## Unitarity techniques in 2D

based on 1304.1798 with B. Hoare, V. Forini, 1405.**** and work in progress

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May $28^{\text {th }}, 2014$

## Outline

(1) Motivation

(2) Method

(3) Features
(4) Applications
(5) Conclusion and outlook

## Motivation

- Combine a powerful technique with the special properties of $1+1$ dimensions.
- Improve perturbative computations in integrable non-linear sigma models.
- Understand the connection between cut constructibility and integrability.
- Perform non-trivial checks of quantum integrability for classically integrable string backgrounds.
- Compute overall scalar functions for symmetry-determined S-matrices.
- Moving towards the perturbative computation of off-shell quantities.


## The method

Standard unitarity in 4d [Bern, Dixon, Dunbar, Kosower, 1994]


Glue together the two amplitudes and uplift the integral with

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For $\mathrm{L}=1$


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\Rightarrow
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For $L=1$



## s-channel


$T_{M N}^{R S}\left(p, p^{\prime}\right) T_{R S}^{P Q}\left(p, p^{\prime}\right)$
t-channel
$\frac{1}{2} T_{M R}^{S P}(p, p) T_{S N}^{R Q}\left(p, p^{\prime}\right)$
$+\frac{1}{2} T_{M R}^{P S}\left(p, p^{\prime}\right) T_{S N}^{Q R}\left(p^{\prime}, p^{\prime}\right)$


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The result

$$
T^{(1)}=\frac{\theta}{2 \pi}(T(4) T-T ® T)+\frac{i}{2} T ® T+\frac{1}{16 \pi}\left(\frac{1}{m^{2}} \widetilde{T} \underset{\leftarrow}{\oplus} T+\frac{1}{m^{\prime 2}} T \underset{\rightarrow}{\oplus} \widetilde{T}\right)
$$

## Yang-Baxter equation



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Expanding the S-matrix perturbatively

$$
S=\square+(1 L+\cdots
$$

## Yang-Baxter equation



Expanding the S-matrix perturbatively

Tree-level YB


## Yang-Baxter equation



Expanding the S-matrix perturbatively

Tree-level YB

$$
a_{\alpha}+a_{\alpha}+a_{\alpha}=a_{\alpha}+\mathscr{a}_{\alpha}+\mathscr{a}_{\alpha}
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One loop YB


## Yang-Baxter equation



Expanding the S-matrix perturbatively

Tree-level YB


One loop YB - Homogeneous part


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- It is a trivial solution of the homogeneous part of YB.
- For the full result to satisfy YB equation it should be a solution to the homogeneous part of YB.
- It is again $\propto \mathbb{1}$ in all the cases but one: $\operatorname{AdS}_{3} \times S^{3} \times S^{3} \times S^{1}$. Why?


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"Unwanted" contribution

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## Applications

The technique has been applied to several two-dimensional integrable models:

- Bosonic relativistic models (generalized Sine-Gordon); [Hollowood, Miramontes, Park, 1994; Bakas, Park, Shin, 1995]
- Fermionic relativistic models (Pohlmeyer reduced theories for GS string in $A d S_{5} \times S_{5}$ and truncations) ; [Grigoriev, Tseytlin, 2007; Mikhailov, Schafer-Nameki, 2007]
- Non-relativistic models (worldsheet scattering for non-linear sigma model in $A d S_{5} \times S^{5}$ and $A d S_{3} \times S^{3} \times M^{4}$ ).[Metsaev, Tseytlin, 1998; Pesando, 1998; Rahmfeld, Rajaraman, 1998]


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Non-linear sigma model


## Worldsheet scattering in $A d S_{5} \times S^{5}$

Lagrangian of the non-linear sigma model [Metsaev, Tseytin, 1998]

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\mathcal{L}=\frac{\sqrt{\lambda}}{4 \pi} \int d^{2} \sigma \sqrt{-h} h^{a b} G_{M N}(X) \partial_{a} X^{M} \partial_{b} X^{N}+\text { fermions }
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- Uniform light-cone gauge fixing [Frolov, Plefka, Zamaklar 2006]

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H_{w s}=\int d \sigma \mathcal{H}_{w s}=-\int d \sigma P_{-} \equiv E-J
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- Decompactification limit to define asymptotic states

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## Conclusion and outlook

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- The one-loop S-matrix can be expressed as a combination of products of TL S-matrices.
- It surely reproduces the logarithmic dependence (checked also at two loops).
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Future directions

- We found hints of a relation with YB equation which deserves further analysis.
- Rational terms beyond one loop.
- Study of off-shell objects via unitarity.

