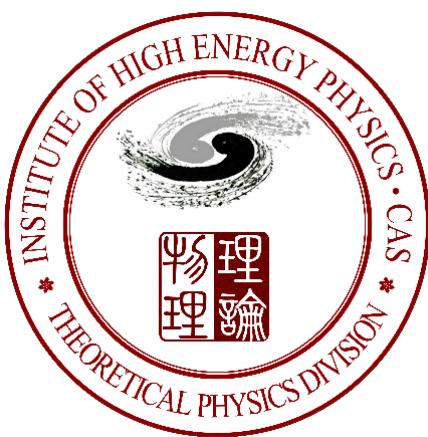


QCD phase diagram and CEP

Mei Huang

Theoretical Physics Division

Institute of High Energy Physics, CAS



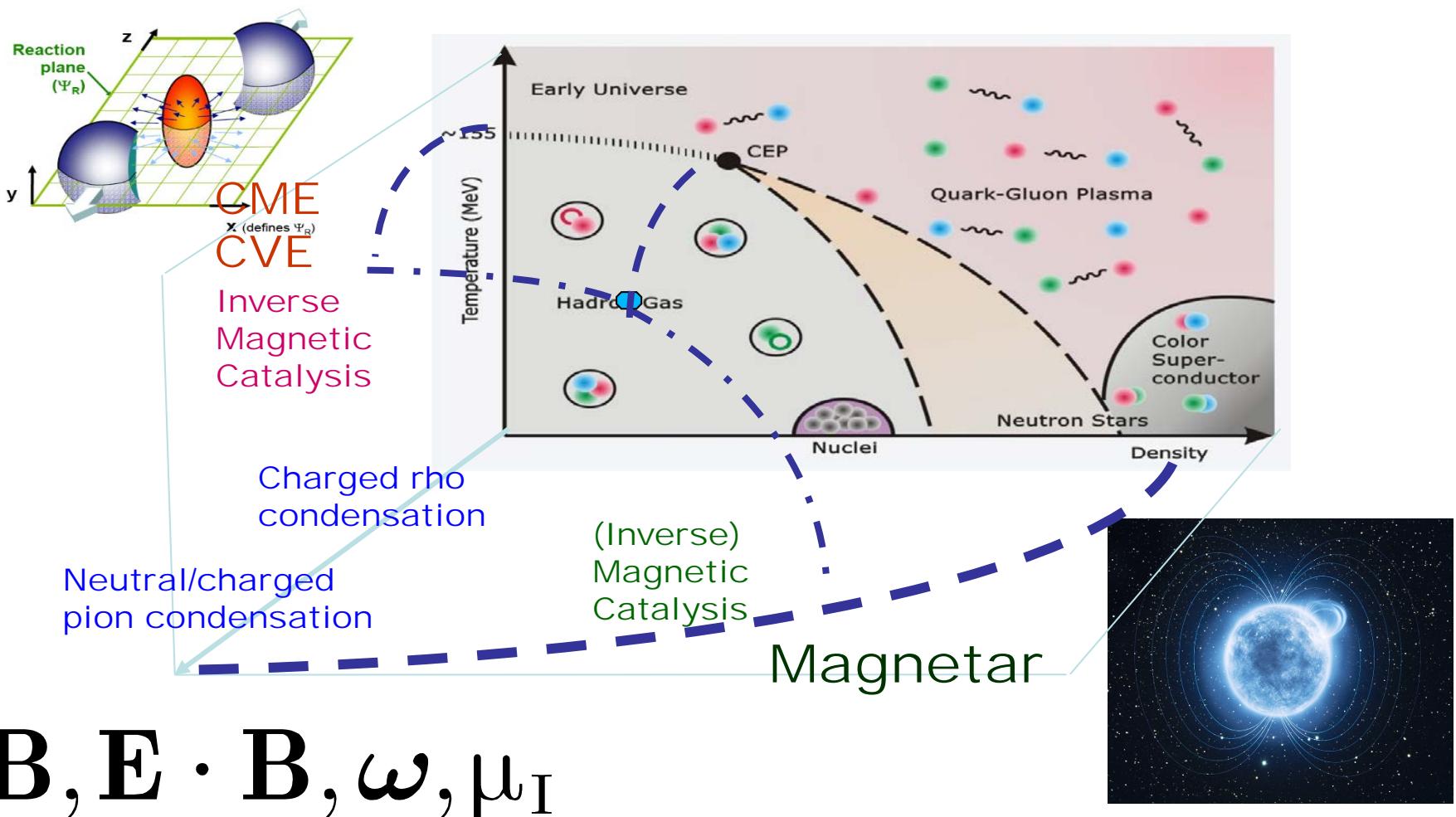
Content

I. Introduction on QCD phase diagram

II. Searching for the QCD CEP

III. Conclusion and discussion

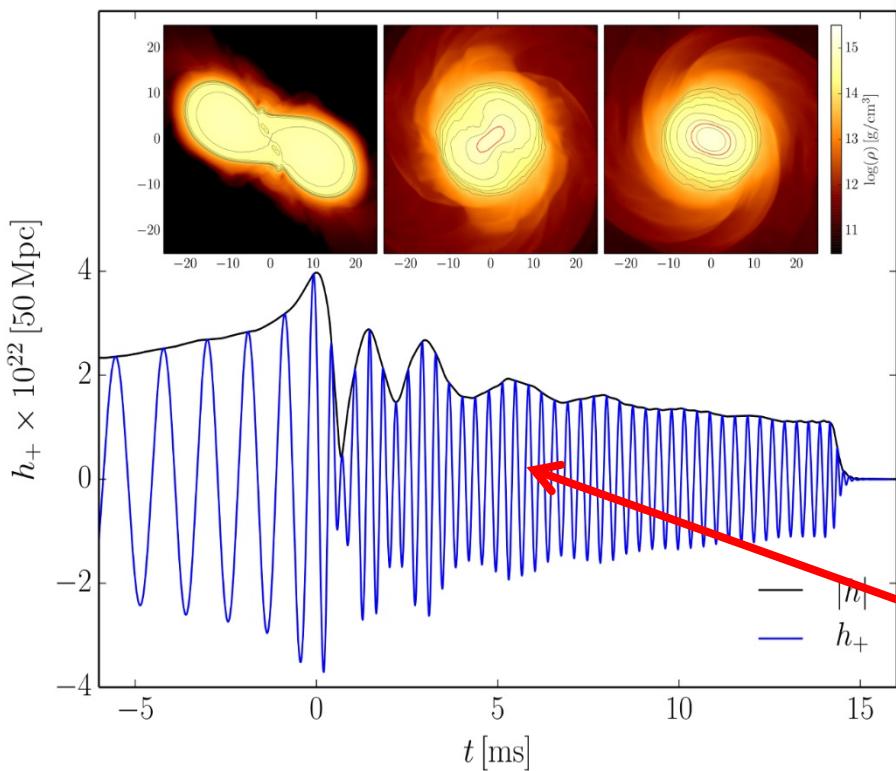
Explored QCD phase diagram now by theorists



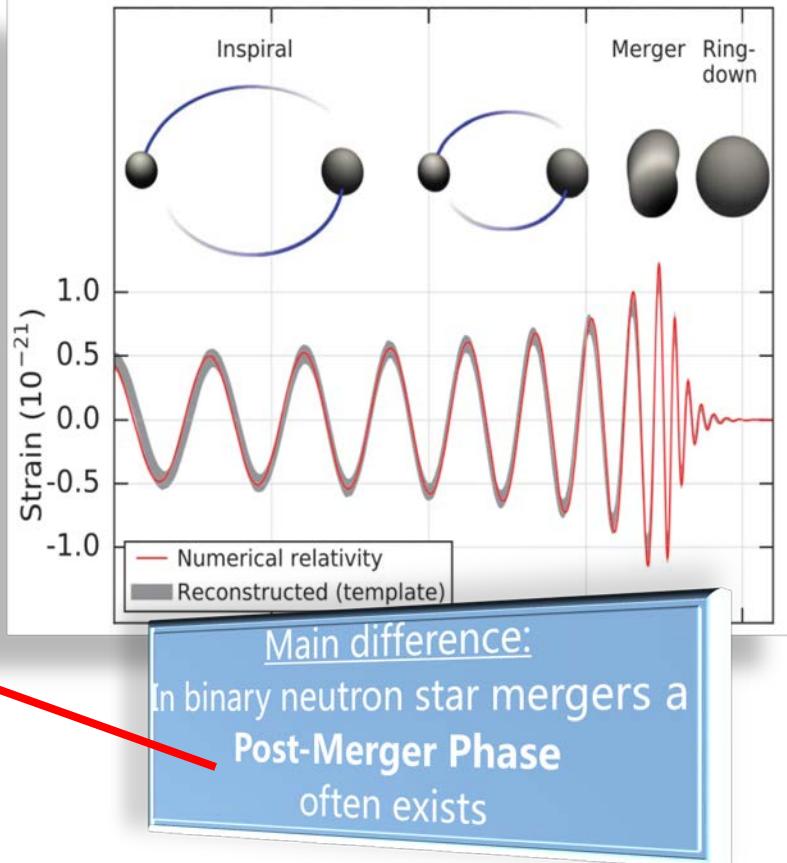
Gravitational Waves from Neutron Star Mergers

GW170817

Neutron Star Collision (Simulation)



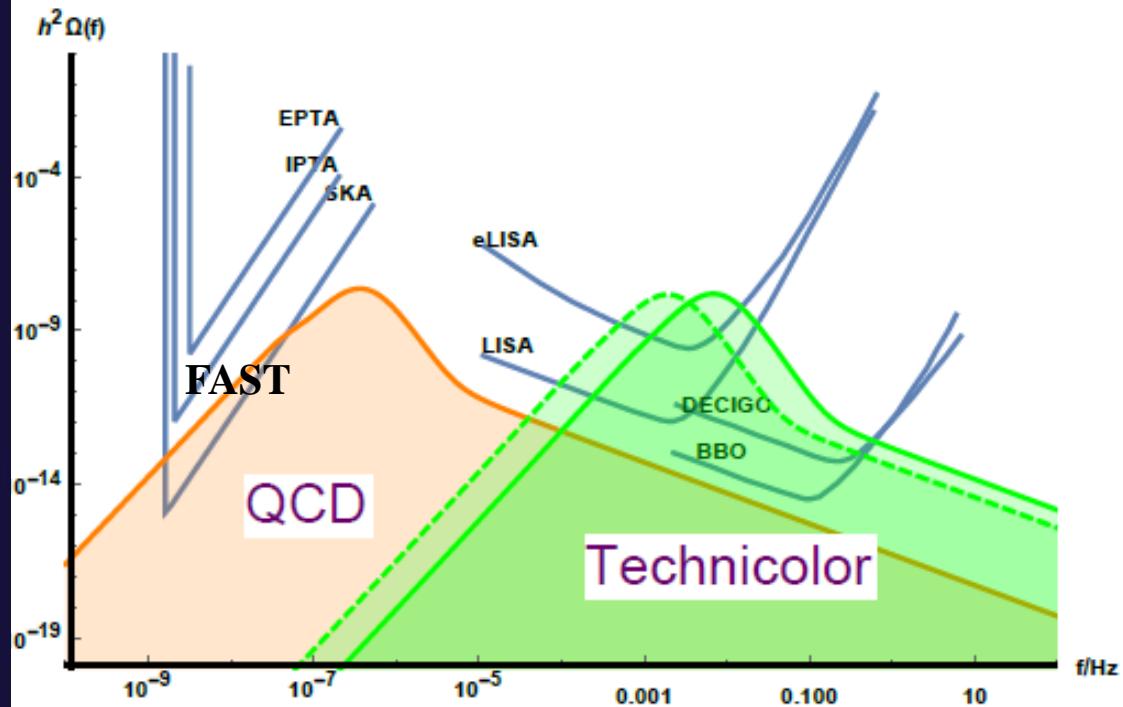
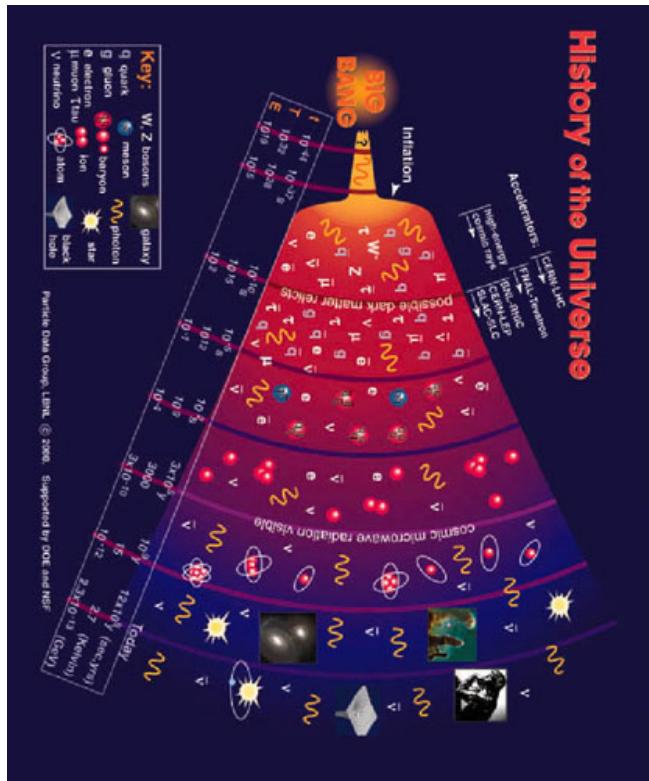
Collision of two Black Holes



Gravitation wave from QCD & electroweak phase transitions in the early universe

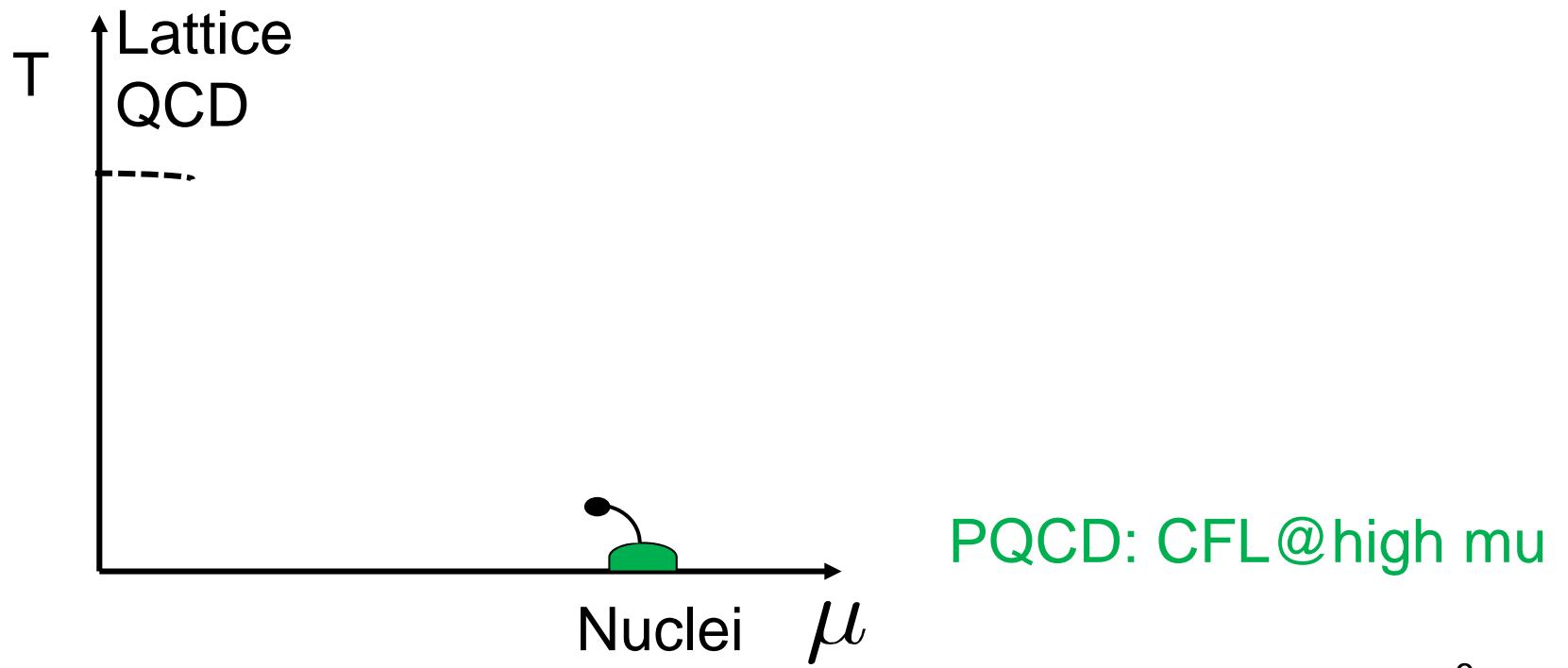
1st order phase transition for pure gluon system!

Yidian Chen, Mei Huang, Qishu Yan,
arXiv:1712.03470, JHEP05(2018)178



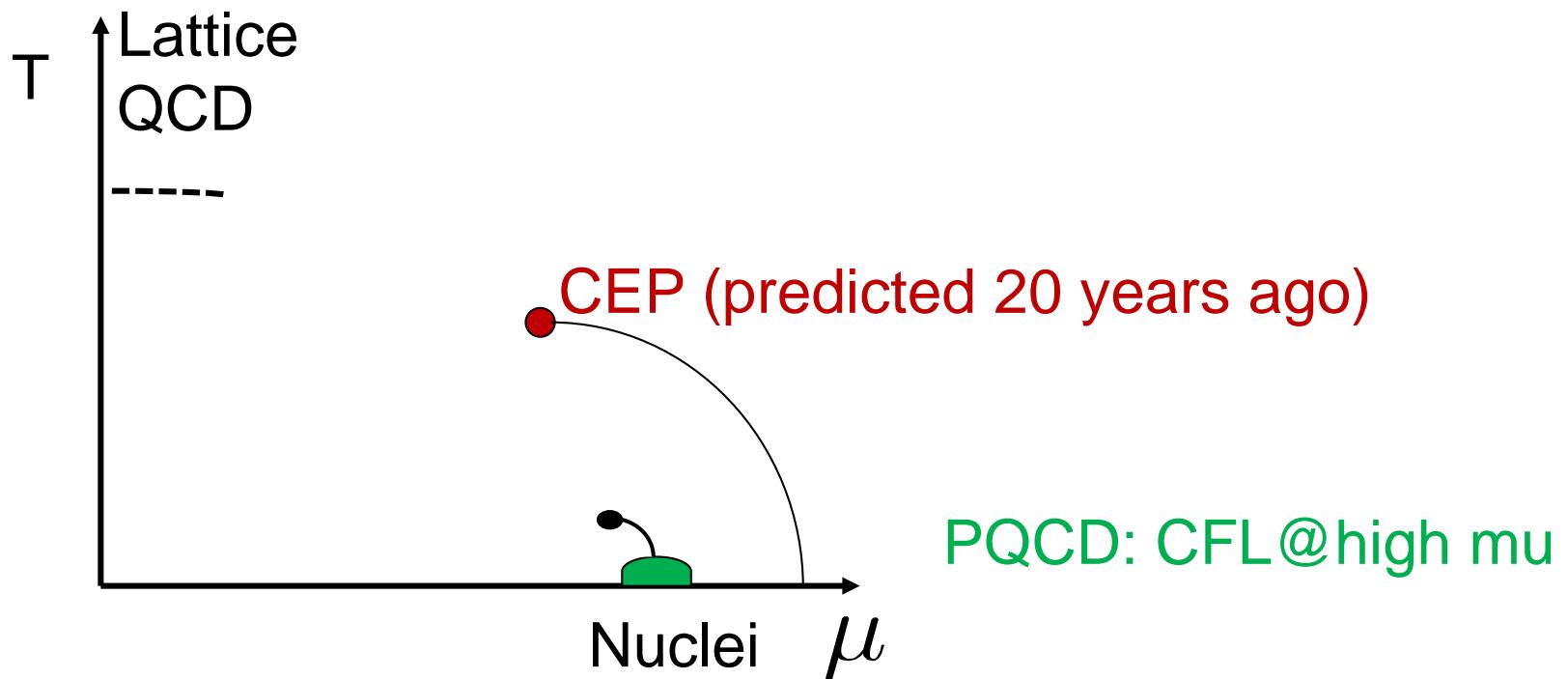
Confirmed QCD phase diagram

PQCD: QGP@High T



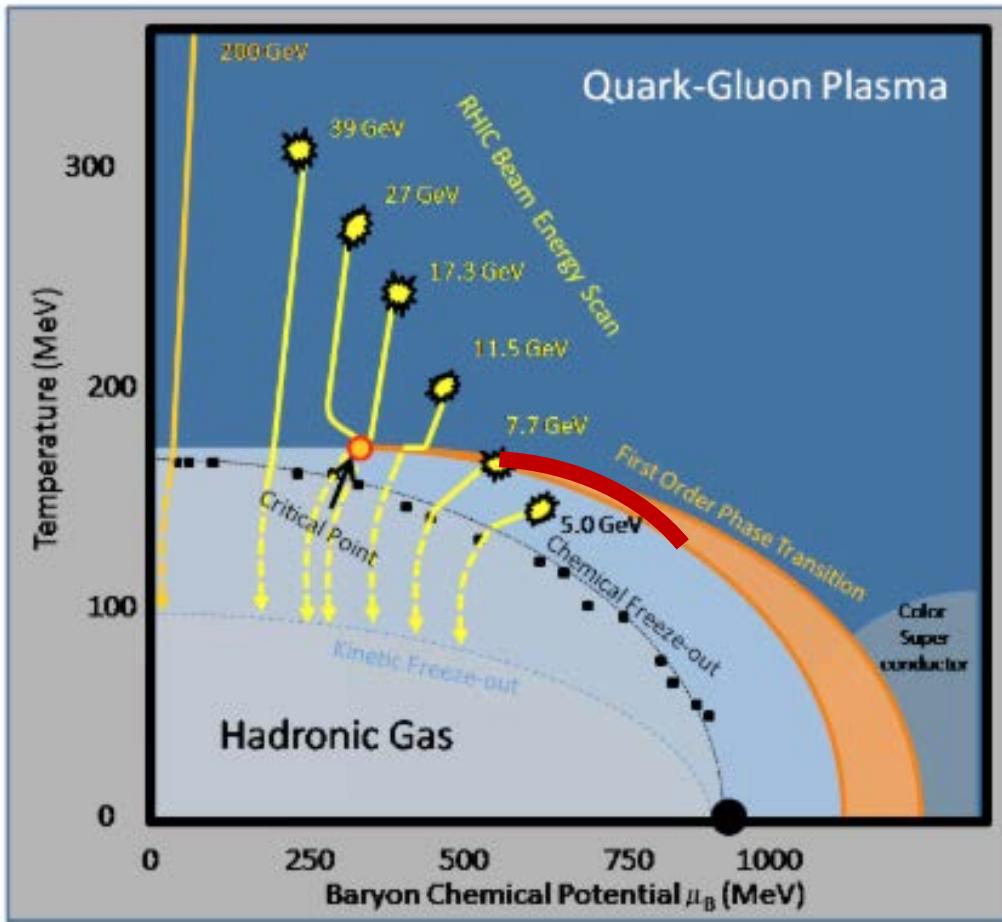
Searching for the QCD CEP

PQCD: QGP@High T



Locating CEP is essential for the QCD phase diagram!

Locating the QCD CEP



- BES @ RHIC
- NICA @ Dubna
- CBM @ FAIR
- HIAF @ IMP

Chiral and deconfinement phase transitions

CEP is for chiral phase transition!

Chiral phase transition:

quark-antiquark condensate (for m=0)

Chiral symmetry breaking: $\langle \bar{\psi}\psi \rangle \neq 0$

Chiral symmetry restoration: $\langle \bar{\psi}\psi \rangle = 0$

Deconfinement phase transition:

referring to the “permanent confinement”

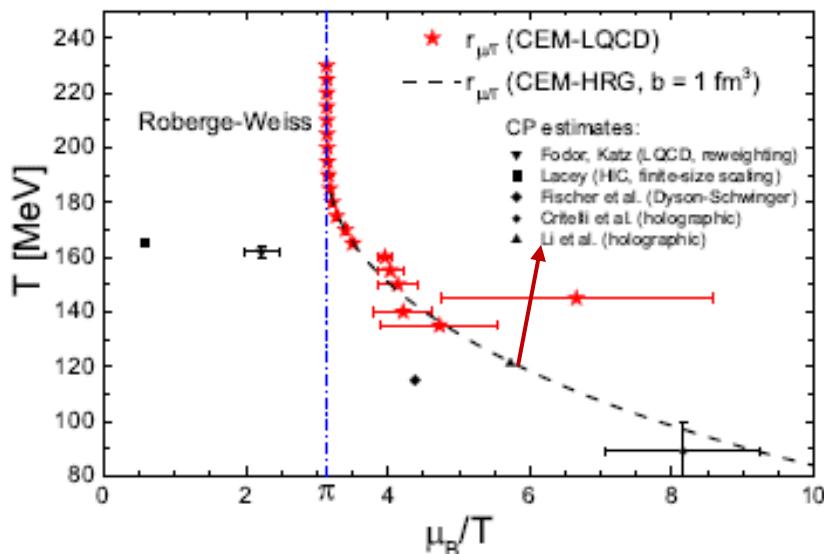
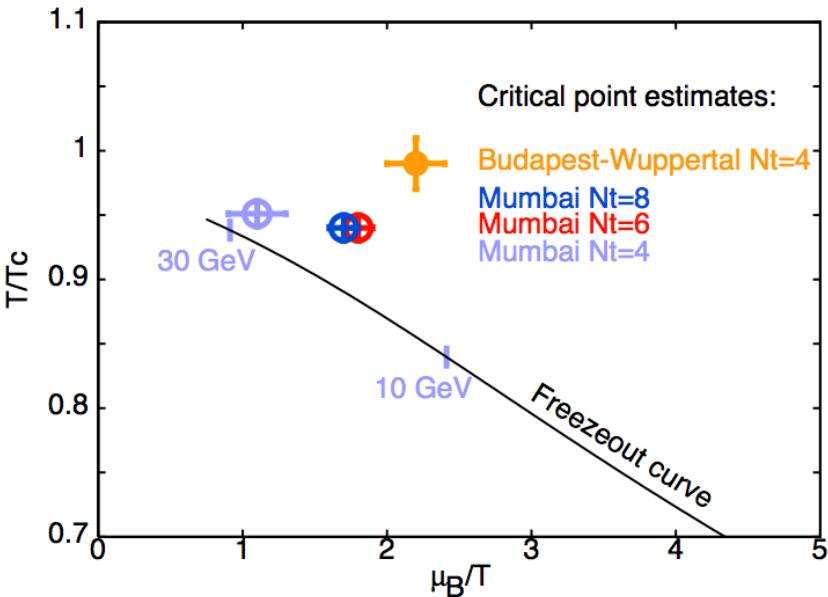
Polyakov loop (for m= infinity)

$$L(\vec{x}) = \frac{1}{N_c} \text{tr } \mathcal{P}(\vec{x}) \text{ with } \mathcal{P}(\vec{x}) = P e^{ig \int_0^\beta dt A_0(t, \vec{x})}$$
$$\langle L(\vec{x}) \rangle \sim \exp(-\beta F_q)$$

Confinement: center symmetric $\langle L \rangle = 0$ $F_q \rightarrow \infty$

Deconfinement: center symmetry breaking $\langle L \rangle \neq 0$, $F_q < \infty$

Location of CEP from Lattice QCD

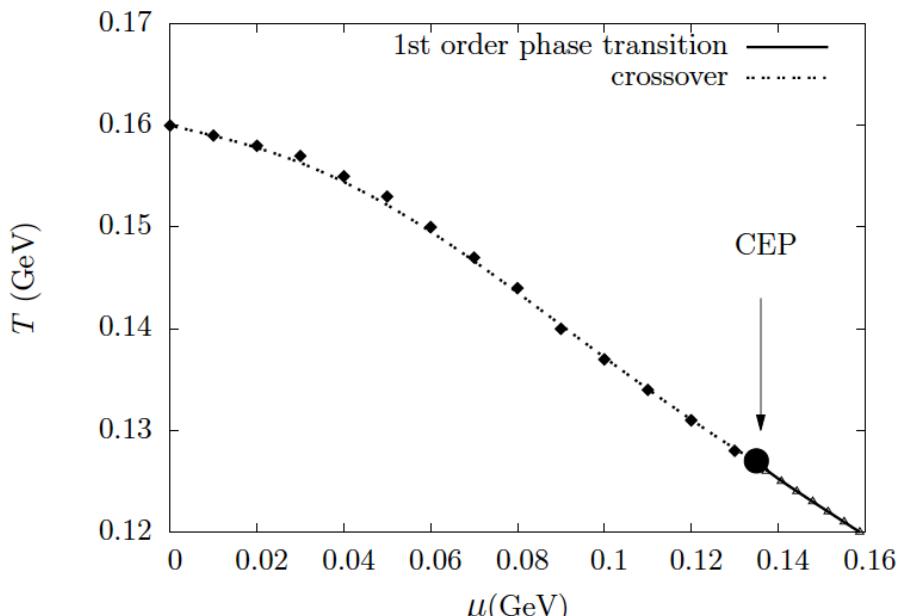


- 1) Fodor&Katz, JHEP 0404,050 (2004).
 $(\mu_E^E, T_E) = (360, 162) \text{ MeV}$
- 2) Gavai&Gupta, NPA 904, 883c (2013)
 $(\mu_E^E, T_E) = (279, 155) \text{ MeV}$
- 3) F. Karsch (CPOD2016)
 $\mu_E^E / T_E > 2$
- 4) V. Vovchenko, J. Steinheimer, O. Philipsen, H. Stoecker, arXiv:1711.01261

$$\mu_B^E / T_E > \pi$$

Latest lattice calculation shows that small baryon number density region for CEP is ruled out!

Location of CEP from DSE



1): Y. X. Liu, et al., PRD90, 076006 (2014).
 $(\mu_B^E, T_E) = (372, 129)$ MeV

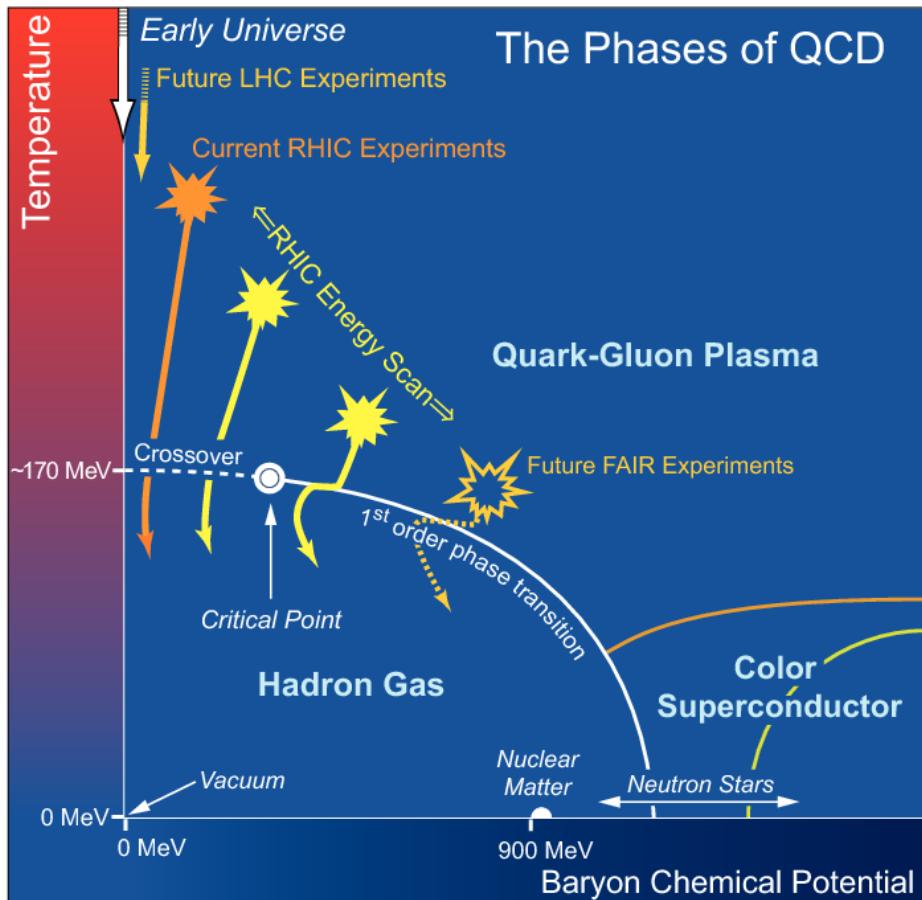
2): Hong-shi Zong et al., JHEP 07, 014 (2014).
 $(\mu_B^E, T_E) = (405, 127)$ MeV

3): C. S. Fischer et al., PRD90, 034022 (2014).
 $(\mu_B^E, T_E) = (504, 115)$ MeV

$$\mu_B = 3 \mu_q$$

baryon number density region 300-500 MeV

Searching for the QCD CEP



BES Phase-I

$\sqrt{s_{NN}}$ (GeV)	Events (10^6)	Year	$*\mu_B$ (MeV)	$*T_{CH}$ (MeV)
200	350	2010	25	166
62.4	67	2010	73	165
39	39	2010	112	164
27	70	2011	156	162
19.6	36	2011	206	160
14.5	20	2014	264	156
11.5	12	2010	316	152
7.7	4	2010	422	140

Higher Order Fluctuations of Conserved Quantities

$$\chi_n^B = \frac{\partial^n [P/T^4]}{\partial [\mu_B/T]^n} \quad B \rightarrow Q, s$$

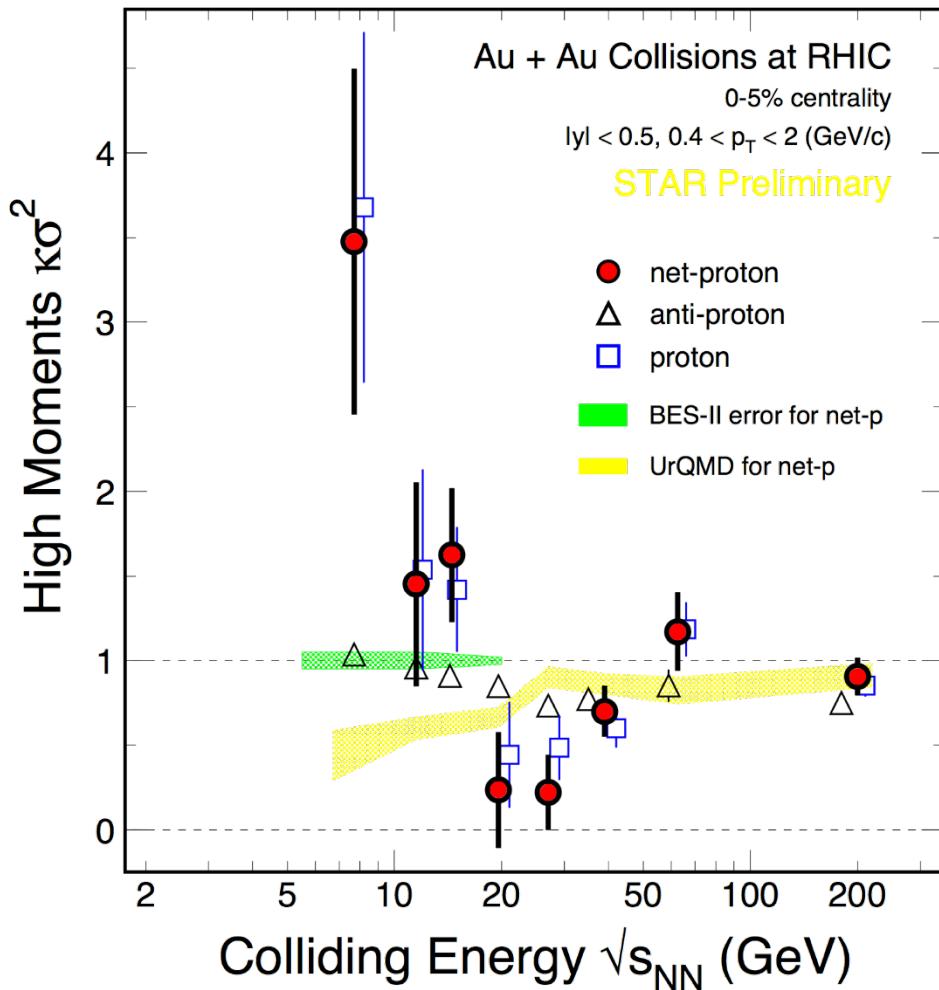
$$C_n^B = VT^3 \chi_n^B$$

$$\frac{\sigma^2}{M} = \frac{C_2^B}{C_1^B} = \frac{\chi_2^B}{\chi_1^B}, \quad S\sigma = \frac{C_3^B}{C_2^B} = \frac{\chi_3^B}{\chi_2^B},$$

$$\frac{S\sigma^3}{M} = \frac{C_3^B}{C_1^B} = \frac{\chi_3^B}{\chi_1^B}, \quad \boxed{\kappa\sigma^2 = \frac{C_4^B}{C_2^B} = \frac{\chi_4^B}{\chi_2^B}.}$$

*S. Ejiri et al, Phys.Lett. B 633 (2006) 275. Cheng et al, PRD (2009) 074505. B. Friman et al., EPJC 71 (2011) 1694. F. Karsch and K. Redlich , PLB 695, 136 (2011).
 S. Gupta, et al., Science, 332, 1525(2012). A. Bazavov et al., PRL109, 192302(12)
 S. Borsanyi et al., PRL111, 062005(13), P. Alba et al., arXiv:1403.4903*

Measurement of Higher Order Fluctuations of Conserved Quantities



Non-monotonic trend is observed for the 0-5% most central Au+Au collisions. Dip structure is observed around 19.6 GeV.

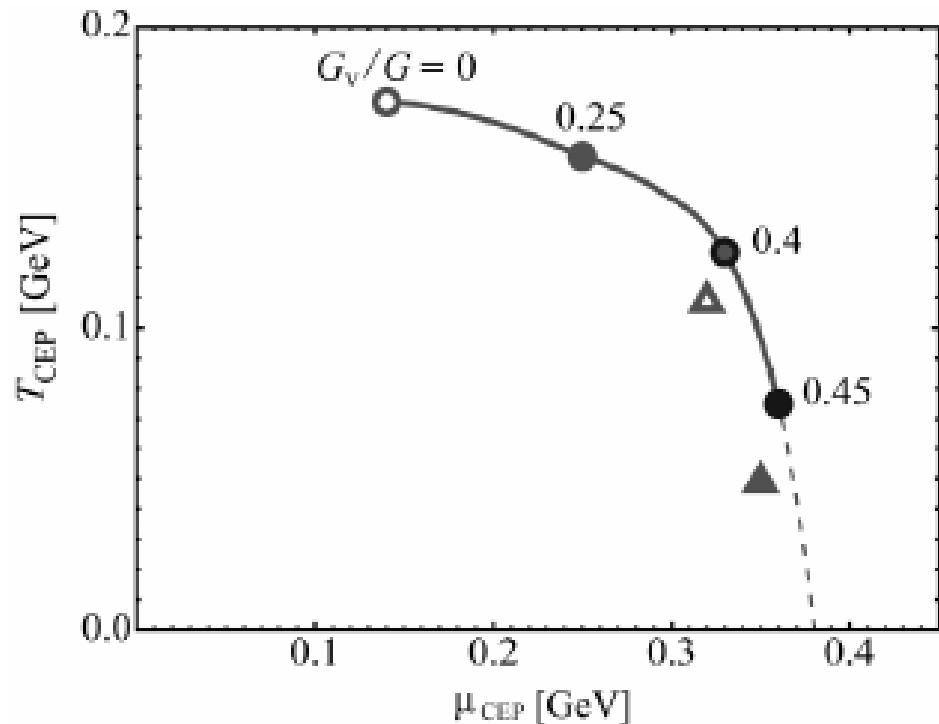
What information about CEP can be read from experimental measurement?

CEP from PNJL-like models

Z.B Li, K.Xu,X.Y.Wang, M.Huang
arXiv:1801.09215

Location of CEP: NJL

NJL, PNJL, Nonlocal NJL,



P.F.Zhuang,M.Huang,
Y.X.Liu,W.J.Fu,Z.Zhang
H.S.Zong,X.Luo,G.Y.Shao.....
J.Deng,J.W.Chen,G.Q.Cao,
X.G.Huang.....

Weise,
Klevansky,
Hatsuda,Kunihiro,
Fukushima,
Redlich,Sasaki,
Ratti,

.....

$$\mu_B = 3 \mu_q$$

Hell, Kashiwa, Weise

Journal of Modern Physics, 2013, 4, 644-650

from small to high baryon number density region

NJL model

$$\mathcal{L}_{NJL} = \bar{\psi}(i\gamma_\mu D^\mu - m)\psi + G_S[(\bar{\psi}\psi)^2 + (\bar{\psi}i\gamma_5\vec{\tau}\psi)^2] - \textcircled{G_V}(\bar{\psi}\gamma_\mu\psi)^2$$

PNJL model K.Fukushima

Shift the location of CEP

$$\frac{\mathcal{U}(\Phi, \bar{\Phi}, T)}{T^4} = -\frac{a(T)}{2}\bar{\Phi}\Phi + b(T) \ln[1 - 6\bar{\Phi}\Phi + 4(\bar{\Phi}^3 + \Phi^3) - 3(\bar{\Phi}\Phi)^2]$$

Mimic gluodynamics

$$a(T) = a_0 + a_1\left(\frac{T_0}{T}\right) + a_2\left(\frac{T_0}{T}\right)^2 \quad b(T) = b_3\left(\frac{T_0}{T}\right)^3$$

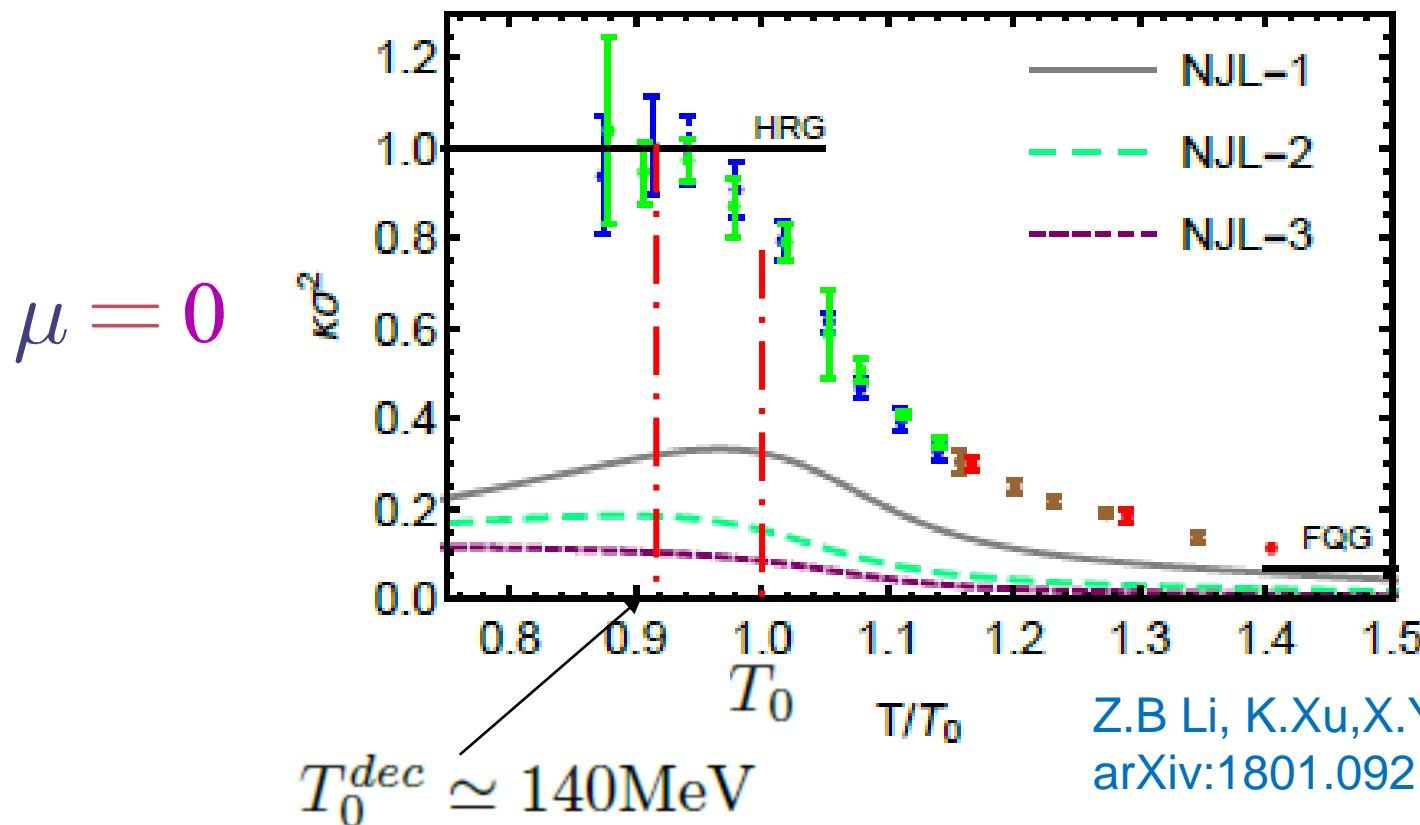
muPNJL model

$$T_0(N_f, \mu_i) = T_\tau e^{-\frac{1}{\alpha_0 f(N_f, \mu_i)}}$$

$$f(N_f, \mu_i) = \frac{11N_c - 2N_f}{6\pi} - \kappa \frac{16N_f}{\pi} \frac{\mu^2}{T_\tau^2}$$

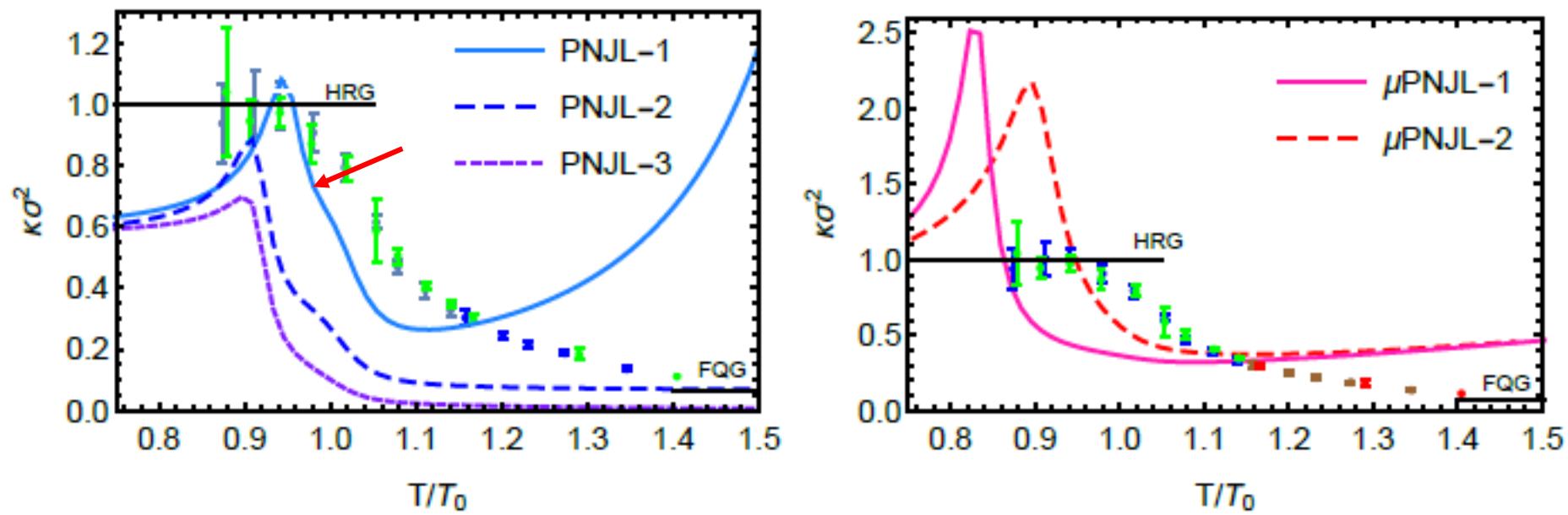
**Observation from lattice result: 0.8 of BNF at chiral phase transition, 1 at hadron-QGP transition
Unexpected results: Dominant contribution (80-90%) from gluodynamics to baryon number fluctuation!!!**

Lattice result: A. Bazavov et al
Phys. Rev. D95 no. 5, (2017) 054504



a physical “confinement-deconfinement” transition

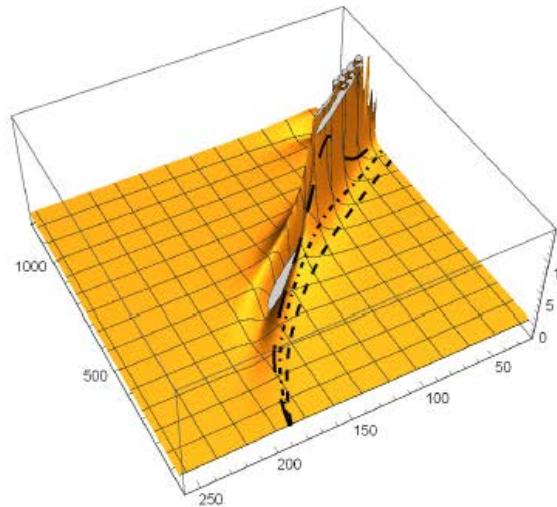
Lattice results for BNF at mu=0 can constrain models



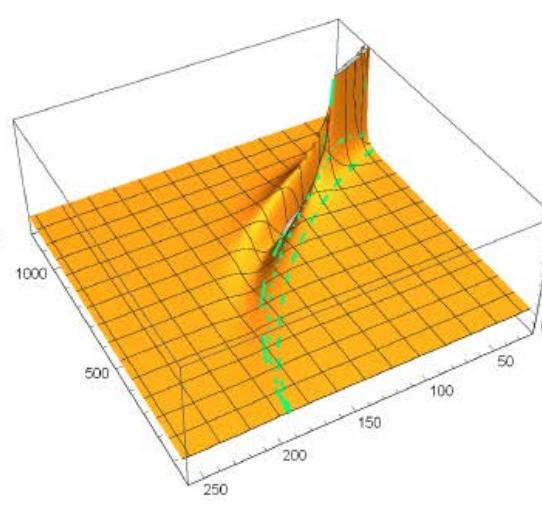
Z.B Li, K.Xu,X.Y.Wang, M.Huang
arXiv:1801.09215

Shifting the location of CEP in the NJL model

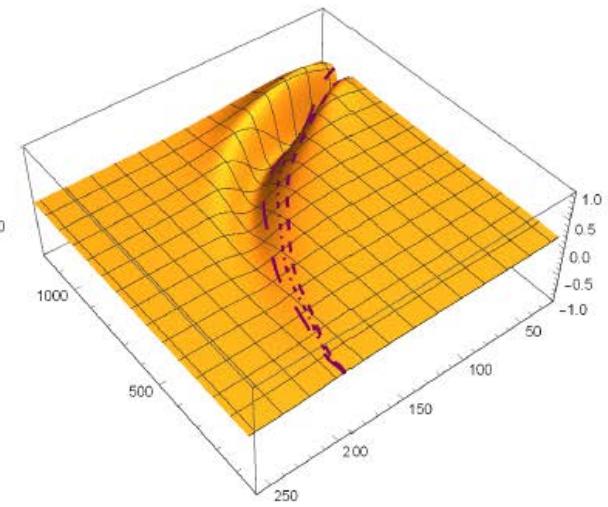
$$G_V = -0.5G_S$$



$$G_V = 0$$



$$G_V = 0.67G_S$$



NJL-1

NJL-2

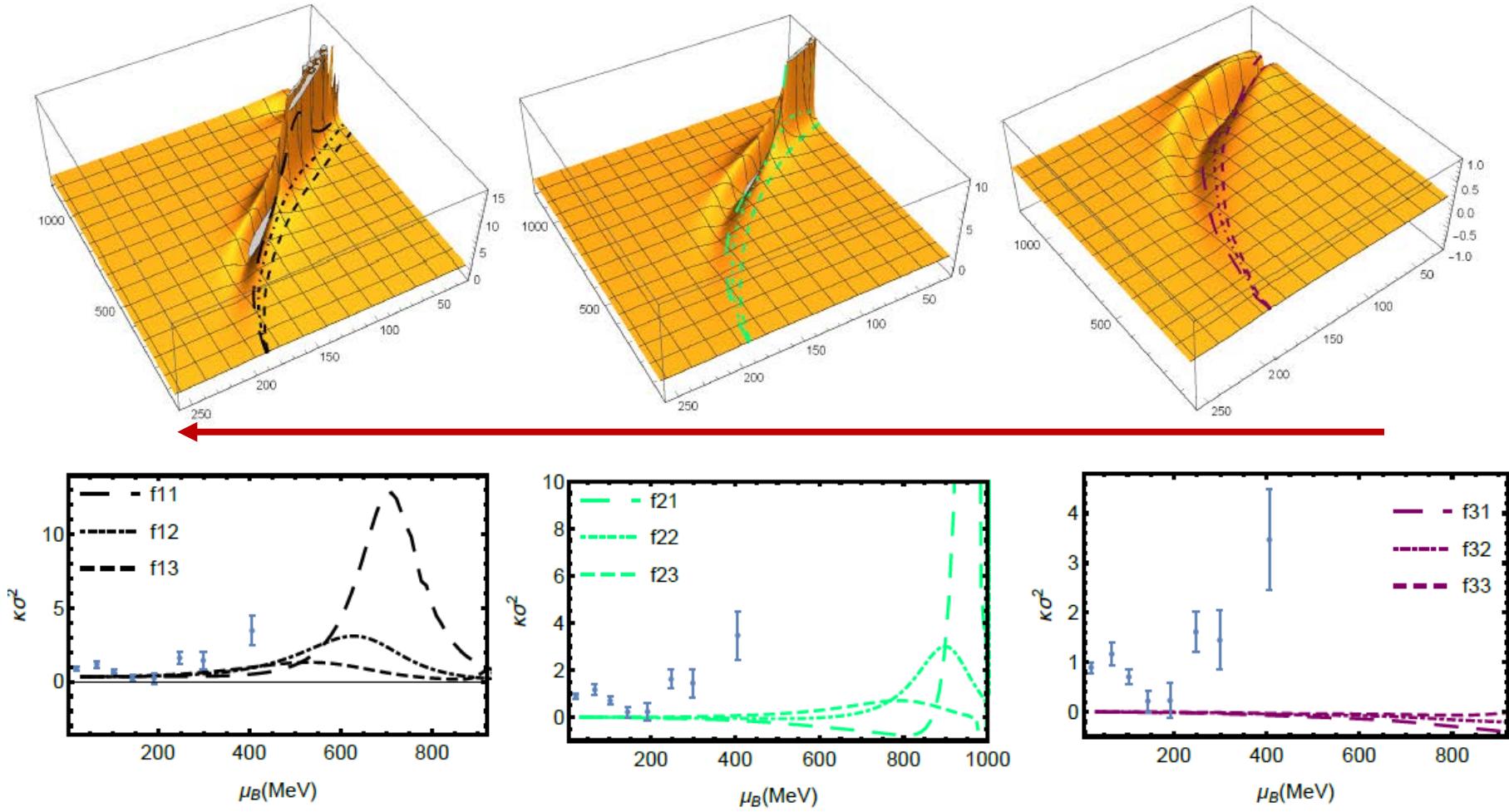
NJL-3

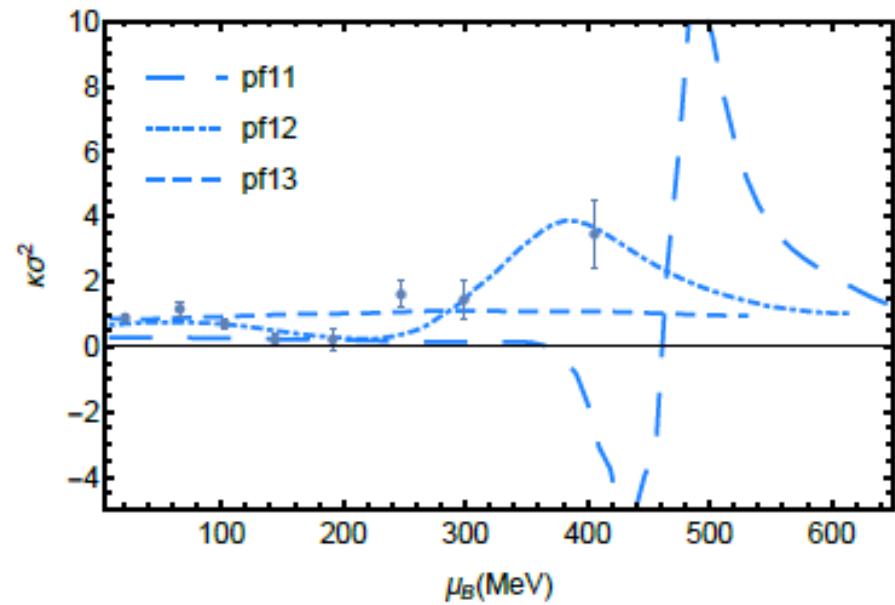
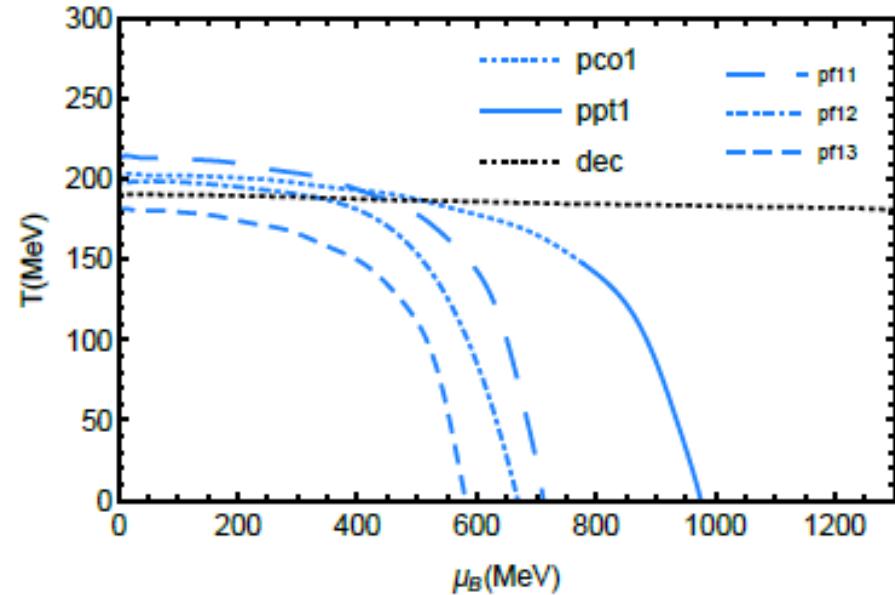
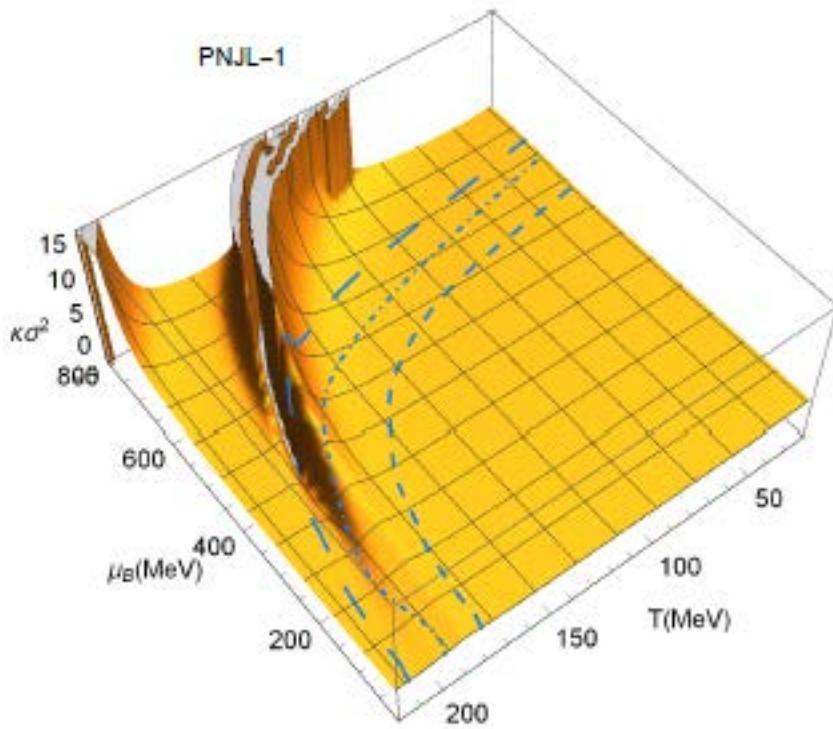


Z.B Li, K.Xu,X.Y.Wang, M.Huang
arXiv:1801.09215

Results from NJL model

1. Location of CEP peak determines the location of the peak for BNF along freeze-out line;
2. If no CEP, no structure for BNF along freeze-out line.



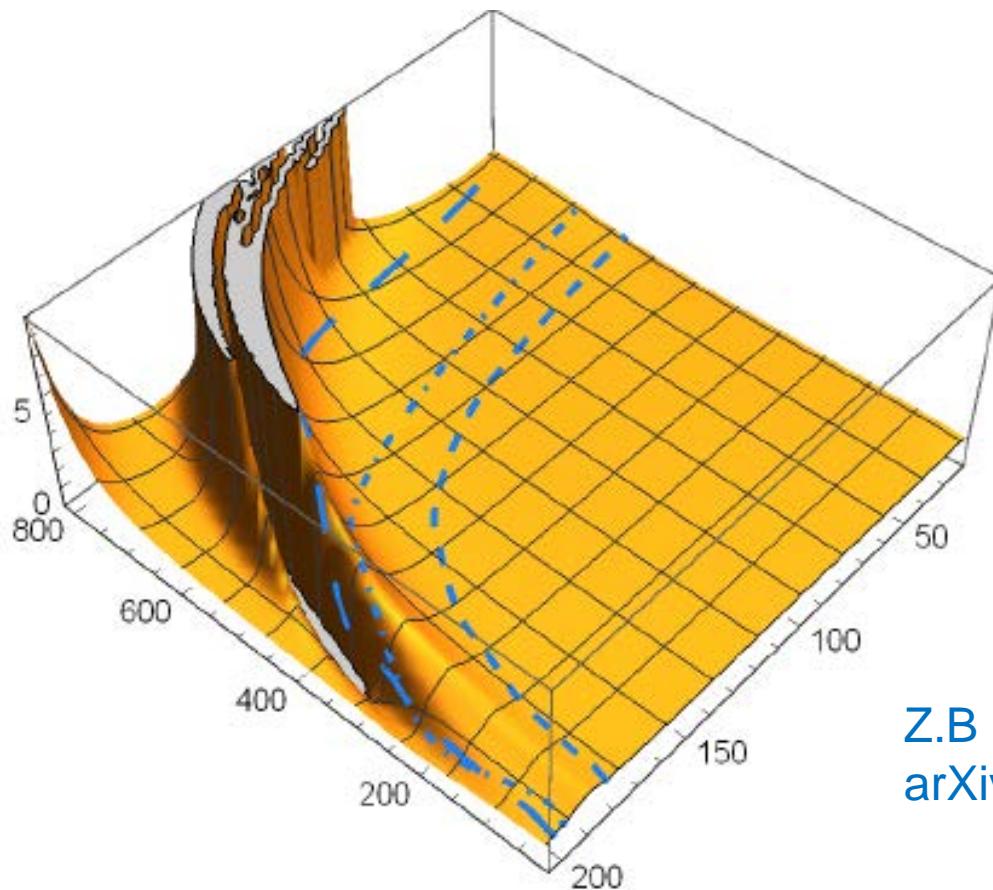


Freeze-out starts from
back-ridge of the phase
boundary, →
forming the dip structure

Z.B Li, K.Xu,X.Y.Wang, M.Huang
arXiv:1801.09215

Freeze-out line crosses the foot of CEP mountain
====> forming the peak structure

Peak structure is a clean signature for CEP!!!



Z.B Li, K.Xu,X.Y.Wang, M.Huang
arXiv:1801.09215

A realistic PNJL model

A. Bhattacharyya, S. K. Ghosh, S. Maity, S. Raha,
B. R. Ray, K. Saha and S. Upadhyaya, arXiv:1609.07882.

NJL part:

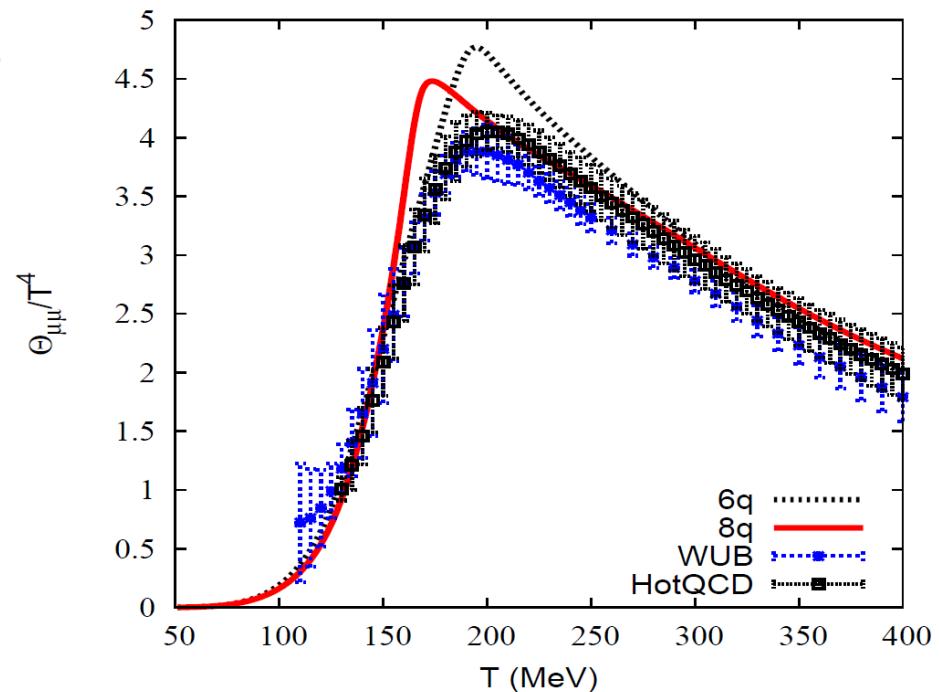
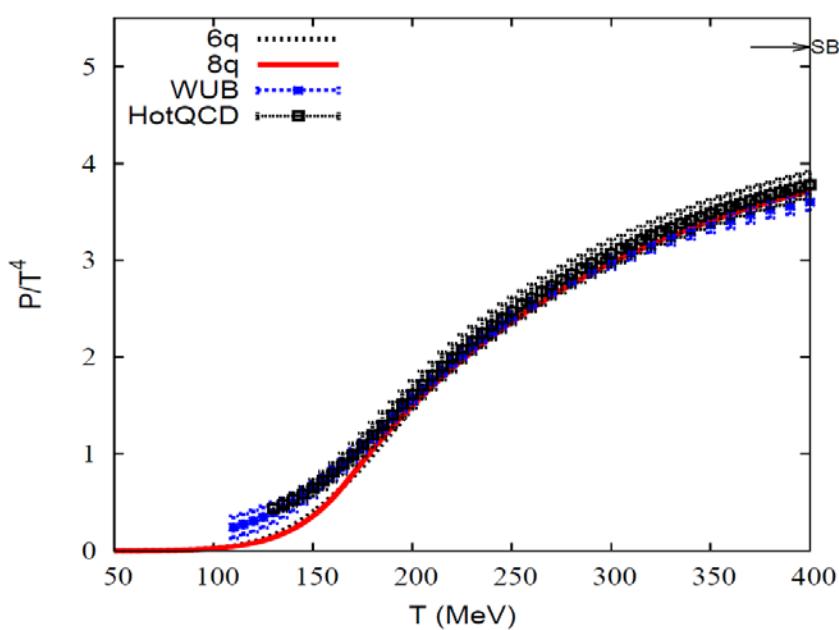
$$\begin{aligned}\Omega = & g_S \sum_f \sigma_f^2 - \frac{g_D}{2} \sigma_u \sigma_d \sigma_s + 3 \frac{g_1}{2} \left(\sum_f \sigma_f^2 \right)^2 + 3g_2 \sum_f \sigma_f^4 - 6 \sum_f \int \frac{d^3 p}{(2\pi)^3} E_f \Theta(\Lambda - |\vec{p}|) \\ & - 2T \sum_f \int \frac{d^3 p}{(2\pi)^3} \ln[1 + 3(\Phi + \bar{\Phi}) e^{-(E_f - \mu_f)/T}] e^{-(E_f - \mu_f)/T} + e^{-3(E_f - \mu_f)/T}] \\ & - 2T \sum_f \int \frac{d^3 p}{(2\pi)^3} \ln[1 + 3(\Phi + \bar{\Phi}) e^{-(E_f + \mu_f)/T}] e^{-(E_f + \mu_f)/T} + e^{-3(E_f + \mu_f)/T}] \\ & + U'(\Phi, \bar{\Phi}, T)\end{aligned}$$

Polyakov Loop:

$$\frac{U'}{T^4} = \frac{U}{T^4} - \kappa \ln[J(\Phi, \bar{\Phi})] \quad \frac{U}{T^4} = -\frac{b_2(T)}{2} \bar{\Phi} \Phi - \frac{b_3}{6} (\Phi^3 + \bar{\Phi}^3) + \frac{b_4}{4} (\Phi \bar{\Phi})^2$$

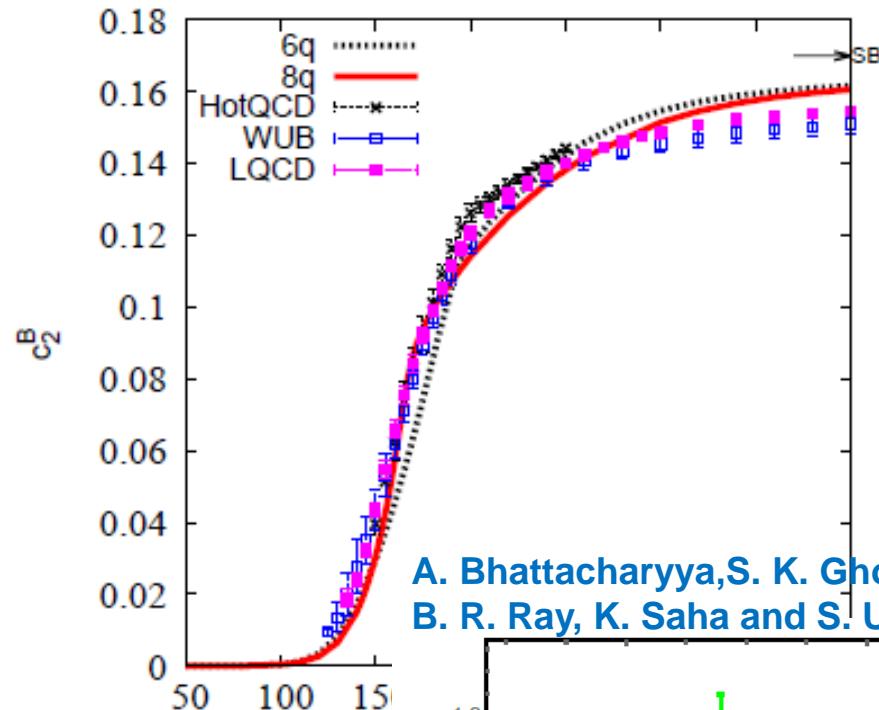
$$J = \left(\frac{27}{24\pi^2} \right) (1 - 6\Phi\bar{\Phi} + 4(\Phi^3 + \bar{\Phi}^3) - 3(\Phi\bar{\Phi})^2)$$

$$b_2(T) = a_0 + a_1 \frac{T_0}{T} \exp(-a_2 \frac{T}{T_0})$$

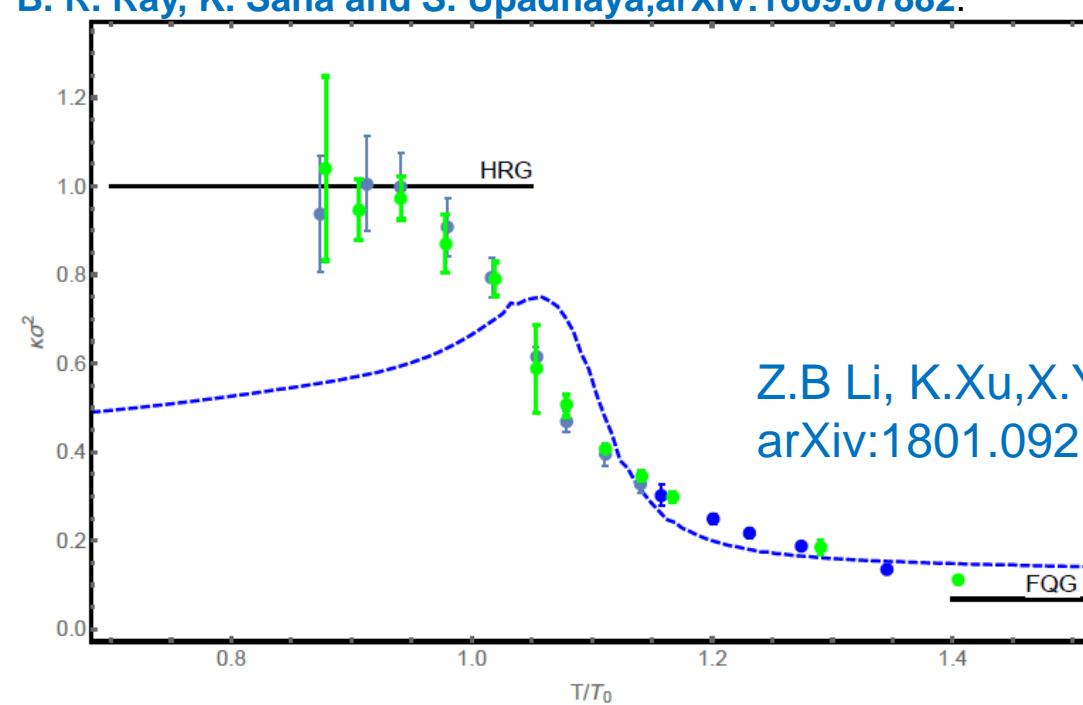
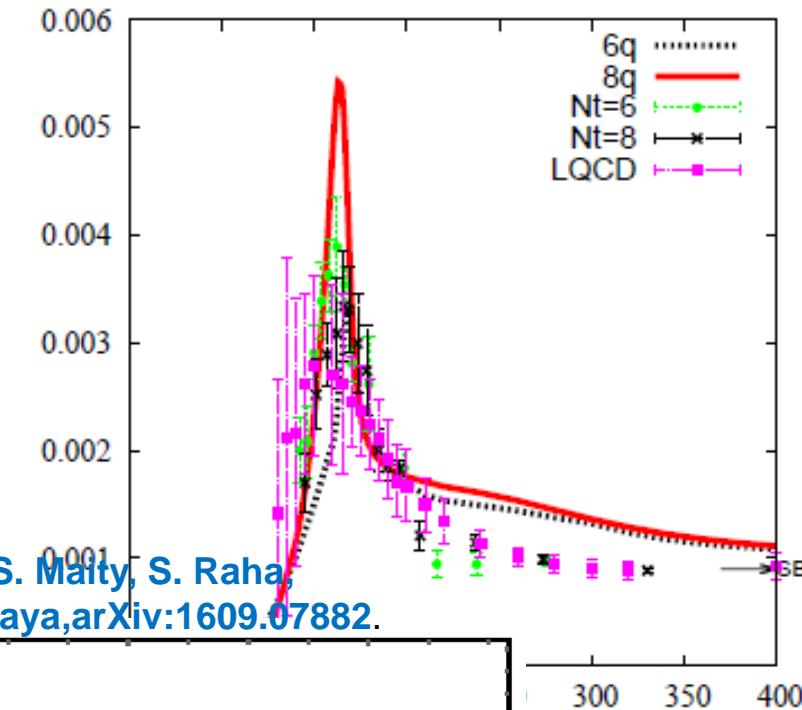


Parameters are fitted to pressure density lattice result at zero baryon chemical potential,

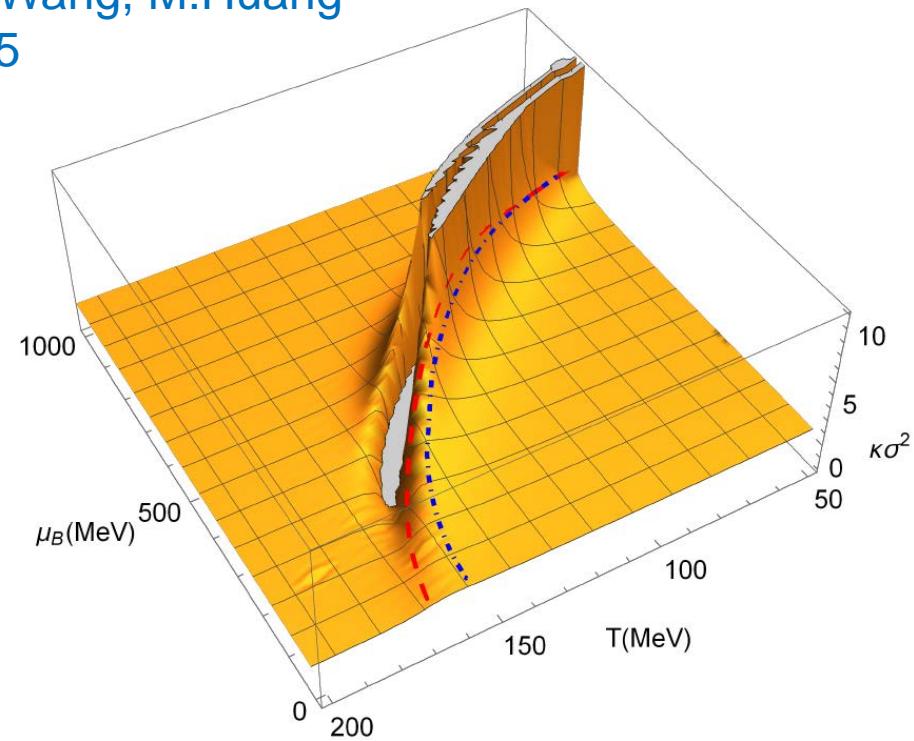
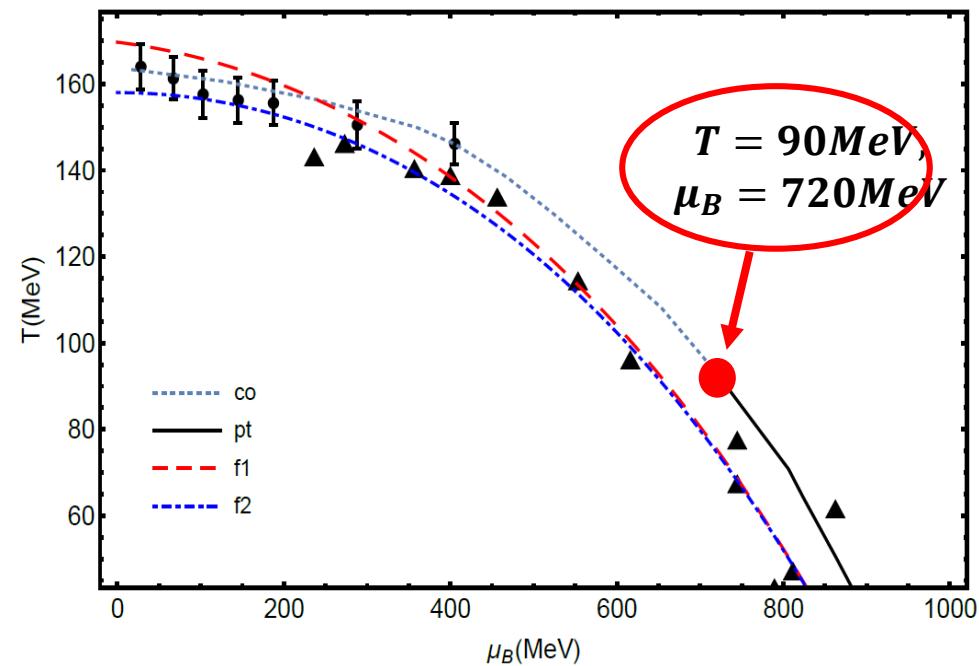
- 1) $T_c=154$ MeV;
- 2) EOS: p,e,s, trace anomaly;
- 3) Baryon number fluctuations



A. Bhattacharyya, S. K. Ghosh, S. Maity, S. Raha,
B. R. Ray, K. Saha and S. Upadhyaya, arXiv:1609.07882.



Z.B Li, K.Xu,X.Y.Wang, M.Huang
arXiv:1801.09215



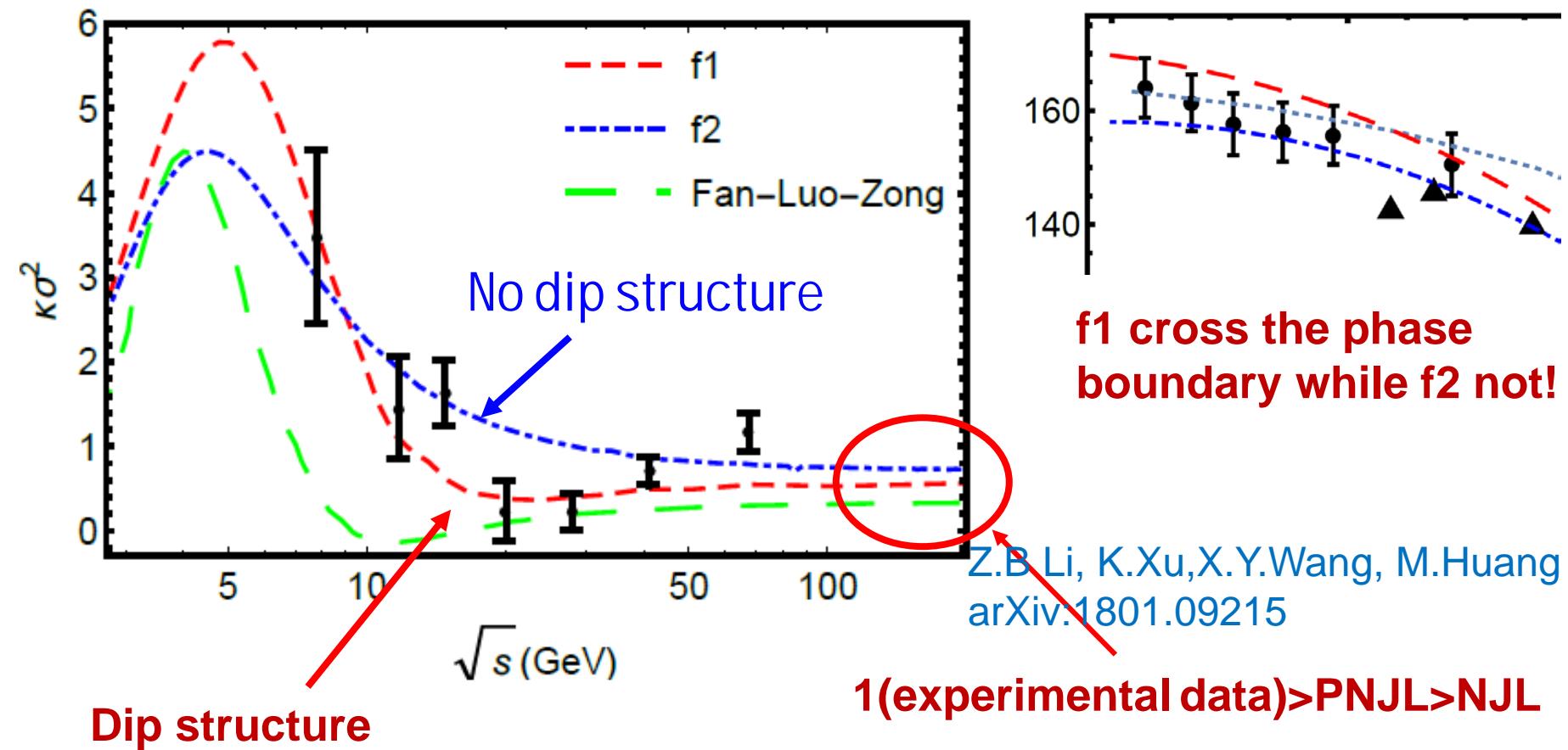
Freeze-out data: BES-I data and other experimental data

Two freeze-out lines:

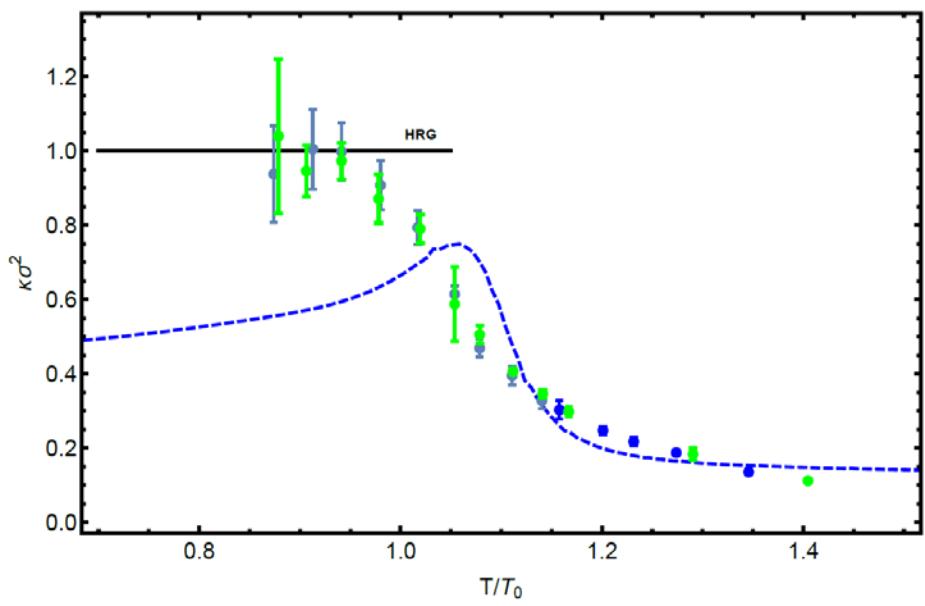
f1: $T(\mu) = 0.158 - 0.14\mu^2 - 0.04\mu^4 - 0.013(0.948 - \mu)^2$

f2: $T(\mu) = 0.158 - 0.14\mu^2 - 0.04\mu^4$

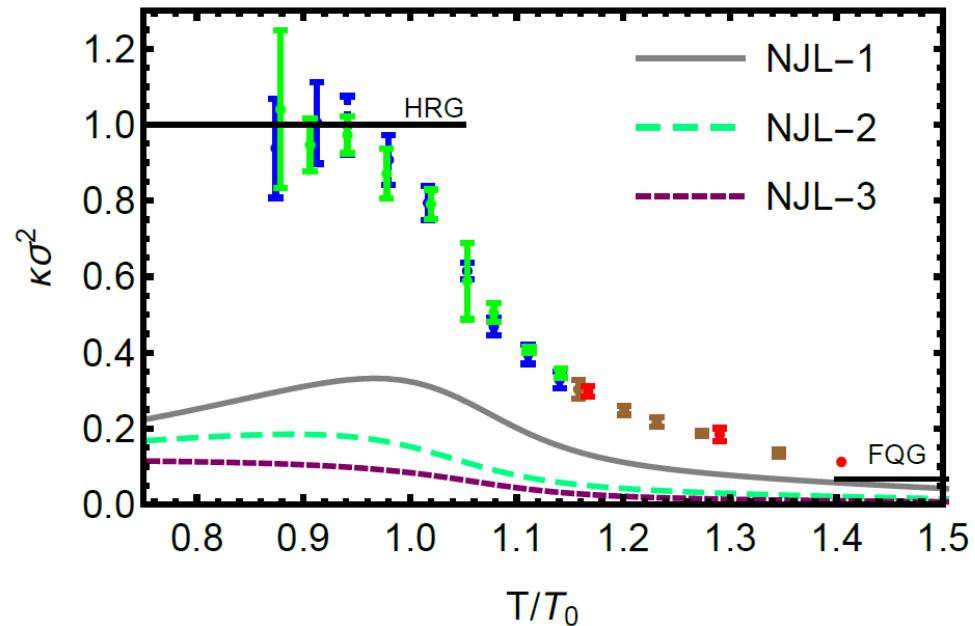
Kurtosis along freeze-out lines



1. Agree well with BES-I data! --->equilibrium result can describe the experimental data!!!
2. The dip structure is sensitive to the relation between the freeze-out line and the phase boundary



PNJL model

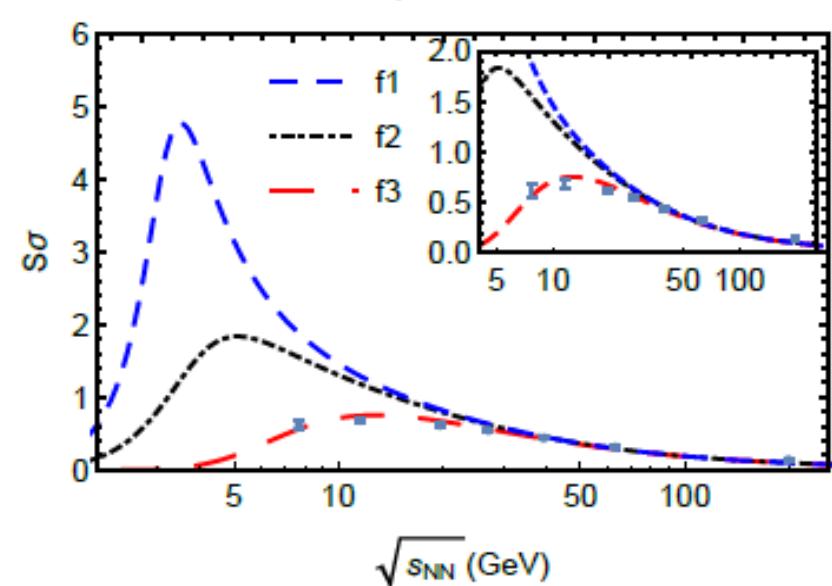
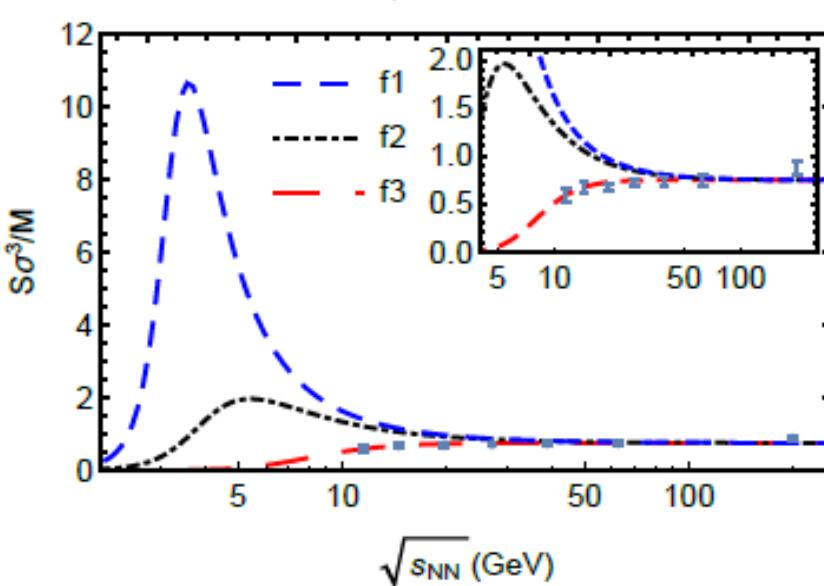
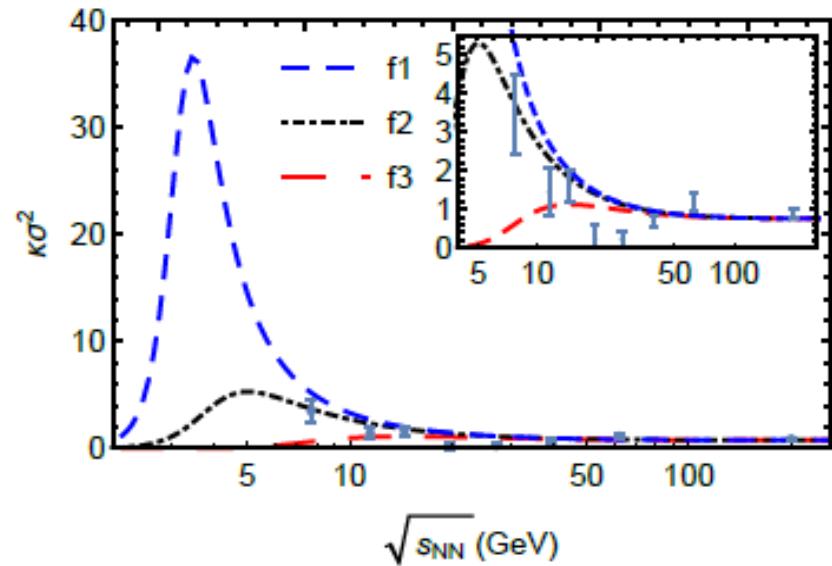
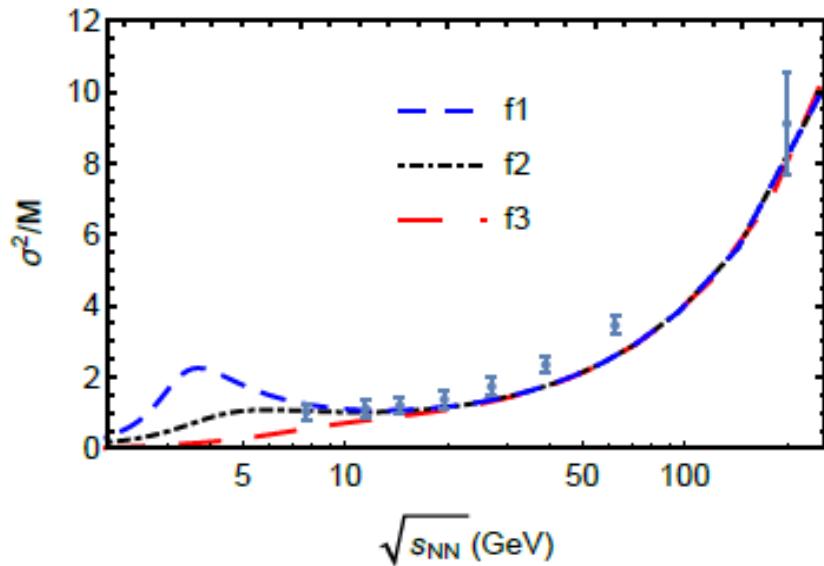


NJL model

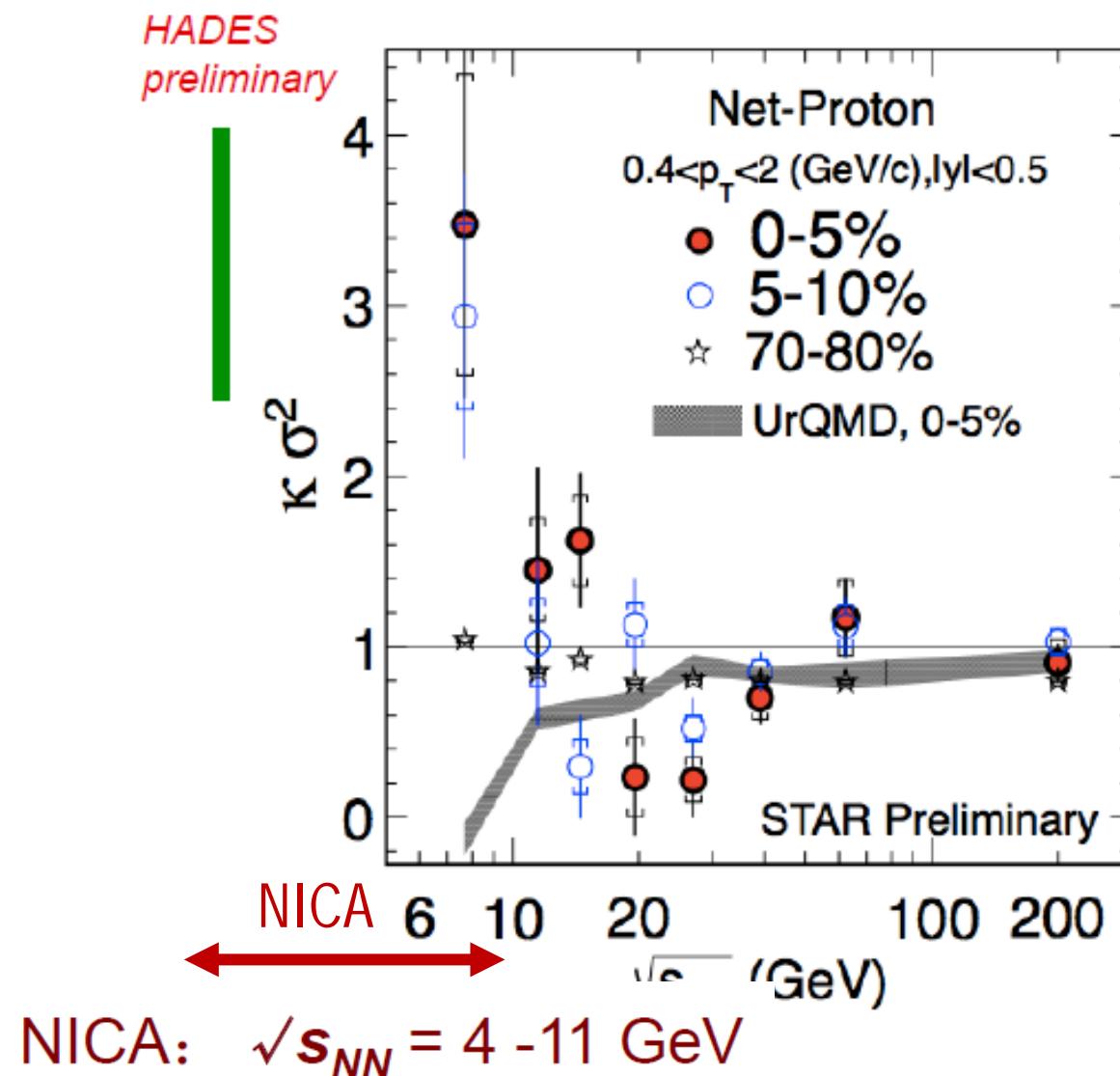
- Contribution from gluon plays a important role!
- The difference of BNF between PNJL model and lattice data below T_c is due to the fact of no real confinement in the PNJL model! (To be improved in the future!)

BES Phase-I measurement meets with HQCD model

Zhibin Li, Yidian Chen, Danning Li, M.H., arXiv:1706.02238



Peak structure is solely determined by CEP!!! A clean signature for CEP!!! From BES-I and HADES, peak structure is expected to show up in the collision energy of 5-6 GeV!!!

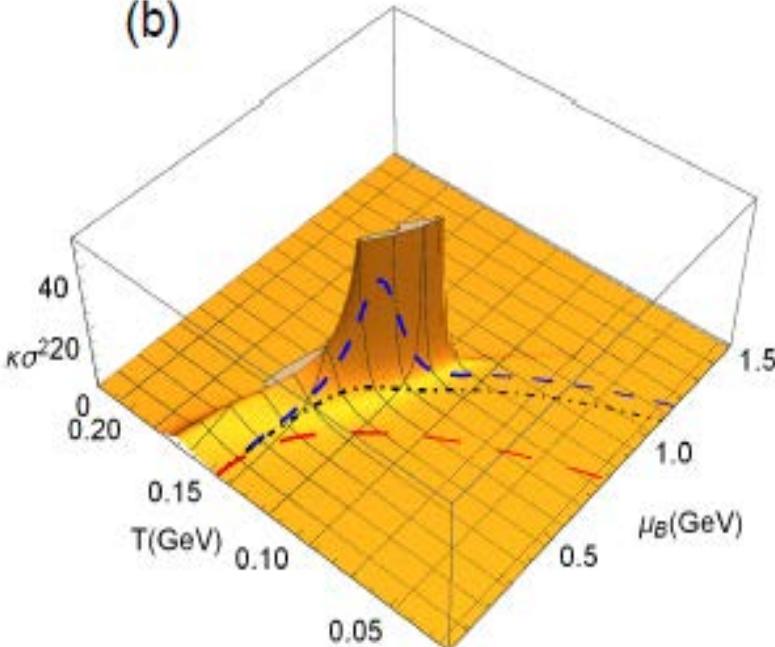


Warning: The precise location of the CEP measured might not be the same as real QCD predicted

Finite size effect, freeze out, evolution of the system

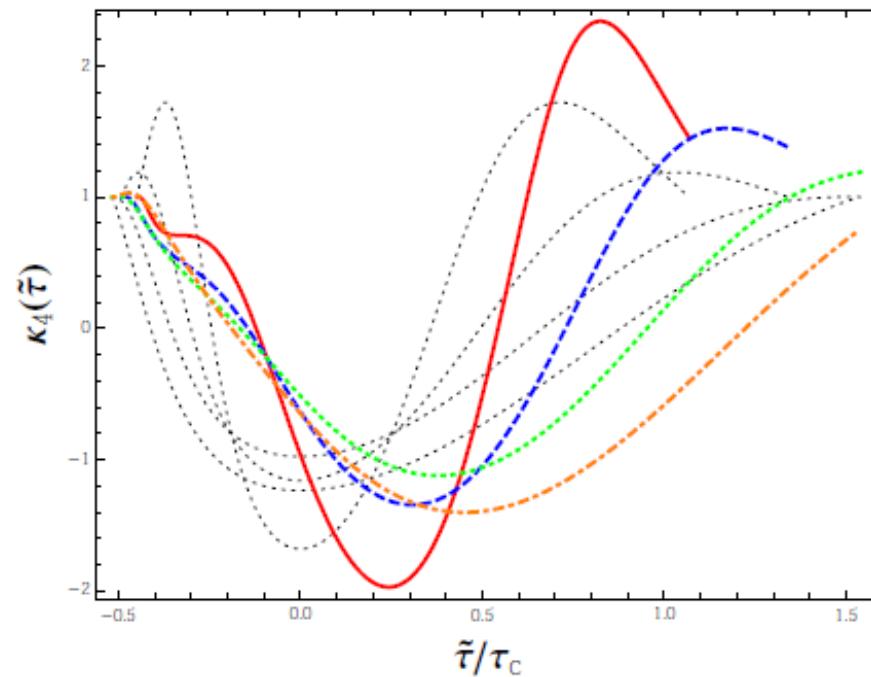
these effects may shift the location of CEP,

(b)



Out-of equilibrium, without
the constraint of stability
condition

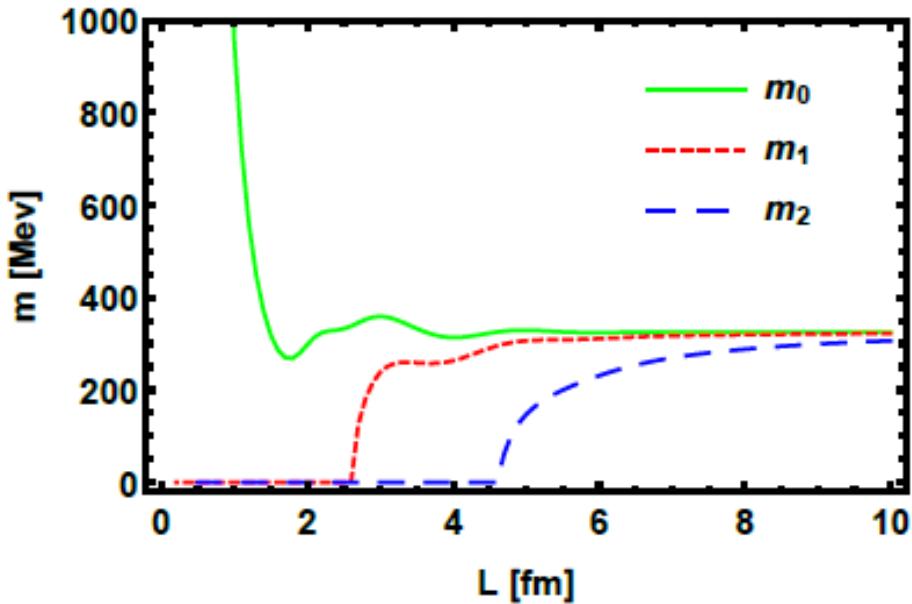
-> Sign change ?



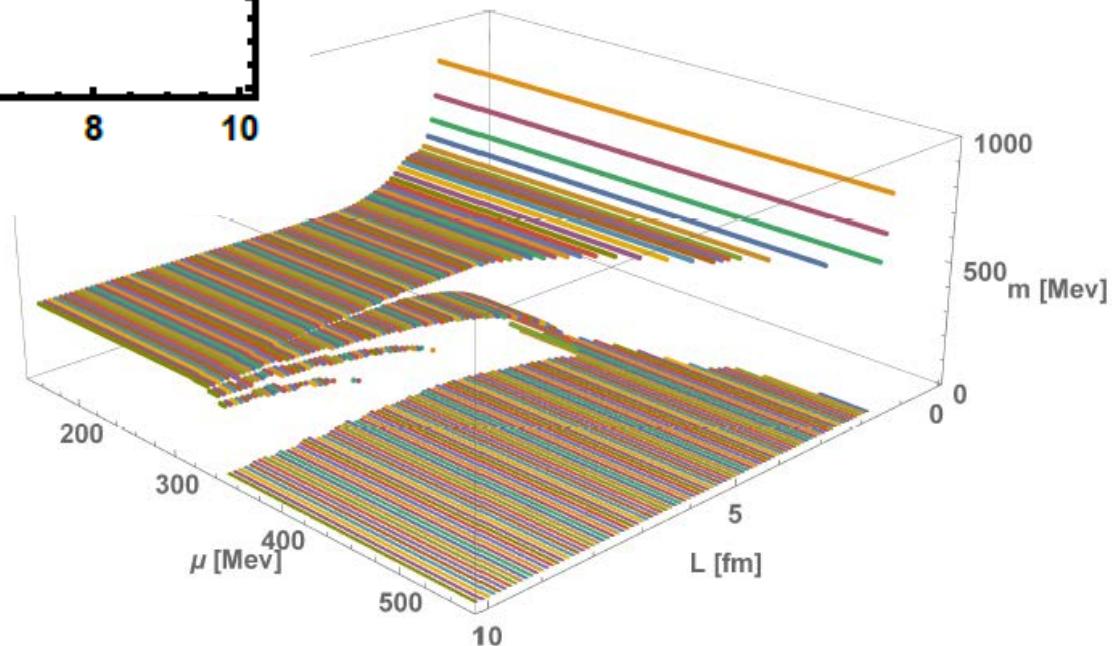
The sign and magnitude
at freeze-out
is most important!

S.Mukherjee, R.Venugopalan, Y.Yin,
Phys.Rev.Lett. 117 (2016) no.22, 222301

**L>5fm: Finite size effect is negligible
L<3 fm: Finite size effect is significant!**



K.Xu, M.Huang, to appear



III. Conclusion and Outlook

- Contribution from gluodynamics is dominant for BNF;
- The peak of BNF along freeze-out line is solely related to the CEP;
- The BES-I measurement of BNF can be described by a realistic PNJL model in equilibrium;
- CEP at small baryon number densities are ruled out both from lattice results and BES-I measurement!

Thanks for your attention!