

Effects of coupling to breakup channels in reactions induced by weakly bound and halo nuclei.

Juan Pablo Fernández García

Istituto Nazionale di Fisica Nucleare

LNS



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Introduction

Weakly bound nuclei (${}^6,7\text{Li}$): Low break-up threshold.

$$S_\alpha = 1.47(2.47) \text{ MeV}$$

Cluster structure (${}^6,7\text{Li} \rightarrow \alpha + d(t)$)



Coupling to breakup channels expected to be important.

Halo nuclei (${}^6\text{He}$): Low break-up threshold.

$$S_{2n} = 0.98 \text{ MeV}$$

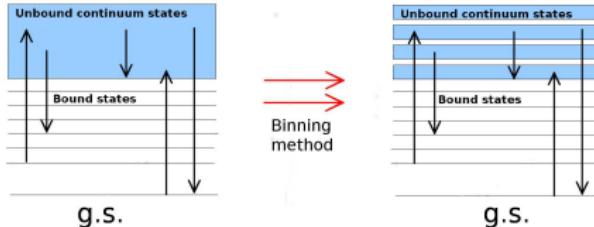
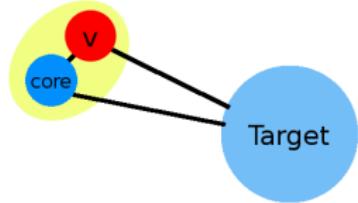
Halo structure (${}^6\text{He} \rightarrow \alpha + n + n$)



Coupling to breakup channels expected to be important.

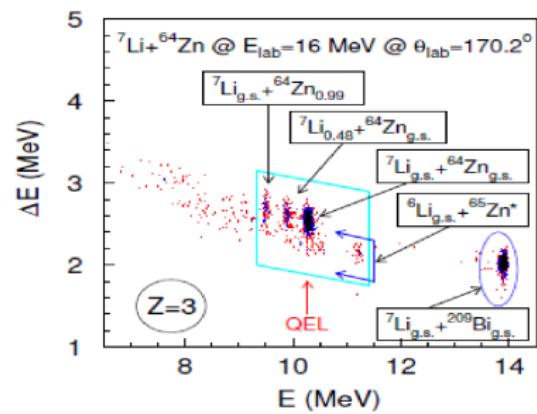
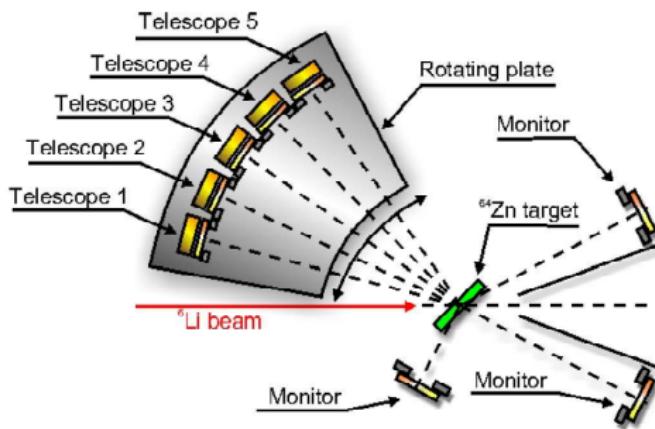
Continuum-Discretized Coupled-Channels (CDCC) formalism

- Uses Coupled-Channel (CC) method to solve the scattering problem.
- For weakly bound/halo nuclei it is important to introduce the continuum (unbound) states of the projectile.
- Unbound states of the projectile included by means of a discretization procedure (binning).
- Describes the reaction using an effective 3-body model.
- The breakup process is treated as an inelastic excitation of the projectile.



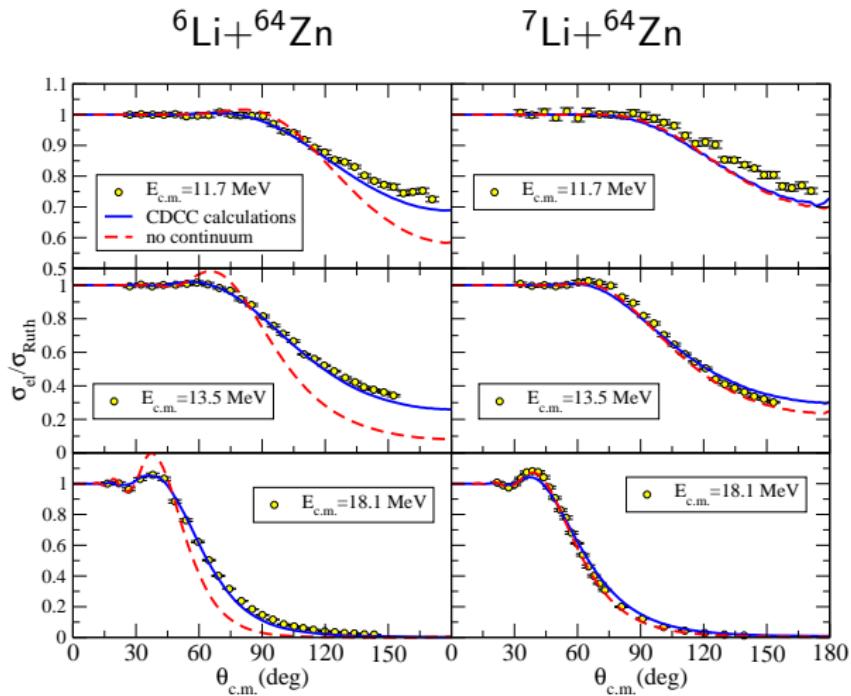
Experiment $^{6,7}\text{Li} + ^{64}\text{Zn}$ at LNS laboratory, Italy

- $^{6,7}\text{Li} + ^{64}\text{Zn}$ @ $E_{c.m.} = 11.7, 12.4, 13.5, 15.0, 16.3$ and 18.1 MeV.
- 5 Telescopes ΔE ($10 \mu\text{m}$) + E ($200 \mu\text{m}$).
- Stable beam. Good data quality.



M. Zadro et al. Phys. Rev. C 80, 064610 (2009), M. Zadro et al. Phys. Rev. C 87, 054606 (2013) and to be published

CDCC calculations - Coupling to the continuum effect

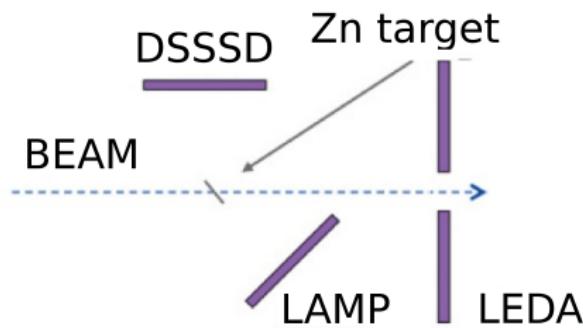


Coupling to breakup channels are more important for ^6Li .
 $S_\alpha(^6\text{Li})=1.47 \text{ MeV}$ $S_\alpha(^7\text{Li})=2.47 \text{ MeV}$

Experiment ${}^6\text{He} + {}^{64}\text{Zn}$ at CRC laboratory, Belgium.

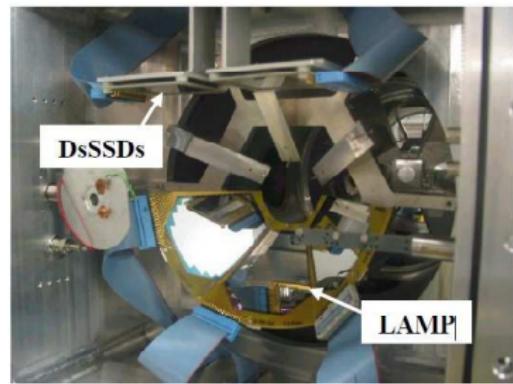
New experimental data above the Coulomb barrier - $E_{c.m.} = 13.5$ and 16.5 MeV.

Angular range between $5 - 120$ degrees.



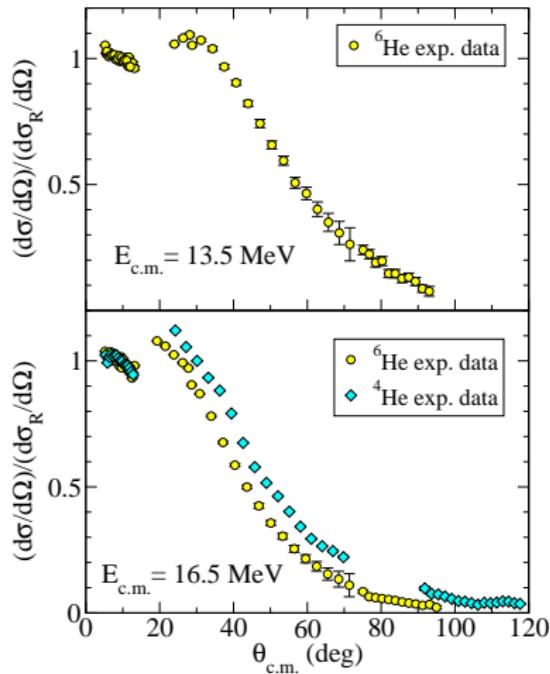
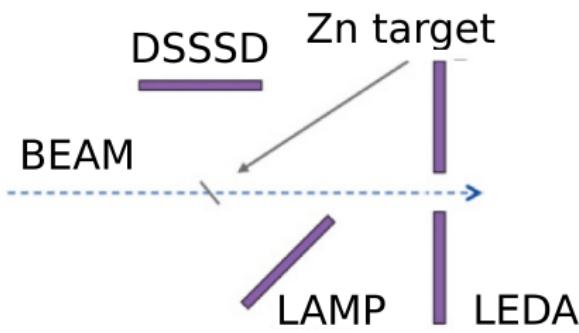
DSSSDs	LAMP	LEDA
$67^\circ - 120^\circ$	$22^\circ - 65^\circ$	$5^\circ - 12^\circ$

A ${}^{64}\text{Zn}$ target of 0.5 mg/cm^2 was used and rotated 45° .



Experiment ${}^6\text{He} + {}^{64}\text{Zn}$ - $E_{c.m.} = 13.5$ and 16.5 MeV.

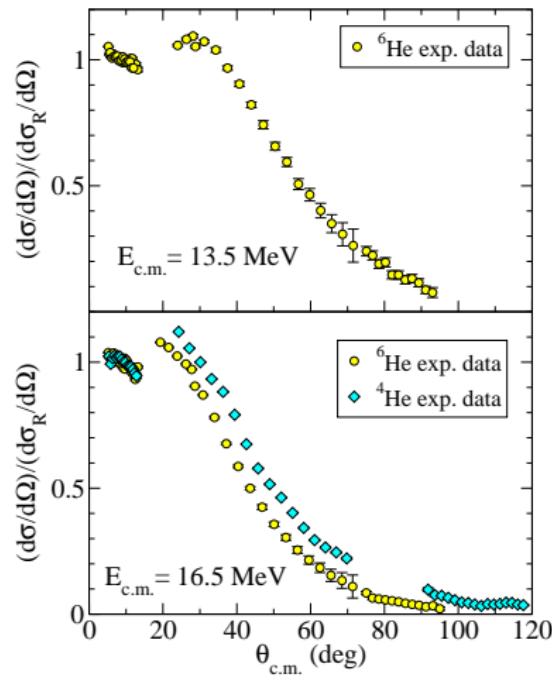
New experimental data above the Coulomb barrier
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Experiment ${}^6\text{He} + {}^{64}\text{Zn}$ - $E_{c.m.} = 13.5$ and 16.5 MeV.

Angular range between 5 -120 degrees.

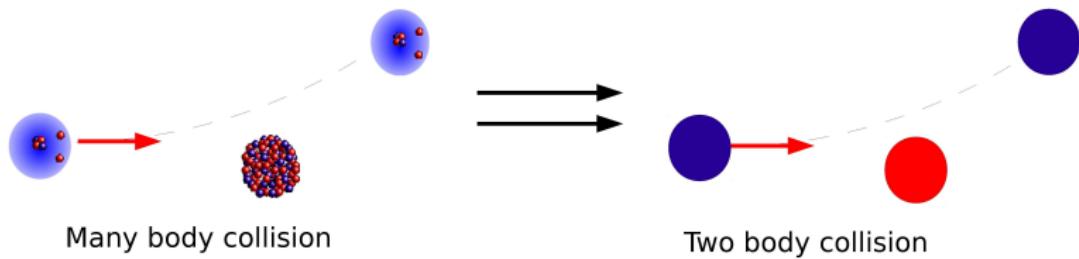
- Deviation from Rutherford.
- Diminution of Fresnel peak.
- Expected important nuclear and/or Coulomb couplings.



Optical model analysis

Describes the interaction projectile-target with an effective average potential $U(R) = V(R) + iW(R)$

The absorption of the elastic channel is represented by an imaginary part in the nuclear potential



- $U(R)$ phenomenological - Woods-Saxon form.

$$U(R) = V_0 f + iW_0 f$$

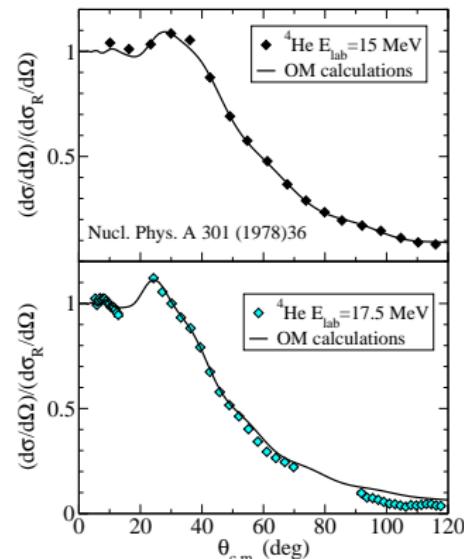
$$f(R, R_x, a_x) = \frac{1}{e^{\frac{R-R_x}{a_x}} + 1}$$

Optical model analysis - ${}^4\text{He} + {}^{64}\text{Zn}$

$$U(R) = U_\alpha(R) + W_L(R) \left\{ \begin{array}{l} U_\alpha(R) = V(R) + iW(R) \text{ -- Core potential} \\ W_L(R) = df(R)/dR \text{ -- Long range effect} \end{array} \right.$$

$E_{lab} = 15 \text{ MeV}^a$	$V_0(\text{MeV})$	$r_0(\text{fm})$	$a_0(\text{fm})$
	131.1	1.04	0.594
$E_{lab} = 17.5 \text{ MeV}$	$V_0(\text{MeV})$	$r_0(\text{fm})$	$a_0(\text{fm})$
	22.38	1.2	0.43
$W_0(\text{MeV})$	$r_i(\text{fm})$	$a_i(\text{fm})$	
	11.53	1.27	0.358
	23.86	1.05	0.43

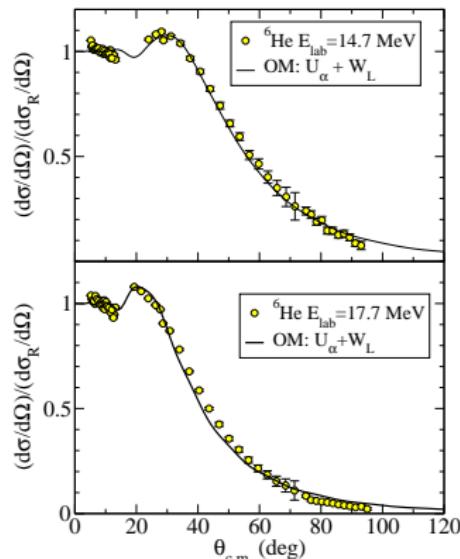
^a T. B. Robinson and V. R. W. Edwards Nucl. Phys. A 301 (1978) 36



Optical model analysis - ${}^6\text{He} + {}^{64}\text{Zn}$

$$U(R) = U_\alpha(R) + W_L(R) \left\{ \begin{array}{l} U_\alpha(R) = V(R) + iW(R) \text{ -- Core potential} \\ W_L(R) = df(R)/dR \text{ -- Long range effect} \end{array} \right.$$

$E_{lab} = 15 \text{ MeV}$	$W_L(\text{MeV})$	$r_L(\text{fm})$	$a_L(\text{fm})$
	2.62	1.2	0.71
$E_{lab} = 18 \text{ MeV}$	$W_L(\text{MeV})$	$r_L(\text{fm})$	$a_L(\text{fm})$
	2.28	1.2	0.94



Optical model analysis - ${}^6\text{He} + {}^{64}\text{Zn}$

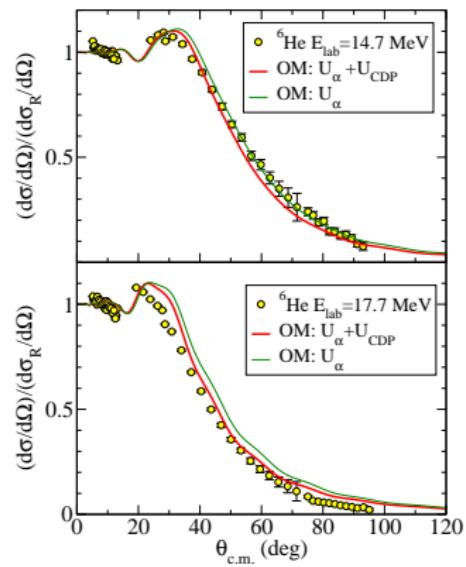
$$U(R) = U_\alpha(R) + U_{CDP}(R)$$

$$U_{CDP} = -\frac{4\pi}{9} \frac{Z_t^2 e^2}{\hbar \cdot v} \frac{1}{(r - a_0)^2 r} \int_{\varepsilon_b}^{\infty} d\varepsilon \frac{dB(E1, \varepsilon)}{d\varepsilon} F(r, \varepsilon)$$

$$F(r, \varepsilon) = g\left(\frac{r}{a_0} - 1, \xi\right) + i f\left(\frac{r}{a_0} - 1, \xi\right)$$

M. V. Andrés, J. Gómez-Camacho and M. A. Nagarajan, Nucl. Phys. A

583, 817 (1995)



Optical model analysis - ${}^6\text{He} + {}^{64}\text{Zn}$

$$U(R) = U_\alpha(R) + W_L(R) + U_{CDP}(R)$$

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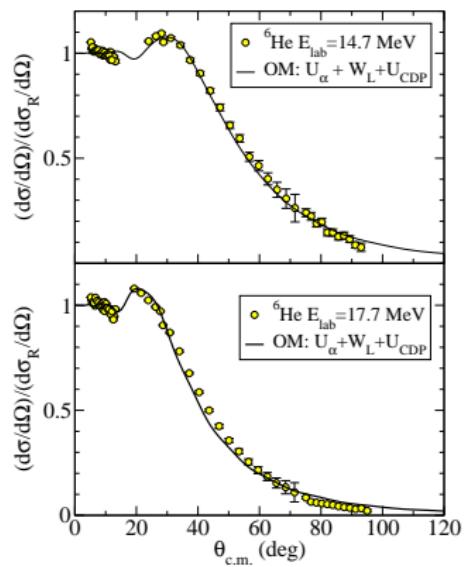
M. V. Andrés, J. Gómez-Camacho and M. A. Nagarajan, Nucl. Phys. A

583, 817 (1995)

$E_{lab} = 15 \text{ MeV}$ $a_L = 0.71 \text{ fm} \rightarrow 0.69 \text{ fm}$

$E_{lab} = 18 \text{ MeV}$ $a_L = 0.94 \text{ fm} \rightarrow 0.90 \text{ fm}$

Small effects of the dipole Coulomb
couplings.



Optical model analysis - ${}^6\text{He} + {}^{64}\text{Zn}$

$$U(R) = U_\alpha(R) + U_{CDP}(R)$$

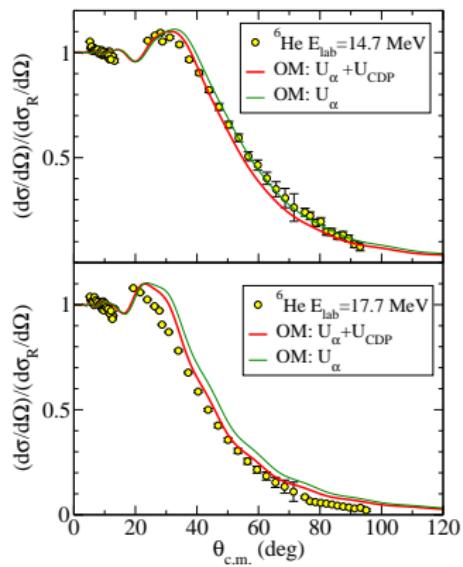
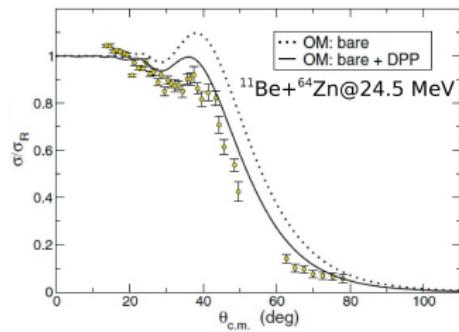
$$U_{CDP} = -\frac{4\pi}{9} \frac{Z_t^2 e^2}{\hbar \cdot v} \frac{1}{(r - a_0)^2 r} \int_{\varepsilon_b}^{\infty} d\varepsilon \frac{dB(E1, \varepsilon)}{d\varepsilon} F(r, \varepsilon)$$

$$F(r, \varepsilon) = g\left(\frac{r}{a_0} - 1, \xi\right) + i f\left(\frac{r}{a_0} - 1, \xi\right)$$

M. V. Andrés, J. Gómez-Camacho and M. A. Nagarajan, Nucl. Phys. A

583, 817 (1995)

Di Pietro et al. Phys. Rev. C 85, 054607 (2012)



Optical model analysis - ${}^6\text{He} + {}^{64}\text{Zn}$

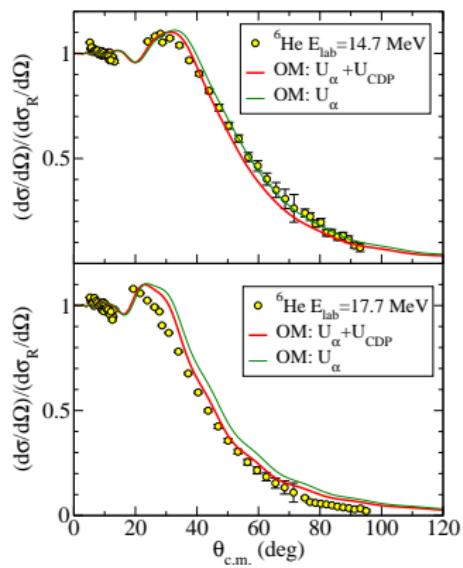
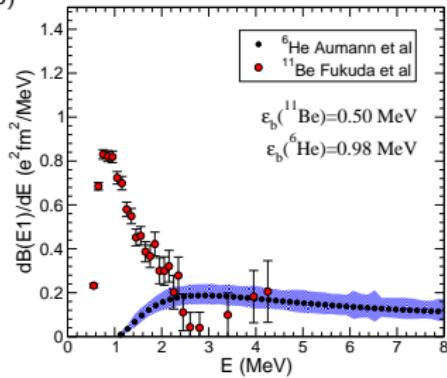
$$U(R) = U_\alpha(R) + U_{CDP}(R)$$

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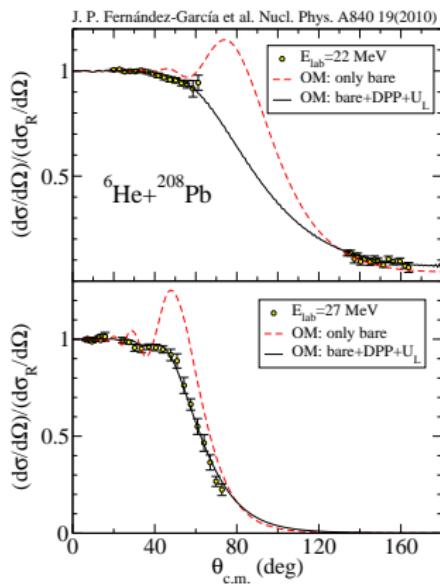
M. V. Andrés, J. Gómez-Camacho and M. A. Nagarajan, Nucl. Phys. A

583, 817 (1995)

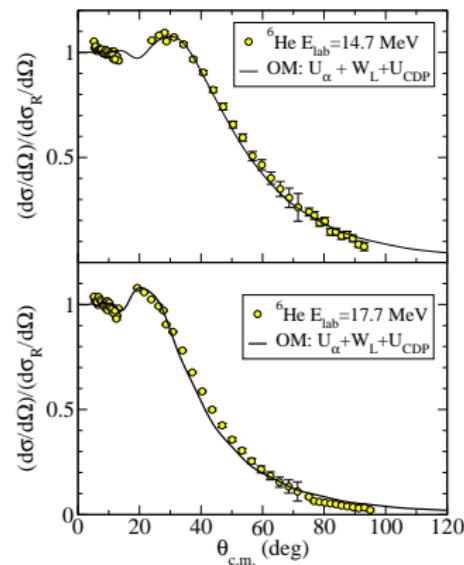


Optical model analysis - ${}^6\text{He} + {}^{64}\text{Zn}$

$$U(R) = U_\alpha(R) + W_L(R) + U_{CDP}(R)$$



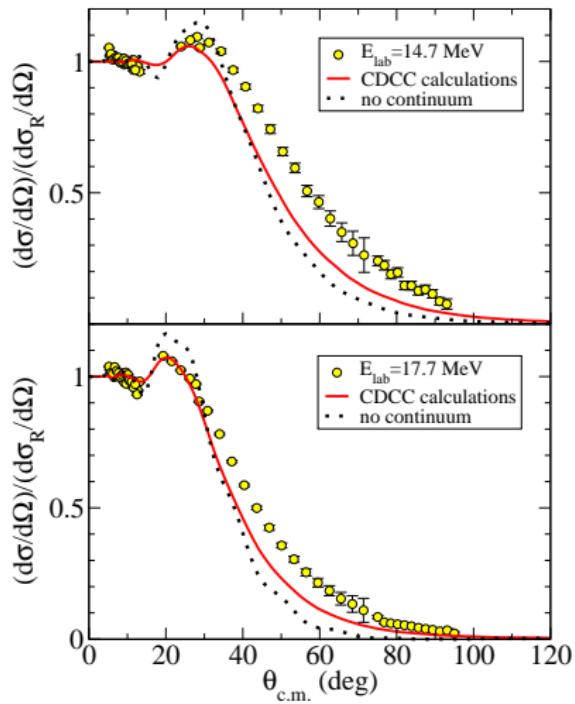
$a_L = 1.9 \text{ fm} \rightarrow 1.0 \text{ fm}$



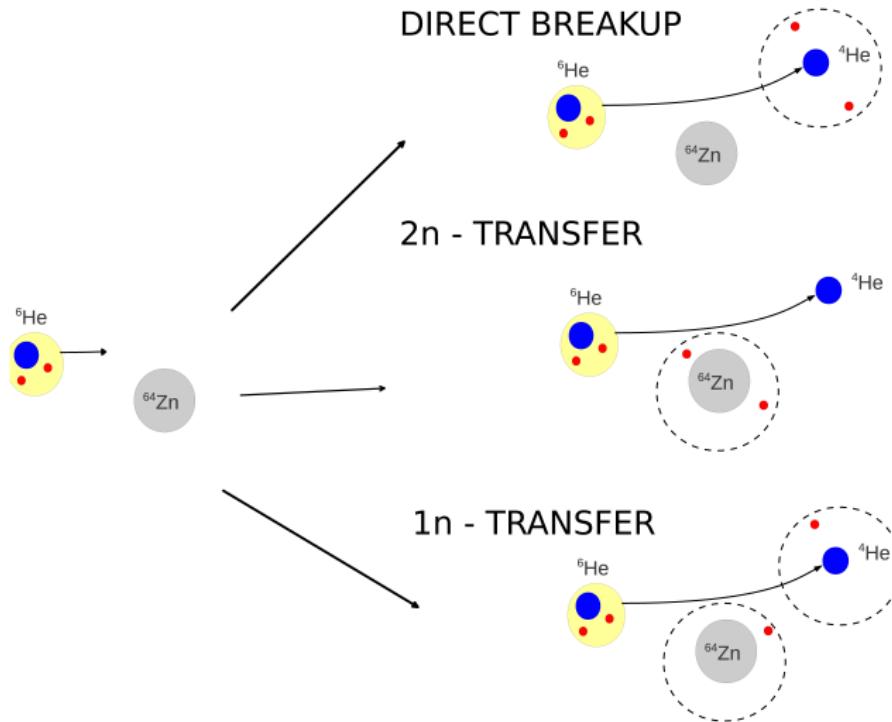
$a_L = 0.94 \text{ fm} \rightarrow 0.90 \text{ fm}$

CDCC calculations - Coupling to the continuum effect

- Effective 3-body model (${}^4\text{He} + 2\text{n} + {}^{64}\text{Zn}$)
- The breakup process is treated as an inelastic excitation of the projectile.
- CDCC calculations underestimate the experimental data.

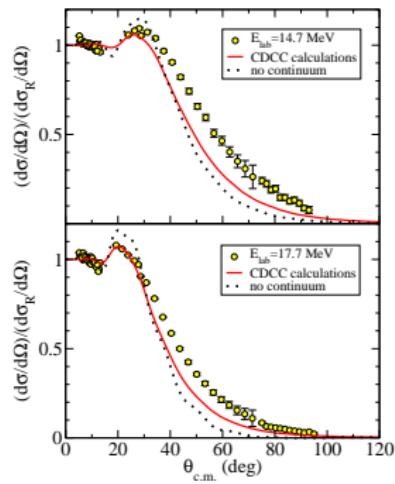


The couplings to breakup channels are important.

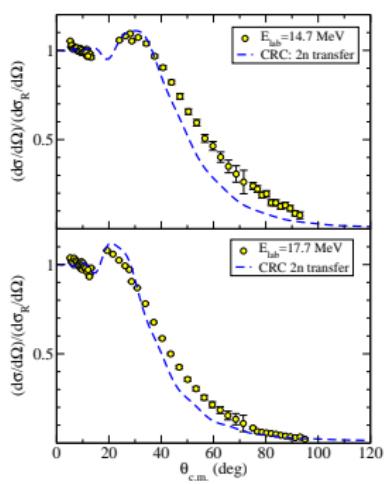
Break-up mechanism - ${}^6\text{He} + {}^{64}\text{Zn}$ 

Break-up mechanism - ${}^6\text{He} + {}^{64}\text{Zn}$

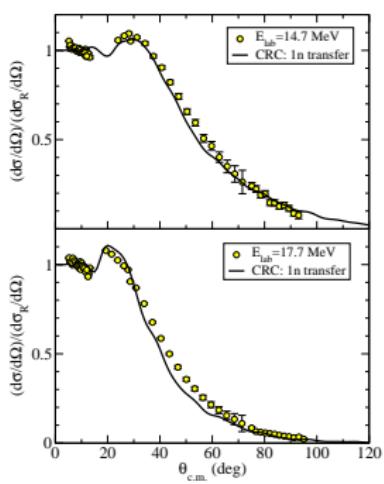
CDCC(DBU)



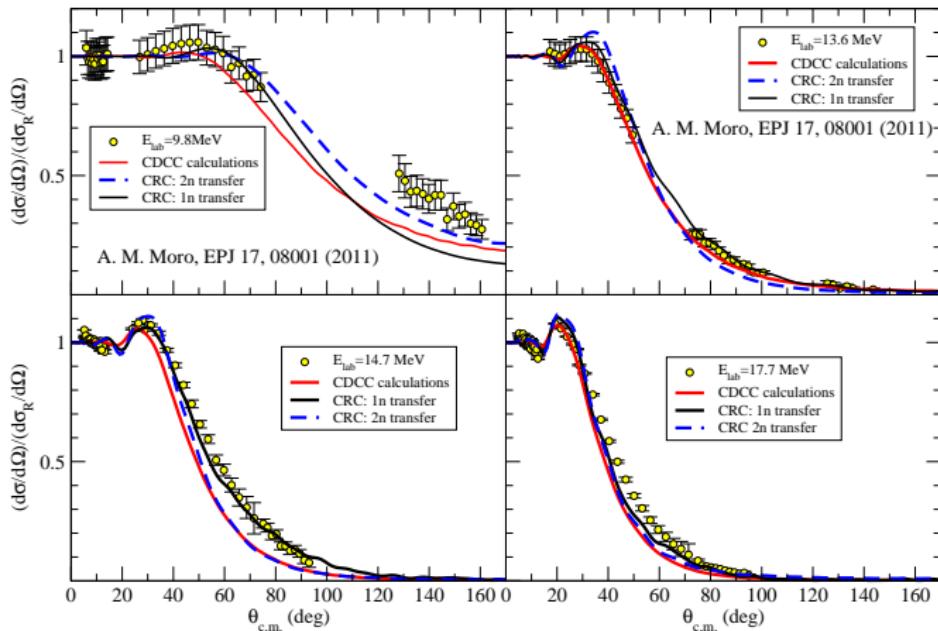
CRC(2n-TC)



CRC(1n-TC)



CRC calculations, based on 1n-transfer mechanism, reproduce the experimental data.

Break-up mechanism - ${}^6\text{He} + {}^{64}\text{Zn}$ 

CRC calculations, based on 1n-transfer mechanism, are in reasonable agreement with the experimental data.

Conclusions

- The reactions $^{6,7}\text{Li} + ^{64}\text{Zn}$ were measured at energies around the Coulomb barrier at LNS, Italy.
- CDCC calculations suggest more important effects of the coupling to the continuum states of the projectile in the reactions induced by ^6Li compared with the ones induced by ^7Li .

Conclusions

- New measurements of ${}^6\text{He} + {}^{64}\text{Zn}$ at energies above the Coulomb barrier have been presented.
- Optical Model calculations have been performed, showing long range effects.
- Small effect of the dipole Coulomb potential has been observed.
- Calculations including couplings to the continuum have been performed within different breakup mechanism.
- We have shown that the CRC method, based on one-neutron transfer mechanism, provides a more suitable approach.

Acknowledgment

Acknowledgment

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