

# *Holographic description of total hadronic cross sections at high energies*

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theory and experiment (QCD@Work2018)

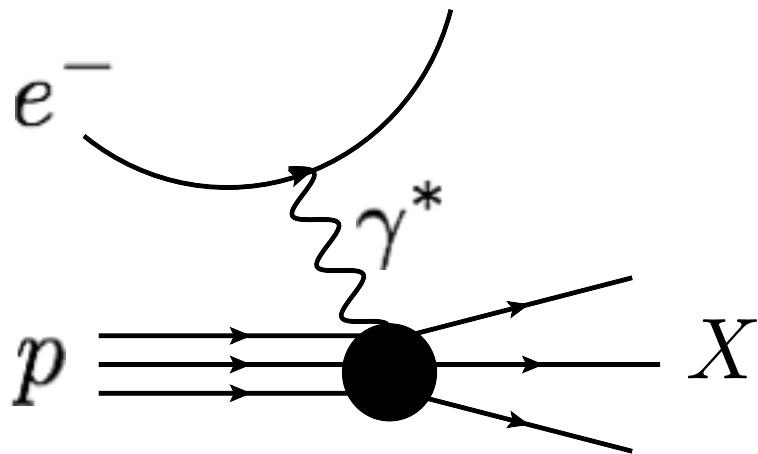
June 25, 2018 @ Matera, Italy

# Outline

1. Introduction of previous analysis on DIS at small  $x$  in holographic QCD
2. Motivation
3. Analysis on total hadronic cross sections
4. Summary

# DIS structure functions

$$\frac{d^2\sigma}{dxdQ^2} = \frac{2\pi\alpha^2}{xQ^4} \left[ \left\{ 1 + (1-y)^2 \right\} F_2(x, Q^2) - y^2 F_L(x, Q^2) \right]$$



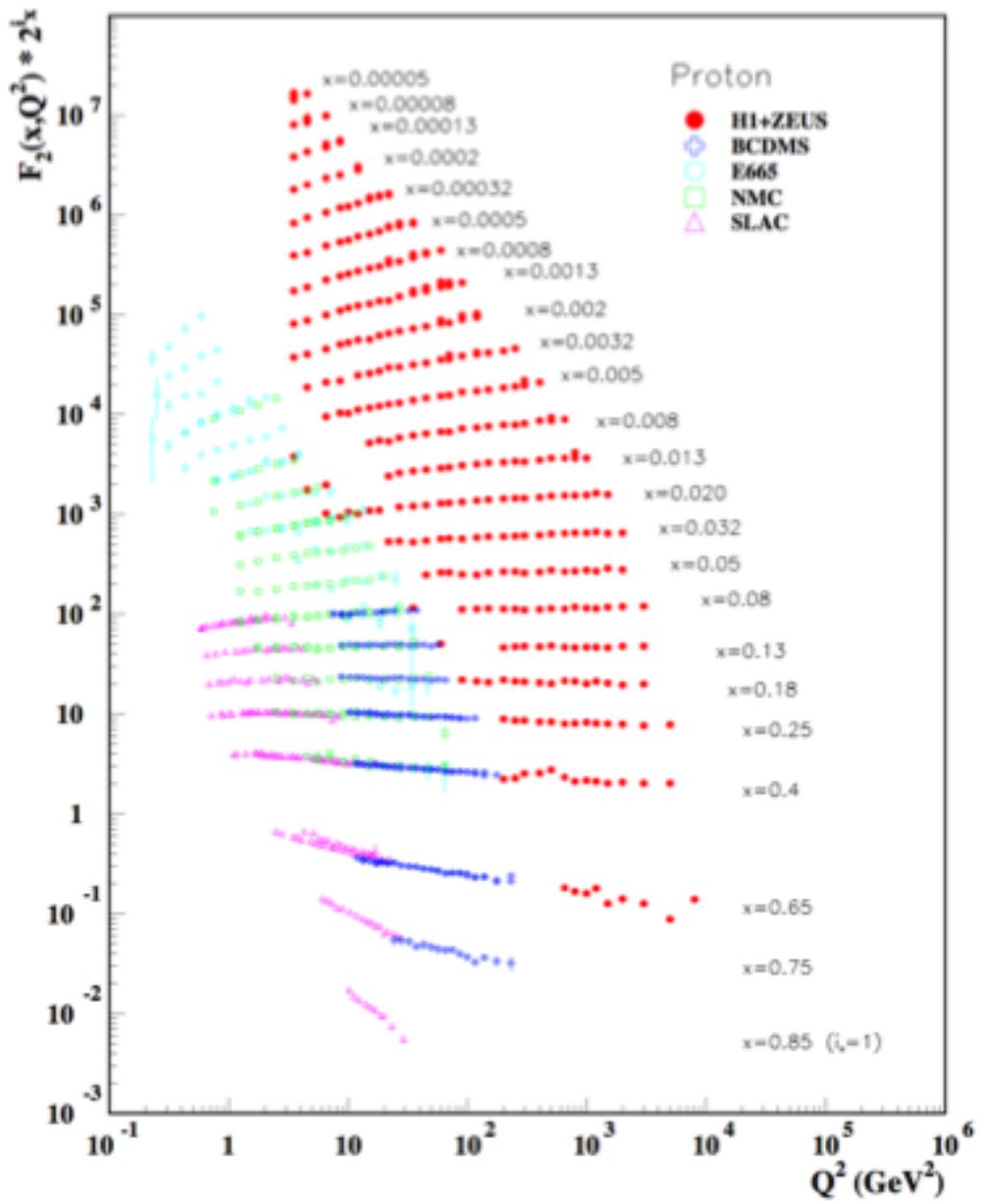
$$F_2(x, Q^2) = \frac{Q^2}{4\pi^2\alpha} \sigma_{tot}(x, Q^2)$$

$$F_L(x, Q^2) = \frac{Q^2}{4\pi^2\alpha} \sigma_L(x, Q^2)$$

- Structure functions are physical quantities which have information on the internal structure of hadrons.
- They depend on two kinematic variables, Bjorken-x and photon 4-momentum squared  $Q^2$ .

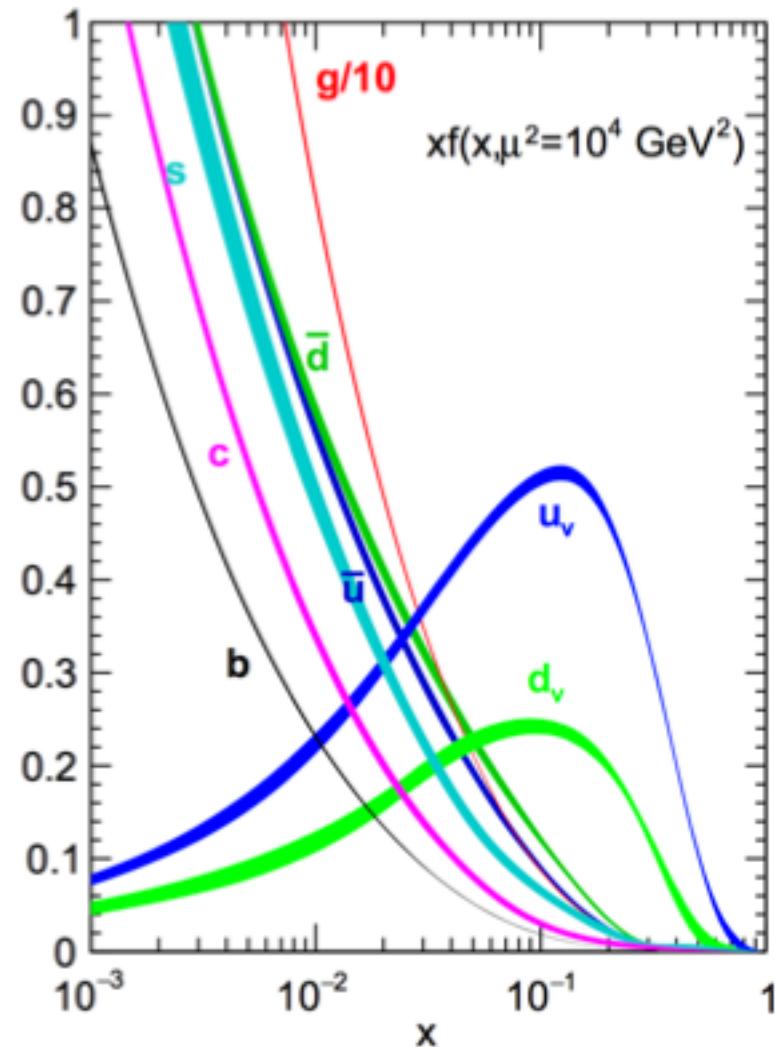
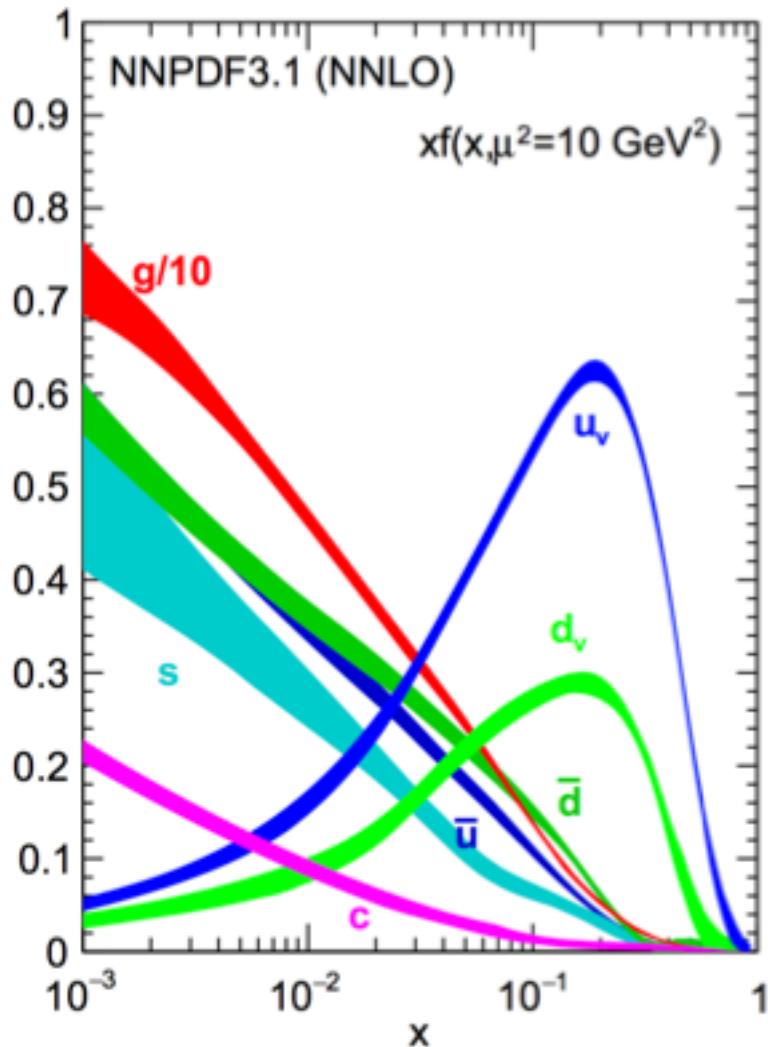
# $F_2$ structure function of proton

- Proton  $F_2$  has been measured in the very wide kinematic range so far.
- Structure functions are basically nonperturbative physical quantities.



# Proton parton distribution functions (PDFs)

NNPDF Collaboration (2017)

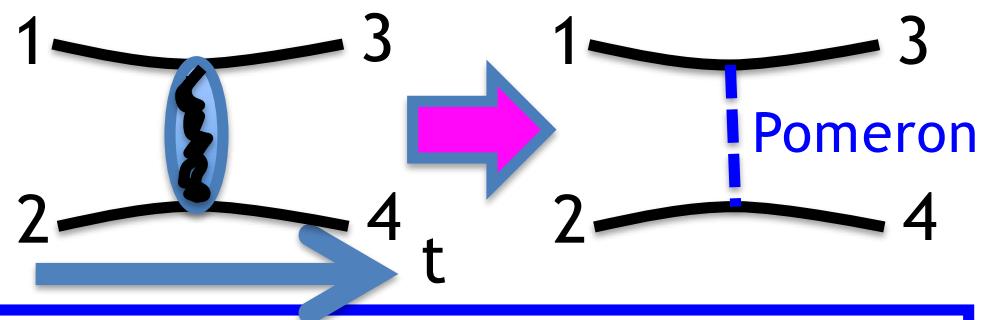


# Pomeron exchange picture

- A description of high energy scattering before QCD
- Pomeron: a color singlet gluonic object
- Total cross sections for high energy two-to-two scattering can be well reproduced by this picture

$$s = (p_1 + p_2)^2 = (p_3 + p_4)^2$$

$$\sigma_{tot}(s) \sim s^{\alpha_0 - 1}$$



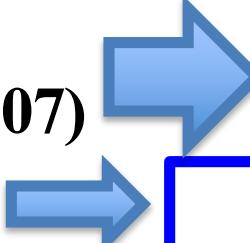
can be expressed by a single parameter (Pomeron intercept)

Therefore,  $F_2$  structure function can be written as

$$F_2(x, Q^2) \sim x^{1-\alpha_0}$$

(this is effective only in the small  $x$  region)

# Holographic description of structure functions

- Polchinski-Strassler (2003)
  - Brower-Polchinski-Strassler-Tan (2007)
  - Brower-Djuric-Sarcevic-Tan (2010)
- 

$$\mathcal{A}(s,t) = 2is \int d^2b e^{iq \cdot b} \int dz dz' P_{13}(z) P_{24}(z') \left\{ 1 - e^{i\chi(s,b,z,z')} \right\}$$

$$F_2(x, Q^2) = \frac{Q^2}{2\pi^2} \int dz dz' P_{13}(z, Q^2) P_{24}(z', Q'^2) \text{Im}[\chi(s, z, z')]$$

$z$  and  $z'$  : 5th coordinate

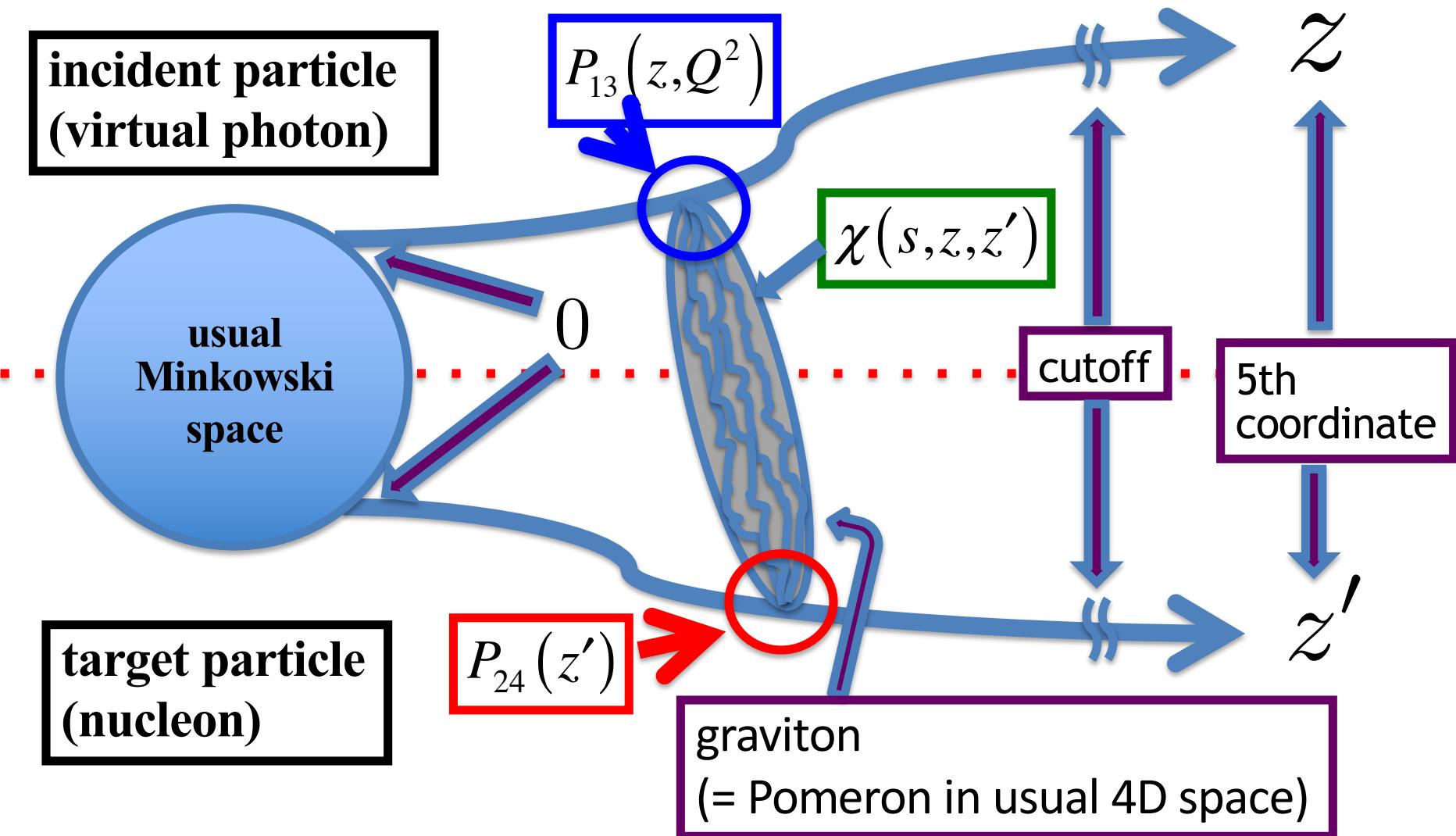
$\chi$  : Pomeron exchange kernel in the AdS space

$P_{13}(z, Q^2)$  : incident particle  
( $Q$ : 4-momentum)

$P_{24}(z', Q'^2)$  : target particle

} overlap functions  
(density distributions  
in the AdS space)

# 5D background spacetime ( $\text{AdS}_5$ )



# Pomeron exchange kernel

Brower-Polchinski-Strassler-Tan (2007)  
 Brower-Strassler-Tan (2009)

$$F_i(x, Q^2) = \frac{g_0^2 \rho^{3/2} Q^2}{32\pi^{5/2}} \int dz dz' P_{13}^{(i)}(z, Q^2) P_{24}(z', Q'^2) (zz') \text{Im}[\chi(s, z, z')] \quad i = 2 \text{ or } L$$

$$\text{Im}[\chi_c(s, z, z')] \equiv e^{(1-\rho)\tau} e^{-\frac{\log^2 z/z'}{\rho\tau}} / \tau^{1/2}$$

$$\tau = \log(\rho zz's/2)$$

mimicking confinement effect

$$\text{Im}[\chi_{\text{mod}}(s, z, z')] \equiv \text{Im}[\chi_c(s, z, z')] + \mathcal{F}(z, z', \tau) \text{Im}[\chi_c(s, z, z_0^2/z')]$$

$$\mathcal{F}(z, z', \tau) = 1 - 2\sqrt{\rho\pi\tau} e^{\eta^2} \text{erfc}(\eta)$$

$$\eta = \left( -\log \frac{zz'}{z_0^2} + \rho\tau \right) / \sqrt{\rho\tau}$$

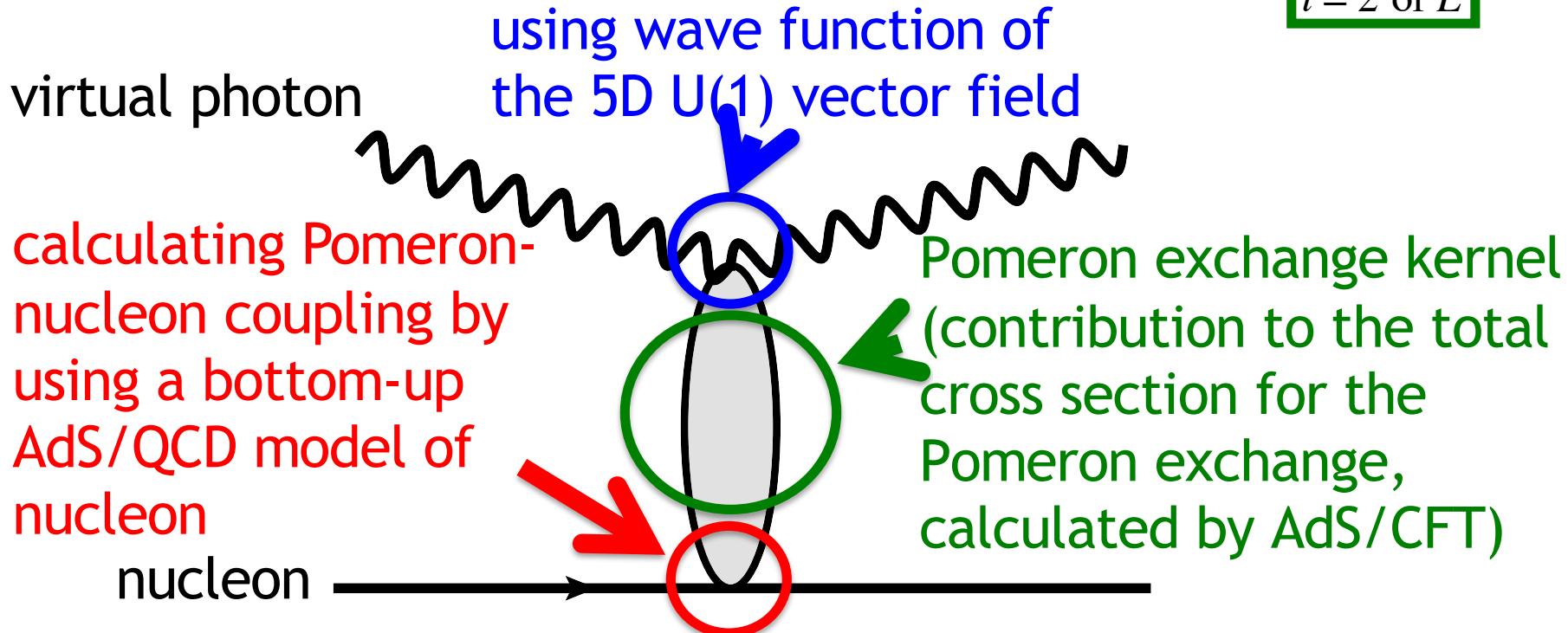
energy dependence magnitude

strength of confinement effect

3 adjustable parameters :  $\rho$ ,  $g_0^2$ ,  $z_0$

# Model setup

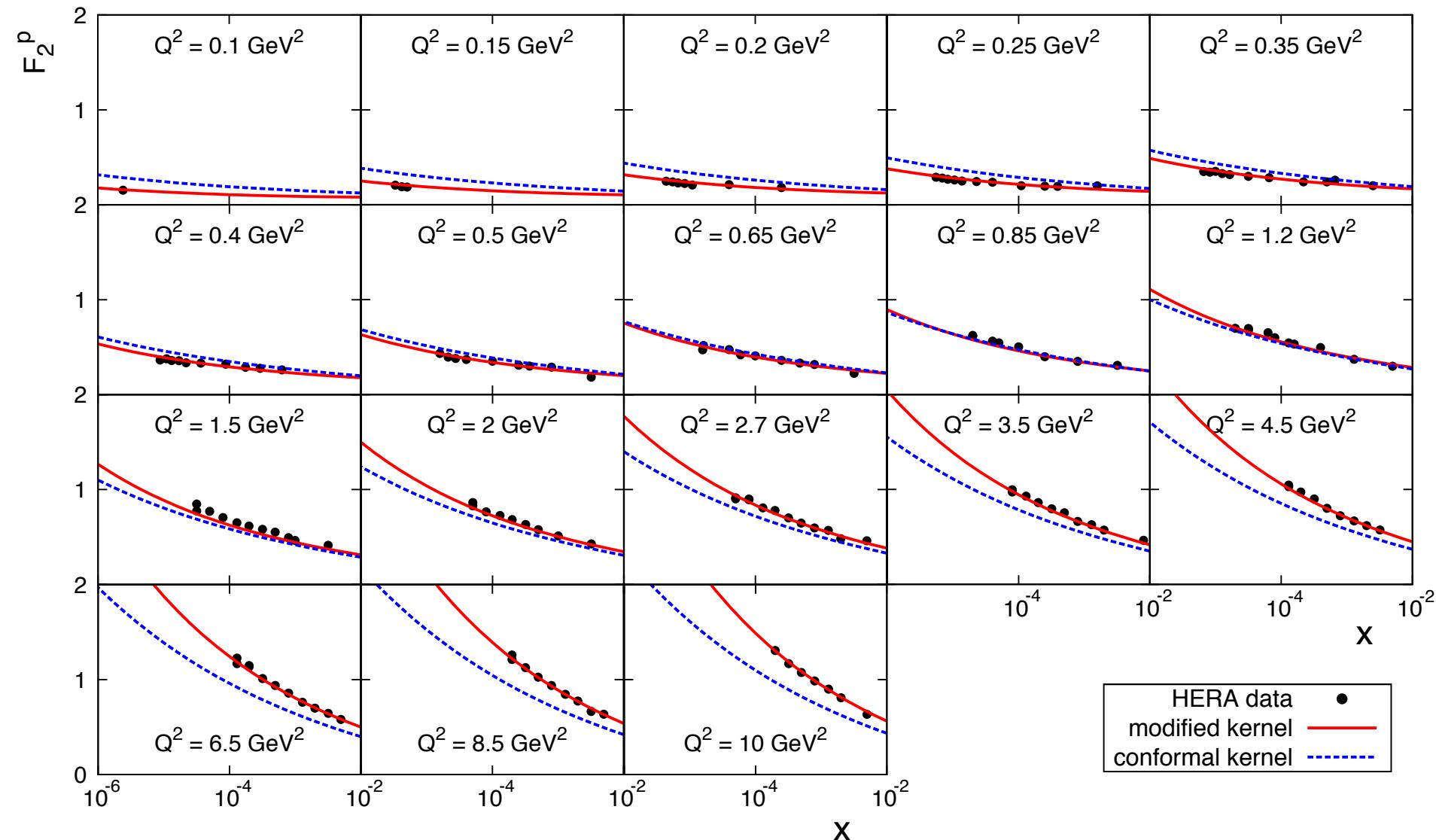
$$F_i(x, Q^2) = \frac{g_0^2 \rho^{3/2} Q^2}{32\pi^{5/2}} \int dz dz' P_{13}^{(i)}(z, Q^2) P_{24}(z', Q'^2) (zz') \text{Im}[\chi(s, z, z')] \quad i = 2 \text{ or } L$$



This framework can be applied to other hadrons' SFs,  
longitudinal SFs, and hadron-hadron scattering.

# Proton structure function

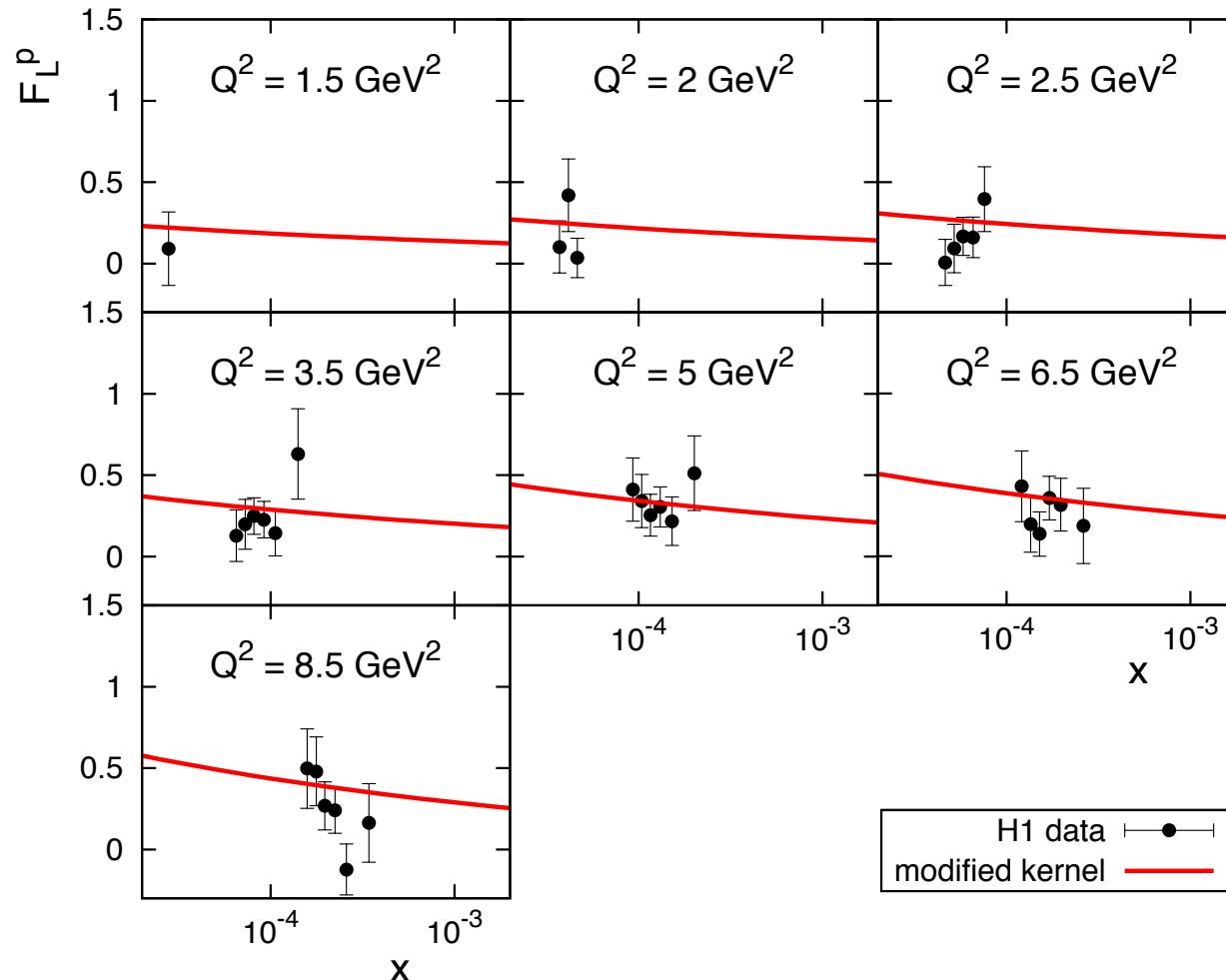
AW-Suzuki (2014)



# Proton longitudinal structure function

- Replacing the density distribution of the probe photon with its longitudinal component

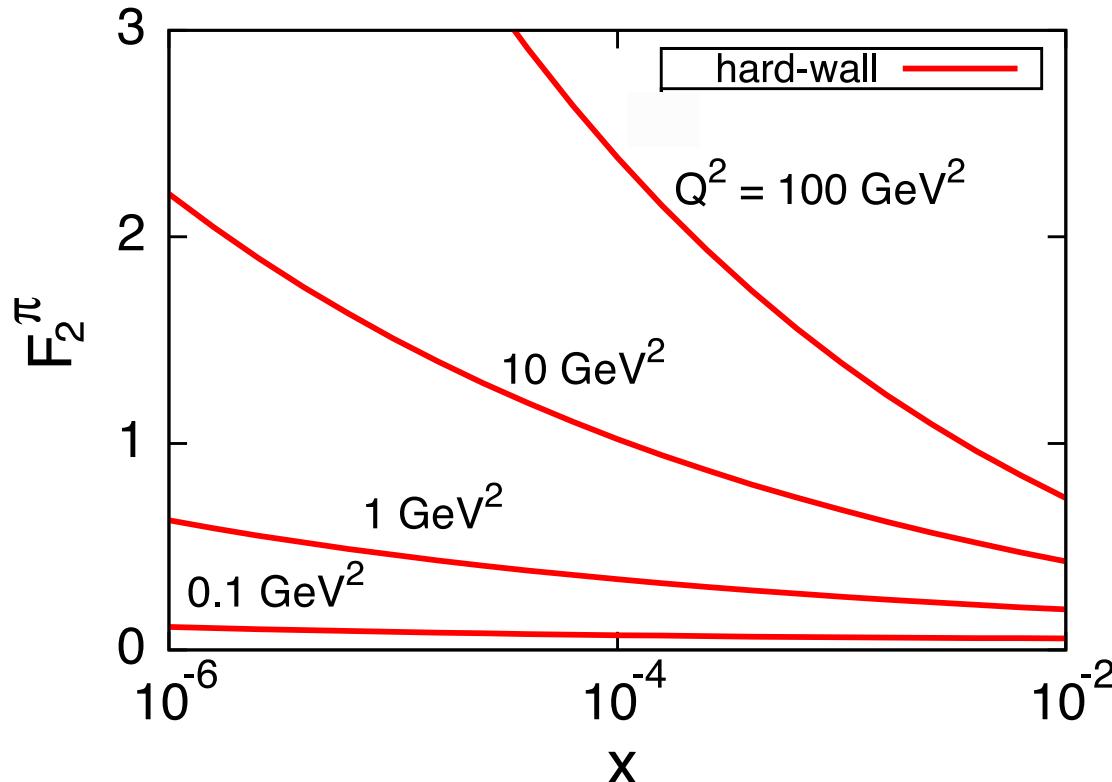
AW-Suzuki (2014)



# Pion structure function

- Replacing the density distribution for target hadron with pion's, which can be calculated by using a bottom-up AdS/QCD model of mesons

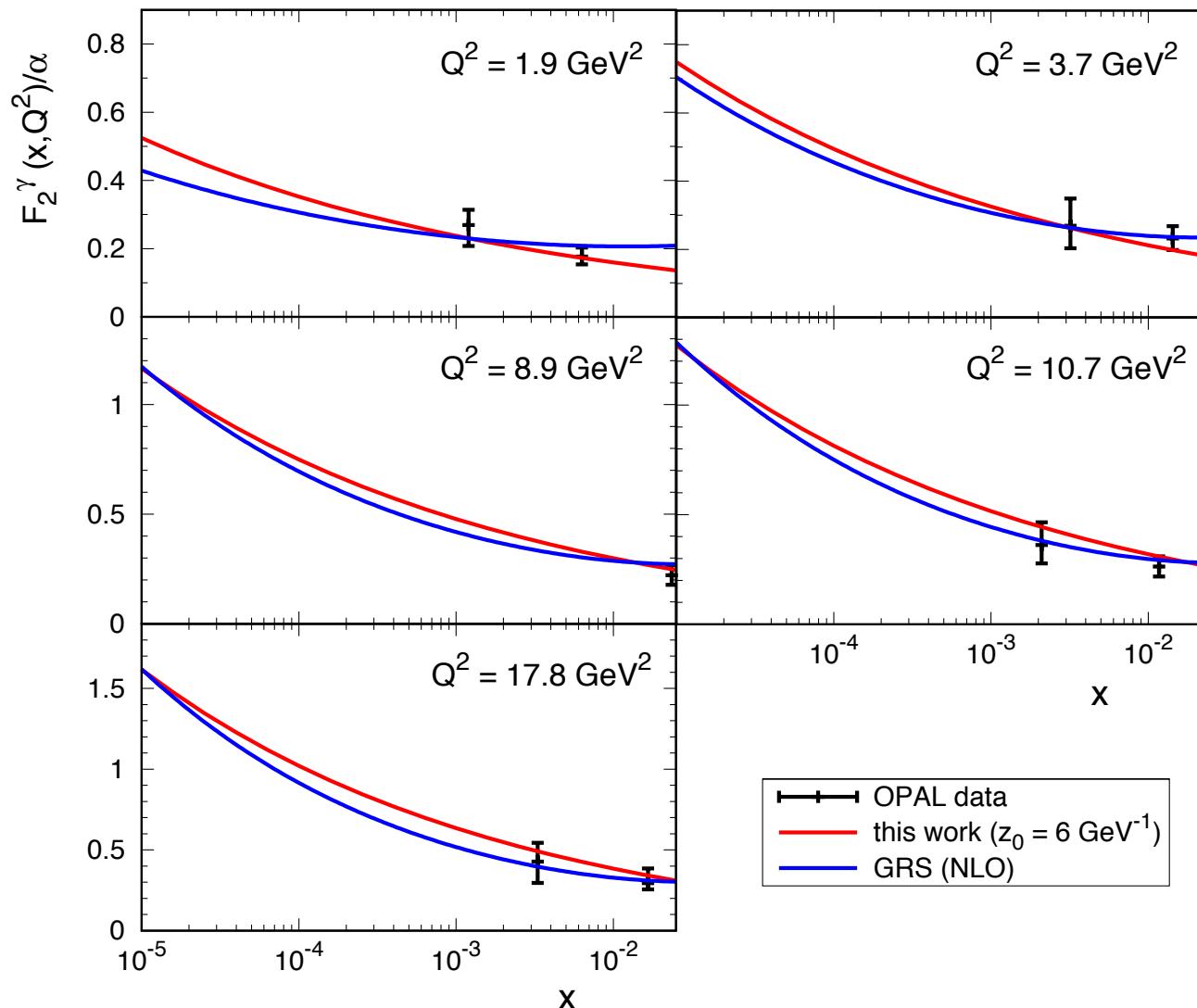
AW-Suzuki (2012)



- There is no experimental data

# Photon structure function

AW-Li (2015)

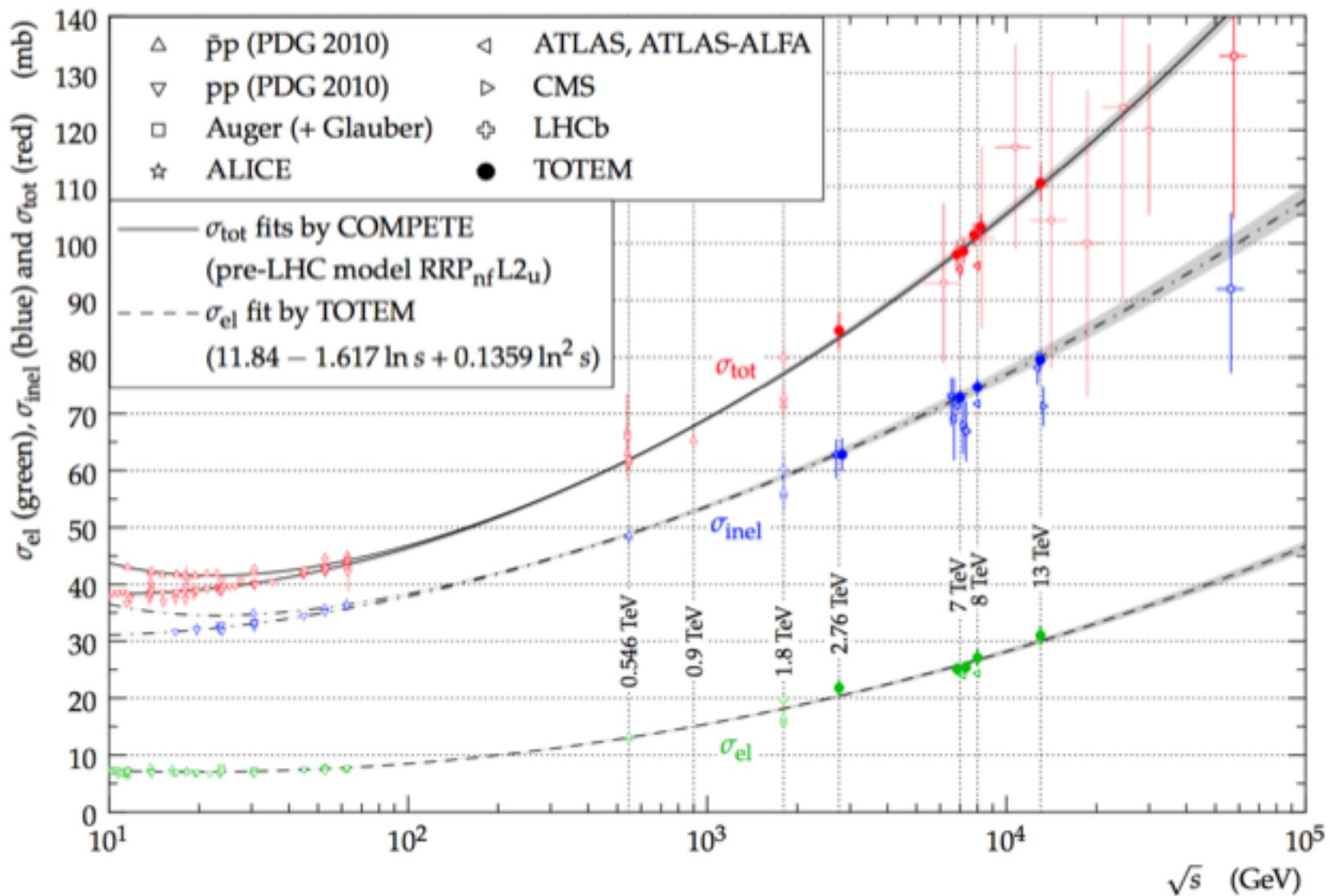


# Motivation (why total cross section?)

- TOTEM 13TeV data have recently come out!
- 13TeV corresponds to  $x \sim 10^{-8}$  which is much smaller than that at HERA ( $x > 10^{-6}$ ).
- In such a very small  $x$  region, is the (single) Pomeron exchange still working well?
- How about the holographic (BPST) Pomeron?
- Via the analysis, we might obtain some information about the applicable limit of the model.

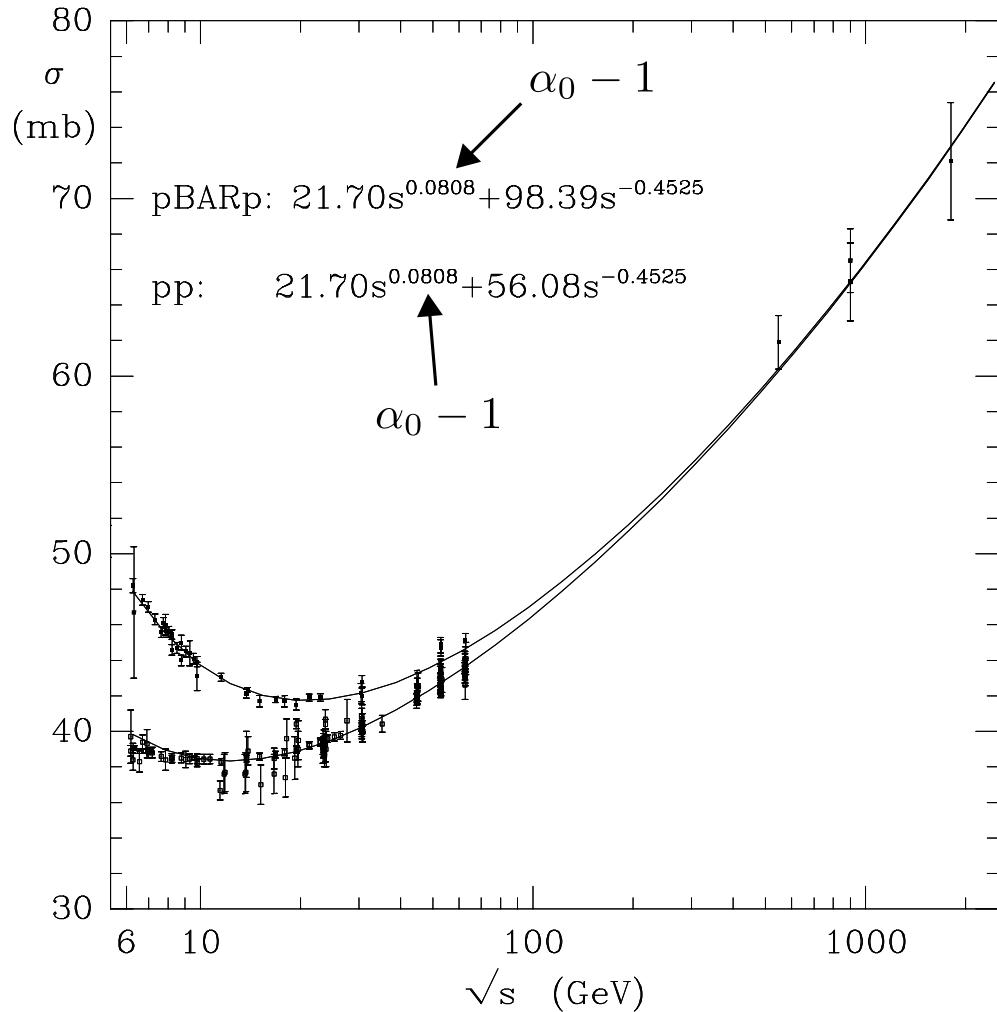
# 13TeV data from TOTEM

TOTEM Collaboration (2017)

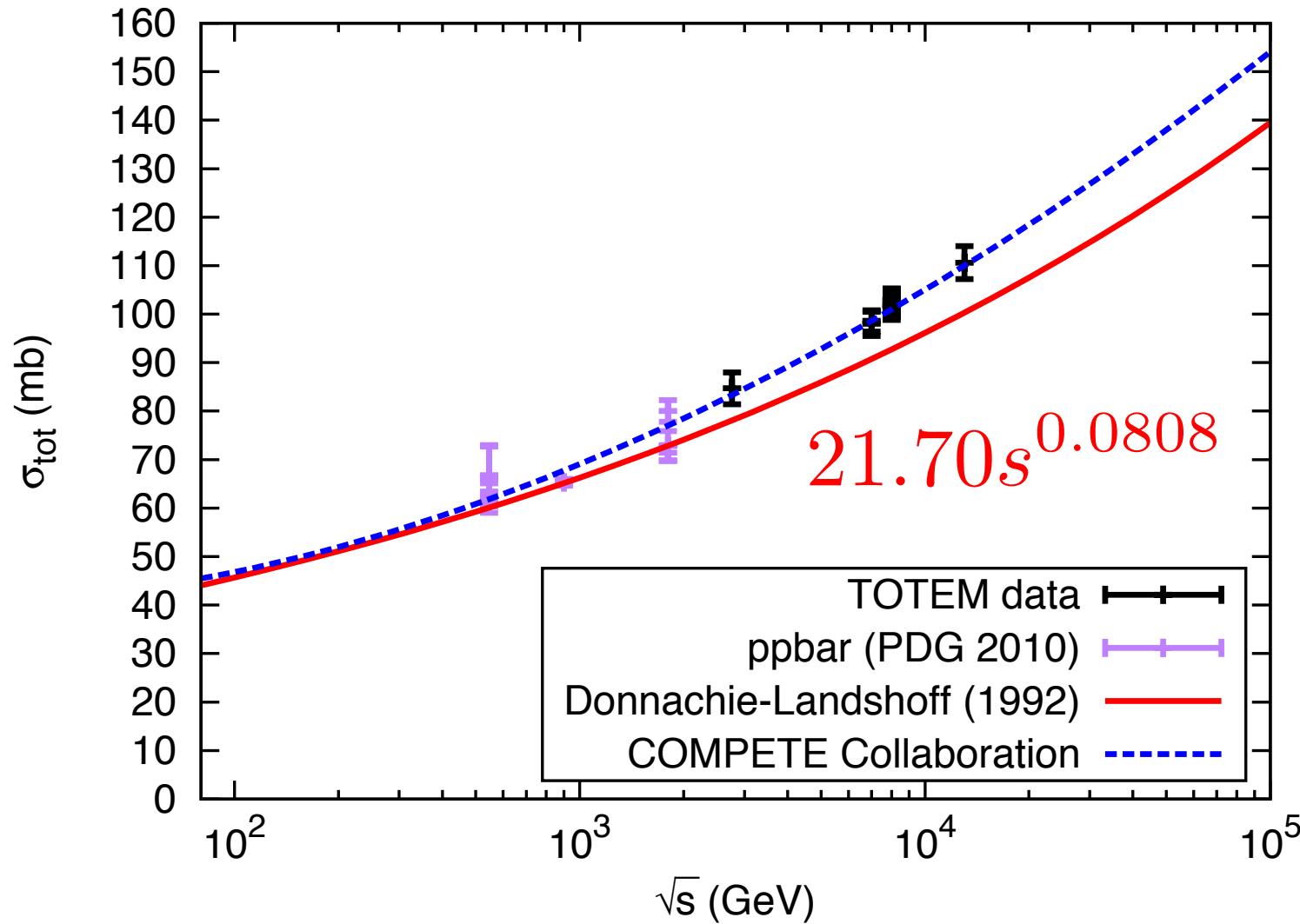


# Total cross sections via Pomeron exchange

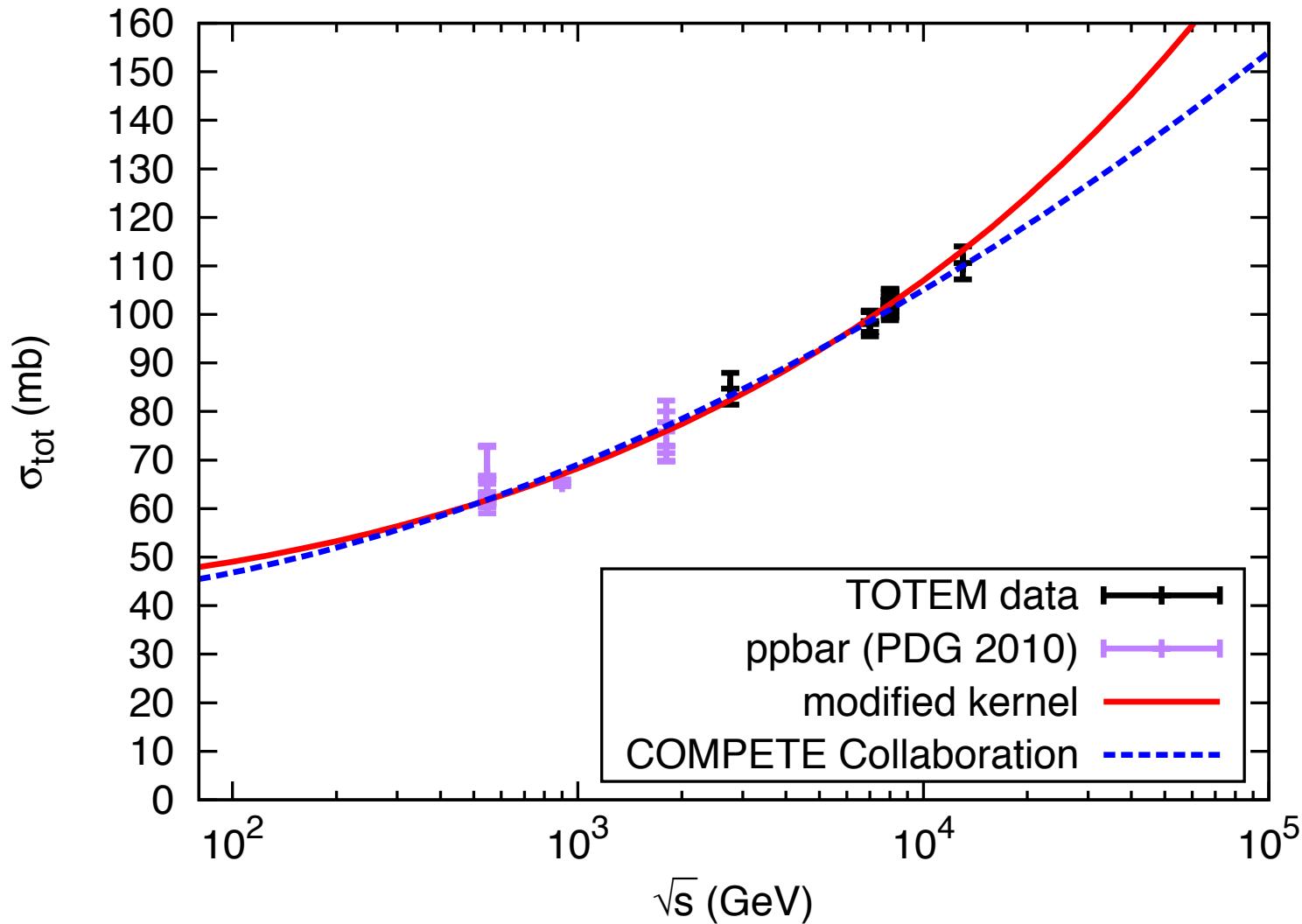
## proton-(anti)proton total cross section



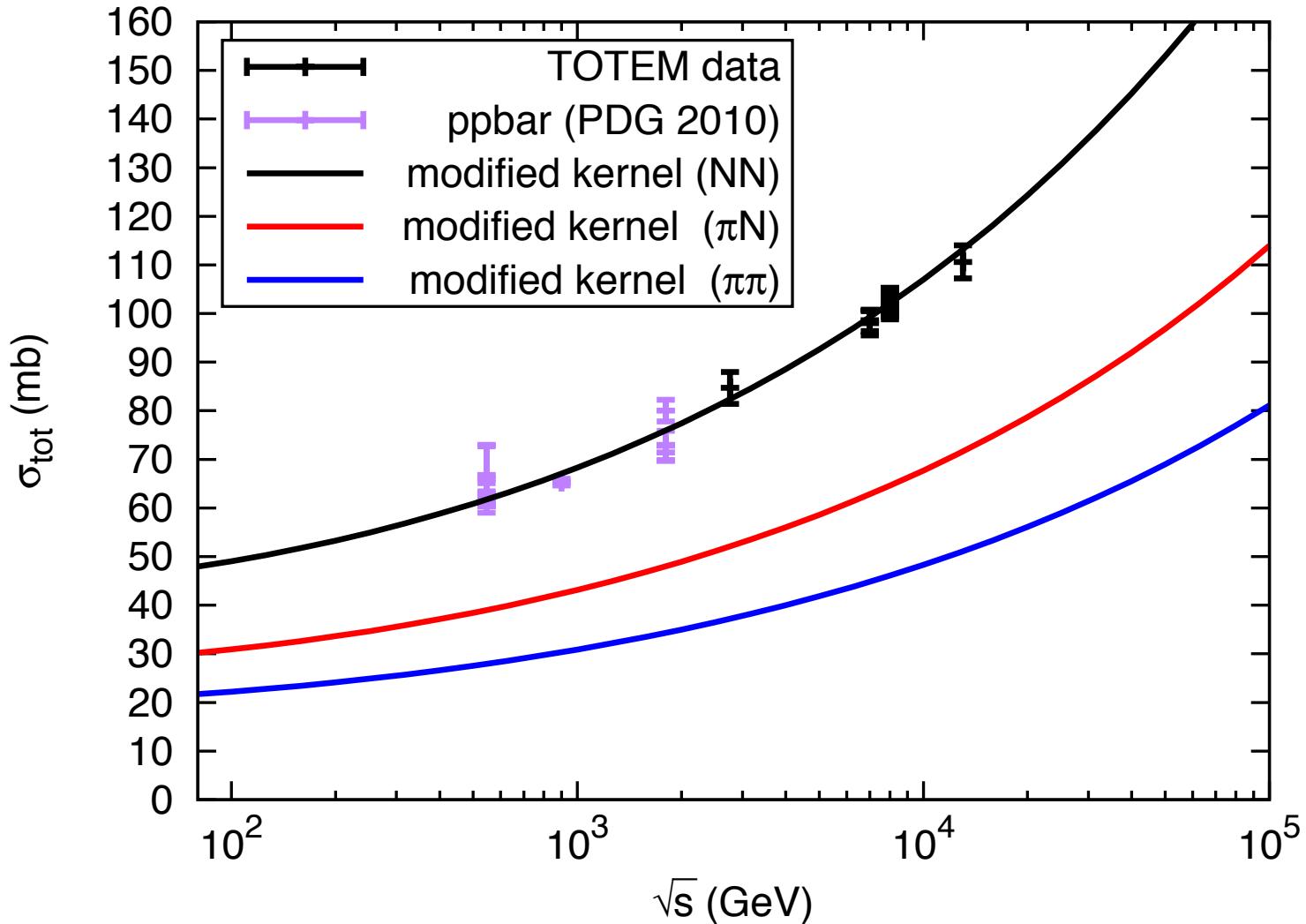
# Can the soft Pomeron reproduce the data?



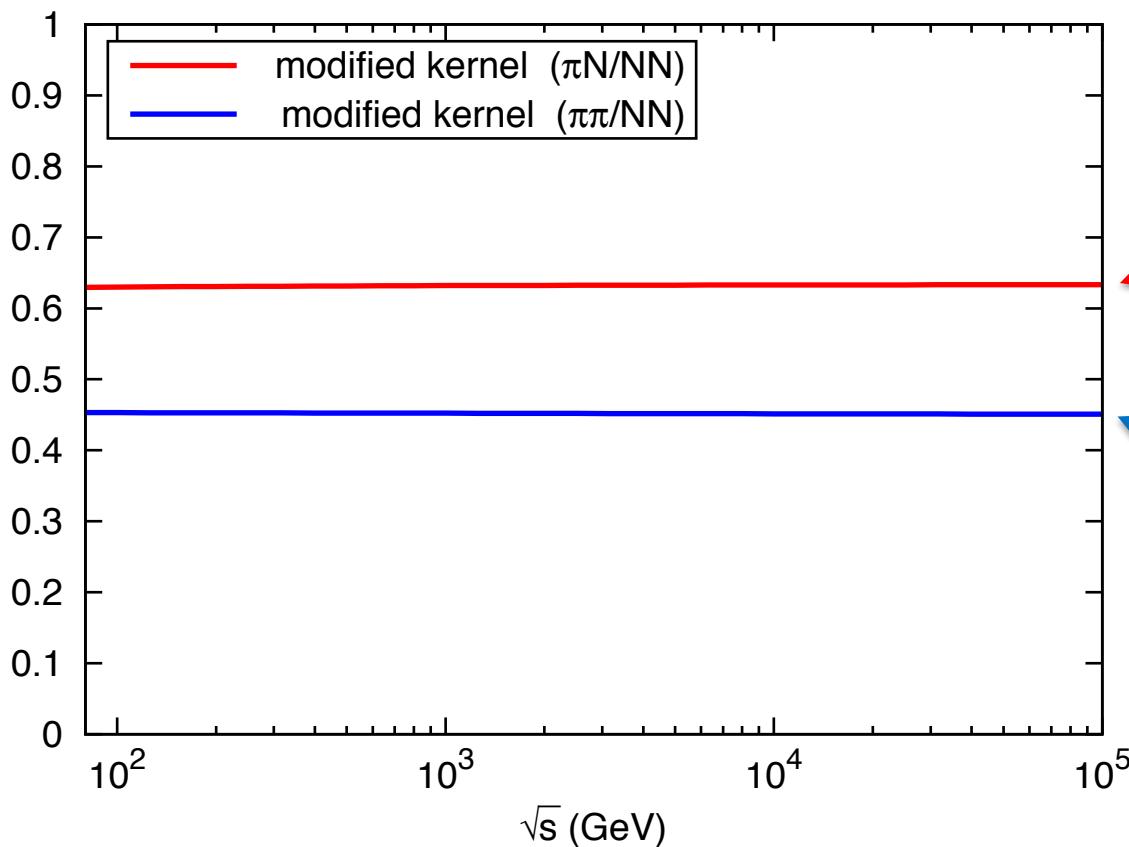
# How about the holographic (BPST) Pomeron?



# $\pi$ -N and $\pi$ - $\pi$ scattering



# Total cross section ratios



PREDICTIONS:

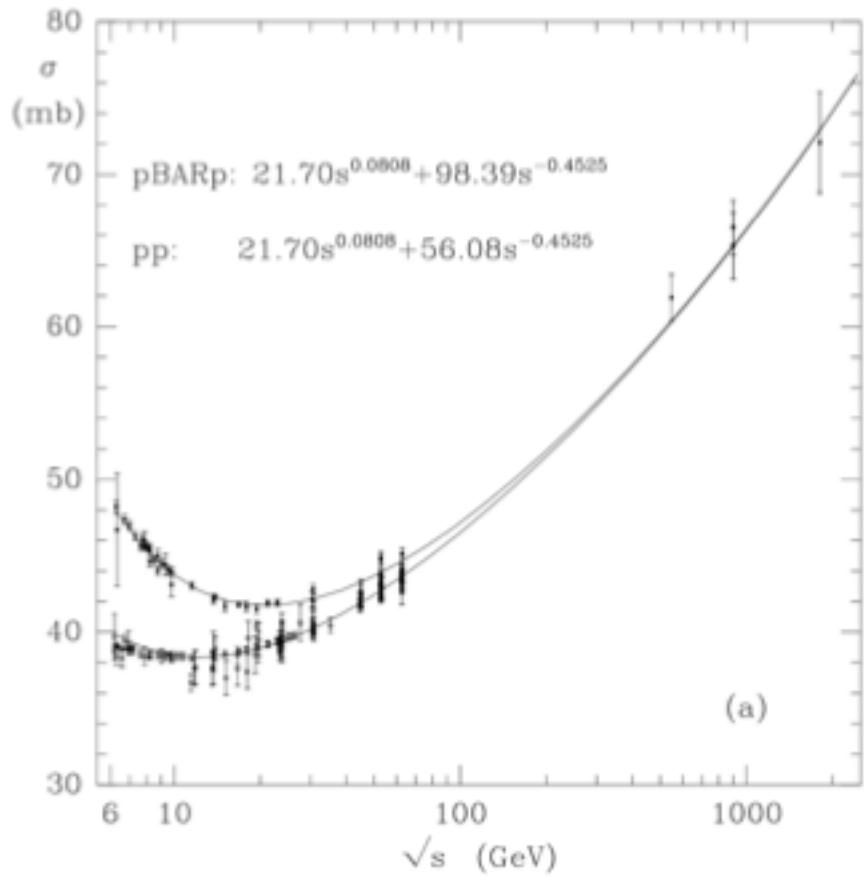
$$\left( \frac{\sigma_{tot}^{\pi N}}{\sigma_{tot}^{NN}} \right) \approx 0.63$$

$$\left( \frac{\sigma_{tot}^{\pi\pi}}{\sigma_{tot}^{NN}} \right) \approx 0.45$$

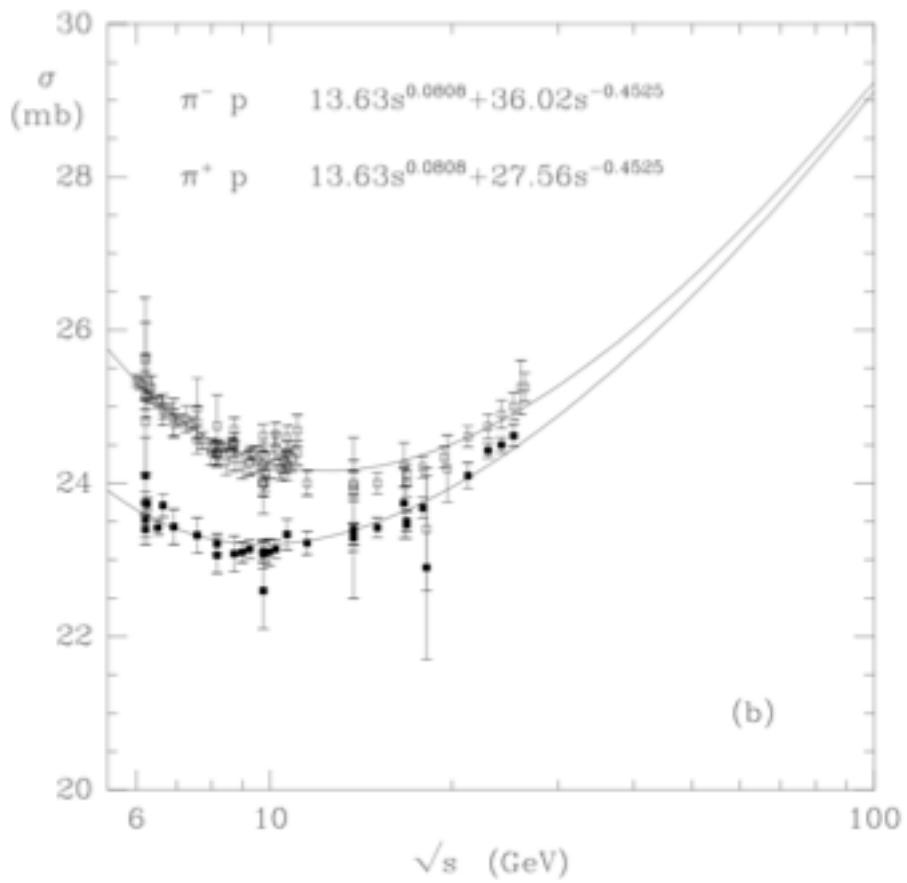
N.B. we do NOT need any additional parameter to calculate the pion involved cross sections!

# p-p and $\pi$ -p scattering by Donnachie-Landshoff

Donnachie-Landshoff (1992)

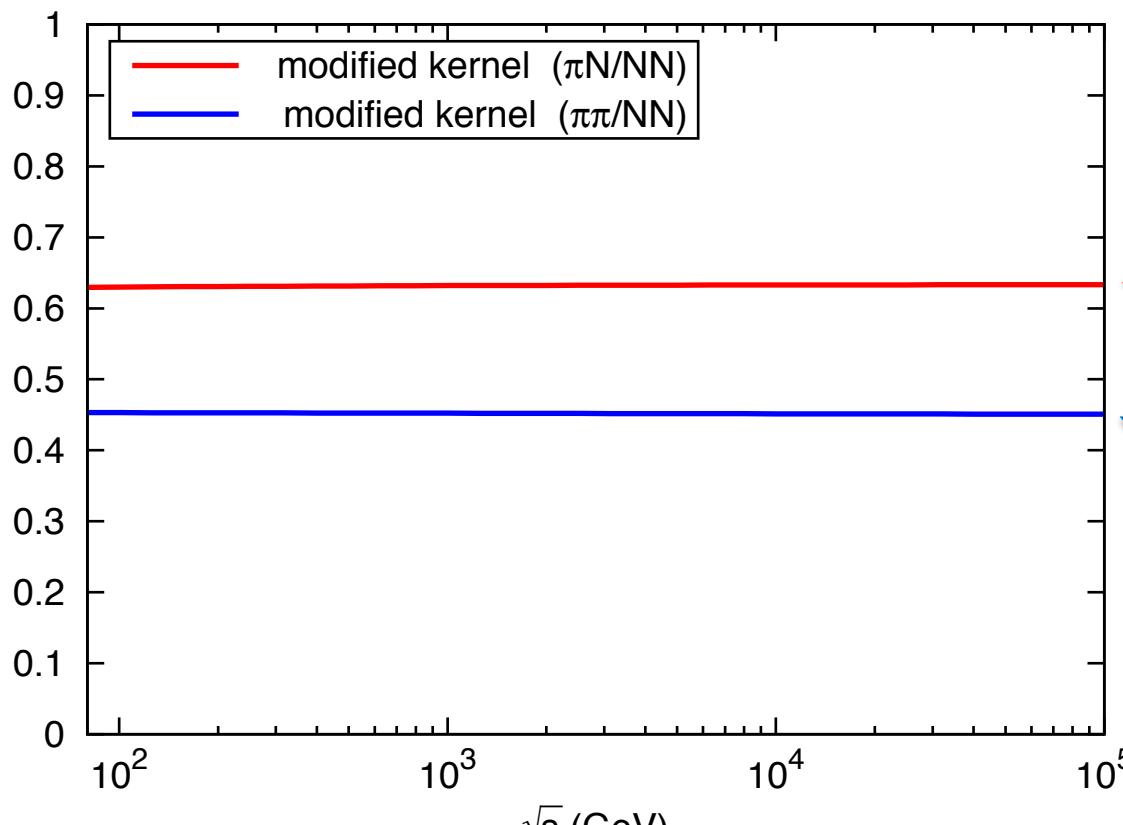


(a)



(b)

# Total cross section ratios



PREDICTIONS:

$$\left( \frac{\sigma_{tot}^{\pi N}}{\sigma_{tot}^{NN}} \right) \approx 0.63$$

$$\left( \frac{\sigma_{tot}^{\pi\pi}}{\sigma_{tot}^{NN}} \right) \approx 0.45$$

$$\left( \frac{\sigma_{tot}^{\pi^\pm p}}{\sigma_{tot}^{pp}} \right)_{DL(1992)} \approx 0.63$$

N.B. we do NOT need any additional parameter to calculate the pion involved cross sections!

# Summary

- We have studied total hadronic cross sections at high energies in the framework of holographic QCD.
- We have shown that all the TOTEM data, including the recent one at 13TeV, can be well reproduced with the model.
- A substantial deviation between our calculation and the empirical fit has been observed in the very large  $s$  region, where  $\text{sqrt}(s) > 10\text{TeV}$ .
- Realization of the saturation effect within the model may be important.