

Implications of LHC Higgs results

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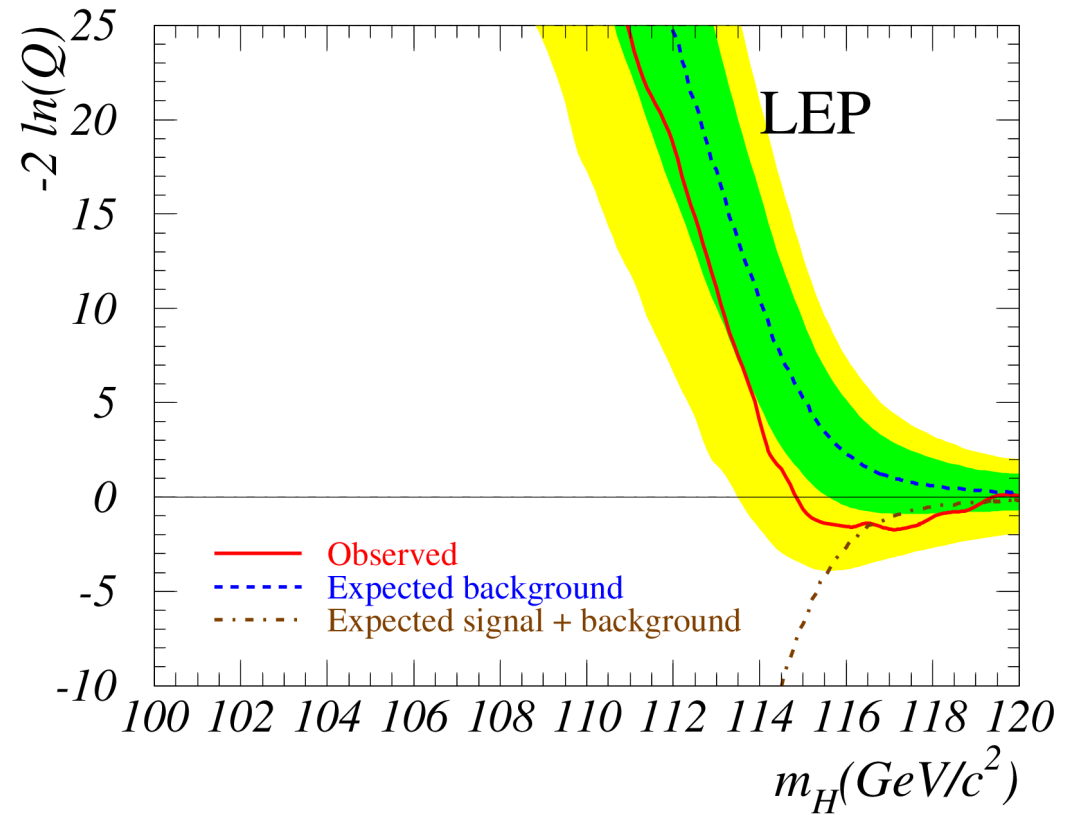
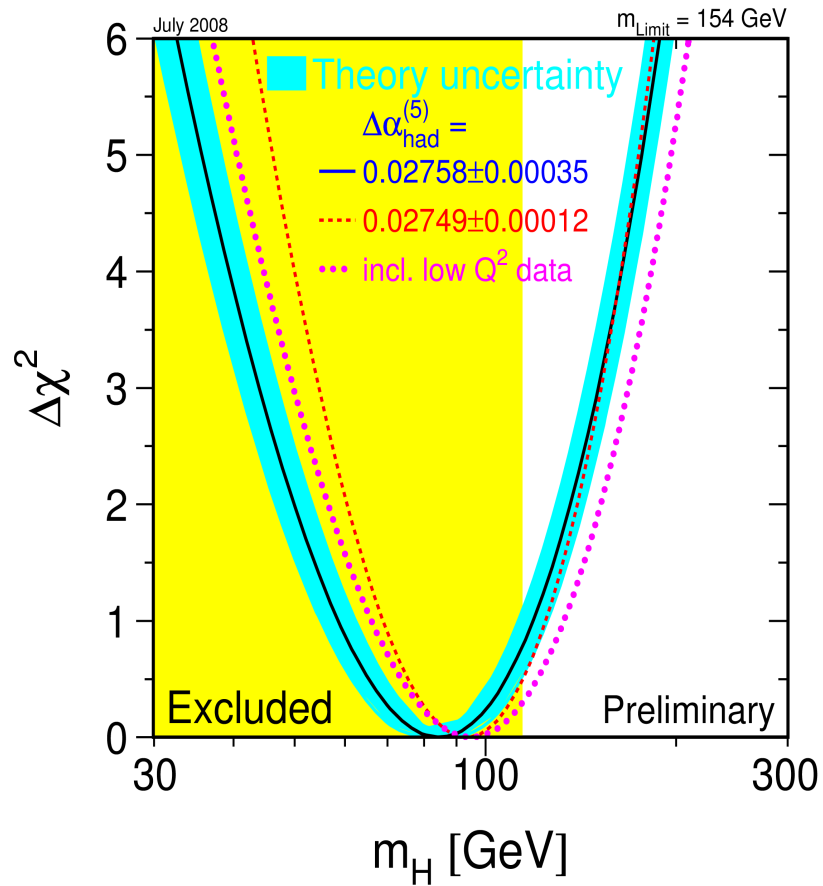
Sezione di Roma III



Outline

- Past and present information on the Higgs boson
- Discussing the hypothesis: $M_h \sim 125 \text{ GeV}$, $\sigma \sim \sigma_{\text{SM}}$ for the SM and the MSSM
- Conclusions

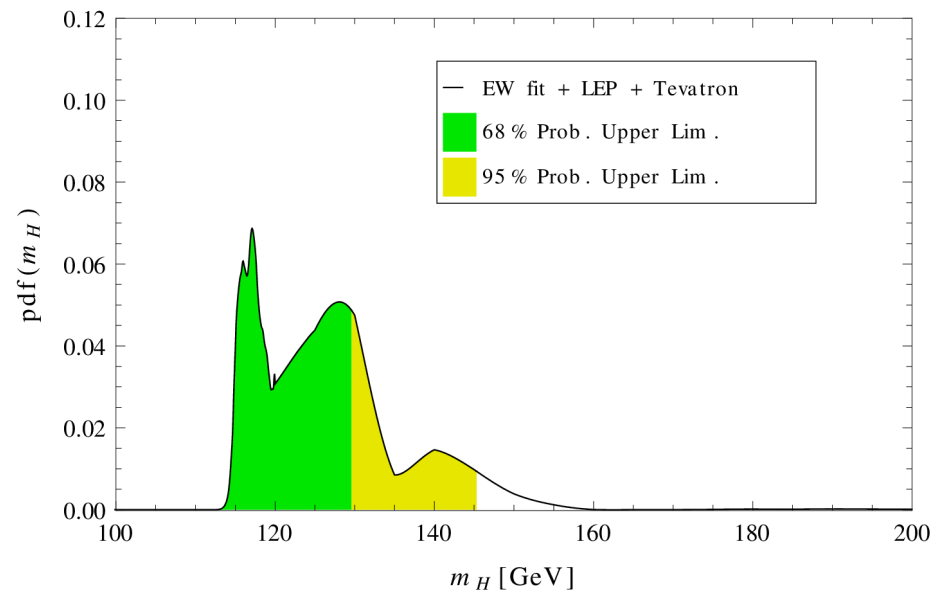
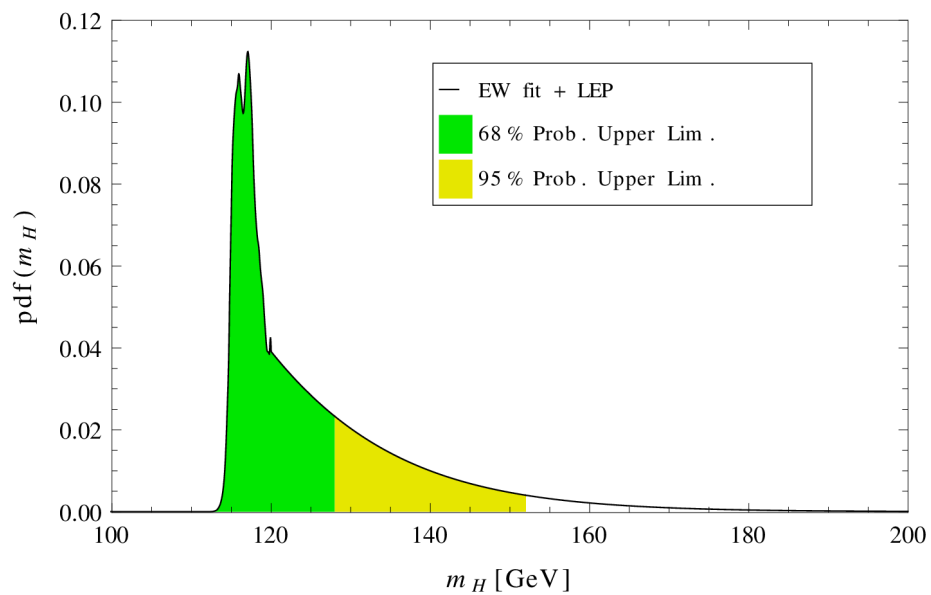
The past: LEP



$$Q = \frac{\mathcal{L}(s + b)}{\mathcal{L}(b)}$$

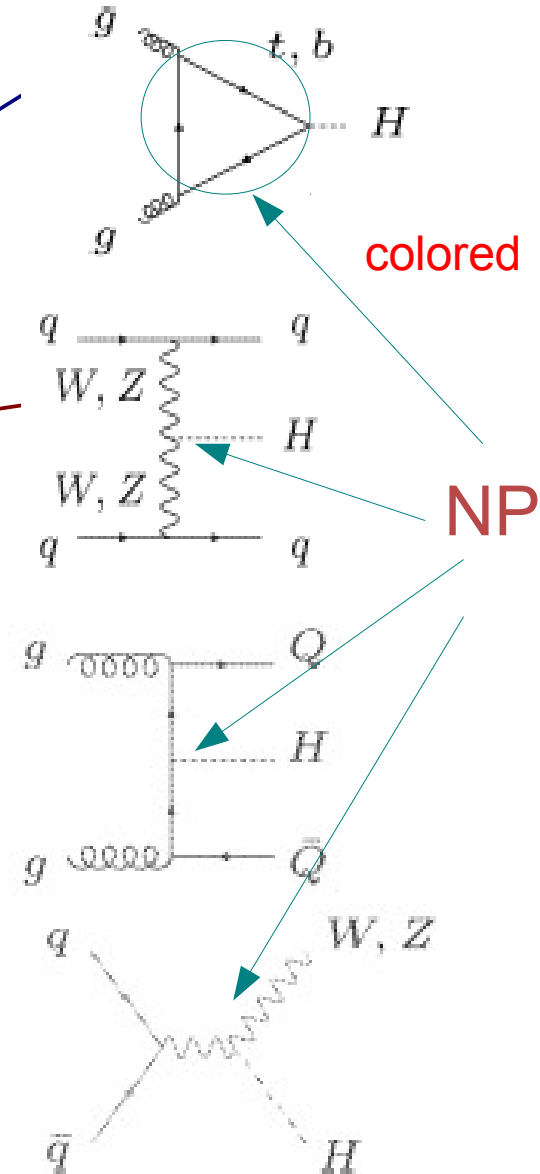
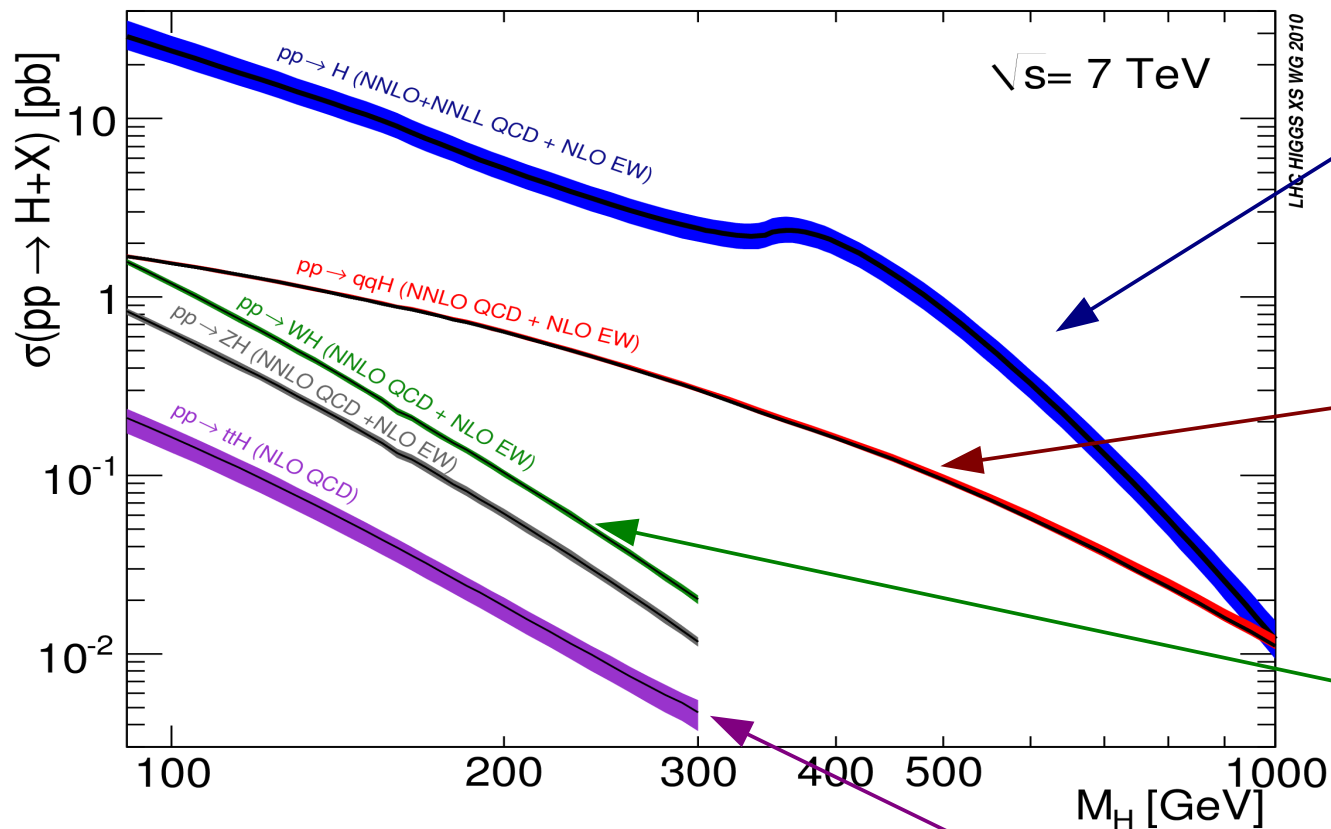
The past: LEP + Tevatron

Combining direct and indirect information:



courtesy of S. Di Vita

The present: LHC, Higgs Production



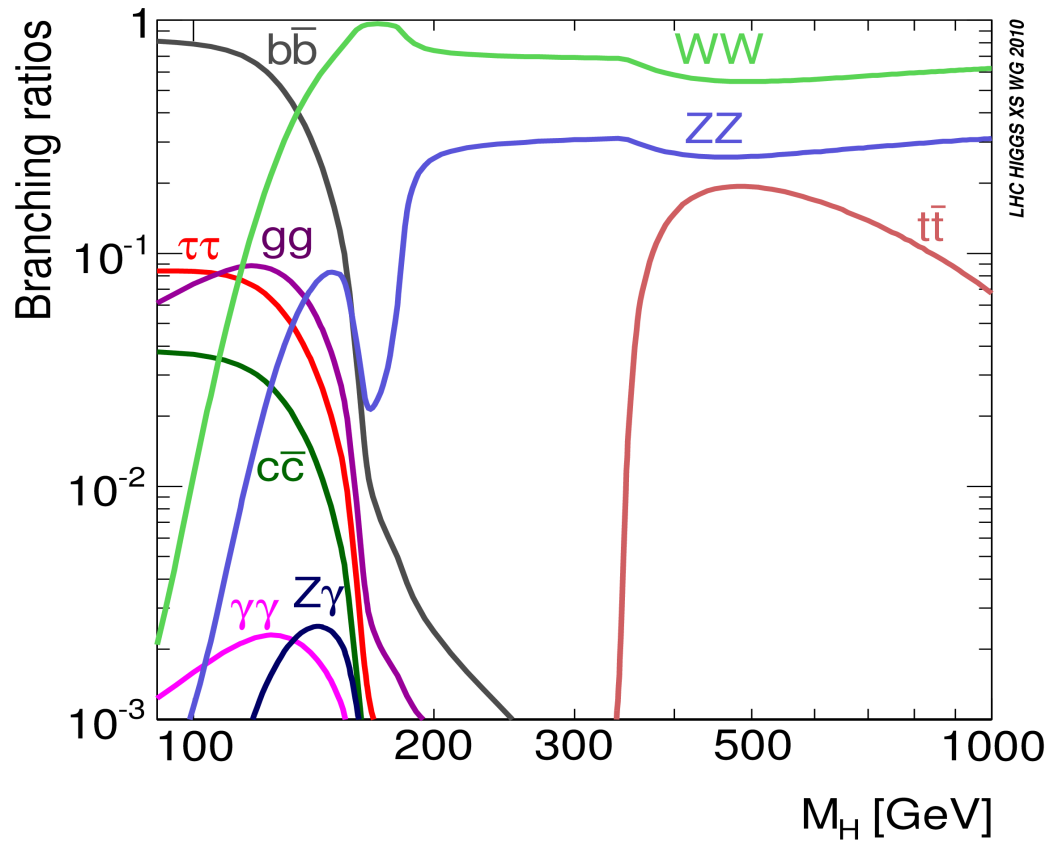
Gluon-fusion process dominant

Weak-boson fusion has a very good-signal/background ratio

Bands include: PDF + α_s + scale uncertainties

Heavy replicas of SM particles contribute to gluon-fusion:
ex. 4th generation

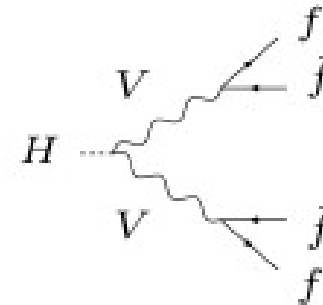
The present: LHC, Higgs Decays



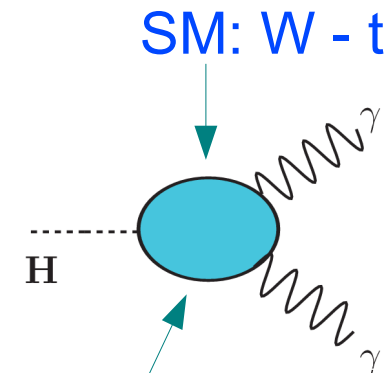
A NP increase in gluon-fusion X-sect. often corresponds to a decrease of $\text{BR}(H \rightarrow \gamma\gamma)$

The $\text{BR}(H \rightarrow \gamma\gamma)$ can increase if NP reduces the other BR's

Golden Channel $V=Z$

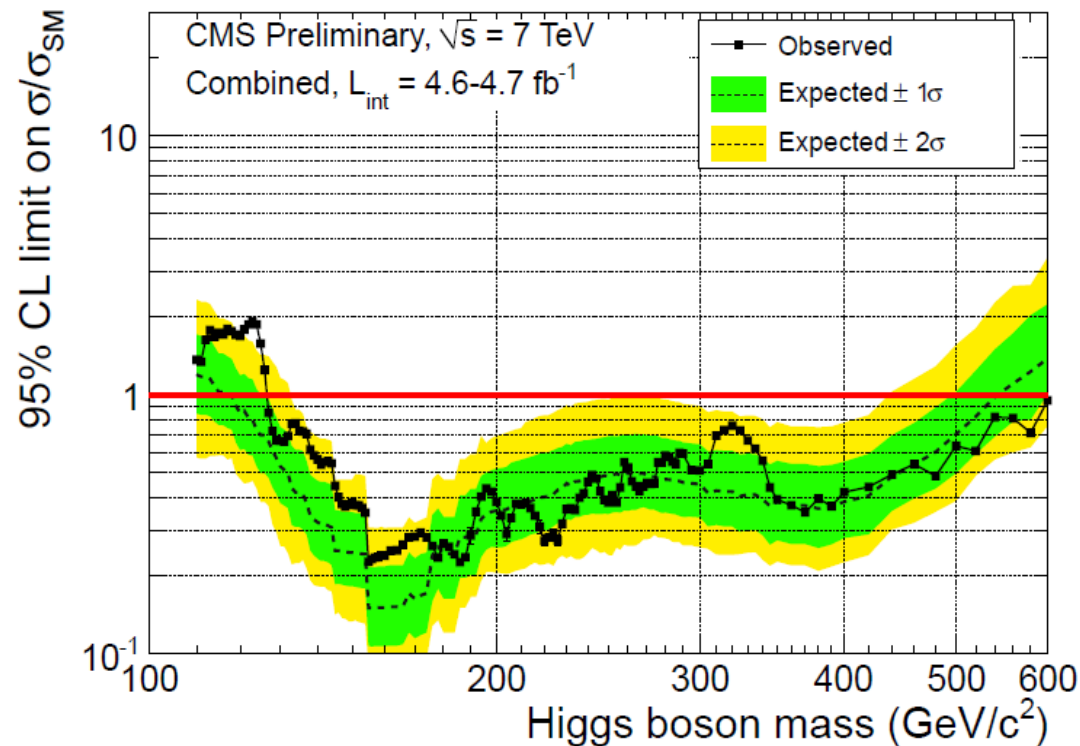
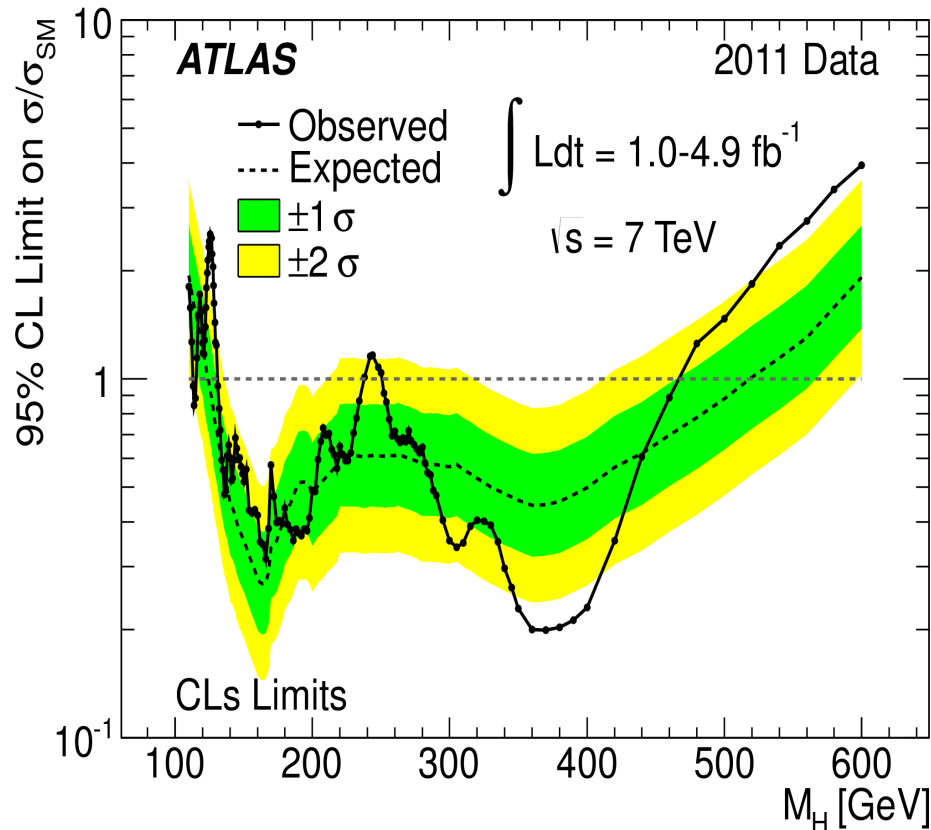


Low Higgs mass



NP: white + colored

The present: LHC, results



Excess of events in $H \rightarrow \gamma\gamma, H \rightarrow ZZ$

ATLAS near $M_H \sim 126 \text{ GeV}$
and near $M_H \sim 245 \text{ GeV}$

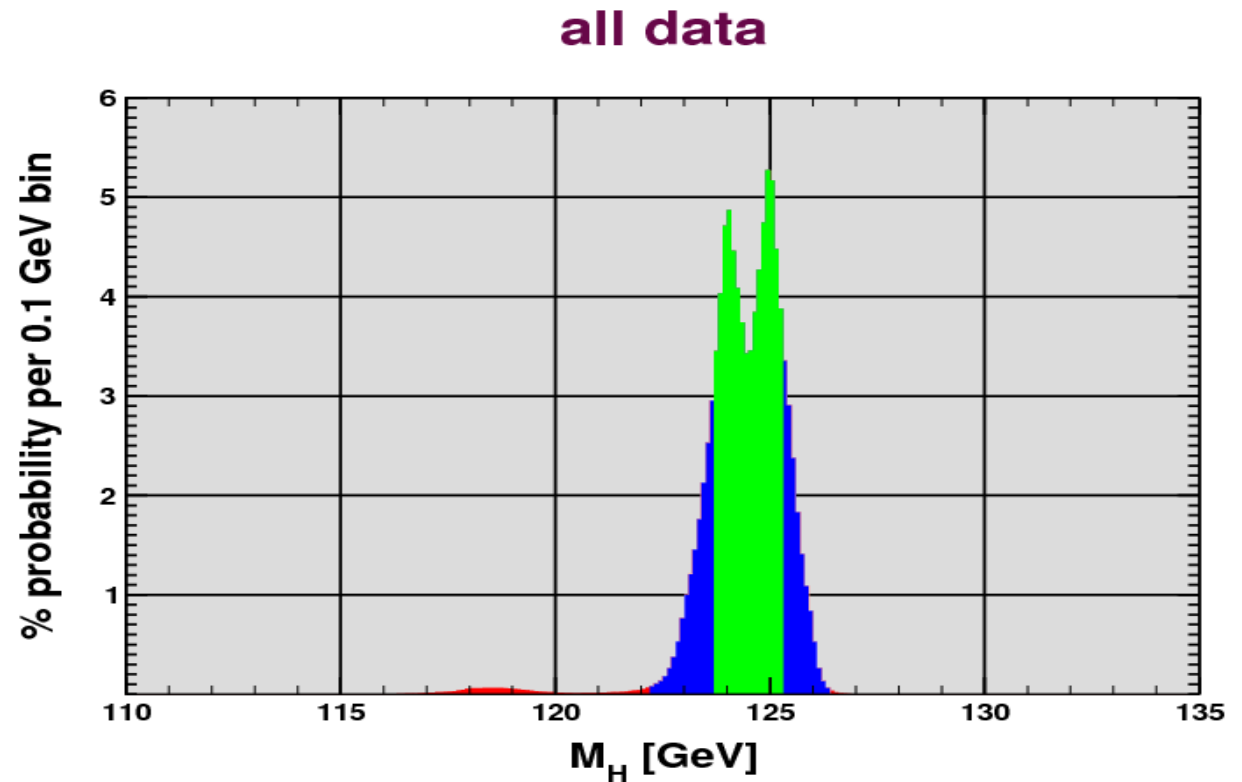
CMS near $M_H \sim 124 \text{ GeV}$
and near $M_H \sim 119.5 \text{ GeV}$

Supported by a broad excess in $H \rightarrow WW$

The present: LHC

This plot should be taken as qualitative.

Q is guessed.
Experiments do not provide likelihoods



Erler, 2012

Working hypothesis: $M_H \sim 125$ GeV, $\sigma \sim \sigma_{\text{SM}}$

but data still allow $M_H > 600$ GeV although this region is cut by EWPT

Reversing the heavy Higgs argument

Specific type of NP could allow a heavy Higgs in the EW fit (“conspiracy”).
Take

$$\sin^2 \theta_{eff}^{lept} \sim \frac{1}{2} \left\{ 1 - \left[1 - \frac{4A^2}{M_Z^2 \hat{\rho} (1 - \Delta \hat{r}_w)} \right]^{1/2} \right\} \quad \begin{aligned} \hat{\rho} &= \rho_0 + \delta\rho \quad (\rho_0^{\text{SM}} = 1, \delta\rho \leftrightarrow (\epsilon_1, T)) \\ \Delta \hat{r}_w &\leftrightarrow (\epsilon_3, S) \end{aligned}$$

$$c_i > 0 \quad \sim \quad (\sin^2 \theta_{eff}^{lept})^\circ + c_1 \ln \left(\frac{M_H}{M_H^\circ} \right) + c_2 \left[\frac{(\Delta\alpha)_h}{(\Delta\alpha)_h^\circ} - 1 \right] - c_3 \left[\left(\frac{M_t}{M_t^\circ} \right)^2 - 1 \right] + \dots$$

To increase the fitted M_H :

$$\left\{ \begin{array}{l} \hat{\rho} > 1 \rightarrow \left\{ \begin{array}{l} \rho_0 > 1 \leftarrow \text{Extra Z} \\ \delta\rho > 0 \leftarrow \text{Isosplitt (s)fermions, Multi Higgs models,} \end{array} \right. \\ \Delta \hat{r}_w < 0 \end{array} \right.$$

↖ Light sleptons

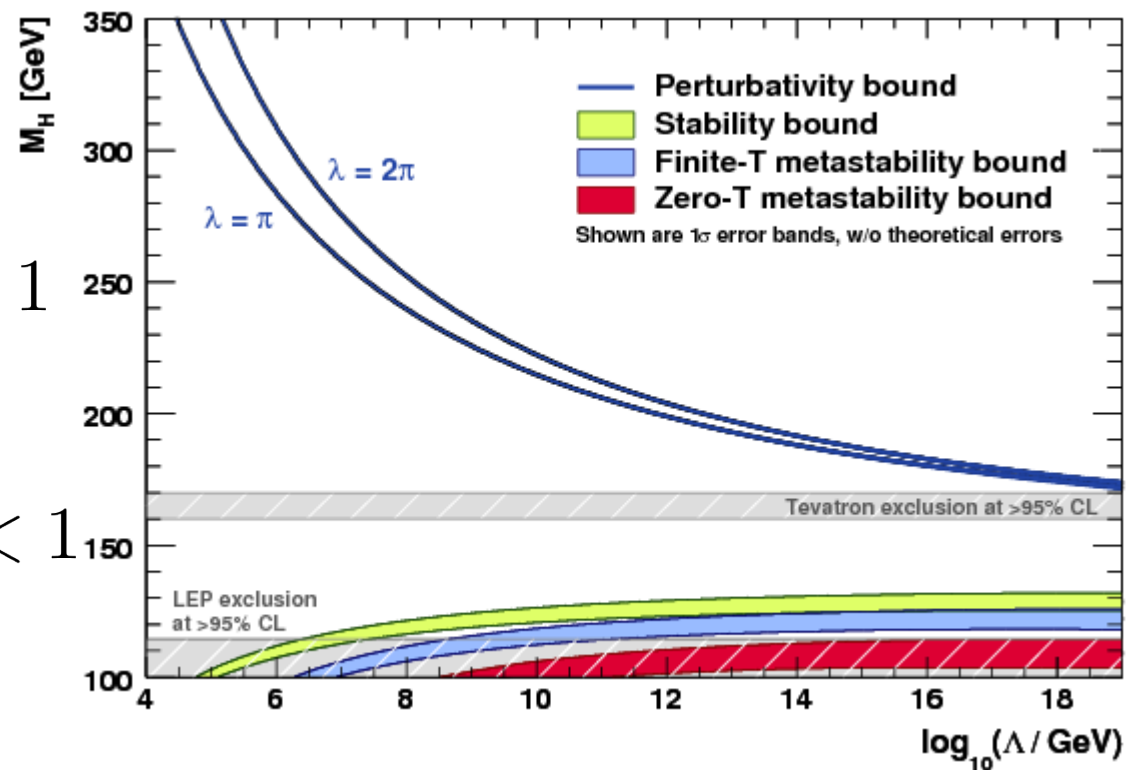
NP (if there) seems to be of the decoupling type
n.b. $M_H > 600$ GeV would point to the conspiracy

Theoretical bounds on the Higgs mass in the SM

$$V_{\text{Higgs}} \approx \frac{\lambda}{4} h^4$$

M_H large $\lambda(m_t) \rightarrow \lambda(\Lambda) \gg 1$
non-perturbative regime

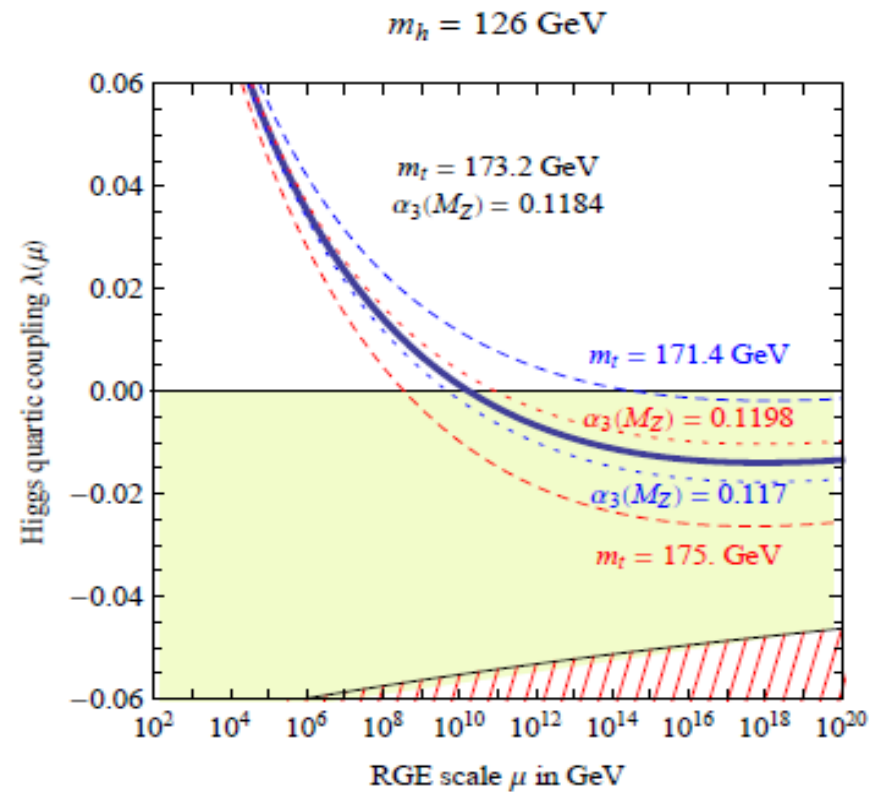
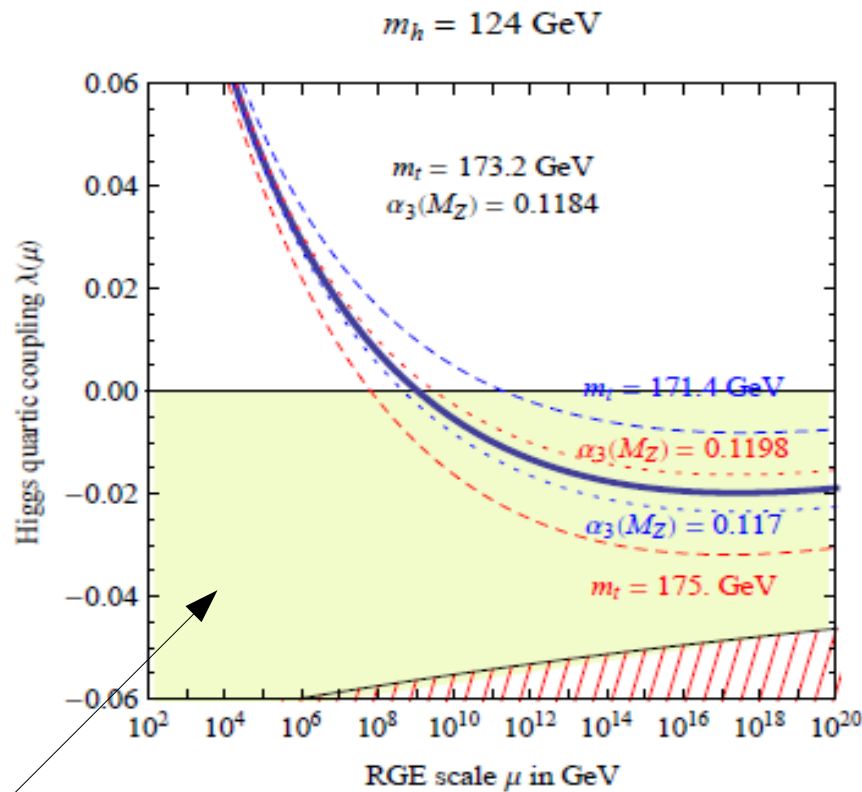
M_H small $\lambda(m_t) \rightarrow \lambda(\Lambda) \ll 1$
vacuum (meta)stability



Ellis et al. 2009

Running depends on m_t, α_s, \dots

$M_H \sim 125$ GeV, no problem with the Landau pole, perturbativity up to the Planck scale



Elias-Miro' et al. 2011

Metastability region

Full stability is at the border. Universe becomes metastable at $\Lambda \sim 10^{10} \text{ GeV}$.

λ never becomes too negative, small probability of quantum tunneling.

Lifetime of the EW vacuum longer than the age of the Universe.

SM ok up to Planck mass.

$M_h \sim 125 \text{ GeV}$ and the MSSM

Higgs sector:

$$H_1 = \begin{pmatrix} H_1^0 \\ H_1^- \end{pmatrix}, \quad H_2 = \begin{pmatrix} H_2^+ \\ H_2^0 \end{pmatrix} \implies h, H, A, H^\pm$$

β, α

Higgs masses: predicted at the tree level in terms of M_A , $\tan \beta$, $M_h < M_Z$

Including radiative corrections: dependence on all SUSY(-breaking) parameters

$(A_t, A_b, \mu \dots)$

$$\begin{array}{lll} M_h \lesssim 135 \text{ GeV} & \xrightarrow{\text{decoupling}} & h \text{ SM-like} \\ M_{A,H,H^\pm} \sim 100 \dots \text{ TeV} & & M_A \sim M_H \sim M_{H^\pm} > \mathcal{O}(200 \text{ GeV}) \end{array}$$

ϕ	$g_{u\bar{u}}^\phi$	$g_{d\bar{d}}^\phi$	g_{VV}^ϕ
h	$\cos \alpha / \sin \beta \rightarrow 1$	$-\sin \alpha / \cos \beta \rightarrow 1$	$\sin(\beta - \alpha) \rightarrow 1$
H	$\sin \alpha / \sin \beta \rightarrow 1 / \tan \beta$	$\cos \alpha / \cos \beta \rightarrow \tan \beta$	$\cos(\beta - \alpha) \rightarrow 0$
A	$1 / \tan \beta$	$\tan \beta$	0

Large $\tan \beta$

$$g_{d\bar{d}}^\phi \uparrow \text{decoupling} \frac{0}{0}$$

delayed decoupling

How easy is to get $M_H \sim 125$ GeV in the MSSM ?

$$M_h^2 \simeq M_Z^2 c_{2\beta}^2 + \frac{3 m_t^4}{4 \pi^2 v^2} \left[\ln \left(\frac{M_S^2}{m_t^2} \right) + \frac{X_t^2}{M_S^2} \left(1 - \frac{X_t^2}{12 M_S^2} \right) \right] + \dots$$



SUSY breaking parameters

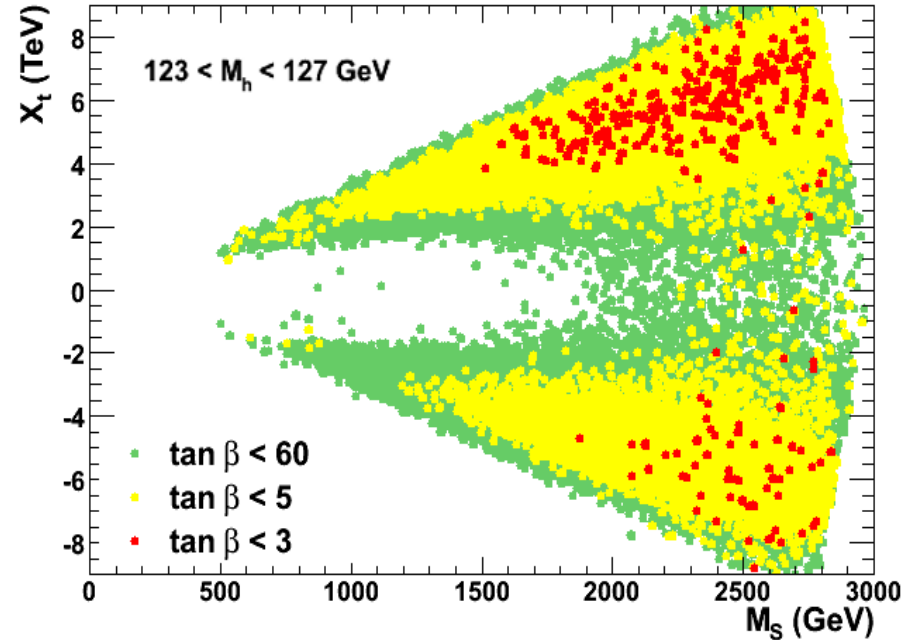
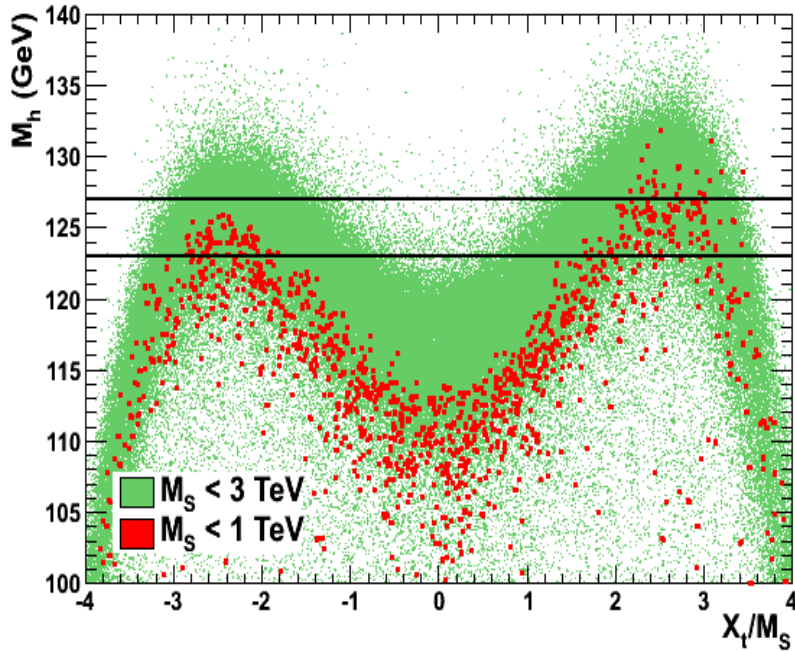
$$X_t = A_t - \mu \cot \beta, \quad M_S = \sqrt{M_{\tilde{t}_1} M_{\tilde{t}_2}}$$

To get $M_H \sim 125$ GeV:

- Large $\tan \beta$, $\tan \beta > 10$ (increase the tree-level)
- Heavy stops, i.e. large M_S (increase the \ln)
- Large stop mixing, i.e. large X_t

The more assumptions we take on the mechanism of SUSY-breaking,
the more difficult becomes to get $M_H \sim 125$ GeV

pMSSM: minimal assumptions on SUSY-breaking parameters



Arbey et al., 2011

22 input parameters varying in the domains:

$$1 \leq \tan \beta \leq 60, \quad 50 \text{ GeV} \leq M_A \leq 3 \text{ TeV}, \quad -9 \text{ TeV} \leq A_f \leq 9 \text{ TeV}, \\ 50 \text{ GeV} \leq m_{\tilde{f}_L}, m_{\tilde{f}_R}, M_3 \leq 3 \text{ TeV}, \quad 50 \text{ GeV} \leq M_1, M_2, |\mu| \leq 1.5 \text{ TeV}.$$

Costrained scenarios:

(no) **GMSB:**

$\tan \beta, \text{sign}(\mu), M_{\text{mess}}, N_{\text{mess}}, \Lambda$

(no) **AMSB:**

$\tan \beta, \text{sign}(\mu), m_0, m_{3/2}$

(yes) **MSUGRA:**

$\tan \beta, \text{sign}(\mu), m_0, m_{1/2}, A_0$

(no) **no-scale:**

$m_0 \approx A_0 \approx 0$

(yes) **VCSSM :**

$m_0 \approx -A_0$

(no) **NMSSM :**

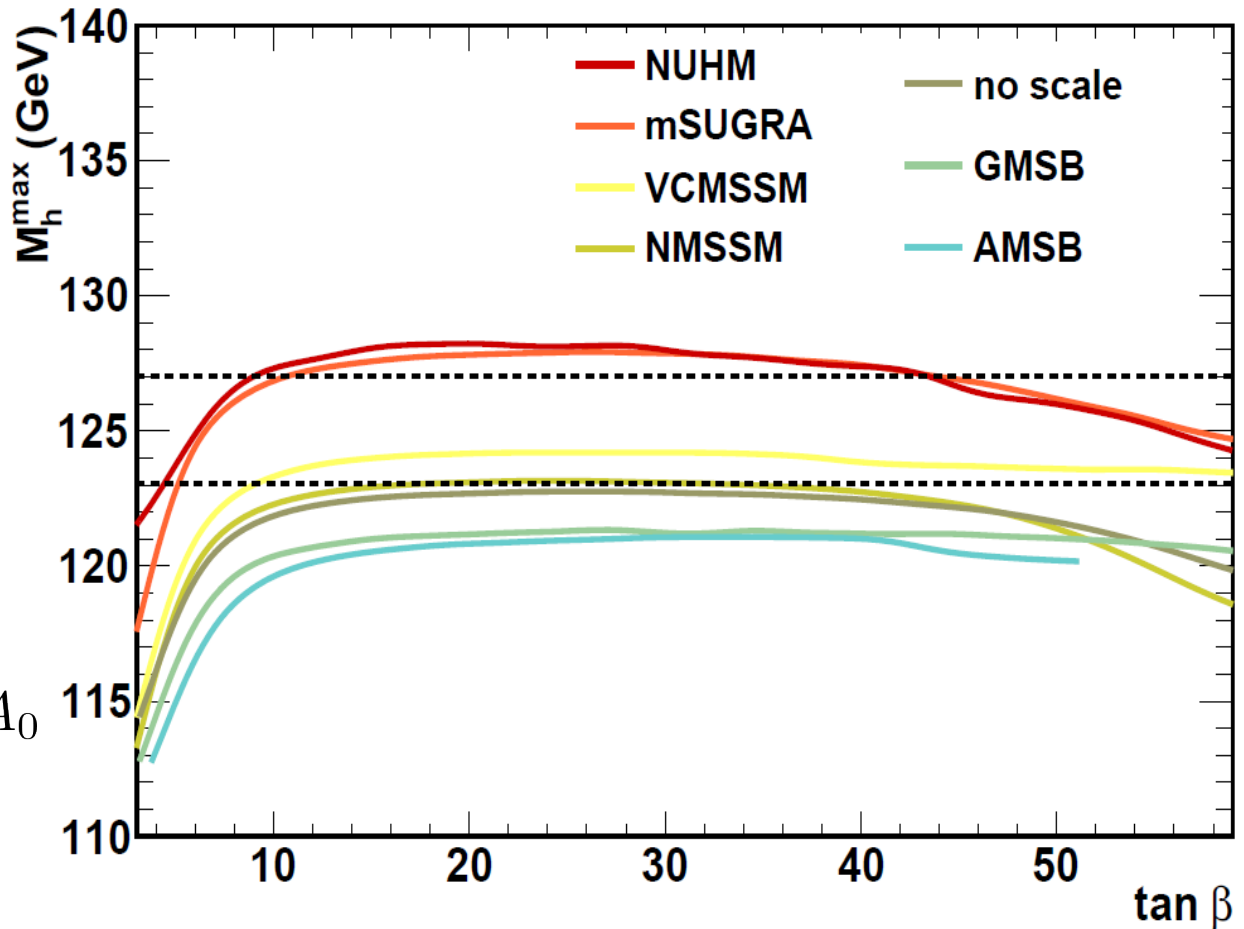
$m_0 \approx 0,$

$A_0 \approx -1/4 m_{1/2}$

(yes) **NUHM:**

non universal m_0

mSUGRA: $50 \text{ GeV} \leq m_0 \leq 3 \text{ TeV}, \quad 50 \text{ GeV} \leq m_{1/2} \leq 3 \text{ TeV}, \quad |A_0| \leq 9 \text{ TeV};$
 GMSB: $10 \text{ TeV} \leq \Lambda \leq 1000 \text{ TeV}, \quad 1 \leq M_{\text{mess}}/\Lambda \leq 10^{11}, \quad N_{\text{mess}} = 1;$
 AMSB: $1 \text{ TeV} \leq m_{3/2} \leq 100 \text{ TeV}, \quad 50 \text{ GeV} \leq m_0 \leq 3 \text{ TeV}.$



Arbey et al., 2011

$\sigma \sim \sigma_{SM}$ and the MSSM

- Squarks and gluinos contribute to the loop-induced gluon fusion production cross section
 - $\sigma(gg \rightarrow h)$ is fully known at NLO QCD (standard + SUSY contributions)
 - $\sigma(gg \rightarrow h)$ implemented in the event generator POWHEG.
E. Bagnaschi, P. Slavich, A. Vicini, G.D. (11)
- a) Interface POWHEG with a mass spectrum generator that provides Higgs masses and couplings.
 - b) Rescale the SM contribution.
 - c) insert the SUSY correction

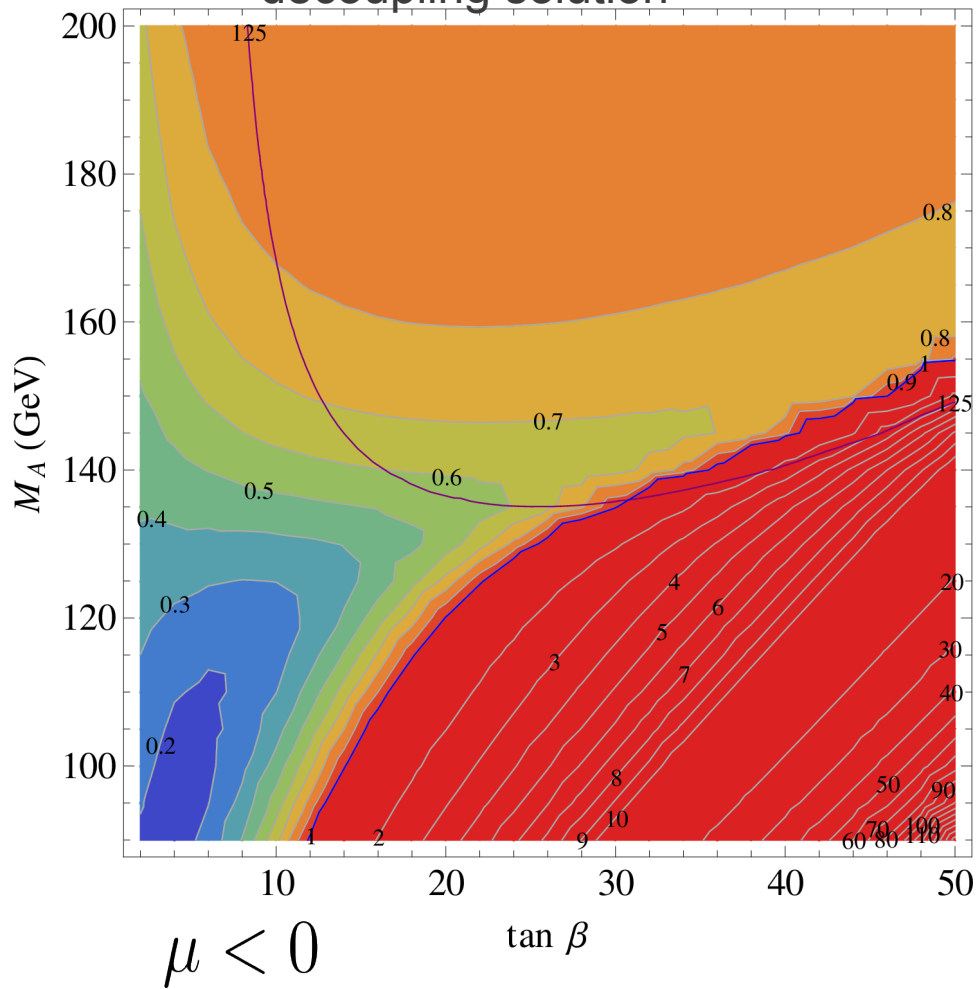
PO(sitive)W(eight)H(ardest)E(mission)G(enerator)

Nason et al. (04--)

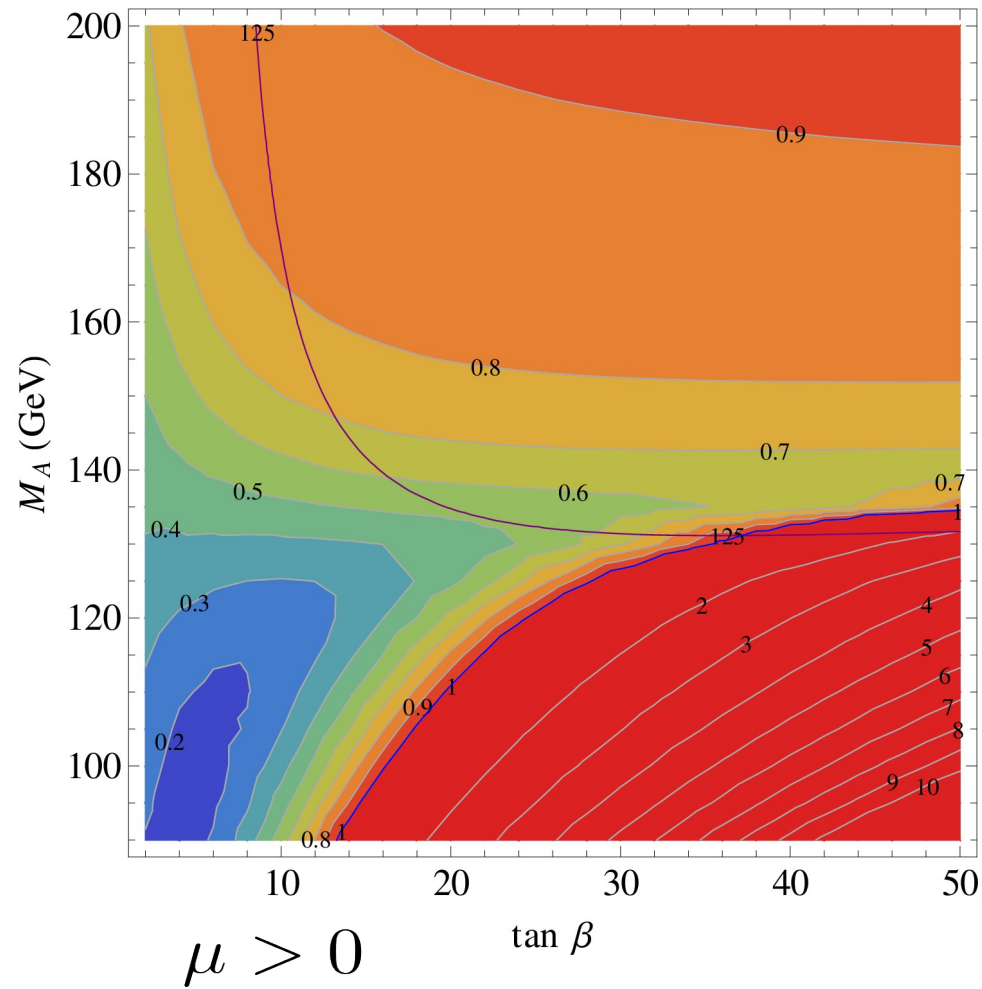
- Matching NLO-QCD matrix elements with Parton Showers
- Generate the hardest emission first, with NLO accuracy, independently of the PS
- Can be interfaces to several SMC programs (HERWIG/PHYTIA)
- Generate events with positive weights
- NLO accuracy of the total cross-section preserved

$$\frac{\sigma(gg \rightarrow h)}{\sigma(gg \rightarrow h_{SM})}$$

decoupling solution



decoupling solution

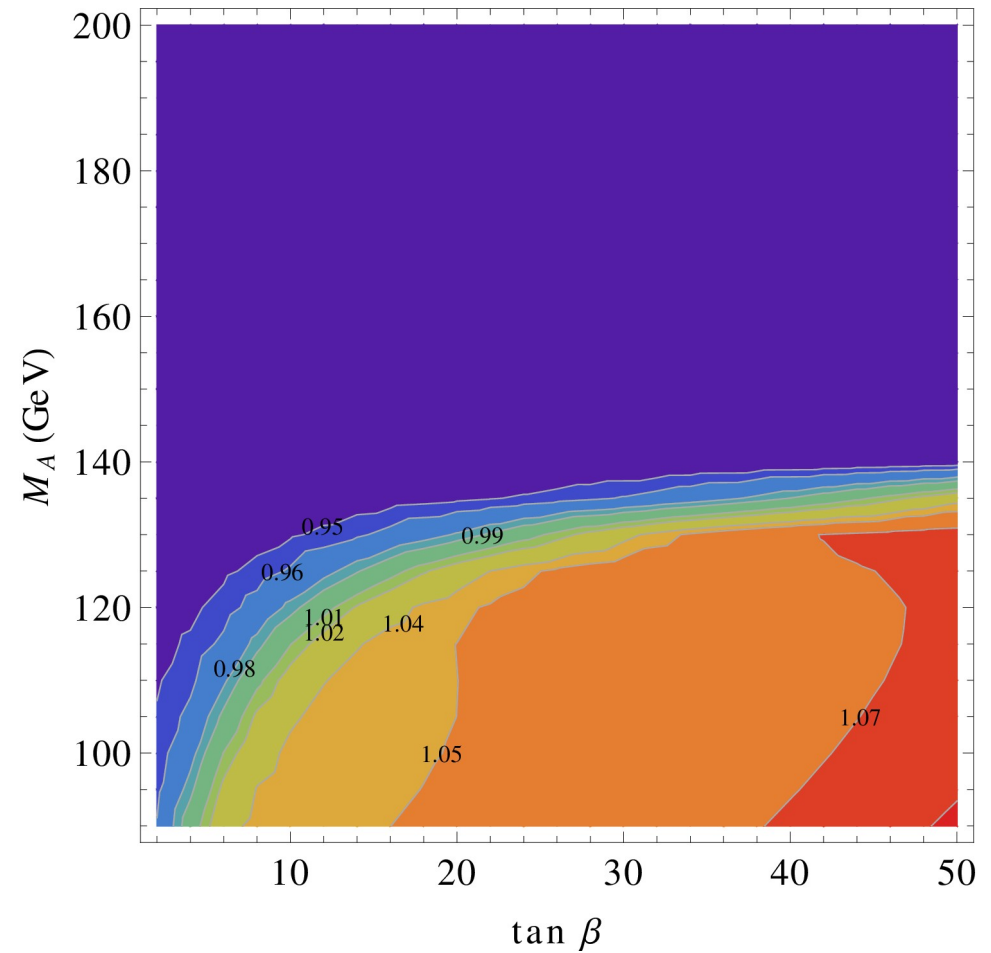
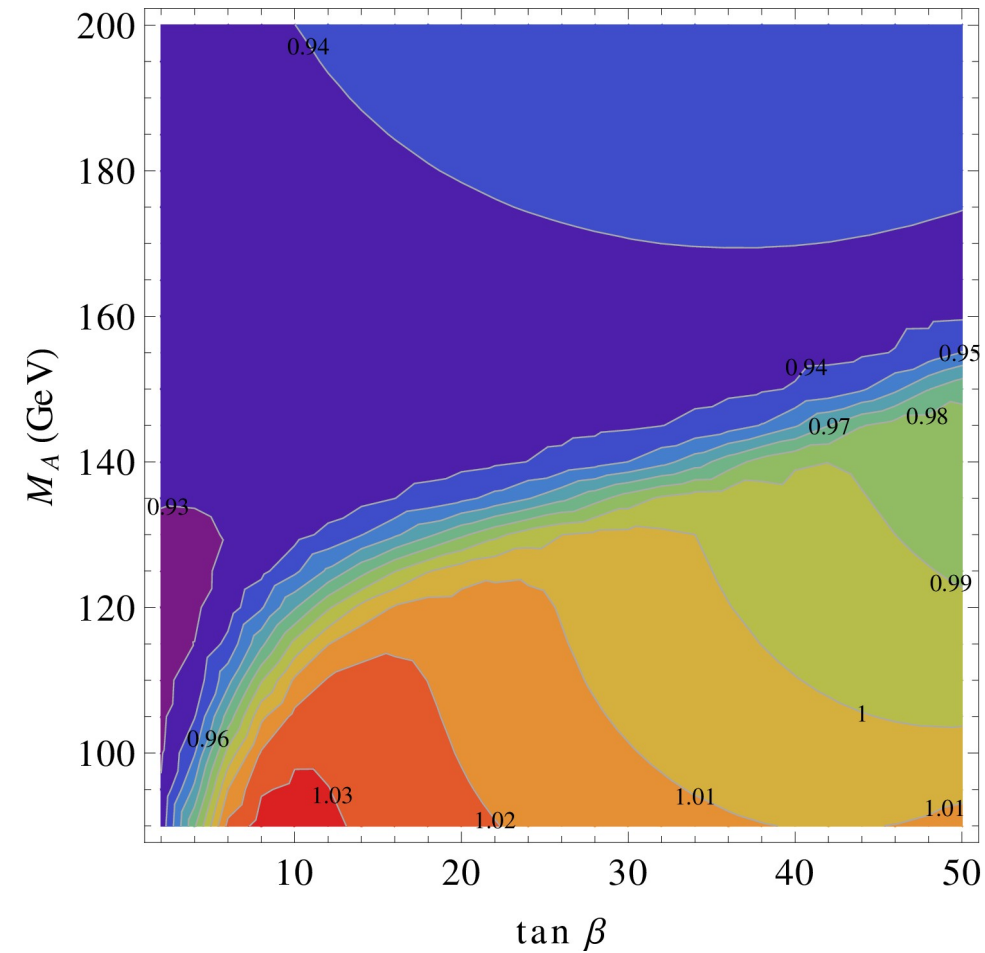


$$m_Q = m_U = m_D = 1000 \text{ GeV}, X_t = A_t - \mu \cot \beta = 2500 \text{ GeV}, M_3 = 800 \text{ GeV}, M_2 = 2 M_1 = 200 \text{ GeV}, |\mu| = 200 \text{ GeV}$$

$$\frac{\sigma(gg \rightarrow h)}{\sigma(gg \rightarrow h_{SM})_{\text{rescaled}}}$$

$\mu < 0$

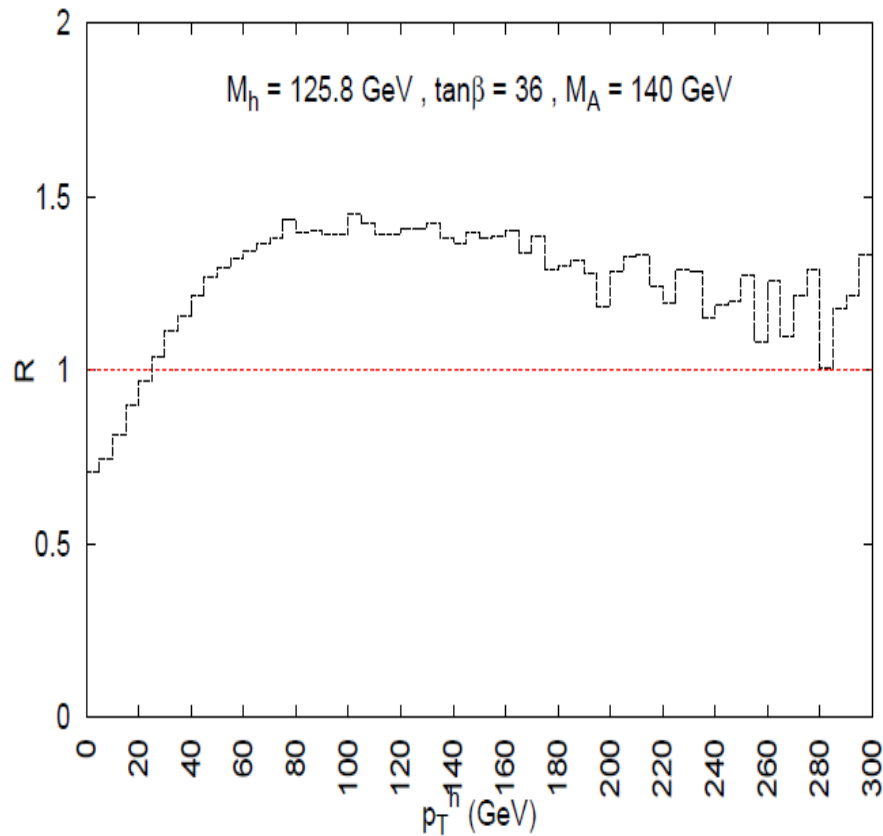
$\mu > 0$



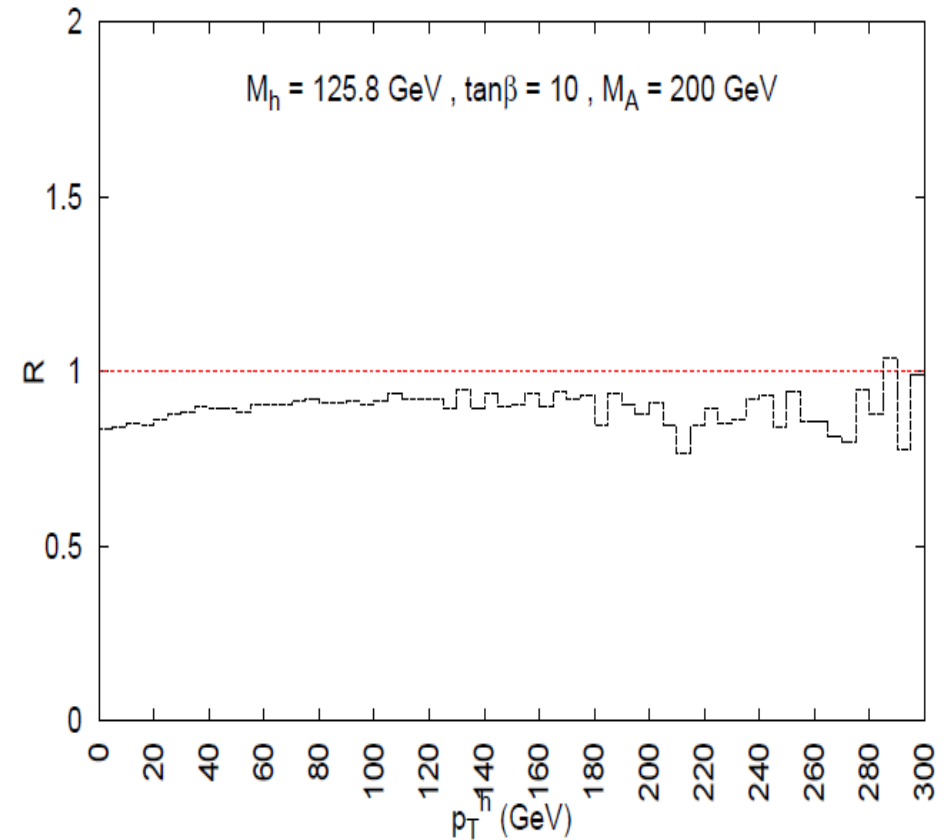
Squarks are heavy: corrections up to 10%

Using the p_t^h to disentangle between SM and MSSM

$$\mu < 0$$



$$\mu > 0$$

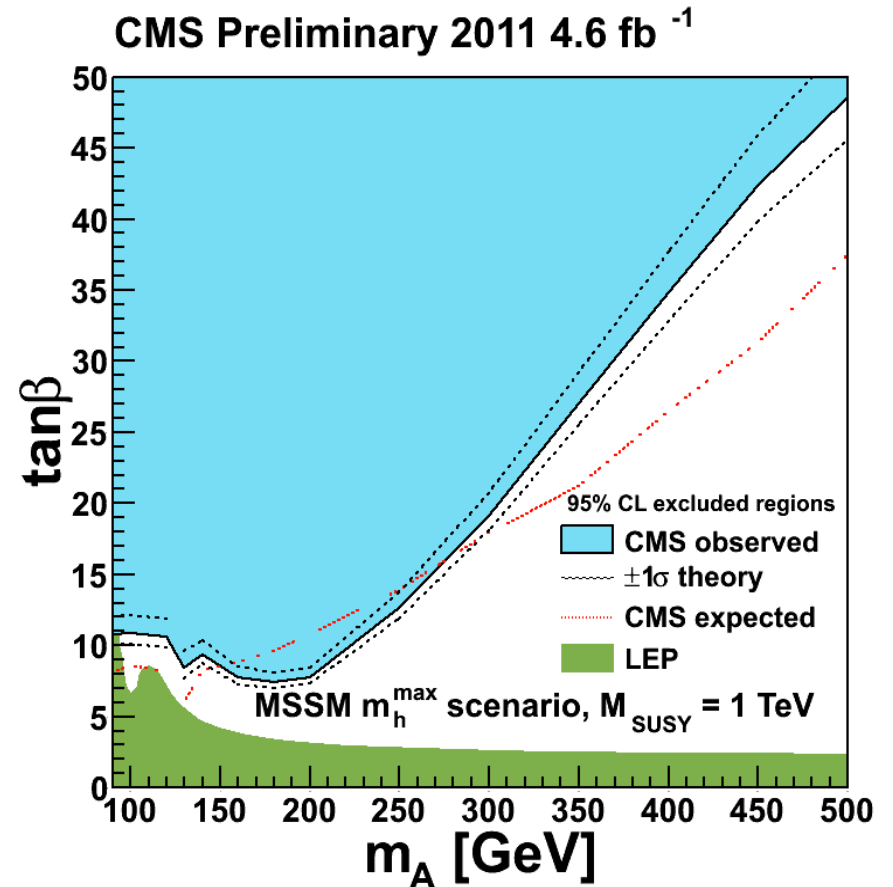
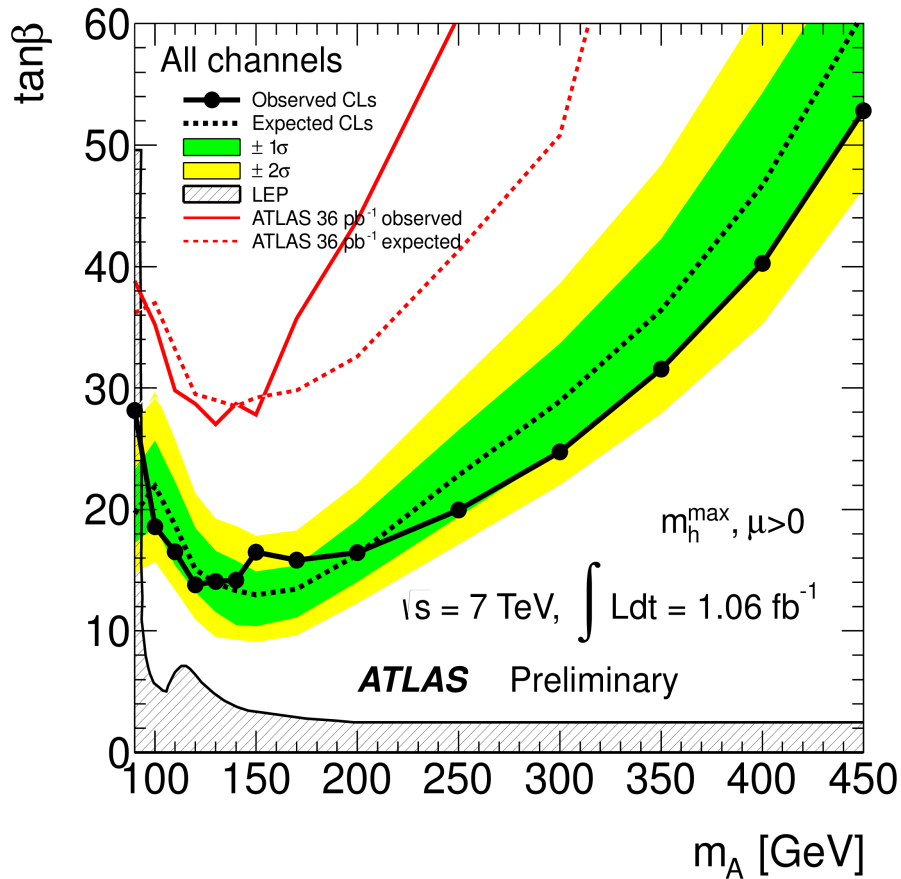


$$R = \text{ratio } \frac{d\sigma}{dp_t^h} \text{ MSSM over SM}$$

obtained using POWHEG + HERWIG

(Bagnaschi et al.)

$\phi \rightarrow \tau\tau$ ($\phi = h, H, A$) kills the non-decoupling solution



The ATLAS, CMS plots represent points in the MSSM parameter space different from ours, the SUSY corrections are not included in these plots, but with these limits

$M_H \sim 125 \text{ GeV}$: Large M_A , to be in the decoupling regime

Conclusions

- It is too early to make any firm statement.
- Personally, I believe that a Higgs boson is in the mass range 116-126(+2):
 $M_h = 121 \pm 5 \text{ GeV}$
- The exact value of the Higgs mass is very important.
A single GeV makes the difference.
- $M_h = 125 \text{ GeV}$ is a very intriguing value.
For the SM it is at the “border” of the stability region.
For the MSSM it is at the ”border” of the mass-predicted region.