# $\eta$ - $\eta'$ mixing in Large- $N_c$ ChPT: discussion, phenomenology, and prospects

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## (preliminary) Work done in collaboration with <u>P. Bickert</u> and S. Scherer (Mainz)

#### I profit as well from arXiv:1504.07742[hep-ph] done with R. Escribano (UAB) and <u>P. Sanchez Puertas</u> (Mainz)



# Outline

- Motivation and Introduction
- $\eta$ - $\eta$ ' mixing in Large-N<sub>c</sub> ChPT
  - strategy and preliminary results
- Conclusions and Prospects

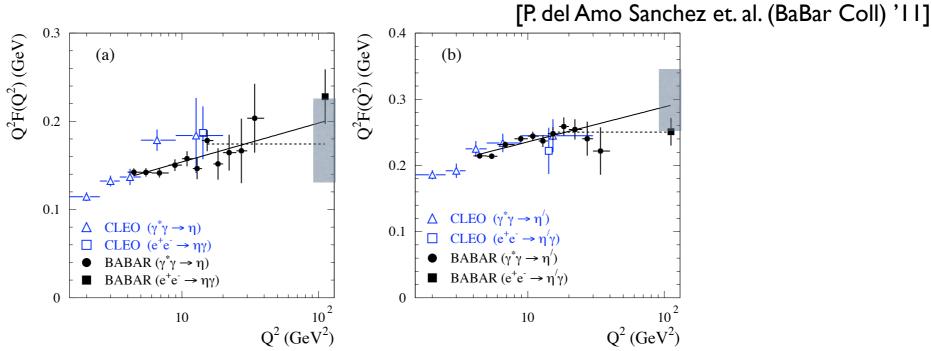
## The $\eta$ - $\eta$ ' system

- Ideal for studying:
  - Symmetries
  - Symmetry breaking in QCD

$$\begin{array}{ll} \eta, \eta' \to 3\pi & \longrightarrow & \text{Quark masses} \\ \eta, \eta' \to 3\pi & \eta' \to \pi\pi\eta & \longrightarrow & \text{Chiral invariant EFT} \\ \eta, \eta' \to 2\gamma \text{ or } \pi\pi\gamma & \longrightarrow & \text{The Chiral Anomaly} \\ \eta, \eta' \to \gamma e^+ e^- & \longrightarrow & \text{E.M. Form Factors} \end{array}$$

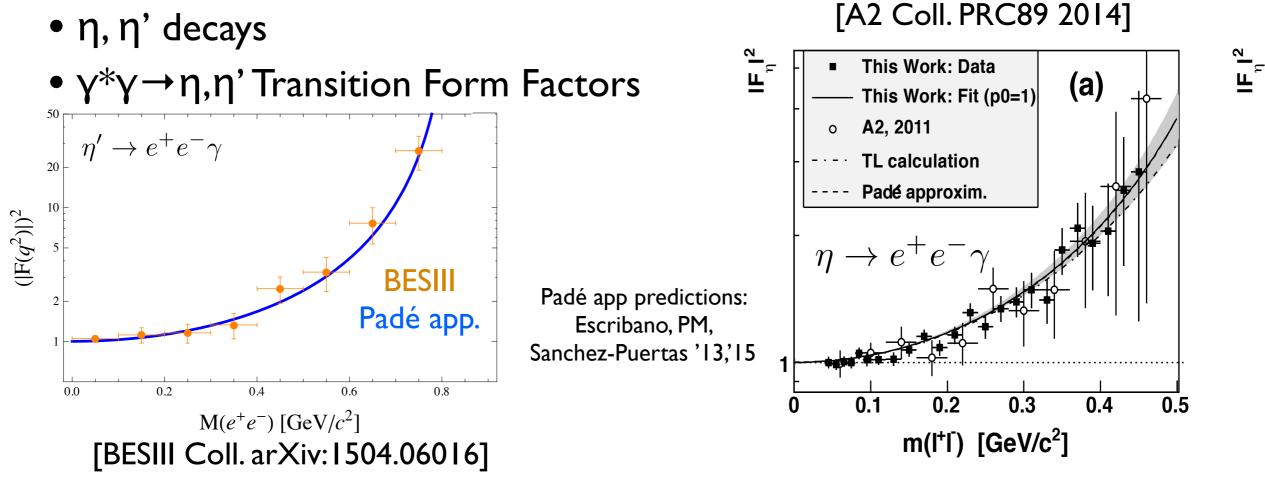
#### • $\eta$ - $\eta$ ' mixing probes

- strange quark content of light pseudoscalar
- gluon dynamics of QCD
- Experimentally related to
  - η, η' decays
  - $\gamma^*\gamma \rightarrow \eta, \eta$ ' Transition Form Factors



#### •η-η' mixing probes

- strange quark content of light pseudoscalar
- gluon dynamics of QCD
- Experimentally related to
  - η, η' decays



- Intense experimental program on  $\eta$ ' physics:
  - At BESIII (remember the talk by S. Fang):

• 
$$\eta' \to \eta \pi \pi$$
  $\eta' \to 3\pi$   $\eta' \to 4\pi$   $\eta' \to \gamma \gamma \pi^0$   $\eta' \to \gamma e^+ e^-$ 

• At MAMI

• Goal I.5xI0<sup>4</sup> h<sup>-1</sup> 
$$\eta$$
':  $\eta' \to \eta \pi \pi$   $\eta' \to \gamma \gamma \pi^0$   $\eta' \to \gamma e^+ e^-$ 

- From EFTs point of view, the mixing is always there: we need a consistent approach to the mixing as well (intermediate result in a broader context of our Collaborative Research Center CRC 1044 in Mainz)
- In this talk: only small portion (interesting by itself) related to the mixing

- A consistent description of the mixing:
  - How?
    - Use the Effective Field Theory of QCD at low energies (Chiral Perturbation Theory) [Weinberg '79, Gasser and Leutwyler '84,'85]
       •adding consistently the η'
  - Construct the <u>most general Lagrangian</u> compatible with all symmetries and establish a <u>power counting</u> to organize the infinite number of terms

•If Nc  $\rightarrow \infty$ , U(I)-anomaly is suppressed and the singlet axial current

is conserved: a ninth Goldstone boson appears if  $m_{uds}=0$  as well ( $\eta$ ')

•Explicit symmetry breaking  $(m_{uds} \neq 0)$ : massive GB +  $\eta$ - $\eta$ ' mixing

•Power counting in Large-Nc ChPT: expansion parameter  $\delta$ 

$$p = \mathcal{O}(\sqrt{\delta}), \quad m = \mathcal{O}(\delta), \quad 1/N_c = \mathcal{O}(\delta)$$
  
•Decay constant  $F \sim \mathcal{O}(1/\sqrt{\delta})$   
•Flavor trace  $\sim \mathcal{O}(\delta)$   
•k-meson vertex  $\sim \mathcal{O}(\delta^{(k-2)/2})$   
•GB propagator  $\sim \mathcal{O}(1/\delta)$   
•GB loop  $\sim \frac{M^2}{F^2} \sim \mathcal{O}(\delta^2) \longrightarrow$  loops are at NNLO

Ready for constructing the Lagrangians

Building blocks:

$$\phi(x) = \sum_{a=0}^{8} \lambda_a \phi_a(x) = \begin{pmatrix} \pi^0 + \frac{1}{\sqrt{3}} \eta_8 + \frac{F}{3} \psi & \sqrt{2}\pi^+ & \sqrt{2}K^+ \\ \sqrt{2}\pi^- & -\pi^0 + \frac{1}{\sqrt{3}} \eta_8 + \frac{F}{3} \psi & \sqrt{2}K^0 \\ \sqrt{2}K^- & \sqrt{2}\bar{K}^0 & -\frac{2}{\sqrt{3}} \eta_8 + \frac{F}{3} \psi \end{pmatrix}$$

$$\begin{split} U &= exp\left(\frac{i\phi}{F}\right) & r_{\mu} = v_{\mu} + a_{\mu} \\ D_{\mu}U &= \partial_{\mu}U - ir_{\mu}U + iUl_{\mu} & l_{\mu} = v_{\mu} - a_{\mu} \\ \chi &= 2BM & M = \text{diag}(\hat{m}, \hat{m}, m_s) & \hat{m} = m_u = m_d \end{split}$$

Expansion in the parameter  $\delta$ :

$$\mathcal{L}_{\rm eff} = \mathcal{L}^{(0)} + \mathcal{L}^{(1)} + \mathcal{L}^{(2)} + \dots$$

At LO:  $O(\delta^{(0)})$ 

[Gasser, Leutwyler '85]

$$\mathcal{L}^{(0)} = \frac{F^2}{4} \langle D_{\mu} U (D^{\mu} U)^{\dagger} \rangle + \frac{F^2}{4} \langle \chi U^{\dagger} + U \chi^{\dagger} \rangle - \frac{1}{2} \tau (\psi + \theta)^2$$

At NLO:  $O(\delta^{(1)})$ 

[Kaiser, Leutwyler '00]

$$\mathcal{L}^{(1)} = L_5 \langle D_\mu U^\dagger D^\mu U(\chi^\dagger U + U^\dagger \chi) \rangle + L_8 \langle \chi^\dagger U \chi^\dagger U + U^\dagger \chi U^\dagger \chi \rangle$$
$$+ \frac{F^2}{12} \Lambda_1 D_\mu \psi D^\mu \psi + i \frac{F^2}{12} \Lambda_2 \bar{\psi} \langle \chi^\dagger U - U^\dagger \chi \rangle \cdots$$
$$\bar{\psi} = \psi + \theta, \qquad D_\mu \psi = \partial_\mu \psi - 2 \langle a_\mu \rangle$$

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Expansion in the parameter  $\delta$ :

$$\mathcal{L}_{eff} = \mathcal{L}^{(0)} + \mathcal{L}^{(1)} + \mathcal{L}^{(2)} + \dots$$

At NNLO:  $O(\delta^{(2)})$ 

$$\mathcal{O}(1/N_c p^2) \quad \mathcal{O}(p^4) \quad \mathcal{O}(N_c p^6)$$

[Herrera-Siklody et al '96] [Kaiser, Leutwyler '00] [Jiang, Ge, Wang '14] [Guo et al '15]

$$\mathcal{L}^{(2,N_c^{-1}p^2)} = -\frac{F^2}{4} v_2^{(2)} \bar{\psi}^2 \langle \chi U^{\dagger} + U\chi^{\dagger} \rangle$$

 $\mathcal{L}^{(2,p^4)} = L_4 \langle D_\mu U^\dagger D^\mu U \rangle \langle \chi^\dagger U + U^\dagger \chi \rangle + L_6 \langle \chi^\dagger U + U^\dagger \chi \rangle^2 + L_7 \langle \chi^\dagger U - U^\dagger \chi \rangle^2$  $+ i L_{18} D_\mu \psi \langle D^\mu U^\dagger \chi - D^\mu U \chi^\dagger \rangle + i L_{25} \bar{\psi} \langle U^\dagger \chi U^\dagger \chi - \chi^\dagger U \chi^\dagger U \rangle + \dots$ 

Expansion in the parameter  $\delta$ :

$$\mathcal{L}_{eff} = \mathcal{L}^{(0)} + \mathcal{L}^{(1)} + \mathcal{L}^{(2)} + \dots$$

At NNLO:  $O(\delta^{(2)})$ 

$$\mathcal{O}(1/N_c p^2) \quad \mathcal{O}(p^4) \quad \mathcal{O}(N_c p^6)$$

[Herrera-Siklody et al '96] [Kaiser, Leutwyler '00] [Jiang, Ge, Wang '14] [Guo et al '15]

$$\mathcal{L}^{(2,N_c p^6)} = C_{12} \langle \chi_+ h_{\mu\nu} h^{\mu\nu} \rangle + C_{14} \langle u_\mu u^\mu \chi_+^2 \rangle + C_{17} \langle \chi_+ u_\mu \chi_+ u^\mu \rangle$$
$$+ C_{19} \langle \chi_+^3 \rangle + C_{31} \langle \chi_-^2 \chi_+ \rangle + \dots$$

$$\begin{split} h_{\mu\nu} &= \nabla_{\mu}u_{\nu} + \nabla_{\nu}u_{\mu}, \\ u_{\mu} &= i\left\{u^{\dagger}(\partial_{\mu} - ir_{\mu})u - u(\partial_{\mu} - il_{\mu})u^{\dagger}\right\}, \\ \chi_{\pm} &= u^{\dagger}\chi u^{\dagger} \pm u\chi^{\dagger}u. \end{split} \qquad \nabla_{\mu}X &= \partial_{\mu}X + [\Gamma_{\mu}, X], \\ \Gamma_{\mu} &= \frac{1}{2}\left\{u^{\dagger}(\partial_{\mu} - ir_{\mu})u + u(\partial_{\mu} - il_{\mu})u^{\dagger}\right\} \end{split}$$

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# $\frac{i}{2} \rightarrow (\delta^{2}) \rightarrow (j \rightarrow i) \rightarrow (\delta^{1}) \rightarrow (j \rightarrow i) \rightarrow (\delta^{1}) \rightarrow (j \rightarrow i) \rightarrow (j \rightarrow$

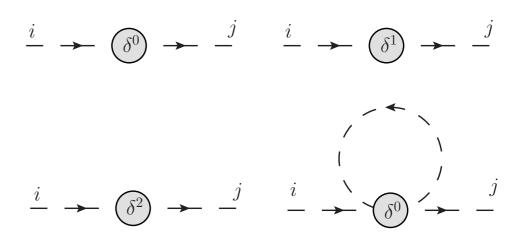
Parametrize first the mixing:

$$\mathcal{L} = \frac{1}{2} \partial_{\mu} \eta_B^T \mathcal{K} \partial^{\mu} \eta_B - \frac{1}{2} \eta_B^T \mathcal{M}^2 \eta_B \qquad \eta_B^T = (\eta_8, \eta_1)$$

(now, relate bare fields to physical  $\eta$ ,  $\eta$ ')

Mixing Lagrangian:

$$\mathcal{L} = \frac{1}{2} \partial_{\mu} \eta_B^T \mathcal{K} \partial^{\mu} \eta_B - \frac{1}{2} \eta_B^T \mathcal{M}^2 \eta_B$$



in terms of

$$\mathcal{K} = \begin{pmatrix} 1 + \delta_8^{(1)} + \delta_8^{(2)} & \delta_{81}^{(1)} + \delta_{81}^{(2)} \\ \delta_{81}^{(1)} + \delta_{81}^{(2)} & 1 + \delta_1^{(1)} + \delta_1^{(2)} \end{pmatrix} \quad \text{and} \quad \mathcal{M}^2 = \begin{pmatrix} M_8^2 & M_{81}^2 \\ M_{81}^2 & M_1^2 \end{pmatrix}$$

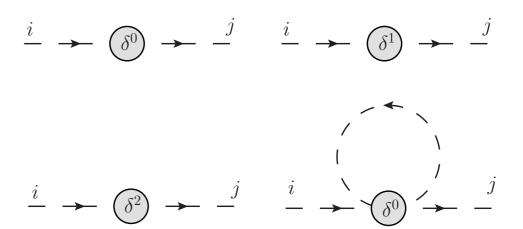
Then, diagonalize K:

$$\eta_{\mathcal{B}} = Z^{1/2} \cdot \hat{\eta} \longrightarrow \hat{\mathcal{M}}^2 = Z^{1/2} \cdot \mathcal{M}^2 \cdot Z^{1/2^T} = \begin{pmatrix} \hat{M}_8^2 & \hat{M}_{81}^2 \\ \hat{M}_{81}^2 & \hat{M}_1^2 \end{pmatrix}$$

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Mixing Lagrangian:

$$\mathcal{L} = \frac{1}{2} \partial_{\mu} \eta_B^T \mathcal{K} \partial^{\mu} \eta_B - \frac{1}{2} \eta_B^T \mathcal{M}^2 \eta_B$$



in terms of

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Then, diagonalize  $\hat{\mathcal{M}}^2 \longrightarrow \hat{\mathcal{M}}^2 = R^T \cdot \mathcal{M}_D^2 \cdot R$ 

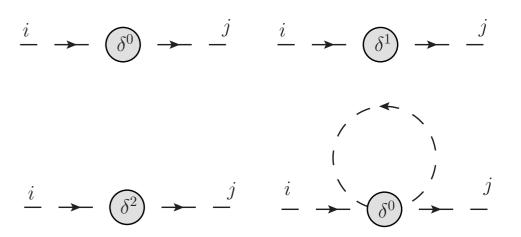
$$\mathcal{M}_D^2 = \begin{pmatrix} M_\eta^2 & 0\\ 0 & M_{\eta'}^2 \end{pmatrix} \qquad \qquad R \equiv \begin{pmatrix} \cos\theta^{(2)} & -\sin\theta^{(2)}\\ \sin\theta^{(2)} & \cos\theta^{(2)} \end{pmatrix}$$

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Mixing Lagrangian:  $\frac{i}{-} \rightarrow (\delta^0) \rightarrow -\frac{j}{-} \rightarrow (\delta^1) \rightarrow -\frac{j}{-}$  $\mathcal{L} = \frac{1}{2} \partial_{\mu} \eta_B^T \mathcal{K} \partial^{\mu} \eta_B - \frac{1}{2} \eta_B^T \mathcal{M}^2 \eta_B$  $\stackrel{i}{\longrightarrow} (\delta^2) \xrightarrow{j} \stackrel{i}{\longrightarrow} (\delta^0) \xrightarrow{j} \stackrel{j}{\longrightarrow} \stackrel{j}{\longrightarrow} (\delta^0)$ From  $\hat{\mathcal{M}}^2 = R^T \cdot \mathcal{M}_D^2 \cdot R$  with  $\mathcal{M}_D^2 = \begin{pmatrix} M_\eta^2 & 0\\ 0 & M_z^2 \end{pmatrix}$  $1)\hat{M}_8^2 = M_n^2 \cos^2 \theta^{(2)} + M_{n'}^2 \sin^2 \theta^{(2)}$  $\sin 2\theta^{(2)} = \frac{2M_{81}^2}{M^2 - M^2}$  $2)\hat{M}_1^2 = M_n^2 \sin^2 \theta^{(2)} + M_{n'}^2 \cos^2 \theta^{(2)}$  $3)\hat{M}_{81}^2 = (M_n^2 - M_{n'}^2)\sin\theta^{(2)}\cos\theta^{(2)}$ 

Mixing Lagrangian:

$$\mathcal{L} = \frac{1}{2} \partial_{\mu} \eta_B^T \mathcal{K} \partial^{\mu} \eta_B - \frac{1}{2} \eta_B^T \mathcal{M}^2 \eta_B$$



Using the Lagrangians, calculate the self-energy diagrams

$$\Sigma_{88}(p^2) = -(\delta_8^{(1)} + \delta_8^{(2)})p^2 + M_8^2,$$
  

$$\Sigma_{81}(p^2) = \Sigma_{18}(p^2) = -(\delta_{81}^{(1)} + \delta_{81}^{(2)})p^2 + M_{81}^2,$$
  

$$\Sigma_{11}(p^2) = -(\delta_1^{(1)} + \delta_1^{(2)})p^2 + M_1^2.$$

(at LO)

$$\sin 2\theta^{(0)} = -\frac{4\sqrt{2}}{3} \frac{\stackrel{\circ}{M_K^2} - \stackrel{\circ}{M_\pi^2}}{M_{\eta'}^2 - M_\eta^2}$$

## Numerical analysis

#### A word on convergence

•Large-Nc ChPT (using input from Bijnens and Ecker '14):

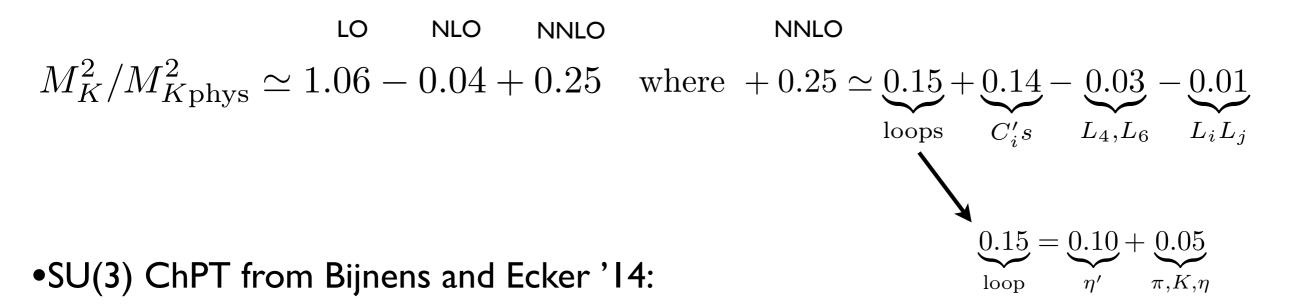
LO NLO NNLO NNLO NNLO NNLO  $F_K/F\pi \simeq 1 + 0.15 + 0.03$ , where  $0.03 \simeq \underbrace{0.05}_{\text{loop}} - \underbrace{0.01}_{C_i} - \underbrace{0.02}_{L_iL_j}$  $\underbrace{0.05}_{\text{loop}} = \underbrace{0.01}_{n'} + \underbrace{0.04}_{\pi,K,n}$ 

•SU(3) ChPT (from Bijnens and Ecker '14):

 $LO \quad \text{NLO} \quad \text{NNLO}$   $F_K/F\pi \simeq 1 + 0.18 + 0.02$ 

#### A word on convergence

•Large-Nc ChPT (using input from Bijnens and Ecker '14):



$${}^{\rm LO}_{K} N {\rm LO}_{K} {\rm NNLO} \label{eq:MK} M_{K}^2 / M_{K}^2 {\rm phys} \simeq 1.11 - 0.07 - 0.04$$

#### A word on convergence

•Large-Nc ChPT (using input from Bijnens and Ecker '14):

$$M_K^2/M_K^2$$
 phys  $\simeq 1.11 - 0.07 - 0.04$ 

#### Results: the mixing angle

General result at NNLO:

$$1)\hat{M}_{8}^{2} = M_{\eta}^{2}\cos^{2}\theta^{(2)} + M_{\eta'}^{2}\sin^{2}\theta^{(2)}$$
$$2)\hat{M}_{1}^{2} = M_{\eta}^{2}\sin^{2}\theta^{(2)} + M_{\eta'}^{2}\cos^{2}\theta^{(2)}$$
$$3)\hat{M}_{81}^{2} = (M_{\eta}^{2} - M_{\eta'}^{2})\sin^{2}\theta^{(2)}\cos^{2}\theta^{(2)}$$

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$$M_{\eta'}^2 + M_{\eta}^2 = \hat{M}_8^2 + \hat{M}_1^2, \quad M_{\eta'}^2 - M_{\eta}^2 = \sqrt{(\hat{M}_8^2 - \hat{M}_1^2)^2 + 4\hat{M}_{81}^4}$$
$$\sin 2\theta^{(2)} = \frac{2\hat{M}_{81}^2}{\mathcal{C}}$$

#### Results: the mixing angle

At LO:

$$M_{\eta'}^2 + M_{\eta}^2 = \hat{M}_8^2 + \hat{M}_1^2, \quad M_{\eta'}^2 - M_{\eta}^2 = \sqrt{(\hat{M}_8^2 - \hat{M}_1^2)^2 + 4\hat{M}_{81}^4}$$

$$\sin 2\theta^{(0)} = -\frac{4\sqrt{2}}{3} \frac{\overset{\circ}{M_K^2} - \overset{\circ}{M_{\pi}^2}}{M_{\eta'}^2 - M_{\eta}^2} \longrightarrow \theta^{(0)} = -19.6^{\circ}$$

#### Results: the mixing angle

At LO:

$$M_{\eta'}^2 + M_{\eta}^2 = \hat{M}_8^2 + \hat{M}_1^2, \quad M_{\eta'}^2 - M_{\eta}^2 = \sqrt{(\hat{M}_8^2 - \hat{M}_1^2)^2 + 4\hat{M}_{81}^4}$$

$$\sin 2\theta^{(0)} = \frac{-4\sqrt{2}(\overset{\circ}{M_{K}^{2}} - \overset{\circ}{M_{\pi}^{2}})}{\sqrt{(2\overset{\circ}{M_{K}^{2}} - 2\overset{\circ}{M_{\pi}^{2}} - M_{0}^{2})^{2} + 32(\overset{\circ}{M_{K}^{2}} - \overset{\circ}{M_{\pi}^{2}})^{2}}} \longrightarrow \theta^{(0)} = -22.0^{\circ}$$

#### Results: the mixing angle

At NLO:

- •We need Low-energy constants from O(p<sup>4</sup>) L<sub>5,8</sub> and OZI  $\Lambda_{1,2}$
- Two strategies:
  - NLO I: Constrained analysis: calculate  $M_{\pi}, M_{K}, M_{\eta}, M_{0}, F_{K}, F_{\pi}$
  - NLO II: Use SU(3) values from Bijnens and Ecker '14 +  $M_{\eta}$ ,  $M_{0}$  (after matching to U(3) and at  $\mu$ =0.77 GeV)

#### Results: the mixing angle

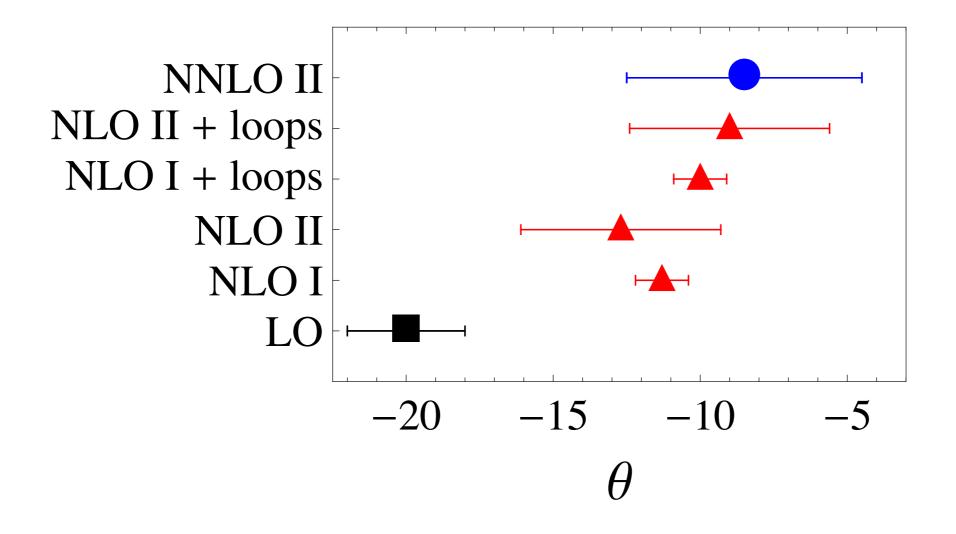
- At NLO + loops:
- •No new LECs ( $L_{5,8}$  and  $\Lambda_{1,2}$ )
- Two strategies:
  - **NLO I + loops: calculate**  $M_{\pi}, M_{K}, M_{\eta}, M_{0}, F_{K}, F_{\pi}$
  - NLO II: Use SU(3) values from Bijnens and Ecker '14 +  $M_{\eta}$ ,  $M_{0}$  (after matching to U(3) and at  $\mu$ =0.77 GeV)

#### Results: the mixing angle

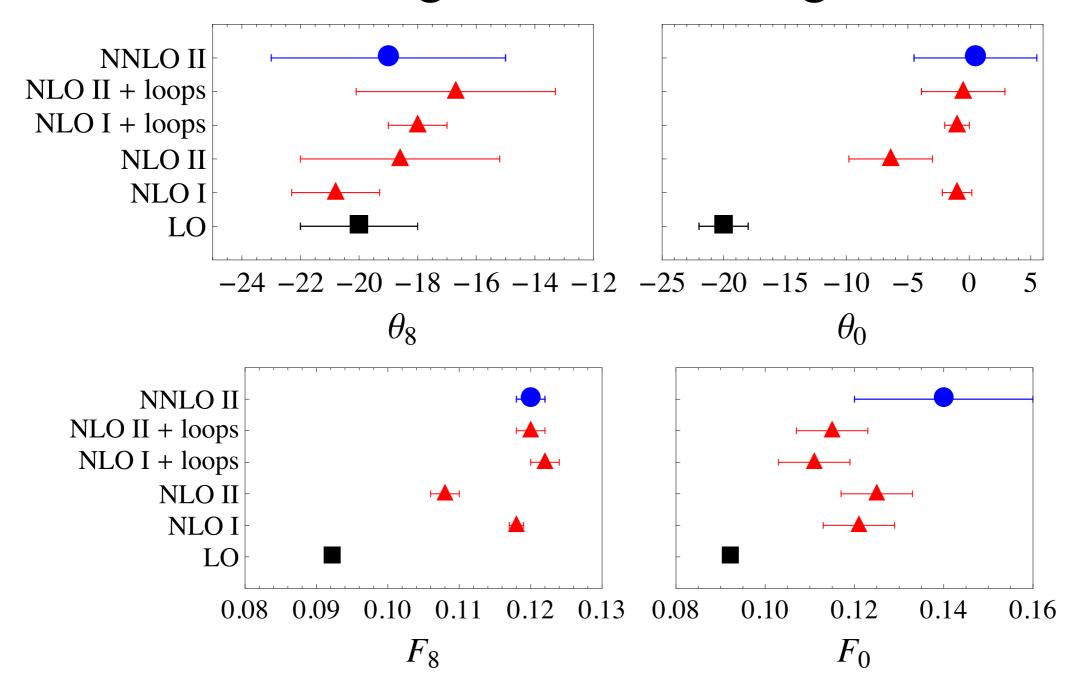
At NNLO:

- NLO + loops + LECs O(p<sup>4</sup>) L<sub>4,6,7,18,25</sub> + LECs O(p<sup>6</sup>) C<sub>12,14,17,19,31</sub> and OZI  $\Lambda_{1,2}$
- Many strategies. Here only one:
  - NNLO II: Use LECs Li and Ci at O(p<sup>6</sup>) from Bijnens and Ecker 'I4 +  $M_{\eta}, M_0$
  - (lattice fits) [Guo et al, '15]

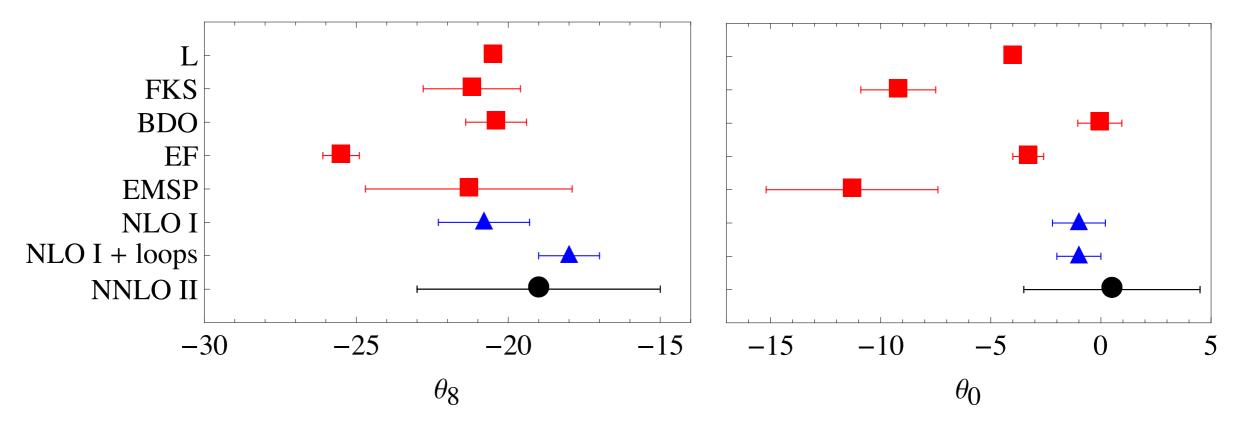
#### Results: the mixing angle



#### Results: mixing in the octet-singlet basis



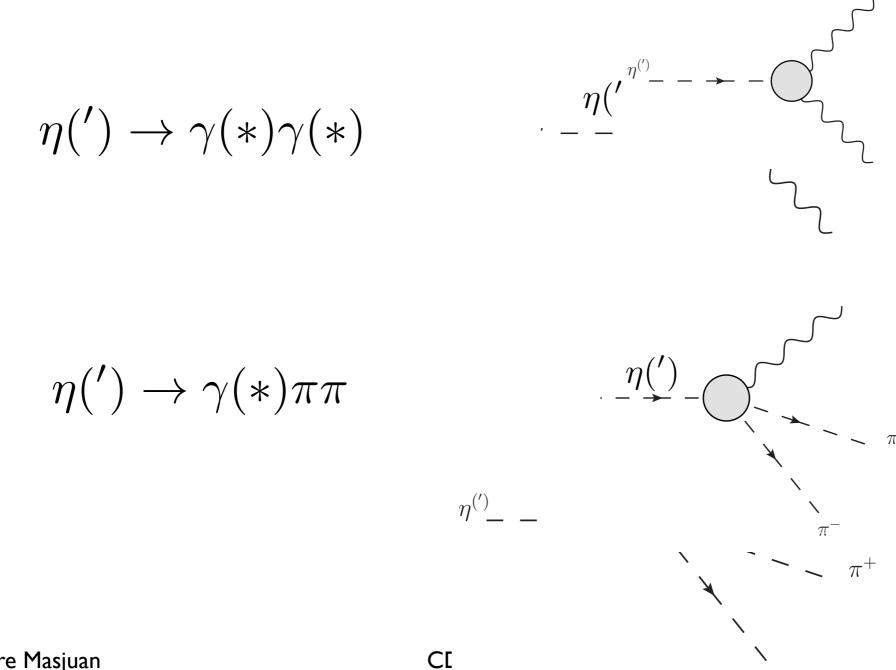
#### comparison with phenomenological determinations



L: Leutwyler '97 FKS: Feldmann, Kroll, Stech '98 BDO: Benayoun, Del Buono, O'Connell, '00 EF: Escribano, Frere '05 EMSP: Escribano, PM, Sanchez-Puertas, '15

## Prospects

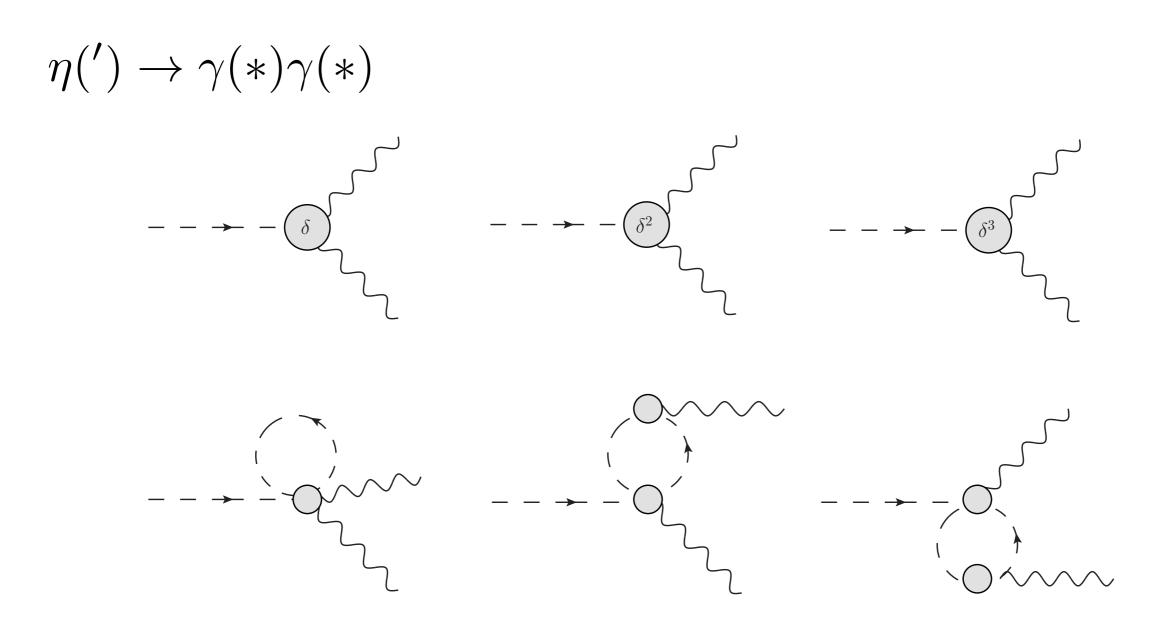
Study of anomalous decays including mixing



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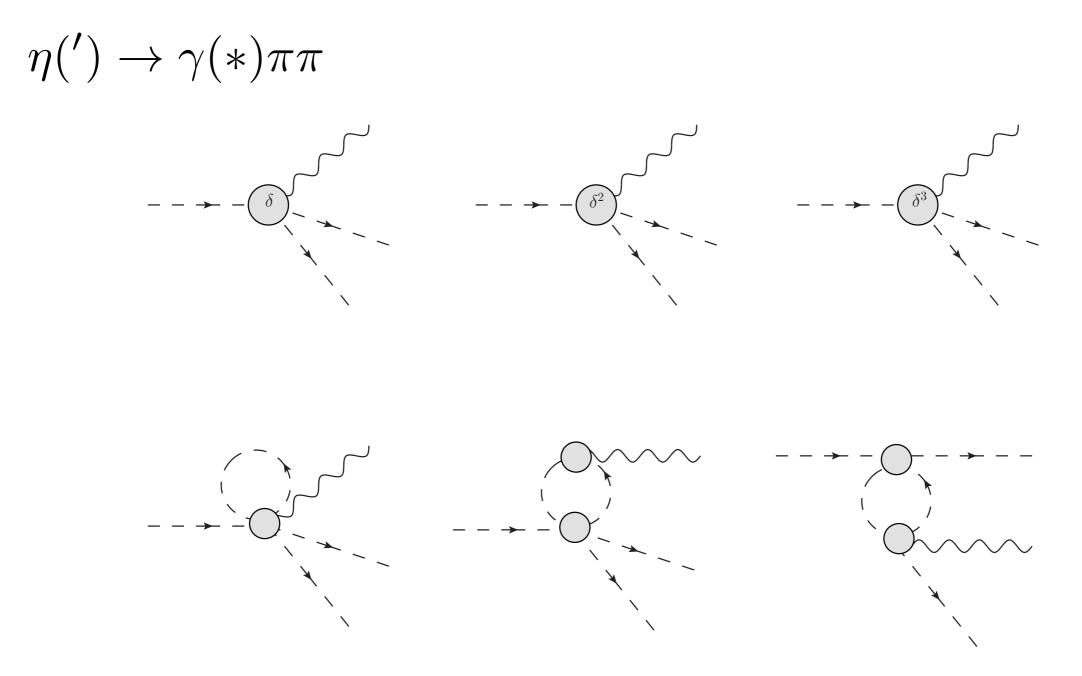
## Prospects

Study of anomalous decays including mixing



## Prospects

Study of anomalous decays including mixing



## Conclusions

•The  $\eta$ - $\eta$ ' systems allows to study symmetries and symmetry breaking in QCD, and is a test of EFTs as well

•Today we explored the  $\eta$ - $\eta$ ' mixing at NNLO in the Large-Nc ChPT framework:

- •we discussed about the relevant Lagrangians
- •we explore the convergence, the role of loops and LECs
- •we provided with preliminary numerical results
  - •while including loops is OK, the proliferation of LECs enlarges the numerical results
- •With this tool at hand, we are now exploring anomalous  $\eta(')$  decays

#### Acknowledgments

- •Thanks to my collaborators:
  - •P. Bickert and S. Scherer for enjoyable collaboration on the calculation of the mixing at NNLO
  - •R. Escribano and P. Sanchez-Puertas for enjoyable collaboration on its phenomenological study
- •Special thanks to the Organizers for the encouragement and support

#### Thanks!

• Notation for the mixing angle:

mixing of mass eigenstates

octet-singlet basis quark-flavour basis mixing angle - $|\eta
angle = \cos \phi_P |\eta_q
angle - \sin \phi_P |\eta_s
angle$  $|\eta
angle = \cos heta_P |\eta_8
angle - \sin heta_P^- |\eta_0
angle$  $|\eta'\rangle = \sin \phi_P |\eta_q\rangle + \cos \phi_P |\eta_s\rangle$  $|\eta'\rangle = \sin\theta_P |\eta_8\rangle + \cos\theta_P |\eta_0\rangle$ 

with

$$egin{aligned} &|\eta_8
angle = rac{1}{\sqrt{6}}(uar{u} + dar{d} - 2sar{s}) & ext{and} &|\eta_q
angle = rac{1}{\sqrt{2}}(uar{u} + dar{d}) \ &|\eta_0
angle = rac{1}{\sqrt{3}}(uar{u} + dar{d} + sar{s}) & |\eta_s
angle = sar{s} \end{aligned}$$

$$\theta_P = \phi_P - \arctan\sqrt{2} \simeq \phi_P - 54.7^\circ$$

- **Assumptions:** no energy dependence
  - $\Gamma_{\eta,\eta'} \ll m_{\eta,\eta'}$
  - no mixing with other pseudoscalars ( $\pi^0$ ,  $\eta_c$ , glueballs)

CD2015, Pisa, June 30th

• Notation for the mixing angles of the decay constants

mixing of decay constants

octet-singlet basis

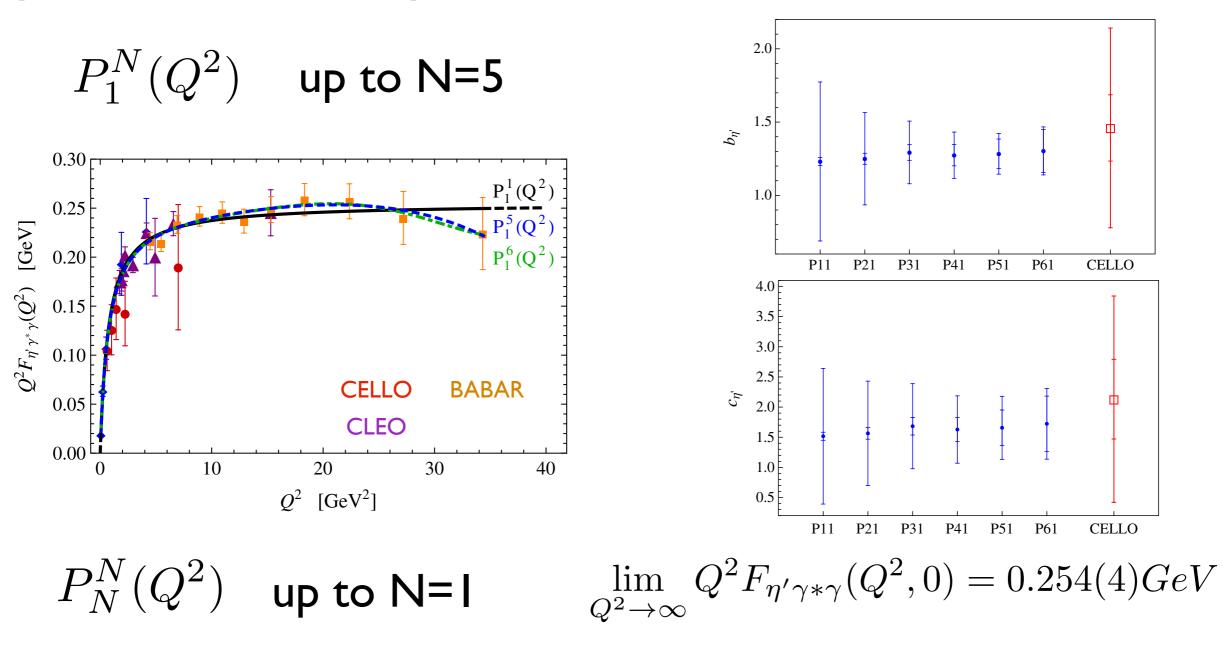
 $\langle 0|A^{a}_{\mu}|P(p)\rangle = if_{P}^{a}p_{\mu}$ with  $A^{a}_{\mu} = \bar{q}\gamma_{\mu}\gamma_{5}\frac{\lambda^{a}}{\sqrt{2}}q$   $f^{a}_{P}(a=8,0;P=\eta,\eta')$ quark-flavour basis  $\langle 0|A^{i}_{\mu}|P(p)\rangle = if^{i}_{P}p_{\mu}$ 

2 mixing angles

## η'-TFF

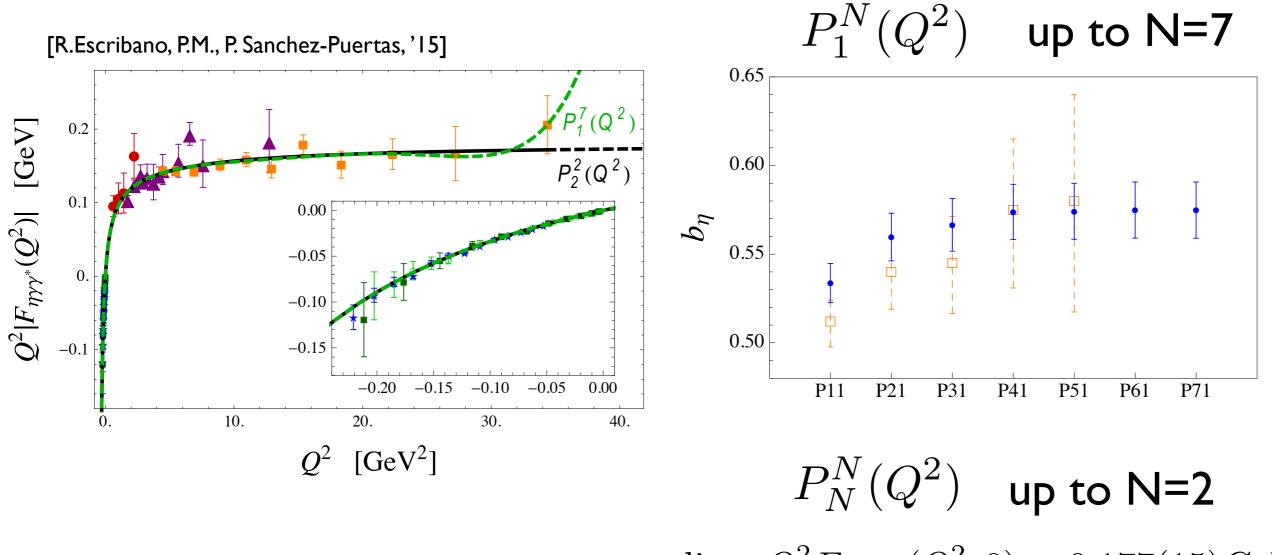
#### Fit to Space-like data: CELLO'91, CLEO'98, L3'98, BABAR'11+ $\Gamma_{\eta' \to \gamma\gamma}$

[R.Escribano, P.M., P. Sanchez-Puertas, '13]



## η-TFF

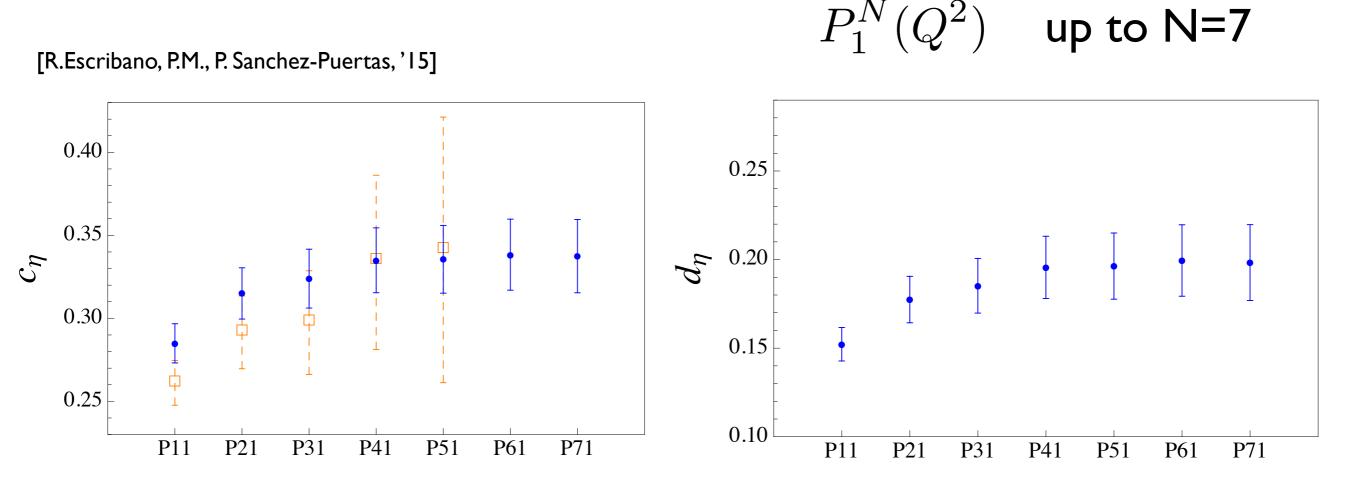
Fit to Space-like data [Cello'91, Cleo'98, BABAR'11]+ $\Gamma_{\eta \to \gamma \gamma}$ + Time-like data [NA60'09, A2'11, A2'13]



 $\lim_{Q^2 \to \infty} Q^2 F_{\eta \gamma * \gamma}(Q^2, 0) = 0.177(15) GeV$ 

# η-TFF

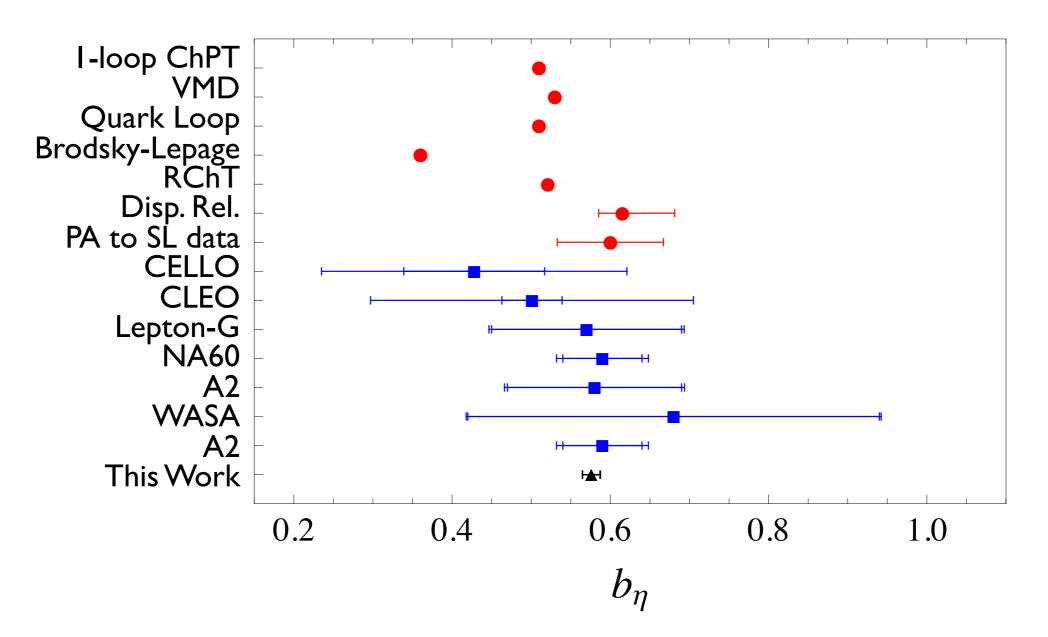
Fit to Space-like data [Cello'91, Cleo'98, BABAR'11]+ $\Gamma_{\eta \to \gamma \gamma}$ + Time-like data [NA60'09, A2'11, A2'13]



# η-TFF

Fit to Space-like data [Cello'91, Cleo'98, BABAR'11]+ $\Gamma_{\eta \to \gamma \gamma}$ + Time-like data [NA60'09, A2'11, A2'13]

[R.Escribano, P.M., P. Sanchez-Puertas, '15]



CD2015, Pisa, June 30th

 $\eta$ - $\eta$ ' mixing in the flavor basis

$$\begin{pmatrix} f_{\eta}^{q} & f_{\eta}^{s} \\ f_{\eta'}^{q} & f_{\eta'}^{s} \end{pmatrix} = \begin{pmatrix} f_{q} \cos[\phi] & -f_{s} \sin[\phi] \\ f_{q} \sin[\phi] & f_{s} \cos[\phi] \end{pmatrix}$$

From the TFFs we can determine  $f_q, f_s, \phi$ 

$$\Gamma_{\eta \to \gamma \gamma} = \frac{9\alpha^2}{32\pi^3} M_{\eta}^3 \left( \frac{C_q \cos[\phi]}{f_q} - \frac{C_s \sin[\phi]}{f_s} \right)^2 \qquad Q$$

$$\Gamma_{\eta' \to \gamma \gamma} = \frac{9\alpha^2}{32\pi^3} M_{\eta'}^3 \left( \frac{C_q \sin[\phi]}{f_q} + \frac{C_s \cos[\phi]}{f_s} \right)^2 \qquad Q$$

$$\lim_{Q^2 \to \infty} Q^2 F_{\eta \gamma \gamma^*}(Q^2) = f_{\eta}^q \frac{10}{3} + f_{\eta}^s \frac{2\sqrt{2}}{3} ,$$
$$\lim_{Q^2 \to \infty} Q^2 F_{\eta' \gamma \gamma^*}(Q^2) = f_{\eta'}^q \frac{10}{3} + f_{\eta'}^s \frac{2\sqrt{2}}{3} .$$

 $\eta$ - $\eta$ ' mixing in the flavor basis

$$\begin{pmatrix} f_{\eta}^{q} & f_{\eta}^{s} \\ f_{\eta'}^{q} & f_{\eta'}^{s} \end{pmatrix} = \begin{pmatrix} f_{q} \cos[\phi] & -f_{s} \sin[\phi] \\ f_{q} \sin[\phi] & f_{s} \cos[\phi] \end{pmatrix}$$

From the TFFs we can determine  $f_q, f_s, \phi$ 

[R.Escribano, P.M., P. Sanchez-Puertas, '15]

[R.Escribano, P.M., P. Sanchez-Puertas, '14]

 $f_q = 1.07(1)f_{\pi}, \quad f_s = 1.39(14)f_{\pi}, \quad \phi = 39.3(1.3)^{\circ}$ 

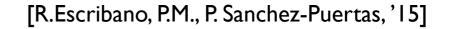
Update of Frere-Escribano '05 with PDG12 using 9 inputs

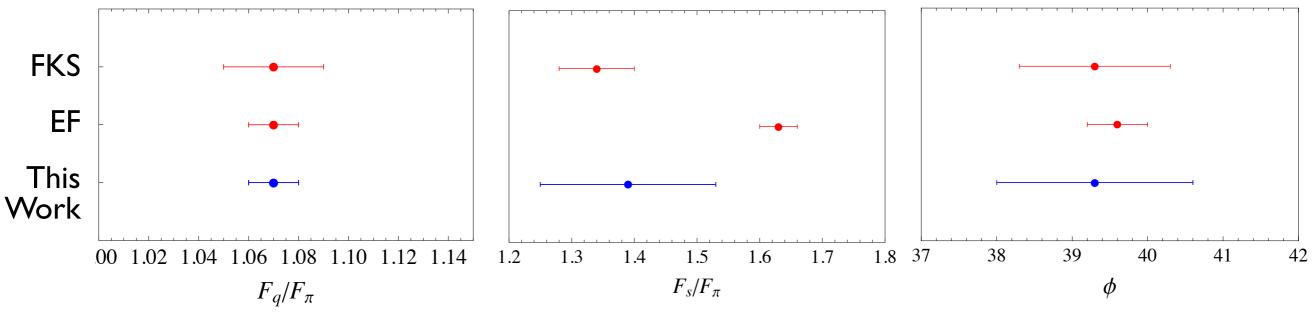
$$f_q = 1.07(1)f_{\pi}, \quad f_s = 1.63(2)f_{\pi}, \quad \phi = 40.4(0.3)^{\circ}$$

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 $\eta$ - $\eta$ ' mixing in the flavor basis

From the TFFs we can determine  $F_q, F_s, \phi$ 





FKS: Feldmann, Kroll, Stech, PLB 449, 339, (1999)

EF: Escribano, Frere, JHEP 0506, 029 (2005) updated in Escribano, P.M, Sanchez-Puertas, 2013.

 $\eta$ - $\eta$ ' mixing in the flavor basis

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#### From the TFFs we can determine $F_q, F_s, \phi$

 $F_q/F_\pi$ 

1.057 1.061 1.065 1.071 1.055 1.086 1.078 1.84 42 1.68 1.53 40 1.39 38 1.26  $() \phi$ 1.14 36 1.04 **L** BaBar 34 32 30 0.14 0.16 0.18 0.20 0.22 0.24 0.26  $\lim_{Q^2 \to \infty} Q^2 F_{\eta \gamma \gamma^*}(Q^2) \text{ (GeV)}$ 

[R.Escribano, P.M., P. Sanchez-Puertas, '15]

TFF at  $q^2 = 112 \text{GeV}^2$ [BABAR PRD'06]

• Is 112GeV<sup>2</sup> not asymptotic yet? Could BELLE crosscheck?

From the TFFs we can determine  $F_q, F_s, \phi$ 

#### and the VPy and J/Y decays used in FKS and EF as inputs

[R.Escribano, P.M., P. Sanchez-Puertas, '15]

( using  $F_{\pi^0} = 131.5 \pm 1.4$  MeV instead of  $F_{\pi^-} = 92.21 \pm 0.14$  MeV )

	Our predictions	Experimental determinations
$g_{ ho\eta\gamma}$	1.55(4)	1.58(5)
$g_{ ho\eta'\gamma}$	1.19(5)	1.32(3)
$g_{\omega\eta\gamma}$	0.56(2)	0.45(2)
$g_{\omega\eta'\gamma}$	0.54(2)	0.43(2)
$g_{\phi\eta\gamma}$	-0.83(11)	-0.69(1)
$g_{\phi\eta^\prime\gamma}$	0.98(14)	0.72(1)
$\frac{J/\Psi \rightarrow \eta' \gamma}{J/\Psi \rightarrow \eta \gamma}$	4.74(60)	4.67(20)