



# Perspectives of dark matter searches with antideuterons

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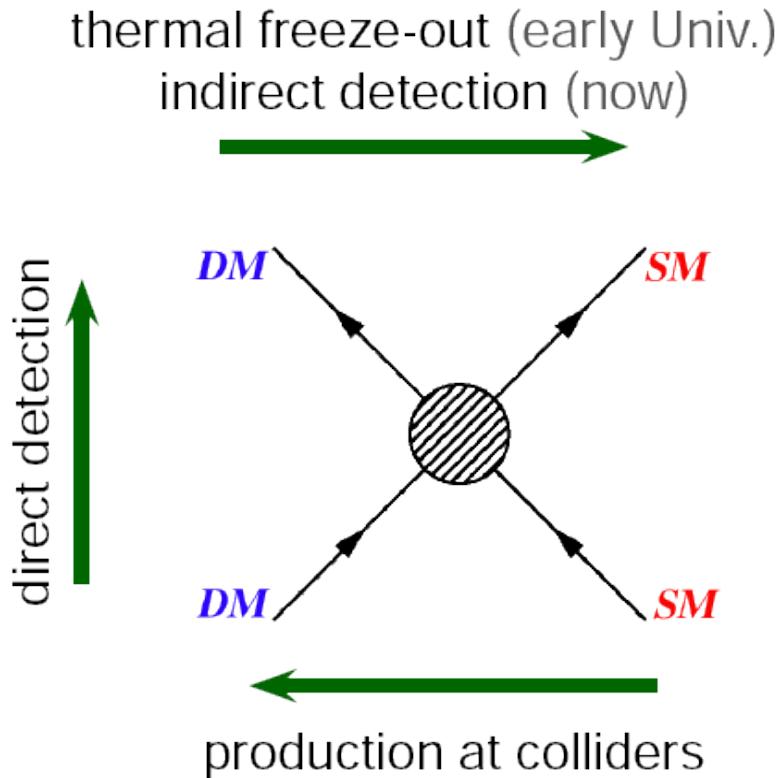
(In collaboration with N.Fornengo and L. Maccione)

Università degli studi di Torino and INFN Torino

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Roma, May 22 2013

# DM indirect detection



Dark matter can be indirectly detected by looking at the products of its annihilation (or decay)

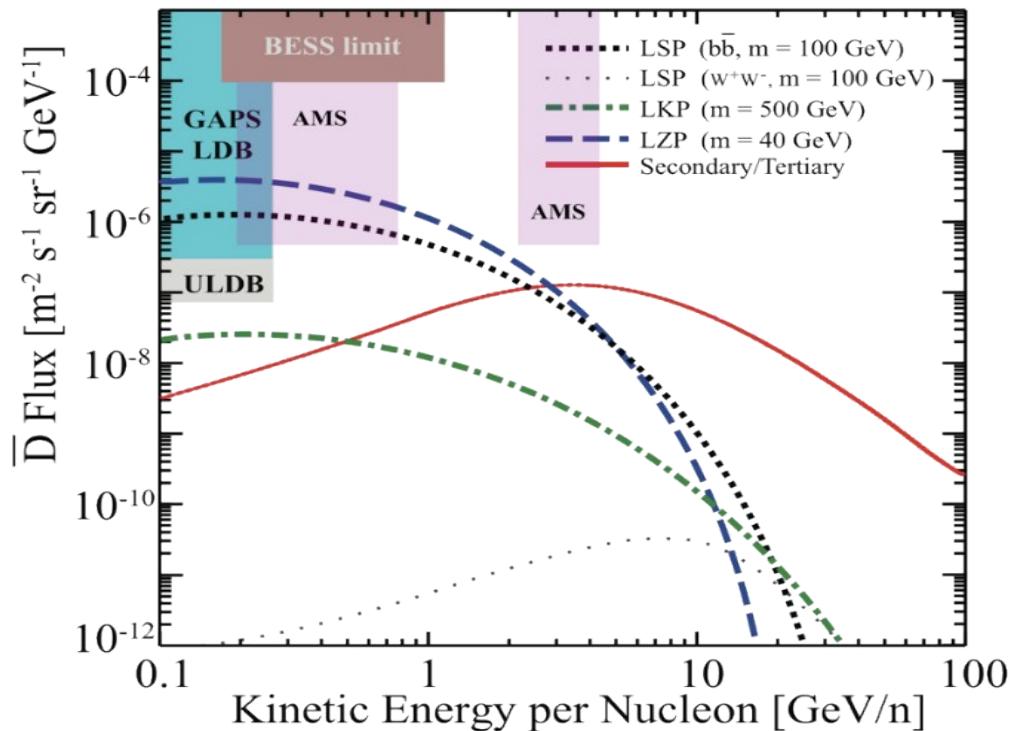
A good DM signature would be represented by an **excess of antimatter** in the cosmic rays flux

Antideuterons seems to be more promising than antiprotons or positrons

[F.Donato, N.Fornengo and P.Salati, Phys. Rev. D 62 (2000)]

# Why antideuterons?

For kinematical reasons, it is very difficult to produce low energy antideuterons in a spallation reaction



When propagating, astrophysical antiprotons lose their energy through inelastic but non-annihilating reactions while antideuterons fragmentate ( $B_d = 2.2 \text{ MeV}$ )

[H.Baer and S.Profumo, JCAP 12 (2005)]

Very large signal-to-background ratio in the low-energy band

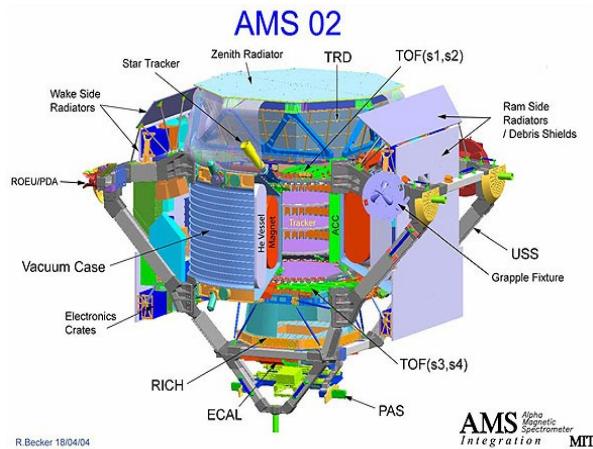
# Experiments

## ► GAPS



- It's a balloon borne experiment, scientific flights are expected to begin in 2016
- In the LDB+ setup, GAPS will fly for 210 days

## ► AMS



- It's a multipurpose experiment that will look for any kind of antimatter signal
- It's operative, on board of the ISS, since 2011
- We will consider, in this analysis, a data taking period of 3 years

# Antideuterons from DM annihilation



We work in a **model independent** framework (pure annihilations channels)

Our spectra are the **building blocks** for any DM model

Modeled with **Pythia 6.4.26** in its default configuration

Different phenomenological models can be used

What can we say about coalescence?

# Event per event coalescence

$$dN_{\bar{d}} = \int d^3 \vec{\Delta} \boxed{F_{\bar{p}, \bar{n}}} \times \boxed{C(\vec{\Delta})}$$

Probability that the two antinucleons coalesce

Probability that the two antinucleons are formed

We do not assume a form for the function  $F_{\bar{p}, \bar{n}}$

$$\vec{\Delta} = \vec{k}_{\bar{p}} - \vec{k}_{\bar{n}}$$
$$C(\vec{\Delta}) = \theta(\Delta^2 - p_0^2)$$

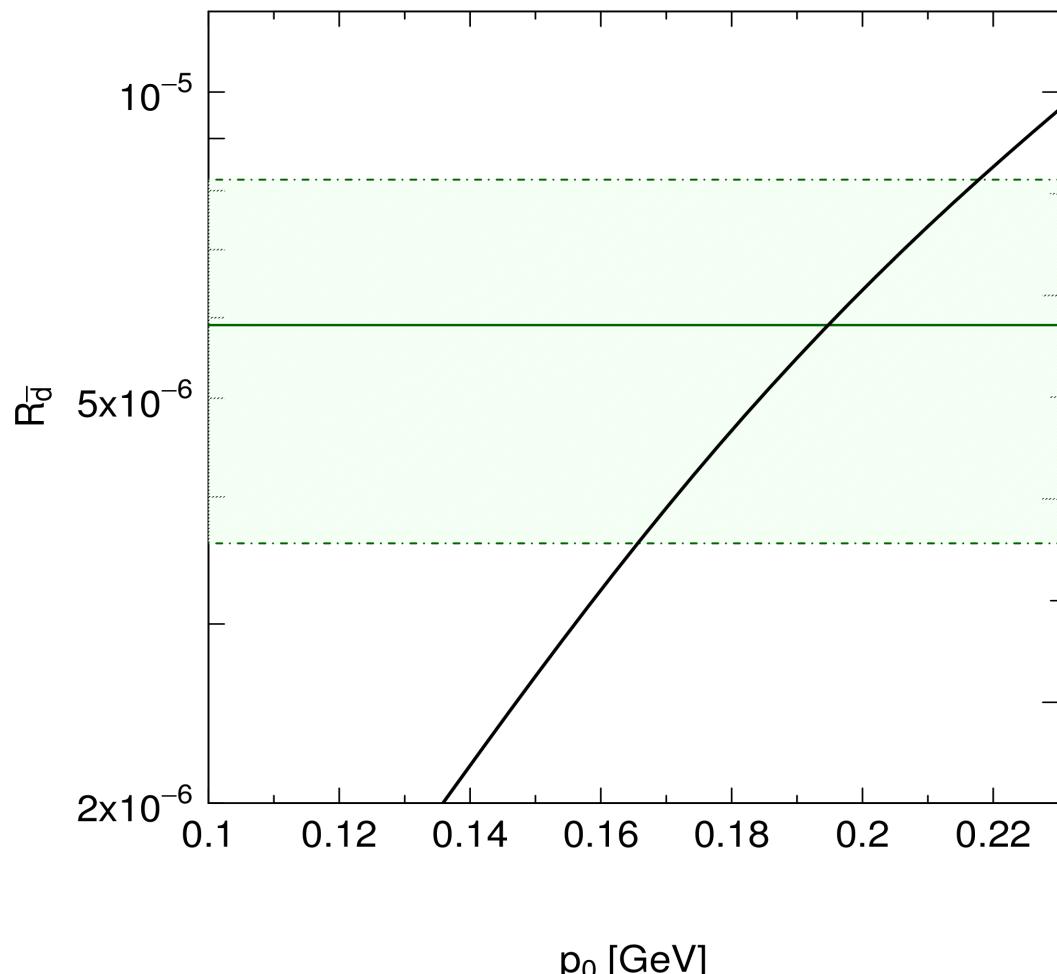
For each event we scan over the final state antinucleons and we consider them **as antideuterons if:**

$$\Delta p < p_0 \text{ and } \Delta r_{\min} < r_0$$

We assume  $r_0 = 2$  fm (radius of the  $\bar{d}$ ) while  $p_0$  **has to be tuned**

# Tuning the model

DM pair annihilation is assumed to be the **analogous to an  $e^+e^-$  collision**

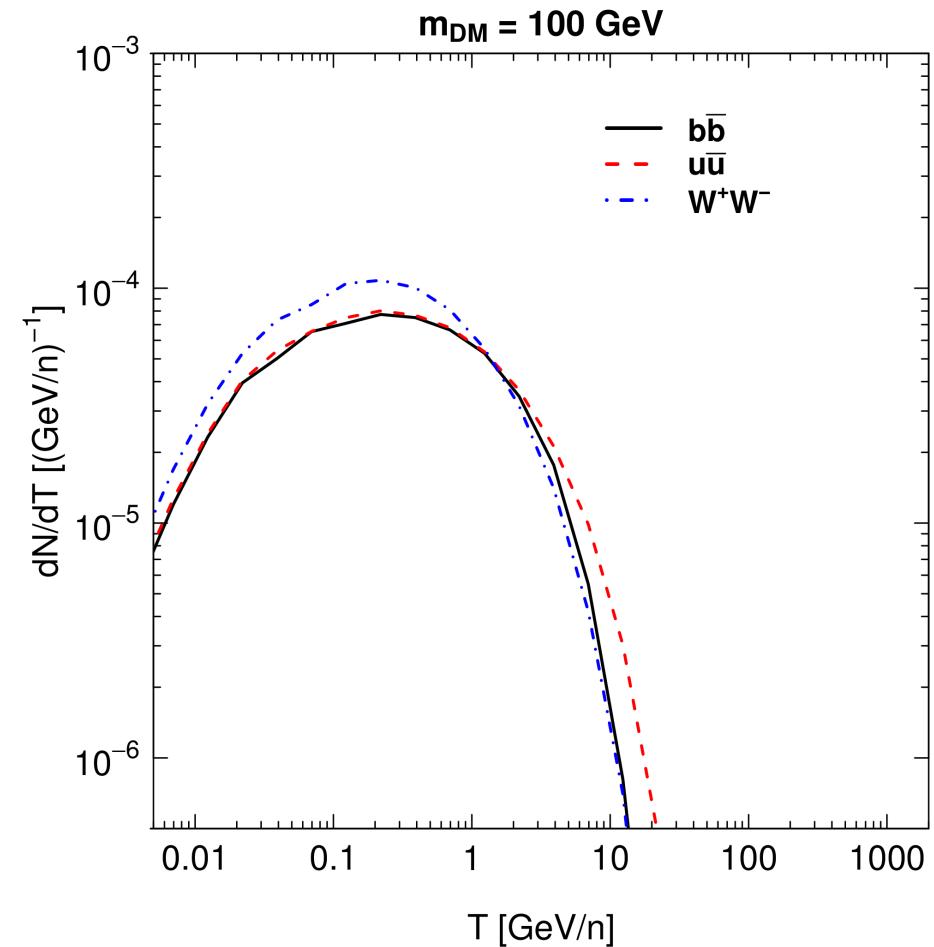
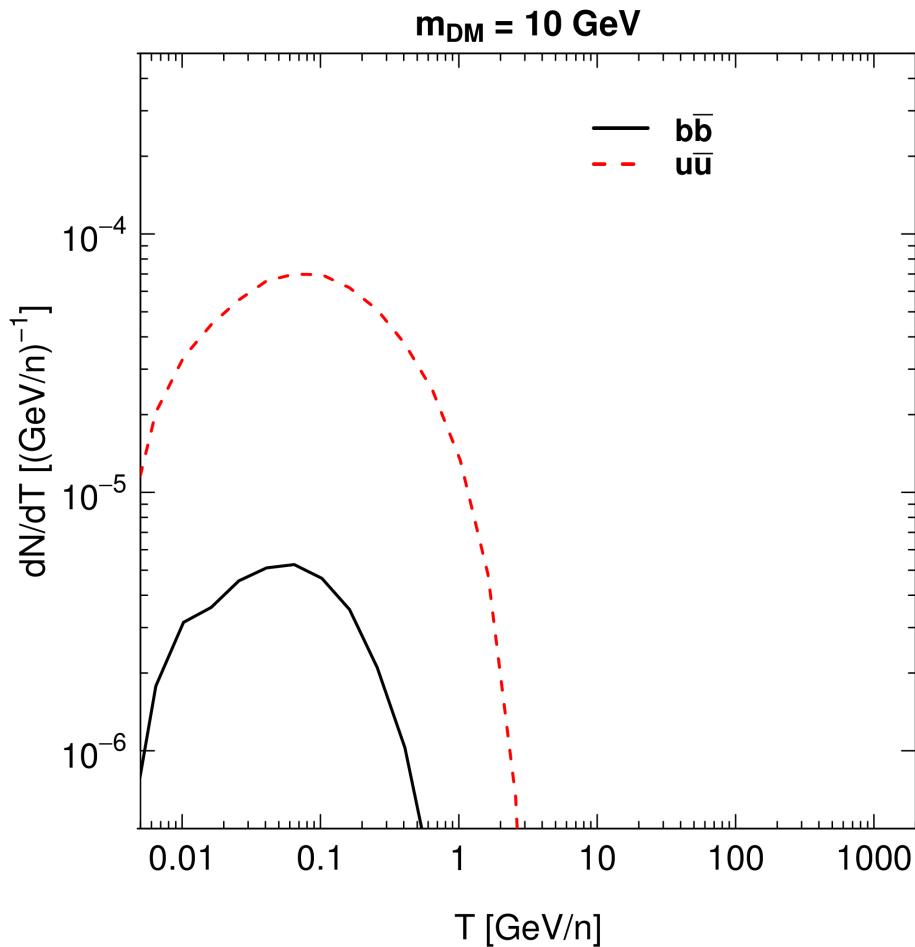


We tune our coalescence  
model in order to reproduce  
**ALEPH** data on  
antideuterons production  
rate at the Z resonance  
 $\sqrt{s}=91.2\text{ GeV}$

[ALEPH collaboration, Phys. Lett. B  
369 (2006) 192]

$$p_0 = (195 \pm 22) \text{ MeV}$$

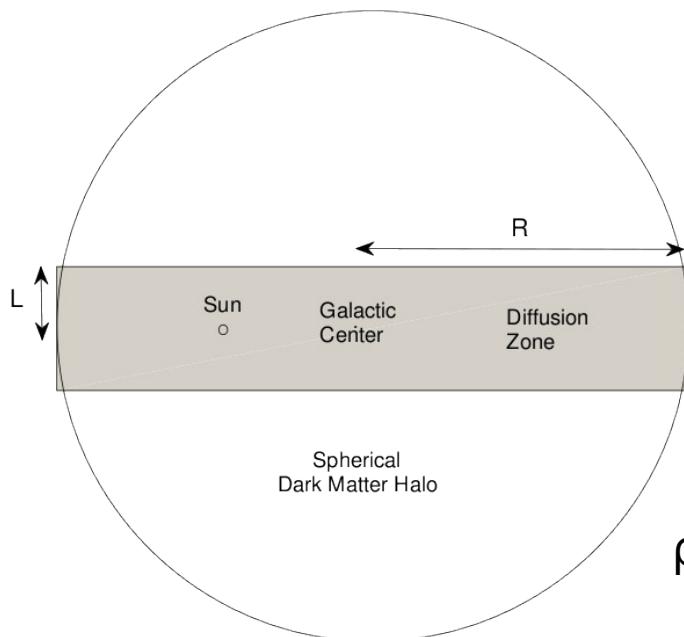
# Injected spectra



# Galactic propagation

$$-\nabla[K(\vec{r}, E)\nabla n_{\bar{d}}(\vec{r}, E)] + \frac{\partial}{\partial z} V_c(z) n_{\bar{d}}(E, \vec{r}) + 2h\delta(z)\Gamma_{ann} n_{\bar{d}}(E, \vec{r}) = q_{\bar{d}}(\vec{r}, E)$$

## Two-zone diffusion model



$$K(x, E) = K_0 \beta \left( \frac{R}{1GV} \right)^\delta$$

$$\vec{V}_c = sign(z) V_c$$

$$q_{\bar{d}}(x, E) = \frac{1}{2} \langle \sigma v \rangle \frac{dN_{\bar{d}}}{dT} \left( \frac{\rho(x)}{m_\chi} \right)^2$$

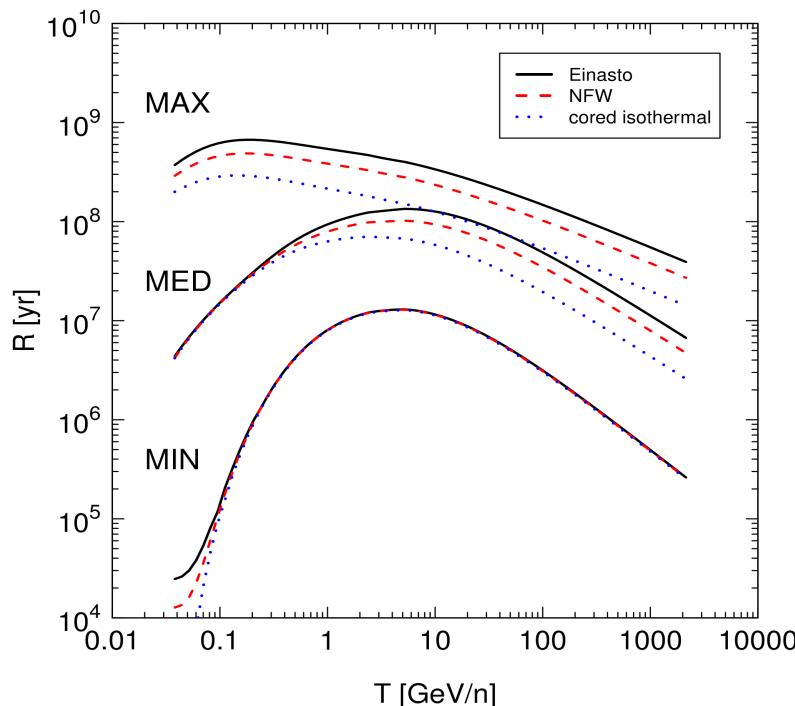
$\rho(r)$  Is the DM halo profile (e.g. Einasto, NFW, isothermal)

**CAVEAT : No energy losses and no reacceleration!**

# Galactic propagation - II

No reacceleration & energy losses → **the solution can be factorized:**

$$\phi_{\overline{d}}(T) = \left( \frac{\nu_{\overline{d}}}{4\pi} \left( \frac{\rho_{\odot}}{m_{\chi}} \right)^2 \frac{1}{2} \langle \sigma v \rangle \frac{dN_{\overline{d}}}{dT} \right) \times R_{\overline{d}}(T)$$



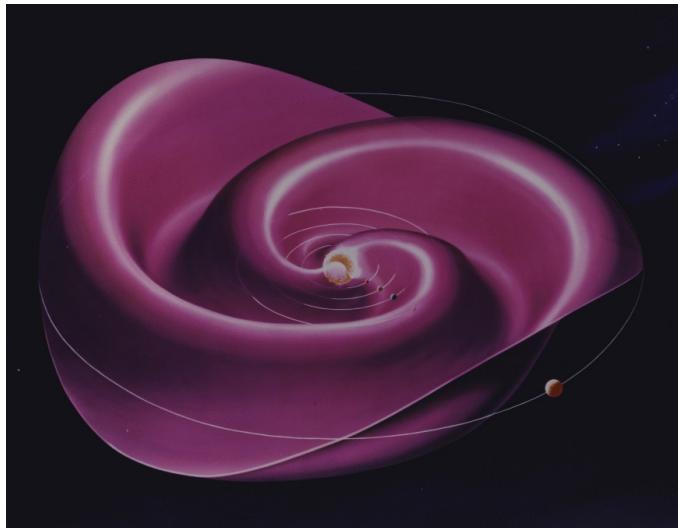
All the astrophysics is here!

	$\delta$	$K_0$ [Kpc <sup>2</sup> /Myr]	L [Kpc]	$V_c$ [Km/s]
MIN	0.85	0.0016	1	13.8
MED	0.70	0.0112	4	12
MAX	0.46	0.0765	15	5

$K_0, V_c$  and  $\delta$  constrained by B/C data

# Interactions with the heliosphere

**How can we model the motion of a charged particle inside the SMF?**



The Sun's magnetic field is a **large rotating spiral**

An **heliospheric current sheet (HCS)**, whose shape varies with time following solar activity, separates field lines directed toward or away from the sun

Generally, this is done by using the **force field approximation**:

$$\Phi_{TOA}(T_{TOA}) = \frac{2mT_{TOA} + T_{TOA}}{2mT_{IS} + T_{IS}} \Phi_{IS}(T_{IS})$$

$$T_{TOA} = T_{IS} - \phi$$

[J. S. Perko, Astron. Astrophys. 184 (1987) 119]

# Charge dependent solar modulation

Given the distribution  $f$  of a CR particle in the phase space:

[E. N. Parker, P&SS 13, 9 (1965)]

$$\frac{\partial f}{\partial t} = -(\vec{V}_{sw} + \vec{v}_d) \cdot \nabla f + \nabla \cdot (K \cdot \nabla f) + \frac{P}{3} (\nabla \cdot \vec{V}_{sw}) \frac{\partial f}{\partial p}$$

convection

drifts

diffusion (random walk)

adiabatic losses

We varies three parameters:

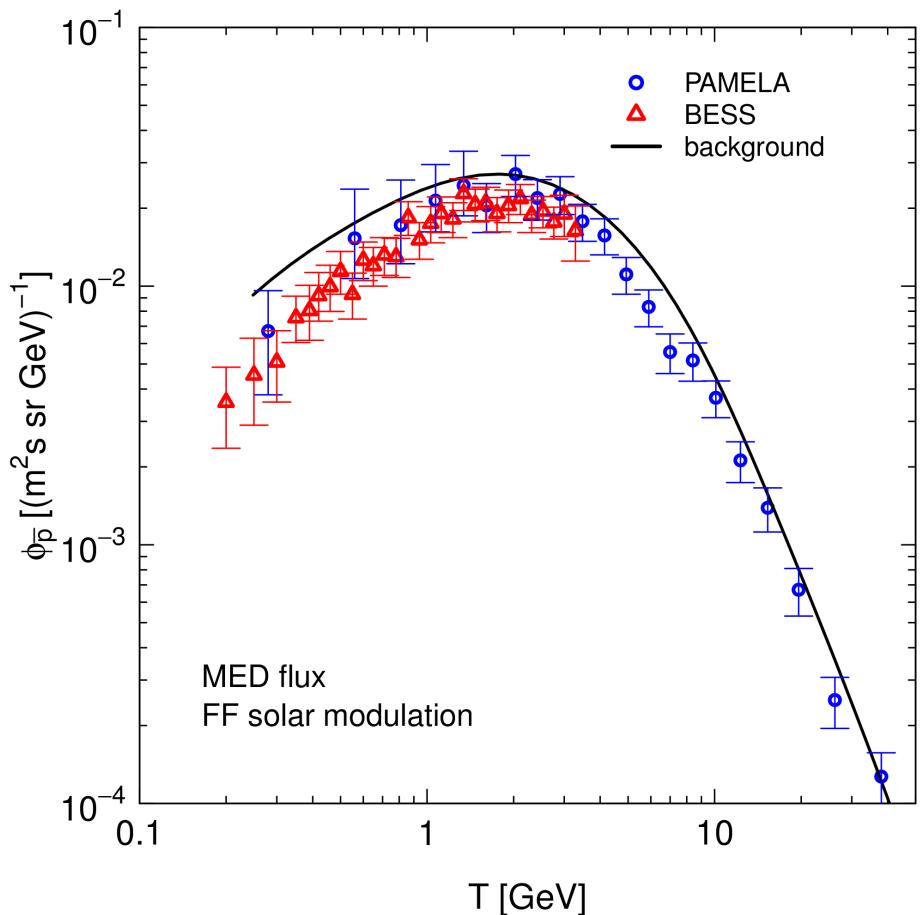
- The **tilt angle  $\alpha$** : it describes the spatial extent of the HCS
- The **mean free path  $\lambda$**  of the CR particle along the magnetic field direction
- The **spectral index  $\delta$**  that relates the solar diffusion to the particle's energy

We exploit the recently developed code **HELIOPROP**

[L. Maccione, Phys.Rev.Lett. 110, 081101(2013), arXiv:1211.6905 [astro-ph.HE]]

# Antiprotons bounds - I

The antiproton flux was recently measured by **PAMELA** and **BESS POLAR II**



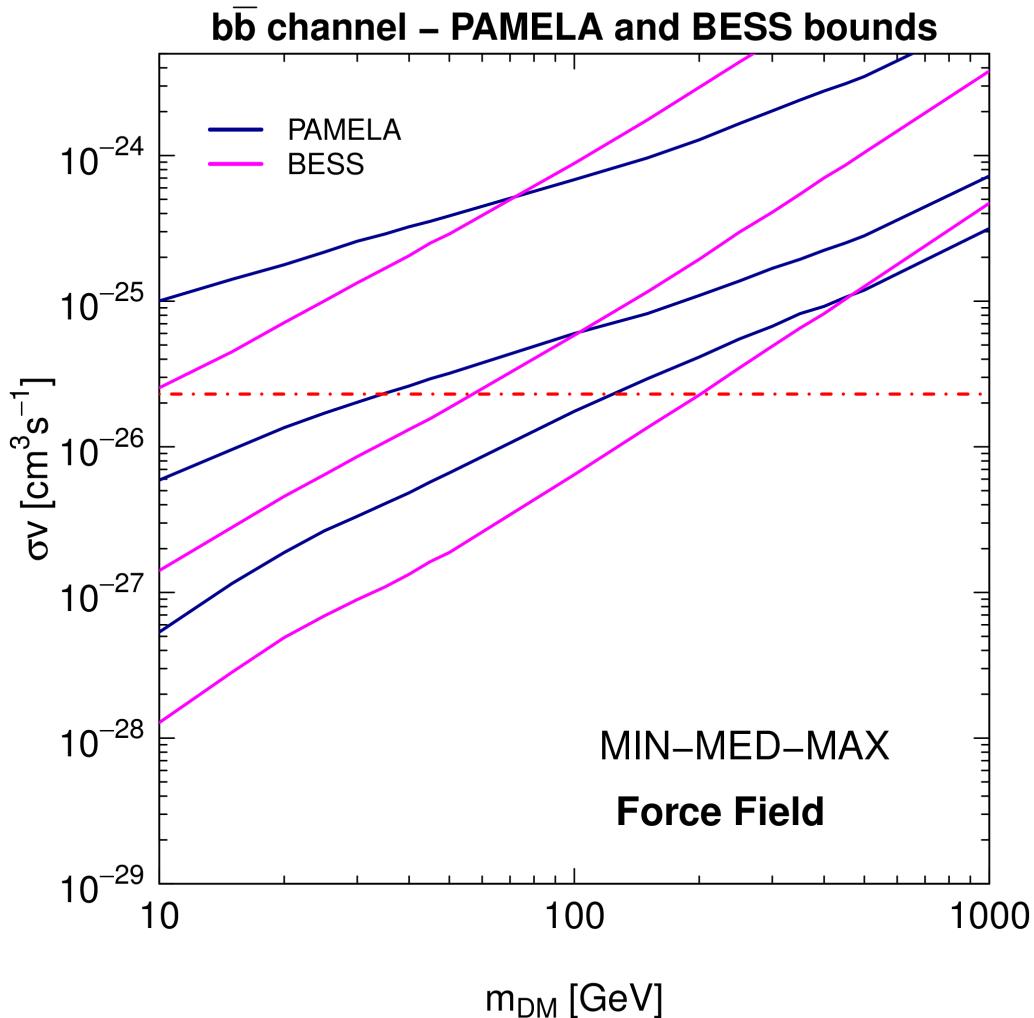
**The measured flux is well fitted by the pure background**

**Very little space for dark matter!**

We calculate the bounds on DM annihilation cross section by imposing an **uncertainty of the 40%** on the calculated astrophysical background

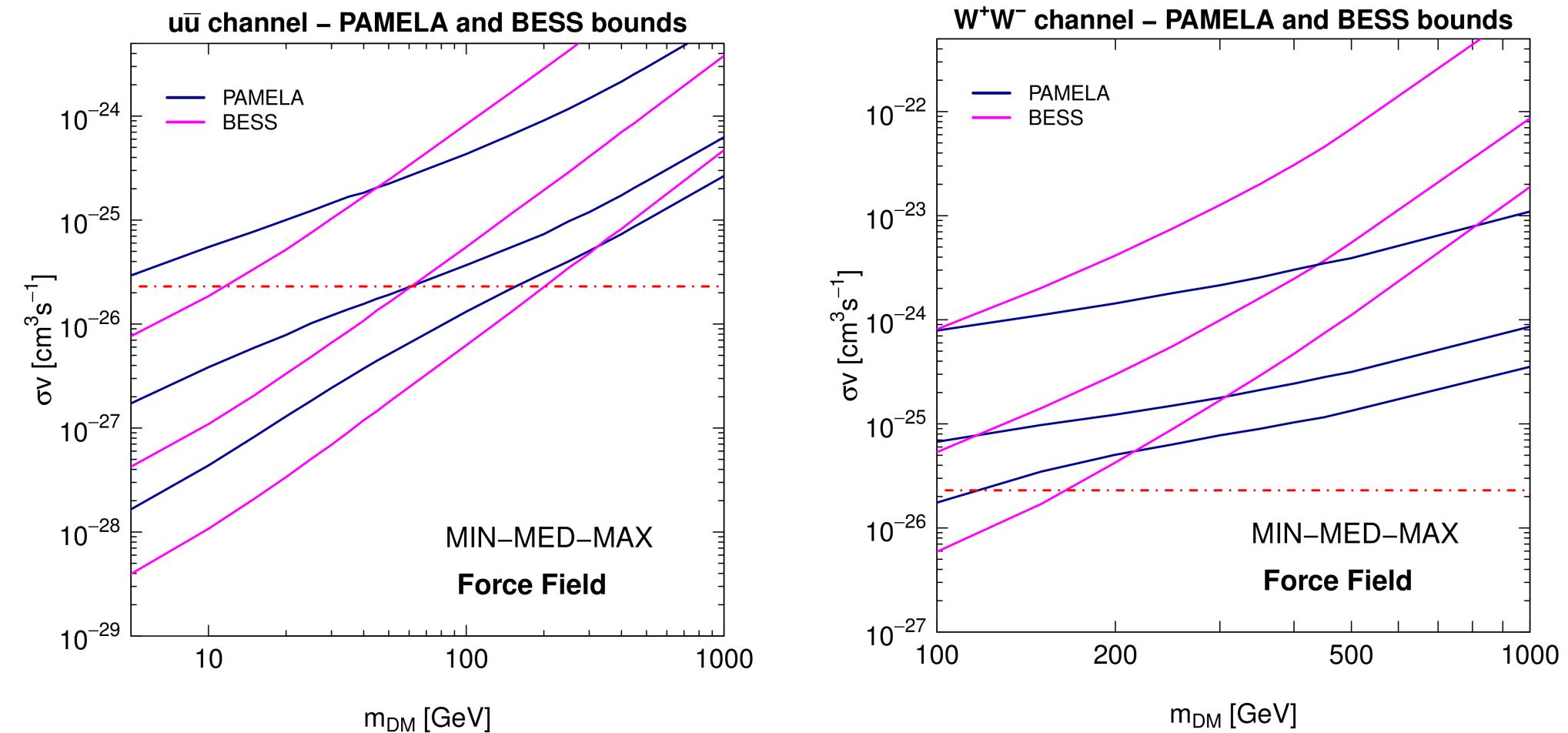
F.Donato, N.Fornengo and D.Maurin, Phys. Rev. D 78 (2008)

# Antiprotons bounds - II



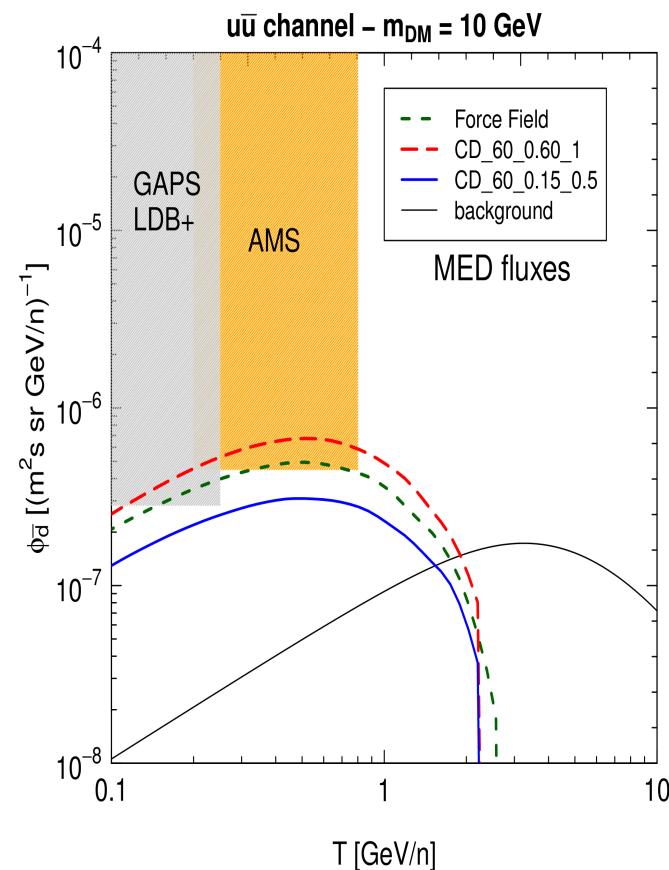
- ▶ Bounds calculated with a 99% c.l.
- ▶ 3 set of curves for 3 propagation models (MIN, MED and MAX)
- ▶ We want to be as conservative as possible -> in the following analysis we will only consider PAMELA bounds

# Antiprotons bounds - III

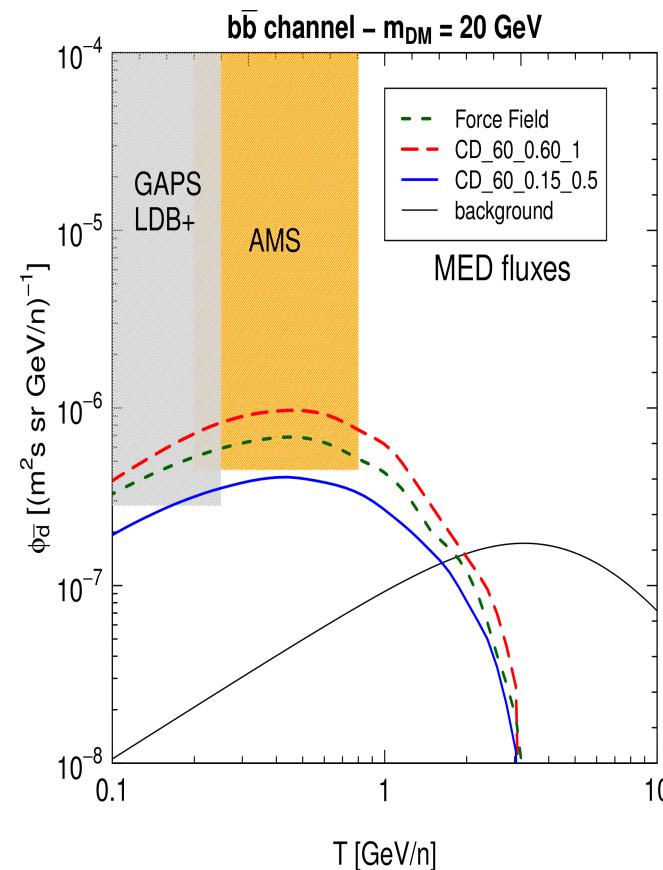


# Fluxes at Earth

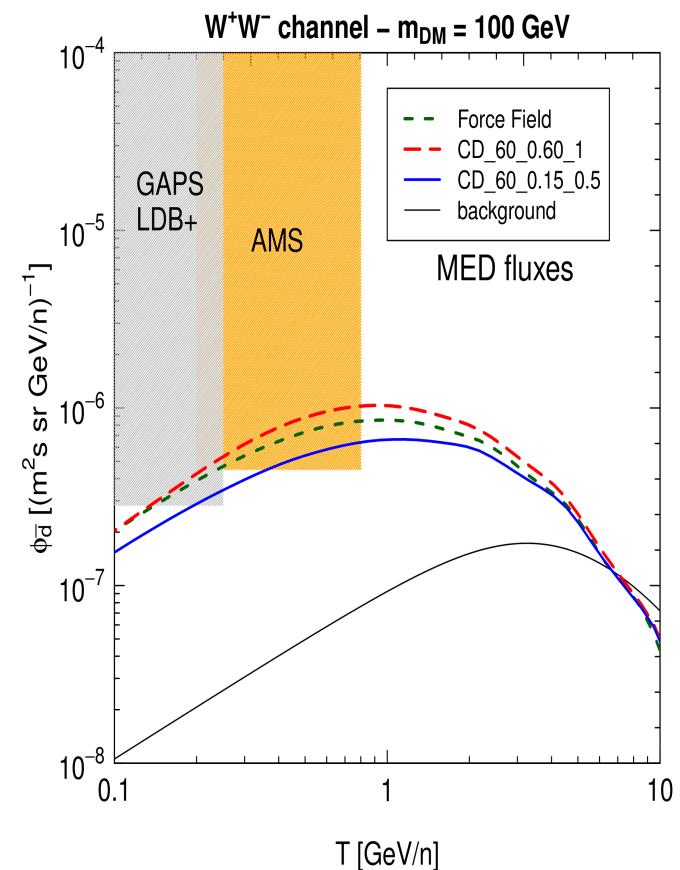
**uu channel**



**bb channel**

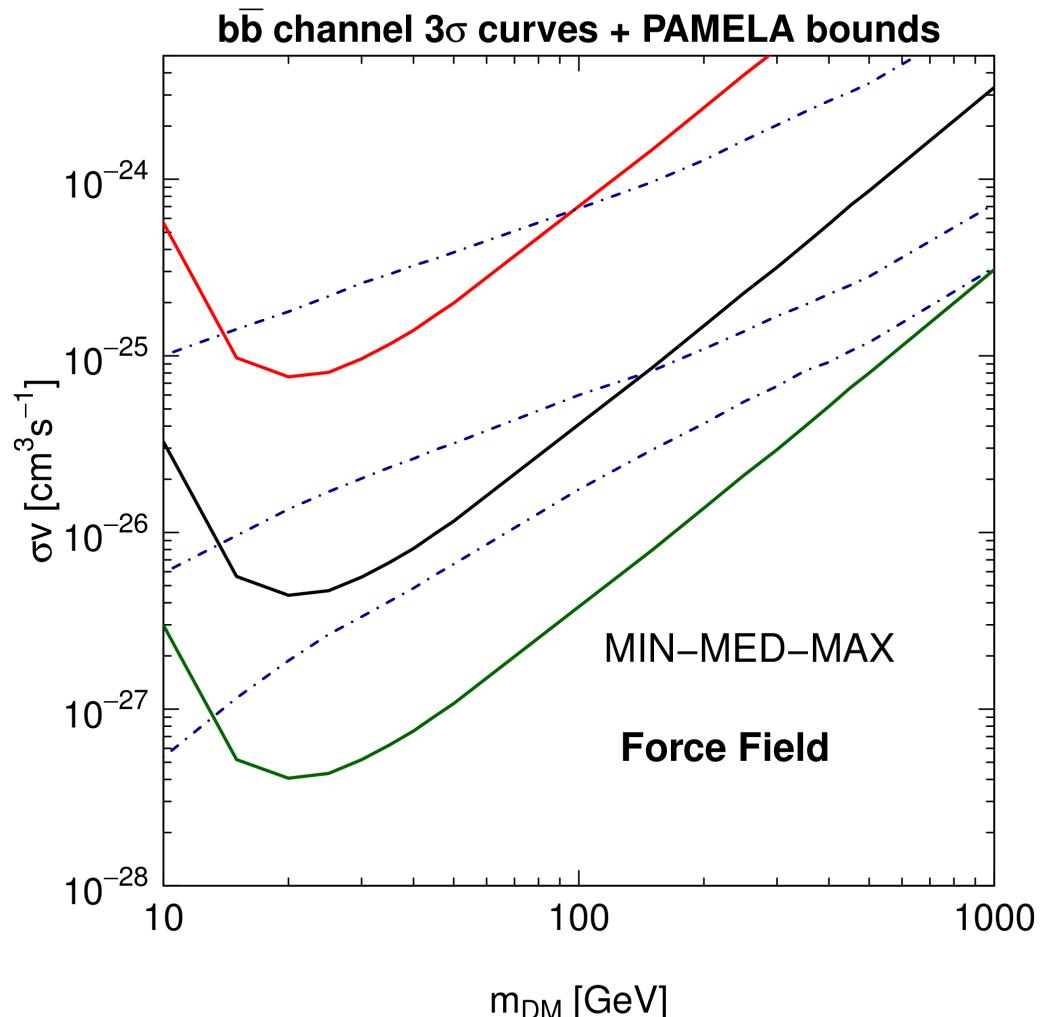


**$W^+W^-$  channel**



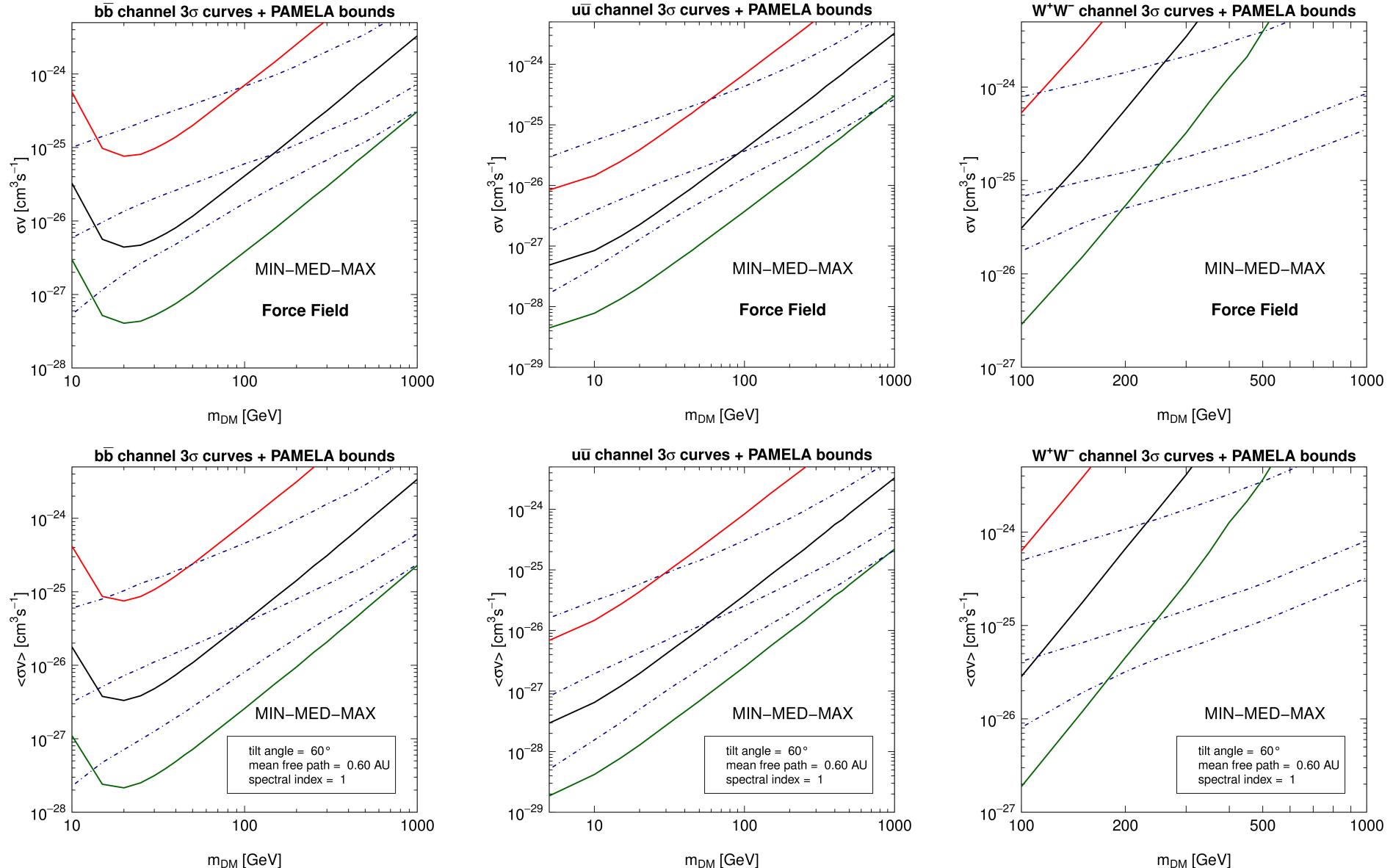
Cross section are ~the maximal permitted by antiprotons bounds

# Experimental reachabilities (GAPS) - I

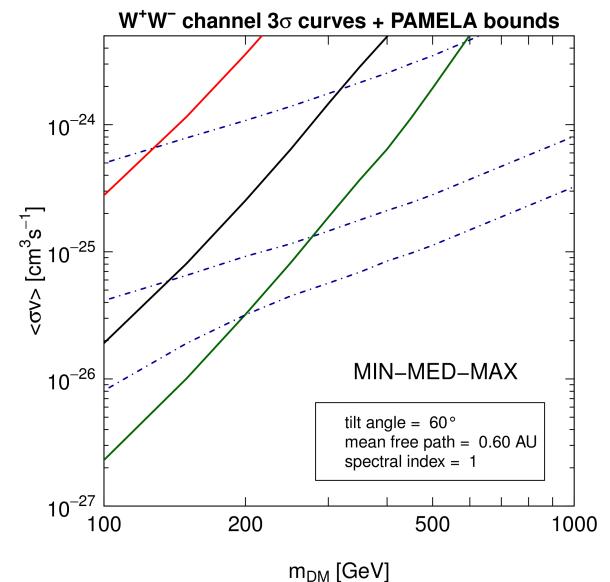
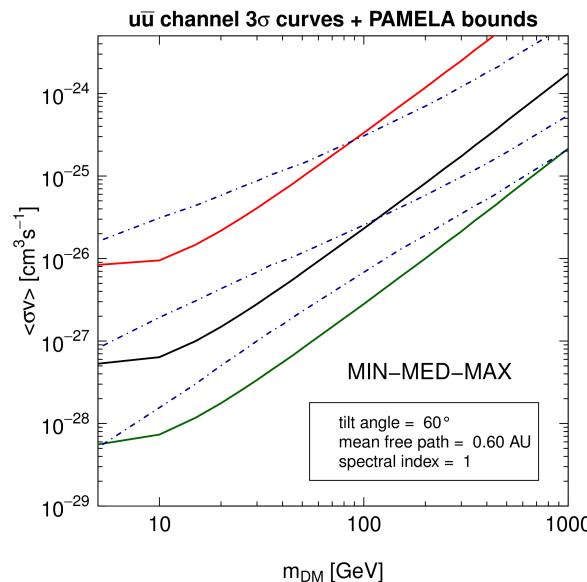
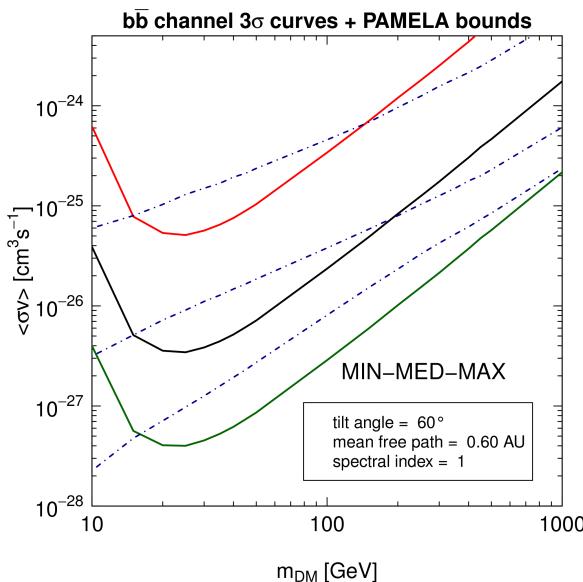
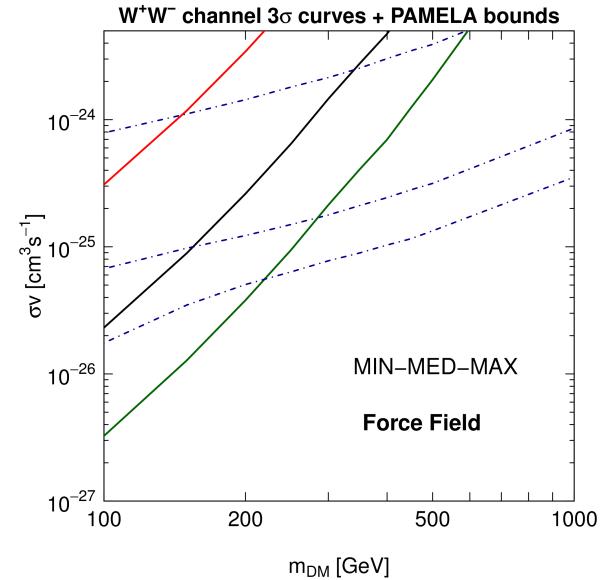
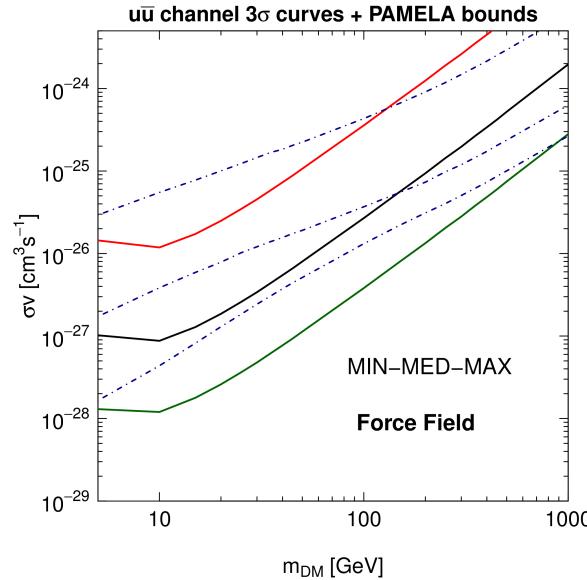
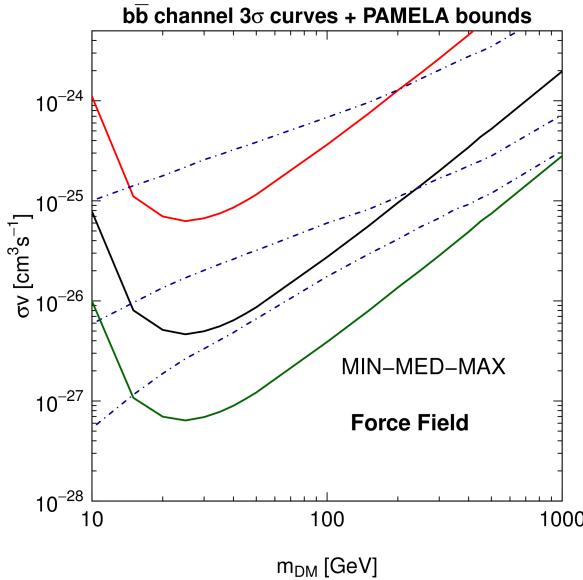


- ▶ Reachability curves for a  $3\sigma$  c.l. (i.e. 1 signal + bkg event)
- ▶ 3 set of curves for 3 propagation models (MIN, MED and MAX)
- ▶ Only PAMELA bounds are considered

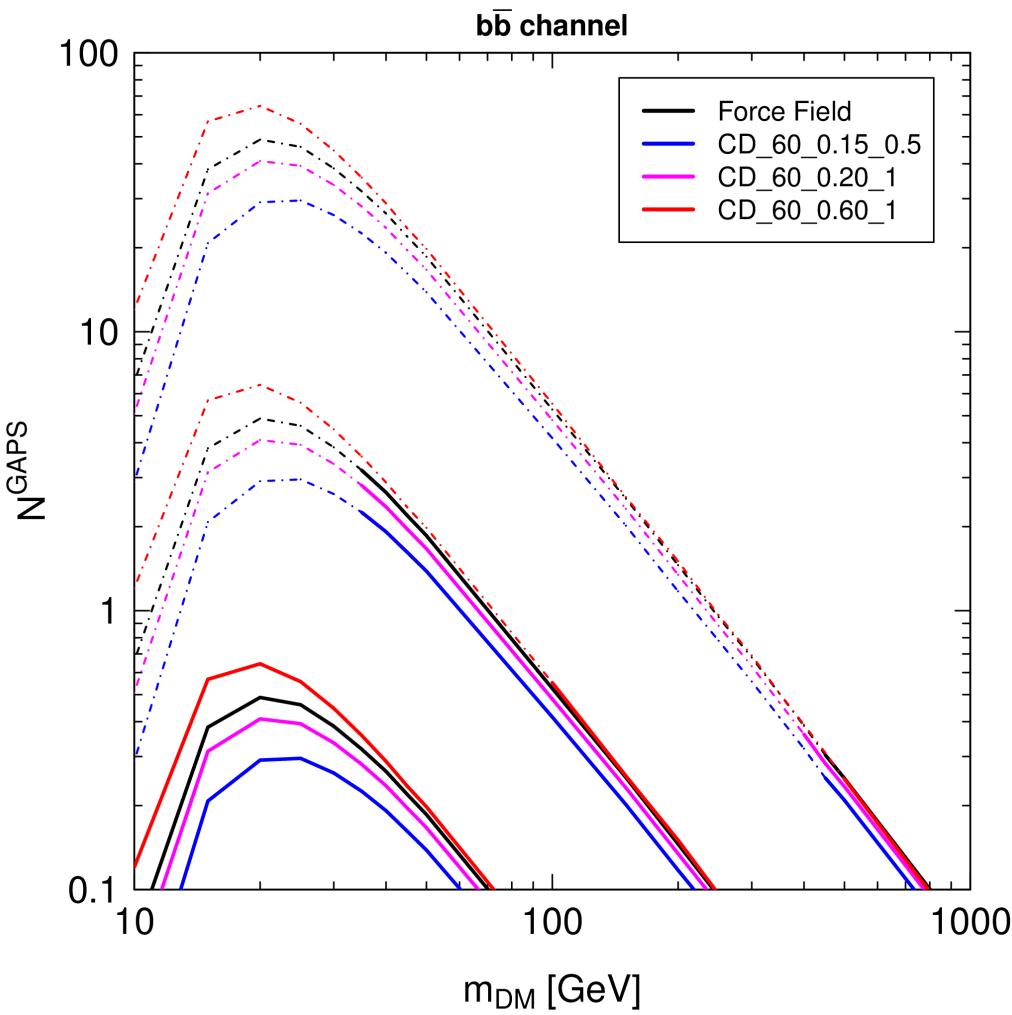
# Experimental reachabilities (GAPS) - II



# Experimental reachabilities (AMS)

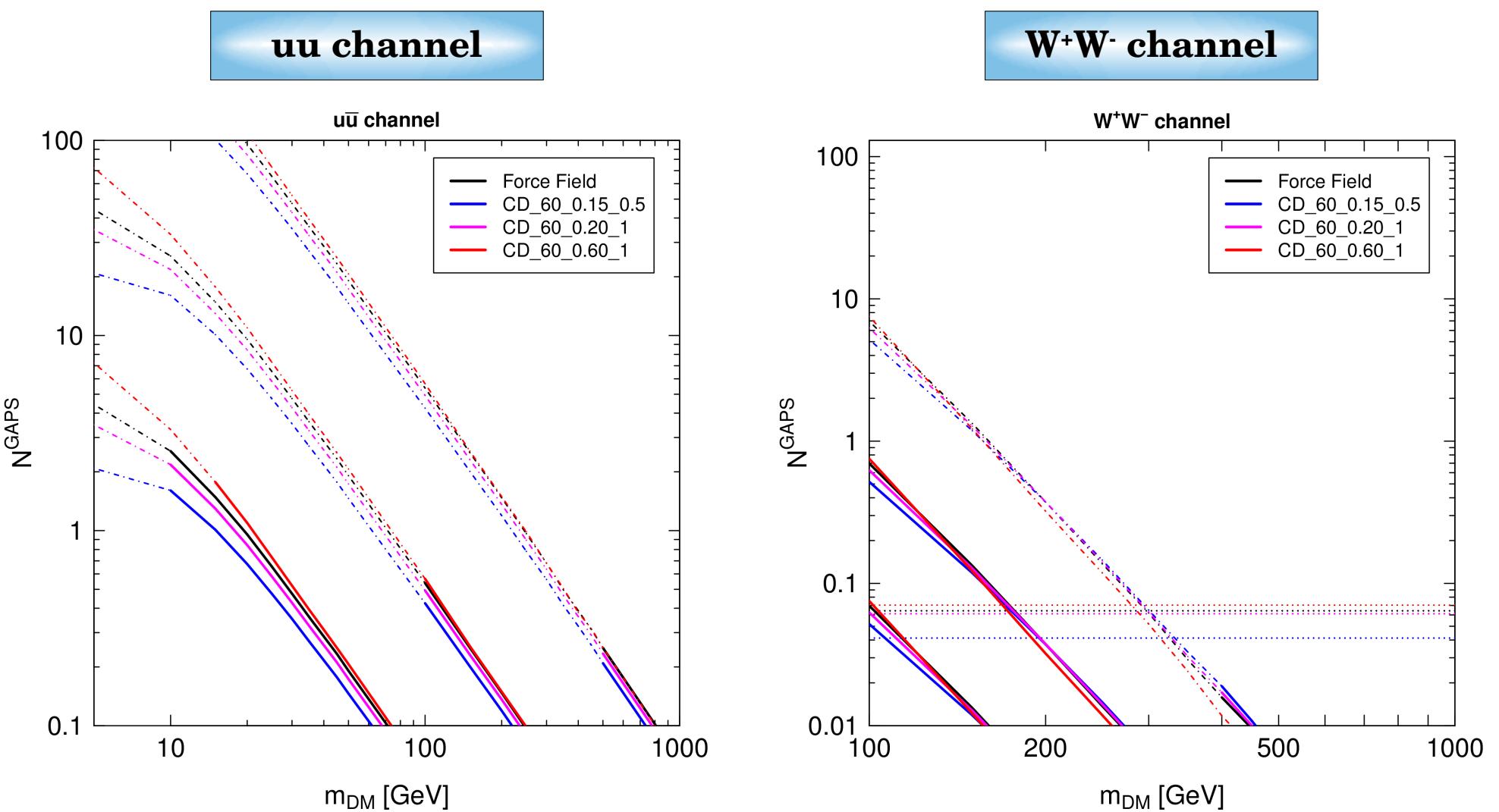


# How many events?



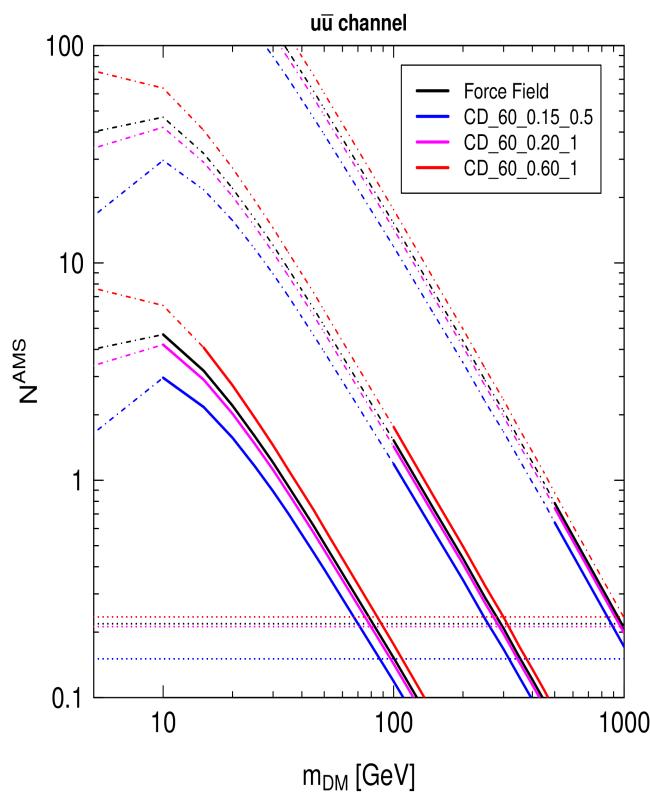
- ▶ Number of signal events for 4 models of solar modulation (Einasto MED configuration)
- ▶ 3 set of curves for 3 values of  $\langle \sigma v \rangle$  (thermal, 0.1 x thermal, 10 x thermal)
- ▶ Solid lines = **allowed** by PAMELA bounds  
Dot-dashed lines = **forbidden** by PAMELA bounds

# How many events? - GAPS

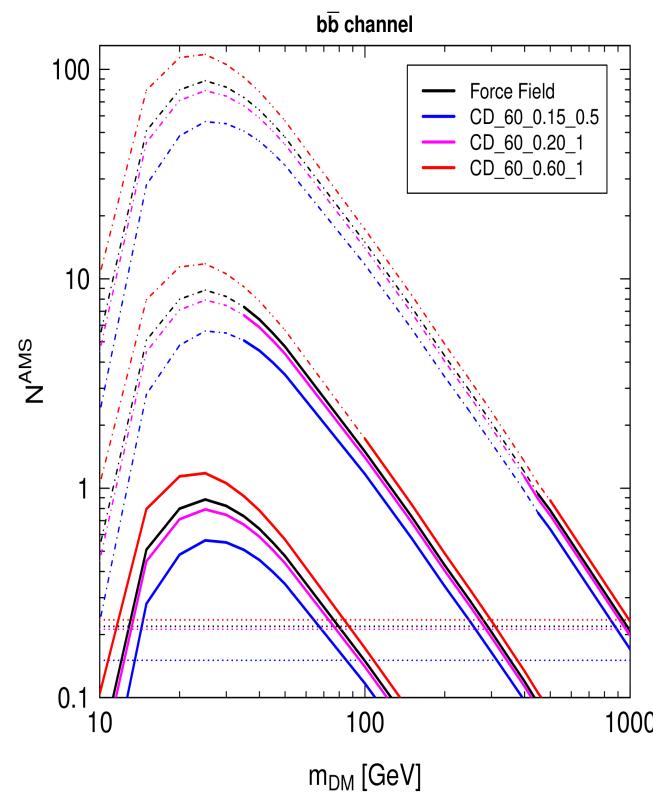


# How many events? - AMS

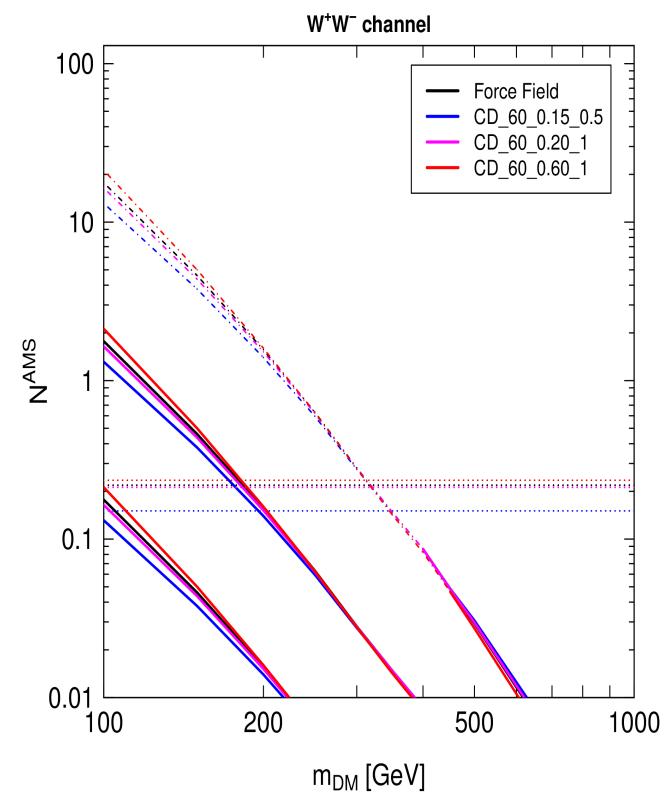
**uu channel**



**bb channel**



**$W^+W^-$  channel**



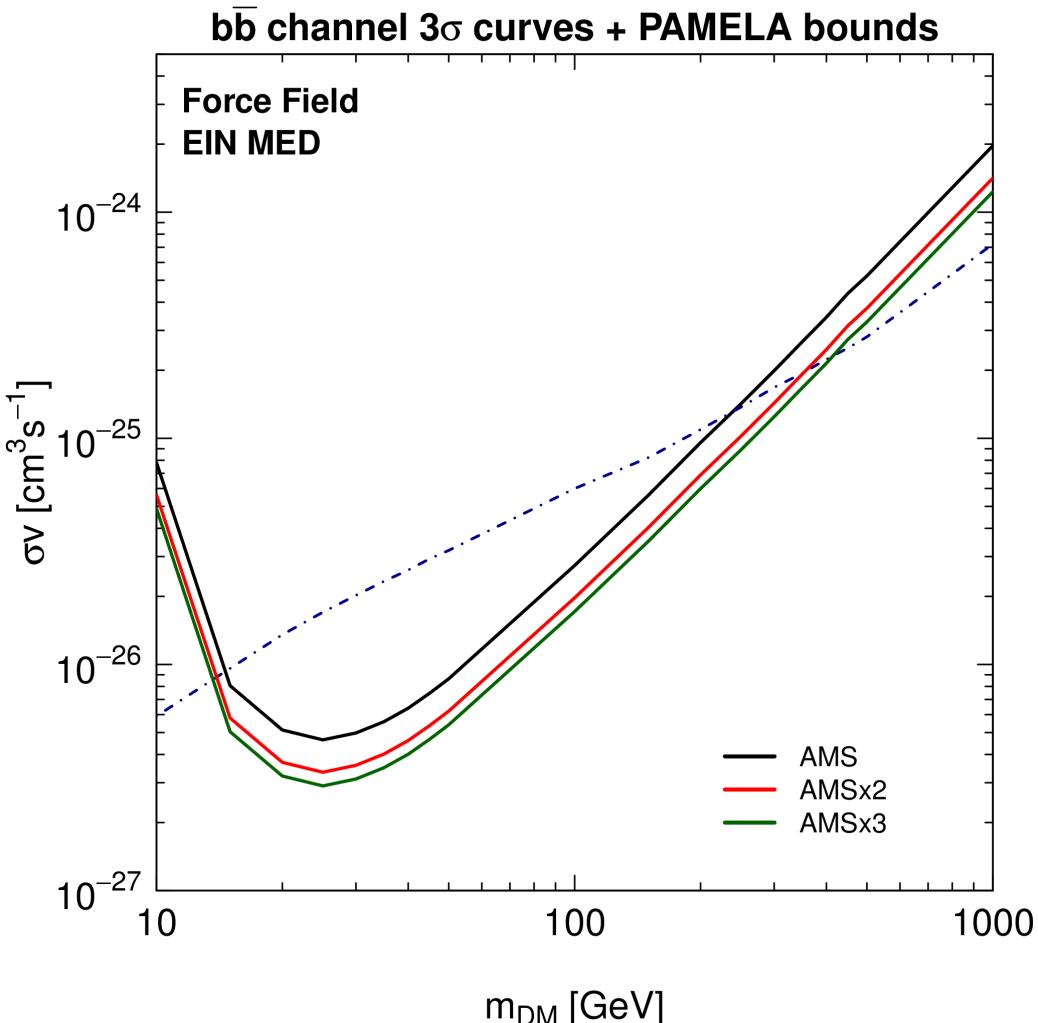
# Conclusions

- ▶ Antideuterons are a very promising, almost background free channel for DM indirect detection
- ▶ In order to have a precise estimate of experimental reachabilities an accurate treatment of solar modulation is needed
- ▶ The bounds imposed by antiprotons experiments such as PAMELA and BESS POLAR II (and, in future, also AMS) are extremely strong. However, some space for DM is still there

**To sum up, antiprotons bounds make the observation of antideuterons at current and future experiments challenging but not impossible**

# **Back-up slides**

# Keeping track of data taking time



- ▶ AMS will fly and take data for a period much longer than the one that we are considering (3 years)
- ▶ Remember however that in a longer time, more background events will be detected -> we'll need more DM events to have a detection!

# CD solar modulation - I

In **spherical coordinates**, the SMF is given by:

$$\vec{B} = AB_0 \left( \frac{r}{r_0} \right)^{-2} \left( \hat{r} - \frac{\Omega r \sin \theta}{V_{SW}} \hat{\varphi} \right)$$

$$\begin{cases} V_{SW} = 400 \text{ Km/s} \\ |B| (1 \text{ AU}) = B_\oplus = 5 \text{ nT} \\ A = H(\theta - \theta') \\ \theta' = \frac{\pi}{2} + \sin^{-1} \left( \sin \alpha \sin \left( \varphi + \frac{\Omega r}{V_{SW}} \right) \right) \end{cases}$$

tilt angle  $\rightarrow [20^\circ, 60^\circ]$

$\Omega$   $\rightarrow$  solar differential rotation rate

The function A takes into account  
the presence of the HCS  
which is related to the tilt angle

The SMF influences also the drifts:  $\vec{v}_d = \nabla \times (K_A \vec{B} / |B|)$

Antisymmetric part of K

# CD solar modulation - II

The solar diffusion depends on the symmetric part of the K tensor:

$K = \text{diag}(K_{\parallel}, K_{\perp r}, K_{\perp \theta})$ . The component parallel to the SMF direction goes like:

$$\begin{cases} k_{\parallel} = K_{\parallel}(\vec{B}) \times \left( \frac{\rho}{1 \text{ GV}} \right)^{\delta} & \text{if } \rho < 4 \text{ GV} \\ k_{\parallel} = K_{\parallel}(\vec{B}) \times \left( \frac{\rho}{1 \text{ GV}} \right)^{1.95} & \text{if } \rho > 4 \text{ GV} \end{cases}$$

[0.5 or 1]

As a function of the mean free path, we have  $K_{\parallel, \perp} = \lambda_{\parallel, \perp} \frac{v}{3}$

$$\lambda_{\parallel} = \lambda_0 \times \lambda_{\parallel}(\vec{B}, \rho)$$

[0.15, 0.20 or 0.60 AU]

# How many events do we need?

We want to determine how many events we need in order to claim for a DM discovery with a certain **confidence level**

The result is given by poissonian statistics:

We need  $N_{\text{crit}}$  events, being  $N_{\text{crit}}$  the smallest  $N$  satisfying (for a  $3\sigma$  c.l.):

$$\sum_{n=0}^{N-1} P(n, b) > 0.997$$

$$P(n, b) = \frac{b^n e^{-n}}{n!}$$

It's a cumulative distribution (the poissonian is discrete!)

Basically, for a  $3\sigma$  c.l. we need 1 event for GAPS and 2 for AMS

# Number of expected events

The number of expected events for a given experiment is calculated as:

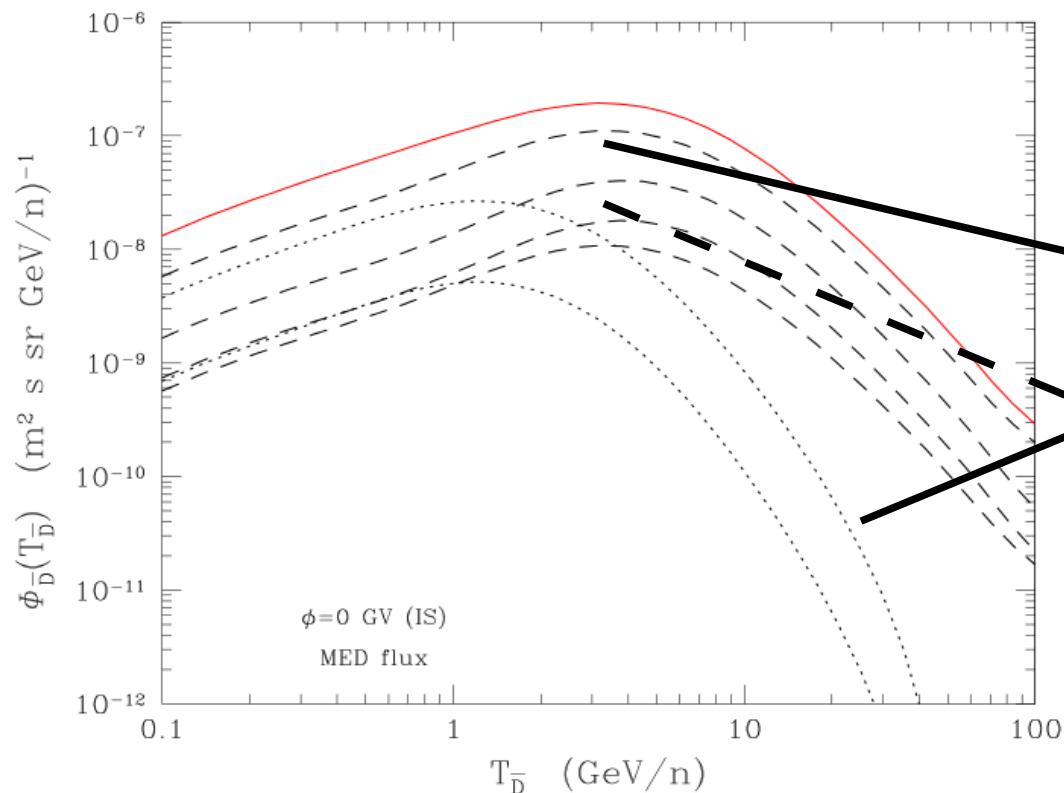
$$N_{\bar{d}} = \int_{E_{k,min}}^{E_{k,max}} dE_k A(E_k) \times \Phi_{TOA}(E_k)$$

Where  $A(E_k)$  is the acceptance of the experiment under scrutiny. In our case we define an effective exposure:

$$N_{\bar{d}} = \left( \int_{E_{k,min}}^{E_{k,max}} dE_k \Phi_{TOA}(E_k) \right) \times \mathcal{E}$$

This exposure  $\mathcal{E}$  can be determined from the experimental sensitivity (i.e. A flux that corresponds to 1 event for GAPS and 2 for AMS )

# Background flux - IS



Given by spallation of CRs on  
the ISM

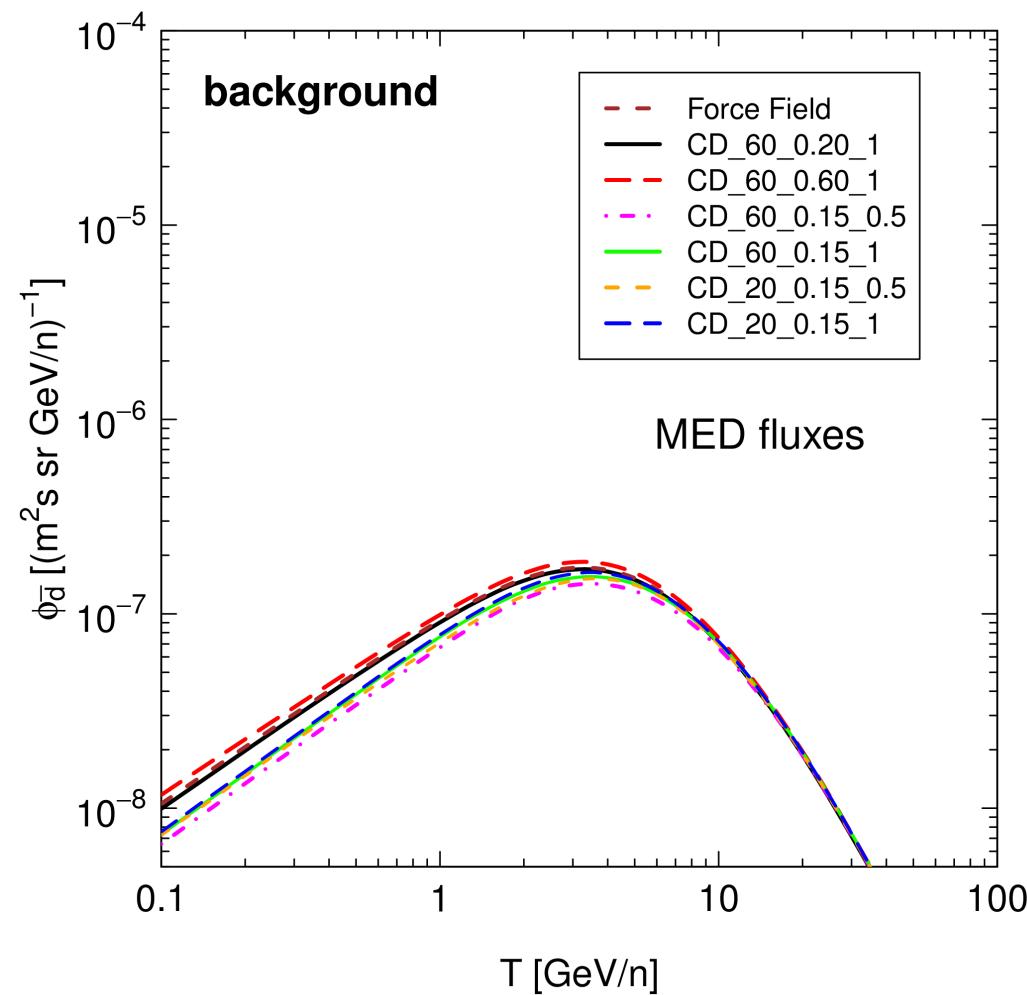
$P/\bar{p}(p) + H$

$P + He$

**F.Donato, N.Fornengo and D.Maurin,  
Phys. Rev. D 78 (2008)**

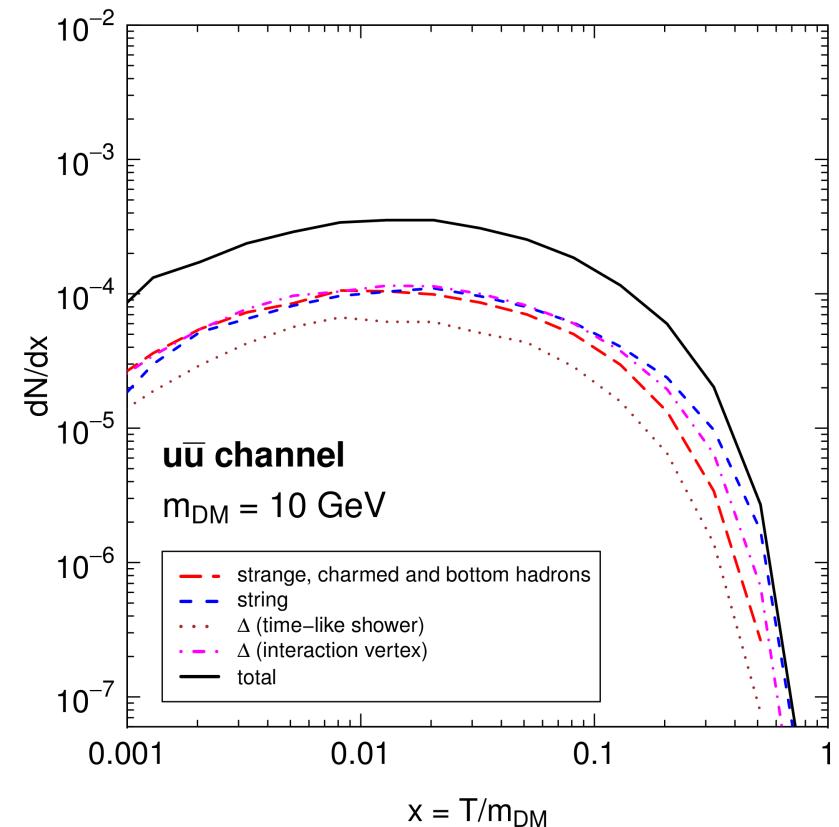
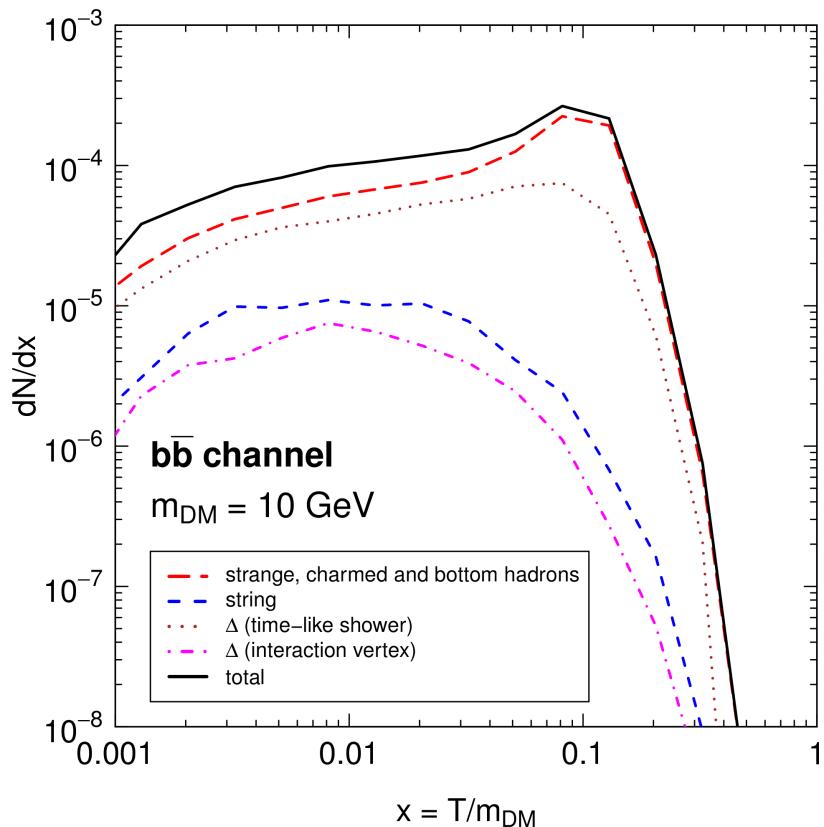
CAVEAT : it is still calculated with the OLD coalescence model

# Background flux - TOA



# Coalescence at low masses

Spectra obtained **without** imposing the condition  $\Delta r < r_0$



For the  $b\bar{b}$  channel, at low masses, the majority of pn pairs come from  
**long living hadrons**