Finite temperature gluon spectral functions from twisted mass lattice QCD

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Physical picture of QCD phases above and below the crossover

Below T_c : Confinement and chiral symmetry breaking
Modelled by Hadron Resonance Gas (Remarkably: with masses taken from $T = 0$! Apparently no other degrees of freedom ?)
Above T_c : (gradual) Deconfinement and chiral symmetry restoration
Modelled by colored degrees of freedom with strong interaction.
(Apparently there are - in addition - remnants of mesonic objects, not-yet melted charmonia, glueballs ?)
Kinetic description : gluon- and quark-like quasi particles (one needs their spectral functions !)
Calculation of transport coefficients possible in terms of quasi particles.
For that the knowledge of the spectral function of the corresponding quasi particles is required.
Can be parametrized by in-medium dispersion relation :
T-dependent mass and a T-dependent width

Spectral function of gluons (and quarks)

This work tries to extract the in-medium gluon spectral function from Euclidean gluon correlation data.

Below T_c, non-positivity of the gluon spectral function demonstrates, that the gluon is not a "particle as usual".

Violation of spectral positivity is an important feature (over and over observed in studies of the gluon propagator) of confinement.

Non-positivity seen also in the Laplace transform ($p_4 \rightarrow$ Euclidean time) of the $T = 0$ gluon propagator, $G(\tau, \vec{p})$.

Non-positivity thoroughly discussed in:

- R. Alkofer and L. von Smekal, Phys. Rept. **353** (2001) 281 [hep-ph/0007355]
- J. M. Cornwall, Mod. Phys. Lett. A **28** (2013) 1330035 [arXiv:1310.7897 (hep-ph)]

In general, spectral function are obtained by analytic continuation of Euclidean correlation functions.

The ill-posedness

Notoriously, an ill-posed problem :

- finite number (actually, only a very small number) of data points $N_{q_4} = N_\tau$ at finite T, unless one uses highly anisotropic lattices
- wanted: a continuous spectral function $\rho(\omega)$
- usually the data is very noisy !
- What in other cases is helpful ? For physical (bound state) particles (light, heavy-light mesons, charmonia) the spectral function is positive semidefinite (giving a number of distinct gauge-invariant states per mass interval).
- Gluons, in contrast, are unphysical particles : violate spectral positivity, *i.e.* $\rho(\omega)$ may irregularly assume positive and negative values.

This complicates our task.
Superconvergent sum rule $\int_0^\infty \rho_T(m^2)dm^2 = 0$ (Reinhard Oehme)

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Gauge potential and propagator from lattice links

$$A_\mu(x+\hat{\mu}/2) = \frac{1}{2iag_0}(U_{x\mu} - U_{x\mu}^\dagger)\,|_{traceless}$$

Fourier transform of the gauge potential on lattice

$$\tilde{A}_\mu^a(q)$$

Fourier transformed gluon propagator : correlator of two Fourier transformed gauge potentials

$$D_{\mu\nu}^{ab}(q) = \left\langle \widetilde{A}_\mu^a(q)\widetilde{A}_\nu^b(-q)\right\rangle .$$

Discrete lattice momenta (finite number)

$$k_\mu a = \frac{\pi n_\mu}{N_\mu}, \quad n_\mu \in (-N_\mu/2, N_\mu/2]\,.$$

Lattice momenta

Lattice momenta are related to physical momenta via

$$q_\mu(n_\mu) = \frac{2}{a} \sin\left(\frac{\pi n_\mu}{N_\mu}\right) .$$

For finite temperature in isotropic lattices, relatively bigger volumes are required to get realistic transition parameters,

$$N_i >> N_\tau$$

(chance to approach the thermodynamic limit).

A reasonable aspect ratio is required, $V^{1/3} = L >> 1/T$

Gluon propagator at zero temperature

At zero temperature $O(4)$ symmetry expected :

Usually momentum cuts (cylinder cut, cone cut etc.) are applied with the aim to minimize hypercubic artefacts (to improve the continuum extrapolation).
A systematic procedure exists to reduce hypercubic artefacts:
"Using NSPT for the Removal of Hypercubic Lattice Artifacts",
Jakob Simeth, Andre Sternbeck, Meinulf GÃ¶ckeler, Holger Perlt, Arwed Schiller,
PoS LATTICE2014 (2015) 294, arXiv:1501.06322
"Discretization Errors for the Gluon and Ghost Propagators in Landau Gauge using NSPT",
Jakob Simeth, Andre Sternbeck, Ernst-Michael Ilgenfritz, Holger Perlt, Arwed Schiller, PoS LATTICE2013 (2014) 459, arXiv:1311.1934

Transversal and longitudinal projectors

At non-zero temperature : no (approx.) $O(4)$ rotational symmetry anymore !

Define transversal and longitudinal polarization tensors

$$P^T_{\mu\nu} = (1-\delta_{\mu 4})(1-\delta_{\nu 4})\left(\delta_{\mu\nu} - \frac{q_\mu q_\nu}{\vec{q}^{\,2}}\right)$$

$$P^L_{\mu\nu} = \left(\delta_{\mu\nu} - \frac{q_\mu q_\nu}{q^2}\right) - P^T_{\mu\nu} \, .$$

This tensor structure defines two propagators :
D_L (longitudinal, electric) and D_T (transversal, magnetic)

$$D^{ab}_{\mu\nu}(q) = \delta^{ab}\left(P^T_{\mu\nu} D_T(q_4^2, \vec{q}^{\,2}) + P^L_{\mu\nu} D_L(q_4^2, \vec{q}^{\,2})\right) \, .$$

Transversal and longitudinal propagators

The explicit expressions for the propagators $D_{T,L}$ read (if $q_4 \neq 0$)

$$D_T(q) = \frac{1}{2N_g} \left\langle \sum_{i=1}^{3} \widetilde{A}_i^a(q)\widetilde{A}_i^a(-q) - \frac{q_4^2}{\vec{q}^{\,2}} \widetilde{A}_4^a(q)\widetilde{A}_4^a(-q) \right\rangle$$

and

$$D_L(q) = \frac{1}{N_g} \left(1 + \frac{q_4^2}{\vec{q}^{\,2}}\right) \left\langle \widetilde{A}_4^a(q)\widetilde{A}_4^a(-q) \right\rangle$$

In the past, propagators were mostly studied as function of spatial $|\vec{q}|$ (and, moreover, restricted to $q_4 = 0$).
However, zero Matsubara frequency data is not sufficient for the present task of extracting the spectral function.

Källan-Lehmann representation

We may relate the gluon correlators for imaginary frequencies q_4 to their spectral function via the Källen-Lehmann representation (at any temperature, $T = 0$ and $T \neq 0$, for each momentum $\mathbf{q}$)

$$D_{T,L}(q_4, \mathbf{q}) = \int_{-\infty}^{\infty} \frac{1}{iq_4 + \omega}\, \rho_{T,L}(\omega, \mathbf{q})\, d\omega$$
$$= \int_{0}^{\infty} \frac{2\omega}{q_4^2 + \omega^2}\, \rho_{T,L}(\omega, \mathbf{q})\, d\omega\ ,$$

with the spectral function being antisymmetric around the origin of real-time frequencies, $\rho(-\omega) = -\rho(\omega)$.
Inverting this relation using the simulated correlator data represents the spectral function depending on the temperature.
Obviously, knowledge of the q_4 dependence becomes now crucial !

Reconstruction method for the spectral function

Our method is a Bayesian Reconstruction method.

Other methods to find the spectral function being used:

- Maximal entropy method (prevented by non-positivity)
- Tikhonov regularization (used by the Coimbra-Leuven group: P. Silva, O. Oliveira, D. Dudal)

A method directly relating data given in the Euclidean time domain to the corresponding spectral function is the

- Gilbert-Backus method, being used
  - by the Mainz group (H. Meyer et al.)
  - by the ITEP group (V. Braguta, A. Kotov, N. Astrakhantsev) and by M. Ulybyshev, Regensburg

Reconstruction from data in the Euclidean time domain

This is in contrast to correlators obtained in Euclidean time domain, *e.g.* for the calculation of viscosity (or electric conductivity),

$$C_{TT}(x_0) = T^{-5} \int d^3\mathbf{x} \, \langle T_{12}(0) T_{12}(x_0, \mathbf{x}) \rangle \, ,$$

where the correlation function can be written in terms of the corresponding spectral function $\rho_{TT}(\omega)$ as follows :

$$C_{TT}(x_0) = T^{-5} \int_0^{\infty} \rho_{TT}(\omega) \, \frac{\cosh \omega(\frac{1}{2T} - x_0)}{\sinh \frac{\omega}{2T}} \, d\omega \, . \tag{1}$$

In our case, the kernel has no explicit temperature dependence !

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Bayesian spectral reconstruction

In quasi-continuum version: we have to reproduce a finite and noisy set of data points by D_i^ρ, an ω-integral over ρ (handled as Riemann sum).
Frequency bins : $N_\omega = O(1000)$ bins.
Lattice data points : only $N_{q_4} \in [4 \ldots 20]$ values (the number depending on temperature for a fixed scale, *i.e.* in a fixed-β setting).

Hence, our task is inverting a bin-discretized Källen-Lehmann relation

$$D_i^\rho = \sum_{l=1}^{N_\omega} K_{il}\, \rho_l\, \Delta\omega_l, \quad i \in [0, N_{q_4}], \quad N_\omega \gg N_{q_4}$$

Using a naive χ^2 fit for the (binwise constant-valued) ρ_l values would yield an infinite number of degenerate solutions.

Bayesian spectral reconstruction

It starts by writing the probability of a test spectral function ρ to be the correct spectral function, given the measured data (D_i) and given further, so called prior information (I). This probability is proportional to the product of two terms

$$P[\rho|D, I] \propto P[D|\rho, I]\; P[\rho|I].$$

This expression follows from the multiplication theorem of conditional probabilities and formally allows the prior information I (in other words, a default model) to influence both factors on the right hand side.

The functional to maximize

The first factor of it,

$$P[D|\rho, I] = \exp[-L]$$

refers to the likelihood probability, where the likelihood L measures the χ^2 distance between the correlator points D_i^ρ (as obtained from the test function ρ) and the actually simulated data D_i

$$L = \frac{1}{2} \sum_{i,j=1}^{N_{q_4}} (D_i - D_i^\rho) C_{ij}^{-1} (D_j - D_j^\rho),$$

where C_{ij} is the usual covariance matrix of the simulated D_i's.
Prior information (I) enters here only implicitely.
The L functional (as functional of ρ_i) possesses $N_\omega - N_{q_4}$ flat directions.
In any Bayesian approach this must be regularised by a prior probability, which is specified in terms of an "entropy" functional.

The standard Bayesian Reconstruction method

The *a priori* probability of ρ is $P[\rho|I] = \exp[-\alpha(S(\omega)]$ specified by some “entropy functional”.

- Maximal entropy method (MEM)
  Here, the Shannon-Jaynes relative entropy plays this role. It is applicable in case of replacing the default model $m(\omega)$ by some freely chosen $\rho(\omega)$

$$S_{\mathrm{SJ}} = \int d\omega\left(\rho(\omega) - m(\omega) - \rho(\omega)\log\left[\frac{\rho(\omega)}{m(\omega)}\right]\right)$$

The prior knowledge enters through the parametrization given by the default spectral density $m(\omega)$ (in binned form).
The coefficient α (multiplying the relative entropy) expresses the importance given to the prior information.
For $\alpha \to \infty$, the most probable $\rho(\omega)$ turns out to be $\rho(\omega) = m(\omega)$, independently of any data.

The improved regulator

Non-positivity is a problem. The Shannon-Jaynes entropy above might be used even if there are regions of positive and negative $\rho(\omega)$, but in our case these regions are not known *a priori* !

- Standard Bayesian method (BR)
  Here, the Shannon-Jaynes relative entropy is replaced by a regulating functional for which - in the absence of simulation data - the most probable $\rho(\omega)$ is again $\rho(\omega) = m(\omega)$.

$$S_{\mathrm{BR}} = \int d\omega \left(1 - \frac{\rho(\omega)}{m(\omega)} + \log\left[\frac{\rho(\omega)}{m(\omega)}\right]\right)$$

  Only the ratio $\rho(\omega)/m(\omega)$ matters here !
  "Bayesian Approach to Spectral Function Reconstruction for Euclidean Quantum Field Theories",
  Yannis Burnier and Alexander Rothkopf,
  Phys. Rev. Lett. 111 (2013) 18200331, arXiv:1307.6106

The improved regulator

- Novel Bayesian method, which accounts for the non-positivity of $\rho(\omega)$; here the generalized entropy functional

$$S^g_{\rm BR} = \int d\omega \Big(-\frac{|\rho(\omega) - m(\omega)|}{h(\omega)} + \log\Big[\frac{|\rho(\omega) - m(\omega)|}{h(\omega)} + 1\Big]\Big).$$

  takes over the role of regulator and relative entropy.
  $r_l = \frac{|\rho_l - m_l|}{h_l}$ is the distance of ρ from the default model m, to be taken relative to h, which encodes the confidence in the default model.
  "Bayesian inference of nonpositive spectral functions in quantum field theory",
  Alexander Rothkopf,
  Phys. Rev. D95 (2017) 056016, arXiv:1611.00482

Getting rid of the factor α

This analytic form of S^g_{BR} allows one to integrate out α in a straight forward fashion, allowing full ignorance about the values of α (putting the corresponding distribution $W[\alpha] = const$):

$$P[\rho|D, I, m] \propto P[D|\rho, I] \int_0^\infty d\alpha \; P[\rho|m, \alpha] \; W[\alpha]$$

Variational problem

Once $m(\omega)$ and $h(\omega)$ are specified, say as $m(\omega) = 0$ and $h(\omega) = 1$, we have to carry out a numerical search for the most probable Bayesian spectrum according to

$$\left.\frac{\delta P[\rho|D, I]}{\delta \rho}\right|_{\rho=\rho^{\text{Bayes}}} = 0,$$

Alternative choices $m(\omega) = \pm 1$ and lowering the confidence in the default model by changing $h(\omega) = 1 \rightarrow h(\omega) = 2$ allow to check the influence of the prior.

A measure of reliability

Then the integrated curvature of the minimization functional with respect to $\rho(\omega)$

$$Q[\rho, m, h] = \mathrm{Log}[P[D|\rho, I] \int_0^\infty d\alpha P[\rho|m, \alpha]]$$

may serve as a measure of reliability (depending on ω).

Outline

Configurations taken from the tmfT collaboration

Aim of the tmfT collaboration : improve the lattice thermodynamics of Wilson fermions by using the twisted-mass improvement.

Main results obtained for $N_f = 2$

- Localization and characterization of the crossover for various light quark masses (or "pion" mass values)
- Equation of State (EoS) for two light flavors
- Unquenching effect on the gluon propagator

Main results obtained for $N_f = 2+1+1$

- Localization and characterization of the crossover for various light quark masses (or "pion" mass values) in presence of s and c quarks with realistic mass
- Equation of State (EoS) including two light flavors and additional s and c quarks with realistic mass (still incomplete)
- T dependence of the topological susceptibility with four flavors

topological susceptibilty

“Topological susceptibility from Nf=2+1+1 lattice QCD at nonzero temperature”,
Anton Trunin, Florian Burger, Ernst-Michael Ilgenfritz, Maria Paola Lombardo, Michael Müller-Preussker,
J.Phys.Conf.Ser. 668 (2016) no.1, 012123, arXiv:1510.02265
(Strangeness in QuarkMatter 2015)

“Topology (and axion’s properties) from lattice QCD with a dynamical charm”,
Florian Burger, Ernst-Michael Ilgenfritz, Maria Paola Lombardo, Michael Müller-Preussker, Anton Trunin,
arXiv:1705.01847 (Quark Matter 2017)

The fermionic action is improved compared with unimproved Wilson fermions coming in four flavours

Fermions are grouped in two twisted doublets.

- The light doublet action (degenerate *u* and *d* quarks) with mass tuned by the twisted mass parameter μ_l

$\kappa_l = \kappa_c(\beta)$ (*i.e.* “maximal twist”)

$$S_f^l[U, \chi_l, \overline{\chi}_l] = \sum_{x,y} \overline{\chi}_l(x)[\delta_{x,y} - \kappa_l D_{\mathrm{W}}(x,y)[U] + 2i\kappa_l a\mu_l \gamma_5 \delta_{x,y} \tau_3]\chi_l(y)$$

The fermionic action is improved compared with unimproved Wilson fermions coming in four flavours

- The heavy doublet action (non-degenerate s and c quarks) with masses tuned by two twisted mass parameters μ_σ and μ_δ

  whereas again $\kappa_h = \kappa_c(\beta)$ (*i.e.* "maximal twist")

$$S_f^h[U, \chi_h, \overline{\chi}_h] = \sum_{x,y} \overline{\chi}_h(x)[\delta_{x,y} - \kappa_h D_W(x,y)[U] + 2i\kappa_h a\mu_\sigma \gamma_5 \delta_{x,y} \tau_1 + 2\kappa_h a\mu_\delta \delta_{x,y} \tau_3] \chi_h(y)$$

In both actions, the τ_i are the Pauli matrices in the respective doublet (*i.e.* flavor) space.

The fermionic action is a Wilson-type action

The term $D_W[U]$ denotes the standard gradient term for Wilson fermions

$$D_W[U] = \frac{1}{2a}[\gamma_\mu(\nabla_\mu + \nabla^*_\mu) - \nabla^*_\mu \nabla_\mu]$$

Simulation algorithm :

“Numerical simulation of QCD with u, d, s and c quarks in the twisted-mass Wilson formulation”,
T. Chiarappa, F. Farchioni, K. Jansen, I. Montvay, E. E. Scholz, L. Scorzato, T. Sudmann, and C. Urbach,
Eur. Phys. J. C50 (2007) 373, arXiv:hep-lat/0606011

The gluonic action

Gauge action: improved Iwasaki action is used

$$S_g[U] = \beta\Big(c_0 \sum_P [1 - \frac{1}{3}\mathrm{Tr}\,(U_P)] + c_1 \sum_R [1 - \frac{1}{3}\mathrm{Tr}\,(U_R)]\Big)\,.$$

with ($c_0 = 3.648$ and $c_1 = -0.331$), where the sum (P) contains all plaquettes and the sum (R) all planar rectangles.

Lattice setting

Configurations taken from simulations of the "twisted mass at finite temperature" (tmfT) collaboration (M. Müller-Preussker et al.).
For this first spectral paper, only the ensembles for $M_\pi \approx 370$ MeV have been analysed for all three lattice spacings (will be extended ...).

ETMC ens. ($T = 0$)	A60.24	B55.32	D45.32
tmfT ens. ($T \neq 0$)	A370	B370	D370
β	1.90	1.95	2.10
a [fm]	0.0936	0.0823	0.0646
m_π [MeV]	364(15)	372(17)	369(15)
$T_{\rm deconf}$ [MeV]	202(3)	201(6)	193(13)
$N_\tau = N_{q_4}$ in range	4-14	10-14	4-20

Table: Properties of the three sets of finite-temperature ensembles used in our study, among them the deconfinement crossover temperature $T_{\rm deconf}$ (defined by the peak of Polyakov loop susceptibility).

Grid sizes for $D370$, *i.e.* $\beta = 2.10$ and $M_\pi \approx 370$ MeV

$D370$										
N_τ	4	6	8	10	11	12	14	16	18	20
T [MeV]	762	508	381	305	277	254	218	191	170	152
N_s	32	32	32	32	32	32	32	32	40	48
N_{meas}	310	400	120	410	420	380	790	610	590	280

Table: Grid sizes and temperatures of the $D370$ ensembles used for the computation of the correlation functions in this work. N_{meas} refers to the number of available correlator measurements (uncorrelated configurations).

Gauge condition : Landau gauge

One essential detail :

Propagators require gauge fixing: we specify the Landau gauge.
This corresponds to the following discretized local condition

$$\nabla_\mu A_\mu = \sum_{\mu=1}^{4} (A_\mu(x+\hat{\mu}/2) - A_\mu(x-\hat{\mu}/2)) = 0 \;,$$

to be imposed on the gauge fields defined in terms of link variables.
This can be achieved by using the freedom of performing suitable gauge transformations acting on the links.

Iterative gauge fixing

This condition may be fulfilled by iteratively applying local gauge transformations g_x

$$U_{x\mu} \overset{g}{\mapsto} U^g_{x\mu} = g^\dagger_x U_{x\mu} g_{x+\mu} \,, \qquad g_x \in SU(3) \,,$$

in order to maximize the gauge functional

$$F_U[g] = \frac{1}{3} \sum_{x,\mu} \mathrm{Tr} \left(g^\dagger_x U_{x\mu} g_{x+\mu} \right) .$$

We apply the convergence criterium

$$\max_x \mathrm{Tr} \left[\nabla_\mu A_{x\mu} \nabla_\nu A^\dagger_{x\nu} \right] < 10^{-13} .$$

This procedure has been carried out by means of the CUDA LGT library (Schröck 2012), adapted by A. Trunin for the use for our lattice configurations.

Outline

Longitudinal gluon correlation functions for $\beta = 2.1$ at zero Matsubara frequency

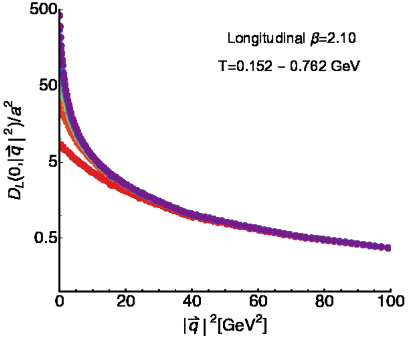


Figure: The longitudinal gluon correlators at $\beta = 2.10$ evaluated for different temperatures $T = 152 \ldots 762$ MeV at **vanishing** imaginary frequency $q_4 = 0$ for finite spatial momenta $|\vec{q}|^2$. The right panel is zoomed in towards the origin.

Longitudinal gluon correlation functions for $\beta = 2.1$ including nonzero Matsubara frequencies

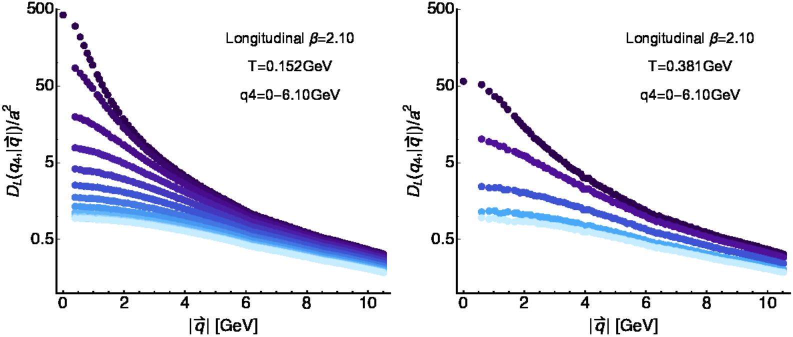


Figure: The longitudinal gluon propagators at $\beta = 2.10$ evaluated for two temperatures (left $T = 152$ MeV, right $T = 381$ MeV) showing the $|\vec{q}|$ dependence at various fixed q_4 values. Darkest colors are assigned to the lowest value of the corresponding parameter q_4, *i.e.* Matsubara $q_4 = 0$.

Longitudinal gluon correlation functions for $\beta = 2.1$ including nonzero Matsubara frequencies

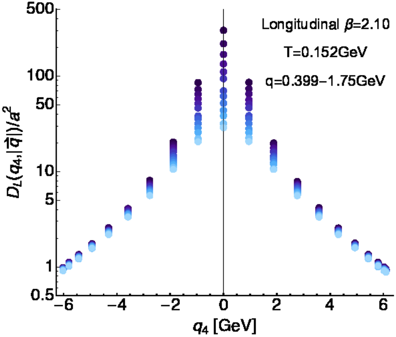
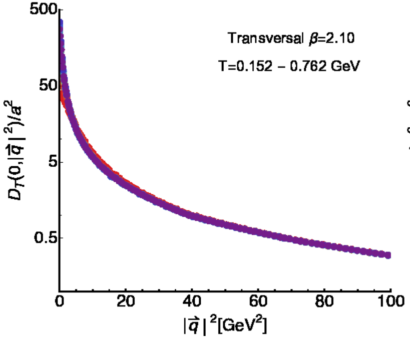


Figure: The longitudinal gluon propagators at $\beta = 2.10$ evaluated for two temperatures (left $T = 152$ MeV, right $T = 381$ MeV) showing the q_4 dependence for fourteen lowest $|\vec{q}|$ momentum values. Darkest colors are assigned to the lowest value of the corresponding 3-momentum $|\vec{q}|$.

Outline

Transversal gluon correlation functions for $\beta = 2.1$ at zero Matsubara frequency



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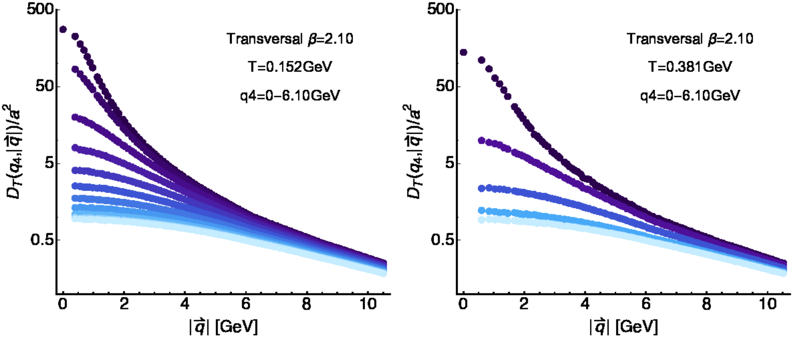
Transversal gluon correlation functions for $\beta = 2.1$ including nonzero Matsubara frequencies



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Transversal gluon correlation functions for $\beta = 2.1$ including nonzero Matsubara frequencies

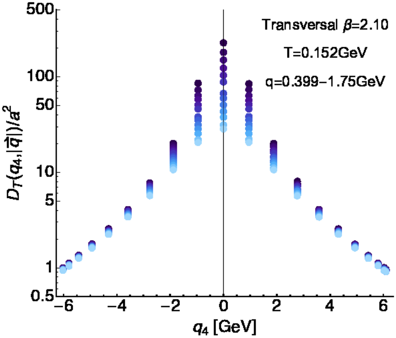


Figure: The transversal gluon propagators at $\beta = 2.10$ evaluated for two temperatures (left $T = 152$ MeV, right $T = 381$ MeV)

Outline

Longitudinal spectral function in confinement and deconfinement

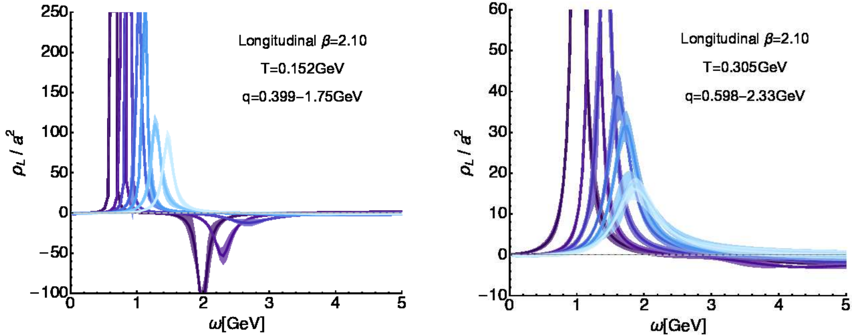
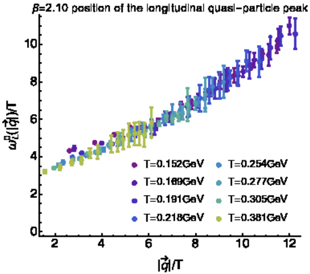


Figure: Reconstructed longitudinal gluon spectra on the $\beta = 2.10$ ensembles for (left) $T = 152$ MeV and (right) $T = 305$ MeV. The different curves refer to seven out of fourteen lowest spatial momenta at which correlator data is available. The y-axis is shifted to allow to see strong negative contributions (appearing in confinement). Error bands arose from varying the default.

Longitudinal quasi-particle peak position as function of momentum



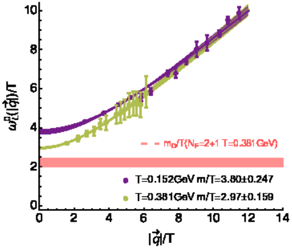


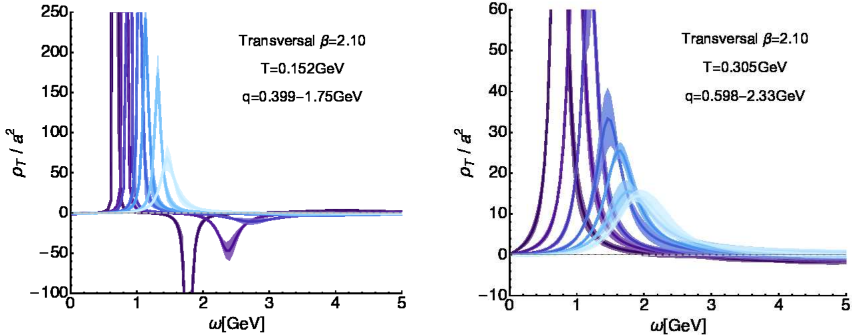
Figure: Left: Momentum dependence of the longitudinal quasi-particle peak position at $\beta = 2.10$ (takes for small momenta a non-zero value). At the lowest temperatures within the hadronic phase one finds a larger intercept than in deconfinement. Right: Fit of the lowest and highest temperature curves with the ansatz $\omega_L^0(|\vec{q}|) = A\sqrt{B^2 + |\vec{q}|^2}$. (Quasiparticle mass defined as $m = AB$.) Debye mass from $N_f = 2 + 1$ lattice QCD given for comparison.

Longitudinal mass

$$\begin{aligned} m_L/T|_{T=0.152\mathrm{GeV}} &= 3.80 \pm 0.25 \\ m_L/T|_{T=0.381\mathrm{GeV}} &= 2.97 \pm 0.16 \end{aligned} \qquad (2)$$

Outline

Transversal spectral function in confinement and deconfinement



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Transversal quasi-particle peak position as function of momentum

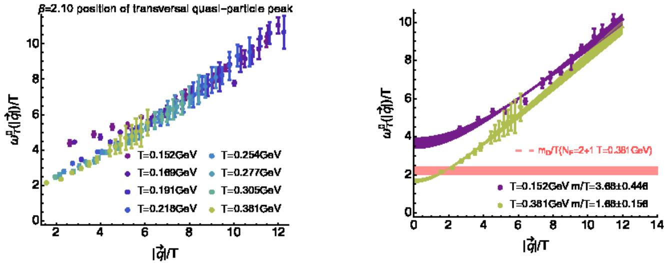

Figure: Left: Momentum dependence of the transversal quasi-particle peak position at $\beta = 2.10$ (takes for small momenta a non-zero value). At the lowest temperatures within the hadronic phase one finds a larger intercept than in deconfinement. Right: Fit of the lowest and highest temperature curves with the ansatz $\omega_L^0(|\vec{q}|) = A\sqrt{B^2 + |\vec{q}|^2}$. (Quasiparticle mass defined as $m = AB$.) Debye mass from $N_f = 2 + 1$ lattice QCD given for comparison.

Transversal mass

$$m_T/T|_{T=0.152\text{GeV}} = 3.68 \pm 0.45$$
$$m_T/T|_{T=0.381\text{GeV}} = 1.68 \pm 0.16 \quad (3)$$

- Confined phase:
  The values for the transversal quasi-particle mass agree with the longitudinal masses (given above) in confinement.
- Deconfined phase :
  The values for the transversal quasi-particle mass differ significantly from the longitudinal masses (given above) in deconfinement.
- The present statistics is not sufficient to study the width of the peak as function of temperature more in detail.

Outline

Outlook

This pioneer study ...

- shall be sytematically extended to lower light quark masses : repeat the investigation for the $m_\pi \approx 200\ \mathrm{MeV}$ ensembles next.
- shall be critically questioned : Influence of quality of gauge fixing.

There should be a ...

- methodical study to copare with other tools of analytical continuation (*e.g.* Tichonov regularization). A collaborative effort together with the Coimbra group is planned.
- Hopefully, in future we will be able to extend the study to the quark spectral function from lattice data, following the recent SDE (proof of principle) investigation
  "Bayesian analysis of quark spectral properties from the Dyson-Schwinger equation",
  Christian S. Fischer, Jan M. Pawlowski, Alexander Rothkopf, and Christian A. Welzbacher, arXiv:1705.03207