Causality constraint on bound states and scattering with zero-range force arXiv:1402.4973 [nucl-th].

do perturbative pions deserve another chance?

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# Outline

## Motivation

Chiral E**F**T of few-nucleon systems

## ≈ Light-by-light scattering sum rules

general principles: unitarity, causality, etc.

# Sero-range force: Bound state, tachyon, K-matrix pole using the sum rules as consistency (causality) criterion

phi^4 theory

≥ (Relativistic) Wigner's inequality

positive effective range parameters

# X Conclusions and outlook

Meson and 1-Baryon Sectors:	Few baryons, "Weinberg counting"	Few baryons, "Kaplan- Savage-Wise (KSW)" counting
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Converges ?	Maybe	Yes!	<b>No!</b> Another conceptual problem: 3N force goes from NLO to LO

## **Example 7 Example 7 Solution Second S**

$$M_{\lambda_1\lambda_2\lambda_3\lambda_4} = \varepsilon_{\lambda_4}^{*\mu_4}(\vec{q}_4) \varepsilon_{\lambda_3}^{*\mu_3}(\vec{q}_3) \varepsilon_{\lambda_2}^{\mu_2}(\vec{q}_2) \varepsilon_{\lambda_1}^{\mu_1}(\vec{q}_1) \mathcal{M}_{\mu_1\mu_2\mu_3\mu_4}$$
  
HELICITY AMPL. FEYNMAN AMPL.

IN THE FORWARD DIRECTION ( t = 0,  $s = 4\omega^2$ , u = -s.):  $\mathcal{M}_{\mu_1\mu_2\mu_3\mu_4} = A(s) g_{\mu_4\mu_2}g_{\mu_3\mu_1} + B(s) g_{\mu_4\mu_1}g_{\mu_3\mu_2} + C(s) g_{\mu_4\mu_3}g_{\mu_2\mu_1}$ ,

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1) CROSSING SYMMETRY (1 <-> 3, 2 <-> 4):

$$M_{+-+-}(s) = M_{++++}(-s), \quad M_{++--}(s) = M_{++--}(-s)$$

#### **AMPLITUDES WITH DEFINITE PARITY UNDER CROSSING:**

$$f^{(\pm)}(s) = M_{++++}(s) \pm M_{+-+-}(s)$$

$$g(s) = M_{++--}(s)$$

#### LbL sum rules

#### 2) CAUSALITY => ANALYTICITY => DISPERSION RELATIONS:

$$\operatorname{Re}\left\{\begin{array}{c}f^{(\pm)}(s)\\g(s)\end{array}\right\} = \frac{1}{\pi} \int_{-\infty}^{\infty} \frac{ds'}{s'-s} \operatorname{Im}\left\{\begin{array}{c}f^{(\pm)}(s')\\g(s')\end{array}\right\},$$

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3) OPTICAL THEOREM (UNITARITY):

$$\operatorname{Im} f^{(\pm)}(s) = -\frac{s}{8} \left[ \sigma_0(s) \pm \sigma_2(s) \right],$$
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$$\operatorname{Re} f^{(+)}(s) = -\frac{1}{2\pi} \int_{0}^{\infty} ds' \, s'^2 \, \frac{\sigma(s')}{s'^2 - s^2} \,, \qquad \sigma = (\sigma_0 + \sigma_2)/2 = (\sigma_{||} + \sigma_{\perp})/2$$
$$\operatorname{Re} f^{(-)}(s) = -\frac{s}{4\pi} \int_{0}^{\infty} ds' \, \frac{s' \, \Delta \sigma(s')}{s'^2 - s^2} \,, \qquad \Delta \sigma = \sigma_2 - \sigma_0$$
$$\operatorname{Re} g(s) = -\frac{1}{4\pi} \int_{0}^{\infty} ds' \, s'^2 \, \frac{\sigma_{||}(s') - \sigma_{\perp}(s')}{s'^2 - s^2} \,,$$

## Light-by-light scattering sum rules

4) "LOW-ENERGY THEOREM": 
$$\mathcal{L}_{EH} = c_1 (F_{\mu\nu} F^{\mu\nu})^2 + c_2 (F_{\mu\nu} \tilde{F}^{\mu\nu})^2,$$
  
 $f^{(+)}(s) = -2(c_1 + c_2)s^2 + O(s^4)$   
 $f^{(-)}(s) = O(s^5)$   
 $g(s) = -2(c_1 - c_2)s^2 + O(s^4)$ 

$$O(s^{1}): \qquad 0 = \int_{0}^{\infty} \frac{\mathrm{d}s}{s} \begin{bmatrix} \sigma_{2}(s) - \sigma_{0}(s) \end{bmatrix} \qquad \begin{array}{c} \text{GERASIMOV \& MOULIN, NPB (1976)} \\ \text{BRODSKY \& SCHMIDT, PLB (1995)} \end{array}$$

$$O(s^2): \qquad c_1 = \frac{1}{8\pi} \int_0^\infty \frac{\mathrm{d}s}{s^2} \,\sigma_{||}(s) \,,$$
$$c_2 = \frac{1}{8\pi} \int_0^\infty \frac{\mathrm{d}s}{s^2} \,\sigma_{\perp}(s)$$

V. P. & VANDERHAEGHEN, PRL (2010)

#### Zero-range force in light of the LbL sum rule

PAUK, V.P. & VANDERHAEGHEN, PLB 2014

Bubble-chain sum:





$$G(s) = -i \int \frac{d^4\ell}{(2\pi)^4} \frac{1}{\left[(p+\ell)^2 - m^2\right](\ell^2 - m^2)} \qquad \text{with } p^2 = s.$$

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 $\lambda > 0$ : no poles  $\lambda < 0$ : one pole and one K-matrix pole

$$T(s) = \frac{1}{K^{-1}(s) - i}$$



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but not the K-matrix pole...

#### **Phase shifts**

Levinson's theorem:  $\delta(0) = \pi N_{\text{bound states}}$ 



Figure 7: Phase shift for different values of  $\tilde{\lambda}$ .

The 90 degree crossing, i.e. the K-matrix pole does not correspond to any S-matrix pole in this case

## Wigner's causality bound

 $\begin{array}{c} r \leq 0 \\ \text{effective} \\ \text{range} \end{array}$ 

WIGNER, PHYS REV (1955)

PHILLIPS & COHEN, PLB (1997); HAMMER & D. LEE, ANN PHYS (2010); ...

#### Wigner's causality bound



#### Non-relativistic limit



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