anselm 85; van Bibber et al. 87]

Light-shining-through-a-wall Searches

- Currently best limits from LSW: **ALPS** (DESY) and **OSQAR** (CERN)

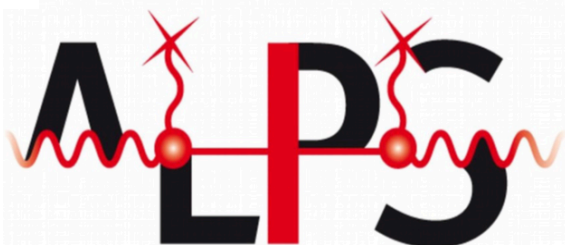
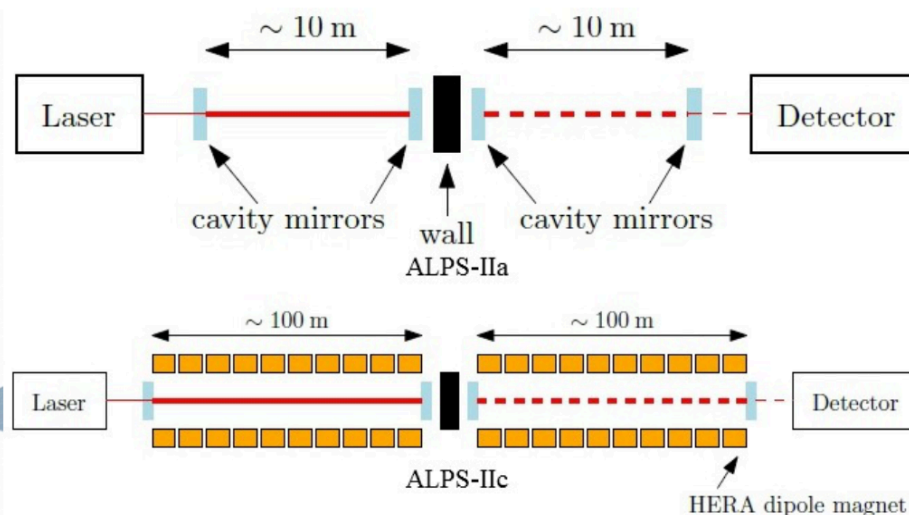
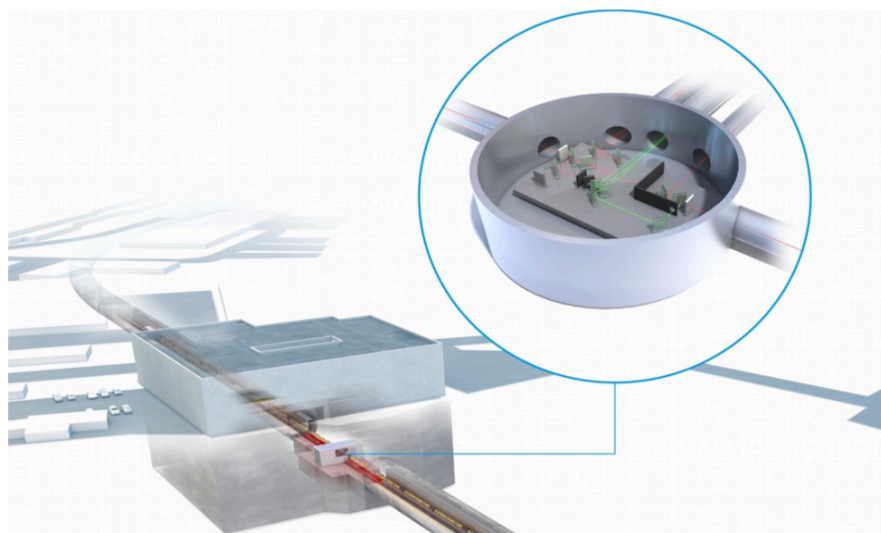


[Balou et al. 14]

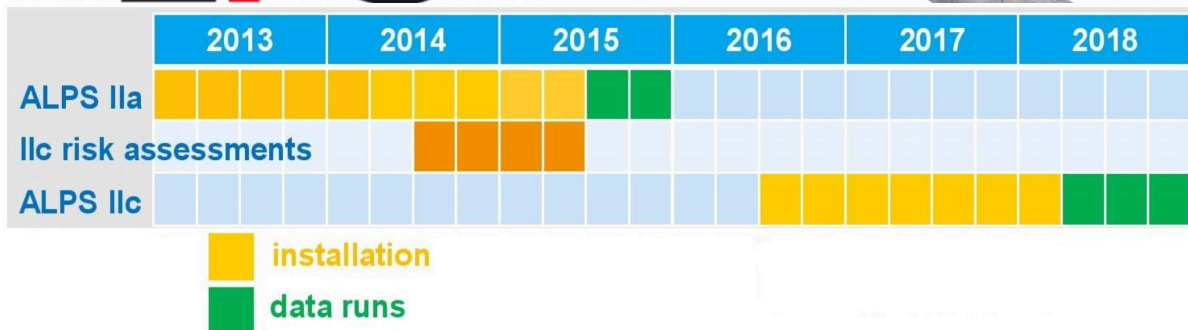


Light-shining-through-a-wall Searches

> **ALPS II** in prepar. at DESY (in collaboration with UHH, AEI, U Mainz)



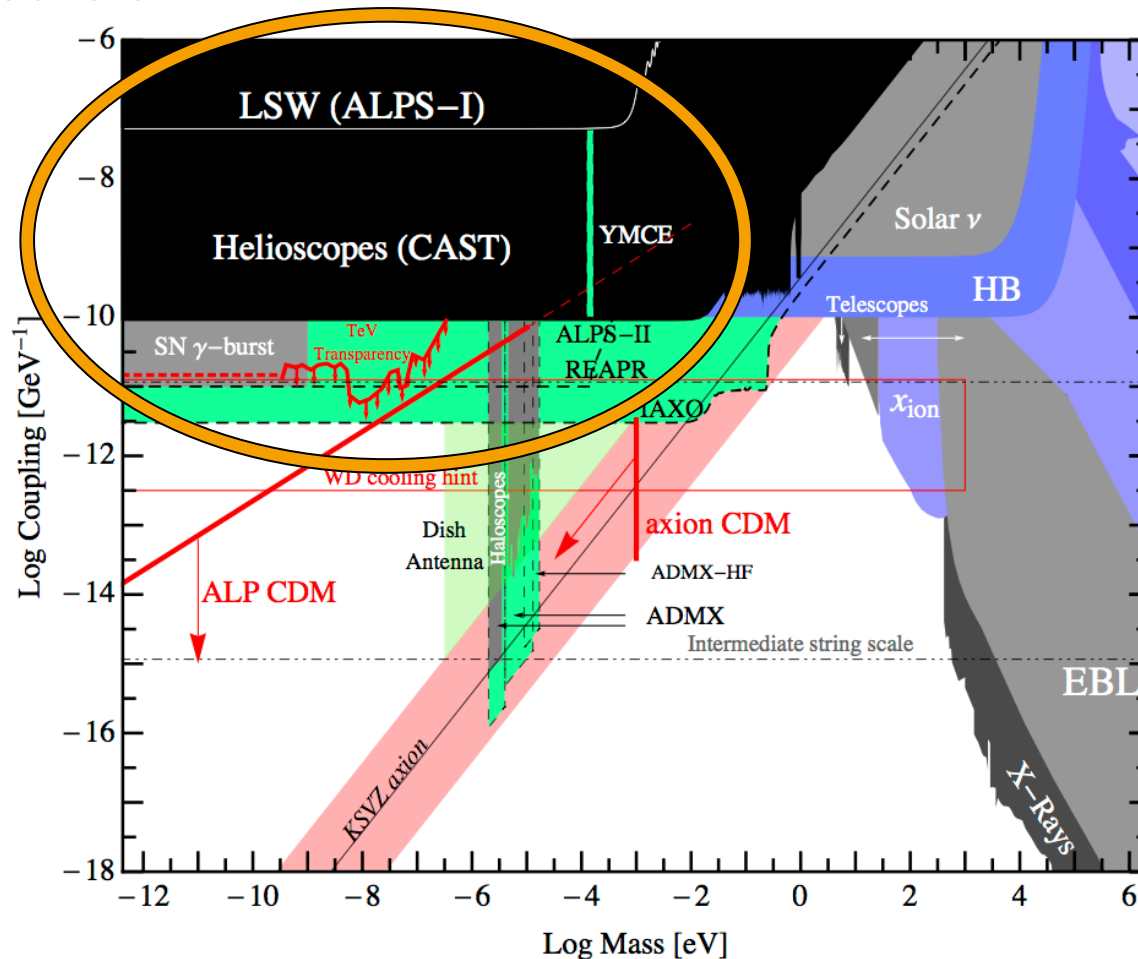
| Parameter | Scaling | ALPS I | ALPS IIc | Sens. gain |
|--|---|-------------------------|---------------------------|------------|
| Effective laser power P_{laser} | $g_{a\gamma} \propto P_{\text{laser}}^{-1/4}$ | 1 kW | 150 kW | 3.5 |
| Rel. photon number flux n_γ | $g_{a\gamma} \propto n_\gamma^{-1/4}$ | 1 (532 nm) | 2 (1064 nm) | 1.2 |
| Power built up in RC P_{RC} | $g_{a\gamma} \propto P_{\text{reg}}^{-1/4}$ | 1 | 40,000 | 14 |
| BL (before& after the wall) | $g_{a\gamma} \propto (BL)^{-1}$ | 22 Tm | 468 Tm | 21 |
| Detector efficiency QE | $g_{a\gamma} \propto QE^{-1/4}$ | 0.9 | 0.75 | 0.96 |
| Detector noise DC | $g_{a\gamma} \propto DC^{1/8}$ | 0.0018 s^{-1} | 0.000001 s^{-1} | 2.6 |
| Combined improvements | | | | 3082 |



[Bähre et al (ALPS II TDR) 13]

Light-shining-through-a-wall Searches

- Crucial test of ALP explanation of excessive HB star energy loss and AGN spectra at VHE



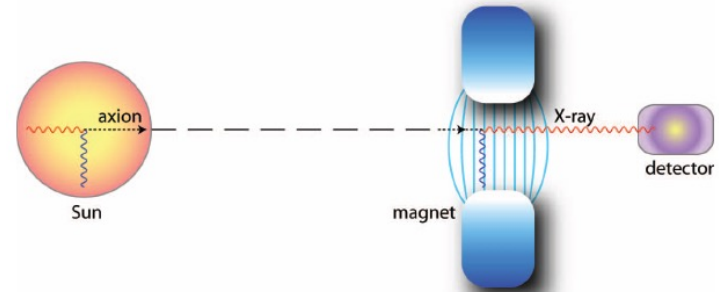
[Essig et al. 1311.0029]

Helioscope Searches

➤ Most sensitive until now: CERN Axion Solar Telescope (CAST)

- Superconducting LHC dipole magnet
- X-ray detectors

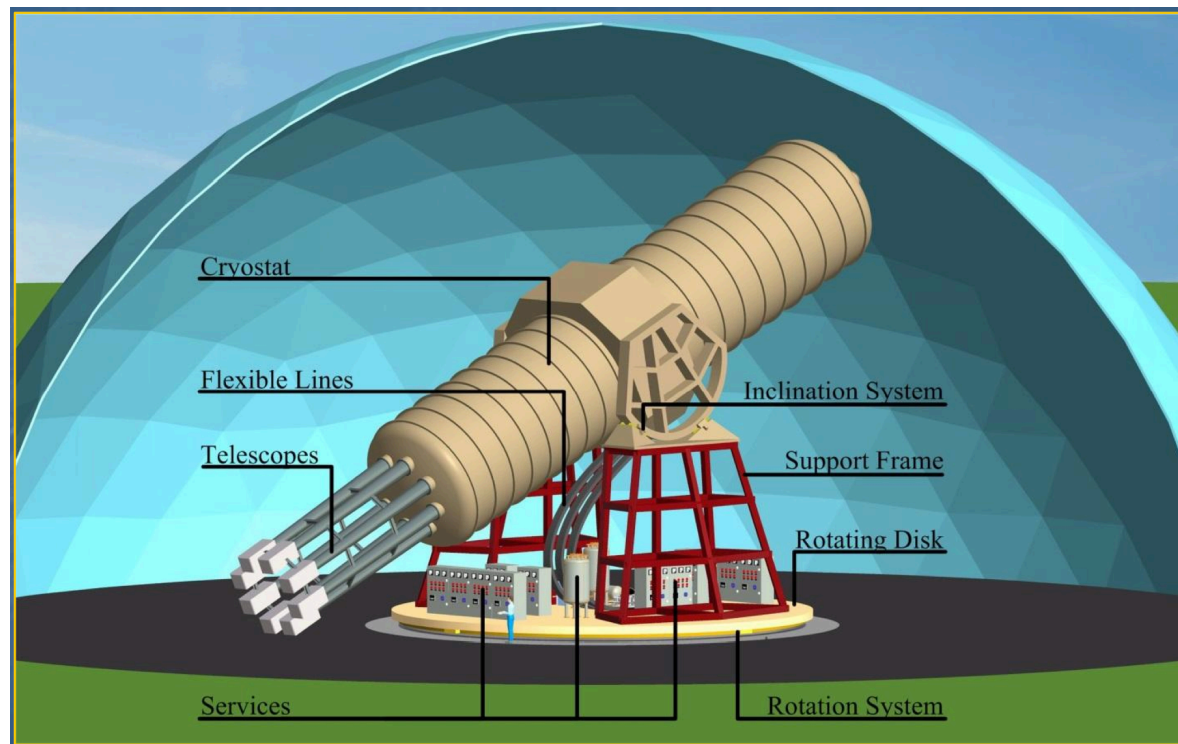
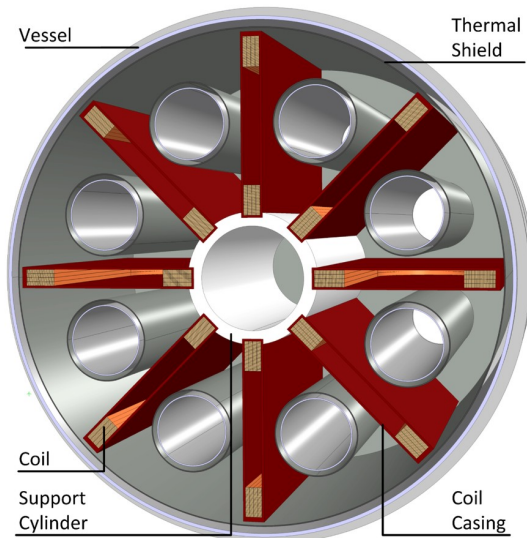
$$P(a \leftrightarrow \gamma) = 4 \frac{(g_{a\gamma} \omega B)^2}{m_a^4} \sin^2 \left(\frac{m_a^2}{4\omega} L_B \right)$$



Helioscope Searches

➤ Proposed successor: [International Axion Observatory \(IAXO\)](#)

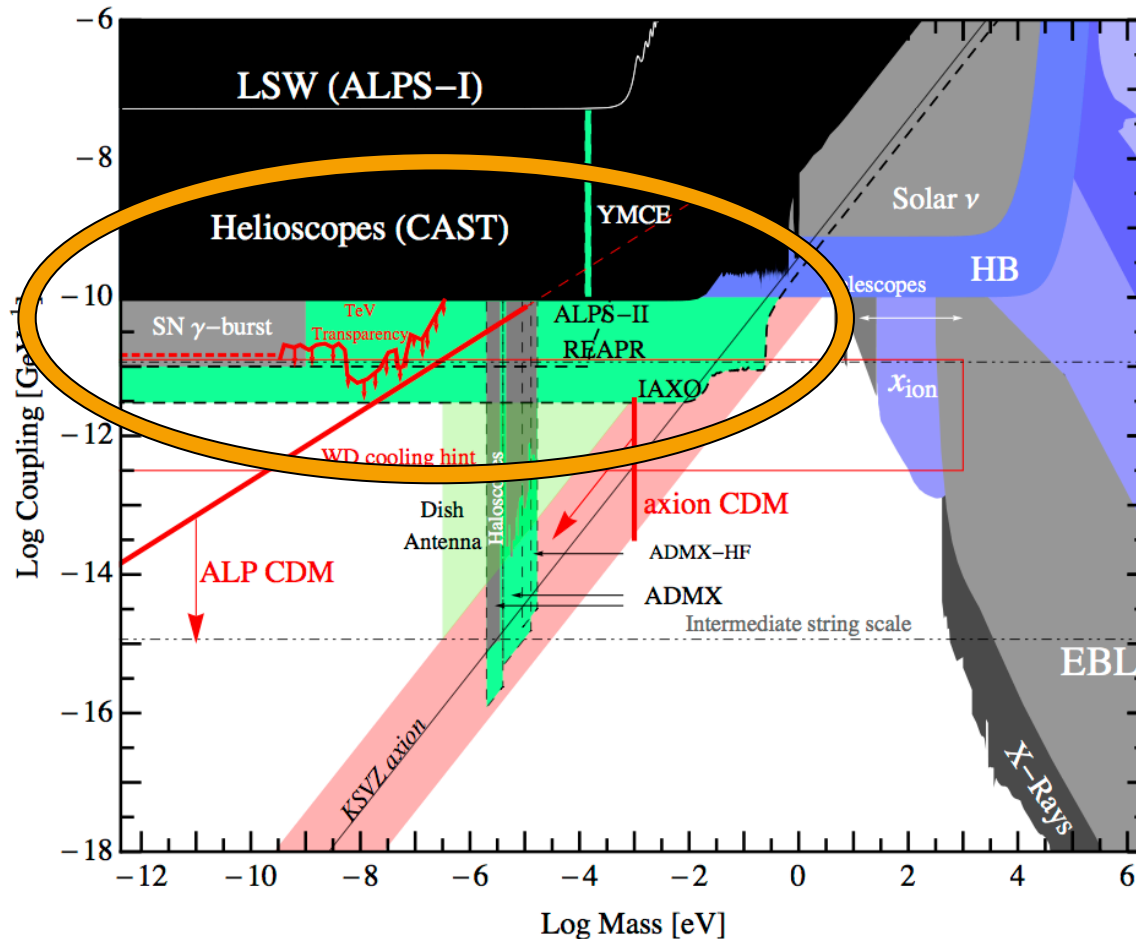
- Dedicated superconducting toroidal magnet with much bigger aperture than CAST
- Extensive use of X-ray optics
- Low background X-ray detectors



[Armengaud et al (IAXO CDR) 1401.3233]

Helioscope Searches

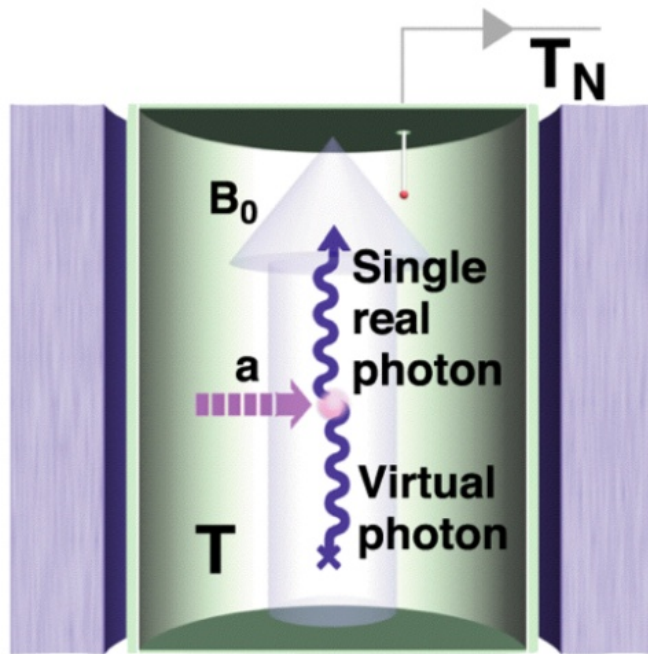
- Crucial test of the axion explanation of the excessive energy losses of RGs, WDs, n star in Cas A and ALP explanation of AGN spectra at VHE



adapted from [Hewett et al 12]

Haloscope Searches: Resonant Cavities

- > Direct detection of axion/ALP dark matter!
- > Axion or ALP DM – photon conversion in microwave cavity placed in magnetic field [Sikivie 83]



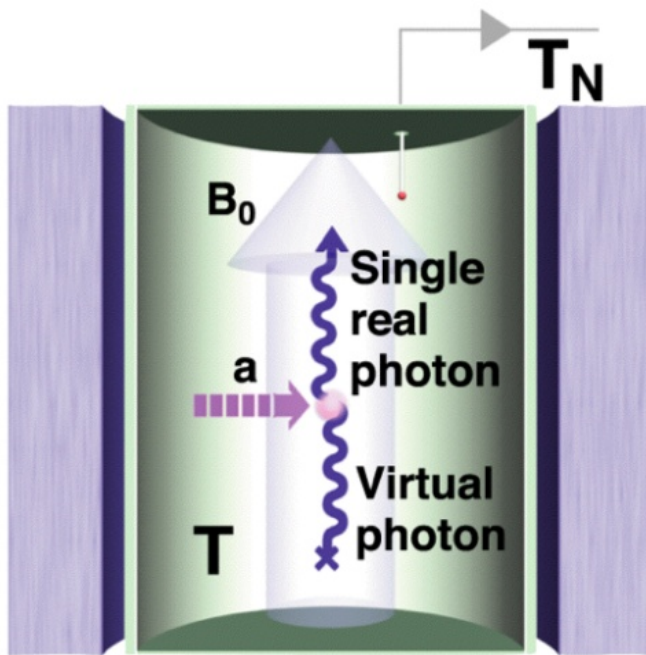
$$P_{\text{out}} \sim g^2 |\mathbf{B}_0|^2 \rho_{\text{DM}} V Q / m_a$$

- > Best sensitivity: mass = resonance frequency $m_a = 2\pi\nu \sim 4 \mu\text{eV} \left(\frac{\nu}{\text{GHz}} \right)$

Haloscope Searches: Resonant Cavities

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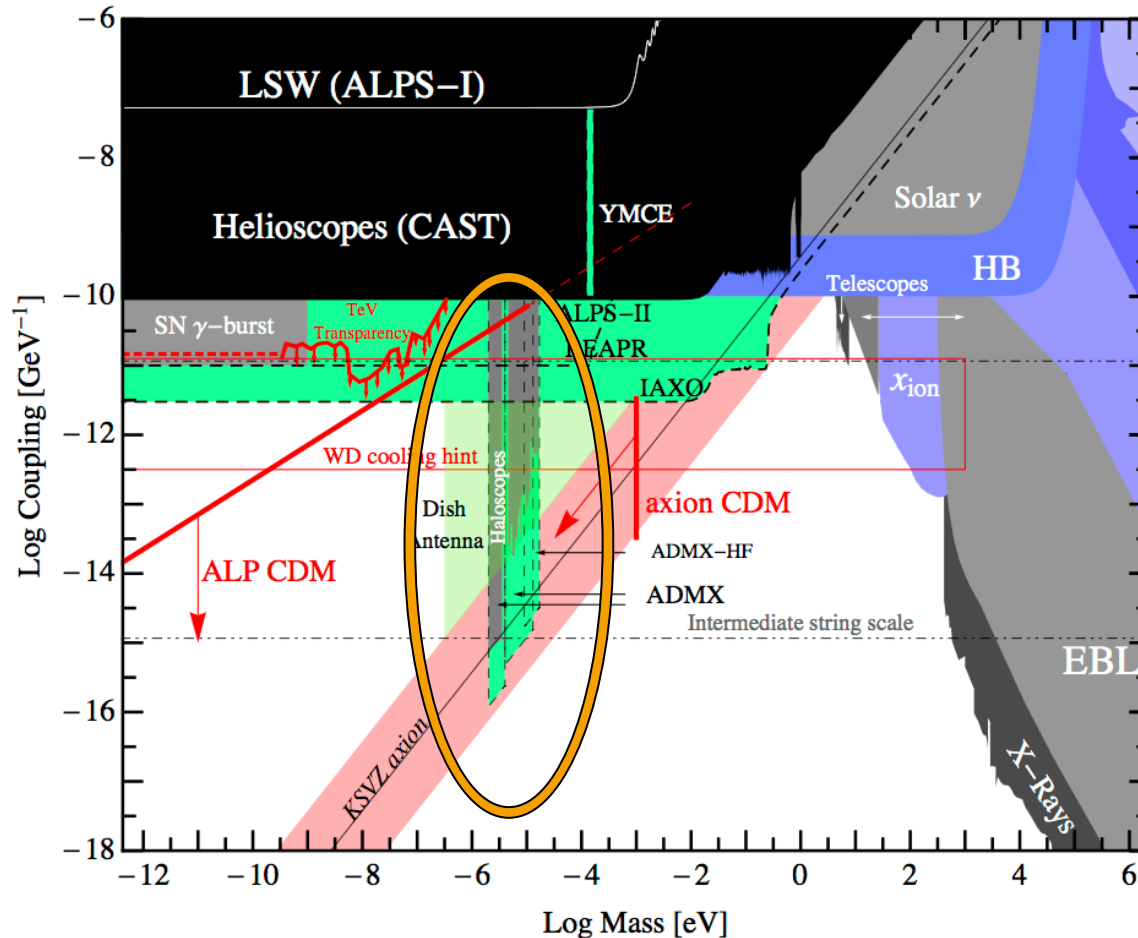
[Sikivie 83]



- > Best sensitivity: mass = resonance frequency $m_a = 2\pi\nu \sim 4 \mu\text{eV} \left(\frac{\nu}{\text{GHz}} \right)$
- > Ongoing: [ADMX](#) (Seattle), exploiting high Q cavity in 8 T SC solenoid

Haloscope Searches: Resonant Cavities

- ADMX able to probe about 1.5 decades in axion/ALP mass:



adapted from [Hewett et al 12]

Haloscopes: MR Searches for Oscillating EDMs

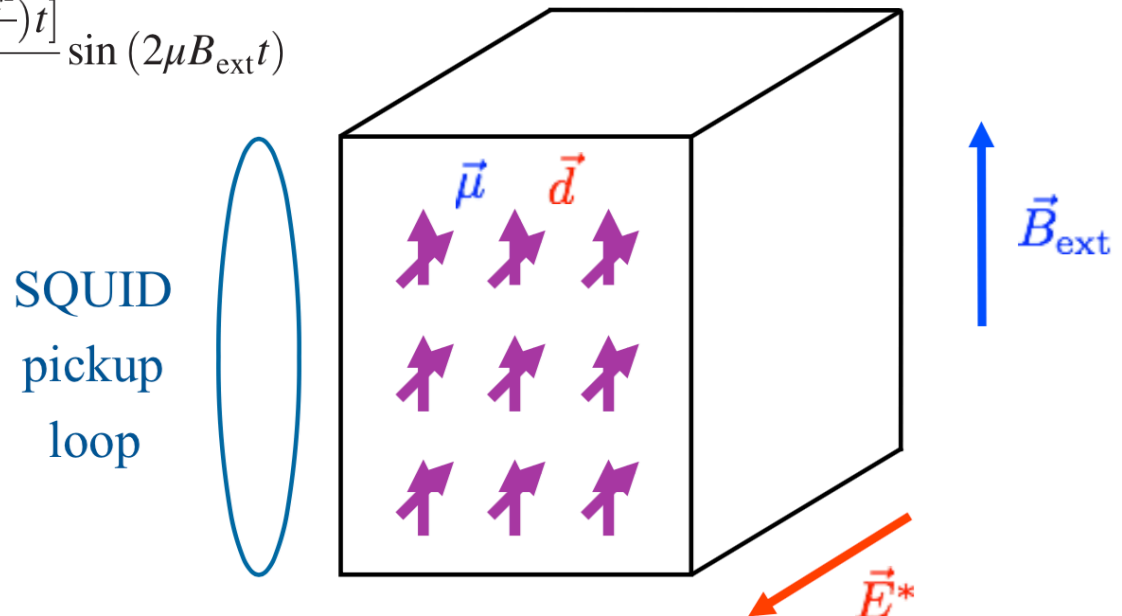
- > Axion DM gives all nucleons oscillating electric dipole moments (EDMs)

$$d_N \equiv g_{Ad} A(t) \sim e \frac{m_u m_d}{(m_u + m_d) m_N^2} \frac{A(t)}{f_A} \sim 10^{-16} \frac{A(t)}{f_A} e \text{ cm}$$

$$\frac{A(t)}{f_A} \sim \frac{\sqrt{\rho_{\text{DM}}}}{m_A f_A} \cos(m_A t) \sim \frac{\sqrt{\rho_{\text{DM}}}}{m_\pi f_\pi} \cos(m_A t) \sim 10^{-19} \cos(m_A t)$$

- > EDMs cause precession of nuclear spins in a nucleon spin polarized sample in the presence of an electric field
- > Resulting transverse magnetisation can be searched for exploiting magnetic resonance (MR) techniques [Budker, Graham, Ledbetter, Rajendran, Sushkov]

$$M(t) \approx n p \mu E^* \epsilon_S d_n \frac{\sin\left[\left(\frac{2\mu B_{\text{ext}} - m_a c^2}{\hbar}\right)t\right]}{\frac{2\mu B_{\text{ext}} - m_a c^2}{\hbar}} \sin(2\mu B_{\text{ext}} t)$$



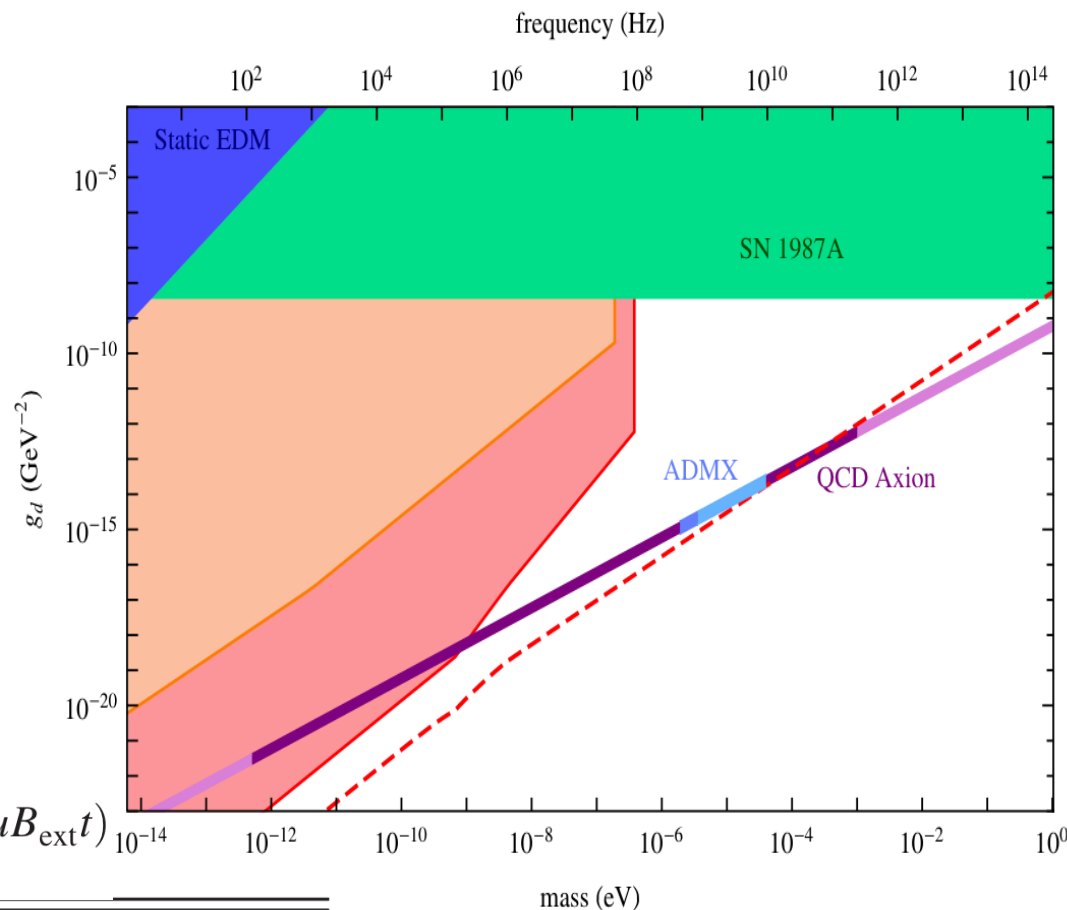
Haloscopes: MR Searches for Oscillating EDMs

> Cosmic Axion Spin Precession Experiment (CASPER) in preparation in Mainz

- Probes very light axion, from GUT to Planck scale SSB (anthropic axion)
- Sensitive in a mass/symmetry breaking scale region complementary to ADMX

$$M(t) \approx np\mu E^* \epsilon_S d_n \frac{\sin\left[\left(\frac{2\mu B_{\text{ext}} - m_a c^2}{\hbar}\right)t\right]}{\frac{2\mu B_{\text{ext}} - m_a c^2}{\hbar}} \sin(2\mu B_{\text{ext}} t)$$

| | n | E^* | p | T_2 | Max B_{ext} |
|---------|---------------------------|------------------------------|-----------|-------|----------------------|
| Phase 1 | 10^{22} cm^{-3} | $3 \times 10^8 \text{ V/cm}$ | 10^{-3} | 1 ms | 10 T |
| Phase 2 | 10^{22} cm^{-3} | $3 \times 10^8 \text{ V/cm}$ | 1 | 1 s | 20 T |



[Budker et al. 14]

$$\mathcal{L} \ni -\frac{i}{2} g_d a \bar{N} \sigma_{\mu\nu} \gamma_5 N F^{\mu\nu}$$



Haloscopes: MR Searches for the Axion/ALP Wind

- > The axion/ALP nucleon coupling

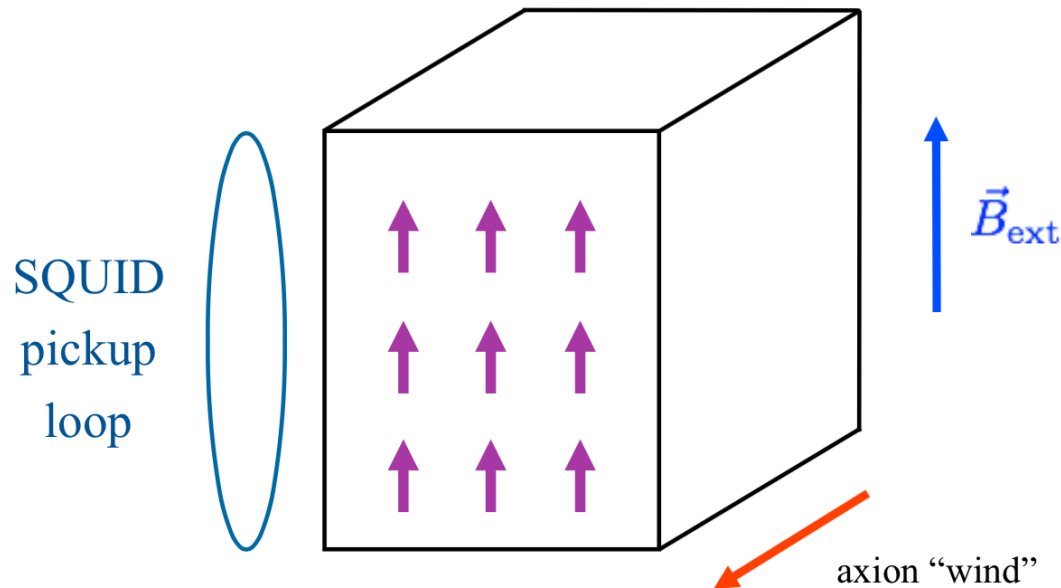
$$g_{\text{aNN}} (\partial_\mu a) \bar{N} \gamma^\mu \gamma_5 N \implies H_N \supset g_{\text{aNN}} \vec{\nabla} a \cdot \vec{S}_N$$

will lead to a spin precession about the axion/ALP DM wind

- > Resulting magnetisation

[Graham, Rajendran 13]

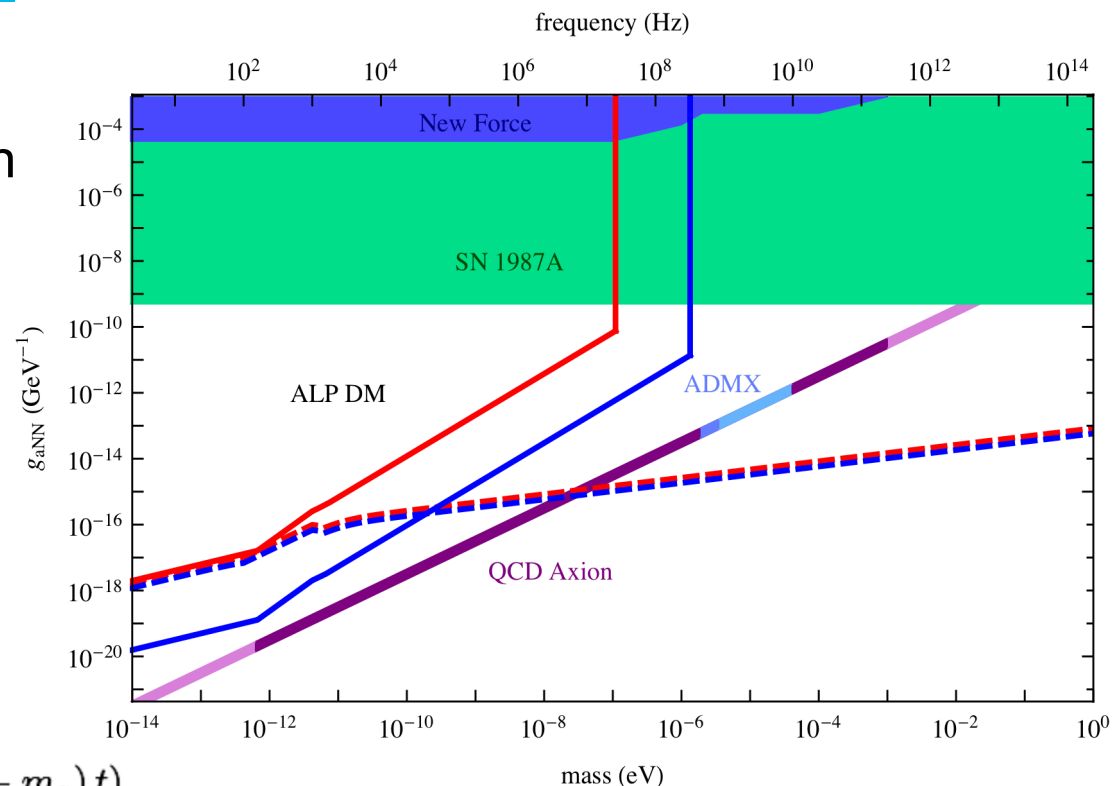
$$M(t) \approx np\mu \left(g_{\text{aNN}} \sqrt{2\rho_{\text{DM}}} v \right) \frac{\sin((2\mu B_{\text{ext}} - m_a)t)}{2\mu B_{\text{ext}} - m_a} \sin(2\mu B_{\text{ext}} t)$$



Haloscopes: MR Searches for the Axion/ALP Wind

> Cosmic Axion Spin Precession Experiment (CASPEr) in preparation in Mainz

- Sensitivity does not reach axion prediction



$$M(t) \approx np\mu \left(g_{aNN} \sqrt{2\rho_{DM}v} \right) \frac{\sin((2\mu B_{\text{ext}} - m_a)t)}{2\mu B_{\text{ext}} - m_a} \sin(2\mu B_{\text{ext}}t)$$

[Graham, Rajendran 13]

| | Element | Density (n) | Magnetic Moment (μ) | T_2 | Max. B | Magnetometer Sensitivity |
|----|---------------|--|------------------------------|-------|--------|--|
| 1. | Xe | $1.3 \times 10^{22} \frac{1}{\text{cm}^3}$ | $0.35 \mu_N$ | 100 s | 10 T | $10^{-16} \frac{\text{T}}{\sqrt{\text{Hz}}}$ |
| 2. | ^3He | $2.8 \times 10^{22} \frac{1}{\text{cm}^3}$ | $2.12 \mu_N$ | 100 s | 20 T | $10^{-17} \frac{\text{T}}{\sqrt{\text{Hz}}}$ |

$$\mathcal{L} \supset g_{aNN} (\partial_\mu a) \bar{N} \gamma^\mu \gamma_5 N$$



Haloscopes: ESR/MR Searches for the Axion/ALP Wind

- > The axion/ALP electron coupling

[Krauss et al. 85; Barbieri et al. 89]

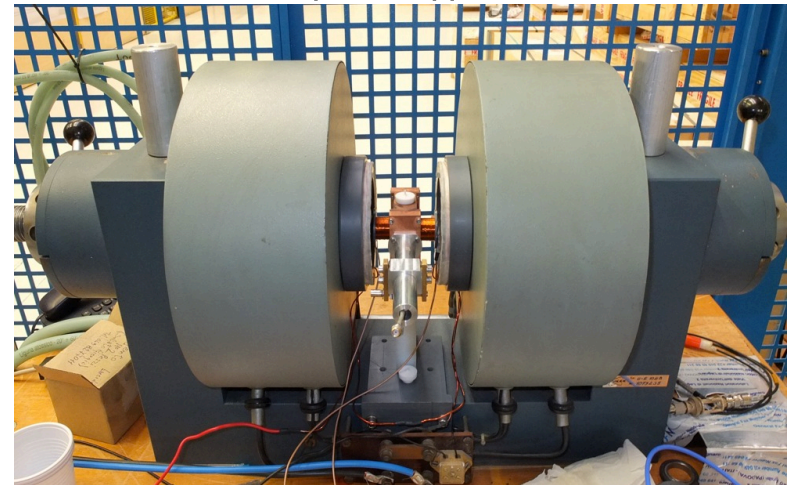
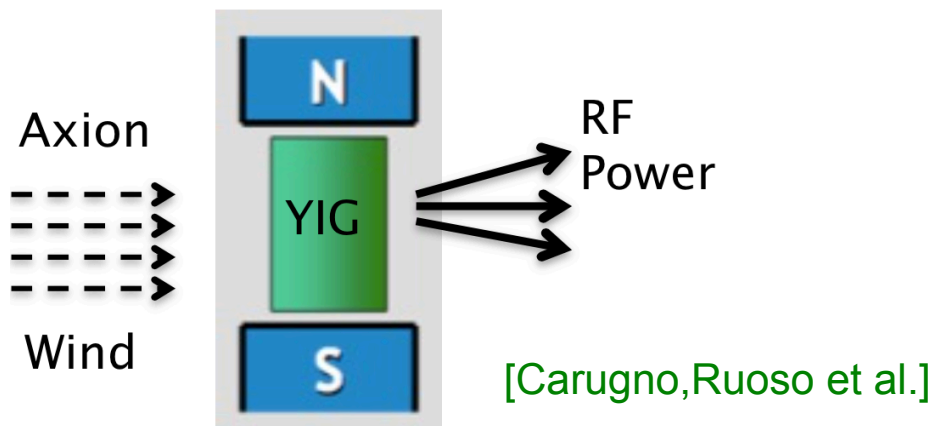
$$\mathcal{L} \supset g_{aee} \partial_\mu a (\bar{e} \gamma_5 \gamma^\mu e)$$

will also lead to a spin precession about the axion/ALP DM wind

- > Larmor frequency and thus sensitivity extended to higher masses by factor

$$\mu_B / \mu_N \sim m_N / m_e \sim 10^3$$

- > **QUAX** (QUaerere AXions) in preparation by INFN (Legnaro, Padua, Torino), Birmingham, Moscow aims to exploit magnetic resonance (MR) inside a magnetized material (Electron Spin Resonance (ESR))



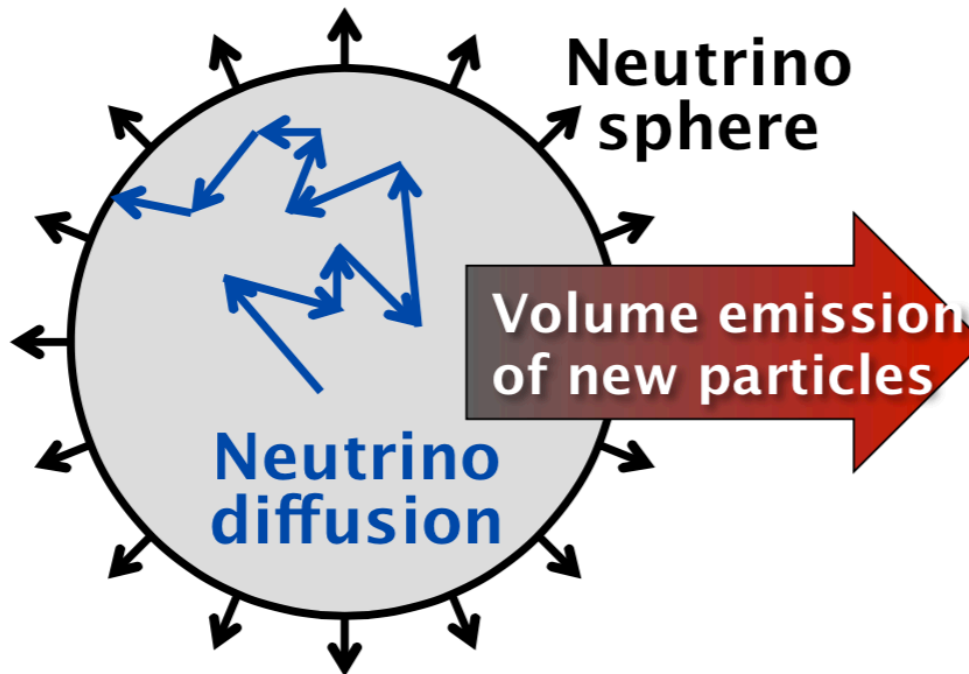
Summary

- > Strong physics case for axion and ALPs:
 - Axion and ALPs occur naturally as NG bosons from breaking of well motivated symm.
 - Solution of strong CP problem
 - Candidates for dark matter
 - Explanation of astrophysical hints (energy losses of stars; AGN spectra)
- > Large parts in axion and ALPs parameter space can be tackled in the upcoming decade by a number of terrestrial experiments:
 - Light-shining-through-a-wall experiments ([ALPS II](#), ...)
 - Helioscopes ([IAXO](#), ...)
 - Haloscopes ([ADMX](#), [CASPEr](#), [QUAX](#), ...)
- > Stay tuned!



Backup: Axion and Neutrino Signal from CC Supernovae

- Late-time neutrino signal most sensitive observable
 - Early neutrino burst powered by accretion, not sensitive to volume energy loss
 - Late-time neutrino signal associated with neutrino diffusion. Emission of axions would steal energy from the late-time neutrino burst and shorten it

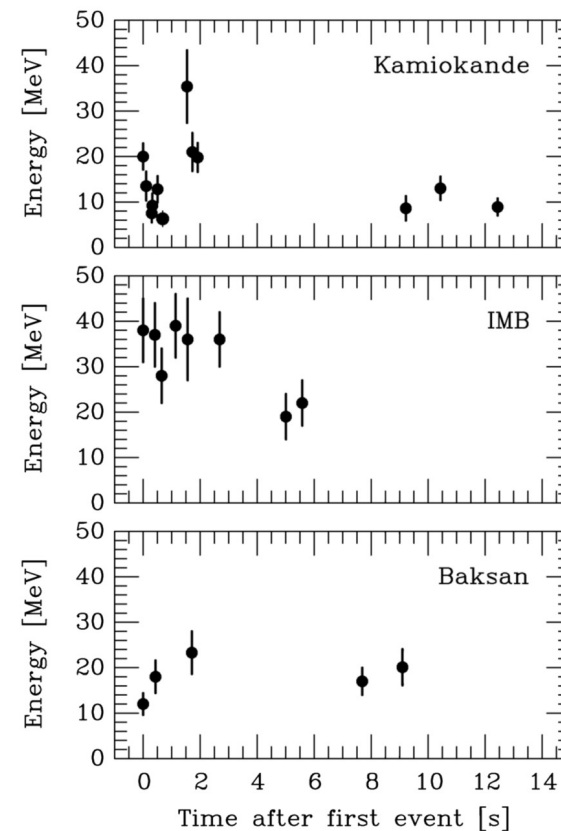
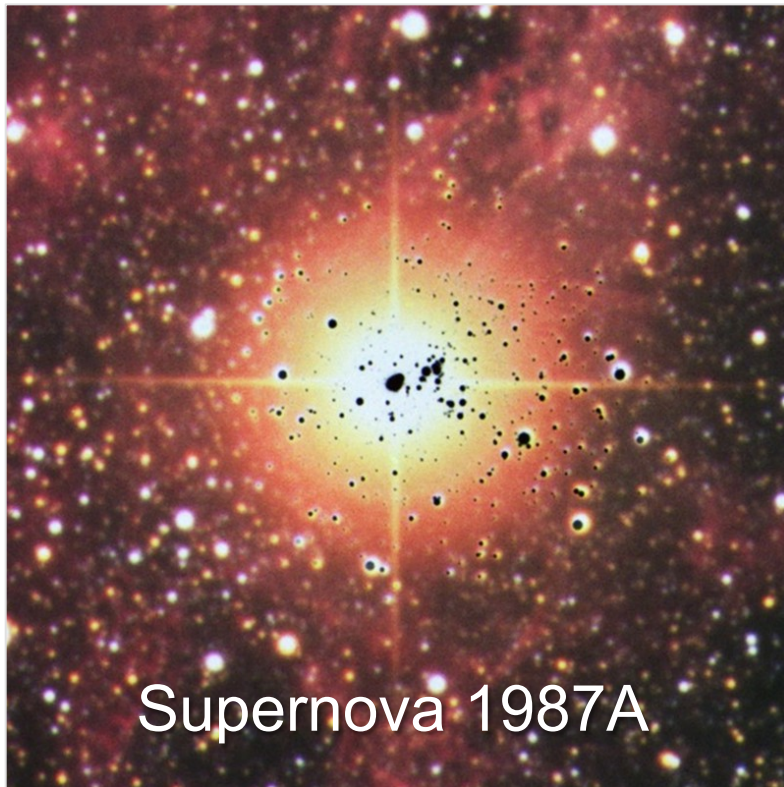


[Raffelt 14]

Backup: Axion and Neutrino Signal from CC Supernovae

➤ Late-time neutrino signal most sensitive observable

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Backup: Axion and Neutrino Signal from CC Supernovae

- > Late-time neutrino signal most sensitive observable
 - Early neutrino burst powered by accretion, not sensitive to volume energy loss
 - Late-time neutrino signal associated with neutrino diffusion. Emission of axions would steal energy from the late-time neutrino burst and shorten it
- > Limit on energy loss due to nucleon bremsstrahlung of axions/ALPs in SN 1987 A: [Raffelt 96,08]

$$g_{aN} \equiv \frac{C_{aN} m_N}{f_a} \lesssim 5 \times 10^{-10}, \quad \text{for } m_a \lesssim \text{MeV}.$$

- > Fundamental particle physics opportunity for neutrino telescopes if galactic core-collapse supernova is observed!
- > Sensitive in a range of symmetry breaking scale where other experiments have difficulties!

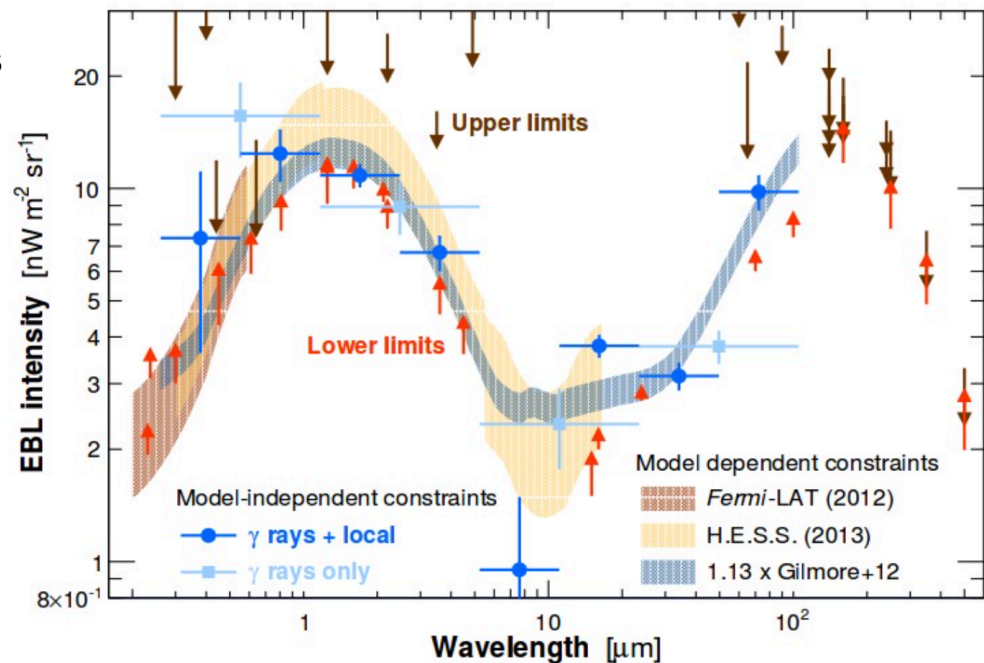


Backup: No Cosmic Pair Production Anomaly?

- Recent claim: no indication of pair production anomaly in VHE AGN spectra [Biteau, Williams 15]
- But then pair production anomaly replaced by EBL anomaly

Deconvolution of the EBL [arXiv:1502.04166]

- Approximate EBL with sums of Gaussians
- Use $O(100)$ spectra
- Tension with models 3-5 σ
- Tension with lower limits from galaxy counts at NIR
- Tension with Fermi-LAT estimate on EBL
- Additionally: too large H_0 , too small z for 1424+240



[Horns 15]



Backup: Isocurvature Constraints on Axion/ALPs

- > If $f_a > \max(H_I/(2\pi), \epsilon_{\text{eff}} E_I)$, quantum fluctuations $\delta\bar{\theta}_a = H_I/(2\pi f_a)$ of the axion/ALP lead to isocurvature (= entropy) fluctuations in CMB, with nearly scale-invariant power spectrum

$$\mathcal{P}_i(k) \simeq \left(\frac{\Omega_a}{\Omega_{\text{DM}}}\right)^2 \begin{cases} \frac{H_I^2}{\pi^2 f_a^2 \langle \bar{\theta}_a \rangle^2}, & \text{for } \langle \bar{\theta}_a \rangle^2 \gg \left(\frac{H_I}{2\pi f_a}\right)^2 \\ 2, & \text{for } \langle \bar{\theta}_a \rangle^2 \ll \left(\frac{H_I}{2\pi f_a}\right)^2 \end{cases}$$

- > Non-observation rules out existence of axion/ALP, unless

$$H_I \lesssim 10^{13} \text{ GeV}$$

- > Detection of $\mathcal{P}_t/\mathcal{P}_s \equiv r = 0.20^{+0.07}_{-0.05}$ by BICEP2 would have implied

$$H_I \simeq \frac{1}{4} \sqrt{A_s r \pi} M_{\text{Pl}} = 1.1 \times 10^{14} \text{ GeV} \left(\frac{r}{0.2}\right)^{1/2}$$

- > $f_a > 1.8 \times 10^{13} \text{ GeV}$ strongly disfavored
[Fox et al. hep-th/0409059; Higaki et al. 1403.4186;
Marsh et al. 1403.4216; Visinelli, Gondolo 1403.4594]

Andreas Ringwald | The Hunt for Axions, Palazzo Franchetti, Istituto Veneto di Scienze, Lettere ed Arti

