Unitarity techniques in 2D

based on 1304.1798 with B. Hoare, V. Forini, 1405.**** and work in progress

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May 28th, 2014

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Outline













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Motivation

- Combine a powerful technique with the special properties of 1+1 dimensions.
- Improve perturbative computations in integrable non-linear sigma models.
- Understand the connection between cut constructibility and integrability.
- Perform non-trivial checks of quantum integrability for classically integrable string backgrounds.
- Compute overall scalar functions for symmetry-determined S-matrices.
- Moving towards the perturbative computation of off-shell quantities.

The method



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The method



Generalized unitarity in 4d [Bern, Dixon, Kosower, 1998; Britto, Cachazo, Feng, 2004]

$$\mathcal{A}^{L} = \sum_{i} c_{i} \underbrace{\mathcal{I}_{i}^{(L)}} \longrightarrow \text{Known basis of L-loop scalar integrals}$$

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The method



Generalized unitarity in 4d [Bern, Dixon, Kosower, 1998; Britto, Cachazo, Feng, 2004]



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The method





Glue together the two amplitudes and uplift the integral with $i\pi\delta^+(p^2-m^2) \rightarrow \frac{1}{n^2-m^2-i\epsilon}$

Generalized unitarity in 2d [Engelund, McKeown, Roiban, 2013]



	Method
s-channel	t-channel
$M \xrightarrow{p} R \xrightarrow{p'} Q$ $A^{(0)} \xrightarrow{l_2} S \xrightarrow{p} P$	$M P p$ $A^{(0)} p$ R^{1} $A^{(0)} p'$

 $T^{RS}_{MN}(p,p')T^{PQ}_{RS}(p,p')$

 $p'_{N} p'_{Q}$ $\frac{1}{2}T^{SP}_{MR}(p,p)T^{RQ}_{SN}(p,p')$ $+\frac{1}{2}T^{PS}_{MR}(p,p')T^{QR}_{SN}(p',p')$



u-channel

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Method		
s-channel	t-channel	u-channel
$M \xrightarrow{p} R \xrightarrow{p'} Q$ $M \xrightarrow{p'} A^{(0)} \xrightarrow{1_1} A^{(0)} \xrightarrow{p'} Q$ $T_{MN}^{RS}(p,p') T_{RS}^{PQ}(p,p')$	$M \qquad P \\ p \qquad p' \qquad p' \\ N \qquad Q \\ \frac{1}{2}T_{MR}^{SP}(p,p)T_{SN}^{RQ}(p,p') \\ 1 = PS \left(-D \right) = O^{SP}(-D -D)$	$M \qquad Q \\ p' \qquad P' \qquad$
	$\pm \overline{2} I_{MR}(P,P) I_{SN}(P,P)$	

The result

$$T^{(1)} = \frac{\theta}{2\pi} (T \textcircled{w} T - T \textcircled{s} T) + \frac{i}{2} T \textcircled{s} T + \frac{1}{16\pi} (\frac{1}{m^2} \widetilde{T} \textcircled{t} T + \frac{1}{m'^2} T \textcircled{t} \widetilde{T})$$

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Expanding the S-matrix perturbatively

$$\overbrace{\texttt{S}} = \underline{\qquad} + \overbrace{\texttt{T}} + \overbrace{\texttt{H}} + \cdots$$

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Expanding the S-matrix perturbatively

$$\overline{S} =$$
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Tree-level YB



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Yang-Baxter equation



Expanding the S-matrix perturbatively

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Expanding the S-matrix perturbatively

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Tree-level YB



One loop YB - Homogeneous part



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The result

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- Satisfies the inhomogeneous YB equation for any TL S-matrix (proven).
- Gives the one-loop rational terms up to an overall phase (observed).

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- In all the cases we considered this coefficient is $\propto 1$, as required by integrability.
- It is a trivial solution of the homogeneous part of YB.

A B F A B F

Image: A matrix

The result

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• Satisfies the inhomogeneous YB equation for any TL S-matrix (proven).
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- For the full result to satisfy YB equation it should be a solution to the homogeneous part of YB.
- It is again ∝ 1 in all the cases but one: AdS₃ × S³ × S³ × S¹. Why?

"Unwanted" contribution

$$rac{1}{4\pi}\left(rac{p^2}{m}+rac{p'^2}{m'}
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- *p* and *p'* are the spatial component of the ingoing (and outgoing) two-momenta.
- Similar to an external leg correction for non-relativistic theory.
- The only theory with this contribution has three-point interactions

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with
$$= \frac{i(1+h^{-1}\Sigma_1(p))}{\mathbf{p}^2 - m^2 - h^{-1}\Sigma_0(p)} + \dots$$

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$$\mathbf{p}_{\mathcal{F}^{(0)}} = \Sigma_0(p) + \Sigma_1(p)(\mathbf{p}^2 - m^2) + \mathcal{O}((\mathbf{p}^2 - m^2)^2)$$

Applications

The technique has been applied to several two-dimensional integrable models:

- Bosonic relativistic models (generalized Sine-Gordon); [Hollowood, Miramontes, Park, 1994; Bakas, Park, Shin, 1995]
- Fermionic relativistic models (Pohlmeyer reduced theories for GS string in $AdS_5 \times S_5$ and truncations); [Grigoriev, Tseytlin, 2007; Mikhailov, Schafer-Nameki, 2007]
- Non-relativistic models (worldsheet scattering for non-linear sigma model in $AdS_5 \times S^5$ and $AdS_3 \times S^3 \times M^4$).[Metsaev, Tseytlin, 1998; Pesando, 1998; Rahmfeld, Rajaraman, 1998]

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Non-linear sigma model



Worldsheet scattering in $AdS_5 \times S^5$

Lagrangian of the non-linear sigma model [Metsaev, Tseytlin, 1998]

$$\mathcal{L} = \frac{\sqrt{\lambda}}{4\pi} \int d^2 \sigma \sqrt{-h} \, h^{ab} G_{MN}(X) \, \partial_a X^M \partial_b X^N + \text{fermions}$$

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Uniform light-cone gauge fixing [Frolov, Plefka, Zamaklar 2006]

$$H_{ws} = \int d\sigma \mathcal{H}_{ws} = -\int d\sigma P_{-} \equiv E - J$$

• Decompactification limit to define asymptotic states

$$-rac{P_+}{2} < \sigma < rac{P_+}{2} \qquad,\qquad P_+ o \infty$$

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Conclusion and outlook

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- We developed a technique to perform perturbative computations in 1+1 dimensions at one loop.
- The one-loop S-matrix can be expressed as a combination of products of TL S-matrices.
- It surely reproduces the logarithmic dependence (checked also at two loops).
- In all the integrable theories we studied it reproduces also the rational terms.

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Future directions

- We found hints of a relation with YB equation which deserves further analysis.
- Rational terms beyond one loop.
- Study of off-shell objects via unitarity.

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