## Causality constraint on bound states and scattering with zero-range force arXiv: I 402.4973 [nucl-th].

do perturbative pions deserve another chance?

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4SFB쿨
THE LOW-ENERGY FRONTIER OF THE STANDARD MODEL
@ 8th Chiral Dynamics Workshop, Pisa, Italy, June 29, 2015

## Outline

## Motivation

Chiral EFT of few-nucleon systems
« Light-by-light scattering sum rules
general principles: unitarity, causality, etc.
$\propto^{\circ}$ Zero-range force:
Bound state, tachyon, K-matrix pole
using the sum rules as consistency (causality) criterion phi^4 theory
$\geq$ (Relativistic) Wigner's inequality
positive effective range parameters

X Conclusions and outlook

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| Converges ? | Maybe | Yes! | No! <br> Another conceptual problem: 3 N force goes from NLO to LO |

## $\approx$ Light by light scattering

$$
M_{\lambda_{1} \lambda_{2} \lambda_{3} \lambda_{4}}=\varepsilon_{\lambda_{4}}^{* \mu_{4}}\left(\vec{q}_{4}\right) \varepsilon_{\lambda_{3}}^{* \mu_{3}}\left(\vec{q}_{3}\right) \varepsilon_{\lambda_{2}}^{\mu_{2}}\left(\vec{q}_{2}\right) \varepsilon_{\lambda_{1}}^{\mu_{1}}\left(\vec{q}_{1}\right) \mathcal{M}_{\mu_{1} \mu_{2} \mu_{3} \mu_{4}}
$$

HELICITY AMPL.
FEYNMAN AMPL.

## IN THE FORWARD DIRECTION $\left(t=0, \quad s=4 \omega^{2}, \quad u=-s.\right)$ :

$$
\begin{gathered}
\mathcal{M}_{\mu_{1} \mu_{2} \mu_{3} \mu_{4}}=A(s) g_{\mu_{4} \mu_{2}} g_{\mu_{3} \mu_{1}}+B(s) g_{\mu_{4} \mu_{1}} g_{\mu_{3} \mu_{2}}+C(s) g_{\mu_{4} \mu_{3}} g_{\mu_{2} \mu_{1}}, \\
M_{++++}(s)=A(s)+C(s) \\
M_{+-+-}(s)=A(s)+B(s) \\
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1) CROSSING SYMMETRY ( $1<->3,2<->4$ ):

$$
M_{+-+-}(s)=M_{++++}(-s), \quad M_{++--}(s)=M_{++--}(-s)
$$

AMPLITUDES WITH DEFINITE PARITY UNDER CROSSING:

$$
\begin{aligned}
& f^{( \pm)}(s)=M_{++++}(s) \pm M_{+-+-}(s) \\
& g(s)=M_{++--}(s)
\end{aligned}
$$

## LbL sum rules

2) CAUSALITY $=>$ ANALYTICITY $=>$ DISPERSION RELATIONS:

$$
\operatorname{Re}\left\{\begin{array}{l}
f^{( \pm)}(s) \\
g(s)
\end{array}\right\}=\frac{1}{\pi} f_{-\infty}^{\infty} \frac{d s^{\prime}}{s^{\prime}-s} \operatorname{Im}\left\{\begin{array}{l}
f^{( \pm)}\left(s^{\prime}\right) \\
g\left(s^{\prime}\right)
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3) OPTICAL THEOREM (UNITARITY):

$$
\begin{aligned}
\operatorname{Im} f^{( \pm)}(s) & =-\frac{s}{8}\left[\sigma_{0}(s) \pm \sigma_{2}(s)\right] \\
\operatorname{Im} g(s) & =-\frac{s}{8}\left[\sigma_{\|}(s)-\sigma_{\perp}(s)\right]
\end{aligned}
$$

$\sigma_{0,2}\left(\sigma_{\|, \perp}\right) \begin{aligned} & \text { ARE CIRCULARLY (LINEARLY) POLARIZED PHOTON-PHOTON FUSION CROSS- } \\ & \text { SECTIONS }\end{aligned}$

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$$
\begin{array}{cl}
\operatorname{Re} f^{(+)}(s)=-\frac{1}{2 \pi} \int_{0}^{\infty} d s^{\prime} s^{\prime 2} \frac{\sigma\left(s^{\prime}\right)}{s^{\prime 2}-s^{2}}, & \sigma=\left(\sigma_{0}+\sigma_{2}\right) / 2=\left(\sigma_{\|}+\sigma_{\perp}\right) / 2 \\
\operatorname{Re} f^{(-)}(s)=-\frac{s}{4 \pi} \int_{0}^{\infty} d s^{\prime} \frac{s^{\prime} \Delta \sigma\left(s^{\prime}\right)}{s^{\prime 2}-s^{2}}, & \Delta \sigma=\sigma_{2}-\sigma_{0} \\
\operatorname{Re} g(s)=-\frac{1}{4 \pi} \int_{0}^{\infty} d s^{\prime} s^{\prime 2} \frac{\sigma_{\| \mid}\left(s^{\prime}\right)-\sigma_{\perp}\left(s^{\prime}\right)}{s^{\prime 2}-s^{2}}, & \\
\text { Vadimir Pascalutsa "Zero range interactions" @ chiral Dynamics Workshop Pisa June 29, 2015 }
\end{array}
$$

## Light-by-light scattering sum rules

4) "LOW-ENERGY THEOREM": $\quad \mathcal{L}_{\mathrm{EH}}=c_{1}\left(F_{\mu \nu} F^{\mu \nu}\right)^{2}+c_{2}\left(F_{\mu \nu} \tilde{F}^{\mu \nu}\right)^{2}$,

$$
\begin{array}{r}
f^{(+)}(s)=-2\left(c_{1}+c_{2}\right) s^{2}+O\left(s^{4}\right) \\
f^{(-)}(s)=O\left(s^{5}\right) \\
g(s)=-2\left(c_{1}-c_{2}\right) s^{2}+O\left(s^{4}\right)
\end{array}
$$

LOW-ENERGY EXPANSION
$O\left(s^{1}\right):$

$$
\begin{array}{ll}
0 & =\int_{0}^{\infty} \frac{\mathrm{d} s}{s}\left[\sigma_{2}(s)-\sigma_{0}(s)\right] \quad \begin{array}{c}
\text { GERASIMOV \& MOULIN, NPB (1976) } \\
\text { BRODSKY \& SCHMIDT, PLB (1995) }
\end{array} \\
c_{1}=\frac{1}{8 \pi} \int_{0}^{\infty} \frac{\mathrm{d} s}{s^{2}} \sigma_{\|}(s), & \text { V.P. \& VANDERHAEGHEN, PRL (2010) }
\end{array} c_{2}=\frac{1}{8 \pi} \int_{0}^{\infty} \frac{\mathrm{d} s}{s^{2}} \sigma_{\perp}(s) \quad . \quad .
$$

## Zero-range force in light of the LbL sum rule

PAUK, V.P. \& VANDERHAEGHEN, PLB 2014

$$
\begin{gathered}
T=V+V G T \\
V=\lambda
\end{gathered}
$$

$$
T(s)=\frac{1}{\lambda^{-1}-G(s)}
$$


$G(s)=-i \int \frac{d^{4} \ell}{(2 \pi)^{4}} \frac{1}{\left[(p+\ell)^{2}-m^{2}\right]\left(\ell^{2}-m^{2}\right)} \quad$ with $p^{2}=s$.

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Bubble-chain sum:

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$$

$$
\text { with } p^{2}=s
$$



$$
\lambda>0: \text { no poles }
$$

$$
\lambda<0 \text { : one pole and one K-matrix pole }
$$

## Light-by-light sum rule as causality criterion

$$
\int_{s_{0}}^{\infty} \mathrm{d} s \frac{\Delta \sigma(s)}{s}=0, \quad \mathcal{L}=\left(D^{\mu} \phi\right)^{*} D_{\mu} \phi-m^{2} \phi^{*} \phi+\frac{\lambda}{4}\left(\phi^{*} \phi\right)^{2}-\frac{1}{4} F^{\mu \nu} F_{\mu \nu}
$$



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To cancel the integral one need to introduce the bound state as the asymptotic state i.e., new channel:


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To cancel the integral one need to introduce the bound state as the asymptotic state i.e., new channel:

but not the K-matrix pole...

## Phase shifts

Levinson's theorem:
$\delta(0)=\pi N_{\text {bound states }}$


Figure 7: Phase shift for different values of $\tilde{\lambda}$.
The 90 degree crossing, i.e. the K-matrix pole does not correspond to any S-matrix pole in this case

## Wigner's causality bound

$r \leq 0$<br>effective<br>range

[^0]
## Wigner's causality bound

$$
r \leq 0
$$

effective range

WIGNER, PHYS REV (1955)
PHILLIPS \& COHEN, PLB (1997);
HAMMER \& D. LEE, ANN PHYS (2010); ..

$$
|\mathbf{k}| \cot \delta(s)=-\frac{1}{a}+\frac{1}{2} \sum_{n=1}^{\infty}(-1)^{n+1} r_{n}|\mathbf{k}|^{2 n}
$$



In the the tachyon (acausal) regime at least one of the effective range parameters is negative.

Therefore our causality criterion yields:

$$
r_{n} \geq 0
$$

## Non-relativistic limit



In non-rel. limi, K-matrix pole disappears and

$$
r=0
$$

in agreement with Wigner's bound

Zero eff. range of 2-body force eventually leads to the problems with the 3-body force
[BEDAQUE, HAMMER, VAN KOLCK (1999)]


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