

HADRONIC UNCERTAINTIES IN ΔA_{CP}

Luca Silvestrini
INFN, Rome

Introduction

Isospin amplitudes from BR's

Unitarity constraints & CP asymmetries

How large can penguins be?

Conclusions

INTRODUCTION

- Some basic facts known for a long time:
- To obtain a good description of $SCS D BR's$ need:
 - final state interactions and corrections to factorization
 - sizable $SU(3)$ breaking
- The SM expectation for direct CPV is $\lesssim 10^{-3}$

See for example Buccella et al. '95

INTRODUCTION II

- Can we envisage a mechanism to enhance the SM prediction for CPV by one order of magnitude to reproduce the exp result

$$\Delta a_{CP}^{\text{dir}} = a_{CP}^{\text{dir}}(K^+K^-) - a_{CP}^{\text{dir}}(\pi^+\pi^-) = (-6.6 \pm 1.6) 10^{-3}?$$

- Can anything analogous to the $\Delta I=1/2$ rule take place in SCS charm decays?

Golden & Grinstein, '89; Brod, Kagan & Zupan '11; Pirtskhalava & Uttayarat '11; Bhattacharya, Gronau & Rosner '12; Cheng & Chiang '12; Brod, Grossman, Kagan & Zupan '12

ISOSPIN & UNITARITY

- Let us start from the basic knowledge:
 - $SU(3)$ breaking is large \Rightarrow use only isospin
 - corrections to factorization are large \Rightarrow use a general parameterization
 - final state interactions are important \Rightarrow implement unitarity & external info on rescattering

Franco, Mishima & LS '12

ISOSPIN AMPLITUDES

$$A(D^+ \rightarrow \pi^+ \pi^0) = \frac{\sqrt{3}}{2} \mathcal{A}_2^\pi, \quad r_{\text{CKM}} = 6.4 \cdot 10^{-4}$$

$$A(D^0 \rightarrow \pi^+ \pi^-) = \frac{\mathcal{A}_2^\pi - \sqrt{2}(\mathcal{A}_0^\pi + ir_{\text{CKM}} \mathcal{B}_0^\pi)}{\sqrt{6}},$$

$$A(D^0 \rightarrow \pi^0 \pi^0) = \frac{\sqrt{2} \mathcal{A}_2^\pi + \mathcal{A}_0^\pi + ir_{\text{CKM}} \mathcal{B}_0^\pi}{\sqrt{3}},$$

A CP-even
B CP-odd

$$A(D^+ \rightarrow K^+ \bar{K}^0) = \frac{\mathcal{A}_{13}^K}{2} + \mathcal{A}_{11}^K + ir_{\text{CKM}} \mathcal{B}_{11}^K,$$

$$A(D^0 \rightarrow K^+ K^-) = \frac{-\mathcal{A}_{13}^K + \mathcal{A}_{11}^K - \mathcal{A}_0^K + ir_{\text{CKM}} \mathcal{B}_{11}^K - ir_{\text{CKM}} \mathcal{B}_0^K}{2},$$

$$A(D^0 \rightarrow K^0 \bar{K}^0) = \frac{-\mathcal{A}_{13}^K + \mathcal{A}_{11}^K + \mathcal{A}_0^K + ir_{\text{CKM}} \mathcal{B}_{11}^K + ir_{\text{CKM}} \mathcal{B}_0^K}{2}.$$

NUMERICAL RESULTS FROM BR's

$$|\mathcal{A}_2^\pi| = (3.08 \pm 0.08) \times 10^{-7} \text{ GeV},$$

$$|\mathcal{A}_0^\pi| = (7.6 \pm 0.1) \times 10^{-7} \text{ GeV},$$

$$\arg(\mathcal{A}_2^\pi/\mathcal{A}_0^\pi) = (\pm 93 \pm 3)^\circ.$$

No $\Delta I=1/2$ rule for D decays, large strong phases

$$|\mathcal{A}_{13}^K - \mathcal{A}_{11}^K - \mathcal{A}_0^K| = (5.0 \pm 0.4) \times 10^{-7} \text{ GeV}$$

Should vanish in the $SU(3)$ limit, but is $O(1)!!$

UNITARITY CONSTRAINTS

$$S = \left(\begin{array}{c|ccc} D \rightarrow D & D \rightarrow \pi\pi & D \rightarrow KK & \dots \\ \hline \pi\pi \rightarrow D & \pi\pi \rightarrow \pi\pi & \pi\pi \rightarrow KK & \dots \\ KK \rightarrow D & KK \rightarrow \pi\pi & KK \rightarrow KK & \dots \\ \vdots & \vdots & \vdots & \ddots \end{array} \right) \equiv \begin{pmatrix} 1 & -i(T)^T \\ -i \text{CP}(T) & S_S \end{pmatrix}$$

implies

$$T^R = S_S (T^R)^*, \quad T^I = S_S (T^I)^*$$

Unfortunately, exp data at the D mass are ambiguous. If unitarity is saturated by the $\pi\pi$ and KK channels alone, then constraints on $\Delta a_{\text{CP}}^{\text{dir}}$ can be derived, otherwise not.

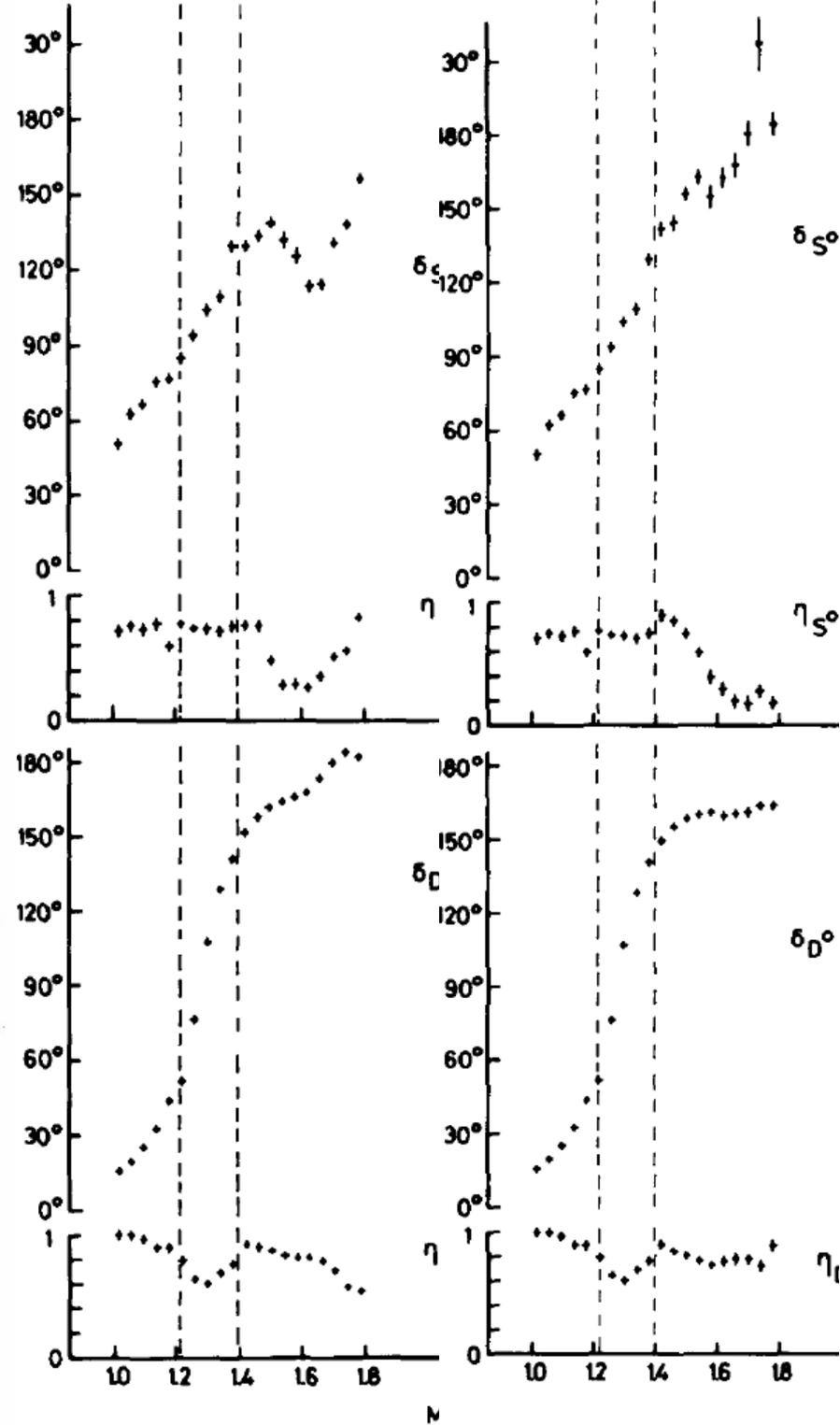
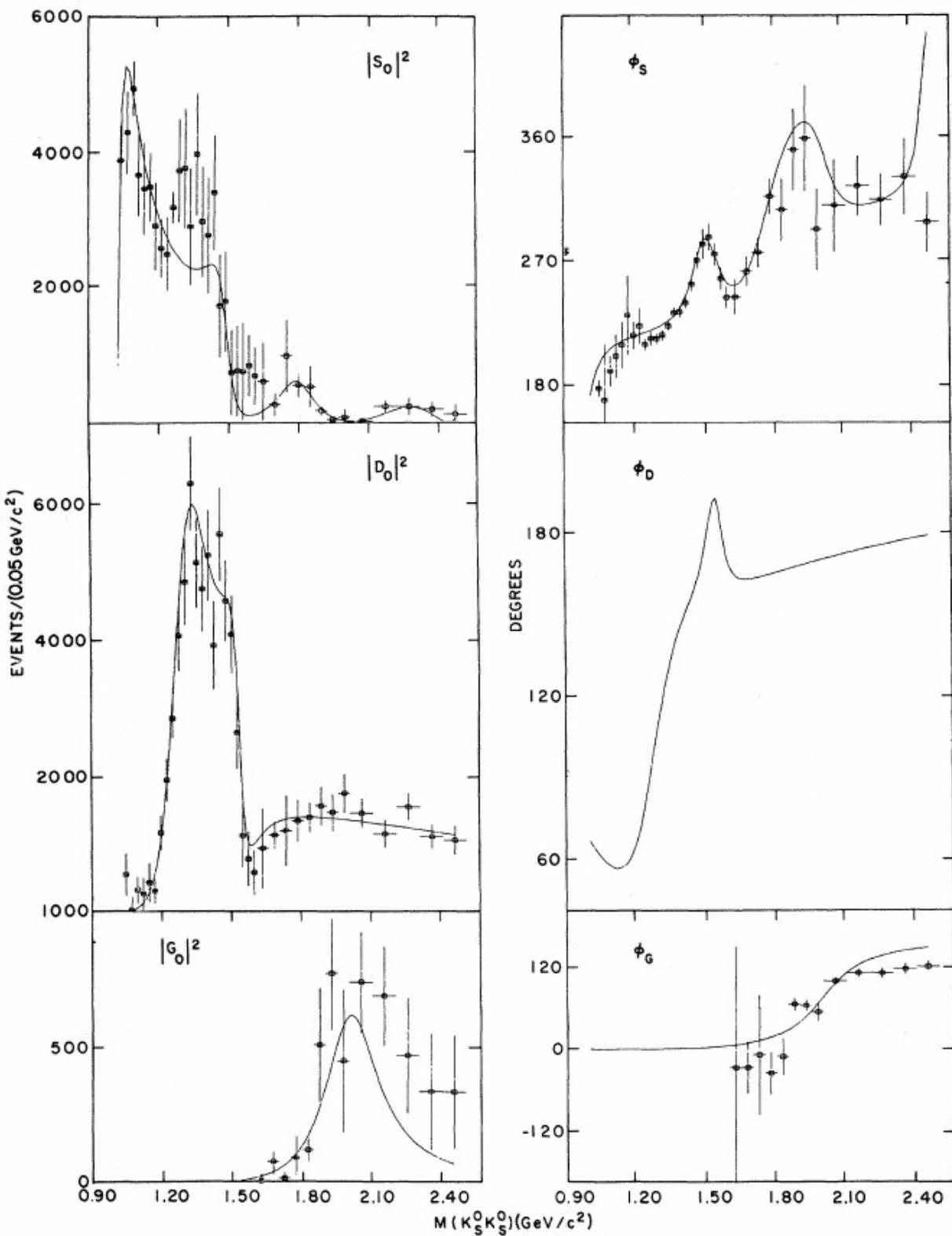


Fig. 4. S

Fig. 3.

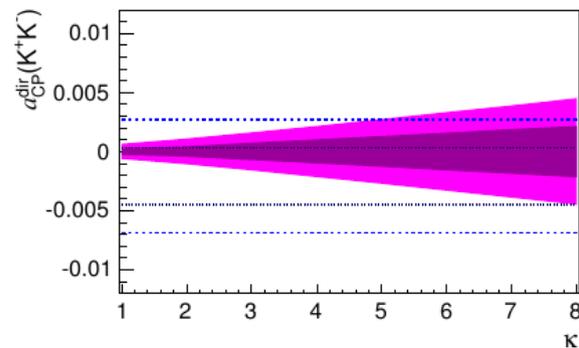
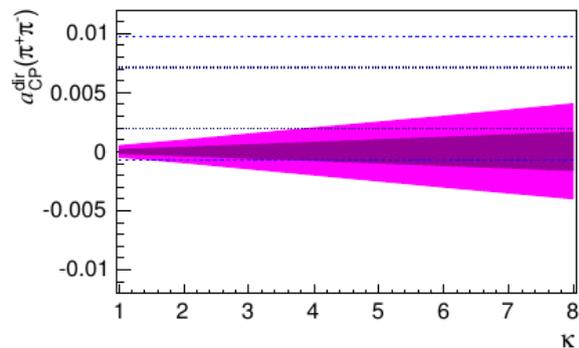
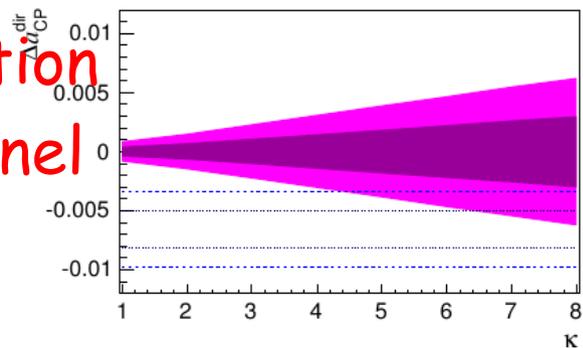
CP ASYMMETRIES

- One can study the CP asymmetries as a function of the upper bound on the size of CPV contributions in the two- and three-channel scenarios. We write

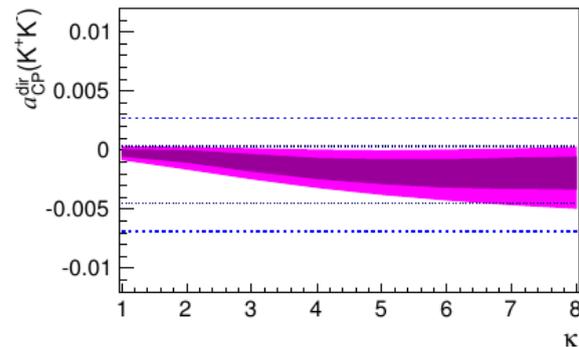
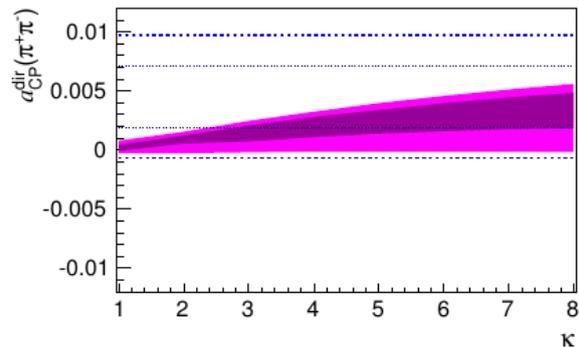
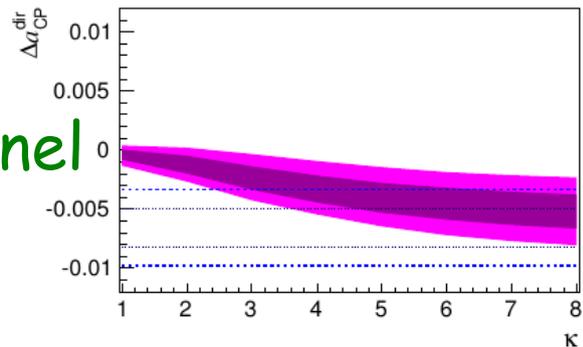
$$\begin{aligned} |\mathcal{B}_0^\pi| &< \kappa |\mathcal{A}_0^\pi|, \\ |\mathcal{B}_0^K - \mathcal{A}_0^K| &< \kappa |\mathcal{A}_0^K|, \\ |\mathcal{B}_{11}^K - (\mathcal{A}_{11}^K - \mathcal{A}_{13}^K)| &< \kappa |\mathcal{A}_{11}^K - \mathcal{A}_{13}^K|, \end{aligned}$$

and consider **predictions** and **fit results** for CP asymmetries

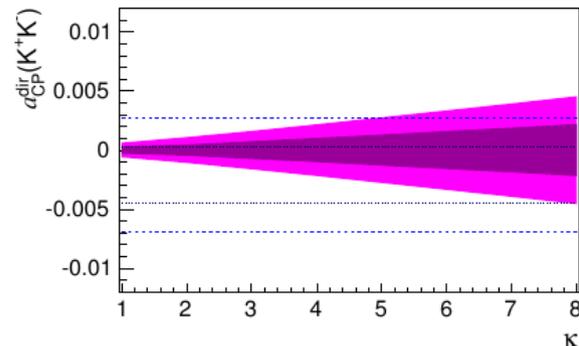
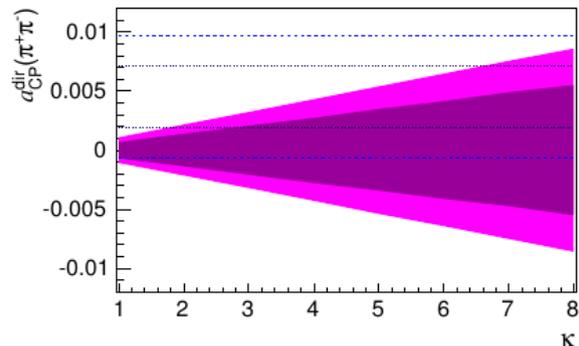
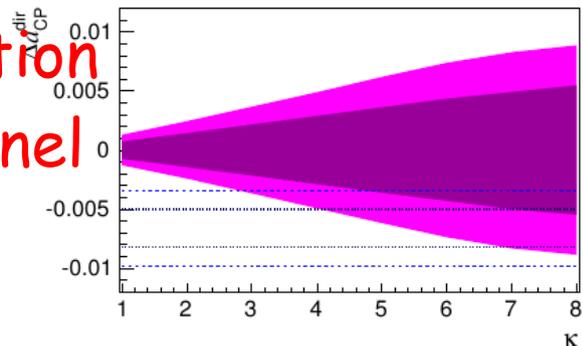
Prediction
2-channel



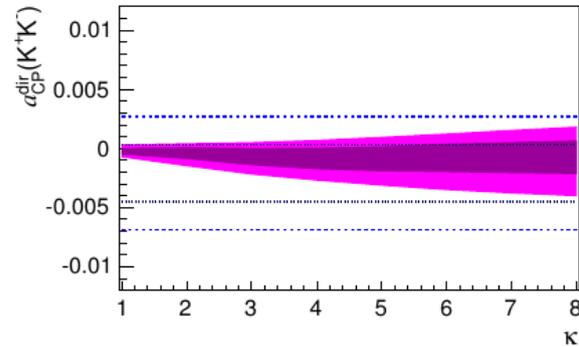
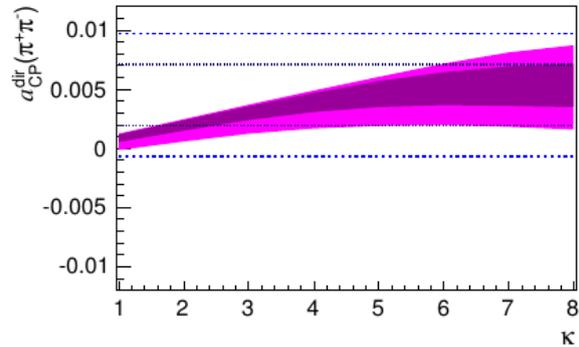
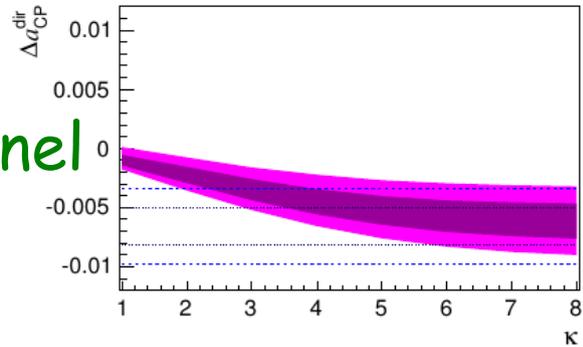
Fit
2-channel



Prediction
3-channel



Fit
3-channel



SuperI

CONCLUSIONS FROM UNITARITY

- The prediction does not reach the exp value within 2σ even for $\kappa=8$ in the 2-channel case
- Without unitarity constraints, the prediction reaches the exp value at the 2σ level for $\kappa>5$, but even for $\kappa=8$ it is still 1σ below
- How large can κ be?
 - translate fit results into RGI parameters
 - compare with K and B

FROM ISOSPIN AMPLITUDES TO RGI PARAMETERS

- The BR fit results can be translated into results for RGI parameters (aka topologies).
Neglecting for simplicity $O(1/N_c^2)$ terms:

$$E_1(\pi) + E_2(\pi) = (1.72 \pm 0.04) \times 10^{-6} e^{i\delta} \text{ GeV},$$

$$E_1(\pi) + A_2(\pi) - P_1^{\text{GIM}}(\pi) = (2.10 \pm 0.02) \times 10^{-6} e^{i(\delta \pm (71 \pm 3)^\circ)} \text{ GeV},$$

$$E_2(\pi) - A_2(\pi) + P_1^{\text{GIM}}(\pi) = (2.25 \pm 0.07) \times 10^{-6} e^{i(\delta \mp (62 \pm 2)^\circ)} \text{ GeV}$$

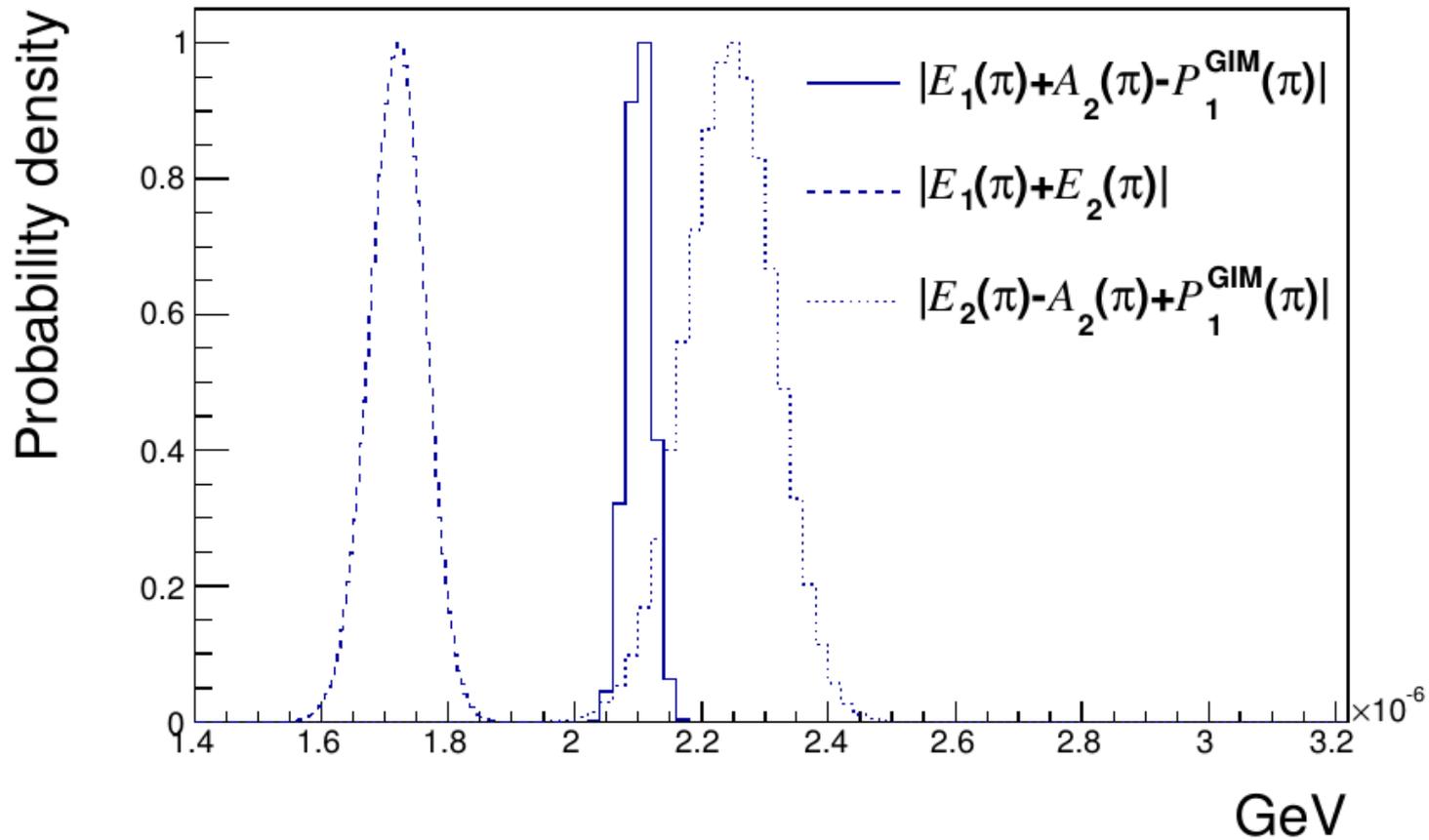
- E_1 does not dominate the amplitudes \Rightarrow we are away from the infinite mass limit
- All amplitudes of same size, w. large phases

THE MEANING OF κ

- The condition $|\mathcal{B}_0^\pi| < \kappa |\mathcal{A}_0^\pi|$, means

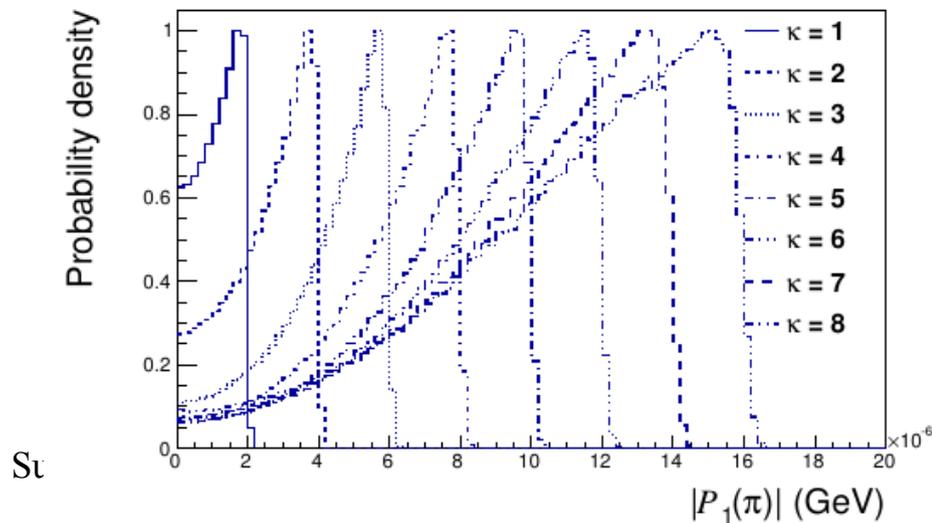
$$|P_1(\pi)| \leq \kappa \left| \frac{2}{3} E_1(\pi) - \frac{1}{3} E_2(\pi) + A_2(\pi) - P_1^{\text{GIM}}(\pi) \right|$$

- κ is the ratio of $|P_1|$ over all other topologies
- Notice that $P_1 \sim P_b - P_s$ while $P_1^{\text{GIM}} \sim P_d - P_s$
- How large should $|P_1|$ be to reproduce $\Delta a_{\text{CP}}^{\text{dir}}$?

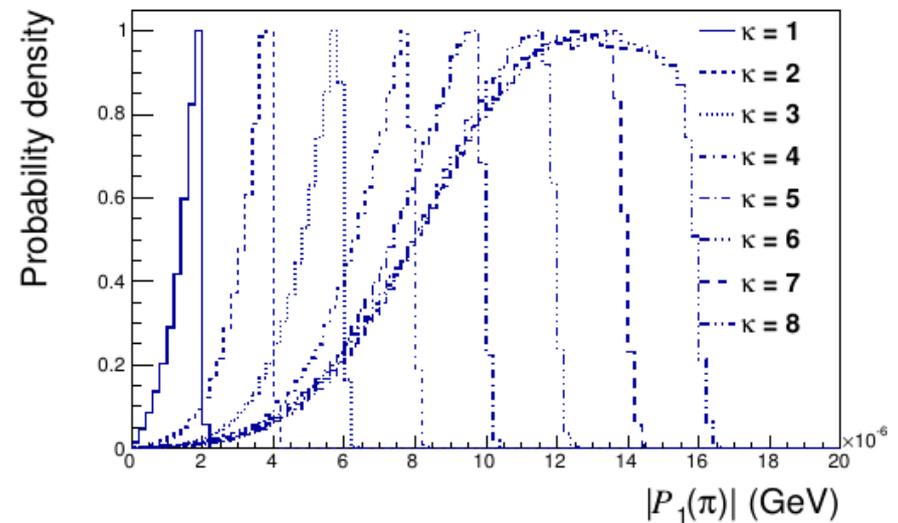


2-channel

3-channel



St



DYNAMICAL ARGUMENTS

- The amplitudes for K , D and $B \rightarrow \pi\pi$ are formally the same, with the obvious flavour and CKM replacements.

- In the Kaon system, one has

$$(P_u - P_c) \sim 3 (P_+ - P_c) \sim 25 (E_1 + E_2)$$

- No enhancement expected for P_u (will be checked on the lattice soon), while P_c and P_+ generate local operators with chirally enhanced matrix elements (SVZ)

DYNAMICAL ARGUMENTS II

- In charm decays, no chiral enhancement is present, so that one expects

$$|P_1| = |P_b - P_s| \leq |E_1|, |E_2|, |A_1|, |A_2|, |P_1^{GIM}|$$

i.e. $\kappa \leq 1$.

- In B decays one is much closer to the infinite mass limit so that $|E_1|$ and $|E_2|$ dominate, with all other contractions power suppressed.

CONCLUSIONS

- We have performed a phenomenological analysis of $\Delta a_{CP}^{\text{dir}}$ with minimal assumptions as a function of $\kappa \sim$ relative size of $|P_1|$
- From the BR we confirm large nonfactorizable contributions, large strong phases and large $SU(3)$ breaking
- Unfortunately, data on $\pi\pi$ scattering at the D mass are not able to fully determine the relevant FSI

CONCLUSIONS II

- In the most conservative scenario (no constraints from unitarity), values of $\kappa > 5$ are needed to reach at 2σ the experimental result
- We cannot find any reasonable dynamical origin for such a large value of $|P_1|$
- If the central value stays with improved errors, we have strong indications of NP