

# A Curious Story of Gravity in the Ultraviolet

Galileo Galilei Institute, Firenze

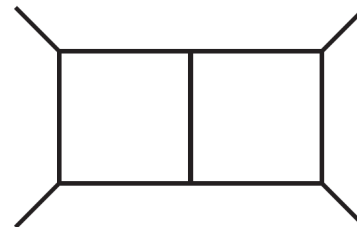
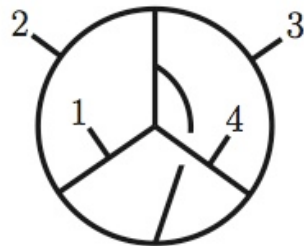
Zvi Bern,

May 18, 2016

UCLA & KITP

**ZB, Clifford Cheung, Huan-Hang Chi, Scott Davies, Lance Dixon, Josh Nohle: arXiv:1507.06118  
& to appear**

**ZB, Huan-Hang Chi, Lance Dixon and Alex Edison: to appear.**



# Quantum Gravity

Often repeated statement:

**“Einstein’s theory of General Relativity is incompatible with quantum mechanics.”**

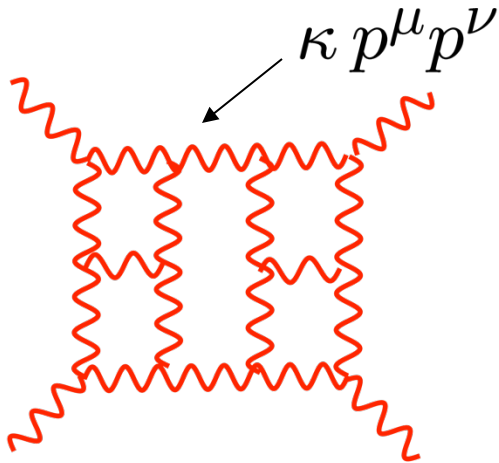
To a large extent this is based on another often repeated statement:

**“All point-like quantum theories of gravity are ultraviolet divergent and non-renormalizable.”**

**Where do these statements come from and are they true?**

# Non-Renormalizability of Gravity?

$$\kappa = \sqrt{32\pi G_N} \quad \leftarrow \text{Dimensionful coupling}$$



**Gravity:** 
$$\int \prod_{i=1}^L \frac{d^D p_i}{(2\pi)^D} \frac{(\kappa p_j^\mu p_j^\nu)}{\text{propagators}}$$

**Gauge theory:** 
$$\int \prod_{i=1}^L \frac{d^D p_i}{(2\pi)^D} \frac{(g p_j^\nu)}{\text{propagators}}$$

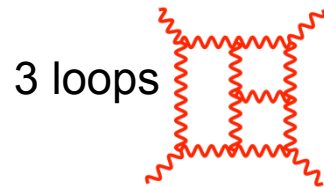
**UV: Large momenta in loop integrals**

- Extra powers of loop momenta in numerator means integrals are badly behaved in the UV and must diverge at some loop order.
- Much more sophisticated power counting in supersymmetric theories but this is basic idea.

- $N = 8$  supergravity is best theory to look at.
- With more supersymmetry expect better UV properties.
- High symmetry implies simplicity.

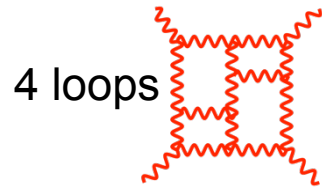
# Feynman Diagrams for Gravity

**Suppose we want to check UV properties of gravity theories:**

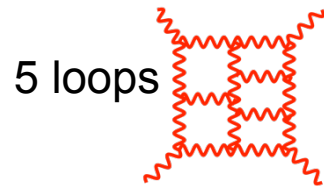


$\sim 10^{20}$   
TERMS

No surprise it has  
never been  
calculated via  
Feynman diagrams.



$\sim 10^{26}$   
TERMS



$\sim 10^{31}$   
TERMS

More terms than  
atoms in your brain!

- Calculations to settle this seemed utterly hopeless!
- Seemed destined for dustbin of undecidable questions.

**Supersymmetry helps, but not enough to make a difference.**

# Where is First Potential $D = 4$ UV Divergence?

<b>3 loops</b> <b><math>N = 8</math></b>	Green, Schwarz, Brink (1982); Howe and Stelle (1989); Marcus and Sagnotti (1985)	<b>X</b>
<b>5 loops</b> <b><math>N = 8</math></b>	Bern, Dixon, Dunbar, Perelstein, Rozowsky (1998); Howe and Stelle (2003,2009)	<b>X</b>
<b>6 loops</b> <b><math>N = 8</math></b>	Howe and Stelle (2003)	<b>X</b>
<b>7 loops</b> <b><math>N = 8</math></b>	Grisaru and Siegel (1982); Bossard, Howe, Stelle (2009); Vanhove; Björnsson, Green (2010); Kiermaier, Elvang, Freedman (2010); Ramond, Kallosh (2010); Biesert et al (2010); Bossard, Howe, Stelle, Vanhove (2011)	<b>?</b>
<b>3 loops</b> <b><math>N = 4</math></b>	Bossard, Howe, Stelle, Vanhove (2011)	<b>X</b>
<b>4 loops</b> <b><math>N = 5</math></b>	Bossard, Howe, Stelle, Vanhove (2011)	<b>X</b>
<b>4 loops</b> <b><math>N = 4</math></b>	Vanhove and Tourkine (2012)	<b>✓</b>
<b>9 loops</b> <b><math>N = 8</math></b>	Berkovits, Green, Russo, Vanhove (2009)	<b>X</b>

ZB, Kosower, Carrasco, Dixon, Johansson, Roiban; ZB, Davies, Dennen, A. Smirnov, V. Smirnov; series of calculations.

**Don't bet on divergence**

**Weird structure.  
Anomaly-like behavior  
of divergence.**

**retracted**

- Conventional wisdom holds that it will diverge soon or later.
- But every detailed prediction either wrong or misleading.

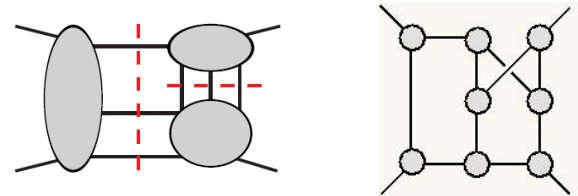
# Our Basic Tools

**We have powerful tools for computing scattering amplitudes and studying their UV properties:**

- **Generalized unitarity method.**

ZB, Dixon, Dunbar, Kosower

ZB, Carrasco, Johansson, Kosower



- **Duality between color and kinematics. Gravity scattering amplitudes directly from gauge-theory ones. Double copy.**

ZB, Carrasco and Johansson (BCJ)

- **Advanced loop-integration technology.**

Chetyrkin, Kataev and Tkachov; Laporta; A.V. Smirnov; V. A. Smirnov; Vladimirov; Marcus, Sagnotti; Czakon; Laporta; etc.

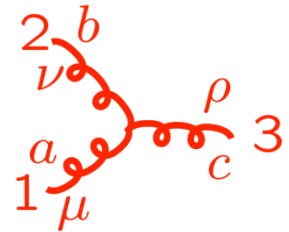
- **I won't explain these tools these but they underlie everything.**
- **Many other tools and advances that I won't discuss here.**

# Duality Between Color and Kinematics

ZB, Carrasco, Johansson (BCJ)

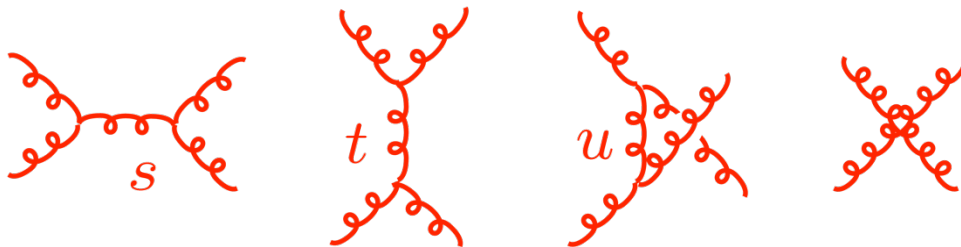
coupling constant  $\rightarrow$  color factor  $\rightarrow$  momentum dependent kinematic factor

$$-g f^{abc} (\eta_{\mu\nu} (k_1 - k_2)_\rho + \text{cyclic})$$



Color factors based on a Lie algebra:  $[T^a, T^b] = i f^{abc} T^c$

Jacobi Identity  $f^{a_1 a_2 b} f^{b a_4 a_3} + f^{a_4 a_2 b} f^{b a_3 a_1} + f^{a_4 a_1 b} f^{b a_2 a_3} = 0$



Use  $1 = s/s = t/t = u/u$   
to assign 4-point diagram  
to others.

$$s = (k_1 + k_2)^2 \quad t = (k_1 + k_4)^2$$

$$u = (k_1 + k_3)^2$$

$$\mathcal{A}_4^{\text{tree}} = g^2 \left( \frac{n_s c_s}{s} + \frac{n_t c_t}{t} + \frac{n_u c_u}{u} \right)$$

Color factors satisfy Jacobi identity:

Numerator factors satisfy similar identity:

$$c_u = c_s - c_t$$

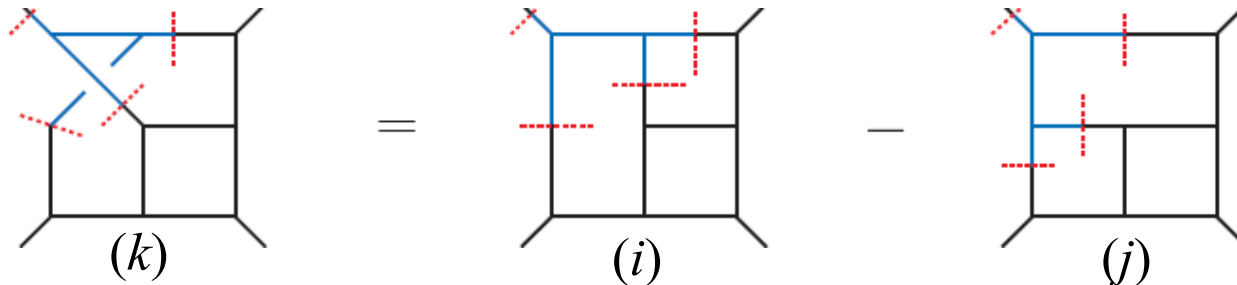
$$n_u = n_s - n_t$$

**Proven at tree level and conjectured at loop level.**

# Duality Between Color and Kinematics

ZB, Carrasco, Johansson (BCJ)

**Conjecture:** kinematic numerators exist with same algebraic properties as color factors, even at loop level.



color factor

$$C_k = C_i - C_j$$

$$n_k = n_i - n_j$$

kinematic numerator

$$n_i \sim k_1 \cdot l_1 k_3 \cdot l_2 \varepsilon_1 \cdot l_3 \varepsilon_2 \cdot k_3 \varepsilon_3 \cdot l_2 \varepsilon_4 \cdot k_3 + \dots$$

If you have a set of duality satisfying kinematic numerators.

**gauge theory  $\longrightarrow$  gravity theory**

simply take

**color factor  $\longrightarrow$  kinematic numerator**

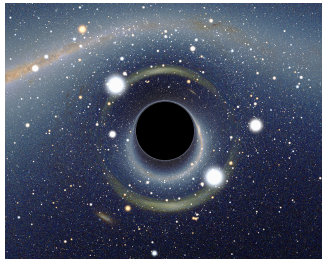
$$C_i \longrightarrow n_i$$

You would never know this from studying the respective Lagrangians. Nor is this understood from string theory.

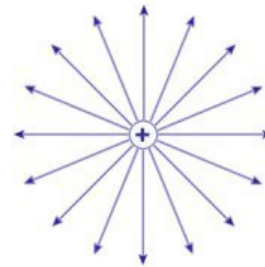
# Applications to Black Hole Physics

Wouldn't it be really cool if every classical solution in gravity could be mapped to a double of gauge theory classical solutions?

**Where to start?** Obviously the coolest place possible: black holes.



**black hole**



**point charge**

Monteiro, O'Connell and White

**Special coordinates: Kerr-Schild coordinates:**

**Schwarzschild  
black hole**

$$g_{\mu\nu} = \eta_{\mu\nu} + \phi k_{\mu} k_{\nu} \quad \phi(r) = \frac{2m}{r}$$

**Coulomb  
point charge**

$$A_{\mu} = \phi k_{\mu} \quad \phi(r) = \frac{Q}{r}$$

$k$  is null

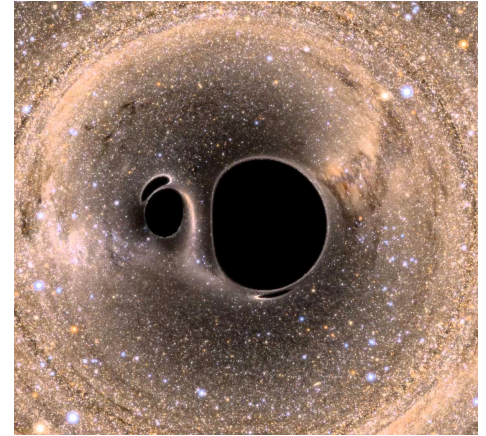
$$\text{Schwarzschild} \sim (\text{Coulomb})^2$$

# Applications to Black Hole Physics

## A variety of other cases:

- **Kerr (rotating) black hole.**
- **Taub-NUT space.**
- **Solutions with cosmological constant.**
- **Radiation from accelerating black hole.**

Luna, Monteiro, O'Connell and White;  
Luna, Monteiro, NicholSEN, O'Connell and White;  
Ridgway and Wise.



It may be possible to extend this to more general cases.  
Need good coordinates!

# UV in Gravity

Most theorists believe that UV properties of quantum field theories of gravity are “well understood”, up to “minor” details, e.g. the precise loop order where divergences occur.

**The main purpose of my talk is to try to convince you that the UV structure of gravity is strange and surprising and most certainly *not* “well understood”.**

- 1. Examples of no divergence even when no known symmetry arguments prevent them. “Enhanced cancellations”.  
Unlike gauge theory.**
- 2. When UV divergences are present in pure (super) gravity, properties are strange and unexpected.**

# Predictions of Ultraviolet Divergences

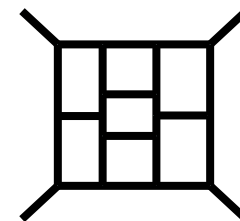
Bossard, Howe, Stelle; Elvang, Freedman, Kiermaier; Green, Russo, Vanhove ; Green and Björnsson ; Bossard , Hillmann and Nicolai; Ramond and Kallosh; Broedel and Dixon; Elvang and Kiermaier; Beisert, Elvang, Freedman, Kiermaier, Morales, Stieberger; Bossard, Howe, Stelle, Vanhove, etc

- **First quantized formulation of Berkovits' pure-spinor formalism.**  
Björnsson and Green
- **Unitarity method.**  
ZB, Davies, Dennen

**Key point:** *all* supersymmetry cancellations are exposed.

**Poor UV behavior, unless new types of cancellations between diagrams exist that are “not consequences of supersymmetry in any conventional sense”**  
Björnsson and Green

- $N = 8$  sugra should diverge at 5 loops in  $D = 24/5$ . ?
- $N = 8$  sugra should diverge at 7 loops in  $D = 4$ . ?
- $N = 4$  sugra should diverge at 3 loops in  $D = 4$ . X
- $N = 5$  sugra should diverge at 4 loops in  $D = 4$ . X

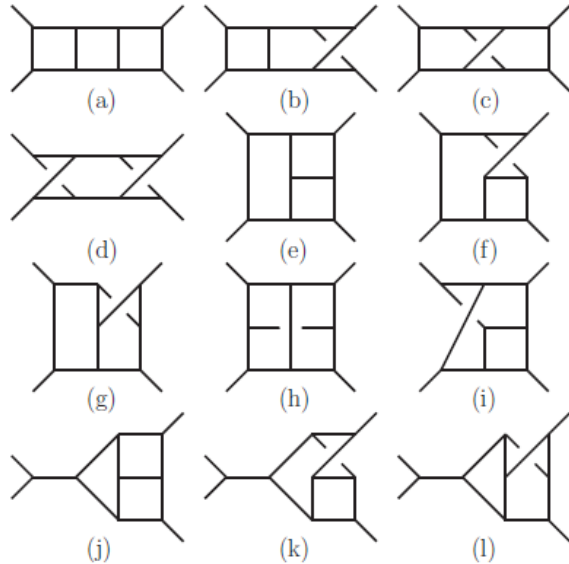


**Consensus agreement from all methods**

**These new types of cancellations do exist: “enhanced cancellations”.**

# $N = 4$ Supergravity UV Cancellation

$$D = 4 - 2\epsilon \quad \text{ZB, Davies, Dennen, Huang}$$



Graph	(divergence)/((12) <sup>2</sup> [34] <sup>2</sup> stA <sup>tree</sup> ( $\frac{\kappa}{2}$ ) <sup>8</sup> )
(a)-(d)	0
(e)	$\frac{263}{768} \frac{1}{\epsilon^3} + \frac{205}{27648} \frac{1}{\epsilon^2} + \left(-\frac{5551}{768} \zeta_3 + \frac{326317}{110592}\right) \frac{1}{\epsilon}$
(f)	$-\frac{175}{2304} \frac{1}{\epsilon^3} - \frac{1}{4} \frac{1}{\epsilon^2} + \left(\frac{593}{288} \zeta_3 - \frac{217571}{165888}\right) \frac{1}{\epsilon}$
(g)	$-\frac{11}{36} \frac{1}{\epsilon^3} + \frac{2057}{6912} \frac{1}{\epsilon^2} + \left(\frac{10769}{2304} \zeta_3 - \frac{226201}{165888}\right) \frac{1}{\epsilon}$
(h)	$-\frac{3}{32} \frac{1}{\epsilon^3} - \frac{41}{1536} \frac{1}{\epsilon^2} + \left(\frac{3227}{2304} \zeta_3 - \frac{3329}{18432}\right) \frac{1}{\epsilon}$
(i)	$\frac{17}{128} \frac{1}{\epsilon^3} - \frac{29}{1024} \frac{1}{\epsilon^2} + \left(-\frac{2087}{2304} \zeta_3 - \frac{10495}{110592}\right) \frac{1}{\epsilon}$
(j)	$-\frac{15}{32} \frac{1}{\epsilon^3} + \frac{9}{64} \frac{1}{\epsilon^2} + \left(\frac{101}{12} \zeta_3 - \frac{3227}{1152}\right) \frac{1}{\epsilon}$
(k)	$\frac{5}{64} \frac{1}{\epsilon^3} + \frac{89}{1152} \frac{1}{\epsilon^2} + \left(-\frac{377}{144} \zeta_3 + \frac{287}{432}\right) \frac{1}{\epsilon}$
(l)	$\frac{25}{64} \frac{1}{\epsilon^3} - \frac{251}{1152} \frac{1}{\epsilon^2} + \left(-\frac{835}{144} \zeta_3 + \frac{7385}{3456}\right) \frac{1}{\epsilon}$

**All three-loop divergences and subdivergences cancel completely!**

**Still no standard-symmetry explanation, despite valiant attempt.**

Bossard, Howe, Stelle; ZB, Davies, Dennen

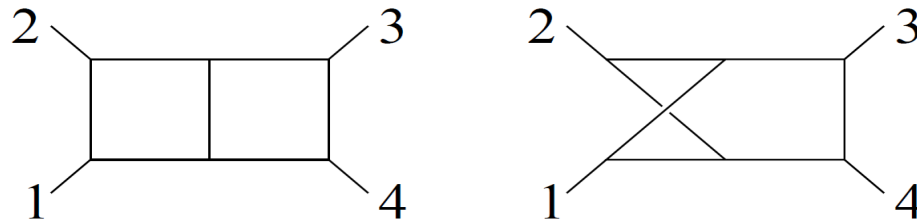
**Prediction based on supergravity imply divergences.  
A nontrivial example of “enhanced cancellations”.**

# Where does new magic come from?

ZB, Davies, Dennen, Huang; Bossard, Howe, Stelle

To analyze we need a simpler example: **Half-maximal supergravity in  $D = 5$  at 2 loops.**

Similar to  $N = 4, D = 4$  sugra at 3 loops, except much simpler.



## Quick summary:

- Finiteness in  $D = 5$  tied to double-copy structure.
- Cancellations in certain forbidden gauge-theory color structures imply hidden UV cancellations in supergravity, even though no standard symmetry explanation.

- Double copy structure implies extra cancellations!
- Quite a nonstandard explanation for a cancellation.

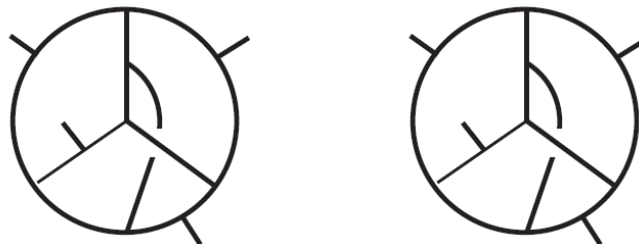
Unfortunately, not easy to extend beyond 2 loops.

# $N = 4$ Supergravity at Four Loops

ZB, Davies, Dennen, Smirnov, Smirnov

We also calculated four-loop divergence in  $N = 4$  supergravity.

$N = 4$  sugra:  $(N = 4 \text{ sYM}) \times (N = 0 \text{ YM})$



Integration uses state-of-the-art software developed for QCD.  
Industrial strength software needed: FIRE5 and special  
purpose C++ code.

# The 4 loop Divergence of $N = 4$ Supergravity

ZB, Davies, Dennen, A.V. Smirnov, V.A. Smirnov

$$\mathcal{M}^{4\text{-loop}} \Big|_{\text{div.}} = \frac{1}{(4\pi)^8} \frac{1}{\epsilon} \left(\frac{\kappa}{2}\right)^{10} \frac{1}{144} (1 - 264\zeta_3) \mathcal{T}$$

dim. reg. UV pole

kinematic factor



$$D = 4 - 2\epsilon$$

It diverges but it has strange properties:

- Contributions to helicity configurations that vanish were it not for a quantum anomaly in  $U(1)$  subgroup of duality symmetry.
- These helicity configuration have vanishing integrands in  $D = 4$ . Divergence is 0/0. Anomaly-like behavior not found in  $N \geq 5$  sugra.

Carrasco, Kallosh, Tseytlin and Roiban

**Motivates closer examination of divergences.**

**Want simpler example: Pure Einstein gravity is simpler.**

# Pure Einstein Gravity

Standard argument for 1 loop finiteness of pure gravity:

't Hooft and Veltman (1974)

$$\cancel{R^2} \quad \cancel{R_{\mu\nu}^2}$$

Divergences vanish by equation of motion and can be eliminated by field redefinition.

$$\cancel{R_{\mu\nu\rho\sigma}^2}$$

In  $D = 4$  topologically trivial space, Gauss-Bonnet theorem eliminates Riemann square term.

$$\int d^4x \sqrt{-g} (R^2 - 4R_{\mu\nu}^2 + R_{\mu\nu\rho\sigma}^2) = 32\pi^2 \chi \quad \text{Euler characteristic}$$

Pure gravity divergence with nontrivial topology:

$$\mathcal{L}^{\text{GB}} = -\frac{1}{(4\pi)^2} \frac{53}{90\epsilon} (R^2 - 4R_{\mu\nu}^2 + R_{\mu\nu\rho\sigma}^2)$$

Capper and Duff (1974)

Tsao (1977); Critchley (1978)

Gibbons, Hawking, Perry (1978)

Goroff and Sagnotti (1986)

Bornsen and van de Ven (2009)

- Euler characteristic vanishes in flat space. 't Hooft and Veltman (1974)
- Dimensional regularization makes it subtle. Capper and Kimber (1980)

This is an “enhanced cancellation”, but here it is well understood.

# The Trace Anomaly

Capper and Duff (1974); Tsao (1977); Critchley (1978); Gibbons, Hawking, Perry (1978);  
Duff and van Nieuwenhuizen (1980); Siegel (1980); Grisaru, Nielsen, Siegel, Zanon (1984);  
Goroff and Sagnotti (1986); Bornsen and van de Ven (2009); Etc.

**The Gauss-Bonnet divergence exactly corresponds to trace anomaly.**

$$D = 4 - 2\epsilon$$

$$\mathcal{L}^{\text{GB}} = -\frac{1}{(4\pi)^2} \frac{1}{360\epsilon} \left( \underset{\substack{\text{graviton} \\ \nearrow}}{4 \cdot 53} + \underset{\substack{\text{scalar} \\ \nearrow}}{1} + \underset{\substack{\text{2 form} \\ \uparrow}}{91} - \underset{\substack{\text{3 form} \\ \nwarrow}}{180} \right) (R^2 - 4R_{\mu\nu} + R_{\mu\nu\rho\sigma}^2)$$

**Gauss-Bonnet**  $\nwarrow$

$$T^\mu{}_\mu = -\frac{1}{(4\pi)^2} \frac{2}{360} \left( 4 \cdot 53 + 1 + 91 - 180 \right) (R^2 - 4R_{\mu\nu}^2 + R_{\mu\nu\rho\sigma}^2)$$

Duff and van Nieuwenhuizen (1980);

**Referred to as trace, conformal, trace or Weyl anomaly.**

# Quantum Inequivalence?

$$D = 4 - 2\epsilon$$

$$D \rightarrow 4$$

$$T^\mu{}_\mu = -\frac{1}{(4\pi)^2} \frac{2}{360} \left( \underset{\substack{\text{graviton} \\ \nearrow}}{4} \cdot \underset{\substack{\text{scalar} \\ \nearrow}}{53} + \underset{\substack{\text{2 form} \\ \nearrow}}{1} + \underset{\substack{\text{3 form} \\ \nwarrow}}{91} - \underset{\substack{\text{Gauss-Bonnet} \\ \nwarrow}}{180} \right) (R^2 - 4R_{\mu\nu}^2 + R_{\mu\nu\rho\sigma}^2)$$

two form dual to scalar

three form not dynamical

$$\partial_\mu \phi \leftrightarrow \varepsilon_{\mu\nu\rho\sigma} H^{\nu\rho\sigma}$$

$$\Lambda^{1/2} \leftrightarrow \varepsilon_{\mu\nu\rho\sigma} H^{\mu\nu\rho\sigma} \quad D = 4$$

**Classically equivalent. But is it quantum mechanically equivalent?**

- Quantum *inequivalence* under duality transformations.  
Duff and van Nieuwenhuizen (1980)
- Quantum equivalence under duality. Gauge artifact.  
Siegel (1980)
- Quantum equivalence of effective action (ignoring trace anomaly).  
Fradkin and Tseytlin (1984)
- Quantum equivalence of susy 1 loop effective action (with Siegel's argument for higher loops)  
Grisaru, Nielsen, Siegel, Zanon (1984)
- Quantum *inequivalence* and boundary modes.  
Finn Larsen and Pedro Lisbao (2015)

# Quantum Inequivalence?

**Need to ask the question in term of physically measurable quantities**

- **Scattering amplitudes good to look at. Cross sections physical.**
- **One loop too trivial. Need to look at two loops.**
- **We will see this question has a lot to do with UV properties.**

# Two-Loop Pure gravity

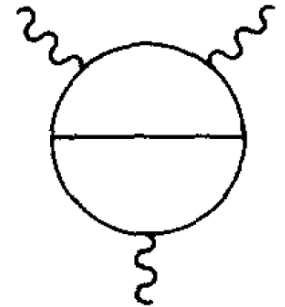
By two loops there is a valid  $R^3$  divergence.

Goroff and Sagnotti (1986); Van de Ven (1992)

Divergence in pure gravity:

$$\mathcal{L}^{R^3} = \frac{209}{2880} \frac{1}{(4\pi)^4} \frac{1}{2\epsilon} R^{\alpha\beta}{}_{\gamma\delta} R^{\gamma\delta}{}_{\rho\sigma} R^{\rho\sigma}{}_{\alpha\beta}$$

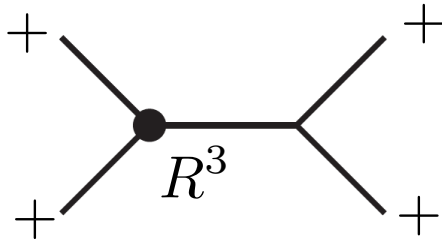
UV divergence



- Based on dimensional regularization.  $D = 4 - 2\epsilon$
- On surface nothing weird going on.
- The Goroff and Sagnotti result is correct in all details.

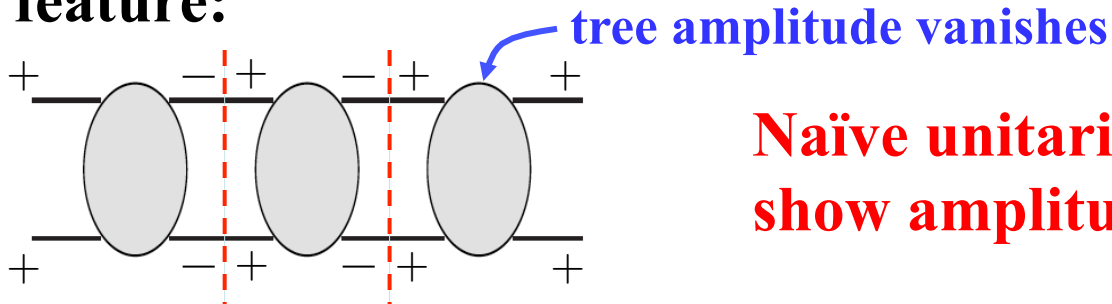
However, a goal of this talk is to show you that UV divergences in pure (super)gravity is subtle and weird, once you probe carefully.

# Two Loop Identical Helicity Amplitude



Pure gravity identical helicity amplitude sensitive to Goroff and Sagnotti divergence.

Curious feature:



Naïve unitarity arguments show amplitude vanishes!

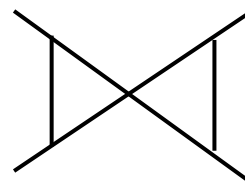
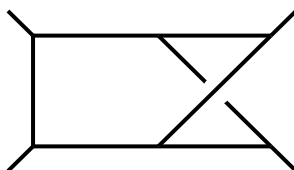
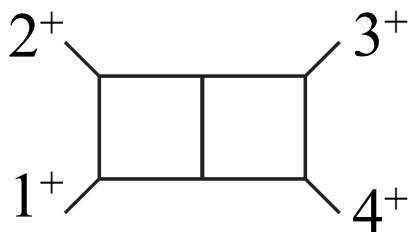
- Gravity amplitude proportional to  $0/0$ , resolved in dim reg.
- Characteristic of quantum anomalies.
- Bardeen and Cangemi pointed out nonvanishing of identical helicity is connected to an anomaly in self-dual sector.

A surprise:

Divergence is *not* generic but tied to anomaly-like behavior.

# Two-Loop Divergence

ZB, Cheung, Chi, Davies, Dixon and Nohle (to appear)



$$209 = 11 \cdot 19$$

$$3431 = 47 \cdot 73$$

Integrating we obtain:

$$\mathcal{M}_4^{2\text{-loop}} \Big|_{\text{bare div.}} = -\frac{1}{\epsilon} \frac{3431}{5400} \mathcal{K} \quad \mathcal{K} = \left( \frac{\kappa}{2} \right)^6 \frac{i}{(4\pi)^4} stu \left( \frac{[12][34]}{\langle 12 \rangle \langle 34 \rangle} \right)^2$$

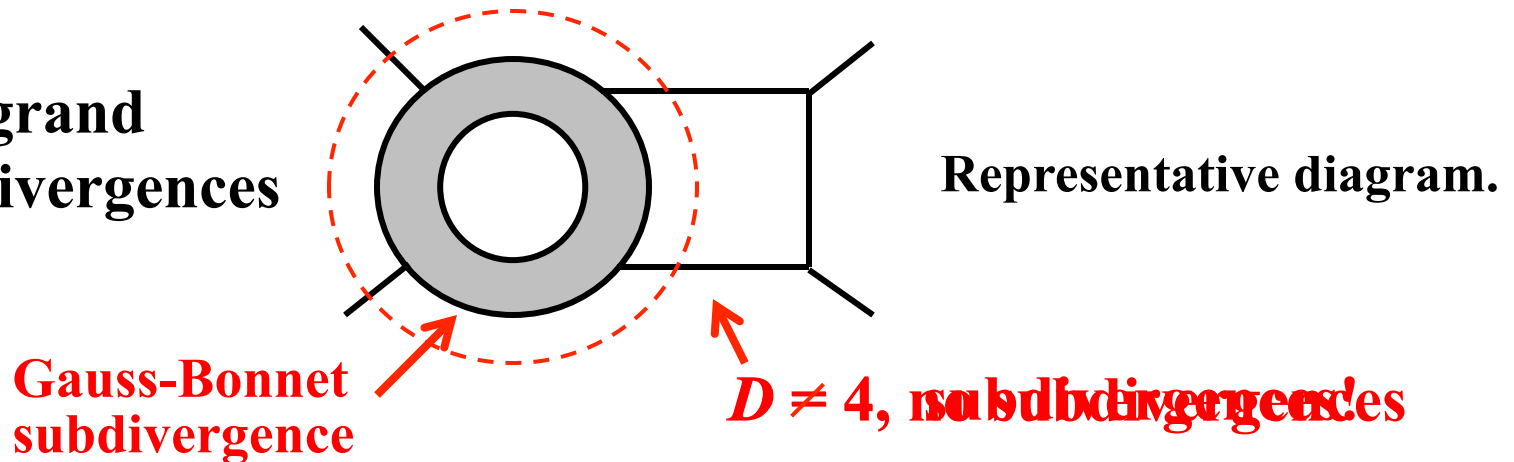
Not the same as the Goroff and Sagnotti result

However, Goroff and Sagnotti subtracted subdivergences integral by integral, following standard procedures.

**Subdivergences? What subdivergences?**  
**There are no one-loop divergences. Right?**

# Subdivergences?

The integrand  
has subdivergences

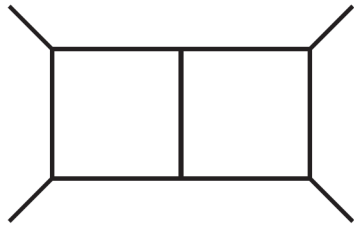


A strange phenomenon: no one loop divergences,  
yet there are one-loop subdivergences to subtract!

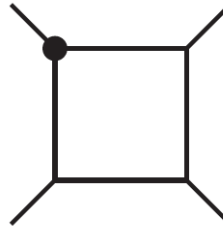
- To match the G&S result we need to subtract subdivergences.
- Using modern methods we can track the pieces.

# Two Loop Identical Helicity Divergence

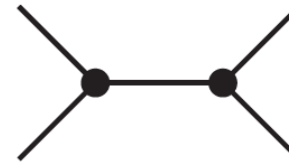
ZB, Cheung, Chi, Davies, Dixon and Nohle



**2 loop bare**



**single GB subtraction**



**double GB subtraction**

$$\begin{aligned}\mathcal{M}_4^{2\text{-loop}} \Big|_{\text{div.}} &= -\frac{1}{\epsilon} \frac{3431}{5400} \mathcal{K} \\ \mathcal{M}_4^{1\text{-loop GB}} \Big|_{\text{div.}} &= \frac{1}{\epsilon} \frac{689}{675} \mathcal{K} \\ \mathcal{M}_4^{\text{tree GB}^2} \Big|_{\text{div.}} &= \frac{1}{\epsilon} \frac{5618}{675} \mathcal{K}\end{aligned}$$

$$\mathcal{M}_4^{\text{total}} \Big|_{\text{div.}} = \frac{1}{\epsilon} \frac{209}{24} \mathcal{K}$$

**Goroff and Sagnotti  
divergence reproduced**

**Trace anomaly plays central role in divergence!**

# Meaning of Divergence?

**What does the divergence mean?**

$$\Lambda^{1/2} \leftrightarrow \varepsilon_{\mu\nu\rho\sigma} H^{\mu\nu\rho\sigma}$$

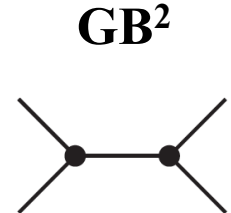
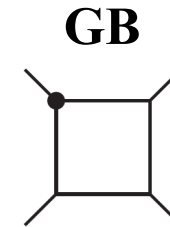
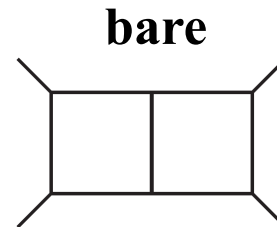
**Adding  $n_3$  3-form field offers good way to understand this:**

- On the one hand, no degrees of freedom in  $D = 4$ , so no change in divergence expected.
- On the other hand, the trace anomaly is affected, so expect change in divergence.
- Note that 3 form proposed as way to dynamically neutralize cosmological constant.

Brown and Teitelboim; Bousso and Polchinski

$1/\epsilon$

bare	$-\frac{3431}{5400} - \frac{199n_3}{30} + 6n_3^2$
GB	$\frac{4 \cdot 53 - 180n_3}{360} \cdot \frac{2 \cdot (13 + 180n_3)}{15}$
GB <sup>2</sup>	$24 \left( \frac{4 \cdot 53 - 180n_3}{360} \right)^2$
total	$\frac{209}{24} - \frac{15}{2}n_3$



**Divergence depends on nondynamical 3-form fields. Quantum inequivalence?**

**But wait: what about finite parts? Need physical quantity!**

# Scattering Amplitudes

## Pure Gravity:

$$\mathcal{M}_G^{(2)}(1^+, 2^+, 3^+, 4^+) = \mathcal{N} \left[ \frac{1}{\epsilon} \frac{209}{24} stu + \frac{117617}{21600} stu \right. \\ \left. + \left( \frac{1}{10} stu - \frac{1}{60} s^3 \right) \log \left( \frac{-s}{\mu^2} \right) + \frac{1}{120} (s^2 + t^2 + u^2) s \log^2 \left( \frac{-s}{\mu^2} \right) + \text{perms} \right]$$

IR singularities  
subtracted and  
independent of 3 form

## Gravity + 3 Form:

$$\mathcal{M}_{G3}^{(2)}(1^+, 2^+, 3^+, 4^+) = \mathcal{N} \left[ \frac{1}{\epsilon} \frac{29}{24} stu + \frac{411617}{21600} stu \right. \\ \left. + \left( \frac{1}{10} stu - \frac{1}{60} s^3 \right) \log \left( \frac{-s}{\mu^2} \right) + \frac{1}{120} (s^2 + t^2 + u^2) s \log^2 \left( \frac{-s}{\mu^2} \right) + \text{perms} \right]$$

divergences differ.  
logarithms identical!

- Value of divergence not physical. Renormalize away.
- 3 form is a Cheshire Cat field: physical scattering unaffected.



- Results consistent with quantum equivalence under duality.
- For carefully defined physically measurable quantities it seems that duality transformations should not alter the physics.

# $N = 1$ Supergravity

ZB, Chi, Dixon, Edison (to appear)

**Divergence violates susy ward identity even though regulator should be supersymmetric! Due to trace anomaly.**

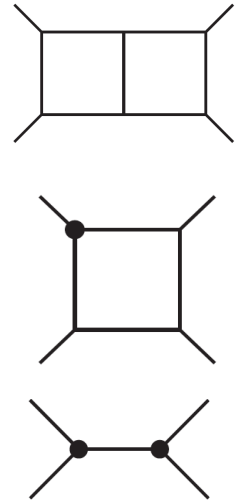
**Result for  $N = 1$  supergravity with 1 matter multiplet**

$$\mathcal{M}_4 \Big|_{\text{div}} = \frac{1}{\epsilon} \frac{81871}{21600} \mathcal{K} + 0 \ln(\mu^2) \mathcal{K}$$

**Very strange, but no stranger than earlier results.**

**Have no fear: no physical effect! Local counterterm eats the divergence restoring susy.**

**Still working on case with no matter multiple, but no reason to expect different outcome.**



# Simple Two-Loop Formula

ZB, Cheung, Chi, Davies, Dixon and Nohle; ZB, Chi, Dixon, Edison (to appear)

**Focus on renormalization scale dependence *not* divergences!**  
**In QCD these are effectively the same. In gravity *not* related!**

Looking at various theories, we wind up with a simple 2 loop formula:

$$\mathcal{M}_4^{(2)} \Big|_{\ln \mu^2} = -\mathcal{K} \frac{N_b - N_f}{8} \ln \mu^2$$

$N_b$  is number of bosonic states.  
 $N_f$  is number of fermionic states.

- Confident this is robust and does not depend on dimensional regularization or details of theory.
- Vanishes at two loops in susy theory, as expected.
- Unless  $\ln \mu^2$  dependence vanishes, theory should still be considered nonrenormalizable.

**Anomaly-like structure leads to a remarkably simple formula for UV properties in any minimally coupled gravity theory! Who ordered this??**

# Summary

1. Gravity integrands from gauge theory. Very powerful tool.
2. Standard view of gravity UV too naive:
  - New phenomenon: “Enhanced” UV cancellations in gravity.
  - So far divergences of pure (super)gravity theories appear to be due to anomalous behavior! No anomaly no divergence.
  - Renormalized scattering amplitudes independent of duality transformations.
3. Focus on renormalization scale dependence rather than divergences. Equivalent in gauge theory, but not in gravity.
4. Simple two-loop formula for renormalization scale behavior in any gravity theory. Who ordered that?

**Expect many more surprises as we probe gravity theories using modern perturbative tools.**

# Extra slides

# Divergences and Duality

ZB, Cheung, Chi, Davies, Dixon and Nohle

$$\mathcal{M}_4 = \left[ \frac{1}{\epsilon} \left( \frac{209}{24} - \frac{15}{2} n_3 \right) - \frac{1}{4} \ln \mu^2 \right] \mathcal{K} + \text{finite}$$

**divergence** (points to  $\frac{1}{\epsilon}$ )      **number of 3 forms** (points to  $n_3$ )      **independent of 3 forms** (points to  $\frac{1}{4}$ )      **Renormalization scale** (points to  $\ln \mu^2$ )

$$\mathcal{K} = \left( \frac{\kappa}{2} \right)^6 \frac{i}{(4\pi)^4} stu \left( \frac{[12][34]}{\langle 12 \rangle \langle 34 \rangle} \right)^2$$

- **Weird that renorm. scale and UV divergence not linked!**  
**Happens because of Gauss-Bonnet subdivergence.**
- **The renormalization scale dependence is robust and almost certainly not an artifact of dimensional regularization.**
- **For carefully defined physically measurable quantities it appears that duality transformations should not alter the physics.**

# Gravity From Gauge Theory

	$n$	$\tilde{n}$
$N = 8$ sugra:	$(N = 4 \text{ sYM})$	$\times (N = 4 \text{ sYM})$
$N = 5$ sugra:	$(N = 4 \text{ sYM})$	$\times (N = 1 \text{ sYM})$
$N = 4$ sugra:	$(N = 4 \text{ sYM})$	$\times (N = 0 \text{ sYM})$

**Spectrum controlled by simple tensor product of YM theories.  
Recent papers show more sophisticated lower-susy cases.**

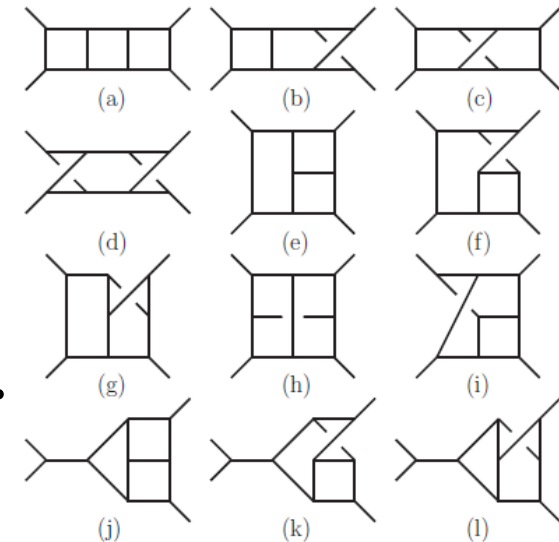
Anastasiou, Bornsten, Duff; Duff, Hughs, Nagy; Johansson and Ochirov;  
Carrasco, Chiodaroli, Günaydin and Roiban; ZB, Davies, Dennen, Huang and Nohle;  
Nohle; Chiodaroli, Günaydin, Johansson, Roiban.

# Enhanced UV Cancellations

ZB, Davies, Dennen (2014)

Suppose diagrams in *all* possible Lorentz covariant representations are UV divergent, but the amplitude is well behaved.

- By definition this is an enhanced cancellation.
- Not the way gauge theory works.



**$N = 4$  sugra**

$n_i \sim s^3 t A_4^{\text{tree}} (p \cdot q)^2 \varepsilon_1 \cdot p \varepsilon_2 \cdot p \varepsilon_3 \cdot q \varepsilon_4 \cdot q + \dots$

**$N = 4$  sugra: pure YM  $\times$   $N = 4$  sYM**

**This diagram is log divergent**

*already log divergent*

3 loop UV finiteness of  $N = 4$  supergravity proves existence of “enhanced cancellation” in supergravity theories.

# Some New Directions in Gravity Loops

If you want to solve a difficult problem get an army of energetic young people to help with new ideas:

- **Better understanding and applications of BCJ duality.**

Chiodaroli, Gunaydin, Johansson and Roiban,; Johansson, Ochirov; O'Connell, Montiero, White; ZB, Davies, Nohle; Boels, Isermann, Monteiro, and O'Connell; Mogull and O'Connell, He, Monteiro, and Schlotterer

- **Scattering equations and double-copy relations.**

Cachazo, He, Yuan

- **Twistor strings now at loop level for  $N = 8$  supergravity.**

Adamo, Casali and Skinner; Geyer, Mason, Monteiro and Tourkine

- **New ideas on unitarity cuts based on Feynman Tree Theorem**

Baadsgaard, Bjerrum-Bohr, Bourjaily, Caron-Huot, Damgaard and Feng

- **Important advances in related string theory amplitudes.**

Carlos Mafra and Oliver Schlotterer

- **Nonplanar analytic hints from Amplituhedron.**

ZB, Hermann, Litsey, Stankowicz, Trnka

- **Awesome equation solver. Millions of equations encountered at 5 loops can be dealt with! Very cool algorithm!**

Schabinger and von Manteuffel