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IN NUCLEAR PHYSICS AND RELATED AREAS



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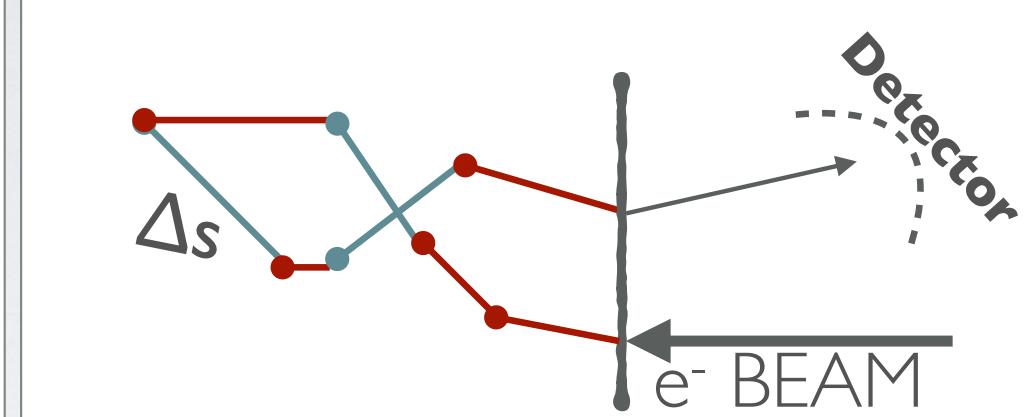
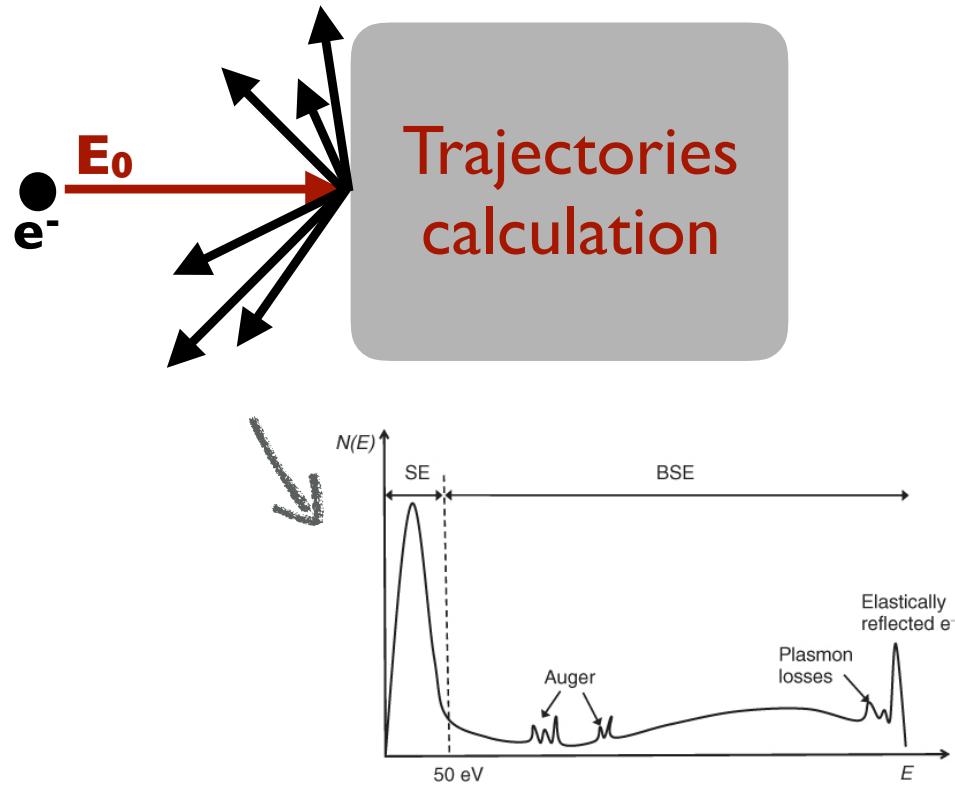
Monte Carlo simulation of Secondary Electron Yield for Noble metals

Martina Azzolini, Nicola M. Pugno, Simone Taioli, Maurizio Dapor

OUTLINE

- ◆ Description of the Monte Carlo computational model to simulate electron transport:
 - ◆ Initial energy distribution
 - ◆ Elastic scattering
 - ◆ Inelastic scattering, ionization and secondary electrons generation
 - ◆ Emission condition
 - ◆ Path length and choice of the interaction kind
- ◆ Comparison between calculated and experimental emission spectra of Cu and secondary electron yield of Cu, Ag and Au.
- ◆ Conclusion and further perspectives

Monte Carlo simulation of electron transport in solid targets



- ◆ Step between interactions
- ◆ Elastic scattering: direction change
- ◆ Inelastic scattering: energy loss, direction change, ionizations and secondary electron generation
- ◆ End of the electron trajectories: kinetic energy lower than threshold, emission of electron (overcome the energy barrier represented by the work function)

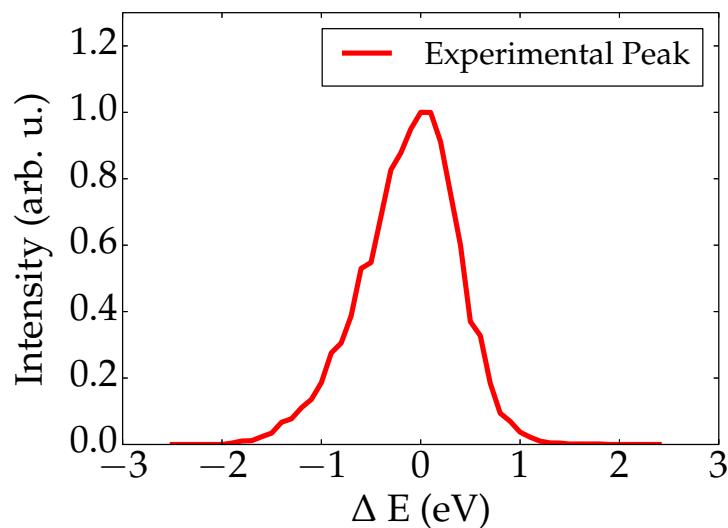
Monte Carlo simulation:

- ◆ Generation of random numbers
- ◆ Calculation of scattering probabilities of interactions
- ◆ Repetition of the calculation

Beam kinetic energy



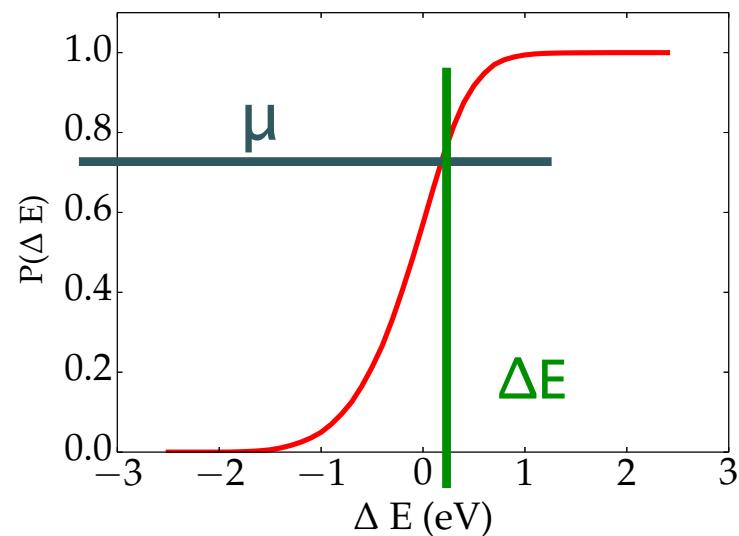
Experimental elastic peak *
(taken on a graphite reference sample)



$$Area = \int_{E_-}^{E_+} f(\Delta E) d(\Delta E)$$

$$\Delta E = ?$$

$$P(\overline{\Delta E}) = \frac{1}{Area} \int_{E_-}^{\overline{\Delta E}} f(\Delta E) d(\Delta E)$$



Generation of a random number μ , uniformly distributed in the interval $[0,1]$

$$\mu = P(\overline{\Delta E})$$

The corresponding value of ΔE is the energy correction to apply

* Experimental data by M. Angelucci, R. Larciprete, R. Cimino

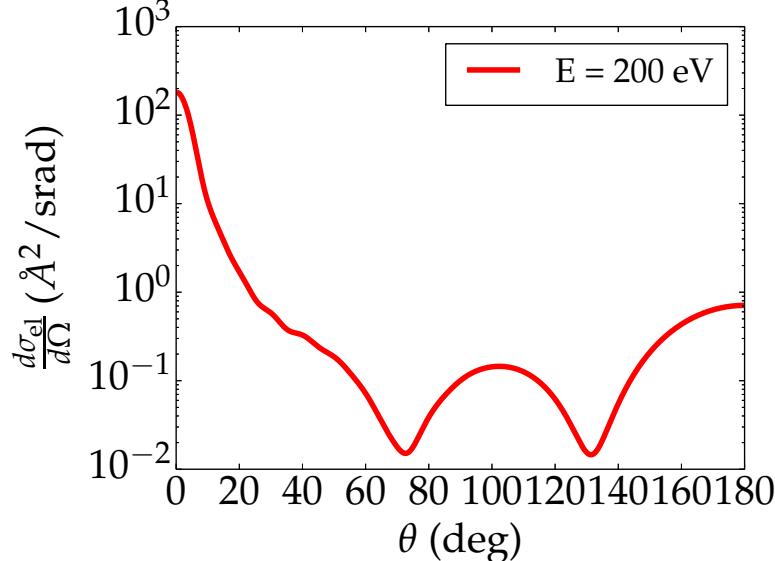
Elastic scattering:

- ◆ Angular deviation

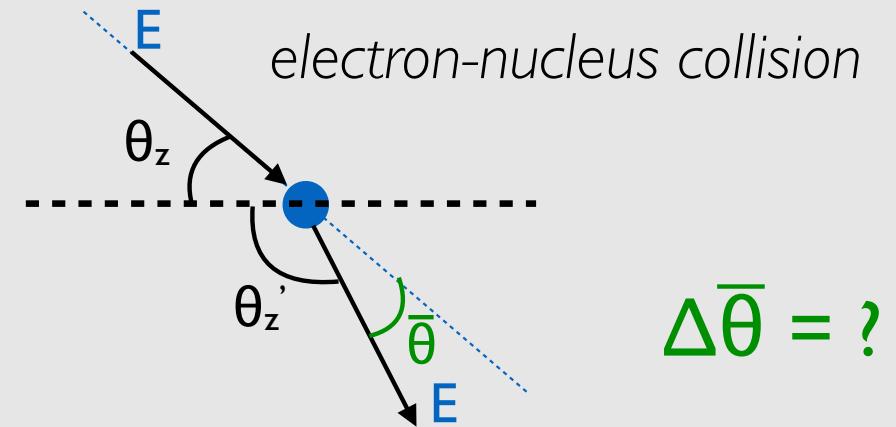
Mott theory with the potential by Salvat.

Differential elastic scattering cross section of copper:

$$\frac{d\sigma_{el}}{d\Omega} = |f|^2 + |g|^2$$

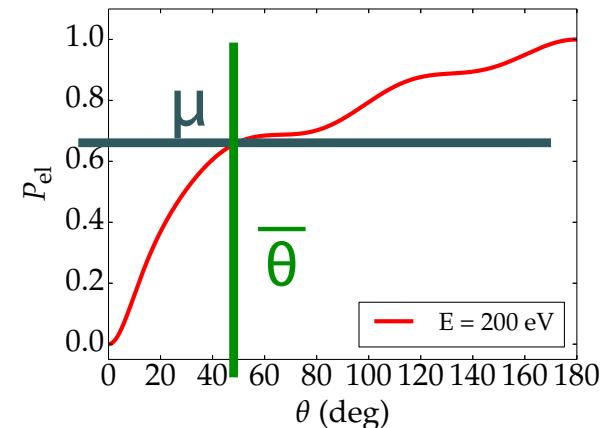


$$\sigma_{el} = 2\pi \int_0^\pi \frac{d\sigma_{el}}{d\theta} \sin(\theta) d\theta$$



Cumulative elastic probability

$$P_{el}(\bar{\theta}, E) = \frac{2\pi}{\sigma_{el}(E)} \int_0^{\bar{\theta}} \frac{d\sigma_{el}(E)}{d\theta} \sin \theta d\theta$$



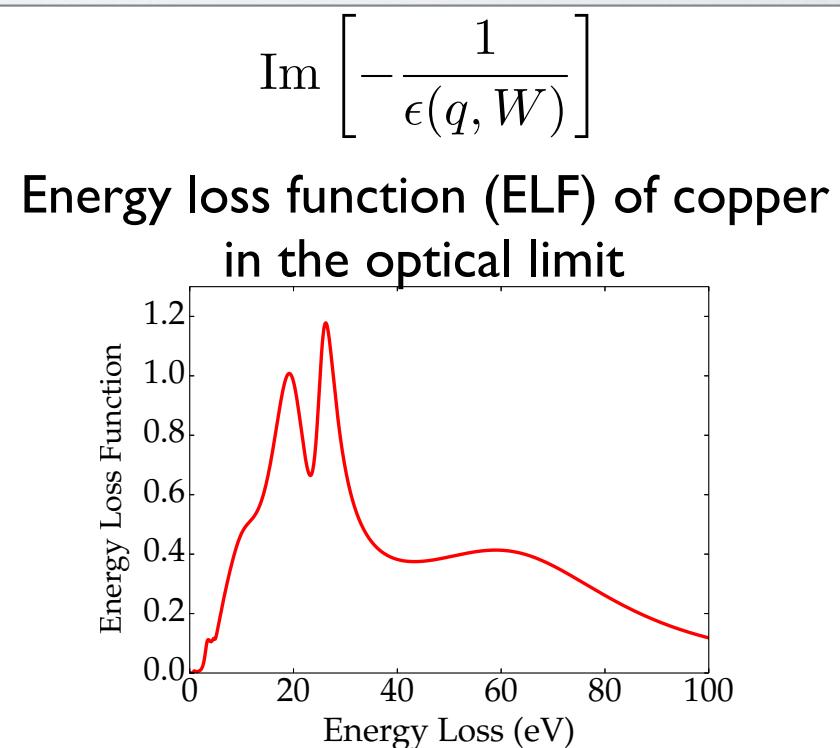
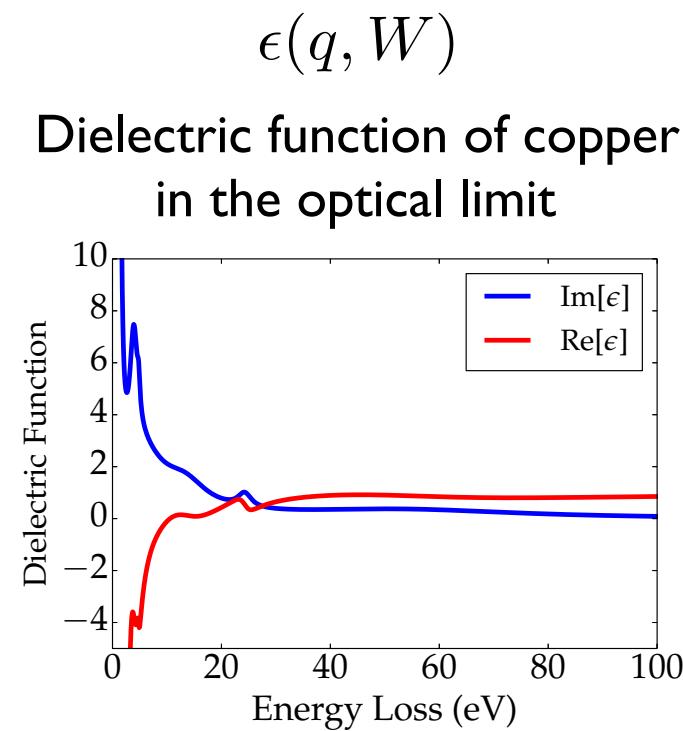
Generation of a random number μ , uniformly distributed in the interval $[0, 1]$

The corresponding value of θ is the trajectory deviation angle

Inelastic scattering

- ◆ Angular deviation
- ◆ Energy loss
- ◆ Ionization + secondary electron generation

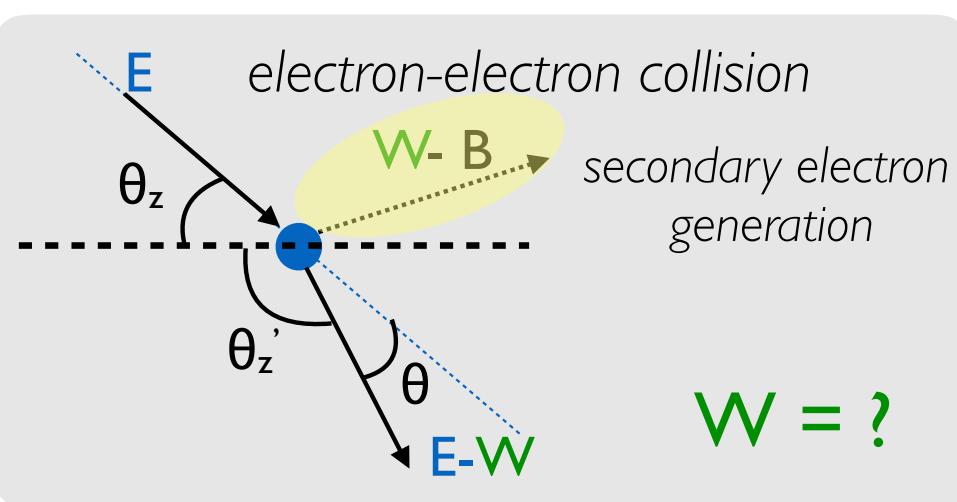
Ritchie dielectric theory and dielectric function obtained by experimental Reflection electron energy loss spectra by Werner et al. It reports the contribution of surface plasmon excitation.



R. H. Ritchie, Plasma losses by fast electrons in thin films, Phys. Rev. 106 (1957) 874.

W. S. Werner, K. Glantschnig, C. Ambrosch-Draxl, Optical constants and inelastic electron-scattering data for 17 elemental metals, Journal of Physical and Chemical Reference Data 38 (4)(2009) 1013–1092.

Inelastic scattering



$$W = ?$$

Differential inelastic scattering cross section

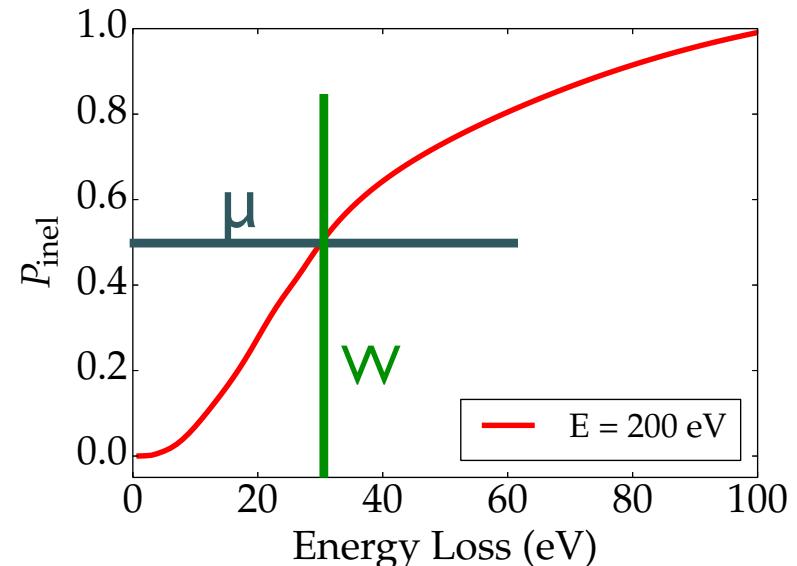
$$\frac{d\sigma_{inel}}{dW} = \frac{1}{N\pi a_0 E} \int_{q_-}^{q_+} Im \left[-\frac{1}{\epsilon(q, W)} \right] \frac{dq}{q}$$

Cumulative inelastic scattering probability:

$$P_{inel}(W, E) = \frac{1}{\sigma_{inel}} \int_0^W \frac{d\sigma_{inel}}{dW'} dW'$$

Generation of a random number μ , uniformly distributed in the interval $[0,1]$.

The corresponding value of W is the energy loss.

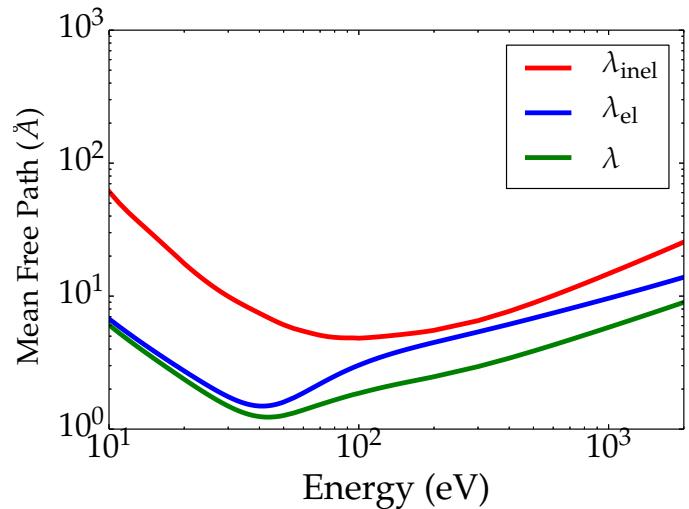


If the energy loss is higher than the first ionization energy (B) of the target atom, a secondary electron is generated. Also its trajectory is calculated.

Step length

λ is the overall mean free path (Cu)

$$\lambda^{-1} = \lambda_{inel}^{-1} + \lambda_{el}^{-1} = N(\sigma_{inel} + \sigma_{el})$$



Δs is the step length between two subsequent collisions

$$\Delta s = -\lambda \ln(\mu)$$

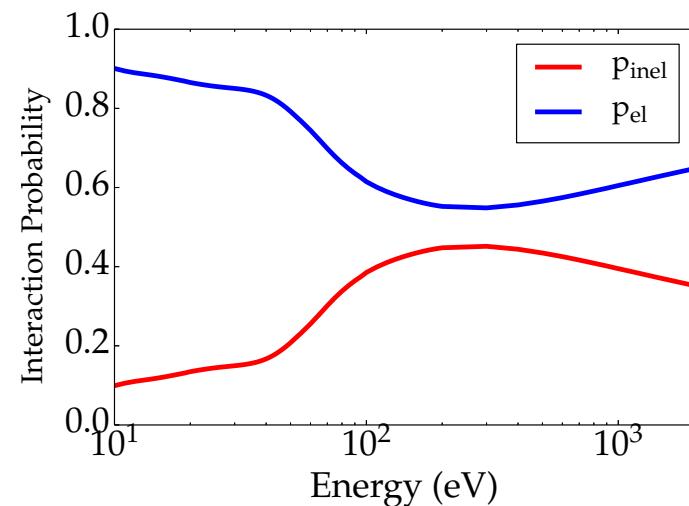
where μ is a random number uniformly distributed in the interval $[0, 1]$

Kind of interaction

Another random number μ (uniformly distributed in the interval $[0, 1]$) is generated. Its value is compared to p_{el} , where:

$$p_{el} = \frac{\lambda}{\lambda_{el}}$$

If $\mu < p_{el}$ the interaction will be elastic, otherwise inelastic.



Electron emission from the surface

The electron has to overcome the energy barrier represented by the work function.

The emission condition to be satisfied is:

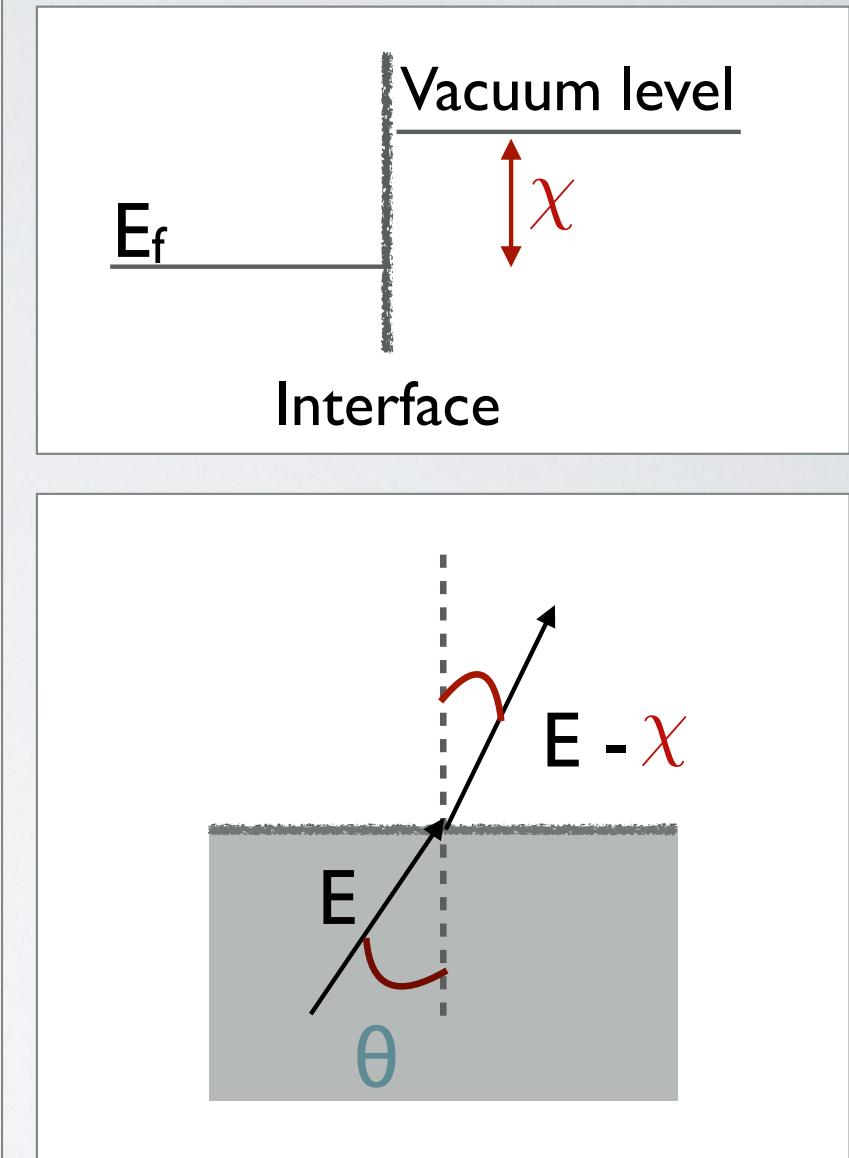
$$E \cos^2 \theta \geq \chi$$

The transmission coefficient is calculated as:

$$T = \frac{4\sqrt{1 - \chi/(E \cos^2 \theta)}}{\left[1 + \sqrt{1 - \chi/(E \cos^2 \theta)}\right]^2}$$

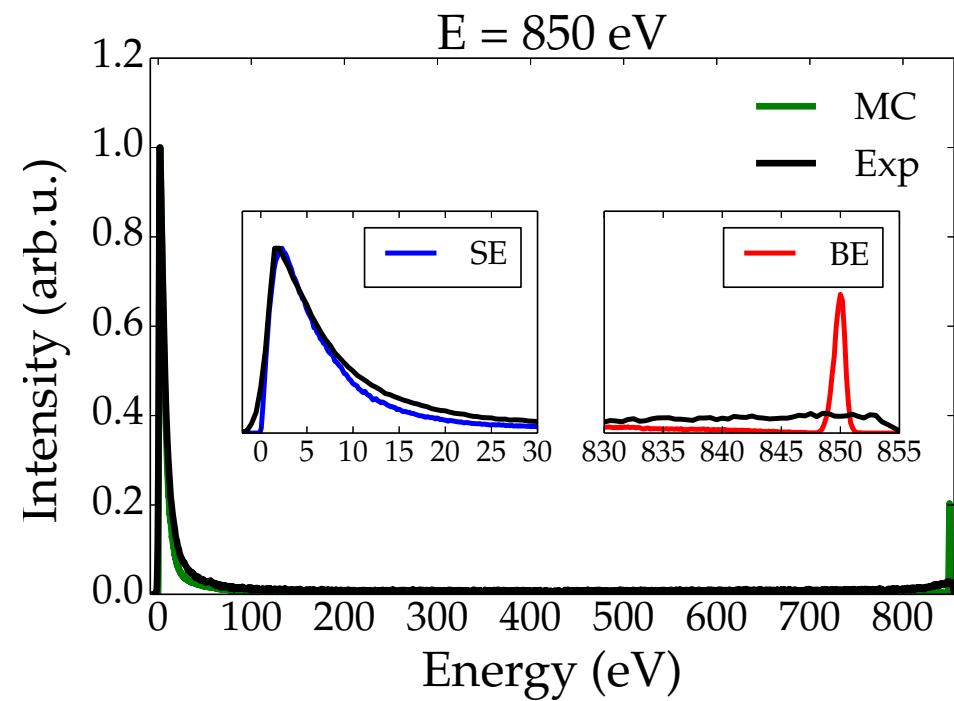
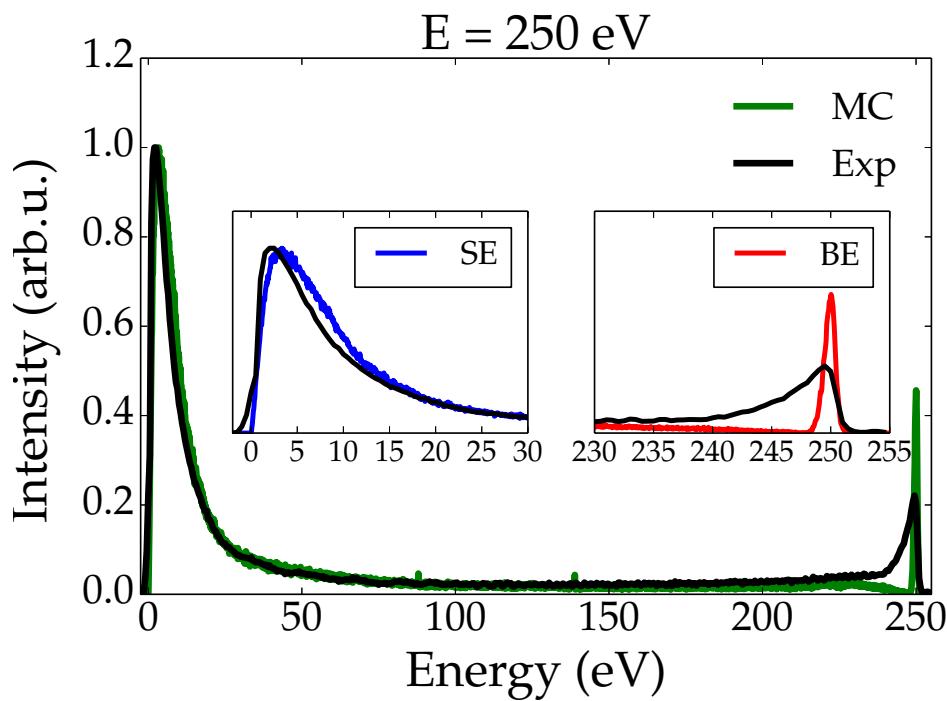
In order to decide if the electron will be emitted, another random number μ (uniformly distributed in the interval $[0,1]$) is generated.

If $\mu < T$ the electron will be emitted and collected and its energy value is stored.



Emission spectra of Cu

10^7 trajectories



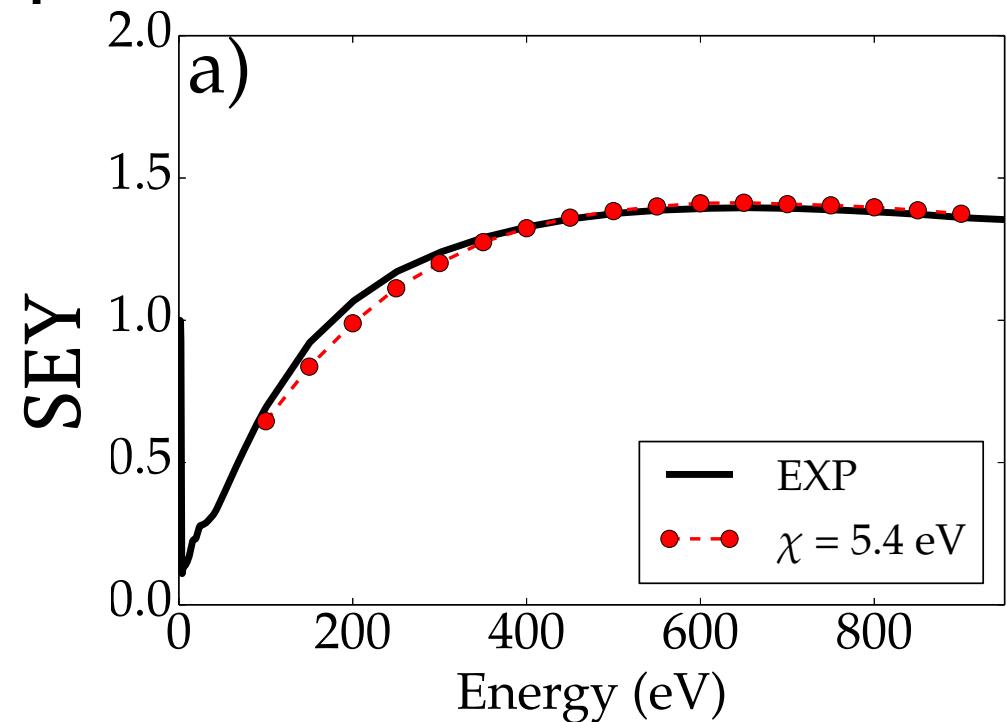
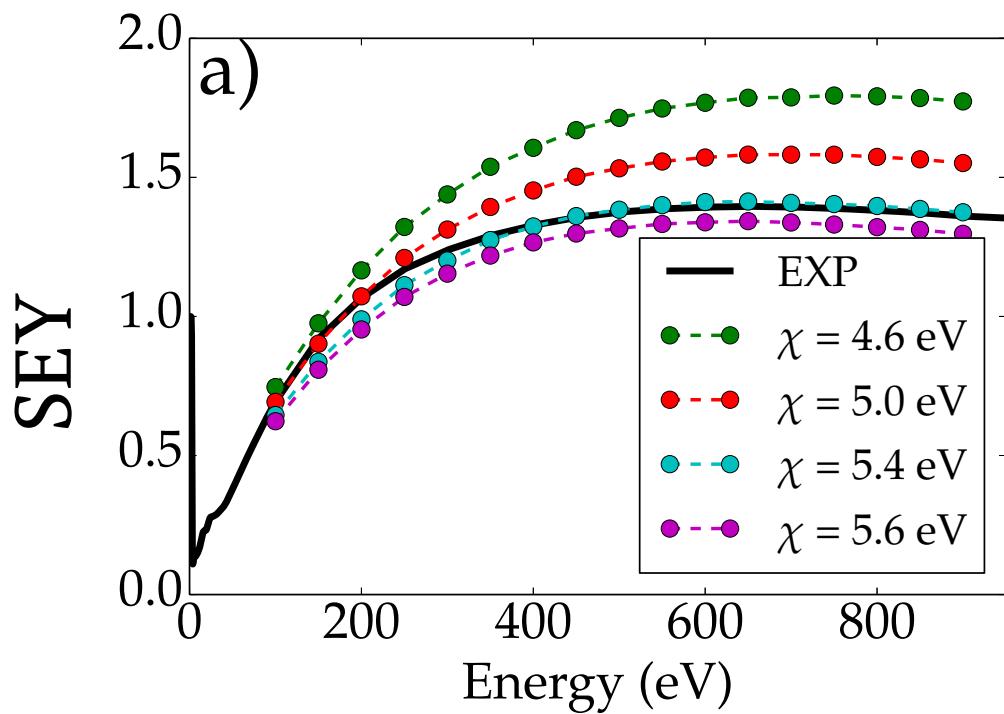
Spectra normalized at a common height of the SE peak.
Preliminary experimental data by M. Angelucci, R. Larciprete, R. Cimino

Secondary electron yield

$$\text{SEY} = \frac{\text{n}^{\circ} \text{ of emitted electrons}}{\text{n}^{\circ} \text{ of total trajectories}}$$

Simulation realized by considering 10^6 electrons in the beam

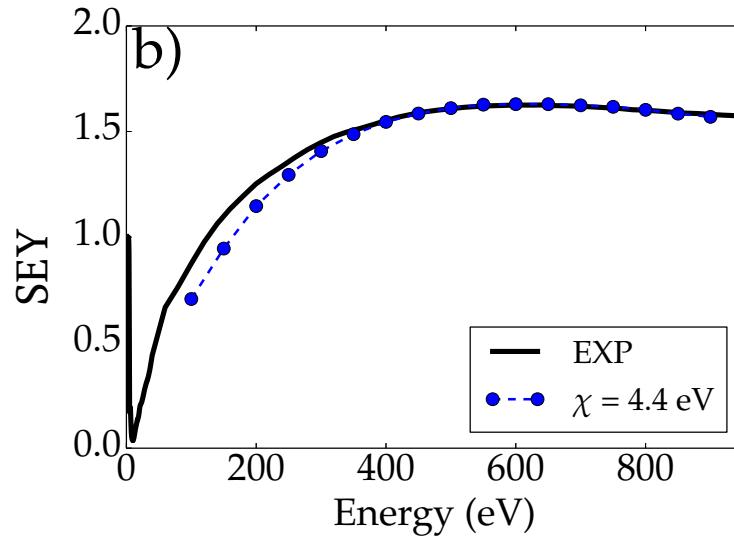
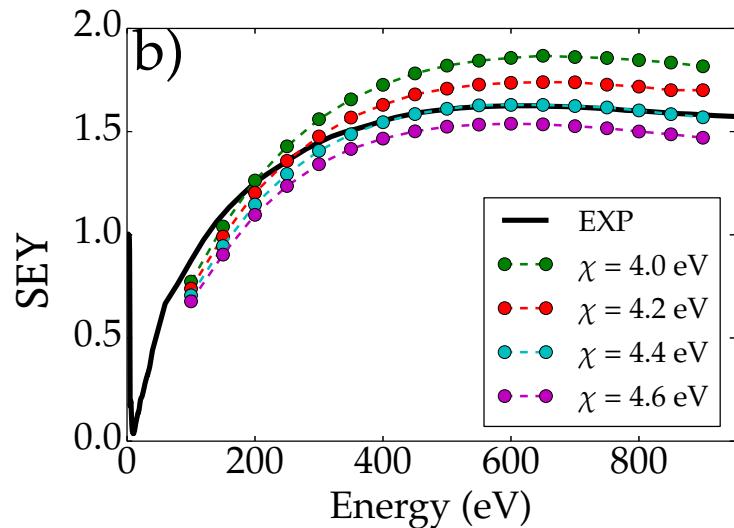
Copper



Experimental $\chi = 4.6 \text{ eV}$

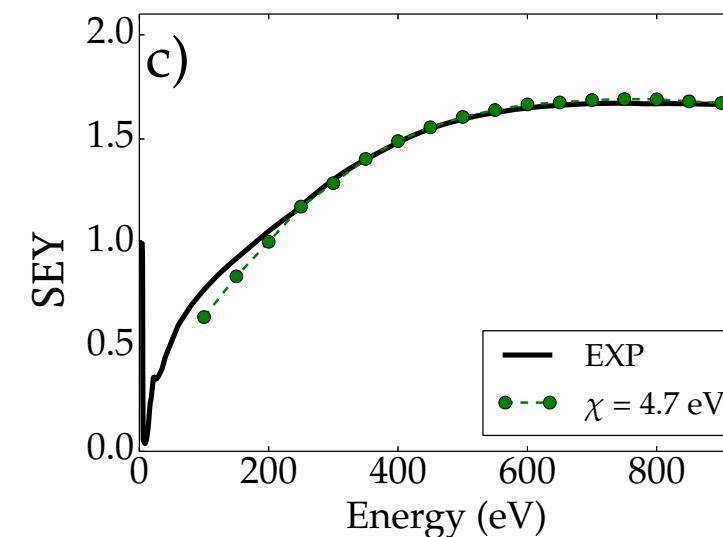
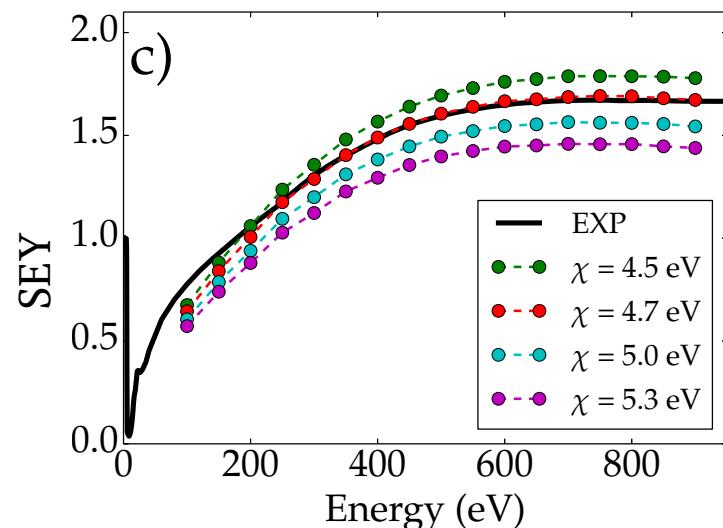
Secondary electron yield

Silver



Experimental $\chi = 4.4 \text{ eV}$

Gold



Experimental $\chi = 5.3 \text{ eV}$

Conclusions

- ◆ We developed a MC code to calculate electron trajectories in metal targets.
- ◆ The ionization and relative secondary electron generations are taken into account.
- ◆ Good agreement between calculations and experimental data of emission spectra and electron yield is obtained.

Next step

- ◆ Test the code by using energy loss function and secondary electron yield realized on the same sample.

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