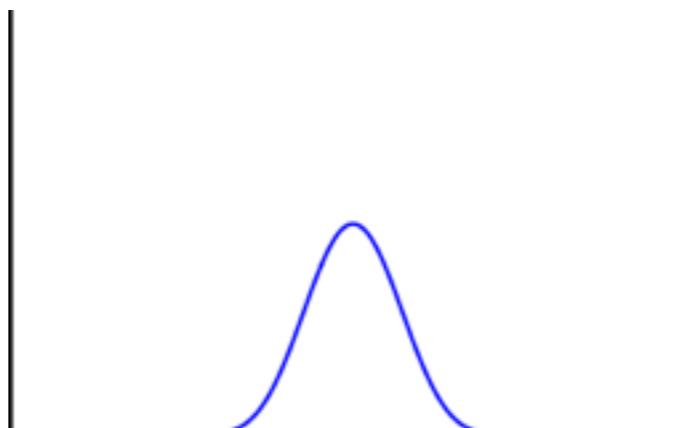


Turbulent strings in AdS/CFT

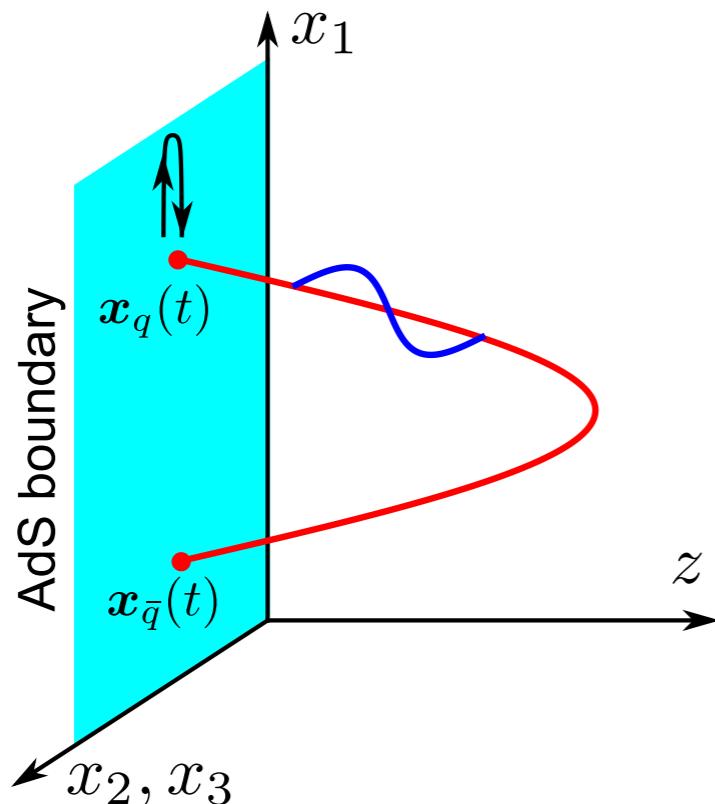
Takaaki Ishii
(University of Crete)

arXiv:1504.02190 with Keiju Murata



Plan

Perturb holographic quark-antiquark potential



Motivations

- AdS turbulence
- Turbulent instability on D7
[Hashimoto-Kinoshita-Oka-Murata]

We solve nonlinear time evolution

c.f.) Cosmic strings in flat space

Contents

1. Introduction
2. Review of the static solution
3. Numerical setup
4. Results
5. Summary

Time-like holographic Wilson loop

[Maldacena, Rey-Yee]

In $\text{AdS}_5 \times \text{S}^5$ $ds^2 = \frac{\ell^2}{z^2} (-dt^2 + dz^2 + d\mathbf{x}^2) + \ell^2 d\Omega_5^2$

Static gauge: $(\tau, \sigma) = (t, z)$

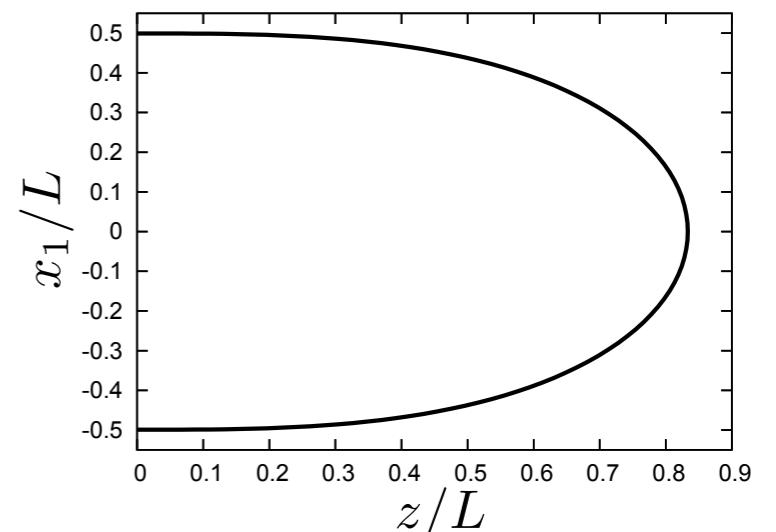
Target space embedding: $x_1 = X_1(z)$

Solution for separation L

$$\begin{aligned} X_1(z) &= \pm z_0 \int_{z/z_0}^1 dw \frac{w^2}{\sqrt{1 - w^4}} \\ &= \pm z_0 [\Gamma_0 + F(z/z_0; i) - E(z/z_0; i)] \end{aligned}$$

\nearrow
 z_0 : string tip
 $\Gamma_0 = 0.599$

$$\frac{L}{2} = z_0 \Gamma_0$$

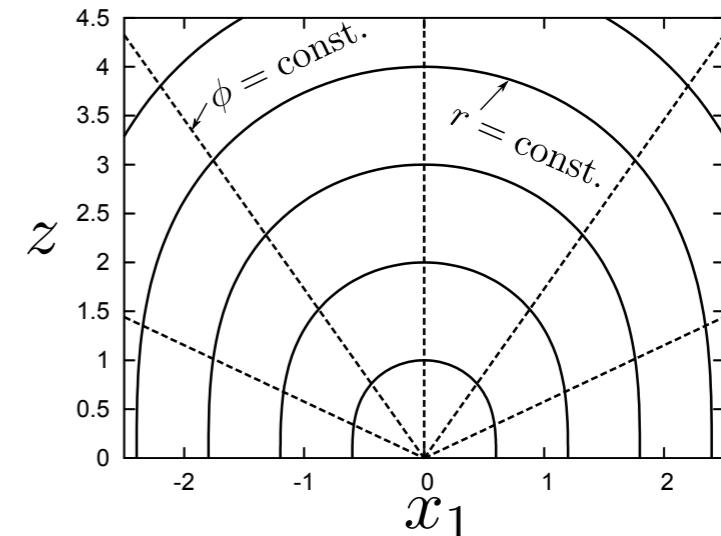


A convenient parametrization

Polar-like coordinates (r, ϕ) in which the static solution is $r=z_0$

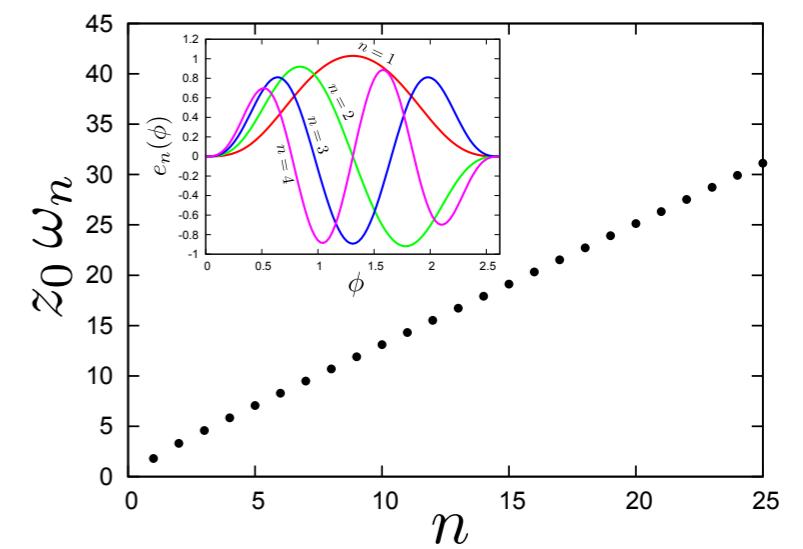
$$z = r f(\phi) = r \operatorname{sn}(\phi; i)$$

$$x_1 = rg(\phi) = r \begin{cases} \phi - E(\operatorname{sn}(\phi; i); i) + \Gamma_0 & (\phi \leq \beta_0/2) \\ \phi + E(\operatorname{sn}(\phi; i); i) - \Gamma_0 - \beta_0 & (\phi > \beta_0/2) \end{cases}$$



We prepare eigenvalues/functions in linearized perturbations

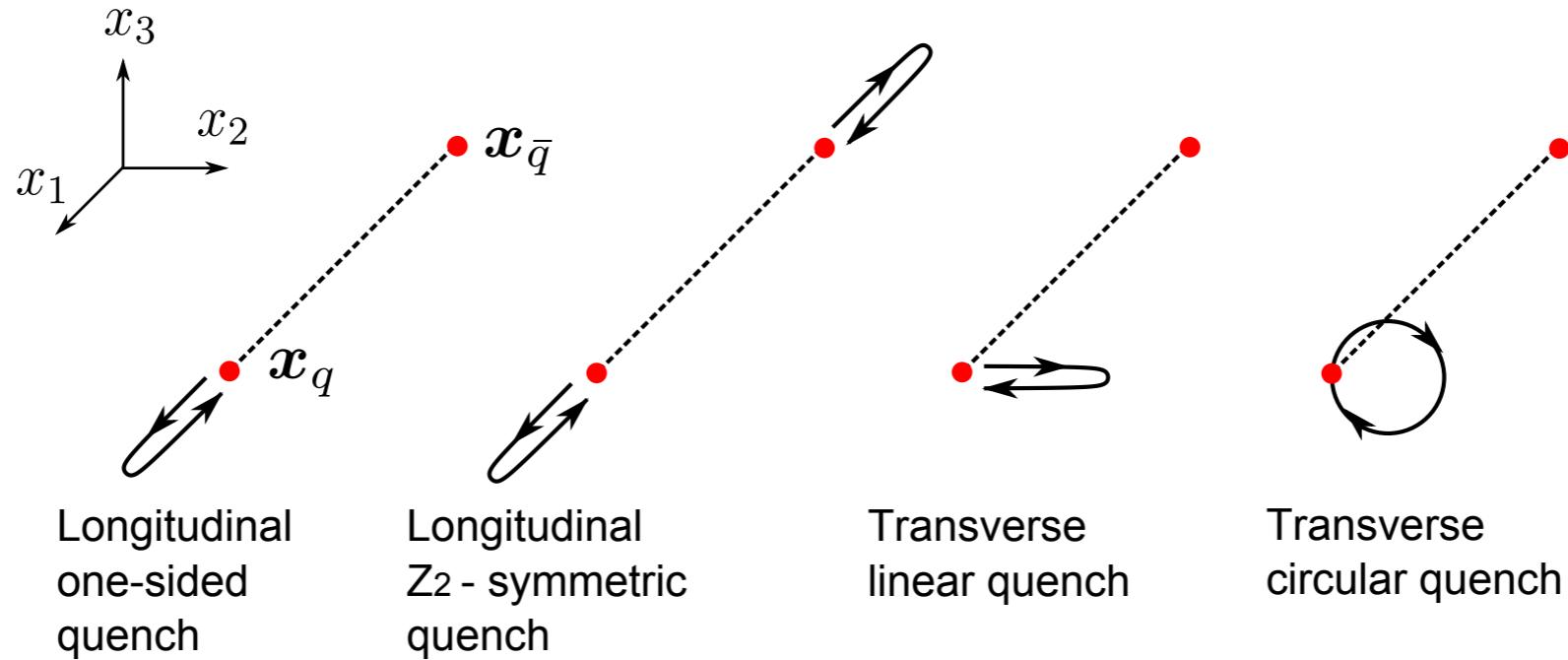
[Callan-Guijosa, Klebanov-Maldacena-Thorn]



Contents

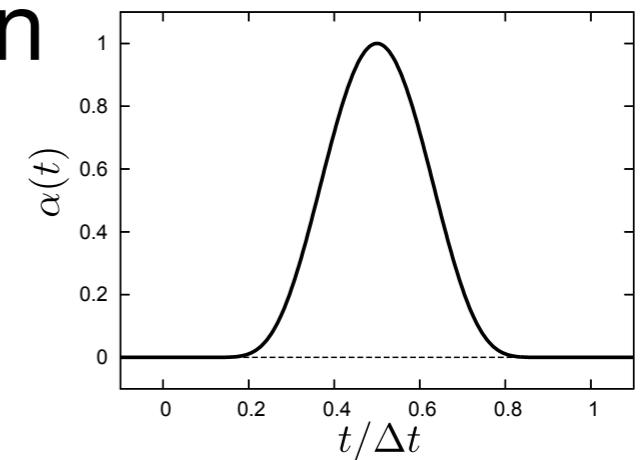
1. Introduction
2. Review of the static solution
3. Numerical setup
4. Results
5. Summary

Perturb the string endpoints



Quench profile: a compact C^∞ function

$$\alpha(t) = \exp \left[2 \left(\frac{\Delta t}{t - \Delta t} - \frac{\Delta t}{t} + 4 \right) \right] \quad (0 < t < \Delta t)$$



Worldsheet double null coordinates

Induced metric $ds_{F1}^2 = -2\gamma_{uv}dudv$

Worldsheet: u, v

Target space: $T(u, v)$, $Z(u, v)$, $X_{1,2,3}(u, v)$

$$\gamma_{uv} = \frac{\ell^2}{Z^2}(-T_{,u}T_{,v} + Z_{,u}Z_{,v} + \mathbf{X}_{,u} \cdot \mathbf{X}_{,v})$$

Equations of motion

$$T_{,uv} = \frac{1}{Z}(T_{,u}Z_{,v} + Z_{,u}T_{,v})$$

$$Z_{,uv} = \frac{1}{Z}(T_{,u}T_{,v} + Z_{,u}Z_{,v} - \mathbf{X}_{,u} \cdot \mathbf{X}_{,v})$$

$$\mathbf{X}_{,uv} = \frac{1}{Z}(\mathbf{X}_{,u}Z_{,v} + Z_{,u}\mathbf{X}_{,v})$$

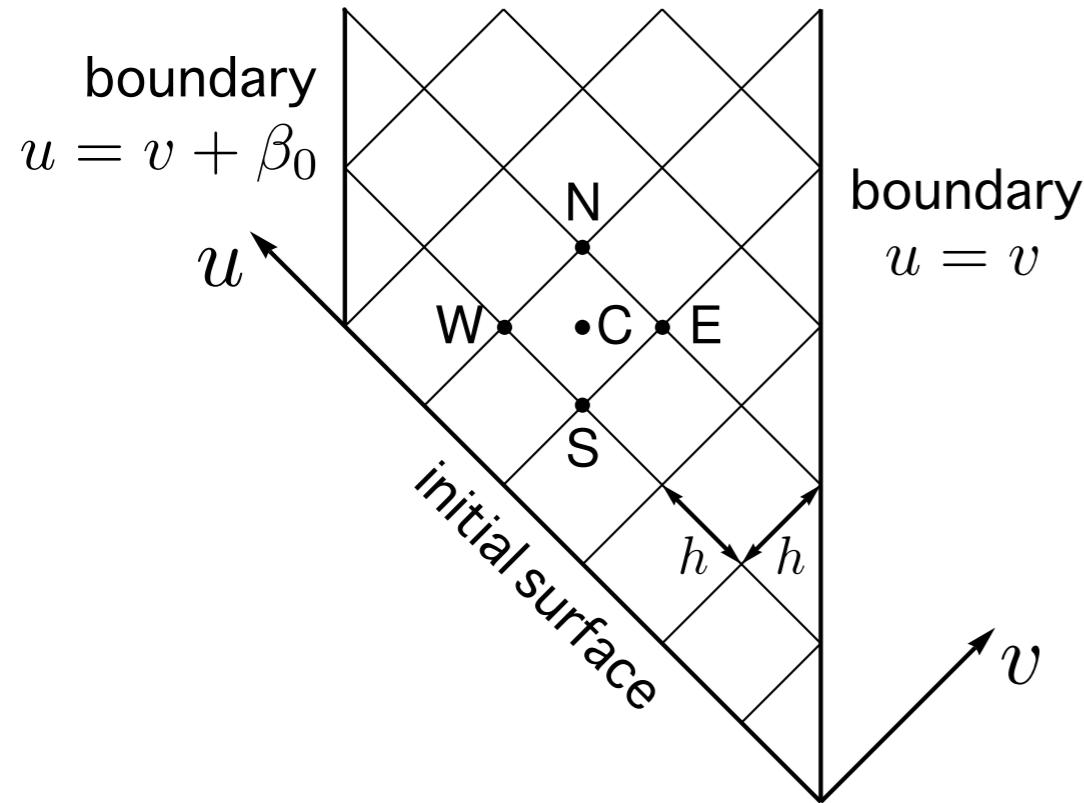
Constraints

$$\gamma_{uu} = \frac{\ell^2}{Z^2}(-T_{,u}^2 + Z_{,u}^2 + \mathbf{X}_{,u}^2) = 0$$

$$\gamma_{vv} = \frac{\ell^2}{Z^2}(-T_{,v}^2 + Z_{,v}^2 + \mathbf{X}_{,v}^2) = 0$$

Discretization

To solve EoMs, we use $O(h^2)$ central finite differential



$$\Psi_{,uv}|_C = \frac{\Psi_N - \Psi_E - \Psi_W + \Psi_S}{h^2}$$

$$\Psi_{,u}|_C = \frac{\Psi_N - \Psi_E + \Psi_W - \Psi_S}{2h}$$

$$\Psi_{,v}|_C = \frac{\Psi_N + \Psi_E - \Psi_W - \Psi_S}{2h}$$

$$\Psi|_C = \frac{\Psi_E + \Psi_W}{2}$$

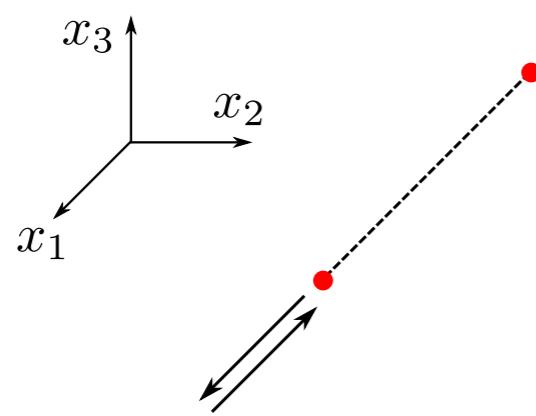
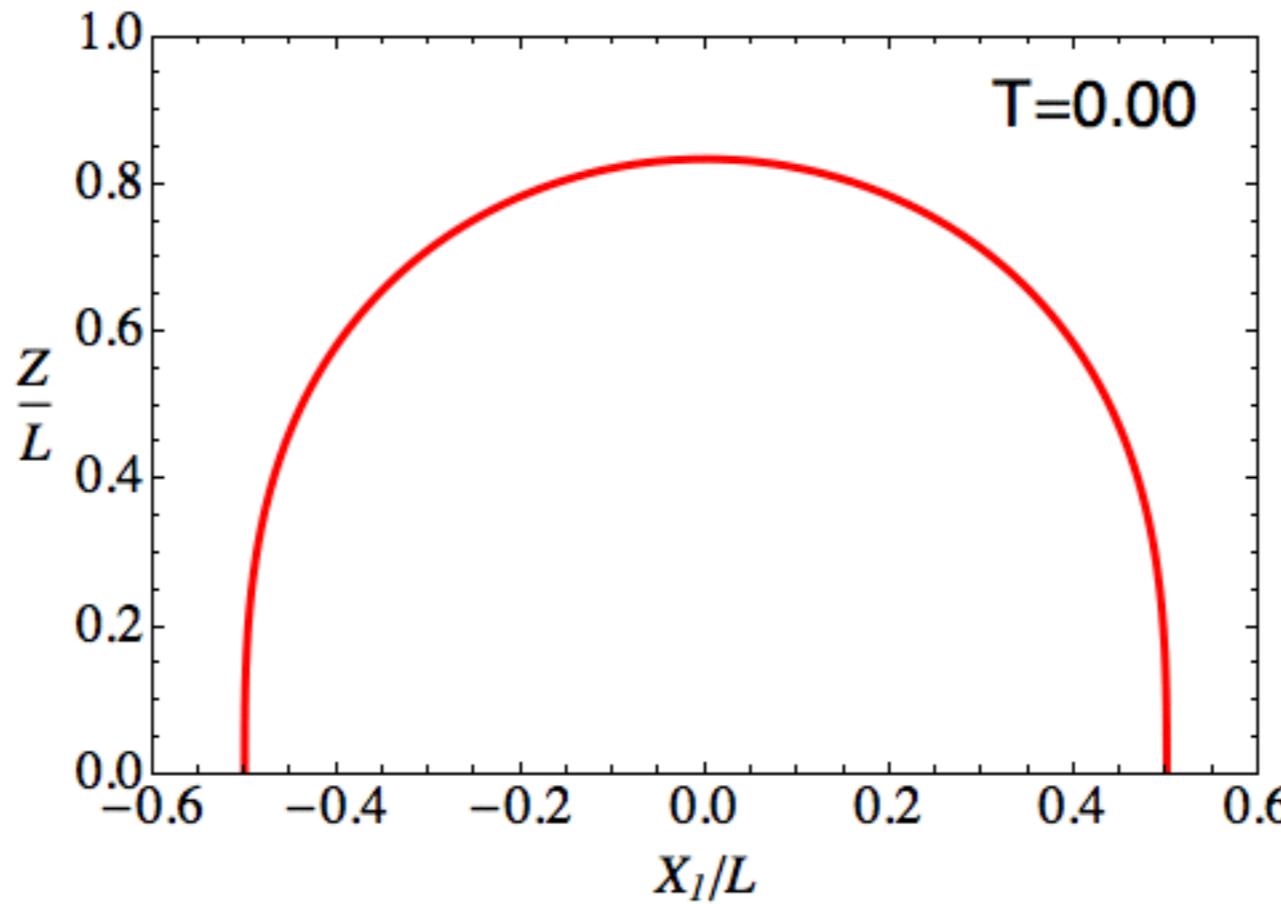
Compute N by using EWS data

Initial data ($v=0$): static solution

Contents

1. Introduction
2. Review of the static solution
3. Numerical setup
4. Results
5. Summary

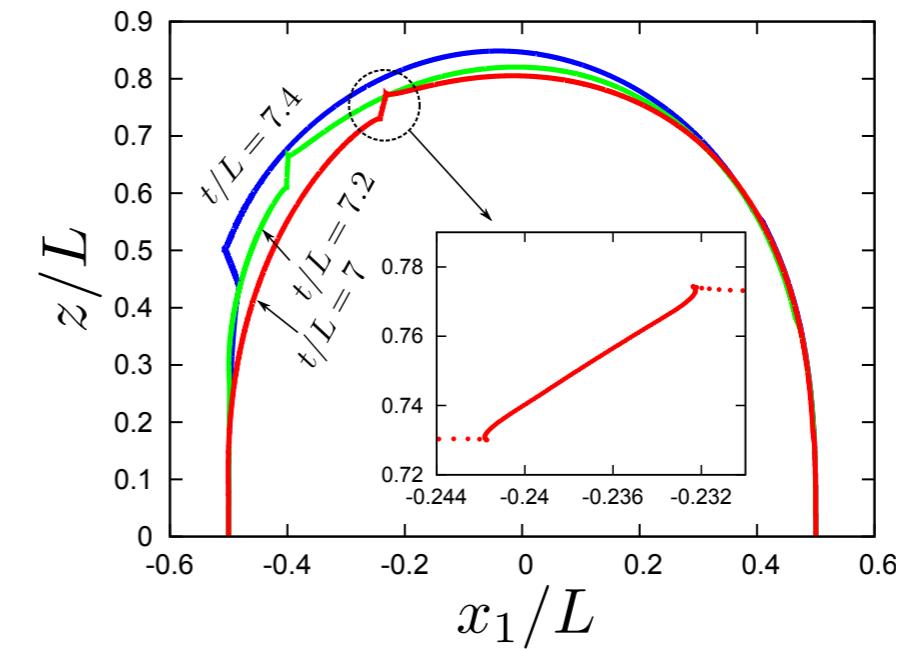
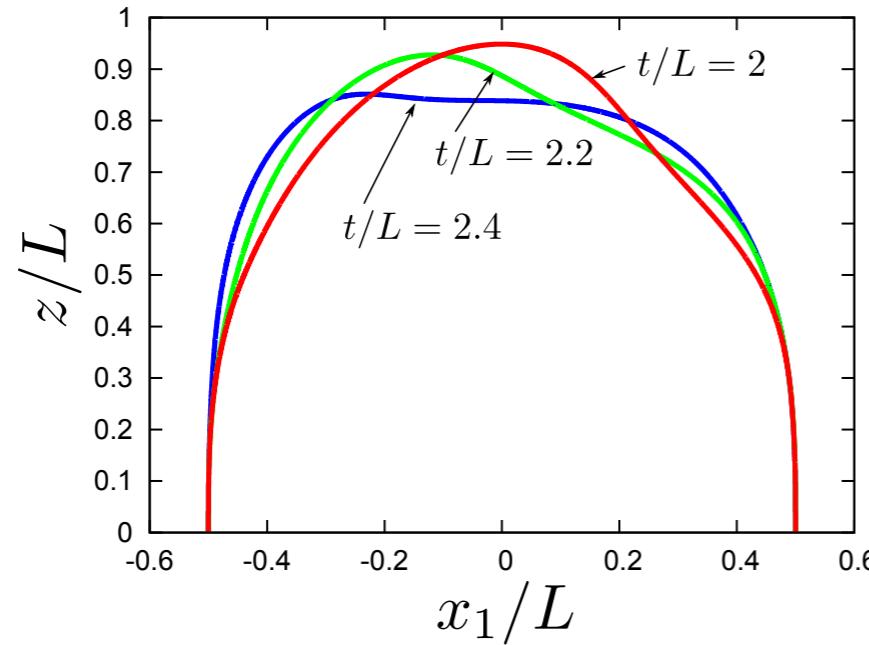
Longitudinal one-sided quench



$\varepsilon=0.03, \Delta t/L=2$

Amplitude: $\varepsilon=\Delta x/L$
Duration: $\Delta t/L$

Cusp formation



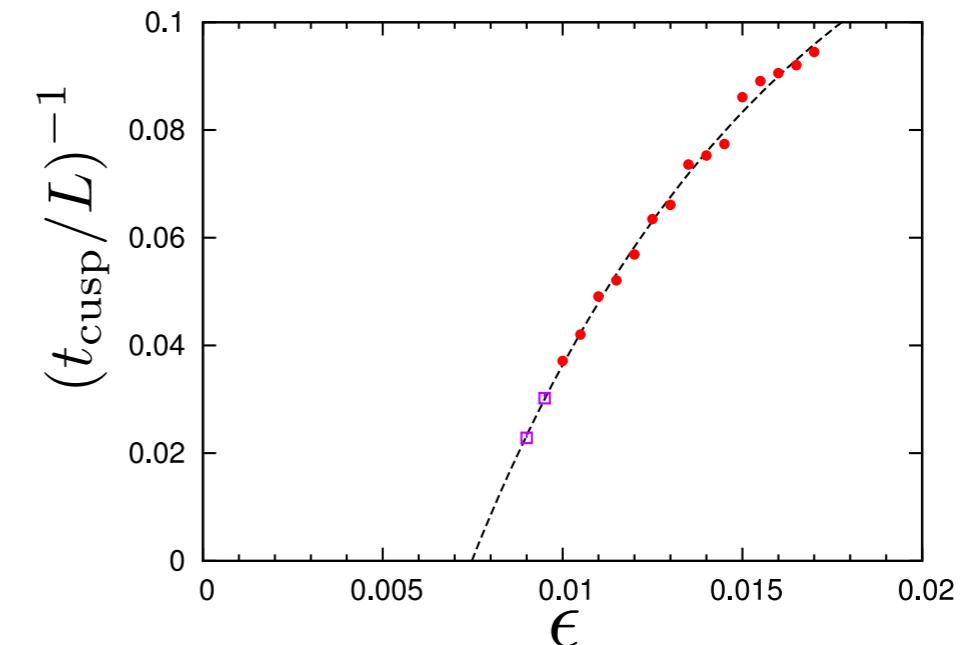
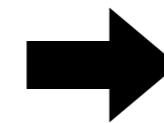
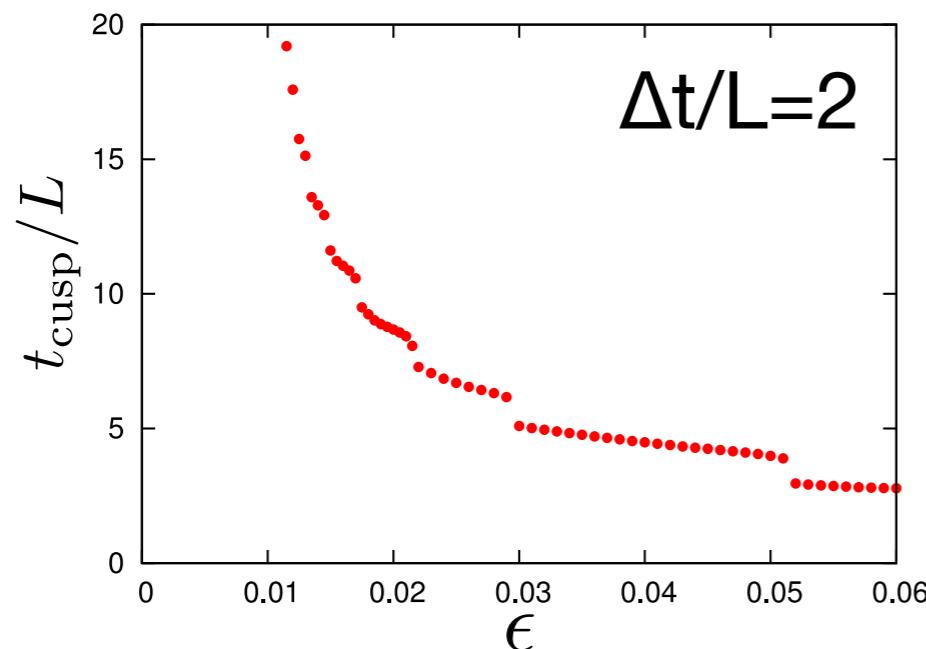
- Cusps are seen in target space (x,z) -coordinates
- Fields on worldsheet (u,v) -coordinates are regular
- Cusps are **created in a pair** (around $t/L \sim 5$)

Cusp detection

The conditions satisfied at a cusp:

$$J_z \equiv T_{,u}Z_{,v} - T_{,v}Z_{,u} = 0$$

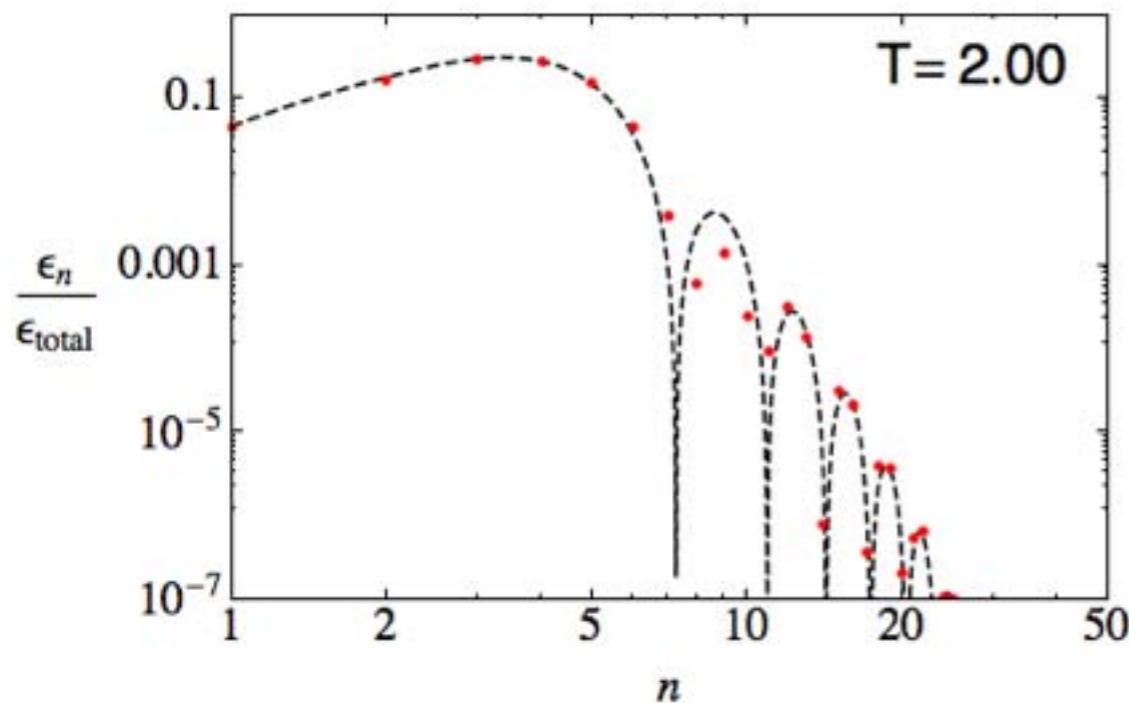
$$J_i \equiv T_{,u}X_{i,v} - T_{,v}X_{i,u} = 0$$



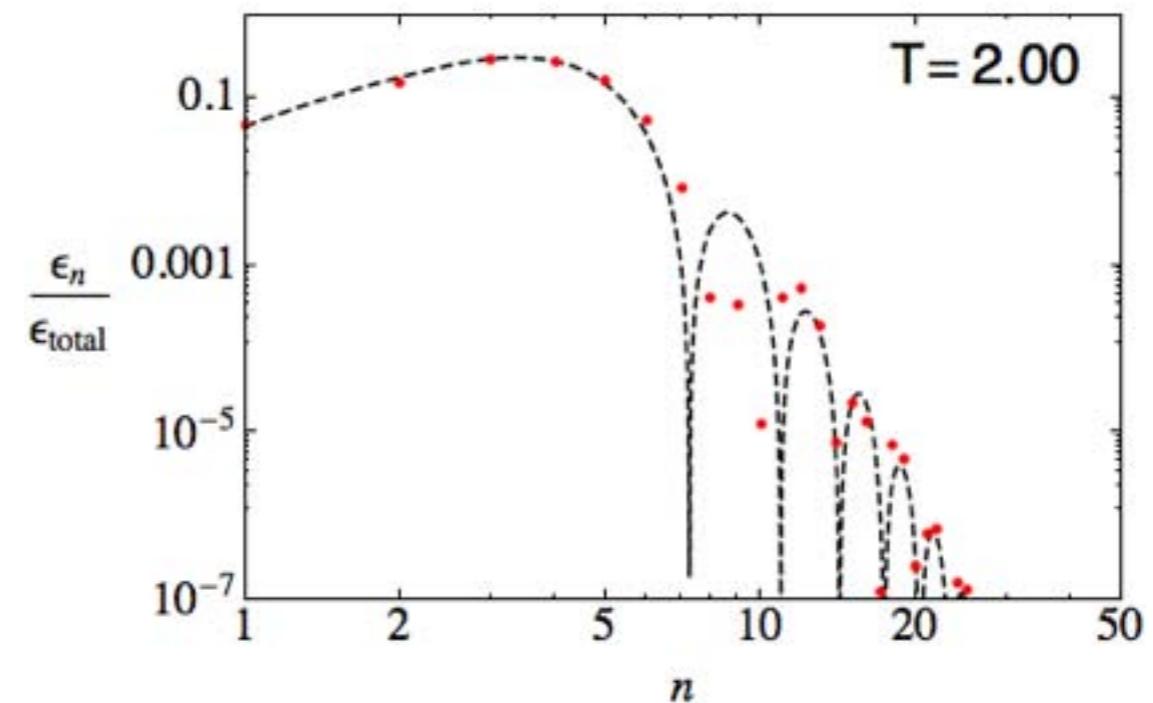
Longer cusp formation
time as ϵ is decreased

Minimal amplitude for
cusps: $\epsilon_{\text{crit}} \sim 0.075$

Energy spectrum (Log-log plot)



$\varepsilon=0.005, \Delta t/L=2$ (no cusp)



$\varepsilon=0.01$ (cusps $T \sim 27$)

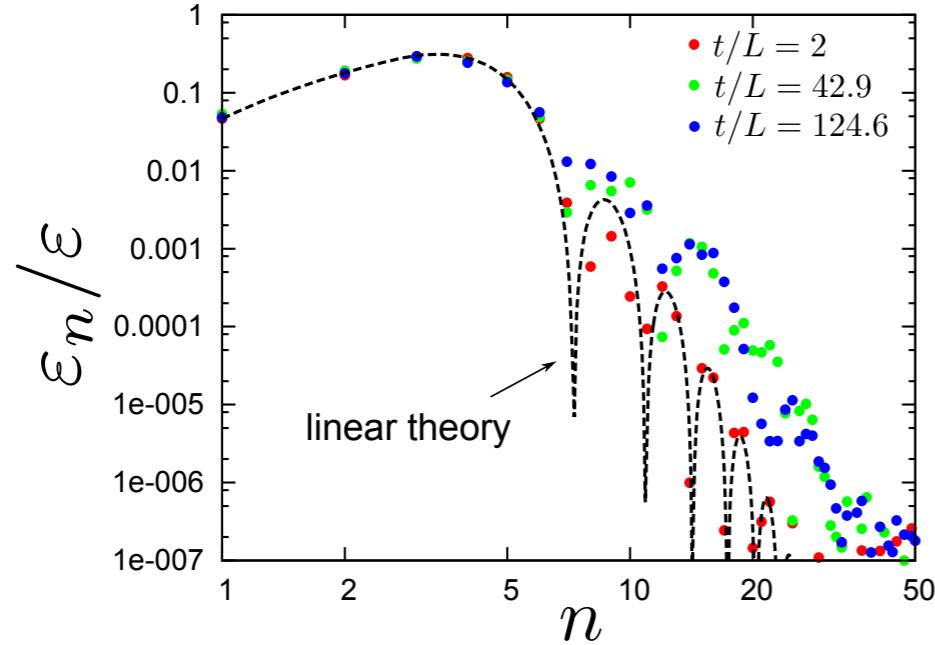
Decompose nonlinear solutions in linear eigenmodes e_n

$$\chi_1 = \sum_{n=1}^{\infty} c_n(t) e_n(\phi)$$

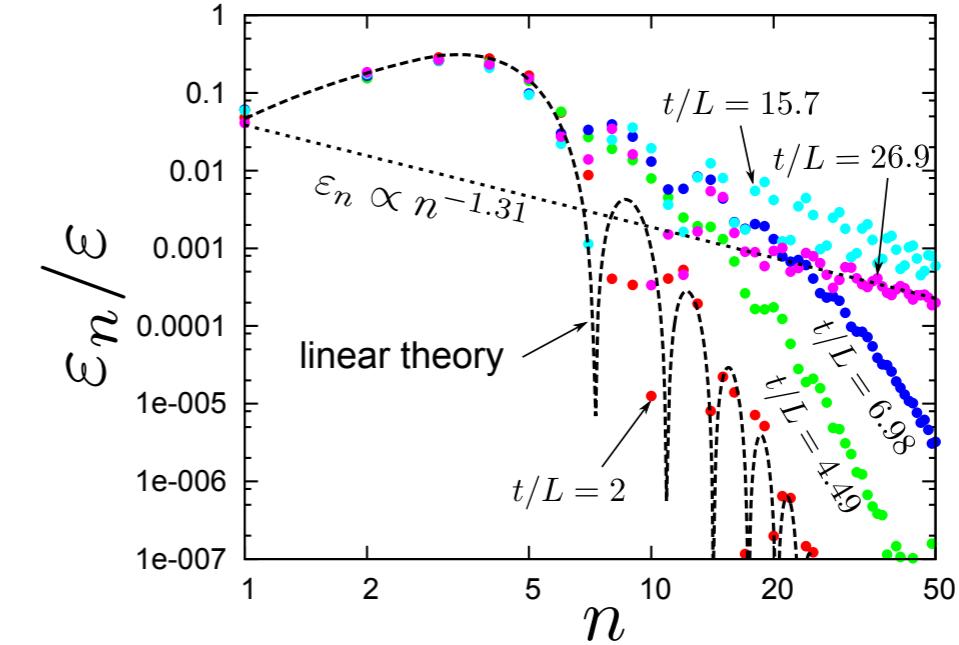
$$\varepsilon_n(t) = \frac{\sqrt{\lambda} z_0}{4\pi} (\dot{c}_n^2 + \omega_n^2 c_n^2)$$

***Dashed lines are in the linearized theory

Energy cascade



$\varepsilon=0.005, \Delta t/L=2$ (no cusp)



$\varepsilon=0.01$ (cusps $T \sim 27$)

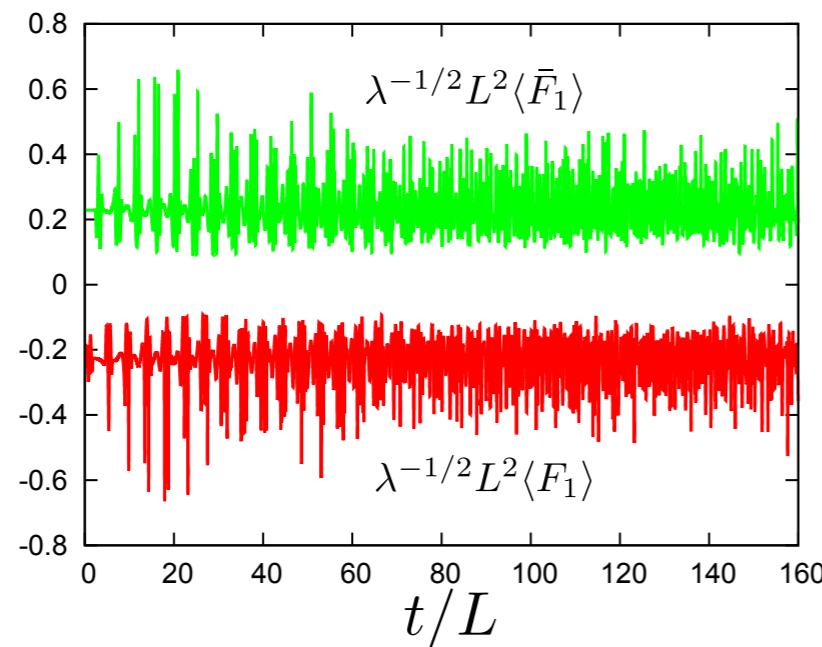
Cusp formation: **direct energy cascade** \rightarrow power law

No cusp: no power law

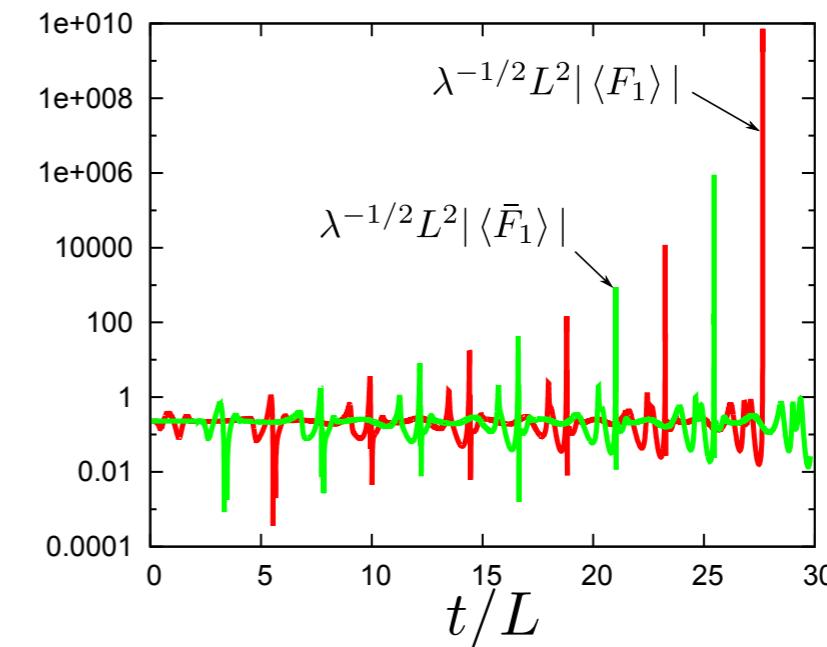
Forces on the endpoints

$$\langle \mathbf{F}(t) \rangle = \frac{\delta S_{\text{on-shell}}}{\delta \mathbf{x}_q}$$

Force diverges when a cusp reaches the boundary



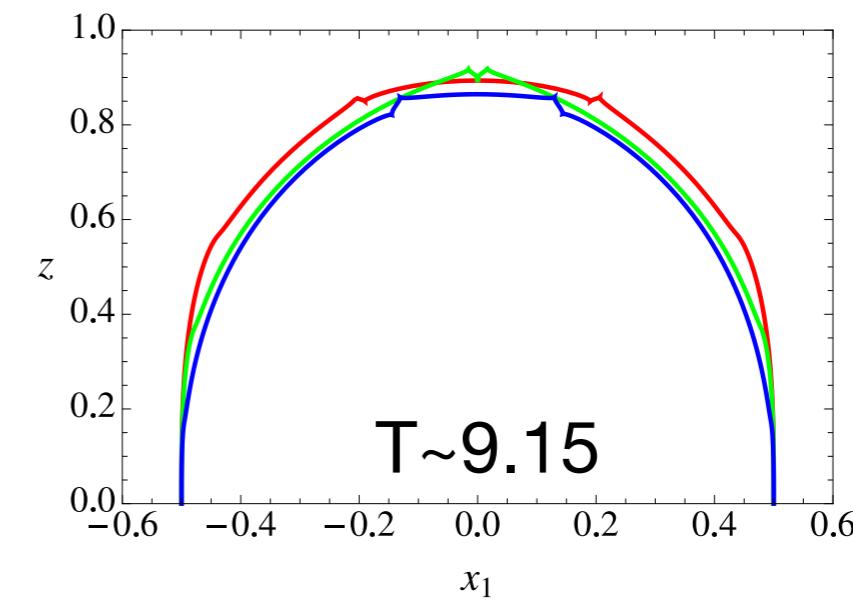
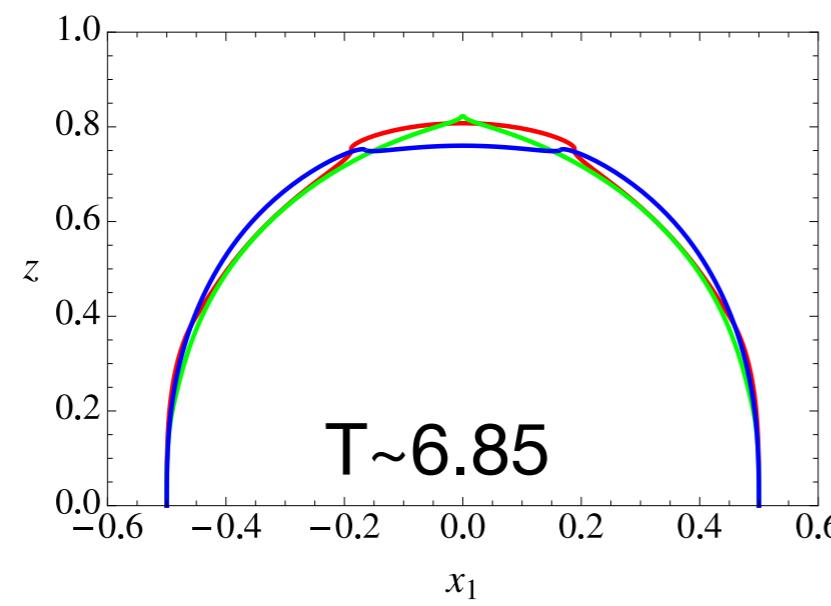
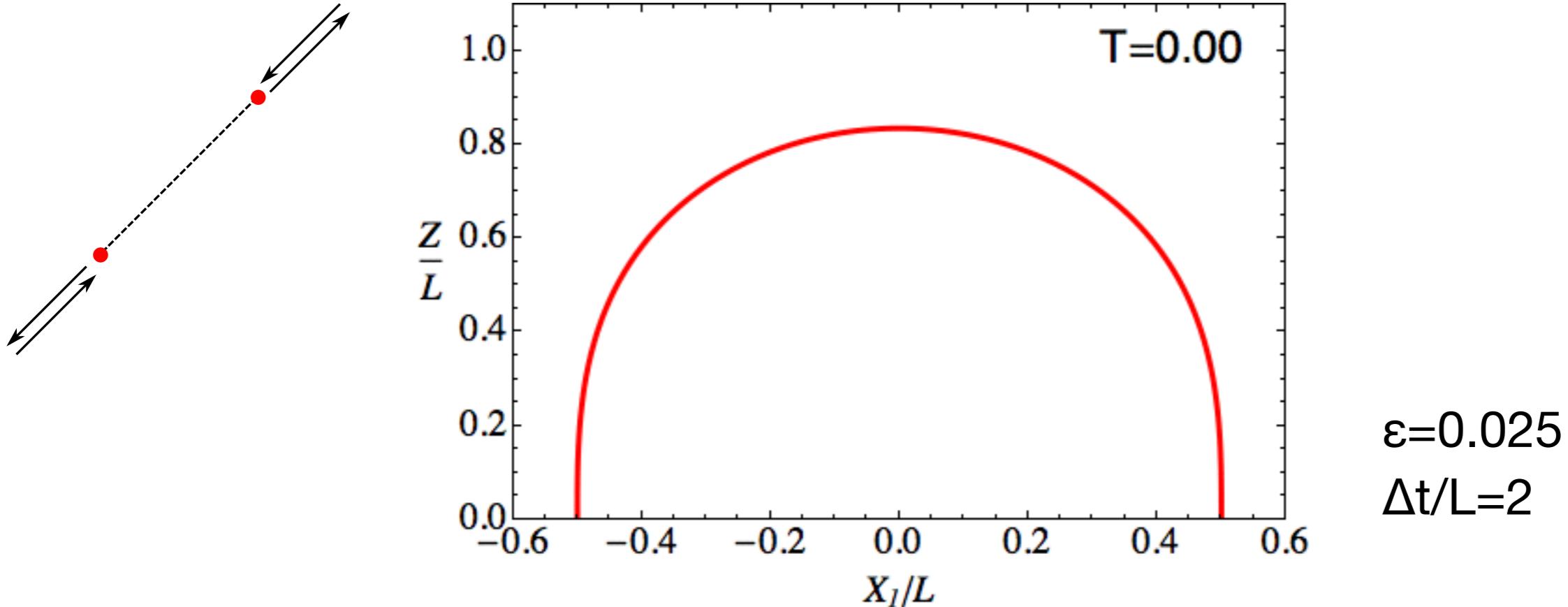
$\epsilon=0.005, \Delta t/L=2$ (no cusp)



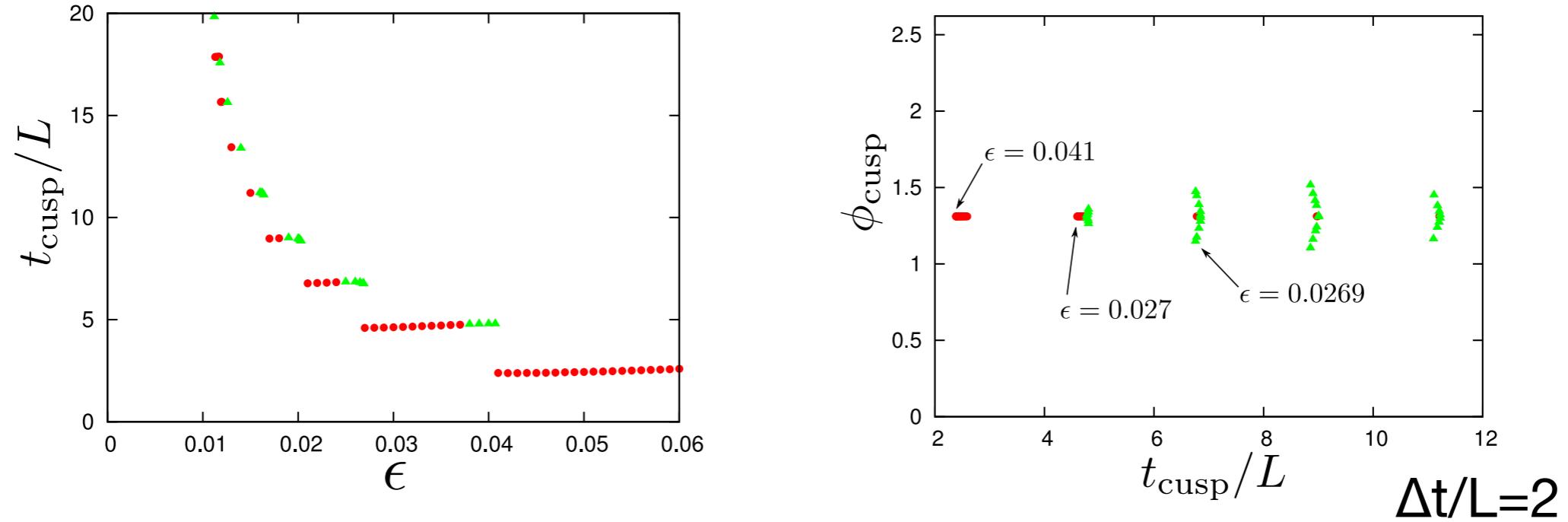
$\epsilon=0.01$ (cusps $T \sim 27$)

***Red: $x=L/2$, green: $x=-L/2$

Z_2 -symmetric quench

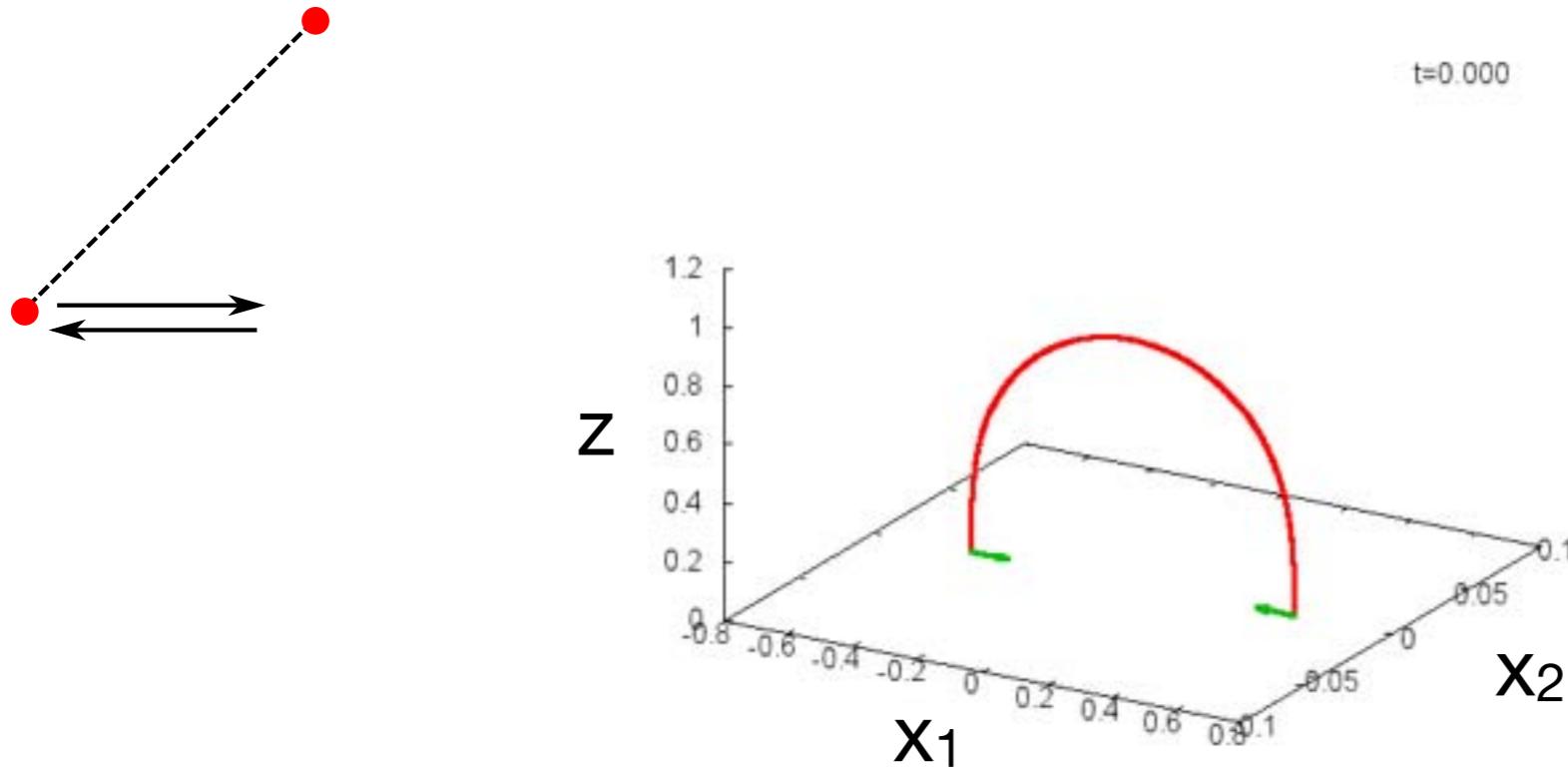


Z_2 -symmetric quench



- More discretized formation times because of wave collisions
- First cusps by such collisions (red \bullet). The cusps are **pair-created and annihilated**.
- Traveling cusps can be formed first (green \blacktriangle)

Transverse linear quench



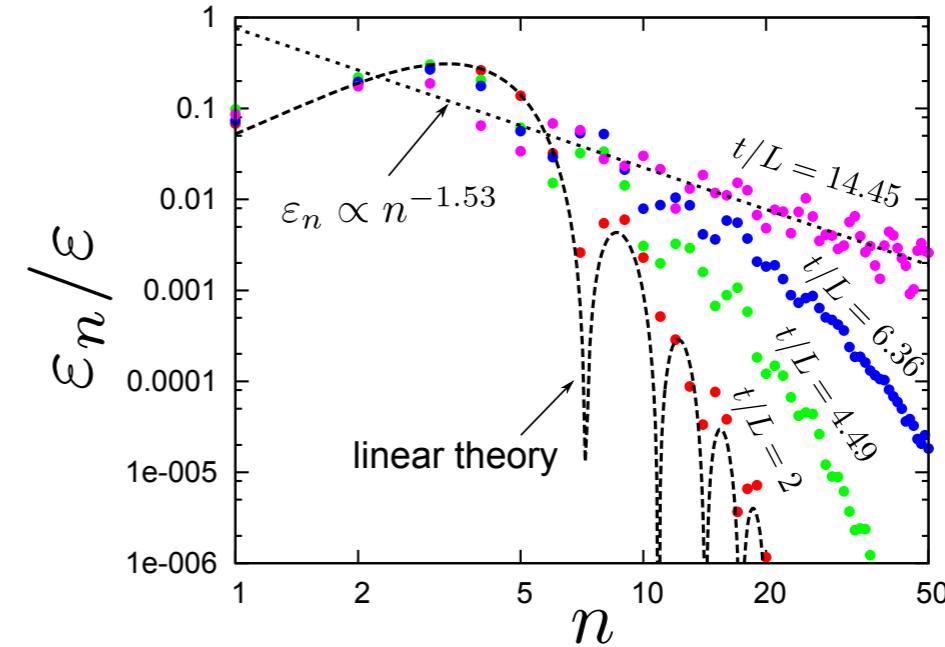
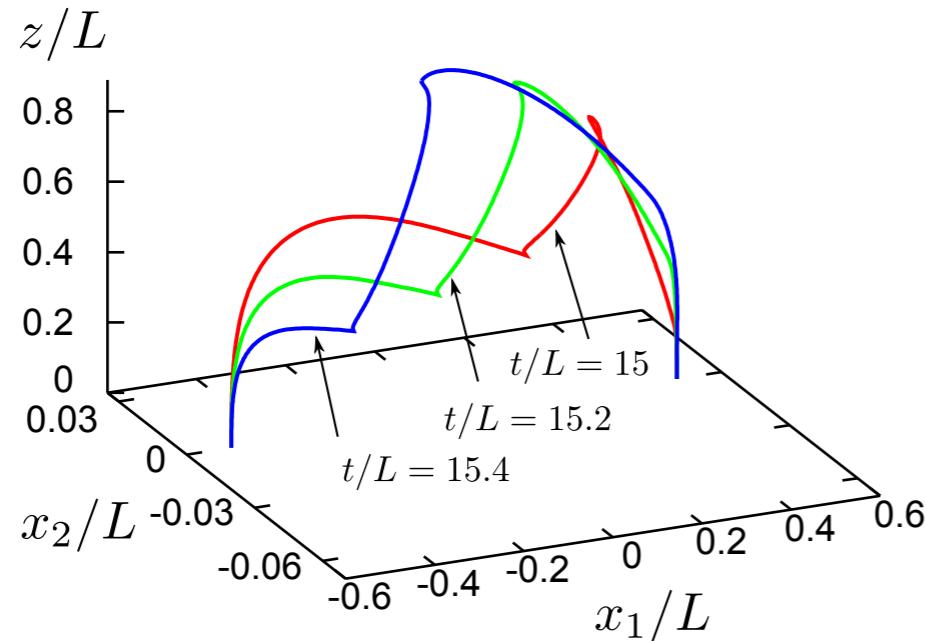
$t=0.000$

$\epsilon=0.03, \Delta t/L=2$

***Green arrows: forces

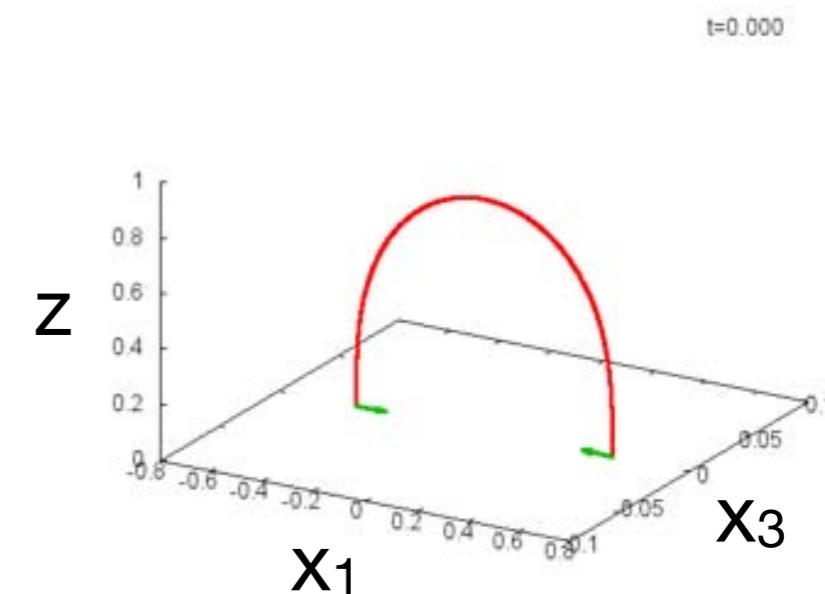
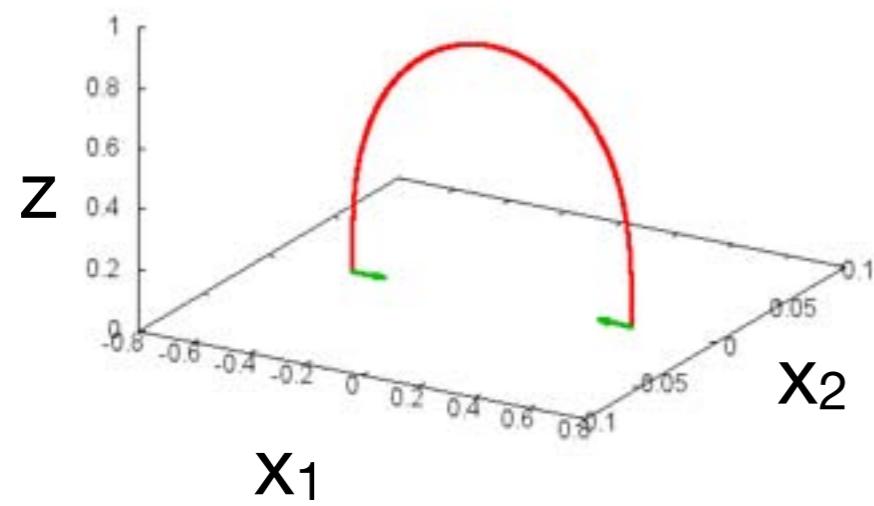
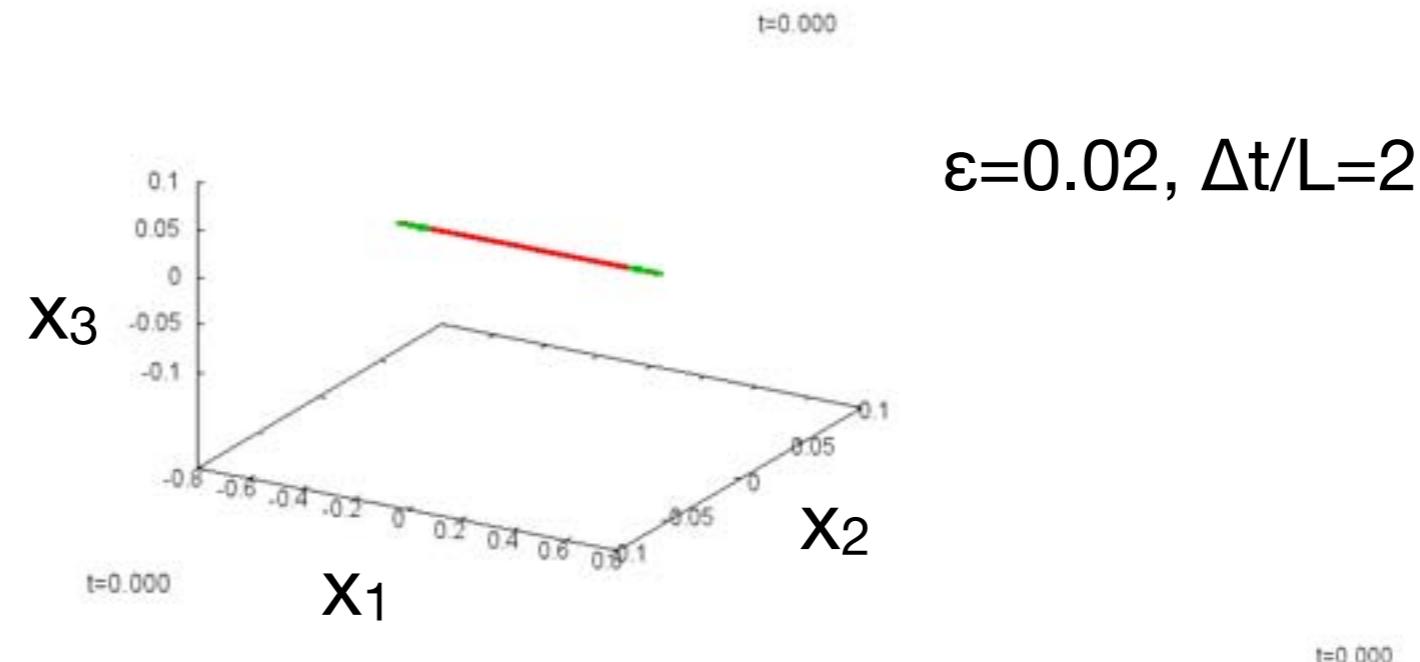
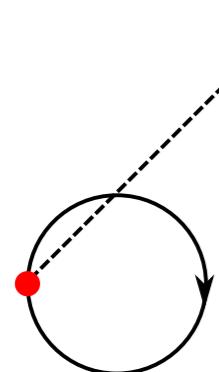
String oscillates in 1+3 dim (t, z, x_1, x_2)

Transverse linear quench



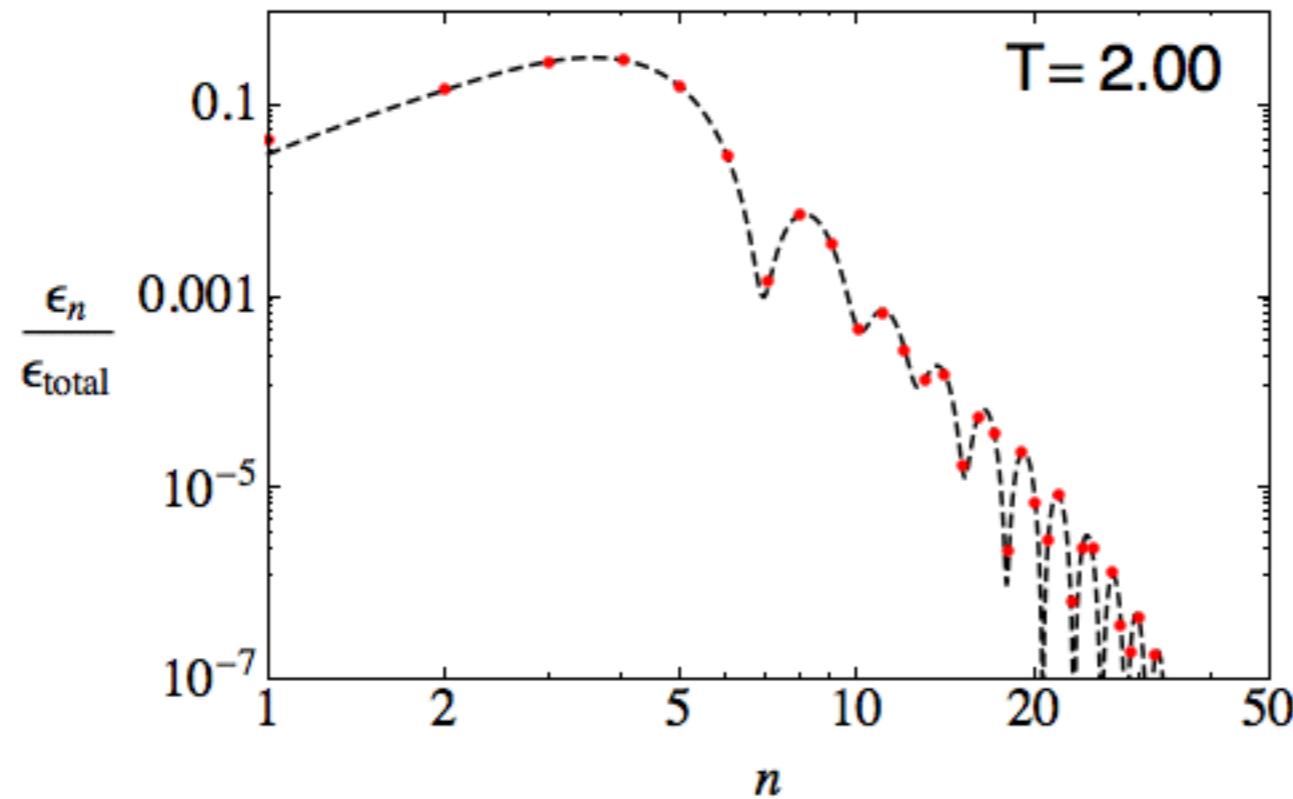
- Cusps are formed at $T \sim 14.45$
- The energy spectrum shows a direct cascade

Transverse circular quench



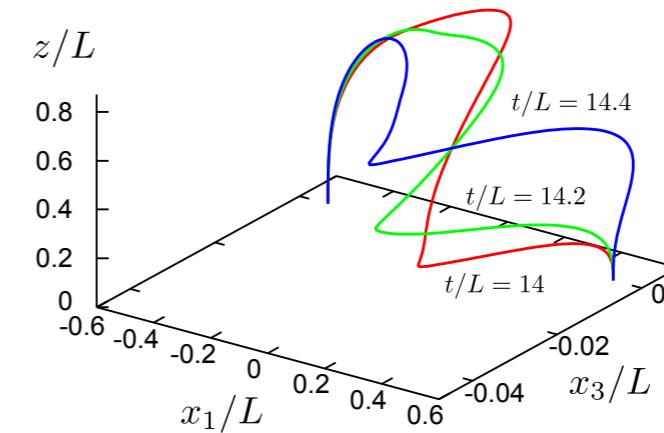
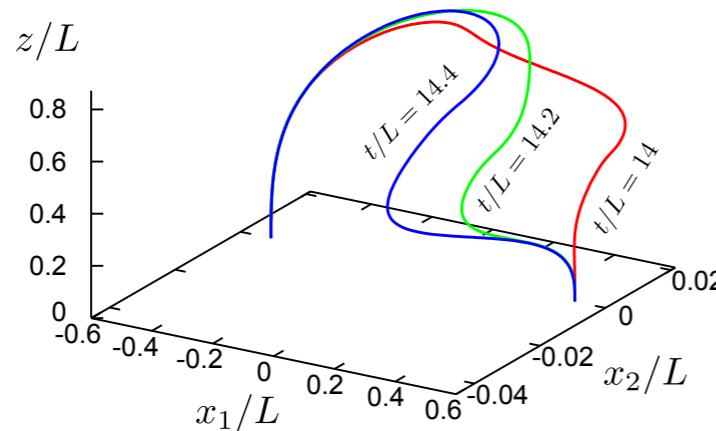
String oscillates in all 1+4 dim (t, z, x_1, x_2, x_3)

Energy spectrum (Log-log plot)

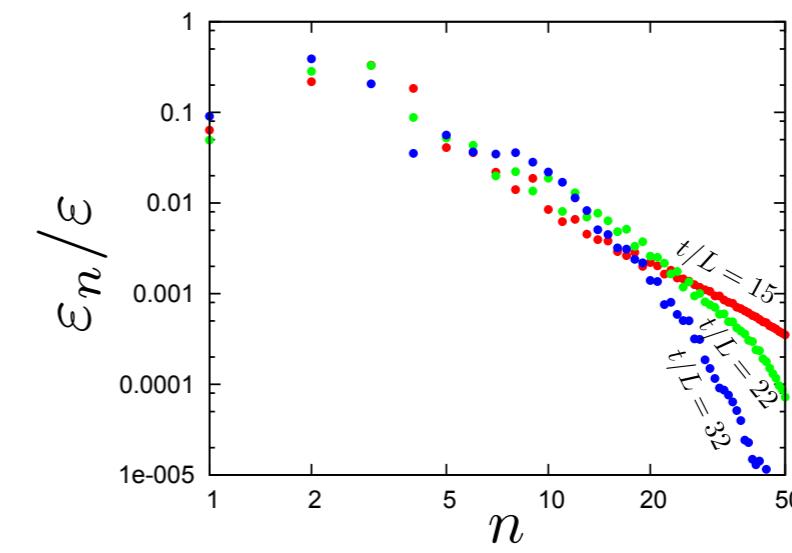
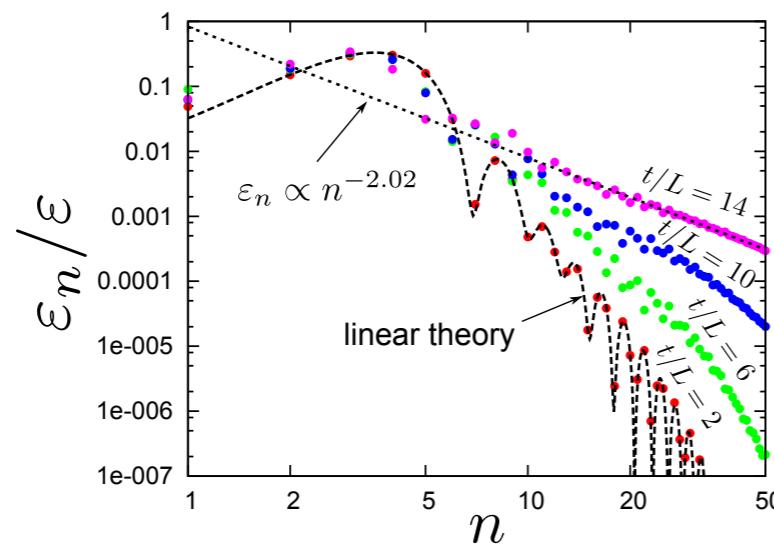


No cusp: no sustainable power law

Transverse circular quench



Cuspy, but not real cusps



Direct cascade \rightarrow inverse cascade

Summary

We computed nonlinear dynamics of the quark-antiquark fundamental string in AdS

- Cusps and turbulent behavior in $\leq 1+3$ dim
- No cusp and an inverse cascade in 1+4 dim

Future works

- Large amplitude/finite temperature
- Non-conformal backgrounds
- and more



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