



# Reconciling cosmology and short-baseline experiments with invisible decay of light sterile neutrinos

Based on arXiv:1404.1794, arXiv:1404.6160

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# Neutrino Oscillations - I

- ▶ neutrino existence proposed by Pauli (1930) to explain  $\beta$  decay
- ▶ first time observed in 1956 by C. Cowan, F. Reines
- ▶ oscillations proposed in 1957 by B. Pontecorvo
- ▶ “massless” until oscillations detected in 1998 (SuperKamiokande)
- ▶  $\nu$  oscillate only if they have different masses (even if very small)  
 $\Rightarrow$  not all of them are massless

Neutrino oscillations: analogous to CKM mixing for quarks, with

$$\nu_\alpha = \sum_{k=1}^3 U_{\alpha k} \nu_k \quad (\alpha = e, \mu, \tau)$$

$\nu_\alpha$  flavour eigenstates,  $U_{\alpha k}$  PMNS mixing matrix,  $\nu_k$  mass eigenstates.

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# Neutrino Oscillations - II

Oscillations sensitive only to mass differences, not to absolute mass scale!

Effective 2 neutrino mixing ( $\Delta m_{21}^2 = m_2^2 - m_1^2$ ,  $\theta_{12}$  mixing angle):

$$P_{\alpha \rightarrow \beta, \alpha \neq \beta} = \sin^2(2\theta_{12}) \sin^2\left(\frac{\Delta m_{21}^2 L}{4E}\right)$$

Current knowledge of the active  $\nu$  mixing:

$$\begin{aligned}\Delta m_{SOL}^2 &= (7.50 \pm 0.20) \cdot 10^{-5} \text{ eV}^2 &= \Delta m_{21}^2 \\ \Delta m_{ATM}^2 &= (2.32_{-0.08}^{+0.12}) \cdot 10^{-3} \text{ eV}^2 &= |\Delta m_{32}^2| \simeq |\Delta m_{31}^2| \\ \sin^2(2\theta_{12}) &= 0.857 \pm 0.024 \\ \sin^2(2\theta_{23}) &> 0.95 \\ \sin^2(2\theta_{13}) &= 0.095 \pm 0.010\end{aligned}$$

[PDG - Beringer et al. (2013)]

CP violation possible only if  $\sin \theta_{13} \neq 0$

CP violating phase still unknown.

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Problem: observed anomalies in short baseline experiments  $\Rightarrow$  deviations from standard 3- $\nu$  description?

A short review: [Fan, Langacker, 2012]

- ▶ *LSND*: search for  $\bar{\nu}_\mu \rightarrow \bar{\nu}_e$ , with  $L/E = 0.4 \div 1.5$  m/MeV. Observed a  $3.8\sigma$  excess of  $\bar{\nu}_e$  events [Aguilar et al., 2001]
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Possible explanation:

$$\Delta m_{SBL}^2 \simeq 1 \text{ eV}^2$$

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# Sterile Neutrino mass

SBL anomaly  $\Rightarrow \Delta m_{SBL}^2 \simeq 1 \text{ eV}^2$  [Giunti et al., 2013]



Existence of an additional neutrino degree of freedom,  
mass around 1 eV, no weak interaction  $\Rightarrow$  *sterile*.



3 active ( $m_i \ll 1 \text{ eV}$ ) + 1 sterile ( $m_s \simeq 1 \text{ eV}$ )  $\nu$  scenario

We must update our mixing paradigm:

$$\nu_\alpha = \sum_{k=1}^{3+1} U_{\alpha k} \nu_k \quad (\alpha = e, \mu, \tau, \textcolor{red}{s})$$

[Giunti et al, 2013]

$$0.82 \leq \Delta m_{SBL}^2 / \text{eV}^2 \leq 2.19$$

(3 $\sigma$ )

$\nu_s$  is mainly  $\nu_4$ :

$$m_s \simeq m_4 \simeq \sqrt{\Delta m_{SBL}^2}$$



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# Additional Neutrino in Cosmology

Sterile  $\nu$  in cosmology: distribution function  $f_s(p) = \frac{\beta_s}{e^{p/\alpha_s T_\nu} + 1}$

Contribution of the  $\nu_s$  to cosmology described with: [Acero, Lesgourgues, 2009]

►  $\Delta N_{\text{eff}} = N_{\text{eff}} - 3.046$ :  $\rho_R = \left[ 1 + \frac{7}{8} \left( \frac{T_\nu}{T_\gamma} \right)^4 N_{\text{eff}} \right] \rho_\gamma$ , it becomes

$$\Delta N_{\text{eff}} = \beta_s \alpha_s^4$$

►  $m_s^{\text{eff}} = (94.1 \text{ eV}) \omega_s = \rho_s / \rho_c^0$ , from which we obtain  $m_s^{\text{eff}} = m_s \beta_s \alpha_s^3$

Constant is given by  $\sum m_i = (94.1 \text{ eV}) \omega_\nu$  for SM neutrinos.

Problem: 2 observables  $(\Delta N_{\text{eff}}, m_s^{\text{eff}})$ , 3 parameters  $(\alpha_s, \beta_s, m_s)$ !

Hp:  $\nu_s$  follows a thermal distribution with  $T_s = \alpha_s T_\nu$  and  $\beta_s = 1$ .

$$\Rightarrow m_{TH}^{\text{eff}} = m_s (\Delta N_{\text{eff}}^{TH})^{3/4}$$

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►  $\Delta N_{\text{eff}} = N_{\text{eff}} - 3.046$ :  $\rho_R = \left[ 1 + \frac{7}{8} \left( \frac{T_\nu}{T_\gamma} \right)^4 N_{\text{eff}} \right] \rho_\gamma$ , it becomes

$$\Delta N_{\text{eff}} = \beta_s \alpha_s^4$$

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Problem: 2 observables  $(\Delta N_{\text{eff}}, m_s^{\text{eff}})$ , 3 parameters  $(\alpha_s, \beta_s, m_s)$ !

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# Parameters

In the following we will study the Universe evolution considering a  $\Lambda$ CDM +  $r_{0.002} + \nu_s$  model with 9 free parameters:

$$\{\omega_{\text{CDM}}, \omega_b, \theta_s, \tau, \ln(10^{10} A_s), n_s\} + r_{0.002} + \{\Delta N_{\text{eff}}, m_s\}$$

$\omega_{\text{CDM}}$  - CDM density today

$\omega_b$  - baryon density today

$\theta_s$  - angular sound horizon

$\tau$  - optical depth to reionization

$\ln(10^{10} A_s)$  - amplitude and

$n_s$  tilt of the primordial power spectrum

$r_{0.002}$  - tensor to scalar ratio at 0.002  $\text{Mpc}^{-1}$

$\Delta N_{\text{eff}}$  effective number of  $\nu_s$

$m_s$  physical mass of  $\nu_s$

Assume:

- ▶  $\sum m_{\nu, \text{active}} = 0.06 \text{ eV}$  (minimal for Normal Hierarchy)
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# Datasets for the CosmoMC analysis

MCMC with CosmoMC with different cosmological data:

- ▶ *Planck*: Planck TT spectra
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- ▶  $H_0$ :  $H_0 = 73.8 \pm 2.4 \text{ km s}^{-1} \text{ Mpc}^{-1}$ , using *Cepheids and SN Ia*.
- ▶ *CFHTLenS*: the CFHTLenS 2D cosmic shear correlation function (from redshifts and shapes of 4.2 million galaxies with  $0.2 < z < 1.3$ ).
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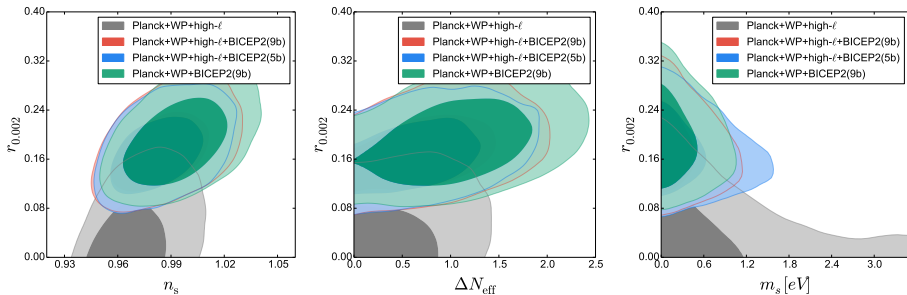
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# Results - I



First tension:  $r_{0.002}$  (with and without BICEP2)

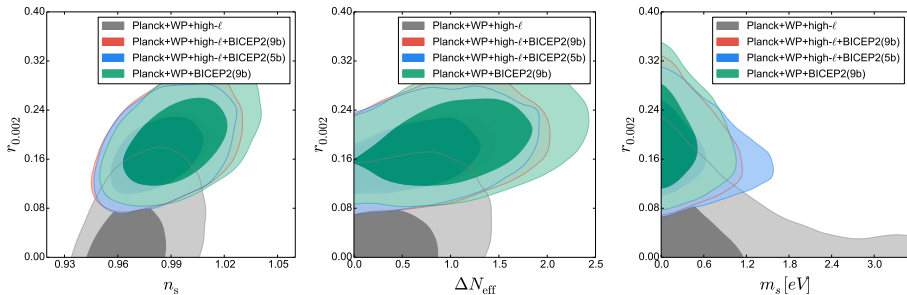
We must wait Planck 2014 data release, with polarization data

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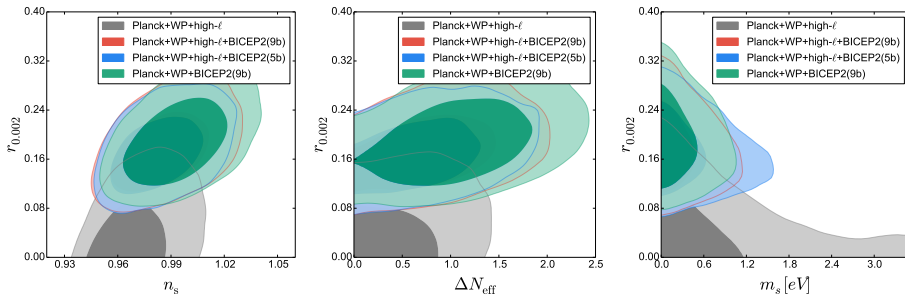
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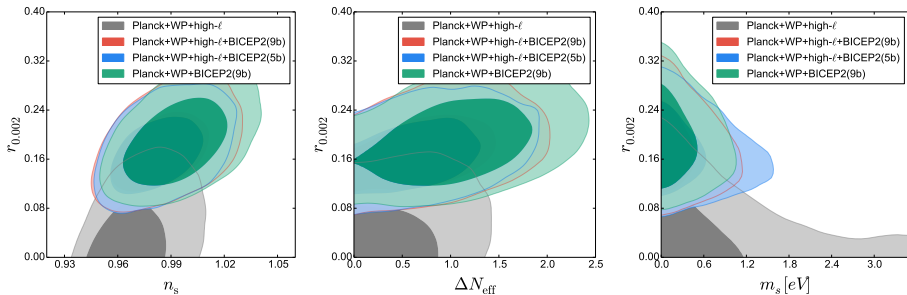
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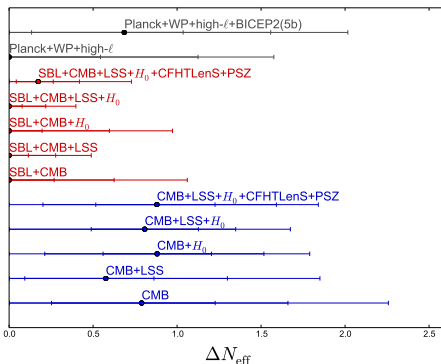
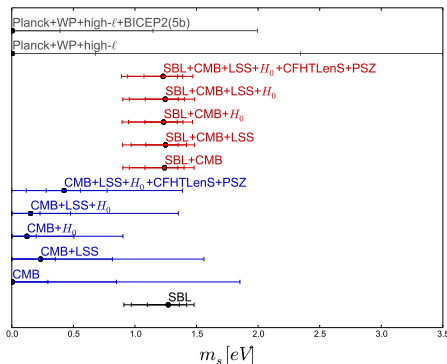
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# Results - II



Second tension:  $m_s$  vs  $\Delta N_{\text{eff}}$  (with and without SBL)

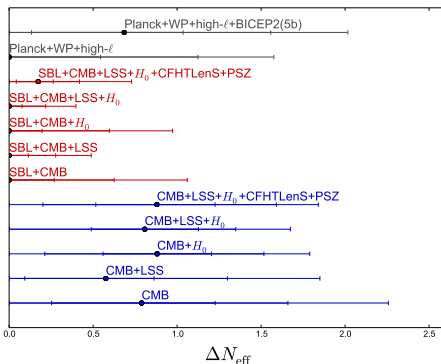
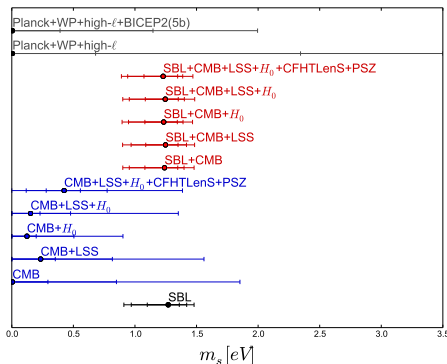
Notice: small  $\Delta N_{\text{eff}}$  if  $m_s \sim 1$  eV

$\Rightarrow \nu_s$  cannot be fully thermalized,  $\Delta N_{\text{eff}} \ll 1 \rightarrow T_s \ll T_\nu$

Notice: CFHTLenS and PSZ data give a preference ( $> 2\sigma$ ) for  $m_s > 0$ , but  $m_s \sim 0.5$  eV and lower than 1 eV at  $> 2\sigma$



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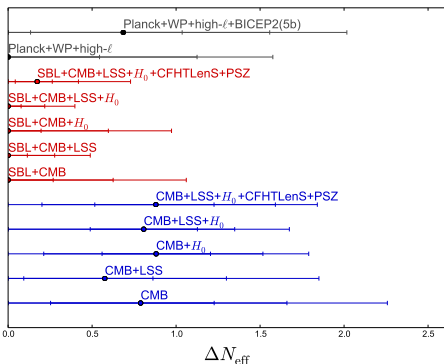
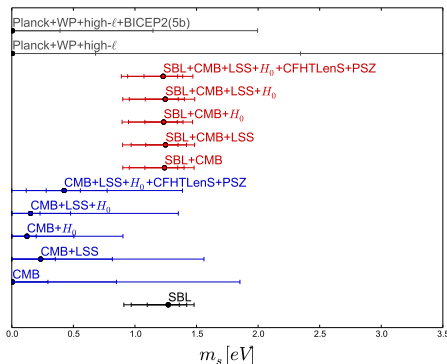
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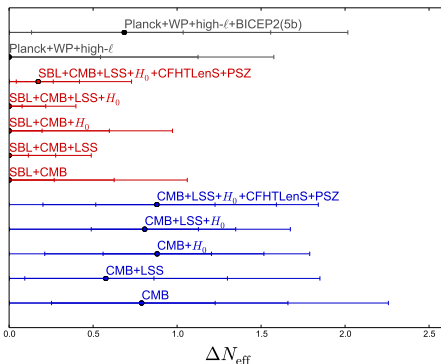
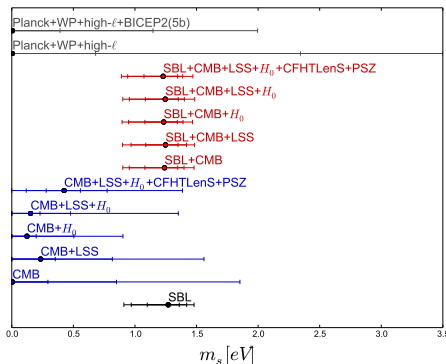
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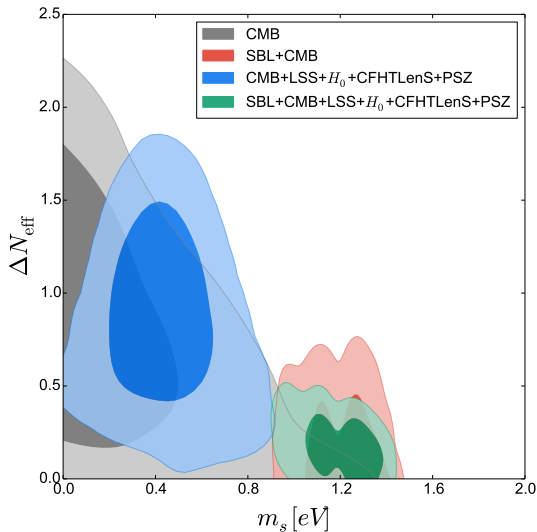
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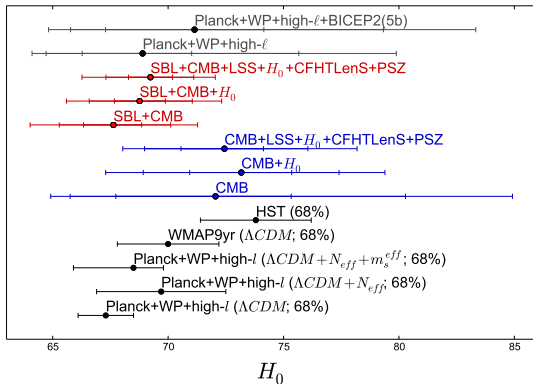
2D marginalized posterior distribution for  $\Delta N_{\text{eff}}$ ,  $m_s$

Comparison:

- CMB only
- complete dataset

with and without SBL data

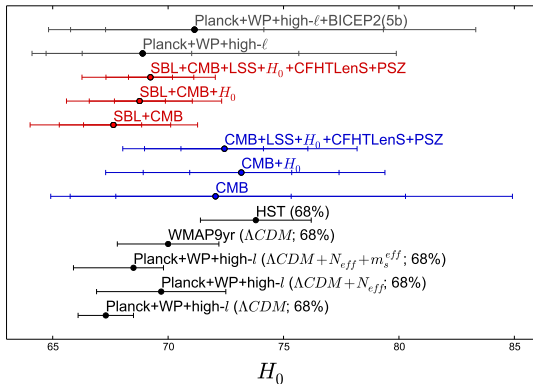
## Results - IV



Third tension:  $H_0$

- ▶ Planck vs local measurements
- ▶ value inferred from CMB highly model-dependent: correlation with  $N_{eff}$ 
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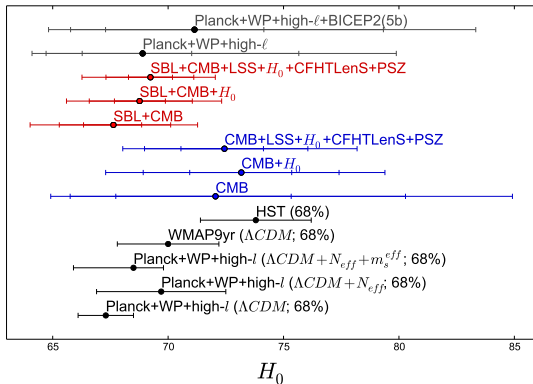
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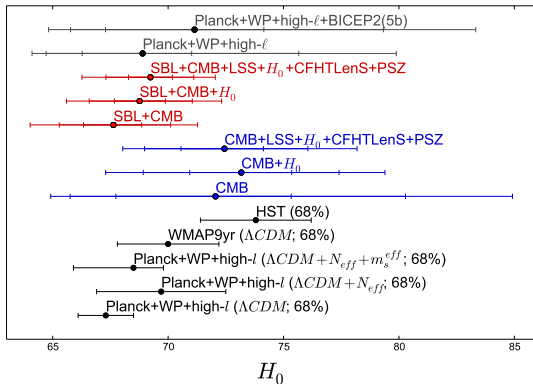
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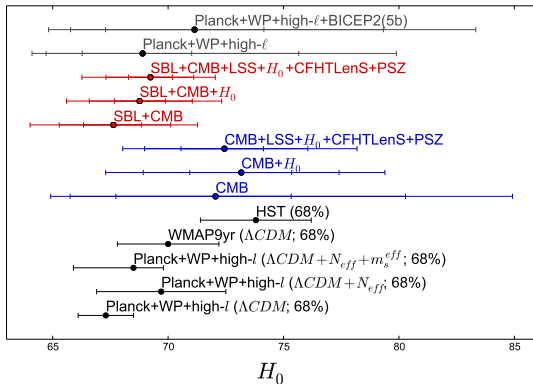


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# Cosmological invisible decay of light sterile neutrinos

Proposed solution for solve the encountered tensions:

$\nu_s$  can decay - lifetime  $\tau_s$  comparable with Age of the Universe  $t_U$

Decay products belong to the sterile sector  $\Rightarrow$  very weak interaction, invisible

Effective number of  $\nu_s$ :  $N_s(t) = \Delta N_{\text{eff}} \cdot e^{-t/\tau_s}$

$\tau_s$  assumed to be constant (no energy dependent).

Decay products have negligible mass: they can be accounted as radiation with effective number  $N_{dp}(t) = \Delta N_{\text{eff}} \cdot (1 - e^{-t/\tau_s})$

Energy distribution of the invisible decay products neglected for simplicity.

Computational problem: time  $t$  depends on the energy density contributions.

Approximation:  $t$  calculated considering matter dominated universe. True except:

- ▶ initial radiation domination - very short
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- ▶ final  $\Lambda$  domination - largest part of  $\nu_s$  has decayed

# Cosmological invisible decay of light sterile neutrinos

Proposed solution for solve the encountered tensions:

$\nu_s$  can decay - lifetime  $\tau_s$  comparable with Age of the Universe  $t_U$

Decay products belong to the sterile sector  $\Rightarrow$  very weak interaction, invisible

Effective number of  $\nu_s$ :  $N_s(t) = \Delta N_{\text{eff}} \cdot e^{-t/\tau_s}$

$\tau_s$  assumed to be constant (no energy dependent).

Decay products have negligible mass: they can be accounted as radiation with effective number  $N_{dp}(t) = \Delta N_{\text{eff}} \cdot (1 - e^{-t/\tau_s})$

Energy distribution of the invisible decay products neglected for simplicity.

Computational problem: time  $t$  depends on the energy density contributions.

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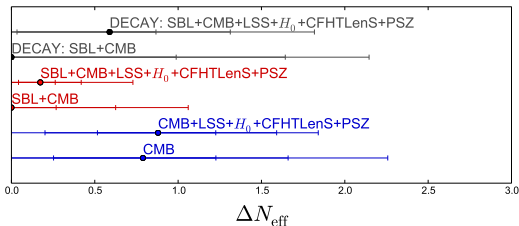
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# Results - I

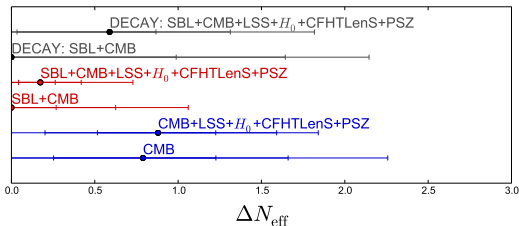


$\Delta N_{\text{eff}} = 1$  is allowed

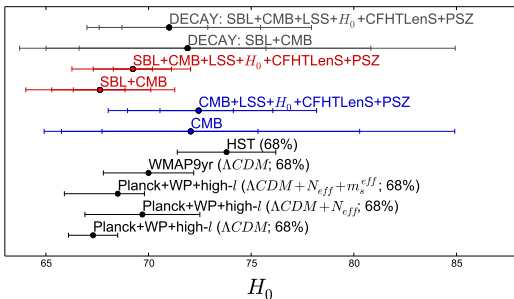
$H_0$  compatible with local measurements (HST)

With sterile neutrino decay,  $\Delta N_{\text{eff}}$  and  $H_0$  are at the same level than the ones without SBL prior

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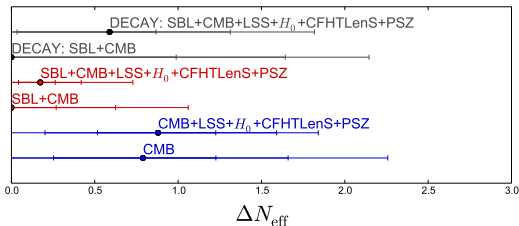
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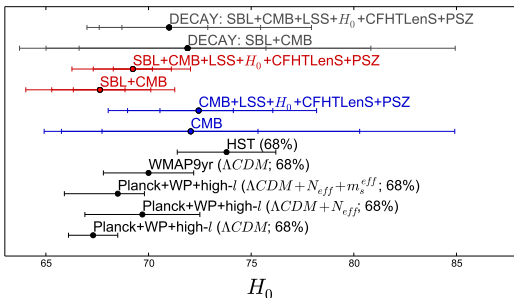
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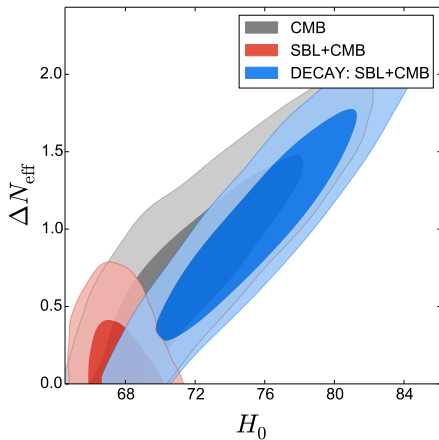
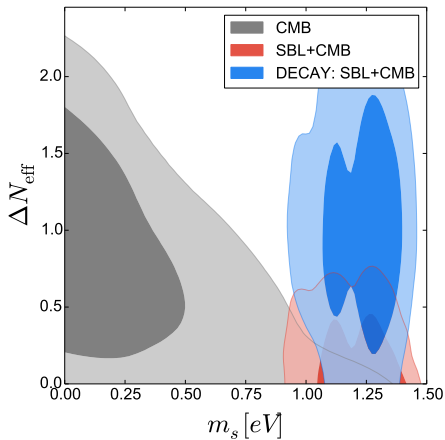
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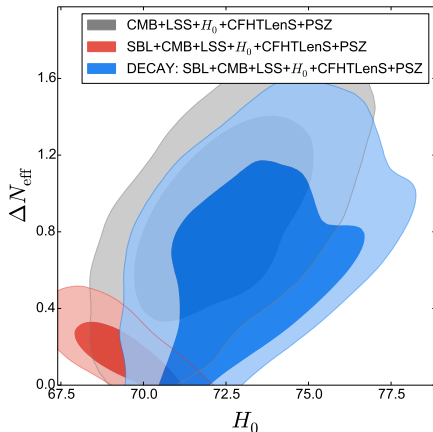
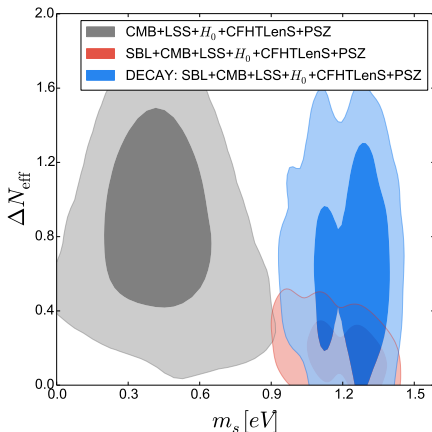
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## Results - II - CMB only



High  $\Delta N_{\text{eff}}$  even with SBL mass  
Correlation between  $\Delta N_{\text{eff}}$  and  $H_0$  recovered

## Results - III - full dataset



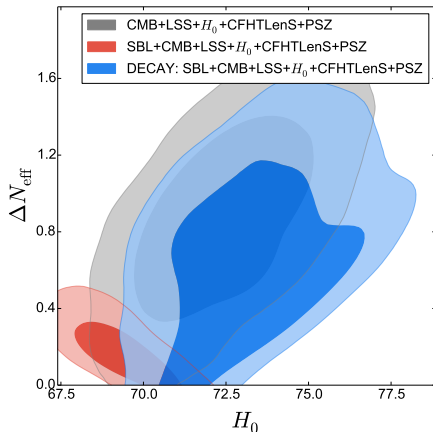
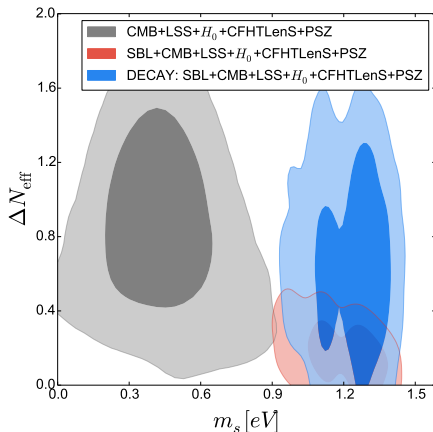
$\Omega_{\nu_s}$  can explain cluster data since it is related both to  $m_s$  and  $\Delta N_{\text{eff}}$ :

$$\Omega_{\nu_s} \propto N_s(t)^{3/4} m_s$$

Correlation between  $\Delta N_{\text{eff}}$  and  $H_0$  recovered

Shape in  $\Delta N_{\text{eff}} - H_0$  plot due to volume effects in the Bayesian analysis

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# Thank you for the attention!

Further details:

[Archidiacono, Fornengo, Gariazzo, Giunti, Hannestad, Laveder, arxiv:1404.1794]

[Gariazzo, Giunti, Laveder, arxiv:1404.6160]



## Correlation between $r_{0.002}$ and $\Delta N_{\text{eff}}$

BICEP2: higher  $r_{0.002}$  that correspond to more large-scale fluctuations.

Primordial power spectrum:

$$\mathcal{P}_k = A_s (k/k_0)^{n_s-1}$$

$k_0$  pivot scale,  $A_s$  amplitude,  $n_s$  tilt

Higher  $r_{0.002}$  can be compensated with an increase of  $n_s \rightarrow$  decrease of large-scale fluctuations

Increase of  $n_s \rightarrow$  increase of small-scale fluctuations ( $k \gg k_0$ )



Effect can be compensated with an increase of  $N_{\text{eff}} \rightarrow$  decrease of small-scale fluctuations due to free streaming of relativistic particles