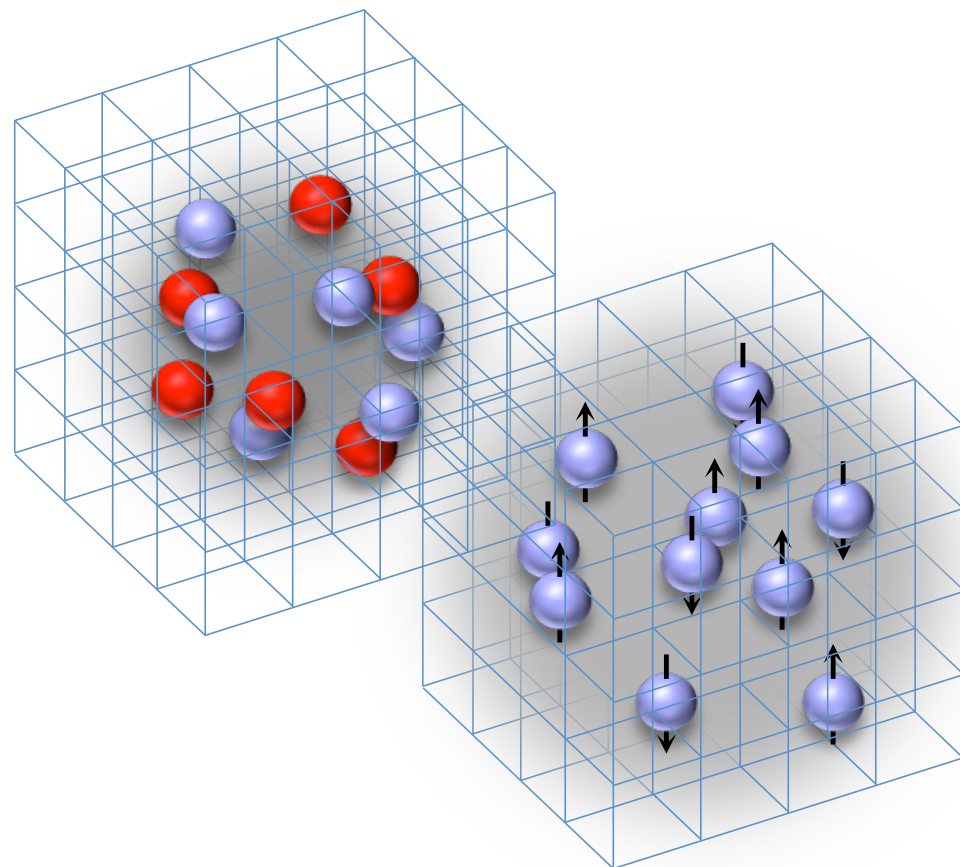




Scattering cluster wave functions on the lattice using the adiabatic projection method



Serdar Elhatisari - Bonn
Michelle Pine - NC State
Dean Lee - NC State
Evgeny Epelbaum - RUB
Hermann Krebs - RUB



Outline

- Introduction
- Adiabatic projection method
- Asymptotic cluster wave functions
- One and three dimensional examples
- Summary



- Objective: **ab initio** calculation of scattering and reactions involving two **clusters**. Processes with alpha-clusters are involved in stellar nucleosynthesis.

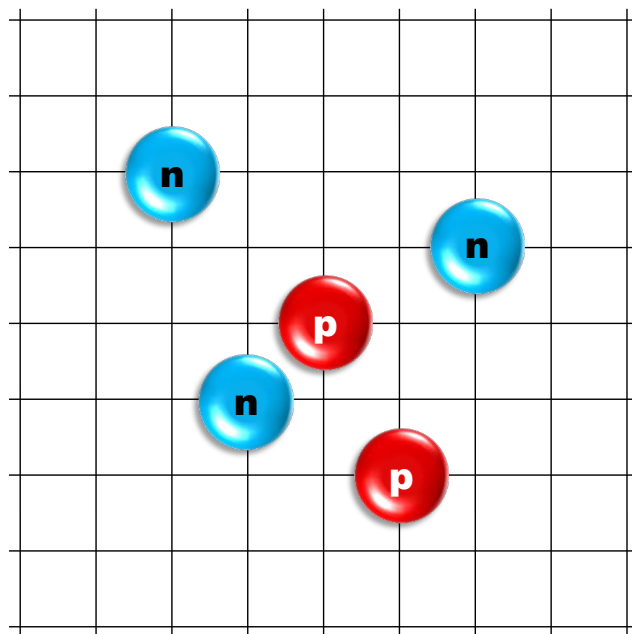
- Objective: **ab initio** calculation of scattering and reactions involving two **clusters**. Processes with alpha-clusters are involved in stellar nucleosynthesis.
 - Example: ${}^4\text{He} + {}^4\text{He} \rightarrow {}^4\text{He} + {}^4\text{He}$ (see next talk)



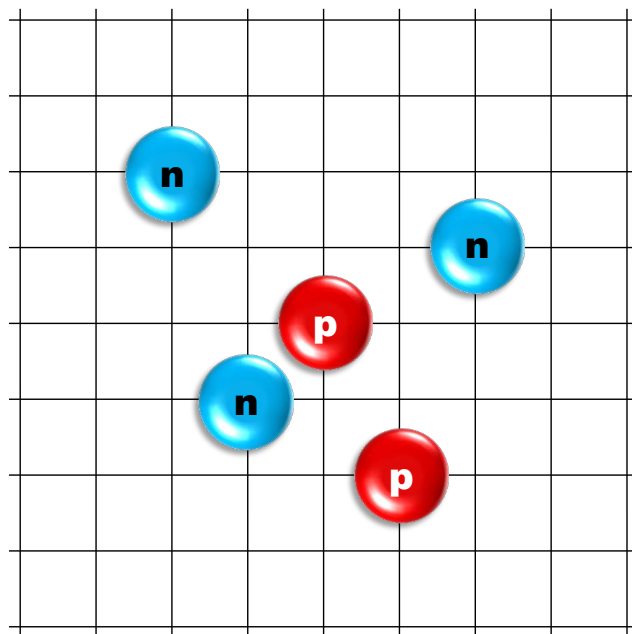
- Objective: **ab initio** calculation of scattering and reactions involving two **clusters**. Processes with alpha-clusters are involved in stellar nucleosynthesis.
 - Example: ${}^4\text{He} + {}^4\text{He} \rightarrow {}^4\text{He} + {}^4\text{He}$ (see next talk)
 - **Lattice effective field theory**

- Objective: **ab initio** calculation of scattering and reactions involving two **clusters**. Processes with alpha-clusters are involved in stellar nucleosynthesis.
 - Example: ${}^4\text{He} + {}^4\text{He} \rightarrow {}^4\text{He} + {}^4\text{He}$ (see next talk)
 - **Lattice effective field theory**

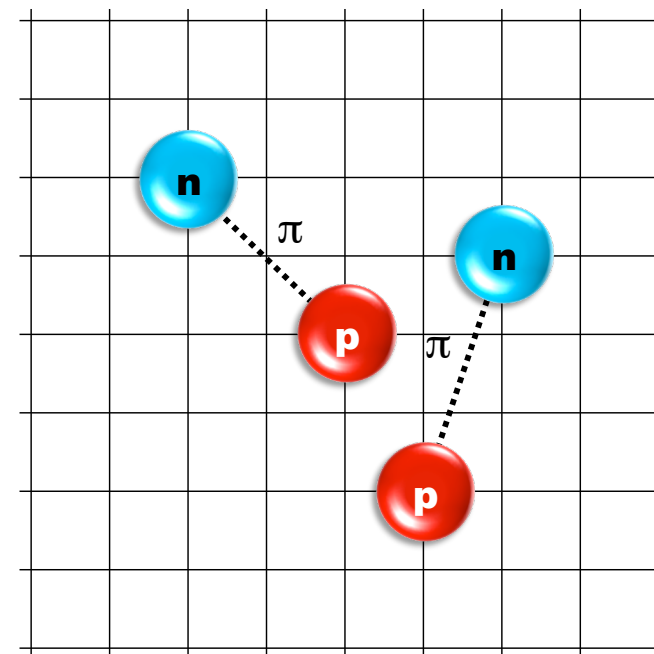
LEFT(π)



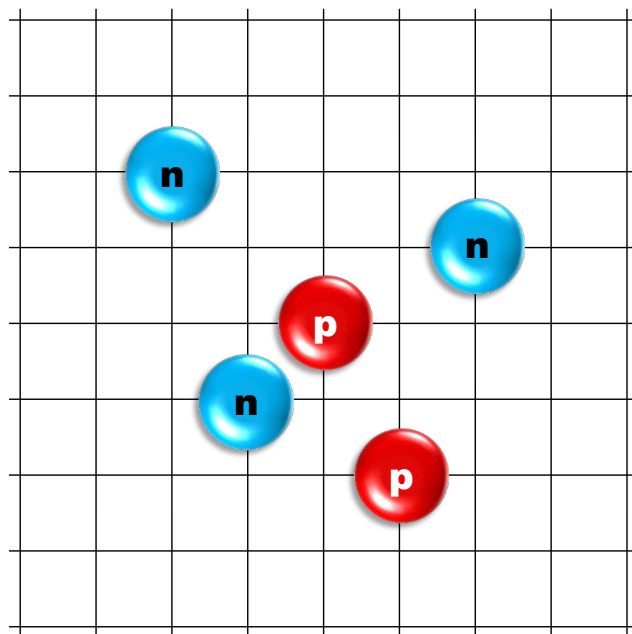
- Objective: **ab initio** calculation of scattering and reactions involving two **clusters**. Processes with alpha-clusters are involved in stellar nucleosynthesis.
 - Example: ${}^4\text{He} + {}^4\text{He} \rightarrow {}^4\text{He} + {}^4\text{He}$ (see next talk)
 - **Lattice effective field theory**

LEFT(π)

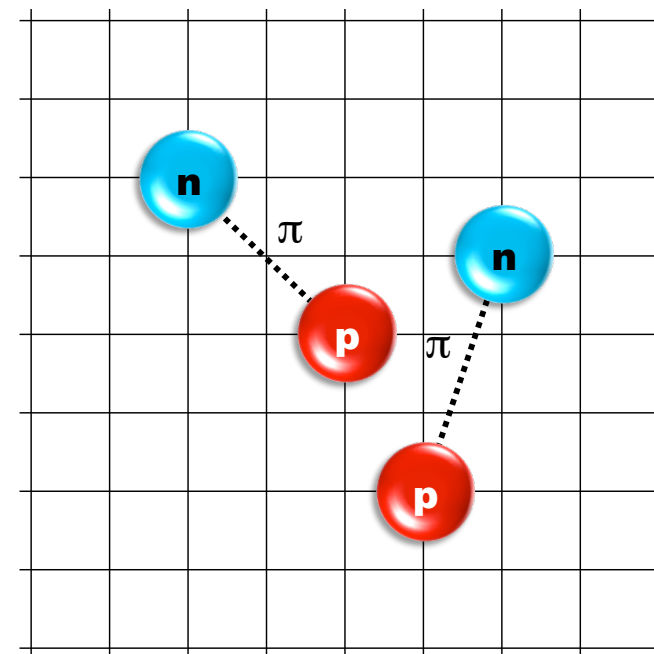
chirale LEFT



- Objective: **ab initio** calculation of scattering and reactions involving two **clusters**. Processes with alpha-clusters are involved in stellar nucleosynthesis.
 - Example: ${}^4\text{He} + {}^4\text{He} \rightarrow {}^4\text{He} + {}^4\text{He}$ (see next talk)
 - **Lattice effective field theory**
 - Direct calculation too computationally demanding: $L^{d(A-1)} \times L^{d(A-1)}$

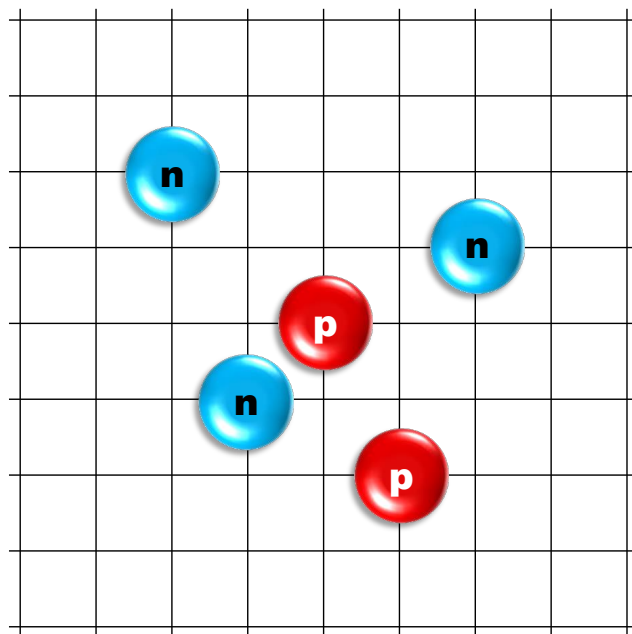
LEFT(π)

chirale LEFT

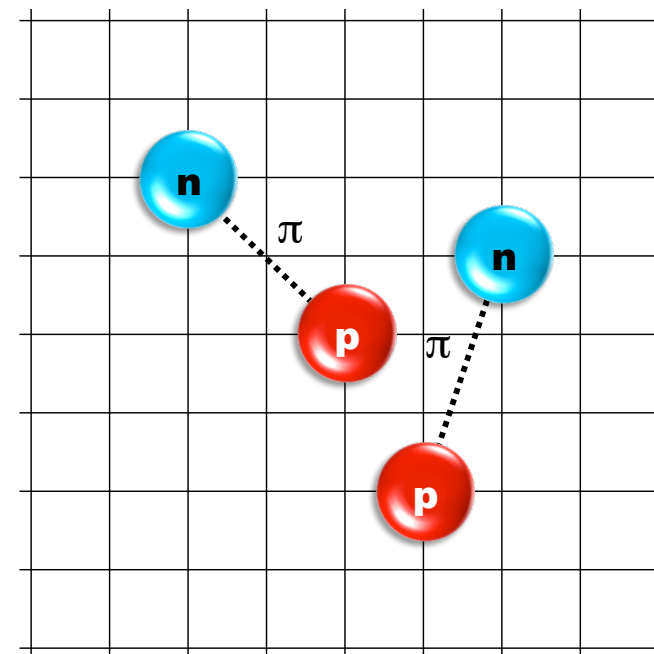


- Objective: **ab initio** calculation of scattering and reactions involving two **clusters**. Processes with alpha-clusters are involved in stellar nucleosynthesis.
- Example: ${}^4\text{He} + {}^4\text{He} \rightarrow {}^4\text{He} + {}^4\text{He}$ (see next talk)
- Lattice effective field theory**
- Direct calculation too computationally demanding: $L^{d(A-1)} \times L^{d(A-1)}$
- Adiabatic projection method**: $L^d \times L^d$

M. Pine, D. Lee and G. Rupak: Eur. Phys. J. A (2013) 49: 151
 G. Rupak and D. Lee, Phys. Rev. Lett. 111, no. 3 (2013), 032502
 S. Elhatisari and D. Lee: Phys. Rev. C 90, no. 6 (2014), 064001

LEFT(π)

chirale LEFT

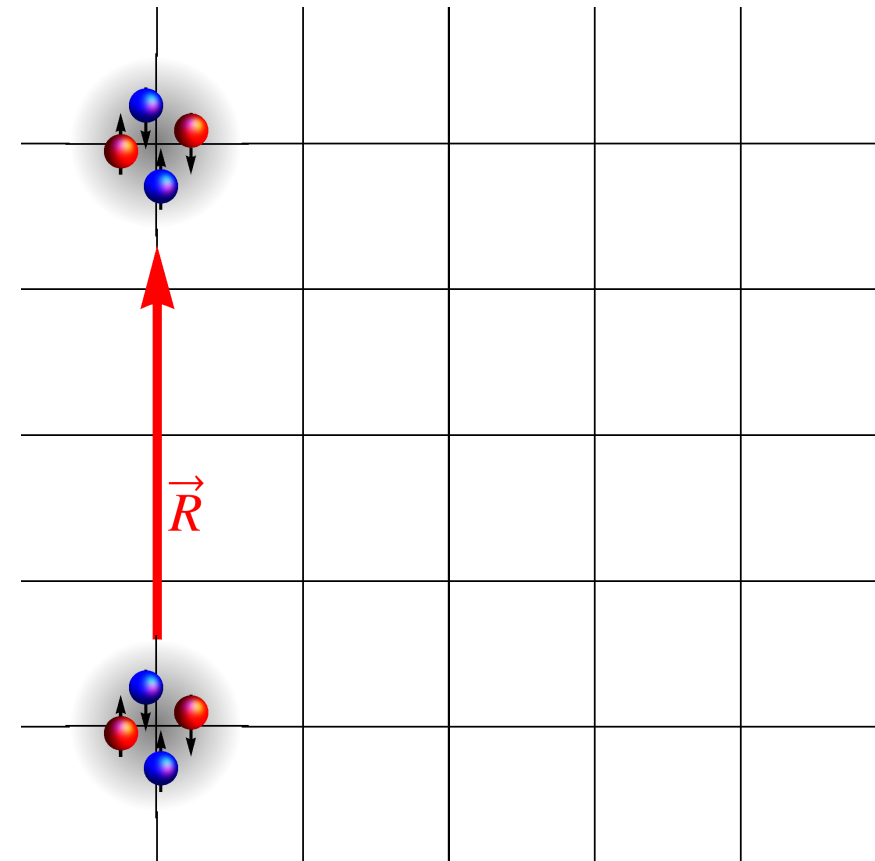




- **First Step:** low-energy cluster Hamiltonian

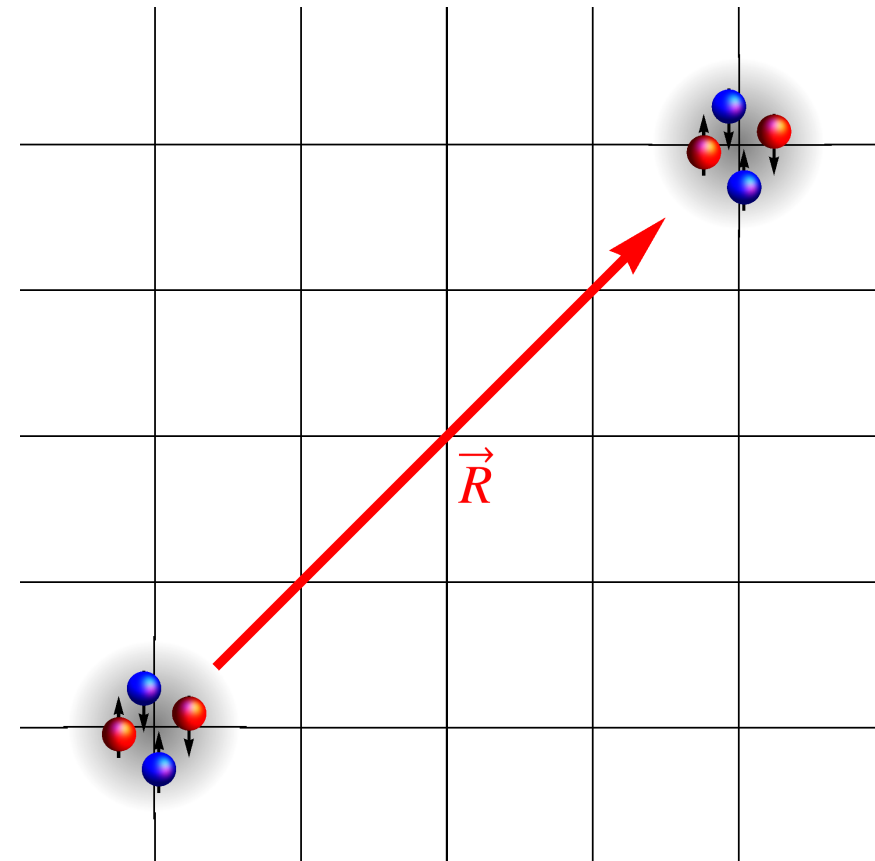
- **First Step:** low-energy cluster Hamiltonian
 - Set of **initial states** labeled by relative cluster separation

$$|\vec{R}\rangle = \sum_{\vec{r}} |\vec{r} + \vec{R}\rangle_1 \otimes |\vec{r}\rangle_2$$



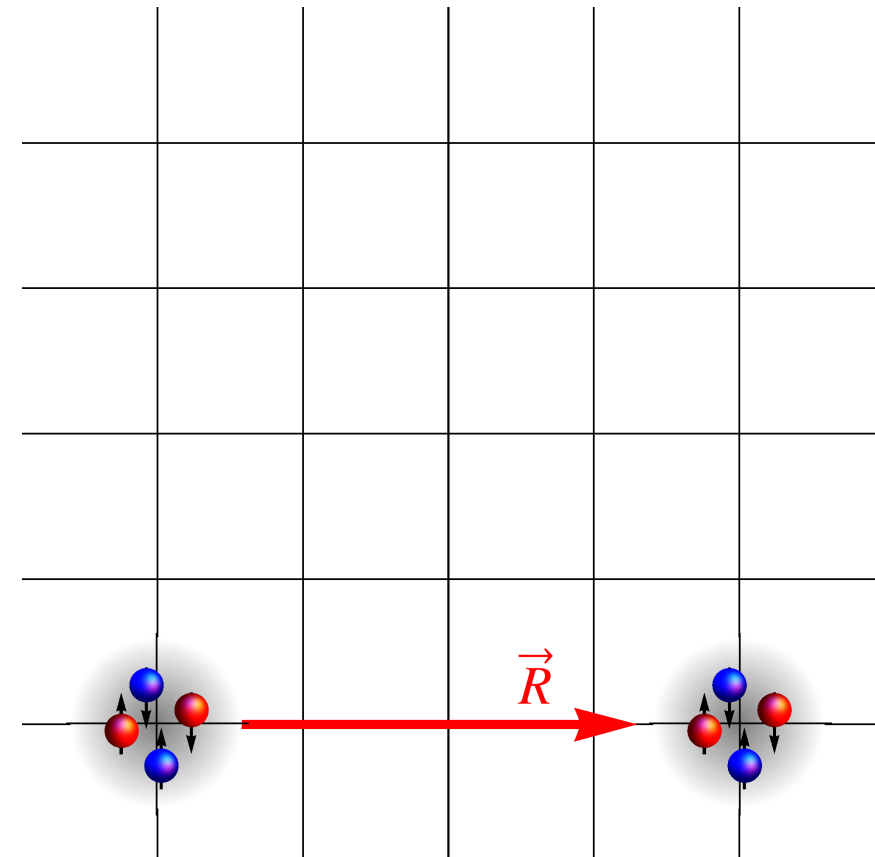
- **First Step:** low-energy cluster Hamiltonian
 - Set of **initial states** labeled by relative cluster separation

$$|\vec{R}\rangle = \sum_{\vec{r}} |\vec{r} + \vec{R}\rangle_1 \otimes |\vec{r}\rangle_2$$



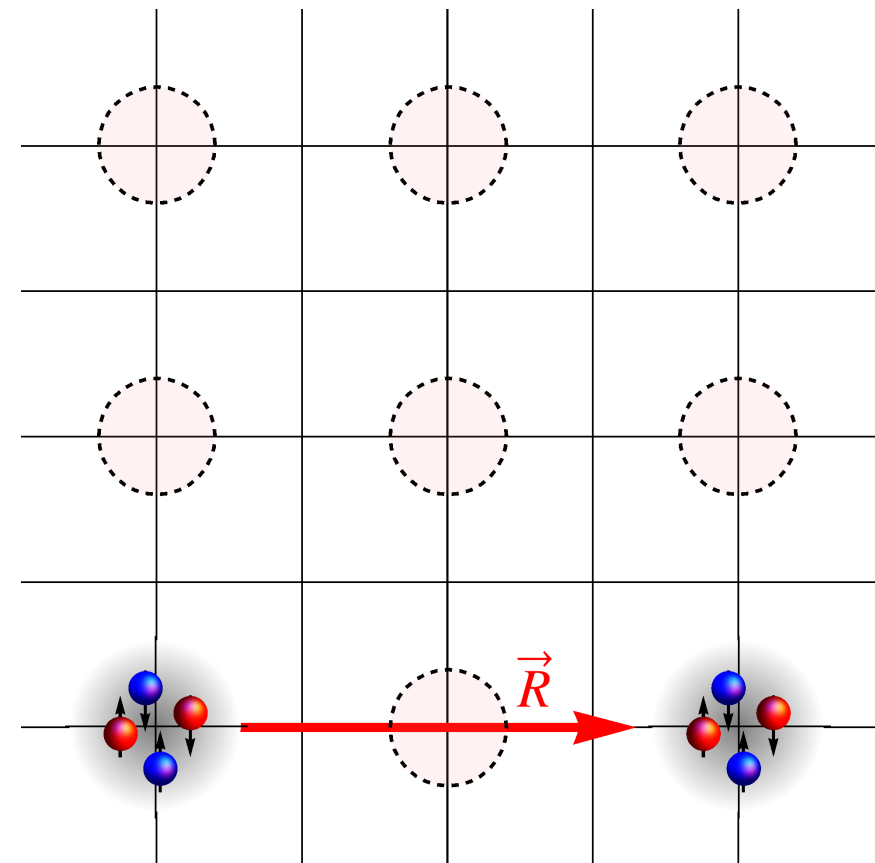
- **First Step:** low-energy cluster Hamiltonian
 - Set of **initial states** labeled by relative cluster separation

$$|\vec{R}\rangle = \sum_{\vec{r}} |\vec{r} + \vec{R}\rangle_1 \otimes |\vec{r}\rangle_2$$



- **First Step:** low-energy cluster Hamiltonian
 - Set of **initial states** labeled by relative cluster separation

$$|\vec{R}\rangle = \sum_{\vec{r}} |\vec{r} + \vec{R}\rangle_1 \otimes |\vec{r}\rangle_2$$

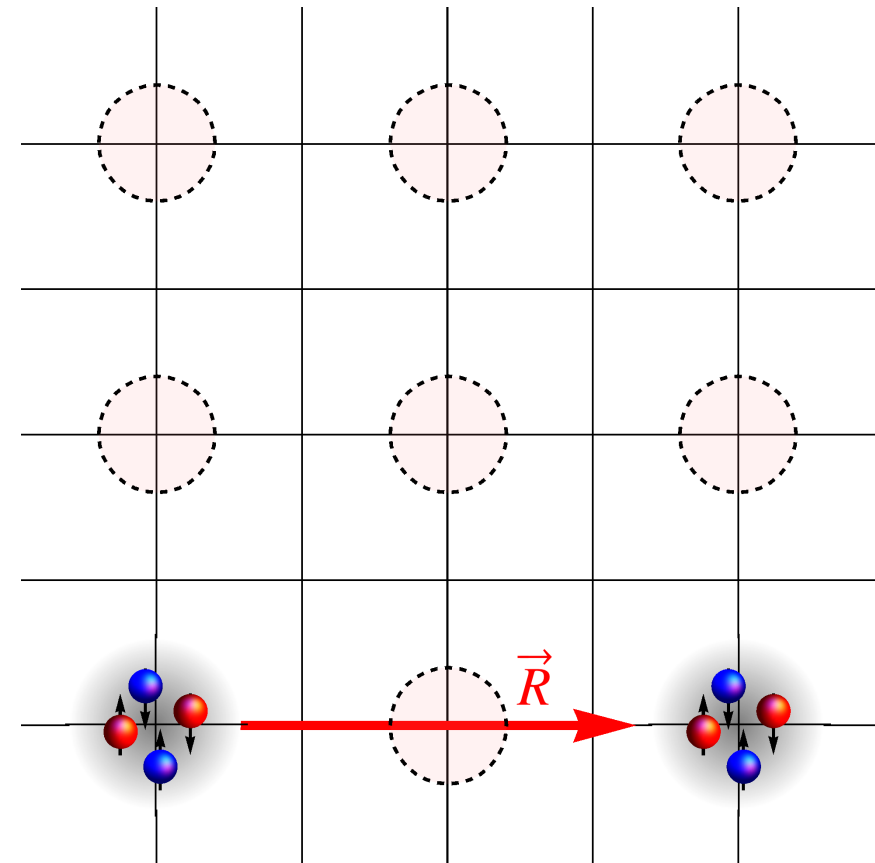


- **First Step:** low-energy cluster Hamiltonian
 - Set of **initial states** labeled by relative cluster separation

$$|\vec{R}\rangle = \sum_{\vec{r}} |\vec{r} + \vec{R}\rangle_1 \otimes |\vec{r}\rangle_2$$

- **Euclidean time propagation** with full microscopic Hamiltonian to calculate dressed cluster states

$$|\vec{R}\rangle_\tau = \exp(-H\tau)|\vec{R}\rangle$$



- **First Step:** low-energy cluster Hamiltonian
 - Set of **initial states** labeled by relative cluster separation

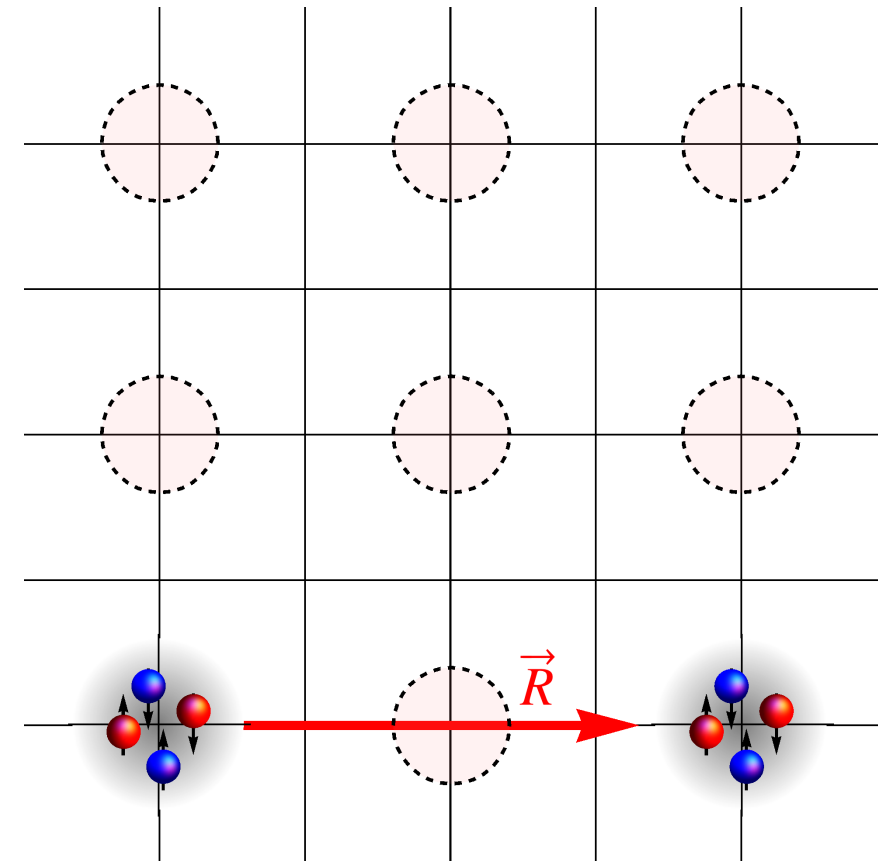
$$|\vec{R}\rangle = \sum_{\vec{r}} |\vec{r} + \vec{R}\rangle_1 \otimes |\vec{r}\rangle_2$$

- **Euclidean time propagation** with full microscopic Hamiltonian to calculate dressed cluster states

$$|\vec{R}\rangle_\tau = \exp(-H\tau)|\vec{R}\rangle$$

- The dressed cluster states span the **low-energy subspace** below the break-up scale of the clusters.
- Adiabatic Hamiltonian:

$$[H_\tau^a]_{\vec{R},\vec{R}'} = \sum_{\vec{R}'',\vec{R}'''} \left[{}_\tau\langle\vec{R}|\vec{R}''\rangle_\tau \right]^{-1/2} \left[{}_\tau\langle\vec{R}''|H|\vec{R}'''\rangle_\tau \right] \left[{}_\tau\langle\vec{R}'''|\vec{R}'\rangle_\tau \right]^{-1/2}$$



- **First Step:** low-energy cluster Hamiltonian
 - Set of **initial states** labeled by relative cluster separation

$$|\vec{R}\rangle = \sum_{\vec{r}} |\vec{r} + \vec{R}\rangle_1 \otimes |\vec{r}\rangle_2$$

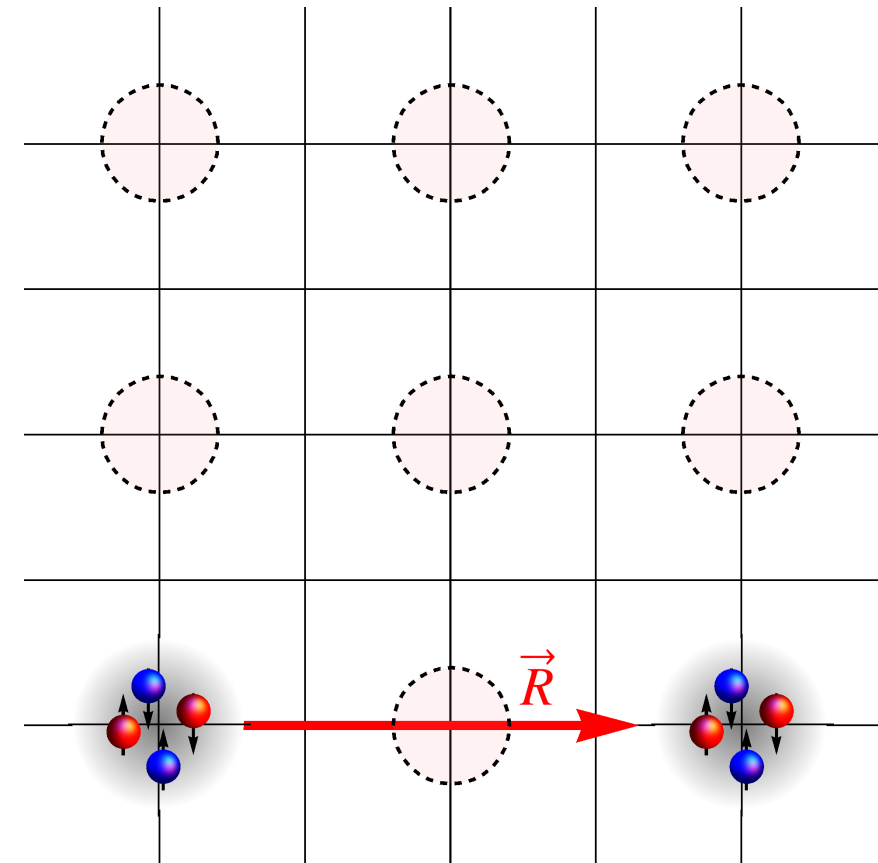
- **Euclidean time propagation** with full microscopic Hamiltonian to calculate dressed cluster states

$$|\vec{R}\rangle_\tau = \exp(-H\tau)|\vec{R}\rangle$$

- The dressed cluster states span the **low-energy subspace** below the break-up scale of the clusters.
- Adiabatic Hamiltonian:

$$[H_\tau^a]_{\vec{R},\vec{R}'} = \sum_{\vec{R}'',\vec{R}'''} \left[{}_\tau\langle\vec{R}|\vec{R}''\rangle_\tau \right]^{-1/2} \left[{}_\tau\langle\vec{R}''|H|\vec{R}'''\rangle_\tau \right] \left[{}_\tau\langle\vec{R}'''|\vec{R}'\rangle_\tau \right]^{-1/2}$$

Matrix elements of
dressed cluster states



- **First Step:** low-energy cluster Hamiltonian
 - Set of **initial states** labeled by relative cluster separation

$$|\vec{R}\rangle = \sum_{\vec{r}} |\vec{r} + \vec{R}\rangle_1 \otimes |\vec{r}\rangle_2$$

- **Euclidean time propagation** with full microscopic Hamiltonian to calculate dressed cluster states

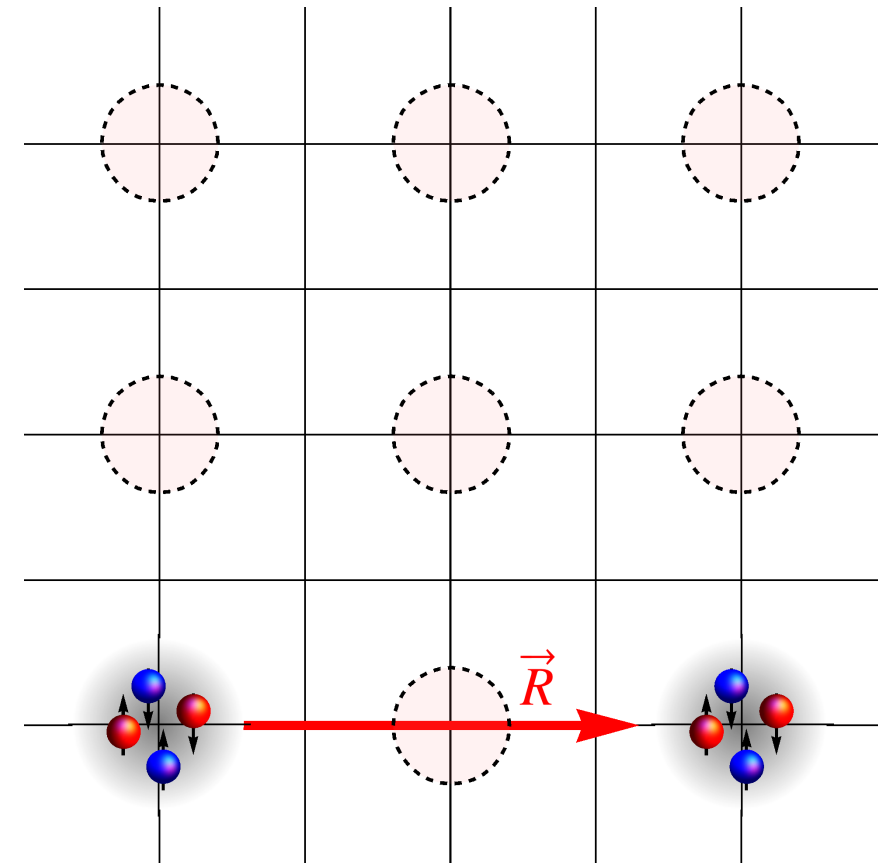
$$|\vec{R}\rangle_\tau = \exp(-H\tau) |\vec{R}\rangle$$

- The dressed cluster states span the **low-energy subspace** below the break-up scale of the clusters.
- Adiabatic Hamiltonian:

$$[H_\tau^a]_{\vec{R}, \vec{R}'} = \sum_{\vec{R}'', \vec{R}'''} \left[{}_\tau \langle \vec{R} | \vec{R}'' \rangle_\tau \right]^{-1/2} \left[{}_\tau \langle \vec{R}'' | H | \vec{R}''' \rangle_\tau \right] \left[{}_\tau \langle \vec{R}''' | \vec{R}' \rangle_\tau \right]^{-1/2}$$

Norm matrix

Matrix elements of
dressed cluster states



- **First Step:** low-energy cluster Hamiltonian
 - Set of **initial states** labeled by relative cluster separation

$$|\vec{R}\rangle = \sum_{\vec{r}} |\vec{r} + \vec{R}\rangle_1 \otimes |\vec{r}\rangle_2$$

- **Euclidean time propagation** with full microscopic Hamiltonian to calculate dressed cluster states

$$|\vec{R}\rangle_\tau = \exp(-H\tau) |\vec{R}\rangle$$

- The dressed cluster states span the **low-energy subspace** below the break-up scale of the clusters.
- Adiabatic Hamiltonian:

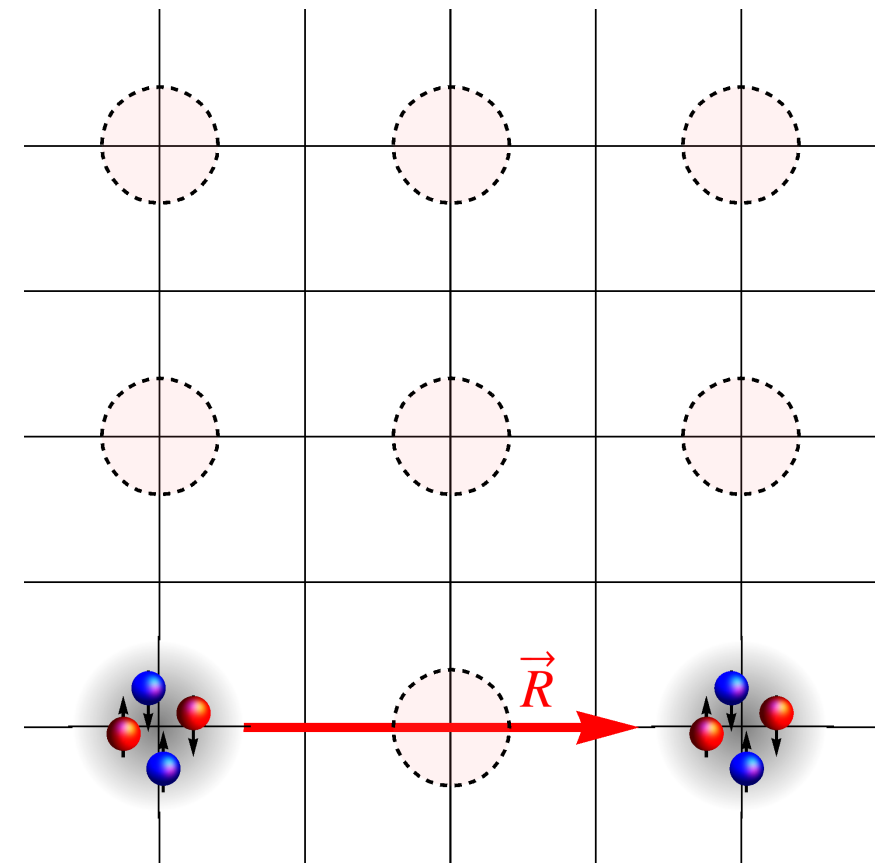
$$[H_\tau^a]_{\vec{R}, \vec{R}'} = \sum_{\vec{R}'', \vec{R}'''} \left[{}_\tau \langle \vec{R} | \vec{R}'' \rangle_\tau \right]^{-1/2} \left[{}_\tau \langle \vec{R}'' | H | \vec{R}''' \rangle_\tau \right] \left[{}_\tau \langle \vec{R}''' | \vec{R}' \rangle_\tau \right]^{-1/2}$$

similar to NCSM

Navratil, Quaglioni, Phys. Rev. C 83, 044609 (2011).

Norm matrix

Matrix elements of
dressed cluster states





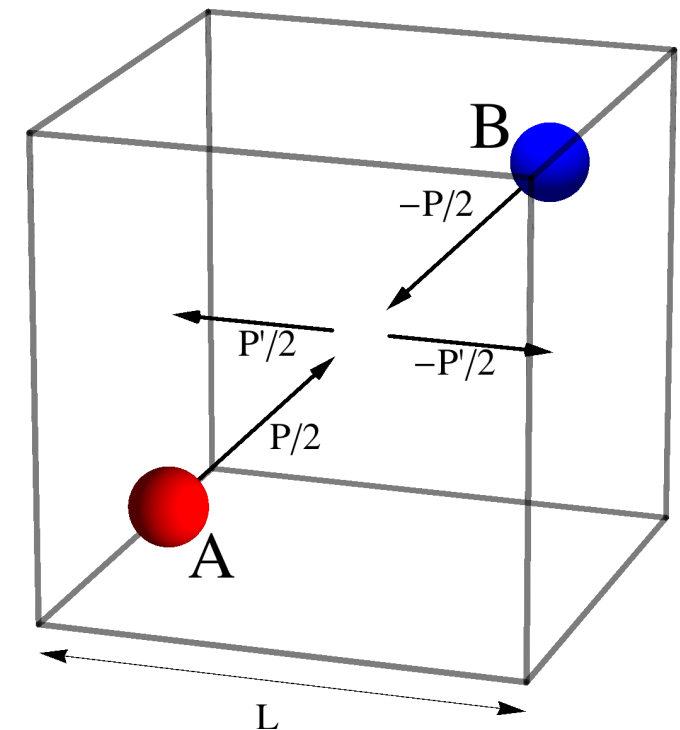
- **2 Step:** Extracting phase shifts

- **2 Step:** Extracting phase shifts
- **Lüscher's method:** Relation between energy levels in a finite periodic box and the infinite volume scattering phase shifts

M.Lüscher, Commun. Math. Phys. 105 (1986), 153

$$p \cot \delta_0(p) = \frac{1}{\pi L} S(\eta), \quad \eta = \frac{p(L)^2 L^2}{4\pi^2}$$

$$S(\eta) = \lim_{\Lambda \rightarrow \infty} \left[\sum_{\vec{k} \in \mathbb{Z}^3} \frac{\theta(\Lambda^2 - \vec{k}^2)}{\vec{k}^2 - \eta} - 4\pi \Lambda \right]$$

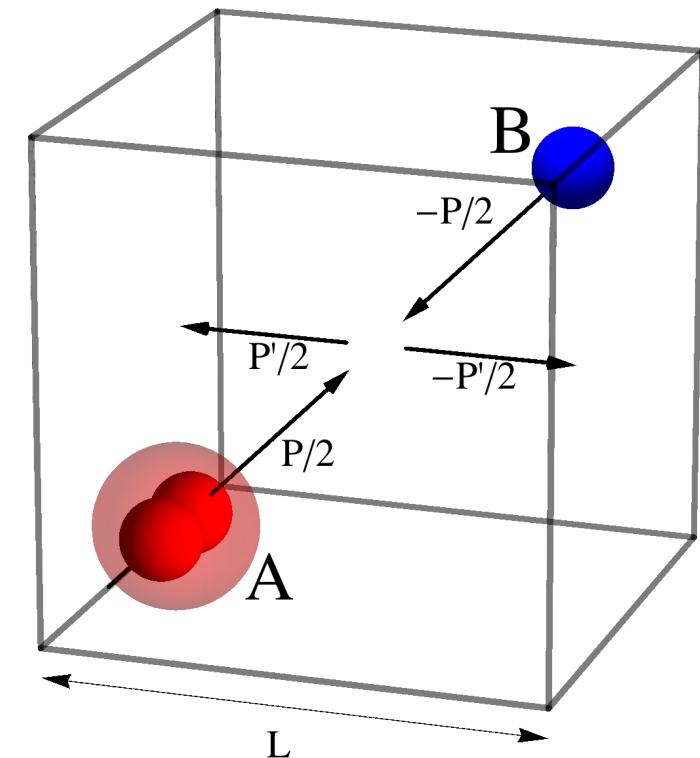


- **2 Step:** Extracting phase shifts
- **Lüscher's method:** Relation between energy levels in a finite periodic box and the infinite volume scattering phase shifts

M.Lüscher, Commun. Math. Phys. 105 (1986), 153

$$p \cot \delta_0(p) = \frac{1}{\pi L} S(\eta), \quad \eta = \frac{p(L)^2 L^2}{4\pi^2}$$

$$S(\eta) = \lim_{\Lambda \rightarrow \infty} \left[\sum_{\vec{k} \in \mathbb{Z}^3} \frac{\theta(\Lambda^2 - \vec{k}^2)}{\vec{k}^2 - \eta} - 4\pi \Lambda \right]$$

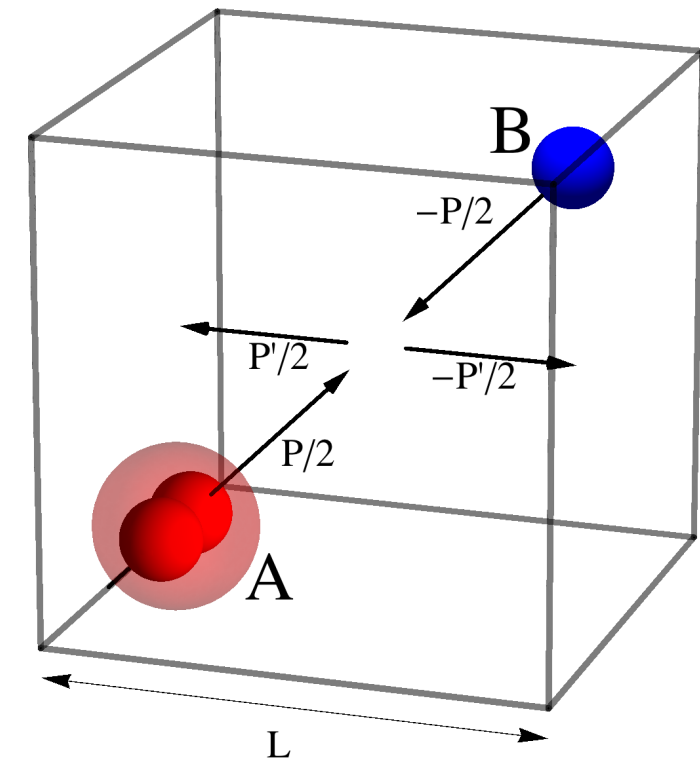


- **2 Step:** Extracting phase shifts
- **Lüscher's method:** Relation between energy levels in a finite periodic box and the infinite volume scattering phase shifts

M.Lüscher, Commun. Math. Phys. 105 (1986), 153

$$p \cot \delta_0(p) = \frac{1}{\pi L} S(\eta), \quad \eta = \frac{p(L)^2 L^2}{4\pi^2}$$

$$S(\eta) = \lim_{\Lambda \rightarrow \infty} \left[\sum_{\vec{k} \in \mathbb{Z}^3} \frac{\theta(\Lambda^2 - \vec{k}^2)}{\vec{k}^2 - \eta} - 4\pi \Lambda \right]$$



- **Topological corrections** due to cluster character

$$E(p, L) = \frac{p^2}{2\mu} - B_1 - B_2 + \tau_1(\eta) \Delta E_1(L) + \tau_2(\eta) \Delta E_2(L)$$

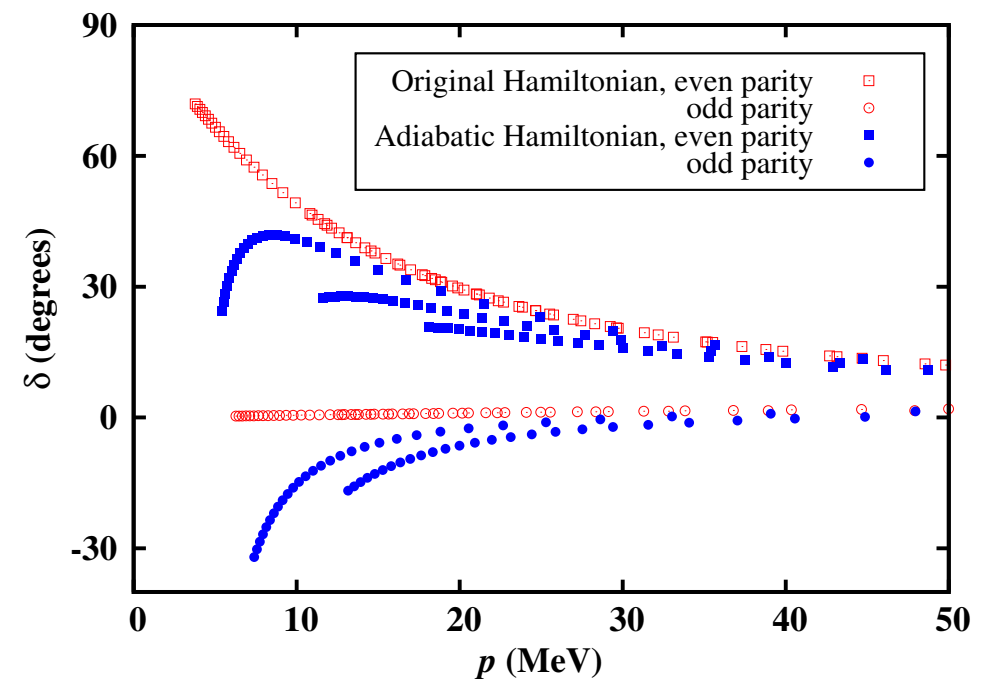
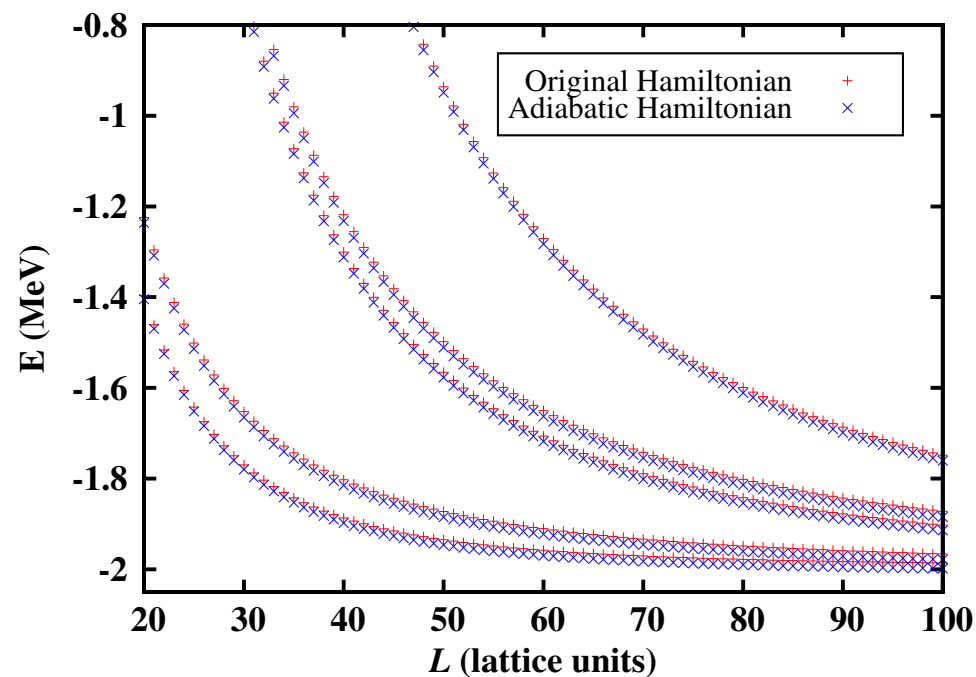
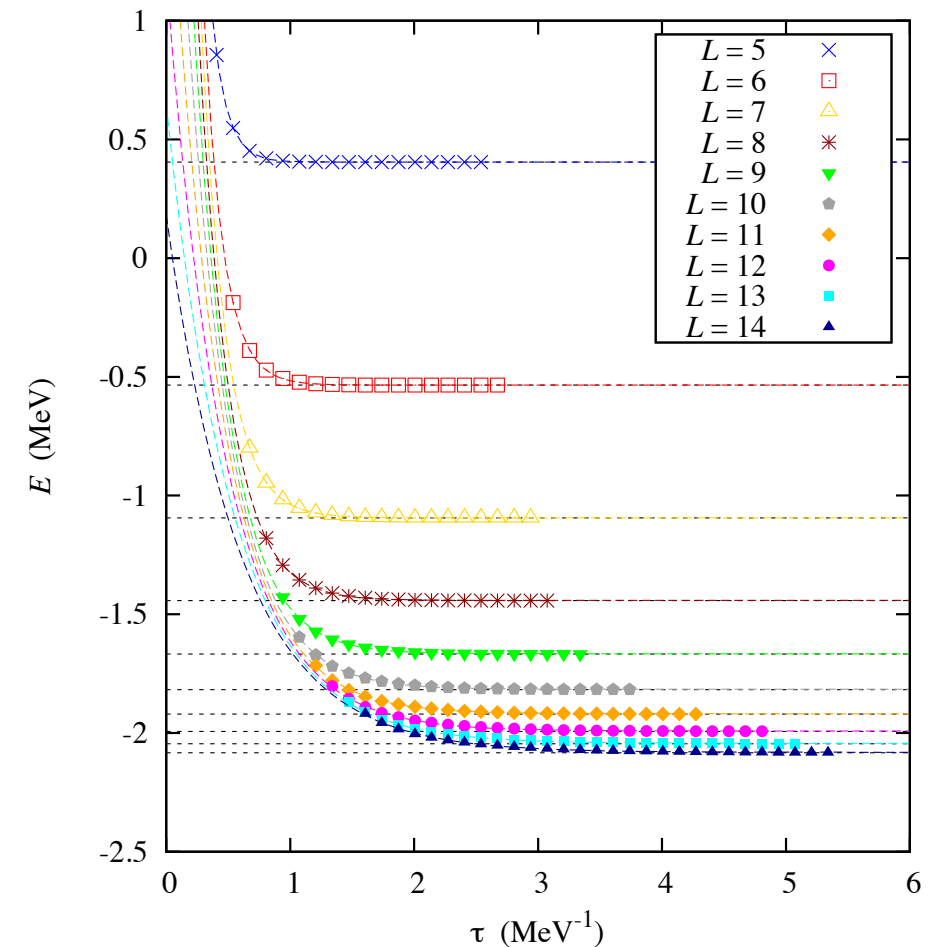
$$\tau(\eta) = \frac{1}{\sum_{\vec{k}} (\vec{k}^2 - \eta)^{-2}} \sum_{\vec{k}} \frac{\sum_{i=1}^3 \cos(2\pi k_i \alpha)}{3(\vec{k}^2 - \eta)^2}$$

S. Bour, S. König, D. Lee, H.-W. Hammer and U.-G. Meißner, Phys. Rev. D 84 (2011), 091503

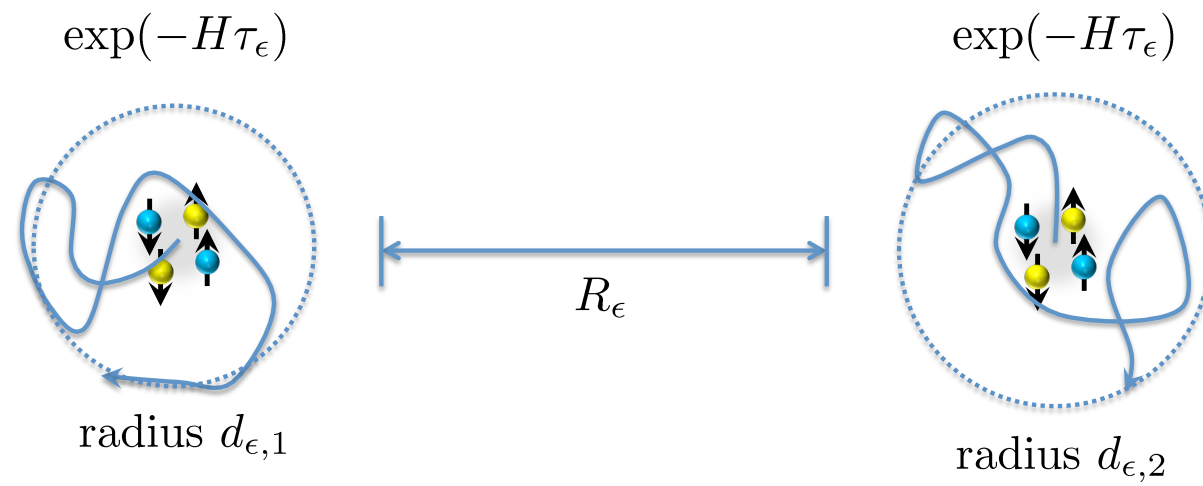
S. Bour, H.-W. Hammer, D. Lee and U.-G. Meißner, Phys. Rev. C 86 (2012), 034003

- There is an exponentially small error in energy levels due to Euclidean time projection
- In larger systems there is a statistical error due to Monte Carlo methods

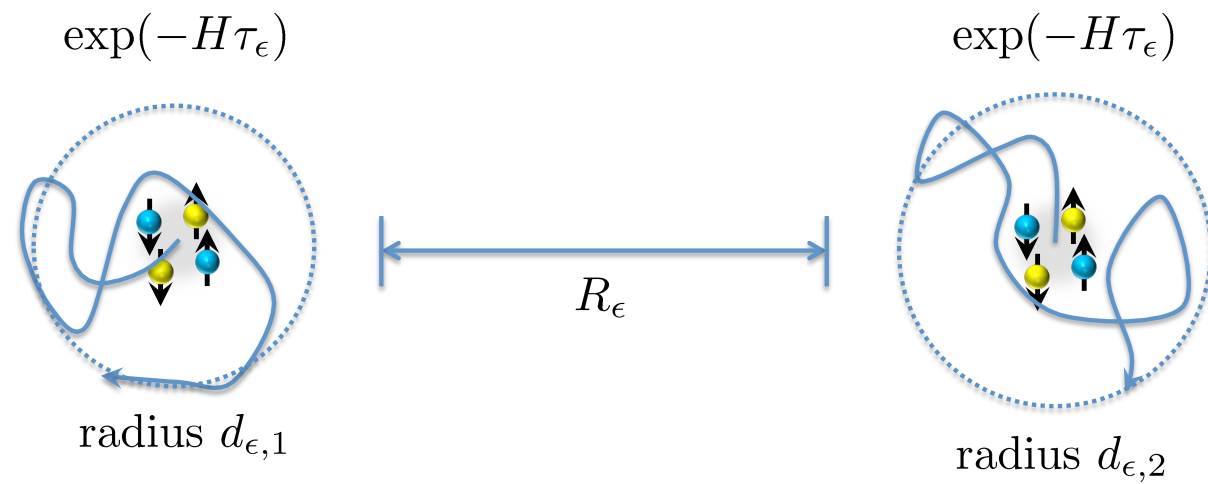
Lüscher's method is unfortunately very sensitive to small errors in energy levels!



Asymptotic cluster wave functions (I)



Asymptotic cluster wave functions (I)

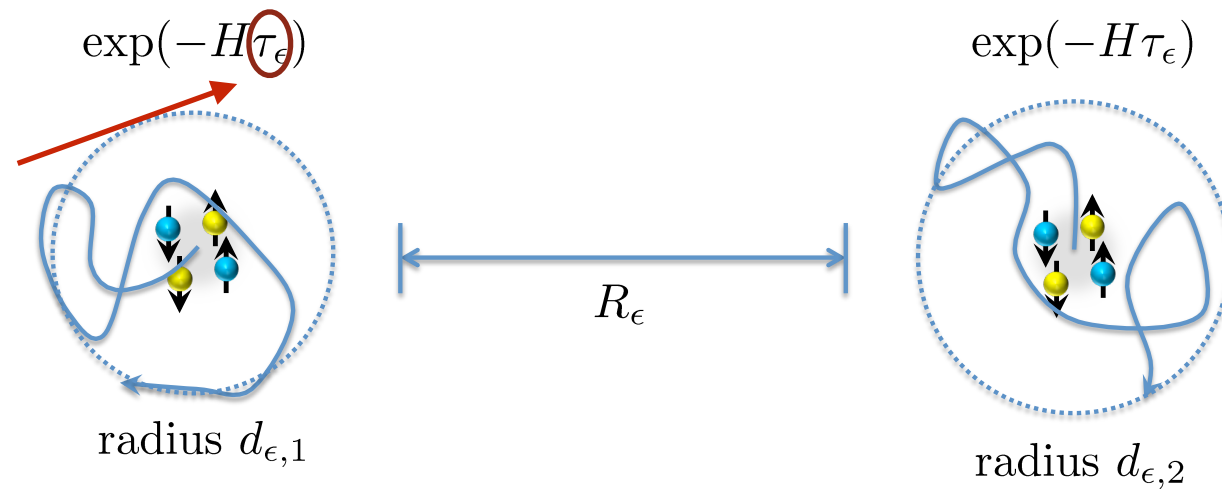


ϵ : relative error

Asymptotic cluster wave functions (I)



time interval such that
the relative contamination
from exited cluster states
is less than ϵ

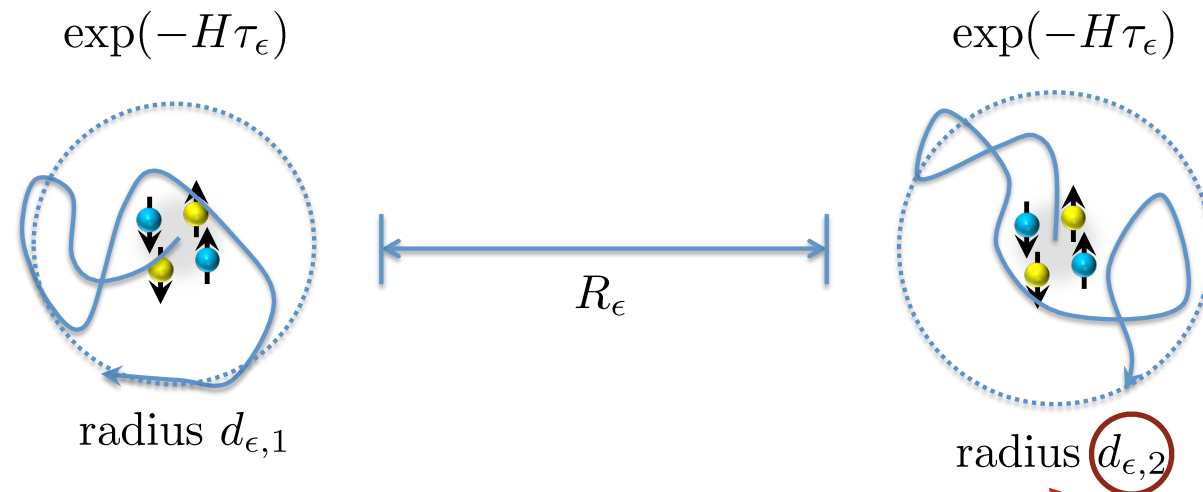


ϵ : relative error

Asymptotic cluster wave functions (I)



time interval such that
the relative contamination
from exited cluster states
is less than ϵ



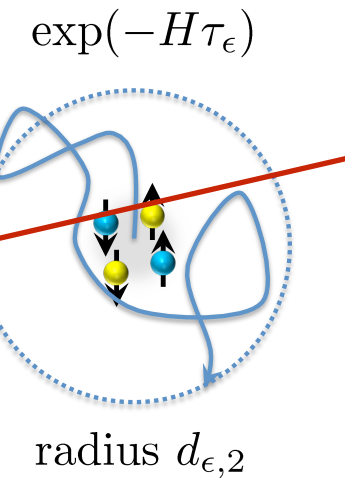
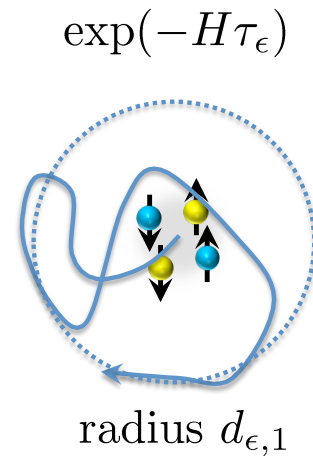
spatial diffusion that
the clusters undergo in time τ_{ϵ}

ϵ : relative error

Asymptotic cluster wave functions (I)



time interval such that
the relative contamination
from exited cluster states
is less than ϵ



asymptotic distance such
that the overlap between
clusters is less than ϵ

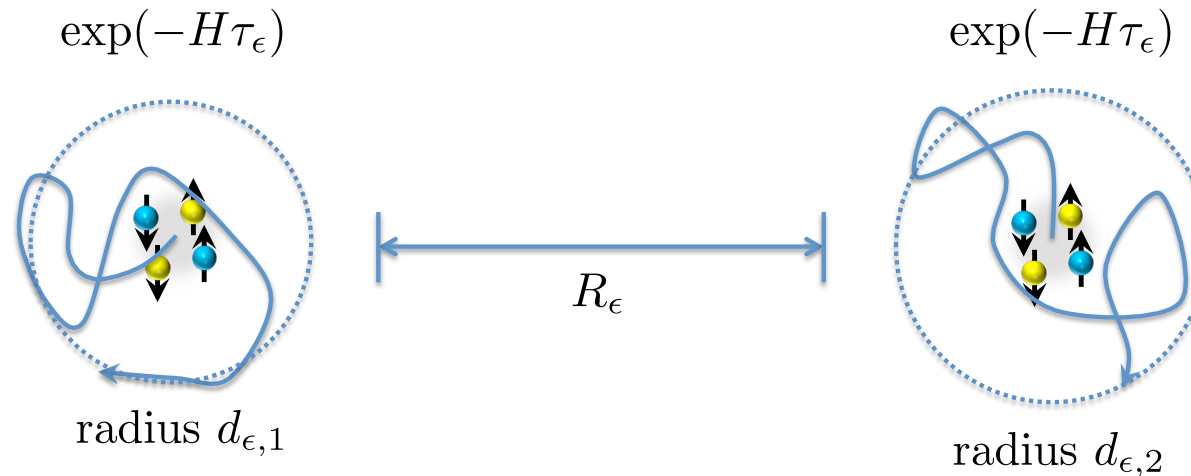
spatial diffusion that
the clusters undergo in time τ_ϵ

ϵ : relative error

Asymptotic cluster wave functions (I)



time interval such that
the relative contamination
from exited cluster states
is less than ϵ



asymptotic distance such
that the overlap between
clusters is less than ϵ

spatial diffusion that
the clusters undergo in time τ_ϵ

ϵ : relative error

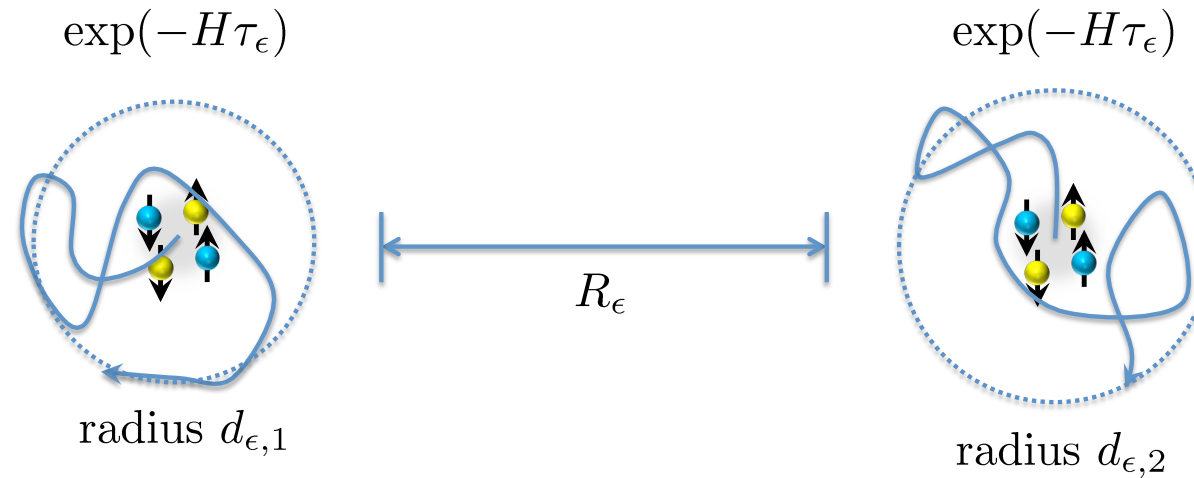
For $L > R_\epsilon \gg d_{\epsilon,1}, d_{\epsilon,2}$ in the asymptotic region $|\vec{R}| > R_\epsilon$
the Hamiltonian is similar to a free lattice Hamiltonian H_{eff}

$$[H_\tau^a]_{\vec{R}, \vec{R}'} = \sum_{\vec{R}'', \vec{R}'''} \left[{}_\tau \langle \vec{R} | \vec{R}'' \rangle_\tau \right]^{-1/2} \left[{}_\tau \langle \vec{R}'' | H | \vec{R}''' \rangle_\tau \right] \left[{}_\tau \langle \vec{R}''' | \vec{R}' \rangle_\tau \right]^{-1/2}$$

Asymptotic cluster wave functions (I)



time interval such that
the relative contamination
from exited cluster states
is less than ϵ



asymptotic distance such
that the overlap between
clusters is less than ϵ

spatial diffusion that
the clusters undergo in time τ_ϵ

ϵ : relative error

For $L > R_\epsilon \gg d_{\epsilon,1}, d_{\epsilon,2}$ in the asymptotic region $|\vec{R}| > R_\epsilon$
the Hamiltonian is similar to a free lattice Hamiltonian H_{eff}

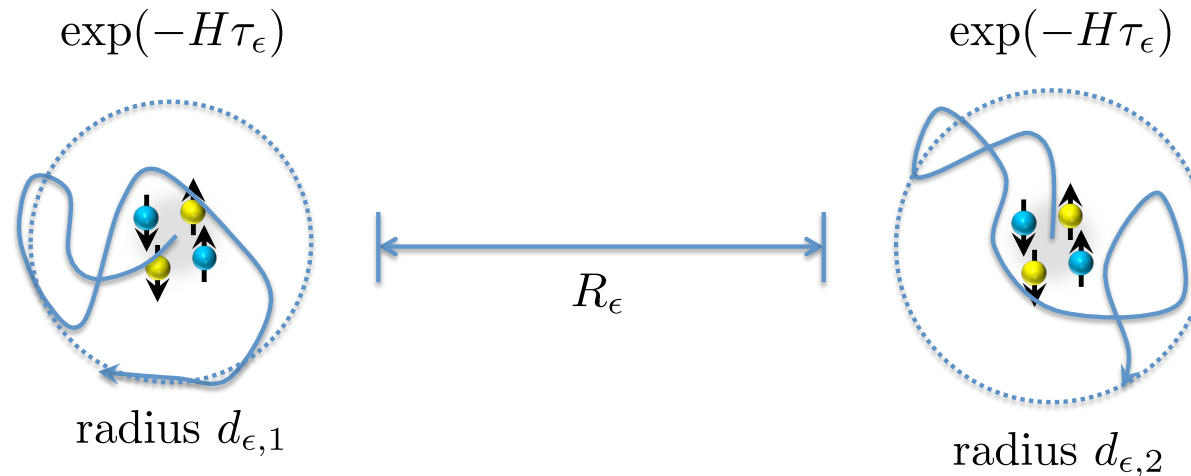
$$[H_\tau^a]_{\vec{R}, \vec{R}'} = \sum_{\vec{R}'', \vec{R}'''} \left[{}_\tau \langle \vec{R} | \vec{R}'' \rangle_\tau \right]^{-1/2} \left[{}_\tau \langle \vec{R}'' | H | \vec{R}''' \rangle_\tau \right] \left[{}_\tau \langle \vec{R}''' | \vec{R}' \rangle_\tau \right]^{-1/2}$$

\uparrow
 $|\vec{R}\rangle_\tau = \exp(-H\tau)|\vec{R}\rangle$

Asymptotic cluster wave functions (I)



time interval such that
the relative contamination
from exited cluster states
is less than ϵ



asymptotic distance such
that the overlap between
clusters is less than ϵ

spatial diffusion that
the clusters undergo in time τ_ϵ

ϵ : relative error

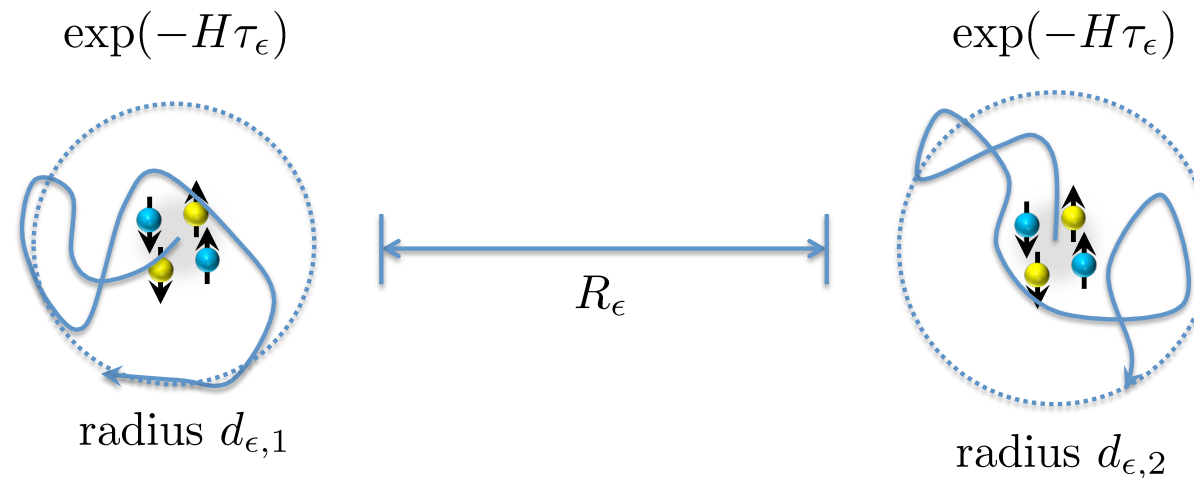
For $L > R_\epsilon \gg d_{\epsilon,1}, d_{\epsilon,2}$ in the asymptotic region $|\vec{R}| > R_\epsilon$
the Hamiltonian is similar to a free lattice Hamiltonian H_{eff}

$$[H_\tau^a]_{\vec{R}, \vec{R}'} = \sum_{\vec{R}'', \vec{R}'''} \left[\langle \vec{R} | e^{-2H_{\text{eff}}\tau} | \vec{R}'' \rangle \right]^{-\frac{1}{2}} \langle \vec{R}'' | e^{-H_{\text{eff}}\tau} H_{\text{eff}} e^{-H_{\text{eff}}\tau} | \vec{R}''' \rangle \left[\langle \vec{R}''' | e^{-2H_{\text{eff}}\tau} | \vec{R}' \rangle \right]^{-\frac{1}{2}}$$

Asymptotic cluster wave functions (I)



time interval such that
the relative contamination
from excited cluster states
is less than ϵ



asymptotic distance such
that the overlap between
clusters is less than ϵ

spatial diffusion that
the clusters undergo in time τ_ϵ

ϵ : relative error

For $L > R_\epsilon \gg d_{\epsilon,1}, d_{\epsilon,2}$ in the asymptotic region $|\vec{R}| > R_\epsilon$
the Hamiltonian is similar to a free lattice Hamiltonian H_{eff}

$$[H_\tau^a]_{\vec{R}, \vec{R}'} = [H_{\text{eff}}]_{\vec{R}, \vec{R}'}$$

Asymptotic cluster wave functions (II)



- The asymptotic cluster wave function can also be used to extract phase shifts

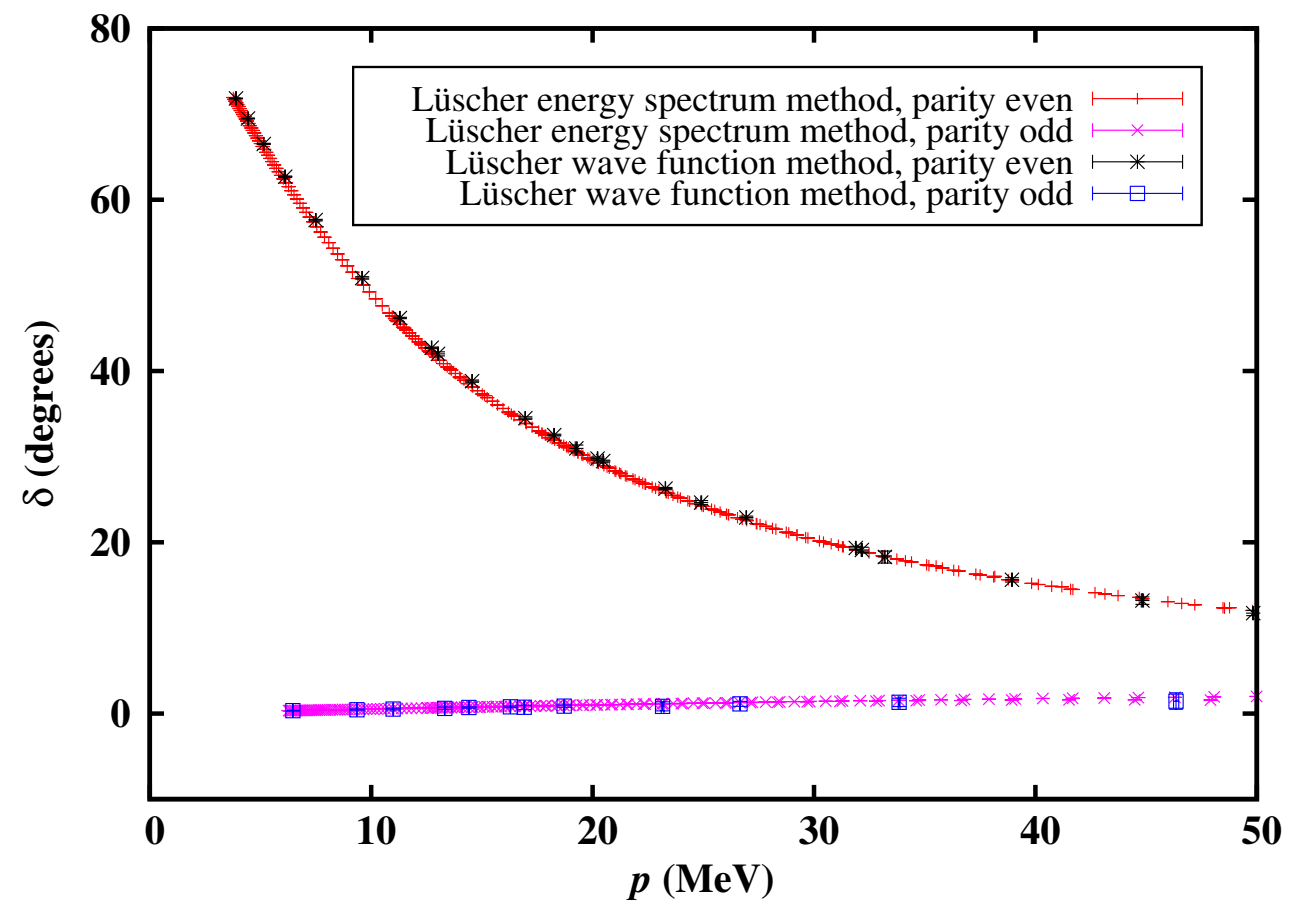
- The asymptotic cluster wave function can also be used to extract phase shifts
 - 1-dim asymptotic form:

$$\Psi_{\ell}^{(p)}(r) = A_{\ell} \cos(pr + \delta_{\ell} - \ell\pi/2)$$

- The asymptotic cluster wave function can also be used to extract phase shifts
- 1-dim asymptotic form:

$$\Psi_{\ell}^{(p)}(r) = A_{\ell} \cos(pr + \delta_{\ell} - \ell\pi/2)$$

check with Lüscher's method applied
to full microscopic Hamiltonian

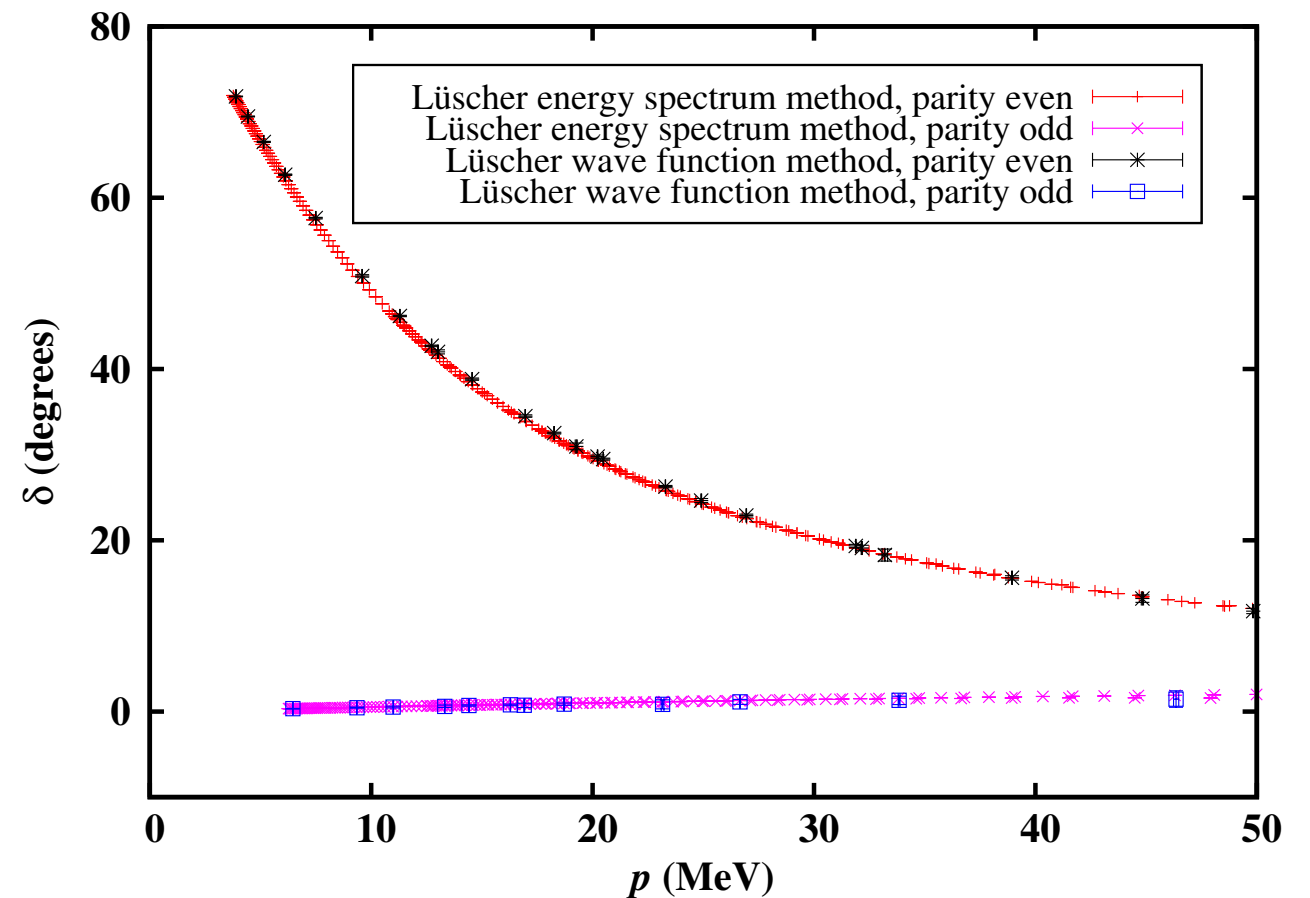


- The asymptotic cluster wave function can also be used to extract phase shifts

- 1-dim asymptotic form:

$$\Psi_{\ell}^{(p)}(r) = A_{\ell} \cos(pr + \delta_{\ell} - \ell\pi/2)$$

check with Lüscher's method applied
to full microscopic Hamiltonian

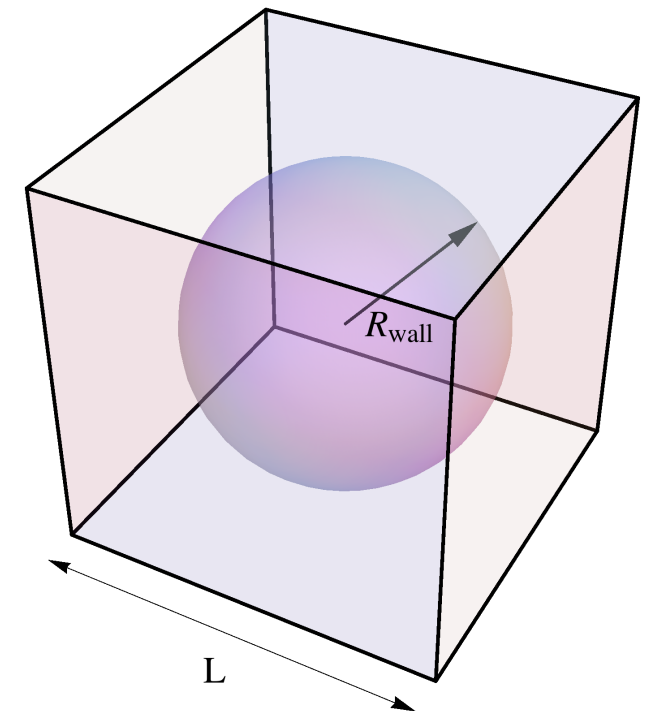


- 3-dim asymptotic form:

$$\Psi_{\ell, m_{\ell}}^{(p)}(\mathbf{r}) = A_{\ell} Y_{\ell, m_{\ell}}(\theta, \phi) [\cos \delta_{\ell}(p) j_{\ell}(pr) - \sin \delta_{\ell}(p) n_{\ell}(pr)]$$

- Impose a hard wall on the relative separation of two point-like particles and fit it to the asymptotic form

Borasoy, Epelbaum, Krebs, Lee, Meißner, EPJA 34 (2007) 185

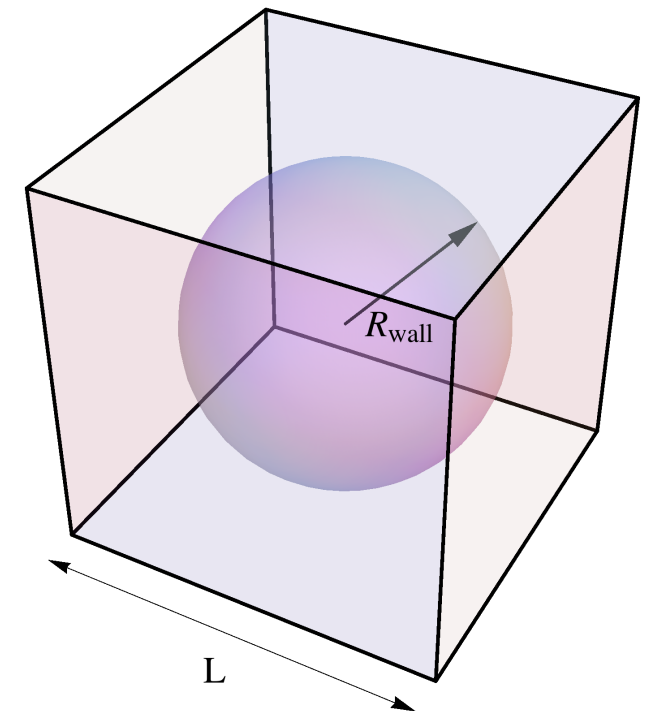


- Impose a hard wall on the relative separation of two point-like particles and fit it to the asymptotic form

Borasoy, Epelbaum, Krebs, Lee, Meißner, EPJA 34 (2007) 185

- 3-parameter fit:

$$\Psi_{\ell}^{(p)}(r) = A_{\ell} \cos(pr + \delta_{\ell} - \ell\pi/2)$$

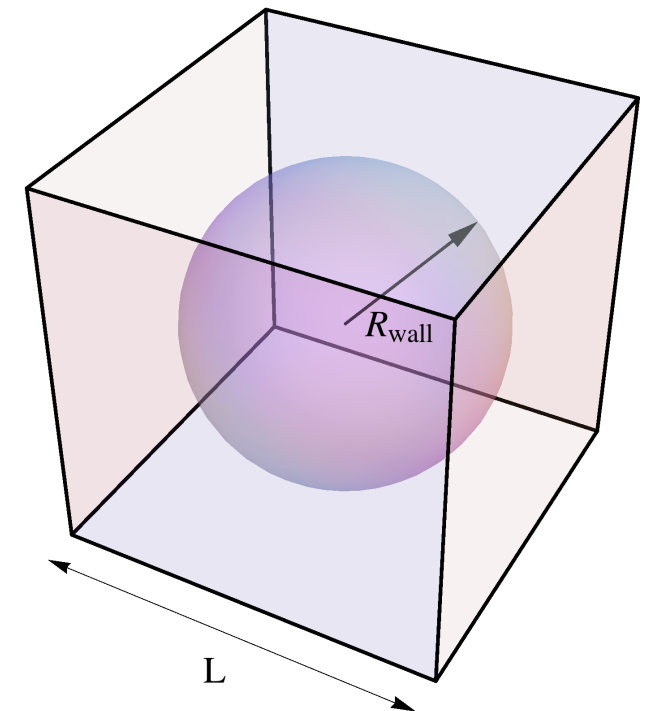


- Impose a hard wall on the relative separation of two point-like particles and fit it to the asymptotic form

Borasoy, Epelbaum, Krebs, Lee, Meißner, EPJA 34 (2007) 185

- 3-parameter fit:

$$\Psi_{\ell}^{(p)}(r) = A_{\ell} \cos(pr + \delta_{\ell} - \ell\pi/2)$$



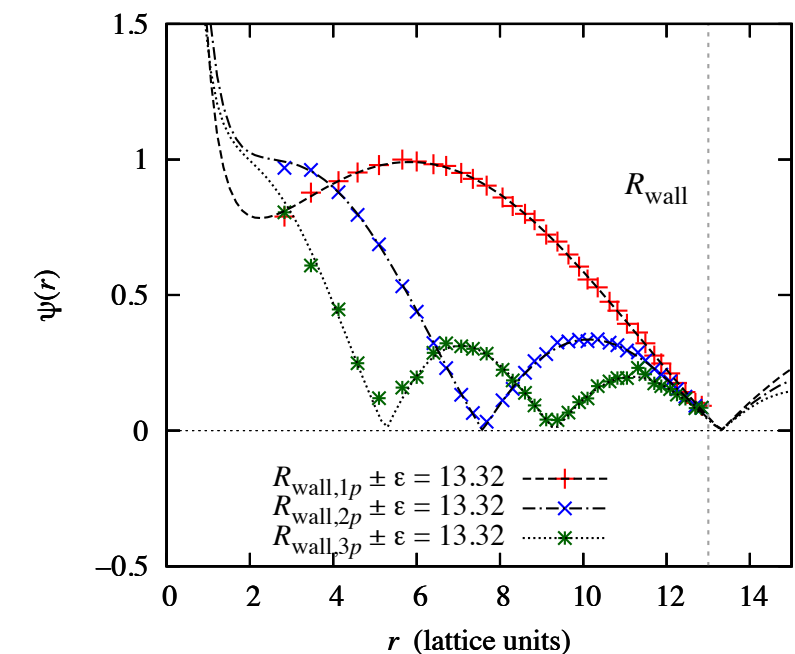
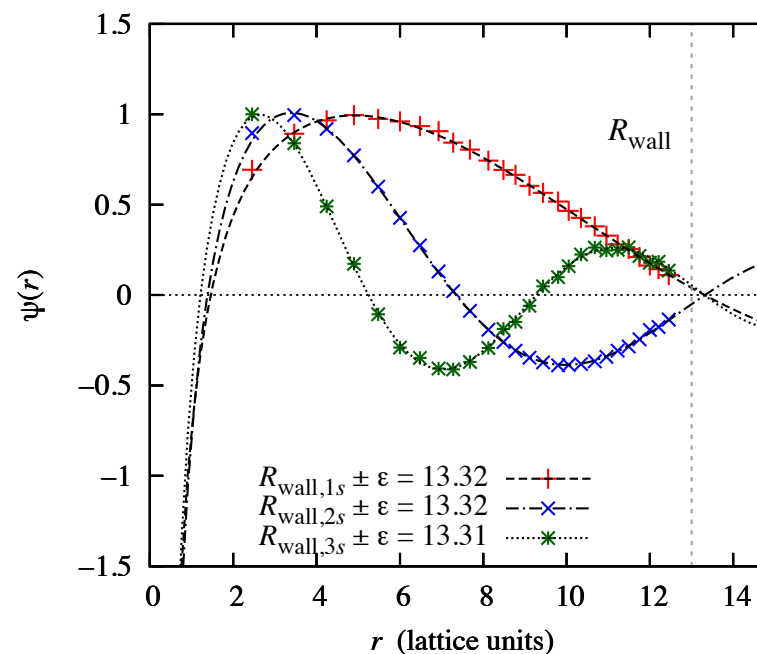
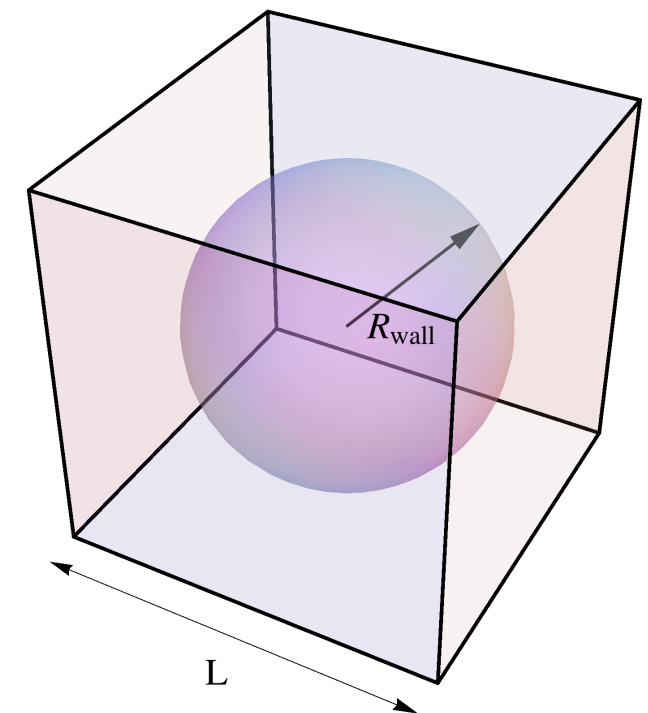
- Impose a hard wall on the relative separation of two point-like particles and fit it to the asymptotic form

Borasoy, Epelbaum, Krebs, Lee, Meißner, EPJA 34 (2007) 185

- 3-parameter fit:

$$\Psi_{\ell}^{(p)}(r) = A_{\ell} \cos(pr + \delta_{\ell} - \ell\pi/2)$$

- Observation: $R_{\text{wall}}' = R_{\text{wall}} + \varepsilon$ changes very little



- Impose a hard wall on the relative separation of two point-like particles and fit it to the asymptotic form

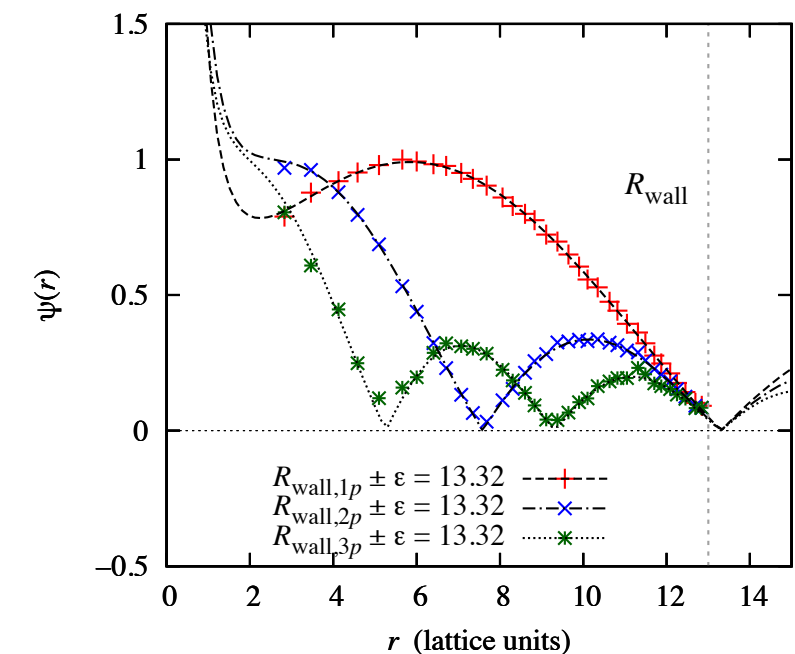
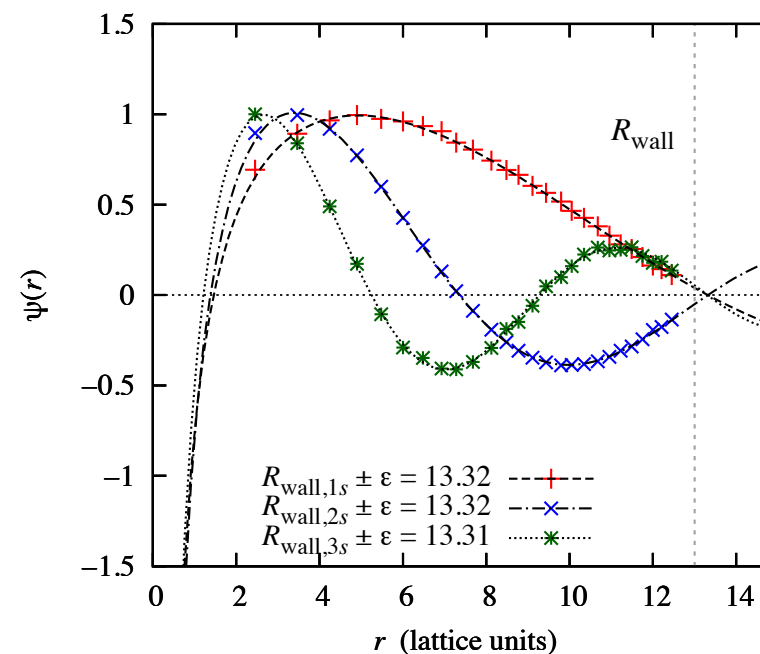
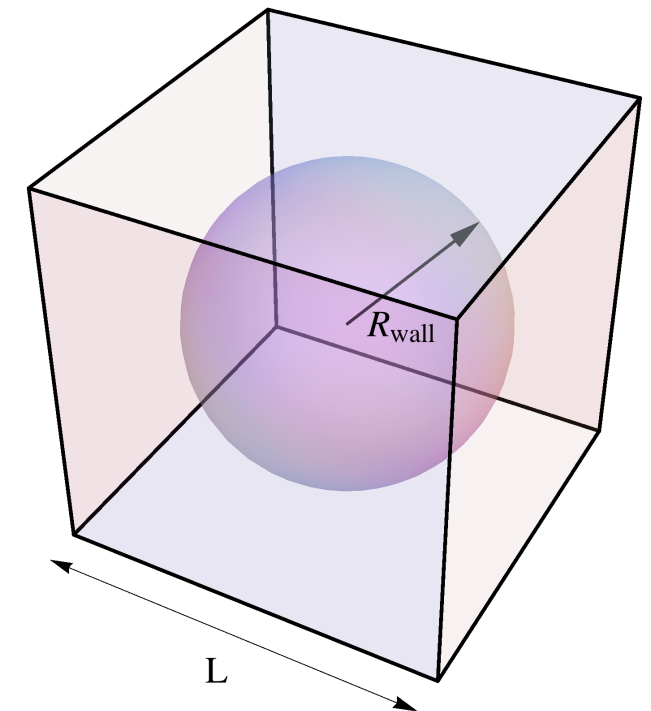
Borasoy, Epelbaum, Krebs, Lee, Meißner, EPJA 34 (2007) 185

- 3-parameter fit:

$$\Psi_{\ell}^{(p)}(r) = A_{\ell} \cos(pr + \delta_{\ell} - \ell\pi/2)$$

- Observation: $R_{\text{wall}}' = R_{\text{wall}} + \varepsilon$ changes very little

- 2-parameter fit:



- Impose a hard wall on the relative separation of two point-like particles and fit it to the asymptotic form

Borasoy, Epelbaum, Krebs, Lee, Meißner, EPJA 34 (2007) 185

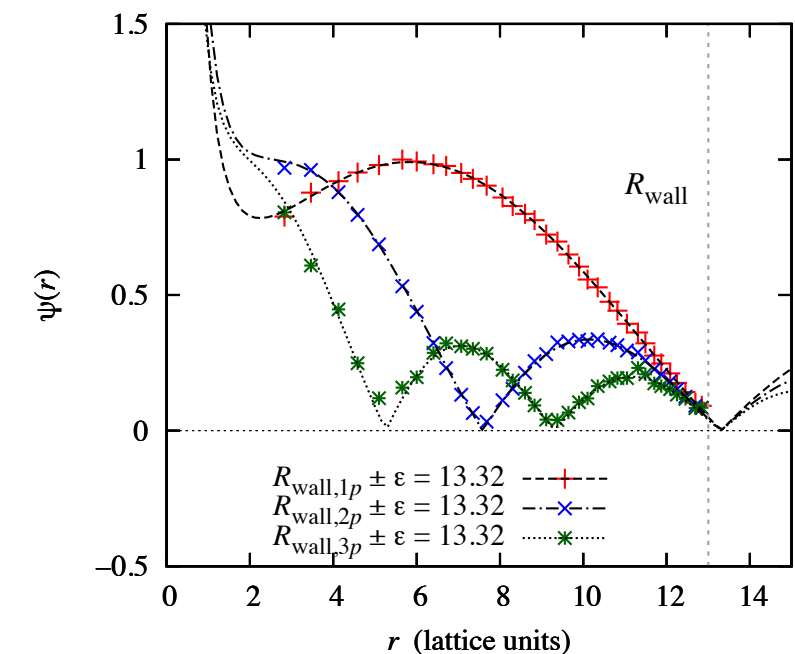
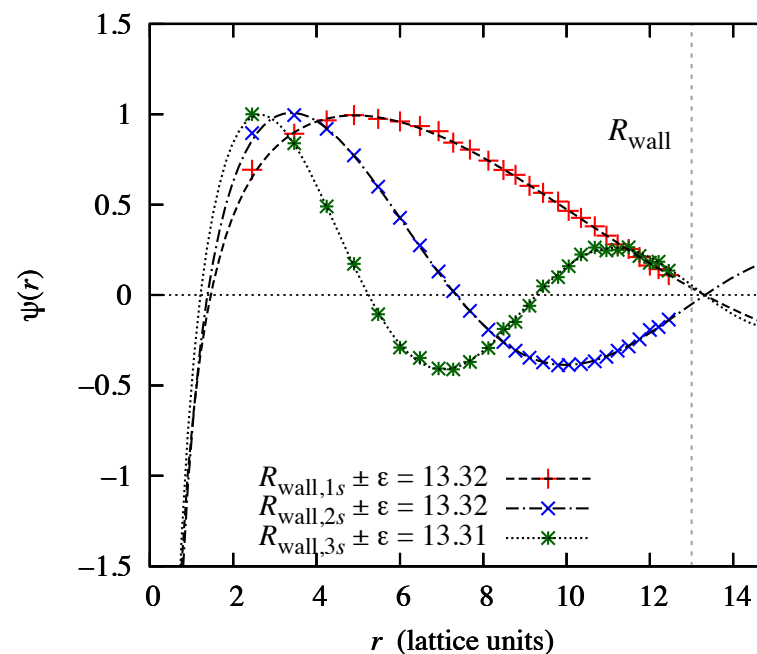
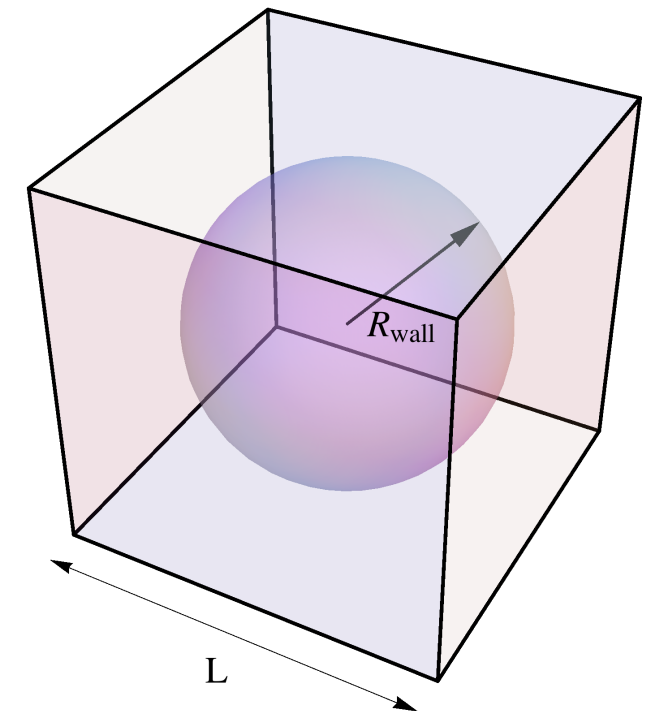
- 3-parameter fit:

$$\Psi_{\ell}^{(p)}(r) = A_{\ell} \cos(pr + \delta_{\ell} - \ell\pi/2)$$

- Observation: $R_{\text{wall}}' = R_{\text{wall}} + \varepsilon$ changes very little

- 2-parameter fit:

- determine R_{wall}' from non-interacting system



- Impose a hard wall on the relative separation of two point-like particles and fit it to the asymptotic form

Borasoy, Epelbaum, Krebs, Lee, Meißner, EPJA 34 (2007) 185

- 3-parameter fit:

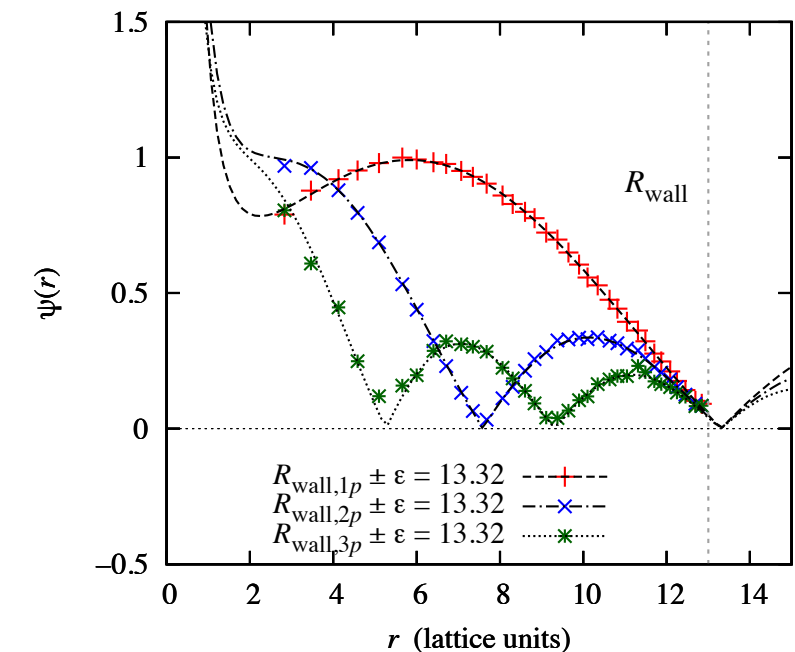
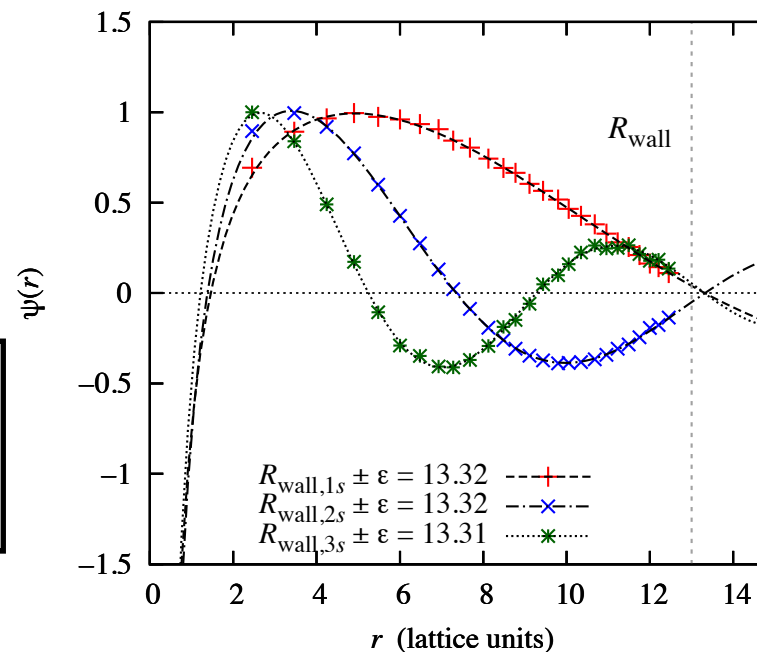
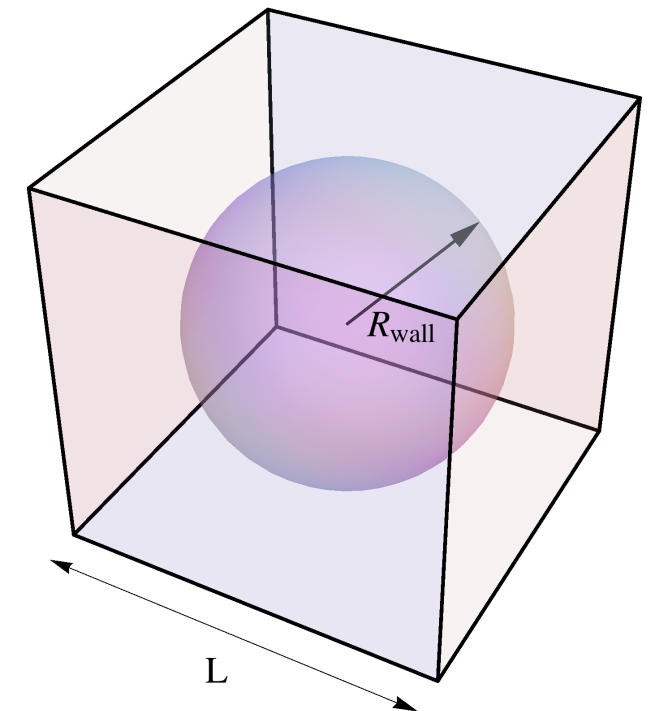
$$\Psi_{\ell}^{(p)}(r) = A_{\ell} \cos(pr + \delta_{\ell} - \ell\pi/2)$$

- Observation: $R_{\text{wall}}' = R_{\text{wall}} + \varepsilon$ changes very little

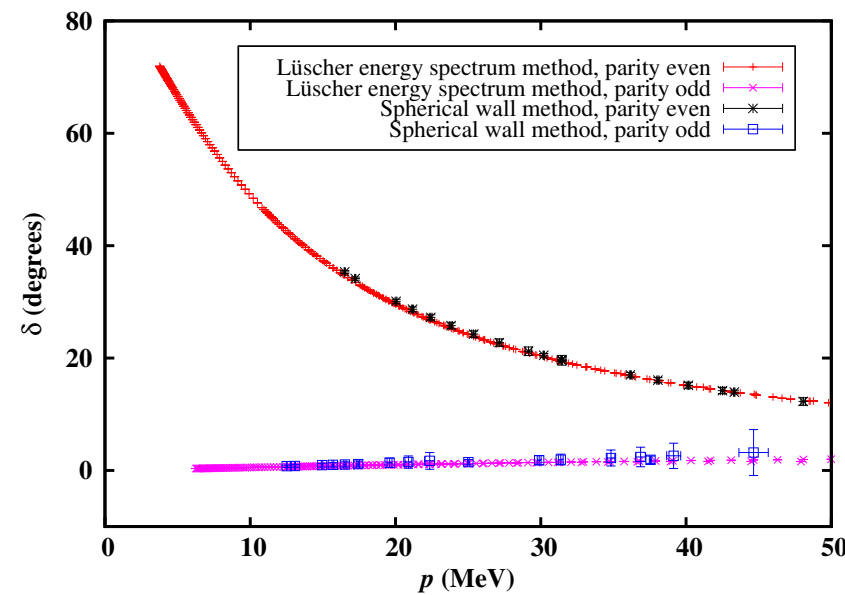
- 2-parameter fit:

- determine R_{wall}' from non-interacting system

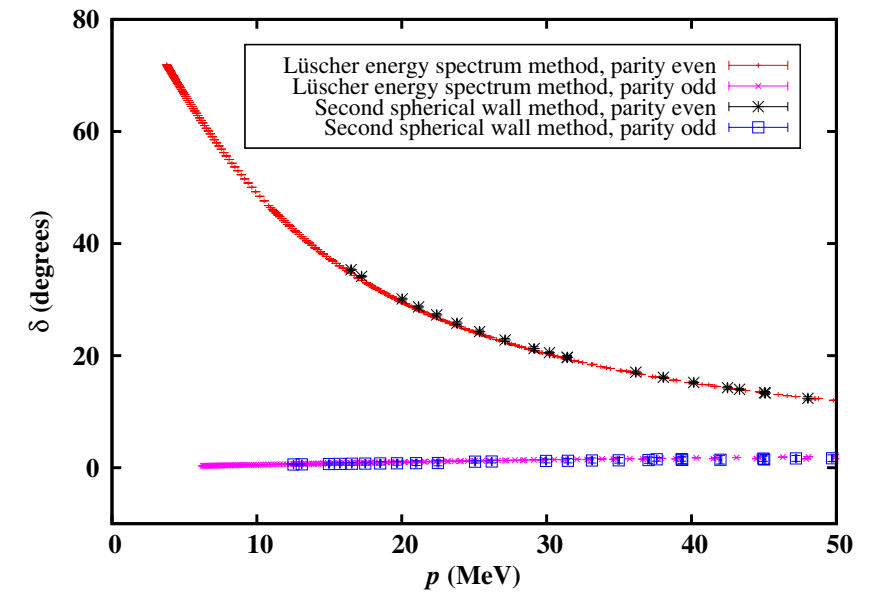
$$\delta_{\ell}(p) = \begin{cases} -pR'_{\text{wall}} + \frac{\pi(\ell+1)}{2} \mod \pi \\ \tan^{-1} \left[\frac{j_{\ell}(pR'_{\text{wall}}/a)}{n_{\ell}(pR'_{\text{wall}}/a)} \right] \end{cases}$$



- particle-dimer in one dimension



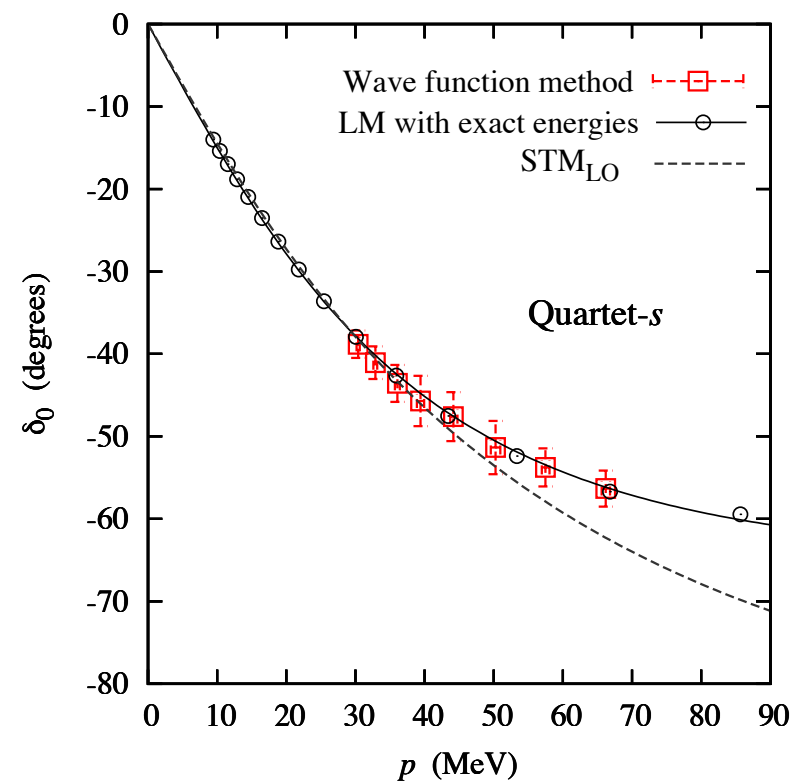
3-parameter fit



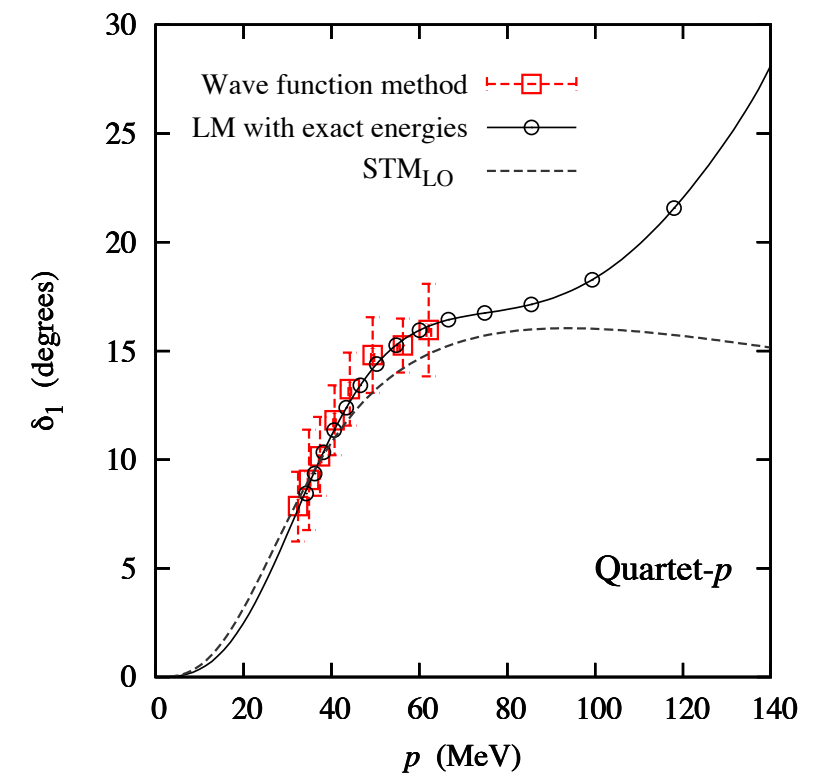
2-parameter fit

- fermion-dimer in three dimensions

Gabbiani, Bedaque, Grißhammer, NPA 675 (2000) 601

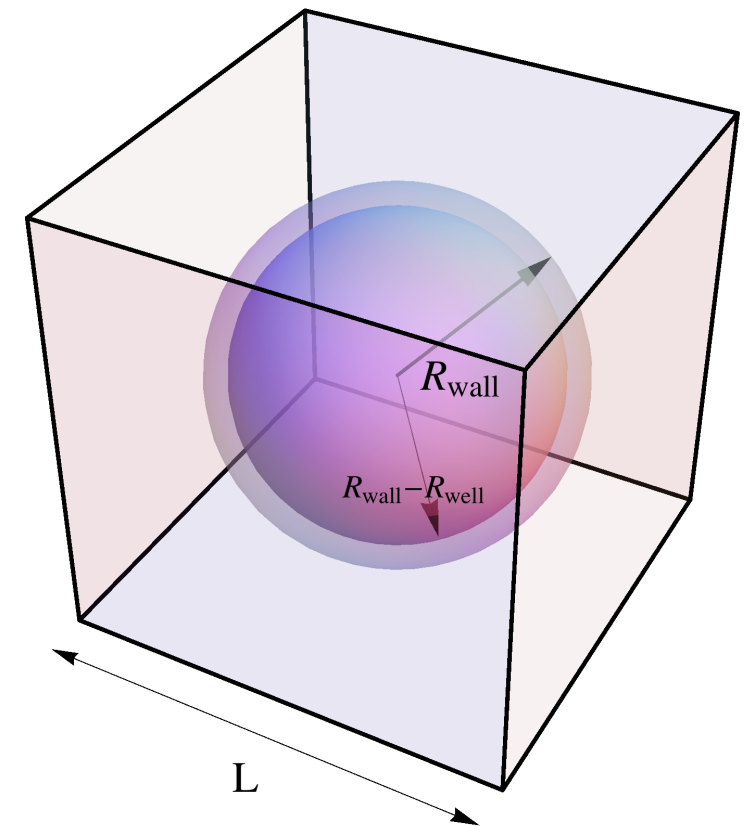


2-parameter fit

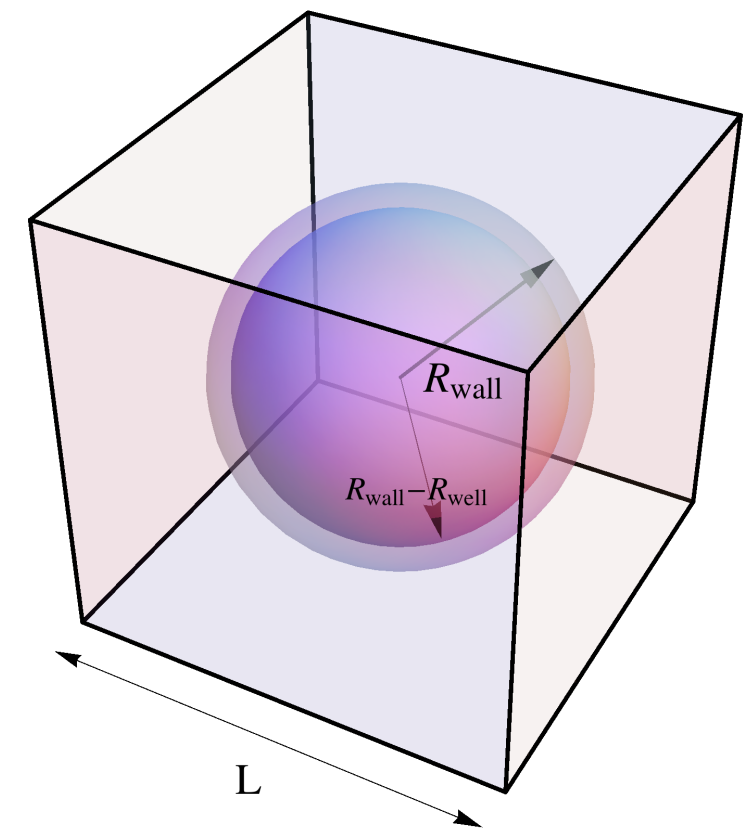


- one disadvantage of the spherical wall method are large R_{wall} and L for low energies

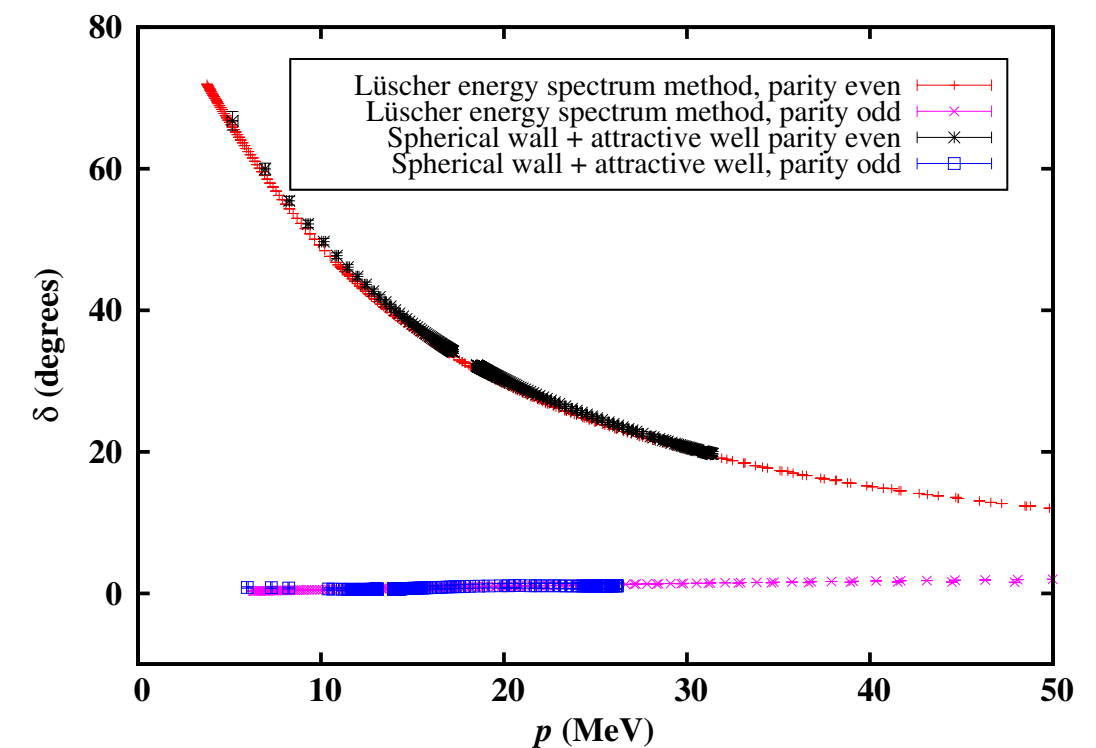
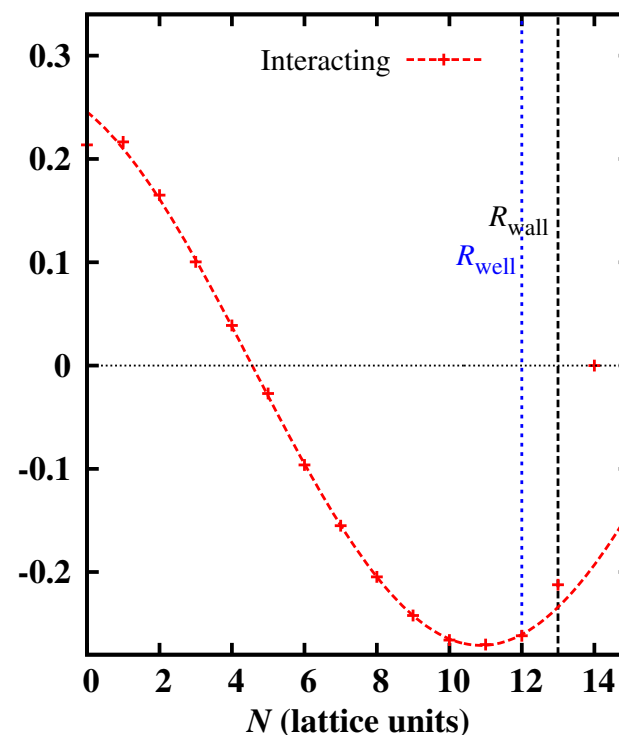
- one disadvantage of the spherical wall method are large R_{wall} and L for low energies
- additional attractive potential in front of the wall boundary to calculate phase shifts for low energies using small lattices



- one disadvantage of the spherical wall method are large R_{wall} and L for low energies
- additional attractive potential in front of the wall boundary to calculate phase shifts for low energies using small lattices



- particle-dimer in one dimension



Summary and Outlook



- The adiabatic Hamiltonian reduces in the asymptotic region to a simple cluster Hamiltonian in a position space basis

- The adiabatic Hamiltonian reduces in the asymptotic region to a simple cluster Hamiltonian in a position space basis
- Thus the asymptotic cluster wave functions have the same asymptotics like point-particle wave function.

- The adiabatic Hamiltonian reduces in the asymptotic region to a simple cluster Hamiltonian in a position space basis
- Thus the asymptotic cluster wave functions have the same asymptotics like point-particle wave function.
- The phase shifts can be directly extracted using simple wave function methods

- The adiabatic Hamiltonian reduces in the asymptotic region to a simple cluster Hamiltonian in a position space basis
- Thus the asymptotic cluster wave functions have the same asymptotics like point-particle wave function.
- The phase shifts can be directly extracted using simple wave function methods
- In future calculations very-high-precision energy calculations are not necessary (and difficult due to Monte Carlo techniques)

- The adiabatic Hamiltonian reduces in the asymptotic region to a simple cluster Hamiltonian in a position space basis
- Thus the asymptotic cluster wave functions have the same asymptotics like point-particle wave function.
- The phase shifts can be directly extracted using simple wave function methods
- In future calculations very-high-precision energy calculations are not necessary (and difficult due to Monte Carlo techniques)

Thank you for your attention!