

1st of July 2015 — Chiral Dynamics (Pisa, Italy)

# Inclusion of isospin breaking effects in lattice simulations

Antonin J. Portelli  
(University of Southampton)



---

# What's new ?

---



---

# What's new ?

---

- ❖ [MILC, 2014] — [Lattice 2014, arXiv:1409.7139]
  - update of quark masses and Dashen's theorem corrections using electro-quenched simulations
  - new insights on finite-volume effects



---

# What's new ?

---

- ❖ [MILC, 2014] — [Lattice 2014, arXiv:1409.7139]
  - update of quark masses and Dashen's theorem corrections using electro-quenched simulations
  - new insights on finite-volume effects
- ❖ [QCDSF, 2015] (pure QCD) — [PRD 91(7), p. 074512]
  - study of the  $\Sigma^0 - \Lambda^0$  system



---

# What's new ?

---

- ❖ [MILC, 2014] — [Lattice 2014, arXiv:1409.7139]
  - update of quark masses and Dashen's theorem corrections using electro-quenched simulations
  - new insights on finite-volume effects
- ❖ [QCDSF, 2015] (pure QCD) — [PRD 91(7), p. 074512]
  - study of the  $\Sigma^0 - \Lambda^0$  system
- ❖ [BMWc, 2014] (EQ) — [to appear]
  - update of quark masses and Dashen's theorem using electro-quenched simulations



---

# What's new ?

---

- ❖ [MILC, 2014] — [Lattice 2014, arXiv:1409.7139]
  - update of quark masses and Dashen's theorem corrections using electro-quenched simulations
  - new insights on finite-volume effects
- ❖ [QCDSF, 2015] (pure QCD) — [PRD 91(7), p. 074512]
  - study of the  $\Sigma^0 - \Lambda^0$  system
- ❖ [BMWc, 2014] (EQ) — [to appear]
  - update of quark masses and Dashen's theorem using electro-quenched simulations
- ❖ [Davoudi & Savage, 2014] — [PRD 90(5), p. 054503]
  - finite-volume corrections to hadron masses in NREFTs



---

# What's new ?

---

- ❖ [BMWc, 2015a] — [Science 347, pp. 1452–1455]
  - new set of  $N_f = 1+1+1+1$  full QCD+QED simulations
  - extensive analytical / numerical study of finite-volume effects
  - high precision computation of the hadron spectrum splittings (continuum, infinite volume and physical point extrapolation)



---

# What's new ?

---

- ❖ [BMWc, 2015a] — [Science 347, pp. 1452–1455]
  - new set of  $N_f = 1+1+1+1$  full QCD+QED simulations
  - extensive analytical / numerical study of finite-volume effects
  - high precision computation of the hadron spectrum splittings (continuum, infinite volume and physical point extrapolation)
- ❖ [BMWc, 2015b] — [arXiv:1502.06921]
  - further discussion of NREFT in finite volume



---

# What's new ?

---

- ❖ [BMWc, 2015a] — [Science 347, pp. 1452–1455]
  - new set of  $N_f = 1+1+1+1$  full QCD+QED simulations
  - extensive analytical / numerical study of finite-volume effects
  - high precision computation of the hadron spectrum splittings (continuum, infinite volume and physical point extrapolation)
- ❖ [BMWc, 2015b] — [arXiv:1502.06921]
  - further discussion of NREFT in finite volume
- ❖ possible summary of all that: [AP, 2015, arXiv:1505.07057]



---

# What's new ?

---

- ❖ [N. Carrasco *et. al*, 2015] — [arXiv:1502.00257]
  - theoretical study of the QED corrections to hadronic processes
  - *cf.* plenary talk tomorrow by V. Lubicz



---

# What's new ?

---

- ❖ [N. Carrasco *et. al*, 2015] — [arXiv:1502.00257]
  - theoretical study of the QED corrections to hadronic processes
  - *cf.* plenary talk tomorrow by V. Lubicz
- ❖ **Stay tuned:** Lattice 2015 (Kobe, Japan) is in two weeks



- ❖ Motivations
- ❖ Lattice QCD+QED
- ❖ Update on electro-quenched results
- ❖ Isospin splittings in the hadron spectrum
- ❖ Summary & outlook



# Motivations



---

# Isospin symmetry breaking

---

- ❖ Isospin symmetric world: up and down quarks are particles with identical physical properties.



# Isospin symmetry breaking

- ❖ Isospin symmetric world: up and down quarks are particles with identical physical properties.
- ❖ Isospin symmetry is explicitly broken by:

- the up and down **quark mass difference**

$$|m_u - m_d|/\Lambda_{\text{QCD}} \simeq 0.01$$

- the up and down **electric charge difference**

$$\alpha \simeq 0.0073$$

|            | up                    | down                  |
|------------|-----------------------|-----------------------|
| Mass (MeV) | $2.3(^{+0.7}_{-0.5})$ | $4.8(^{+0.5}_{-0.3})$ |
| Charge (e) | $2/3$                 | $-1/3$                |

source: [PDG, 2013]



---

# Nucleon mass splitting

---

❖ Well known experimentally:

$$M_n - M_p = 1.2933322(4) \text{ MeV}$$

source: [PDG, 2013]



---

# Nucleon mass splitting

---

- ❖ Well known experimentally:

$$M_n - M_p = 1.2933322(4) \text{ MeV}$$

source: [PDG, 2013]

- ❖ needed for **proton stability**



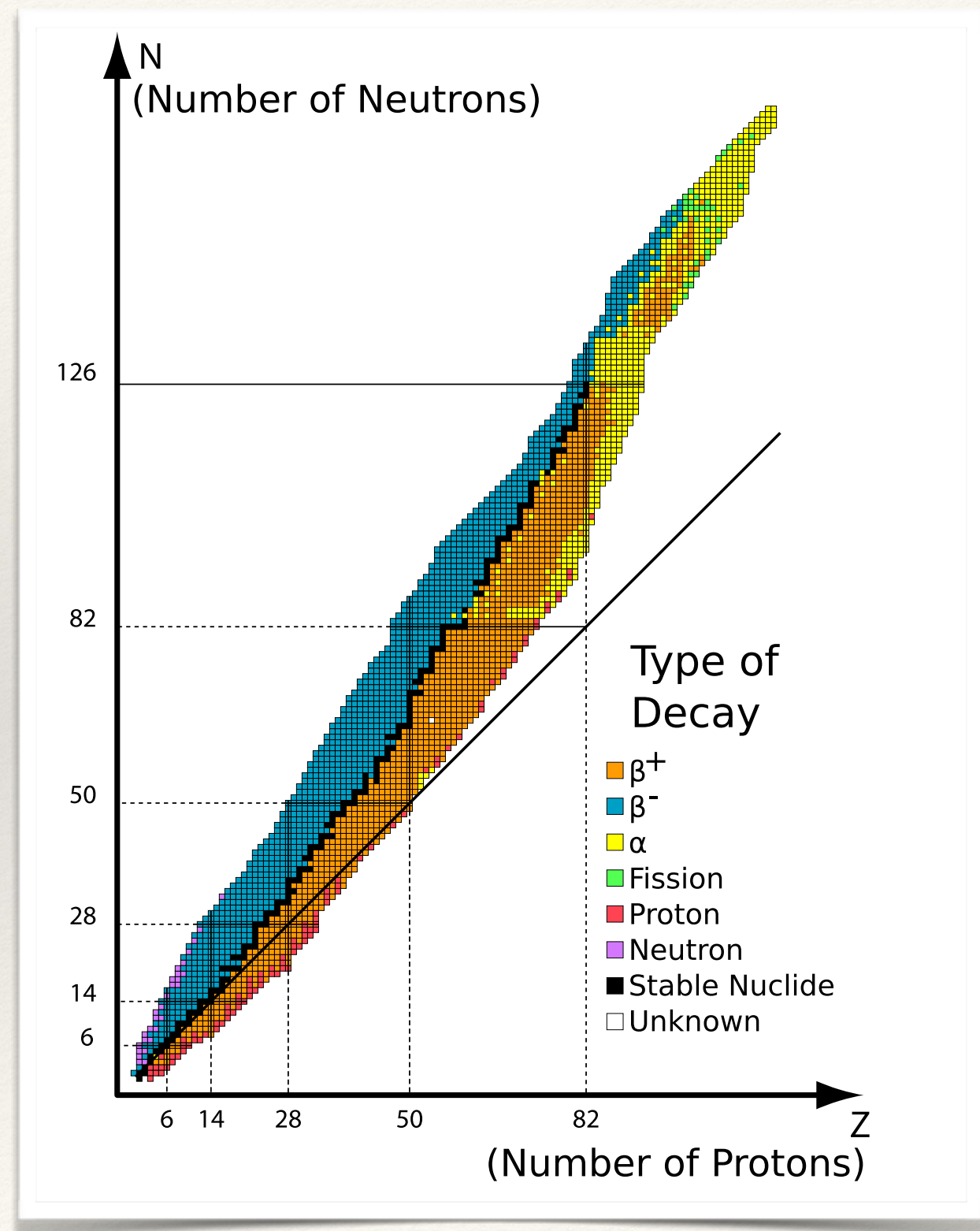
# Nucleon mass splitting

- ❖ Well known experimentally:

$$M_n - M_p = 1.2933322(4) \text{ MeV}$$

source: [PDG, 2013]

- ❖ needed for **proton stability**
- ❖ determines through  $\beta$ -decay the **stable nuclide chart**





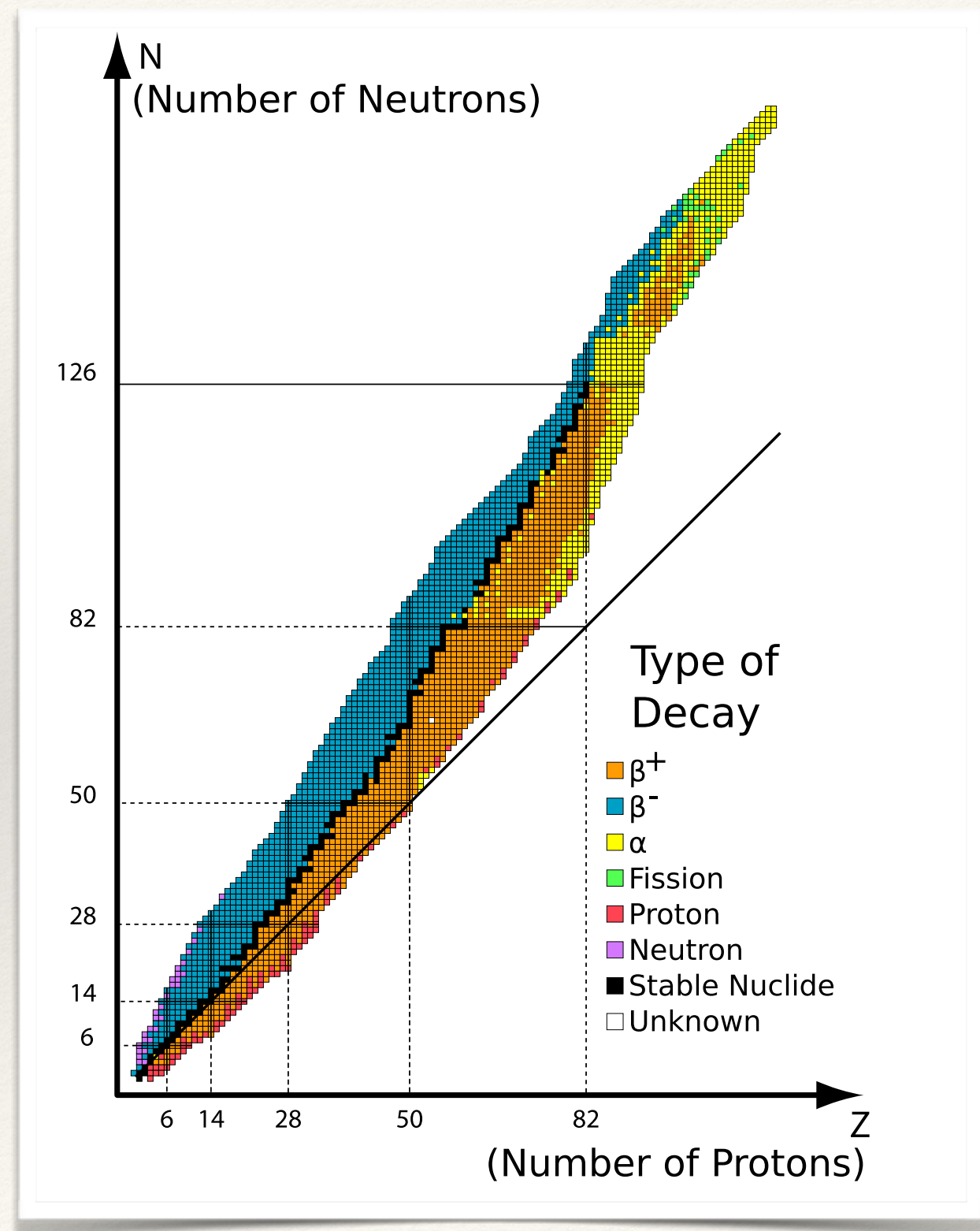
# Nucleon mass splitting

- ❖ Well known experimentally:

$$M_n - M_p = 1.2933322(4) \text{ MeV}$$

source: [PDG, 2013]

- ❖ needed for **proton stability**
- ❖ determines through  $\beta$ -decay the **stable nuclide chart**
- ❖ initial condition for **Big-Bang nucleosynthesis**





---

# Dashen's theorem

---

- ❖ In the SU(3) chiral limit [Dashen, 1969]:

$$\Delta_{\text{QED}} M_K^2 = \Delta_{\text{QED}} M_\pi^2 + \mathcal{O}(\alpha m_s)$$



---

# Dashen's theorem

---

- ❖ In the SU(3) chiral limit [Dashen, 1969]:

$$\Delta_{\text{QED}} M_K^2 = \Delta_{\text{QED}} M_\pi^2 + \mathcal{O}(\alpha m_s)$$

- ❖ How large are the corrections? FLAG parametrisation:

$$\varepsilon = \frac{\Delta_{\text{QED}} M_K^2 - \Delta_{\text{QED}} M_\pi^2}{\Delta M_\pi^2}$$



---

# Dashen's theorem

---

- ❖ In the SU(3) chiral limit [Dashen, 1969]:

$$\Delta_{\text{QED}} M_K^2 = \Delta_{\text{QED}} M_\pi^2 + \mathcal{O}(\alpha m_s)$$

- ❖ How large are the corrections? FLAG parametrisation:

$$\varepsilon = \frac{\Delta_{\text{QED}} M_K^2 - \Delta_{\text{QED}} M_\pi^2}{\Delta M_\pi^2}$$

- ❖  $\varepsilon$  is important to determine light quark mass ratios



# Lattice QCD+QED



---

# Lattice QCD

---

- ❖ Lattice QCD simulation: **Monte-Carlo estimation of discretised QCD functional integrals**



---

# Lattice QCD

---

- ❖ Lattice QCD simulation: **Monte-Carlo estimation of discretised QCD functional integrals**
- ❖ Discretised Yang-Mills action: [K. Wilson, 1974]



---

# Lattice QCD

---

- ❖ Lattice QCD simulation: **Monte-Carlo estimation of discretised QCD functional integrals**
- ❖ Discretised Yang-Mills action: [K. Wilson, 1974]
- ❖ Discretised Dirac action: chiral symmetry must be broken (Nielsen-Ninomiya theorem), **many possible solutions**



---

# Lattice QCD

---

- ❖ Lattice QCD simulation: **Monte-Carlo estimation of discretised QCD functional integrals**
- ❖ Discretised Yang-Mills action: [K. Wilson, 1974]
- ❖ Discretised Dirac action: chiral symmetry must be broken (Nielsen-Ninomiya theorem), **many possible solutions**
- ❖ Fermionic integrals can be performed analytically (Wick's contractions)



---

# Lattice QCD

---

- ❖ Lattice QCD simulation: **Monte-Carlo estimation of discretised QCD functional integrals**
- ❖ Discretised Yang-Mills action: [K. Wilson, 1974]
- ❖ Discretised Dirac action: chiral symmetry must be broken (Nielsen-Ninomiya theorem), **many possible solutions**
- ❖ Fermionic integrals can be performed analytically (Wick's contractions)
- ❖ Gauge integrals are computed stochastically



---

# Lattice QCD

---

- ❖ Lattice QCD simulation: **Monte-Carlo estimation of discretised QCD functional integrals**
- ❖ Discretised Yang-Mills action: [K. Wilson, 1974]
- ❖ Discretised Dirac action: chiral symmetry must be broken (Nielsen-Ninomiya theorem), **many possible solutions**
- ❖ Fermionic integrals can be performed analytically (Wick's contractions)
- ❖ Gauge integrals are computed stochastically
- ❖ **Extremely expensive, but *ab-initio***



---

# Non-compact lattice QED

---

- ❖ Naively discretised **Maxwell action**:

$$S[A_\mu] = \frac{1}{4} \sum_{\mu, \nu} (\partial_\mu A_\nu - \partial_\nu A_\mu)^2$$



---

# Non-compact lattice QED

---

- ❖ Naively discretised **Maxwell action**:

$$S[A_\mu] = \frac{1}{4} \sum_{\mu, \nu} (\partial_\mu A_\nu - \partial_\nu A_\mu)^2$$

- ❖ Pure gauge theory is **free**, it can be solved **exactly**



---

# Non-compact lattice QED

---

- ❖ Naively discretised **Maxwell action**:

$$S[A_\mu] = \frac{1}{4} \sum_{\mu, \nu} (\partial_\mu A_\nu - \partial_\nu A_\mu)^2$$

- ❖ Pure gauge theory is **free**, it can be solved **exactly**
- ❖ Gauge invariance is preserved



---

# Non-compact lattice QED

---

- ❖ Naively discretised **Maxwell action**:

$$S[A_\mu] = \frac{1}{4} \sum_{\mu, \nu} (\partial_\mu A_\nu - \partial_\nu A_\mu)^2$$

- ❖ Pure gauge theory is **free**, it can be solved **exactly**
- ❖ Gauge invariance is preserved
- ❖ No mass gap: **large finite volume effects expected**



---

# Zero-mode subtraction

---

Finite volume: **momentum quantisation**

$$\alpha \int \frac{d^4 k}{(2\pi)^4} \frac{1}{k^2} \cdots \quad \mapsto \quad \frac{\alpha}{V} \sum_k \frac{1}{k^2} \cdots$$



---

# Zero-mode subtraction

---

Finite volume: **momentum quantisation**

$$\alpha \int \frac{d^4 k}{(2\pi)^4} \frac{1}{k^2} \cdots \quad \mapsto \quad \frac{\alpha}{V} \sum_k \frac{1}{k^2} \cdots$$



Possibly IR divergent, but  
not for physical quantities



---

# Zero-mode subtraction

---

Finite volume: **momentum quantisation**

$$\alpha \int \frac{d^4 k}{(2\pi)^4} \frac{1}{k^2} \cdots \quad \longrightarrow \quad \frac{\alpha}{V} \sum_k \frac{1}{k^2} \cdots$$



Possibly IR divergent, but  
not for physical quantities



Contains a straight  $1/0$  !



---

# Zero-mode subtraction

---

- ❖ This problem can be solved by **removing zero modes**



---

# Zero-mode subtraction

---

- ❖ This problem can be solved by **removing zero modes**
- ❖ **Many possible schemes:**  
modification of  $A_\mu(k)$  on a set of measure 0



---

# Zero-mode subtraction

---

- ❖ This problem can be solved by **removing zero modes**
- ❖ **Many possible schemes:**  
modification of  $A_\mu(k)$  on a set of measure 0
- ❖ Different schemes: **different finite volume behaviours**



---

# Zero-mode subtraction

---

- ❖ This problem can be solved by **removing zero modes**
- ❖ **Many possible schemes:**  
modification of  $A_\mu(k)$  on a set of measure 0
- ❖ Different schemes: **different finite volume behaviours**
- ❖ Some more interesting than others



---

# QED<sub>TL</sub> zero-mode subtraction

---

❖ QED<sub>TL</sub>:  $A_\mu(0) = 0$

**Mostly used in all simulations so far**



---

# QED<sub>TL</sub> zero-mode subtraction

---

- ❖ QED<sub>TL</sub>:  $A_\mu(0) = 0$

**Mostly used in all simulations so far**

- ❖ With QED<sub>TL</sub>, the  $T \rightarrow \infty$ ,  $L = \text{cst.}$  limit **can diverge**:

$$\frac{\alpha}{V} \sum_{k \neq 0} \frac{1}{k^2} \cdots \quad \longmapsto \quad \frac{\alpha}{L^3} \int \frac{dk_0}{2\pi} \sum_{\mathbf{k}} \frac{1}{k^2} \cdots$$



---

# QED<sub>TL</sub> zero-mode subtraction

---

- ❖ QED<sub>TL</sub>:  $A_\mu(0) = 0$

**Mostly used in all simulations so far**

- ❖ With QED<sub>TL</sub>, the  $T \rightarrow \infty$ ,  $L = \text{cst.}$  limit **can diverge**:

$$\frac{\alpha}{V} \sum_{k \neq 0} \frac{1}{k^2} \cdots \quad \longmapsto \quad \frac{\alpha}{L^3} \int \frac{dk_0}{2\pi} \sum_{\mathbf{k}} \frac{1}{k^2} \cdots$$

- ❖ QED<sub>TL</sub> **does not have reflection positivity**



---

# QED<sub>TL</sub> finite-volume effects

---

❖ Example — 1-loop QED<sub>TL</sub> [BMWc, 2014]:

$$m(T, L) \underset{T, L \rightarrow +\infty}{\sim} m \left\{ 1 - q^2 \alpha \left[ \frac{\kappa}{2mL} \left( 1 + \frac{2}{mL} \left[ 1 - \frac{\pi}{2\kappa} \frac{T}{L} \right] \right) - \frac{3\pi}{(mL)^3} \left[ 1 - \frac{\coth(mT)}{2} \right] - \frac{3\pi}{2(mL)^4} \frac{L}{T} \right] \right\}$$

up to exponential corrections, with  $\kappa = 2.83729 \dots$



---

# QED<sub>TL</sub> finite-volume effects

---

- ❖ Example — 1-loop QED<sub>TL</sub> [BMWc, 2014]:

$$m(T, L) \underset{T, L \rightarrow +\infty}{\sim} m \left\{ 1 - q^2 \alpha \left[ \frac{\kappa}{2mL} \left( 1 + \frac{2}{mL} \left[ 1 - \frac{\pi}{2\kappa} \frac{T}{L} \right] \right) - \frac{3\pi}{(mL)^3} \left[ 1 - \frac{\coth(mT)}{2} \right] - \frac{3\pi}{2(mL)^4} \frac{L}{T} \right] \right\}$$

up to exponential corrections, with  $\kappa = 2.83729 \dots$

- ❖ **Divergent finite volume effects** with  $T \rightarrow \infty$ ,  $L = \text{cst.}$



---

# QED<sub>TL</sub> finite-volume effects

---

- ❖ Example — 1-loop QED<sub>TL</sub> [BMWc, 2014]:

$$m(T, L) \underset{T, L \rightarrow +\infty}{\sim} m \left\{ 1 - q^2 \alpha \left[ \frac{\kappa}{2mL} \left( 1 + \frac{2}{mL} \left[ 1 - \frac{\pi}{2\kappa} \frac{T}{L} \right] \right) - \frac{3\pi}{(mL)^3} \left[ 1 - \frac{\coth(mT)}{2} \right] - \frac{3\pi}{2(mL)^4} \frac{L}{T} \right] \right\}$$

up to exponential corrections, with  $\kappa = 2.83729 \dots$

- ❖ **Divergent finite volume effects** with  $T \rightarrow \infty$ ,  $L = \text{cst.}$
- ❖ Same behaviour independently discovered by MILC



---

# QED<sub>L</sub> zero-mode subtraction

---

- ❖ QED<sub>L</sub>:  $A_\mu(k_0, \mathbf{0}) = 0$   
inspired from [Hayakawa & Uno, 2008]



---

# QED<sub>L</sub> zero-mode subtraction

---

- ❖ QED<sub>L</sub>:  $A_\mu(k_0, \mathbf{0}) = 0$   
inspired from [Hayakawa & Uno, 2008]
- ❖ QED<sub>L</sub> maintains reflection positivity [BMWc, 2015a]:



---

# QED<sub>L</sub> zero-mode subtraction

---

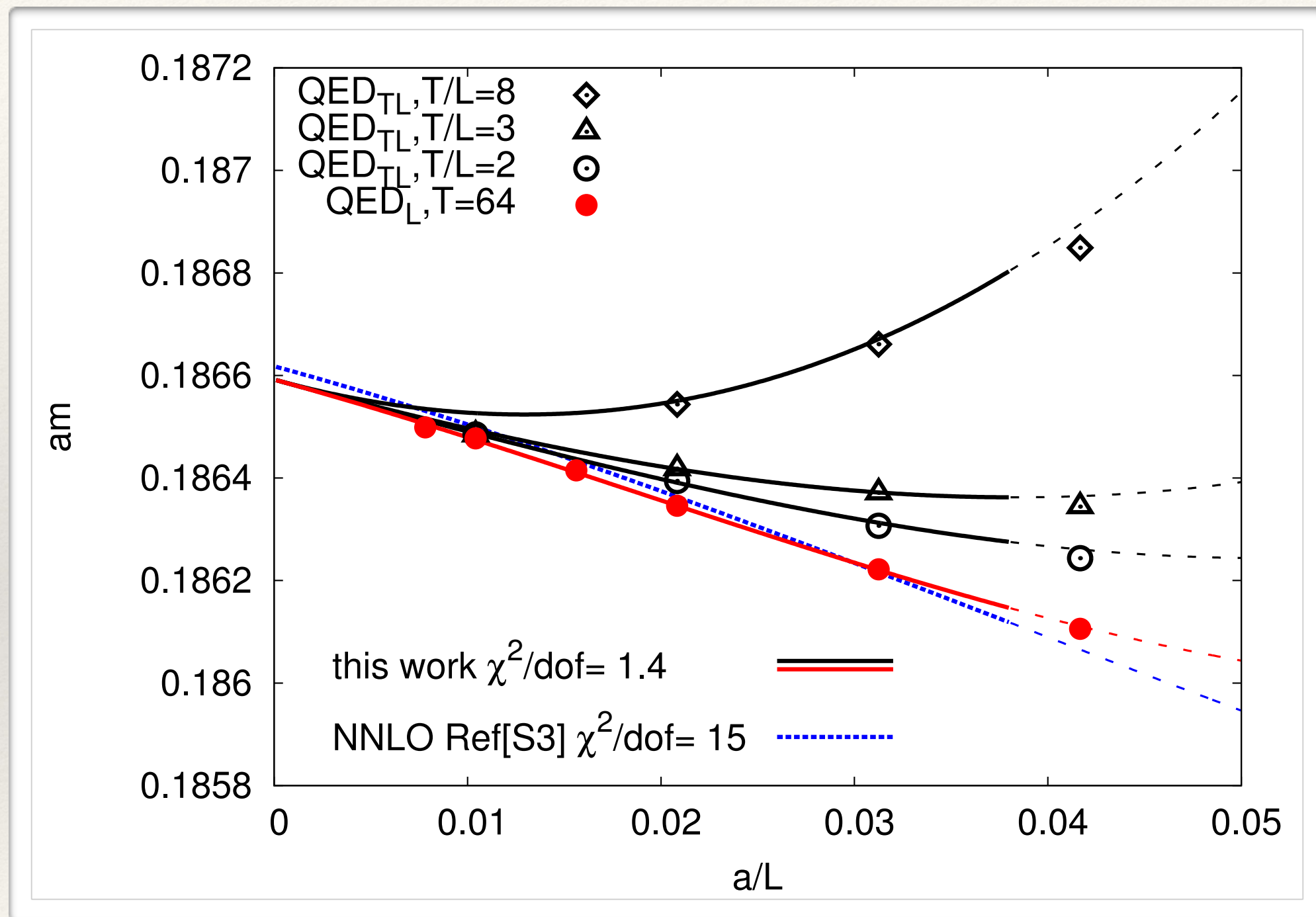
- ❖ QED<sub>L</sub>:  $A_\mu(k_0, \mathbf{0}) = 0$   
inspired from [Hayakawa & Uno, 2008]
- ❖ QED<sub>L</sub> maintains reflection positivity [BMWc, 2015a]:
- ❖ QED<sub>L</sub> finite volume effects:

$$m(T, L) \underset{T, L \rightarrow +\infty}{\sim} m \left\{ 1 - q^2 \alpha \left[ \frac{\kappa}{2mL} \left( 1 + \frac{2}{mL} \right) - \frac{3\pi}{(mL)^3} \right] \right\}$$

**inverse powers of L, independent of T**



# Finite-volume effects



Pure QED simulations (quenched) from [BMWc, 2015a] — [S3]=[Davoudi & Savage, 2014]



---

# Finite volume NRQED

---

- ❖ Anti-particles and particles do not decouple completely because of the missing photon modes



---

# Finite volume NRQED

---

- ❖ Anti-particles and particles do not decouple completely because of the missing photon modes
- ❖ The residual contribution generates a  $O(1/L^3)$  finite volume correction to the self-energy



---

# Finite volume NRQED

---

- ❖ Anti-particles and particles do not decouple completely because of the missing photon modes
- ❖ The residual contribution generates a  $O(1/L^3)$  finite volume correction to the self-energy
- ❖ This contribution is absent from [D & S, 2014], explaining the observed discrepancy



---

# Finite volume NRQED

---

- ❖ Anti-particles and particles do not decouple completely because of the missing photon modes
- ❖ The residual contribution generates a  $O(1/L^3)$  finite volume correction to the self-energy
- ❖ This contribution is absent from [D & S, 2014], explaining the observed discrepancy
- ❖ More details in [BMWc, 2015b]



---

# Finite-volume effects

---

- ❖ What about **composite particles** (QCD + QED)?



---

# Finite-volume effects

---

- ❖ What about **composite particles** (QCD + QED)?
- ❖ [Hayakawa & Uno, 2008]: SU(3) PQChPT



---

# Finite-volume effects

---

- ❖ What about **composite particles** (QCD + QED)?
- ❖ [Hayakawa & Uno, 2008]: SU(3) PQChPT
- ❖ [RBC-UKQCD, 2010]: SU(2) PQChPT + heavy kaons



---

# Finite-volume effects

---

- ❖ What about **composite particles** (QCD + QED)?
- ❖ [Hayakawa & Uno, 2008]: SU(3) PQChPT
- ❖ [RBC-UKQCD, 2010]: SU(2) PQChPT + heavy kaons
- ❖ [Davoudi & Savage, 2014]: NREFTs  
mesons, baryons, nuclei and HVP

$$m(L) \underset{L \rightarrow +\infty}{\sim} m \left\{ 1 - q^2 \alpha \left[ \frac{\kappa}{2mL} \left( 1 + \frac{2}{mL} \right) + \mathcal{O} \left( \frac{1}{L^3} \right) \right] \right\}$$



---

# Finite-volume effects

---

- ❖ What about **composite particles** (QCD + QED)?
- ❖ [Hayakawa & Uno, 2008]: SU(3) PQChPT
- ❖ [RBC-UKQCD, 2010]: SU(2) PQChPT + heavy kaons
- ❖ [Davoudi & Savage, 2014]: NREFTs  
mesons, baryons, nuclei and HVP

$$m(L) \underset{L \rightarrow +\infty}{\sim} m \left\{ 1 - q^2 \alpha \left[ \frac{\kappa}{2mL} \left( 1 + \frac{2}{mL} \right) + \mathcal{O} \left( \frac{1}{L^3} \right) \right] \right\}$$

- ❖ [BMWc, 2015a]: Ward identities: **NLO is universal**



---

# Electro-quenched approximation

---

- ❖ Electro-quenched approximation: **charged valence quarks, but neutral sea quarks**



---

# Electro-quenched approximation

---

- ❖ Electro-quenched approximation: **charged valence quarks**, but **neutral sea quarks**
- ❖ **Non-unitary** theory (partially quenched)



---

# Electro-quenched approximation

---

- ❖ Electro-quenched approximation: **charged valence quarks**, but **neutral sea quarks**
- ❖ **Non-unitary** theory (partially quenched)
- ❖ **Greatly reduce** the computational cost



---

# Electro-quenched approximation

---

- ❖ Electro-quenched approximation: **charged valence quarks, but neutral sea quarks**
- ❖ **Non-unitary** theory (partially quenched)
- ❖ **Greatly reduce** the computational cost
- ❖ Missing contributions are large- $N_c$  and SU(3) flavour suppressed:  $O(10\%)$  of EM effects



---

# Electro-quenched approximation

---

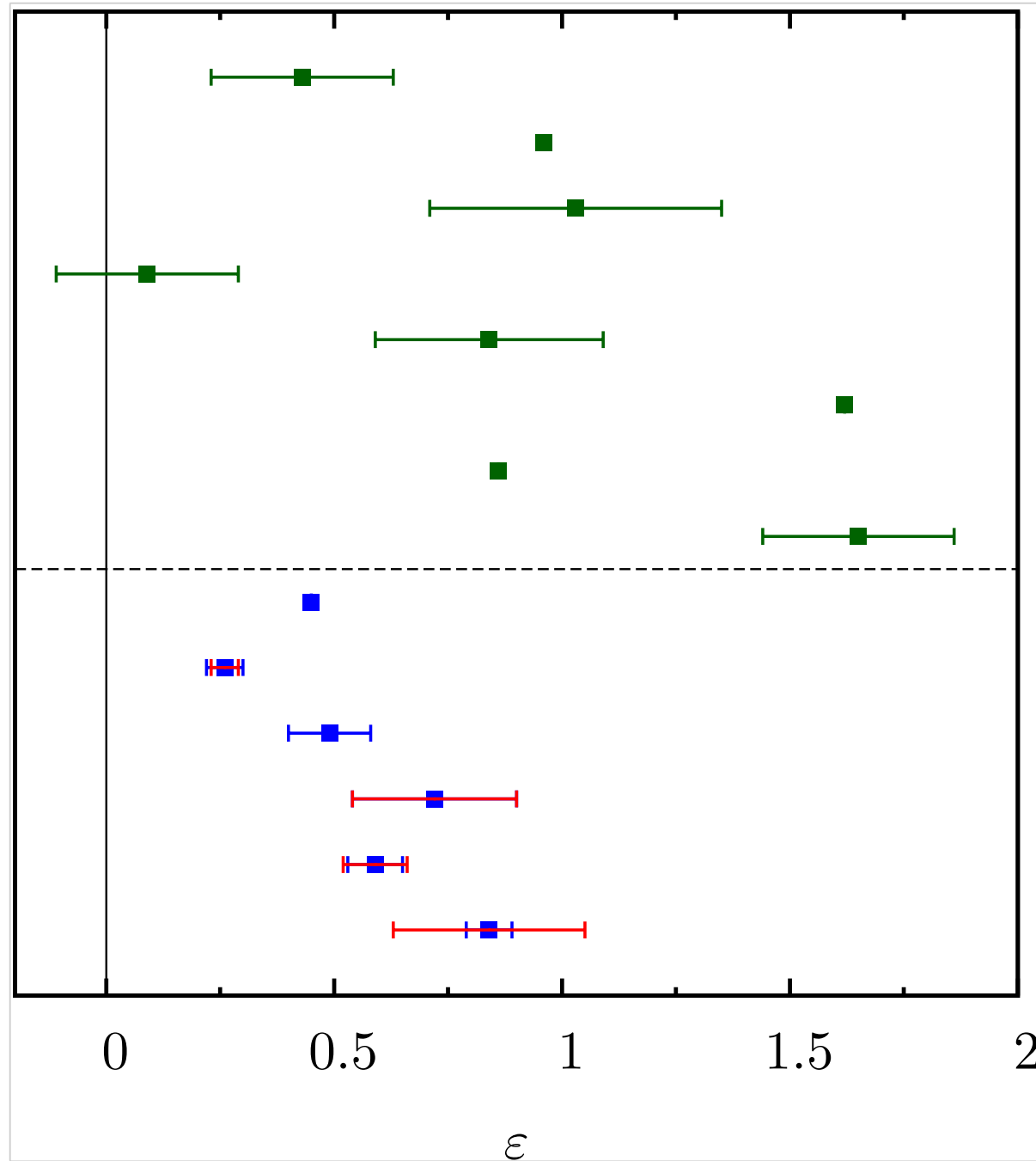
- ❖ Electro-quenched approximation: **charged valence quarks, but neutral sea quarks**
- ❖ **Non-unitary** theory (partially quenched)
- ❖ **Greatly reduce** the computational cost
- ❖ Missing contributions are large- $N_c$  and SU(3) flavour suppressed:  $O(10\%)$  of EM effects
- ❖ In agreement with PQChPT estimates  
[J. Bijnens & N. Danielsson, PRD 75(1), p. 014505, 2007]



# Update on electro-quenched results



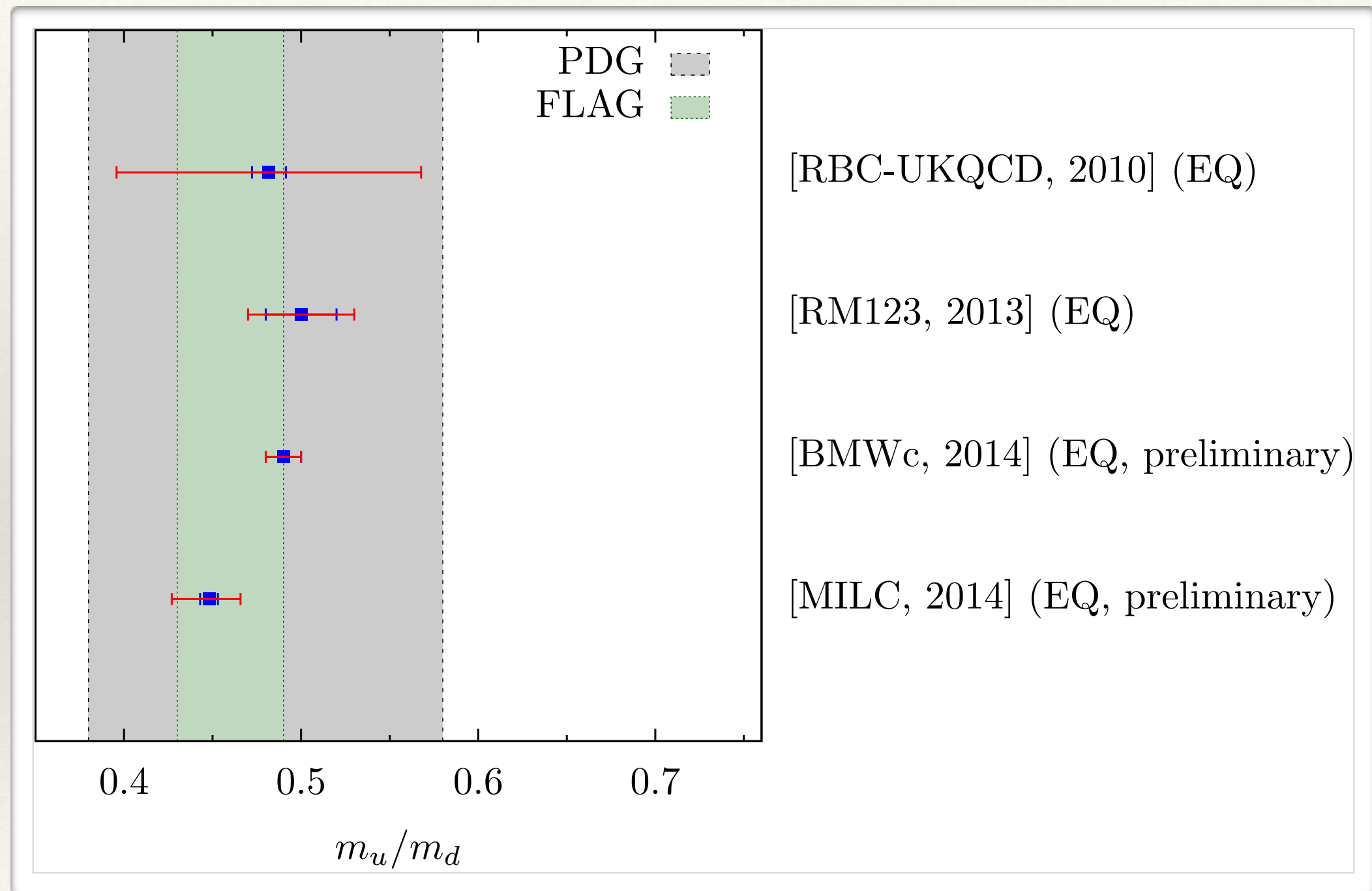
# EQ results for $\varepsilon$



- [Maltman and Kotchan, 1990]
- [Donoghue *et al.*, 1993]
- [Bijnens, 1993]
- [Baur and Urech, 1996]
- [Bijnens and Prades, 1997]
- [Donoghue and Perez, 1997]
- [Gao *et al.*, 1997]
- [Moussallam, 1997]
- [Duncan *et al.*, 1996] (quenched QCD)
- [RBC-UKQCD, 2007]
- [RBC-UKQCD, 2010]
- [RM123, 2013]
- [BMWc, 2014] (EQ, preliminary)
- [MILC, 2014] (preliminary)



# EQ results for light quark masses





# Isospin splittings in the hadron spectrum



---

# [BMWc, 2015a]: mass splitting calculation

---

- ❖ **many smeared sources per configurations ( $O(100)$ )**



---

# [BMWc, 2015a]: mass splitting calculation

---

- ❖ **many smeared sources** per configurations ( $O(100)$ )
- ❖ electric charge renormalisation using **Wilson flow**



---

# [BMWc, 2015a]: mass splitting calculation

---

- ❖ **many smeared sources** per configurations ( $O(100)$ )
- ❖ electric charge renormalisation using **Wilson flow**
- ❖ small extrapolation to the physical point  
(similar to [BMWc, 2013])



---

## [BMWc, 2015a]: mass splitting calculation

---

- ❖ **many smeared sources** per configurations ( $O(100)$ )
- ❖ electric charge renormalisation using **Wilson flow**
- ❖ small extrapolation to the physical point (similar to [BMWc, 2013])
- ❖ Systematic error based on BMW's histogram method. Weights are based on the goodness of the fits, flat and Akaike's information criterion (**overfitting is penalised**)



---

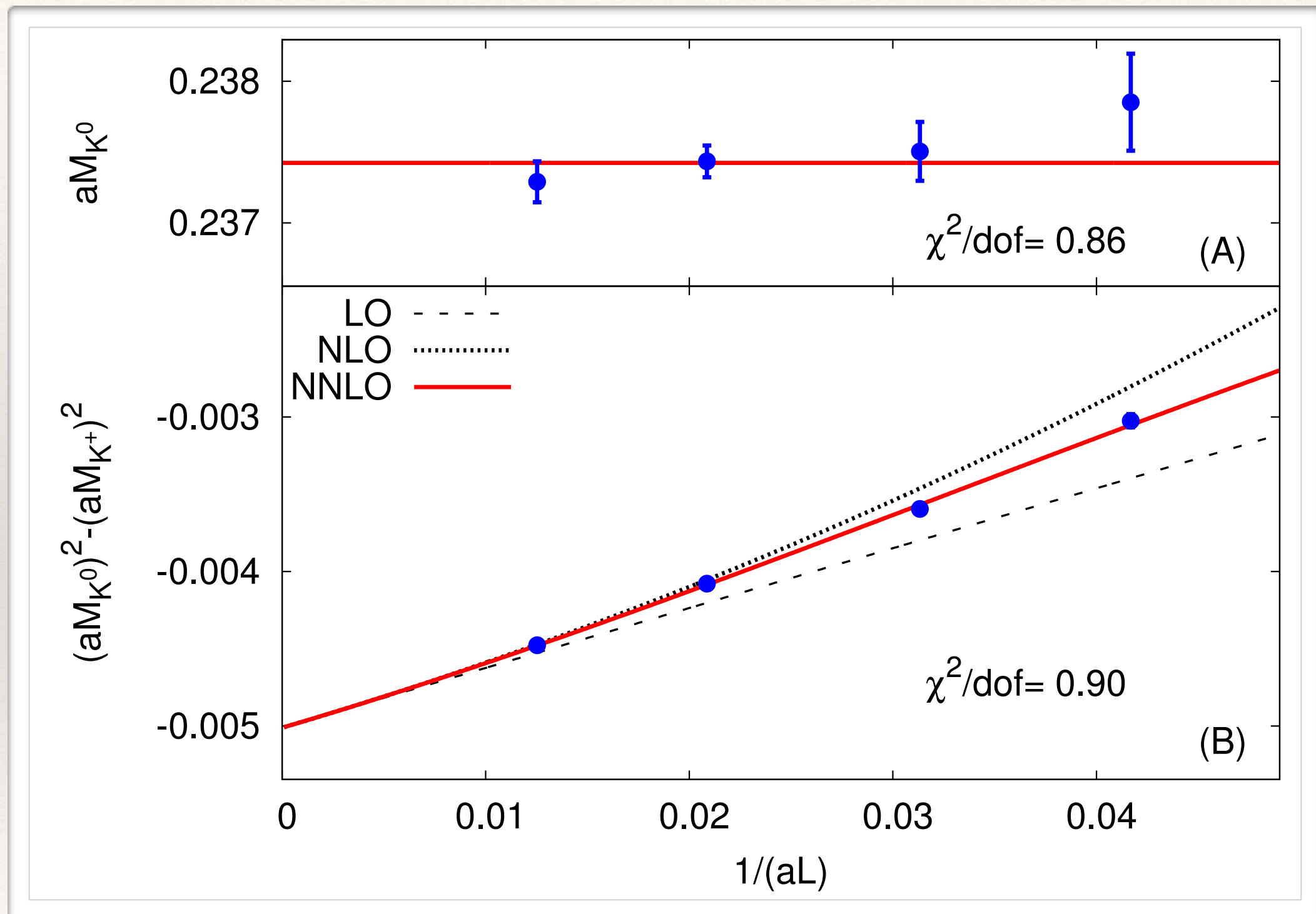
# [BMWc, 2015a]: mass splitting calculation

---

- ❖ **many smeared sources** per configurations ( $O(100)$ )
- ❖ electric charge renormalisation using **Wilson flow**
- ❖ small extrapolation to the physical point (similar to [BMWc, 2013])
- ❖ Systematic error based on BMW's histogram method. Weights are based on the goodness of the fits, flat and Akaike's information criterion (**overfitting is penalised**)
- ❖  **$O(500)$  analyses per mass splitting**

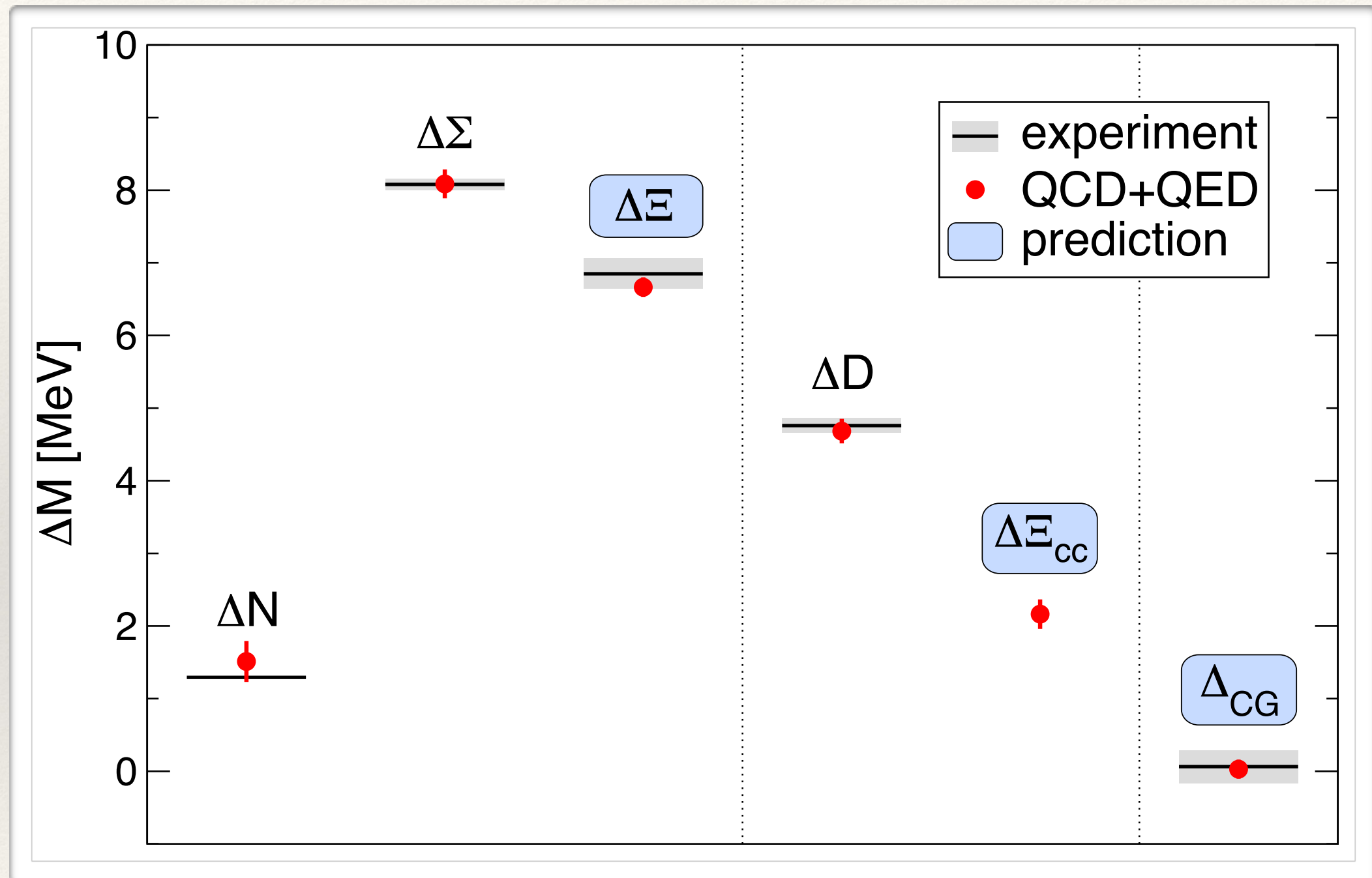


# [BMWc, 2015a]: finite-volume study





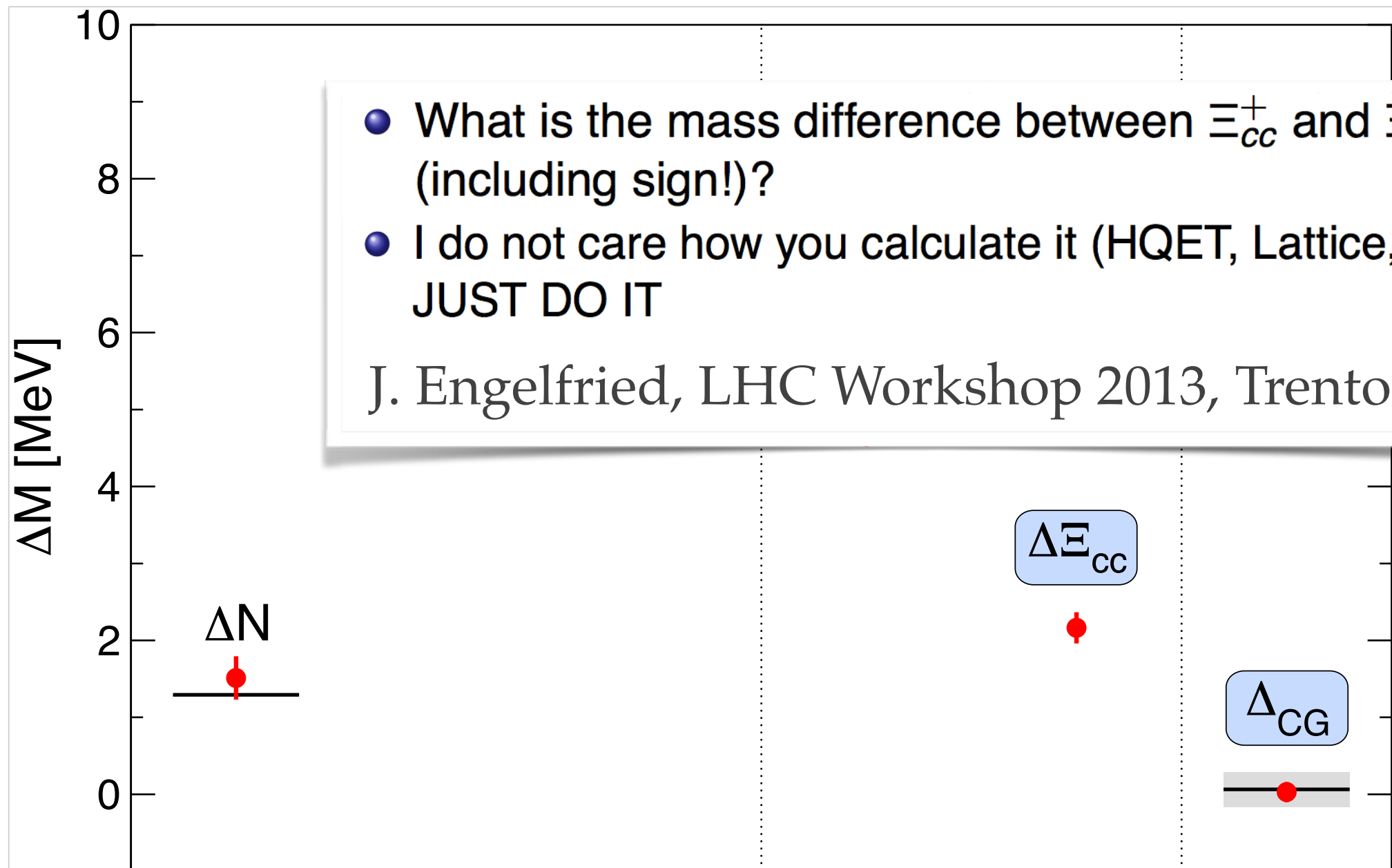
# [BMWc, 2015a]: result summary



$$\Delta_{CG} = \Delta M_N - \Delta M_\Sigma + \Delta M_\Xi \text{ (Coleman-Glashow relation)}$$



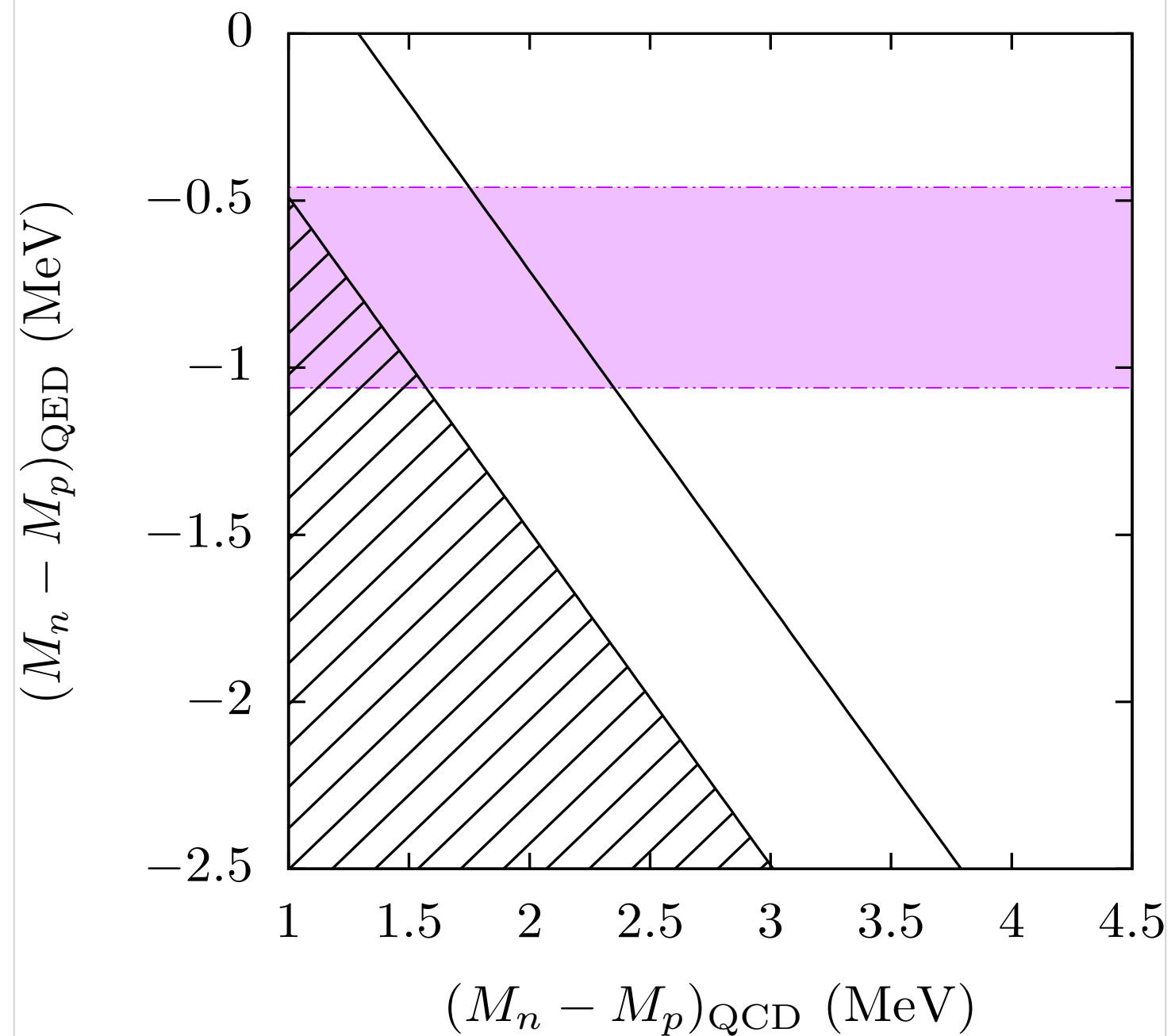
# [BMWc, 2015a]: result summary



$$\Delta_{CG} = \Delta M_N - \Delta M_\Sigma + \Delta M_\Xi \text{ (Coleman-Glashow relation)}$$



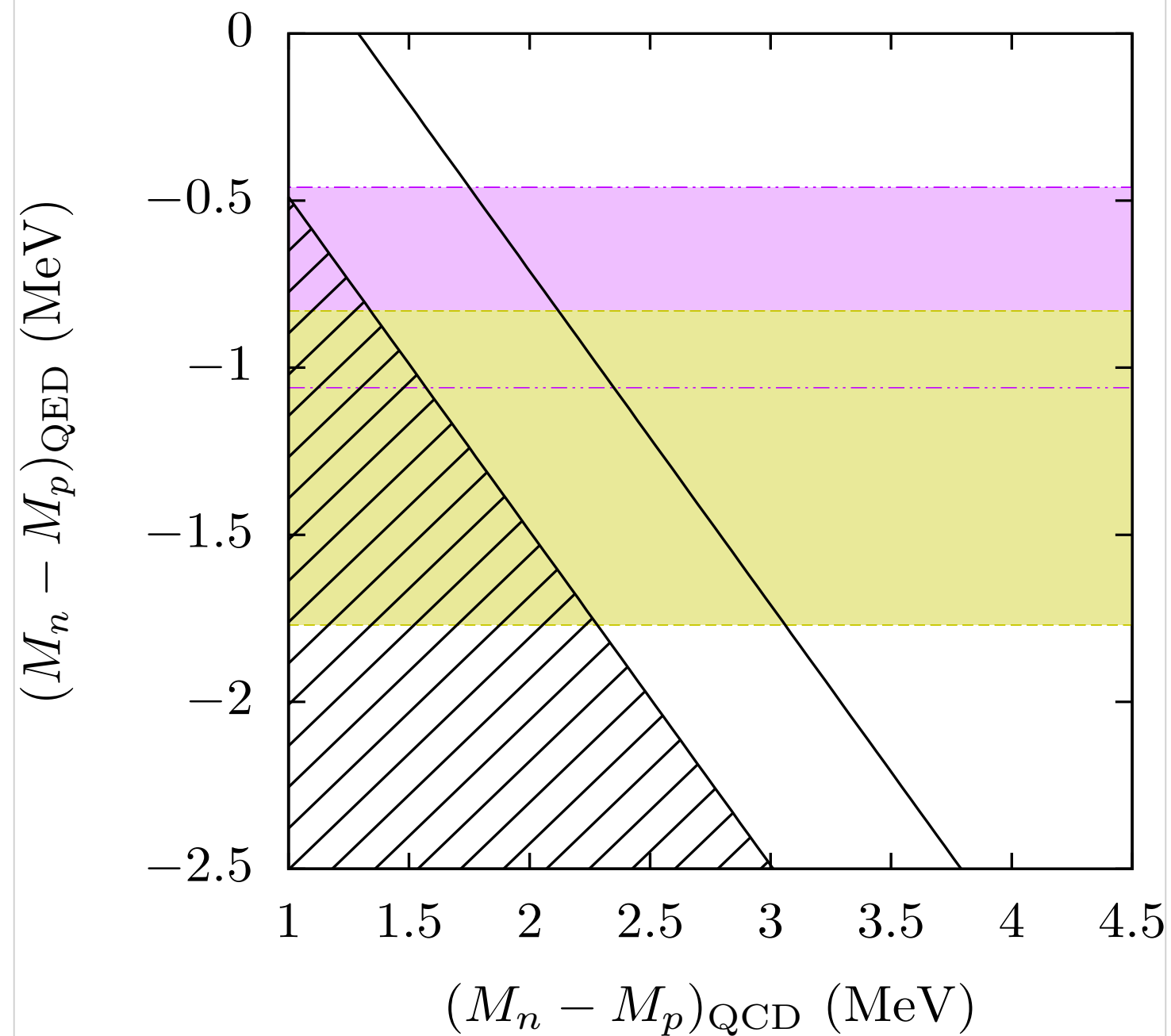
# Results for the nucleon mass splitting



- [Gasser & Leutwyler, 1982]
- no *beta*-decay
- experiment

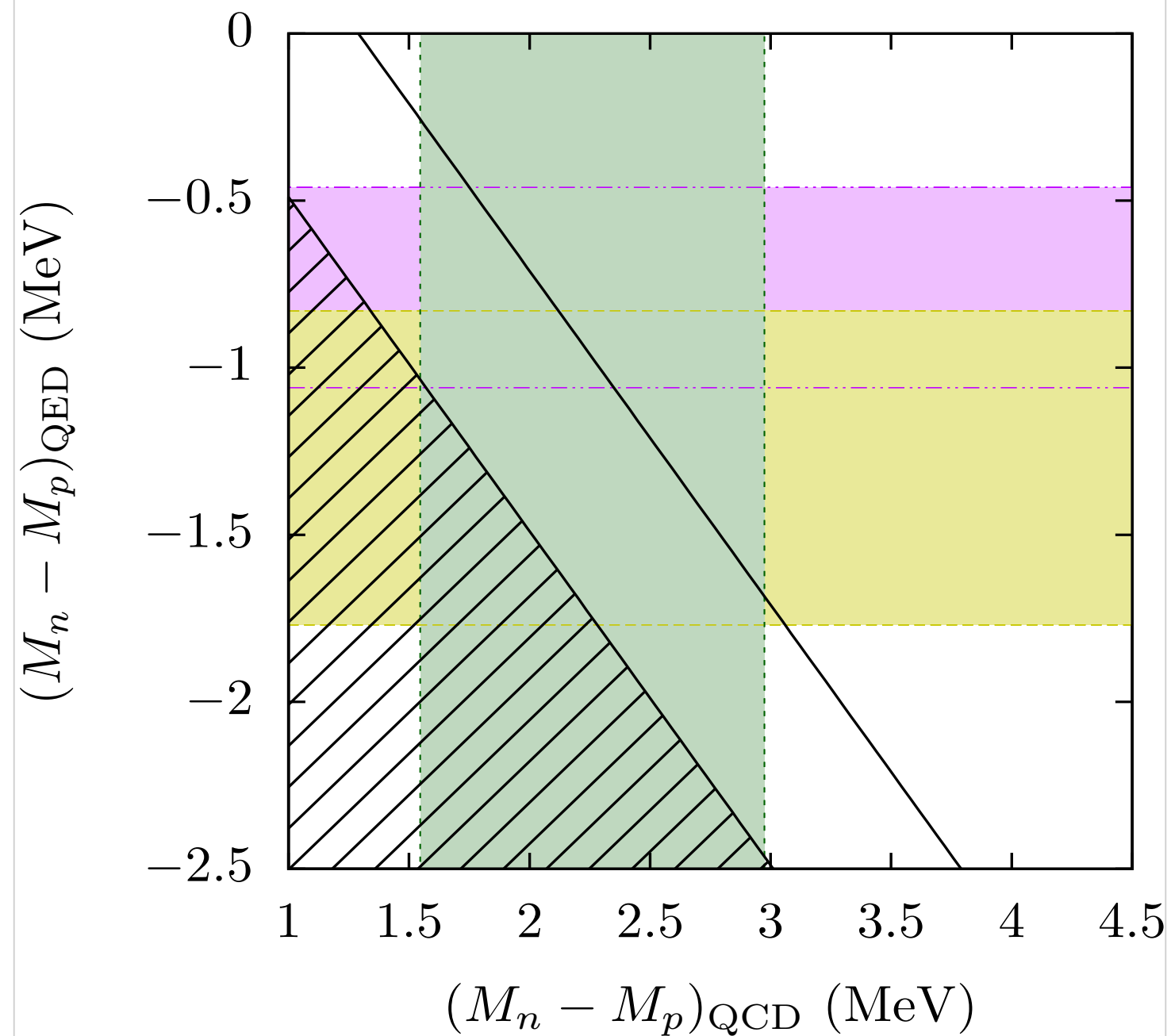


# Results for the nucleon mass splitting



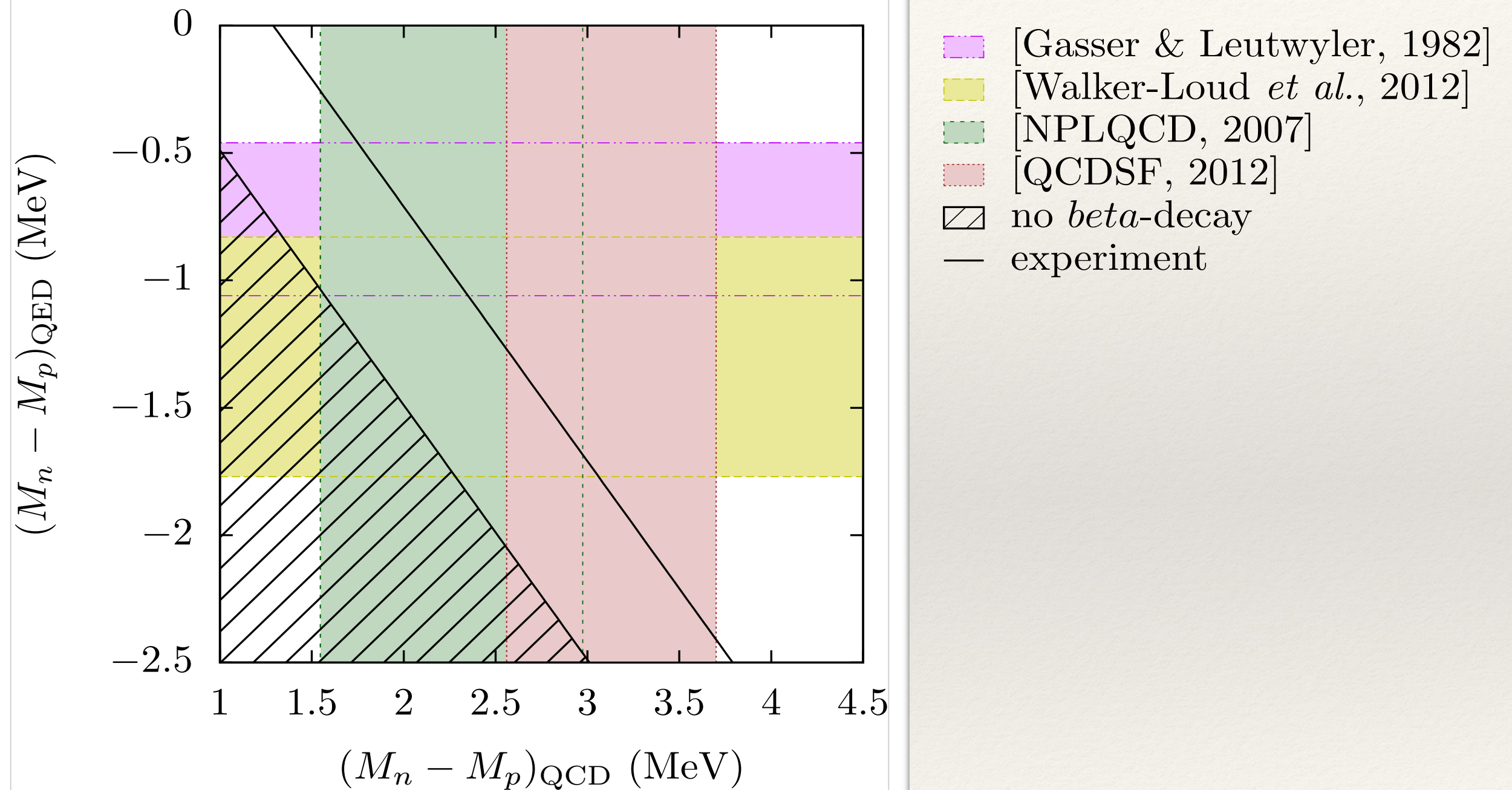


# Results for the nucleon mass splitting



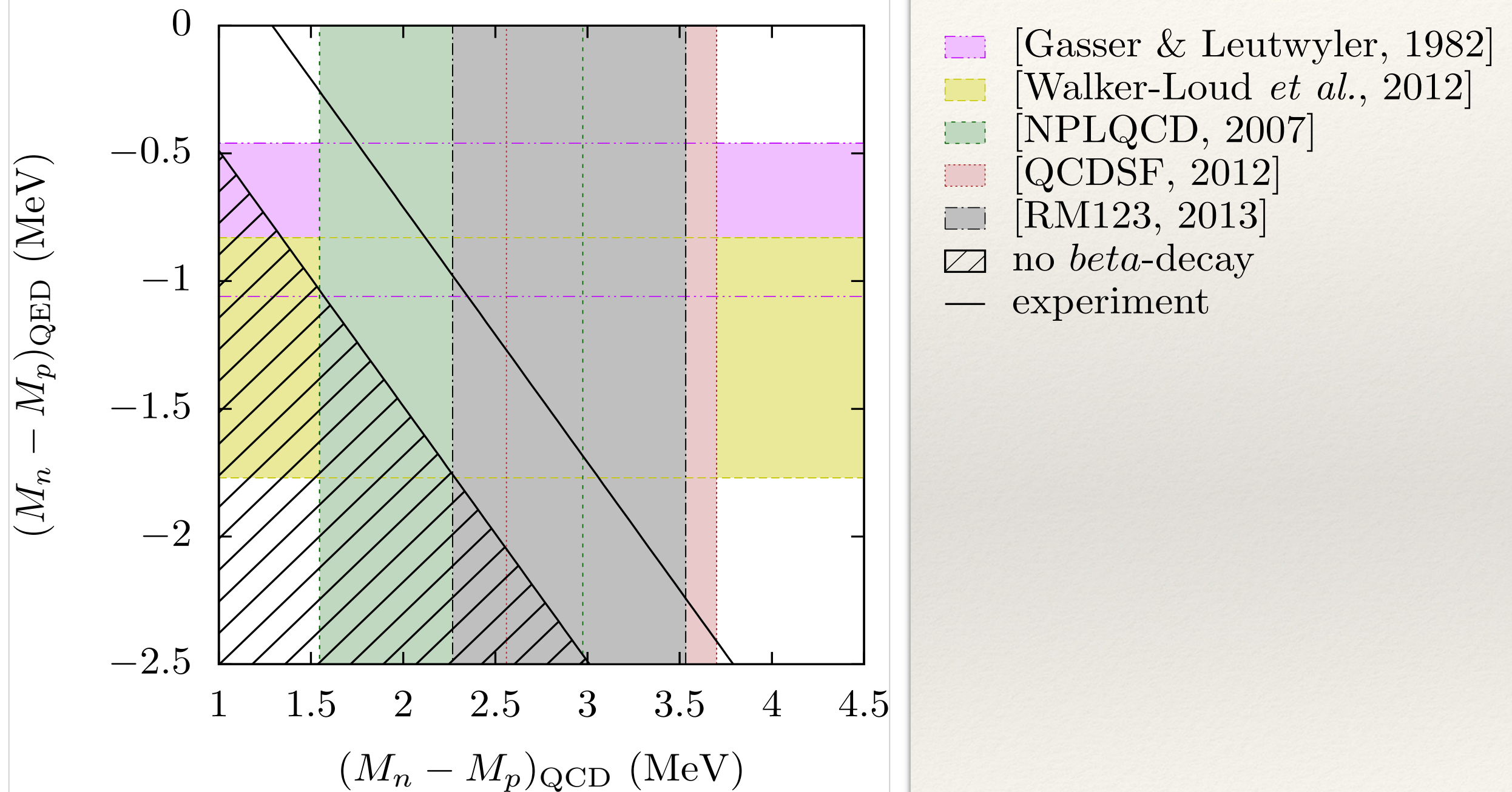


# Results for the nucleon mass splitting



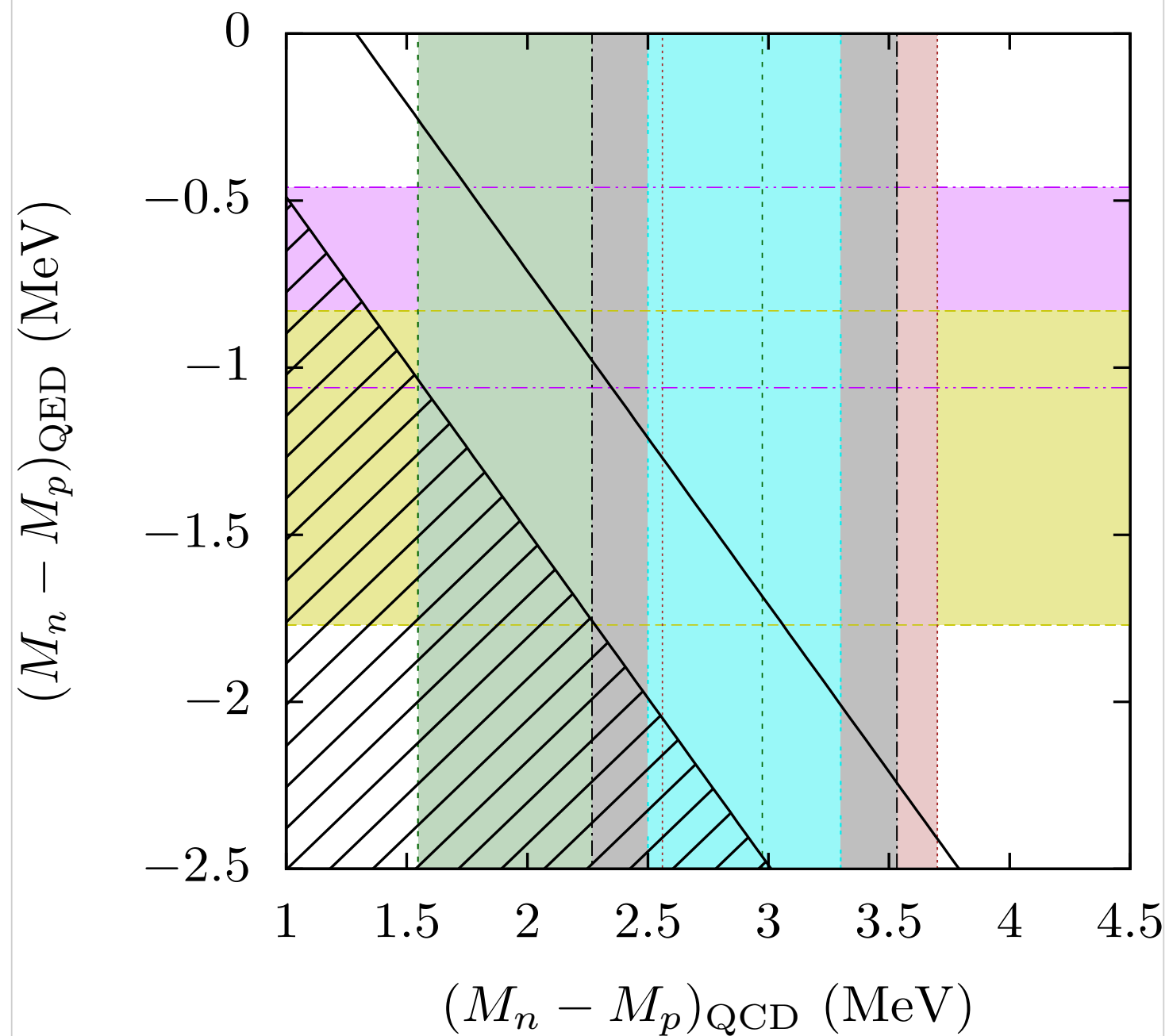


# Results for the nucleon mass splitting





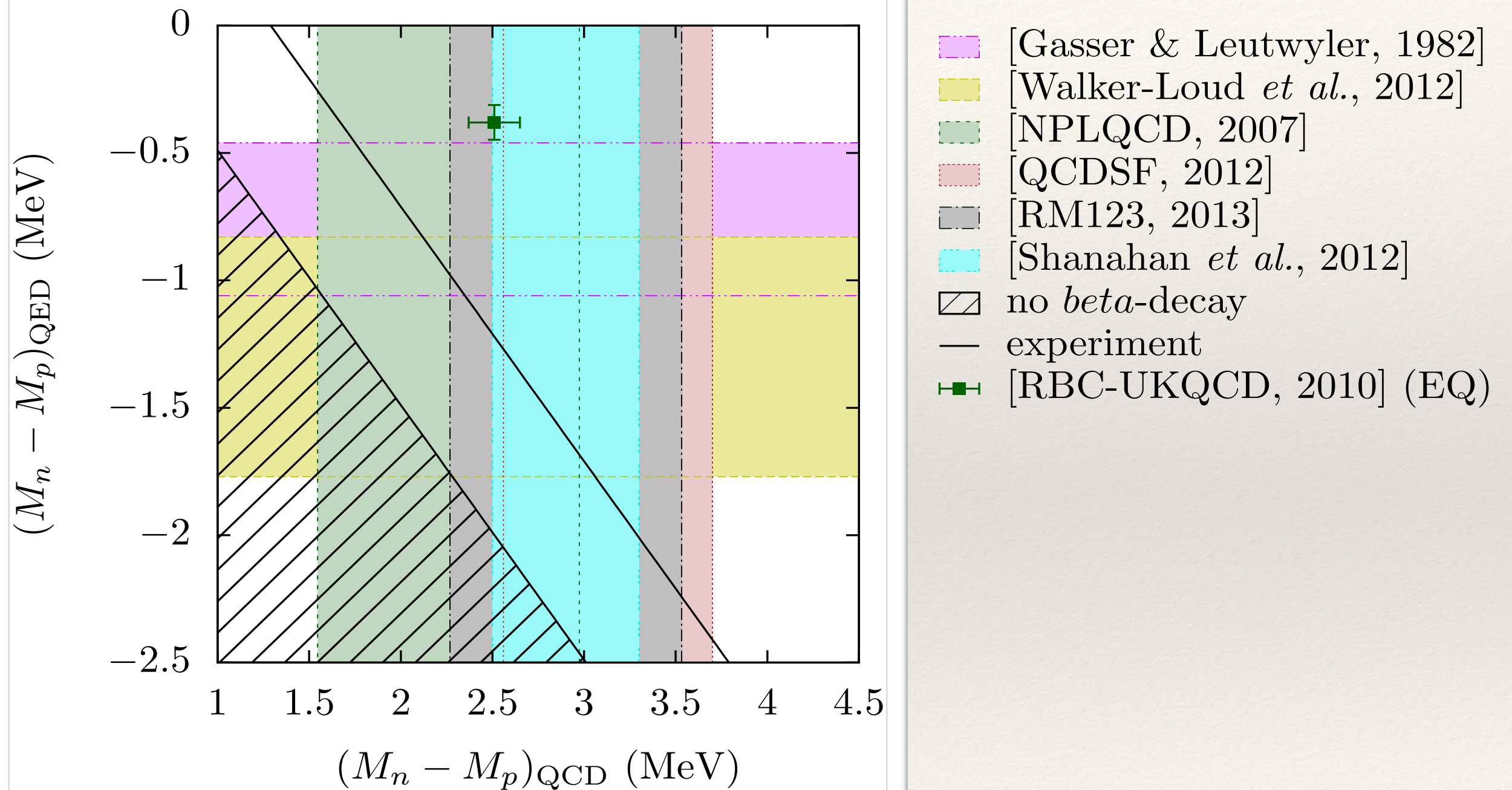
# Results for the nucleon mass splitting



- [Gasser & Leutwyler, 1982]
- [Walker-Loud *et al.*, 2012]
- [NPLQCD, 2007]
- [QCDSF, 2012]
- [RM123, 2013]
- [Shanahan *et al.*, 2012]
- no *beta*-decay
- experiment

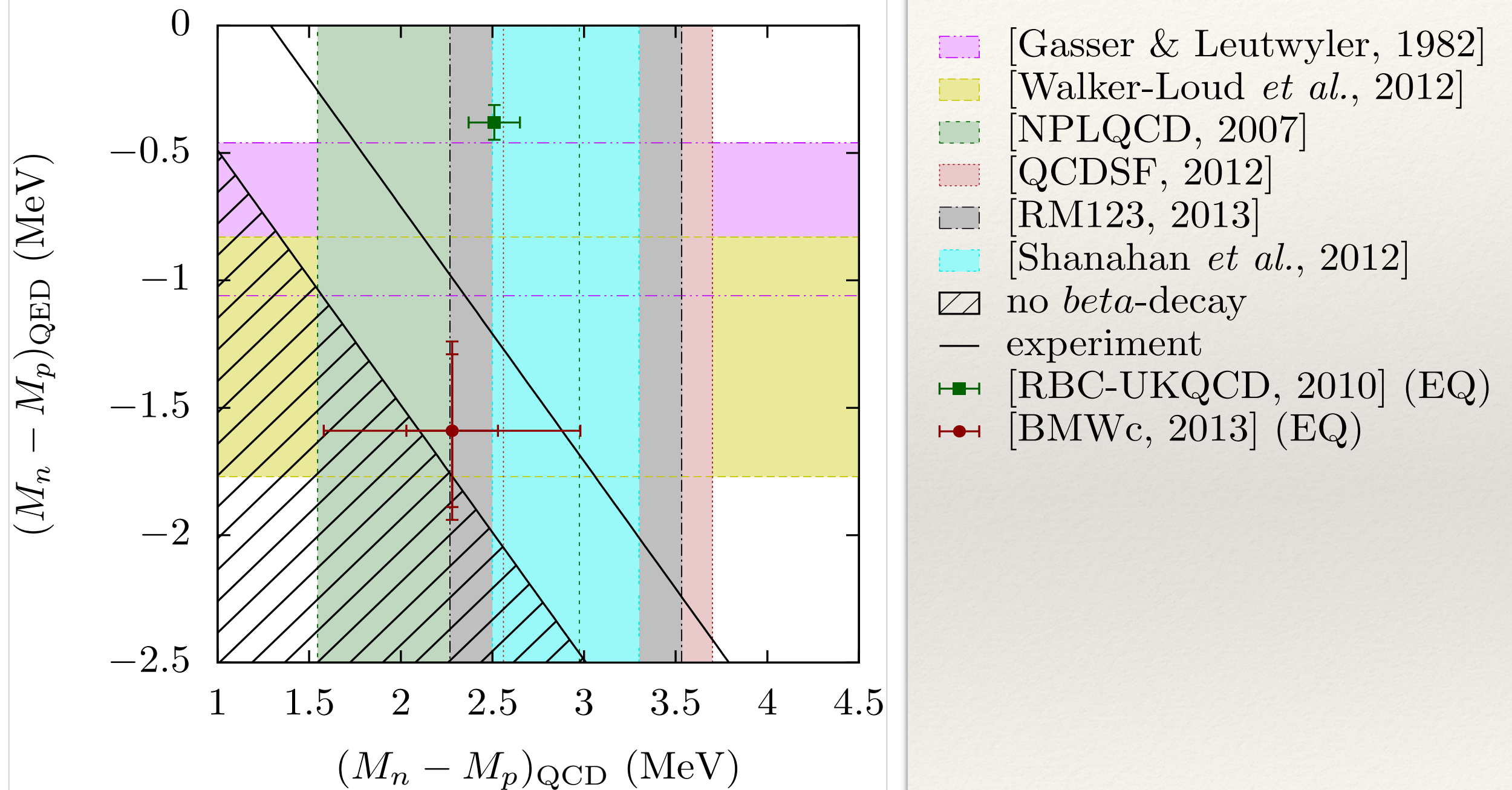


# Results for the nucleon mass splitting



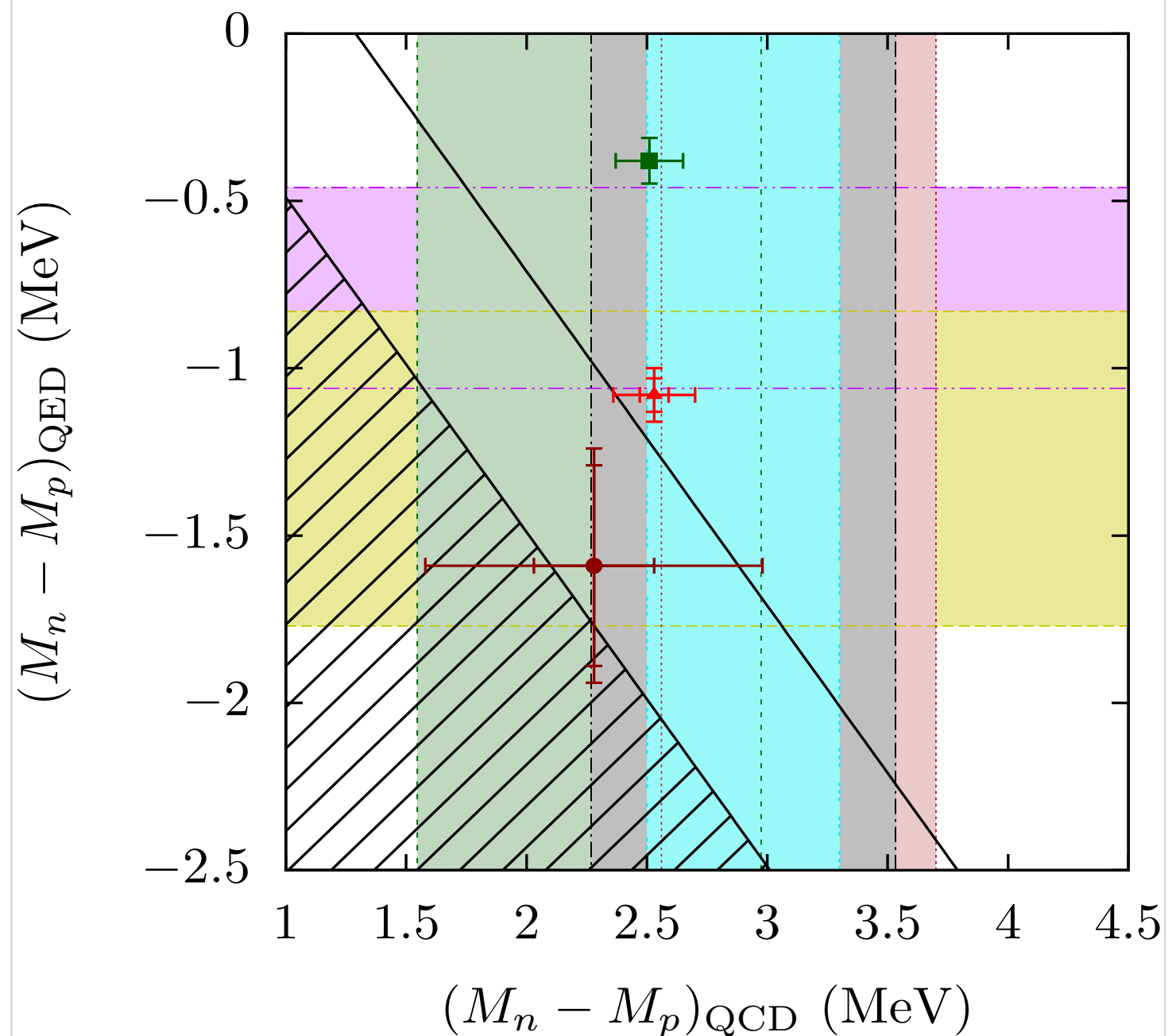


# Results for the nucleon mass splitting





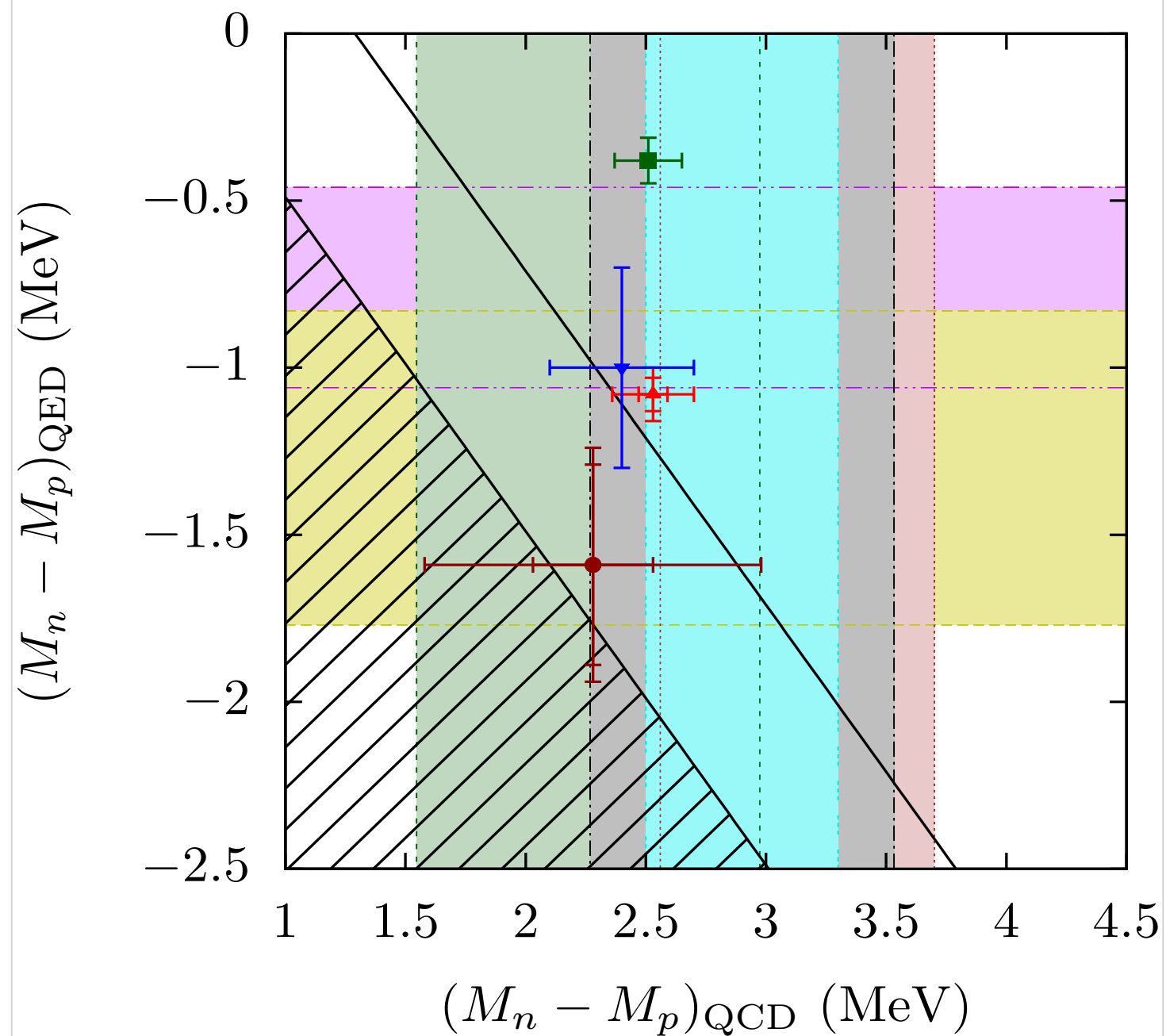
# Results for the nucleon mass splitting



- [Gasser & Leutwyler, 1982]
- [Walker-Loud *et al.*, 2012]
- [NPLQCD, 2007]
- [QCDSF, 2012]
- [RM123, 2013]
- [Shanahan *et al.*, 2012]
- no *beta*-decay
- experiment
- [RBC-UKQCD, 2010] (EQ)
- [BMWc, 2013] (EQ)
- [BMWc, 2015a]



# Results for the nucleon mass splitting



- [Gasser & Leutwyler, 1982]
- [Walker-Loud *et al.*, 2012]
- [NPLQCD, 2007]
- [QCDSF, 2012]
- [RM123, 2013]
- [Shanahan *et al.*, 2012]
- no *beta*-decay
- experiment
- [RBC-UKQCD, 2010] (EQ)
- [BMWc, 2013] (EQ)
- [BMWc, 2015a]
- [QCDSF, 2014]



# Summary & outlook



---

# Summary

---

- ❖ We now have a good understanding of QCD+QED on a finite lattice



---

# Summary

---

- ❖ We now have a good understanding of QCD+QED on a finite lattice
- ❖ Finite-size effects on masses are now **well controlled**



---

# Summary

---

- ❖ We now have a good understanding of QCD+QED on a finite lattice
- ❖ Finite-size effects on masses are now **well controlled**
- ❖ [BMWc, 2015a]: **full simulations of the low-energy SM** with a potential precision of  $O[(N_c m_b^2)^{-1}, \alpha^2] \sim 10^{-4}$



---

# Summary

---

- ❖ We now have a good understanding of QCD+QED on a finite lattice
- ❖ Finite-size effects on masses are now **well controlled**
- ❖ [BMWc, 2015a]: **full simulations of the low-energy SM** with a potential precision of  $O[(N_c m_b^2)^{-1}, \alpha^2] \sim 10^{-4}$
- ❖ The isospin splittings in the hadron spectrum are determined with a **high accuracy and full control of uncertainties**



---

# Summary

---

- ❖ We now have a good understanding of QCD+QED on a finite lattice
- ❖ Finite-size effects on masses are now **well controlled**
- ❖ [BMWc, 2015a]: **full simulations of the low-energy SM** with a potential precision of  $O[(N_c m_b^2)^{-1}, \alpha^2] \sim 10^{-4}$
- ❖ The isospin splittings in the hadron spectrum are determined with a **high accuracy and full control of uncertainties**
- ❖ The nucleon mass splitting is determined as a  $> 5\sigma$  effect



---

# Outlook

---

- ❖ Unquenched computations of the light quark masses and Dashen's theorem corrections



---

# Outlook

---

- ❖ Unquenched computations of the light quark masses and Dashen's theorem corrections
- ❖ QCD+QED decay constants are gauge variant and IR divergent. How to deal with that?  
First lattice attempt: **[plenary talk by V. Lubicz]**



---

# Outlook

---

- ❖ Unquenched computations of the light quark masses and Dashen's theorem corrections
- ❖ QCD+QED decay constants are gauge variant and IR divergent. How to deal with that?  
First lattice attempt: **[plenary talk by V. Lubicz]**
- ❖ Compute corrections to matrix elements  
( $K_{\ell 3}$ ,  $K \rightarrow \pi\pi, \dots$ )



---

# Outlook

---

- ❖ Unquenched computations of the light quark masses and Dashen's theorem corrections
- ❖ QCD+QED decay constants are gauge variant and IR divergent. How to deal with that?  
First lattice attempt: **[plenary talk by V. Lubicz]**
- ❖ Compute corrections to matrix elements  
( $K_{\ell 3}, K \rightarrow \pi\pi, \dots$ )
- ❖ QCD+QED to compute hadronic corrections to anomalous magnetic moments.





---

Thank you!

---



Backup

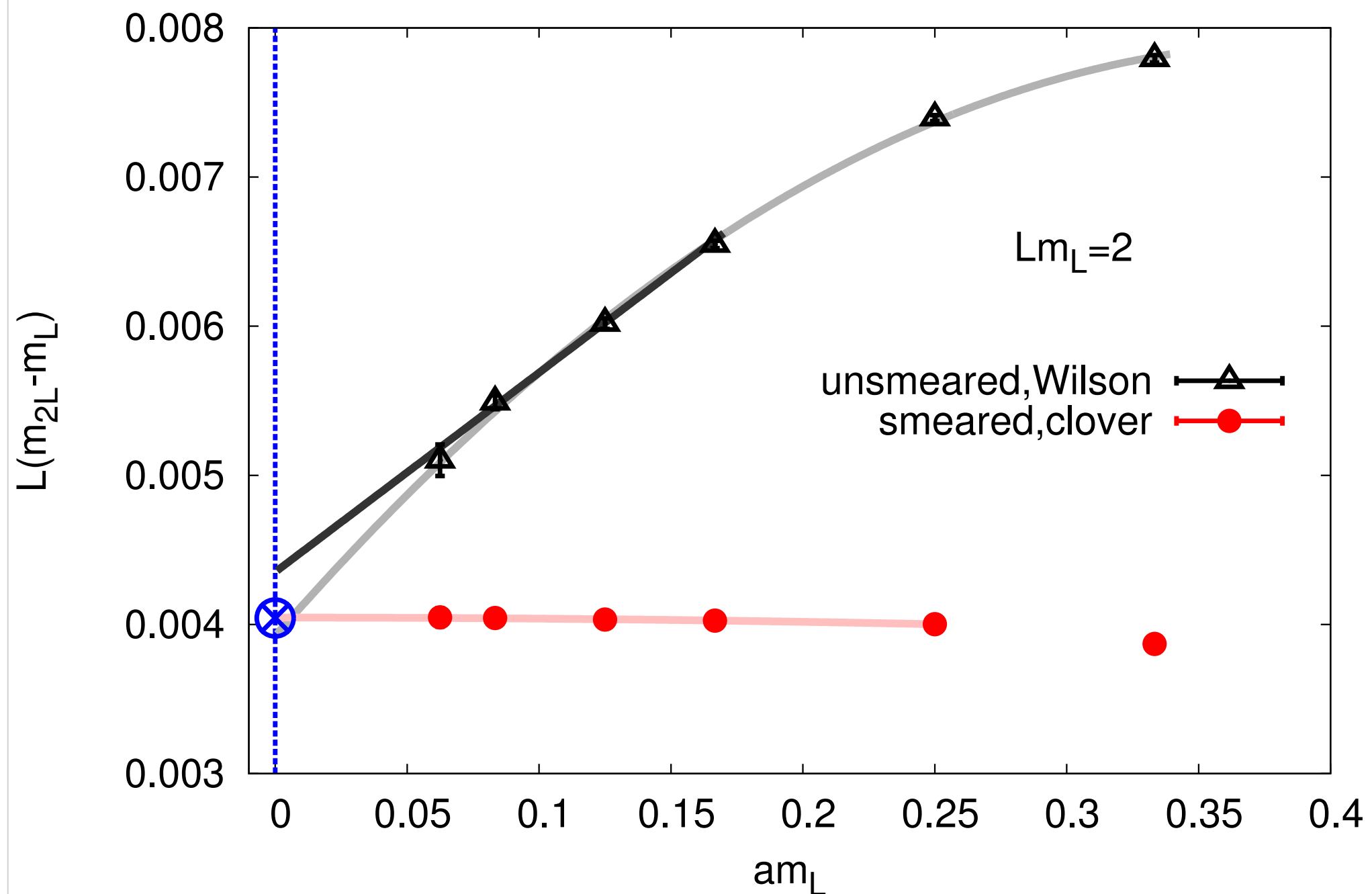


# Full QCD + QED projects

|                     | RBC-UKQCD   | PACS-CS     | QCDSF-UKQCD                | BMWc        |
|---------------------|-------------|-------------|----------------------------|-------------|
| arXiv               | 1006.1311   | 1205.2961   | 1311.4554<br>and Lat. 2014 | 1406.4088   |
| fermions            | DWF         | clover      | clover                     | clover      |
| $N_f$               | 2+1         | 1+1+1       | 1+1+1                      | 1+1+1+1     |
| method              | reweighting | reweighting | RHMC                       | RHMC        |
| $\min(M_\pi)$ (MeV) | 420         | 135         | 250                        | 195         |
| $a$ (fm)            | 0.11        | 0.09        | 0.08                       | 0.06 — 0.10 |
| $\#a$               | 1           | 1           | 1                          | 4           |
| $L$ (fm)            | 1.8         | 2.9         | 1.9 — 2.6                  | 2.1 — 8.3   |
| $\#L$               | 1           | 1           | 2                          | 11          |

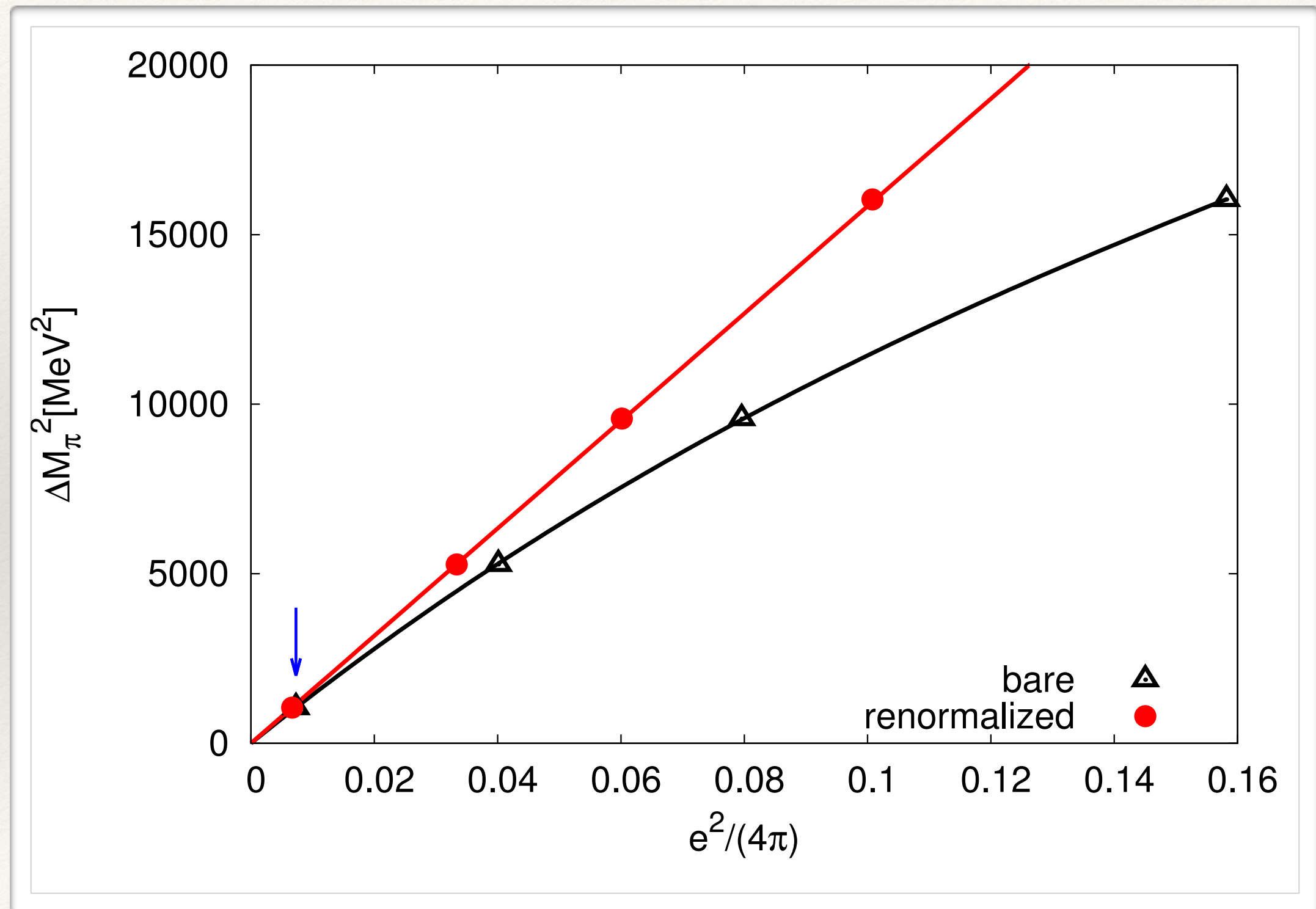


# [BMWc, 2015a]: QED simulations





# [BMWc, 2015a]: charge renormalisation





# [BMWc, 2015a]: charm discretisation effects

