Newton-Cartan trace anomalies

Roberto Auzzi

Università Cattolica del Sacro Cuore, Brescia

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Based on arXiv:1511.08150 and work in progress with G. Nardelli, S. Baiguera and F. Filippini

Introduction: c and a theorem

Zamolodchikov's c-theorem d = 2:

There exists of a c-function which decreases along the RG flow, and it agrees with the central charge at the conformal fixed point $\langle T^{\mu}_{\mu} \rangle = \frac{c}{24\pi} R$

a-theorem in d = 4:

$$\langle T^{\mu}_{\mu} \rangle = c (\text{Weyl})^2 - a E_4 + b R^2 + \tilde{b} \nabla^2 R + d \epsilon^{\mu \nu \rho \sigma} R^{\alpha \beta}_{\mu \nu} R_{\alpha \beta \rho \sigma} ,$$

$$E_4 = R_{\mu\nu\rho\sigma}^2 - 2R_{\mu\nu}^2 + \frac{1}{3}R^2$$

a is a candidate decreasing function !! (Cardy)

a-theorem d = 4

$$d=4\,,\qquad T^\mu_\mu=cW^2-aE_4\,,\qquad a_{UV}>a_{IR}\,,$$

Free theory:
$$a = \frac{1}{90(8\pi)^2}(n_S + \frac{11}{2}n_{F,W} + 62n_V)$$

Tools to study irreversibility properties of RG flows:

- Local Renormalization Group equations (perturbative proof of *a*-theorem by Osborn, 1991)
- Holography
- SUSY, a-maximization
- Dispersion Relations
 Dilaton scattering amplitudes (Komargodski, Schwimmer 2011)

Trace anomalies

Anomalies must satisfy Wess-Zumino consistency conditions:

$$[\Delta^W_\sigma,\Delta^W_{\sigma'}]W=0$$

Eliminates R^2 term at conformal fixed point

Moreover, some of them can be eliminated by local counterterms and are not genuine anomalies, like $\nabla^2 R$

- **type B** anomaly: vanish under Weyl transformation (like σW^2)
- type A anomaly: does not trivially vanish under Weyl, but it is still WZ consistent (like σE₄)

Is there a non-relativistic version?

In order to inspect these issues, we should couple the non-relativistic theory to a non-relativistic version of gravity

Gravity here is just a source for

energy-momentum tensor:

$$\langle T_{\mu\nu} \rangle = \frac{1}{\sqrt{-g}} \frac{\delta W}{\delta g^{\mu\nu}}$$

The kind of non-relativistic gravity depends on the symmetry of the non-relativistic theory: Lifshitz or Schrödinger

With or without boost

$$t \to \lambda^z t$$
, $x \to \lambda x$.

- Without boost: studied in detail for various d, z by Arav, Chapman, Oz: 1410.1831 All anomalies are type B (vanish under Weyl transformation)
- With boost:

Jensen1412.7750, found a type A anomaly Basic tool to classify boost invariant terms in the anomaly: Newton Cartan gravity

Newton-Cartan gravity

- "Spatial metric" $h^{\mu\nu}$, $h_{\mu\nu}$
- 1-form n_{μ} (local time direction),
- background gauge field for particle number symmetry A_μ
- vector field v^{μ}
- Conditions: $n_{\mu}h^{\mu\alpha} = 0$, $n_{\mu}v^{\mu} = 1$ $h^{\mu\alpha}h_{\alpha\nu} = \delta^{\mu}_{\nu} v^{\mu}n_{\nu} = P^{\mu}_{\nu}$

Nearby flat space: introduce δu^{μ} with $\delta u^{\mu}n_{\mu} = 0$ and a transverse metric perturbation $\delta \tilde{h}^{\alpha\beta}n_{\beta} = 0$.

$$\delta W = \int d^3x \sqrt{-g} \left(\frac{1}{2} T_{ij} \delta \tilde{h}_{ij} + j^{\mu} \delta A_{\mu} - \epsilon^{\mu} \delta n_{\mu} - p_i \delta u_i \right) \,,$$

Sources for stress tensor, number current, energy current, momentum density

Newton-Cartan from null reduction

Invariance under local non-relativistic boost transforms v^{μ} , $h_{\mu\nu}$, A_{μ} non-trivially

Null reduction is an instrument to deal efficiently with these symmetries:

$$G_{MN}=\left(egin{array}{cc} 0&n_{\mu}\ n_{
u}&n_{\mu}A_{
u}+n_{
u}A_{\mu}+h_{\mu
u}\end{array}
ight)$$

The anomaly can be written in terms of the extra dimensional curvatures

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Causality

The form $n = n_\mu dx^\mu$ gives the local time direction

Frobenius condition

 $dn \wedge n = 0$,

it is equivalent integrability of time slices

If it not satisfied, no absolute notion of future and past, **causality is lost**

No Frobenius, the anomaly has infinite number of terms

d = 2, *z* = 2 case, [Jensen, 1412.7750]:

$$\mathcal{A} = T_i^i - 2\epsilon^0 - = c (Weyl)^2 - aE_4 + bR^2 + \tilde{b}\nabla^2 R + \dots$$

the same expression as for the relativistic 3+1 anomaly, one should just replace null reduction curvatures+

 $+ \ldots$ which stands for an **infinite number** of terms, such as:

$$W_{ABCP}W^{ABCQ}W^{P}_{MQN}n^{M}n^{N}$$
,
(for n^{A} , we buy 2 more derivatives)

Extra terms have not been yet classified [work in progress]

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With Frobenius, type B

$$dn \wedge n = 0$$
,

The anomaly has just a finite number of terms

$$\mathcal{A}=2\epsilon^0-\mathcal{T}_i^i=c(\mathrm{Weyl})^2+\mathrm{local\ counterterms}$$
 Only one type B anomaly

R.A., S. Baiguera, G. Nardelli: arXiv:1511.08150 I. Arav, S. Chapman, Y. Oz arXiv:1601.06795

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A free scalar in curved NC space

$$\int d^{d+1}x \sqrt{g} \left\{ imv^{\mu} \left(\phi^{\dagger} D_{\mu} \phi - D_{\mu} \phi^{\dagger} \phi \right) - h^{\mu\nu} D_{\mu} \phi^{\dagger} D_{\nu} \phi - \xi R \phi^{\dagger} \phi \right\} \,,$$

$$D_{\mu}\phi = \partial_{\mu}\phi - imA_{\mu}\phi$$
.

Conformal coupling:

$$\xi = \frac{1}{6}$$

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Heat kernel method

Heat kernel method to compute anomaly:

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In flat spacetime:

$$egin{aligned} &\mathcal{K}_{ riangle}(s) = \langle x,t|e^{s riangle}|y,t'
angle = \ &= rac{1}{2\pi}rac{ms}{m^2s^2+rac{(t-t')^2}{4}}rac{1}{(4\pi s)^{d/2}}\exp\left(-rac{(x-y)^2}{4s}
ight) \ & ilde{\mathcal{K}}_{ riangle}(s) = \langle x,t|e^{s riangle}|x,t
angle = rac{2}{m(4\pi s)^{1+d/2}}\,. \end{aligned}$$

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The same spectral dimension as relativistic scalar in d + 2 dimensions !!

Heat kernel in curved spacetime

Seeley-DeWitt expansion:

$$ilde{K}_{M}(s) = rac{1}{s^{d/2+1}} \left(a_{0} + a_{2}s + a_{4}s^{2} + \ldots
ight) \, .$$

 a_4 is proportional to the Weyl anomaly

Result:

$$T_i^i - 2\epsilon^0 = \frac{1}{8m\pi^2} \left(-\frac{1}{360} E_4 + \frac{3}{360} W^2 + \frac{1}{2} \left(\xi - \frac{1}{6} \right)^2 R^2 + \frac{1 - 5\xi}{30} D_A D^A R \right) + \dots,$$

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A non-relativistic a-theorem ?

A free scalar contributes to the non-relativistic *a* anomaly:

 $\frac{1}{8m\pi^2}\frac{1}{360}$

If the UV and IR fixed points contain just scalars, the (conjectured) non relativistic a-theorem would tell that:

$$\sum_{k}^{UV} rac{1}{m_k} \geq \sum_{k}^{IR} rac{1}{m_k}$$

Fermions: Work in progress

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Future directions

- Heat kernel calculation for fermion, vector and Chern-Simon field
- Try to check the conjecture in some examples
- Can we apply Osborn's local RG formalism for a perturbative proof?
- Dispersion relations a la Komargodski and Schwimmer?
- Is there any interesting SUSY story, as in a-maximization ?

• Full classification of anomalies with more derivatives

Thank you