

Meson resonances on the lattice

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Hadron Spectrum Collaboration

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MESON SPECTRUM

PRL 103 262001 (2009) $I = 1$
PRD 82 034508 (2010) $I = 1, K^*$
PRD 83 111502 (2011) $I = 0$
JHEP 07 126 (2011) $c\bar{c}$
PRD 88 094505 (2013) $I = 0$
JHEP 05 021 (2013) D, D_s

BARYON SPECTRUM

PRD 84 074508 (2011) $(N, \Delta)^*$
PRD 85 054016 (2012) $(N, \Delta)_{\text{hyb}}$
PRD 87 054506 (2013) $(N \dots \Xi)^*$
PRD 90 074504 (2014) Ω_{ccc}^*
arXiv:1502.01845 Ξ_{cc}^*

HADRON SCATTERING

PRD 83 071504 (2011) $\pi\pi I = 2$
PRD 86 034031 (2012) $\pi\pi I = 2$
PRD 87 034505 (2013) $\pi\pi I = 1, \rho$
PRL 113 182001 (2014) $\pi K, \eta K$
PRD 91 054008 (2015) $\pi K, \eta K$

“TECHNOLOGY”

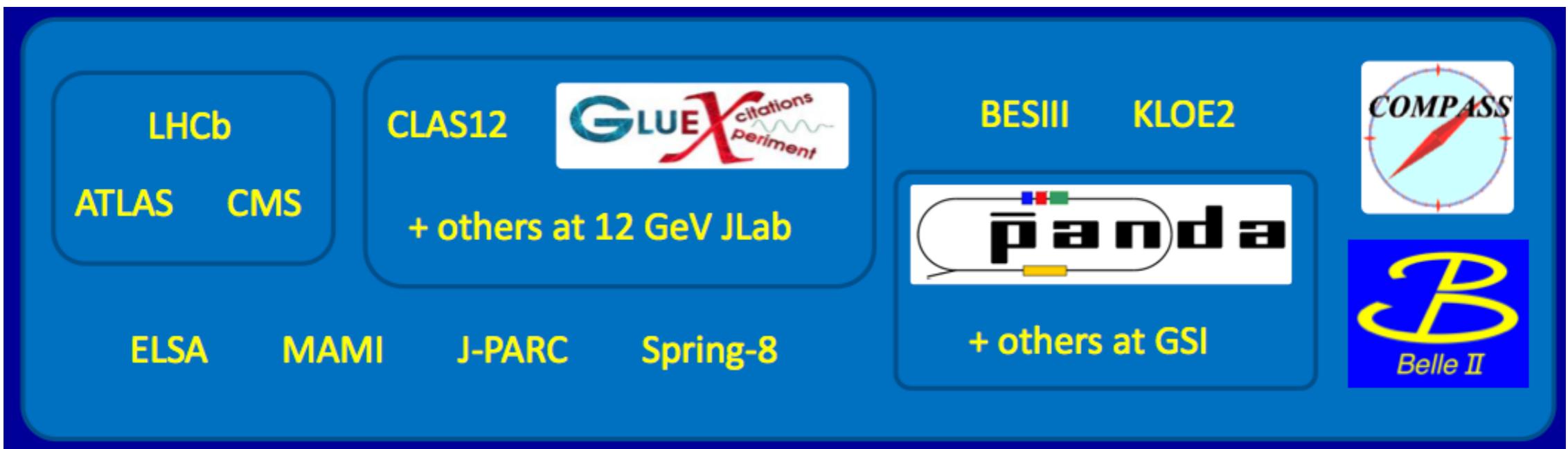
PRD 79 034502 (2009) lattices
PRD 80 054506 (2009) distillation
PRD 85 014507 (2012) $\vec{p} > 0$

MATRIX ELEMENTS

arXiv:1501.07457 $M' \rightarrow \gamma M$
PRD 90 014511 (2014) f_{π^*}

Hadron spectroscopy

- Determination of hadron spectrum of Quantum Chromodynamics (QCD) a central goal in NP
- Several experiments worldwide

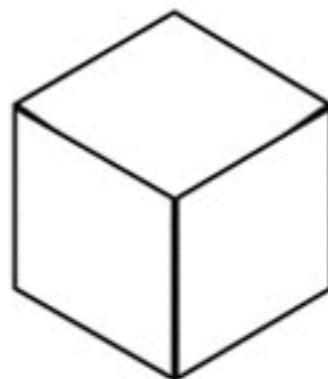


Finite volume QCD & the hadron spectrum

- Compute correlation functions as an average over field configurations

e.g. $\int \mathcal{D}\psi \mathcal{D}\bar{\psi} \mathcal{D}A_\mu \bar{\psi} \Gamma \psi(t) \bar{\psi} \Gamma \psi(0) e^{-\int d^4x \mathcal{L}_{\text{QCD}}(\psi, \bar{\psi}, A_\mu)}$

‘sum’ ‘field correlation’ ‘probability weight’



Field integration within a finite, but continuous, hypercube
Need some kind of ultraviolet regulator....

- Spectrum from two-point correlation functions

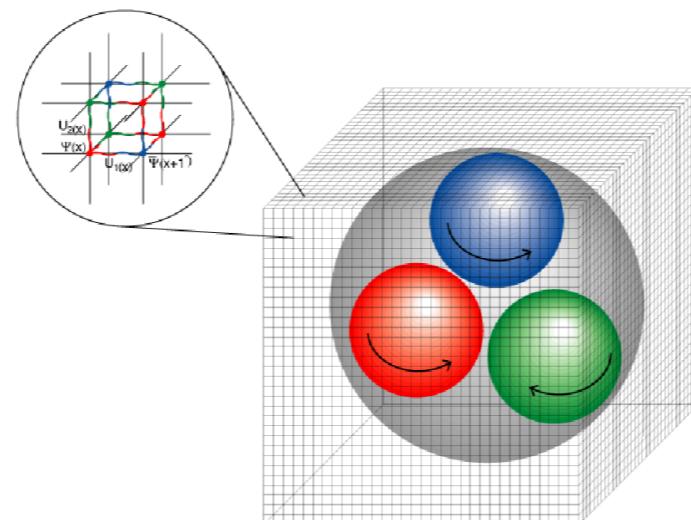
$$\begin{aligned} C(t) &= \langle 0 | \mathcal{O}(t) \mathcal{O}^\dagger(0) | 0 \rangle \\ &= \sum_{\mathbf{n}} e^{-E(\mathbf{n})t} \langle 0 | \mathcal{O}(0) | \mathbf{n} \rangle \langle \mathbf{n} | \mathcal{O}^\dagger(0) | 0 \rangle \end{aligned}$$

Lattice QCD & the hadron spectrum

- Compute correlation functions as a Monte Carlo average over field configurations

e.g. $\int \mathcal{D}\psi \mathcal{D}\bar{\psi} \mathcal{D}A_\mu \bar{\psi} \Gamma \psi(t) \bar{\psi} \Gamma \psi(0) e^{-\int d^4x \mathcal{L}_{\text{QCD}}(\psi, \bar{\psi}, A_\mu)}$

‘sum’ ‘field correlation’ ‘probability weight’



Discretize the action over sites

Serves as an ultraviolet regulator

- Spectrum from two-point correlation functions

$$\begin{aligned} C(t) &= \langle 0 | \mathcal{O}(t) \mathcal{O}^\dagger(0) | 0 \rangle \\ &= \sum_{\mathbf{n}} e^{-E(\mathbf{n})t} \langle 0 | \mathcal{O}(0) | \mathbf{n} \rangle \langle \mathbf{n} | \mathcal{O}^\dagger(0) | 0 \rangle \end{aligned}$$

Excited states from correlators

- how to get at excited QCD eigenstates ?

- optimal operator for state $|\mathfrak{n}\rangle$: $\Omega_{\mathfrak{n}}^\dagger \sim \sum_i v_i^{(\mathfrak{n})} \mathcal{O}_i^\dagger$

for a basis of
meson operators $\{\mathcal{O}_i\}$

- can be obtained (in a variational sense) from the matrix of correlators

$$C_{ij}(t) = \langle 0 | \mathcal{O}_i(t) \mathcal{O}_j^\dagger(0) | 0 \rangle$$

- by solving a generalized eigenvalue problem

$$C(t)v^{(\mathfrak{n})} = C(t_0)v^{(\mathfrak{n})} \lambda_{\mathfrak{n}}(t)$$

‘diagonalize the
correlation matrix’

eigenvalues

$$\lambda_{\mathfrak{n}}(t) \sim e^{-E_{\mathfrak{n}}(t-t_0)}$$

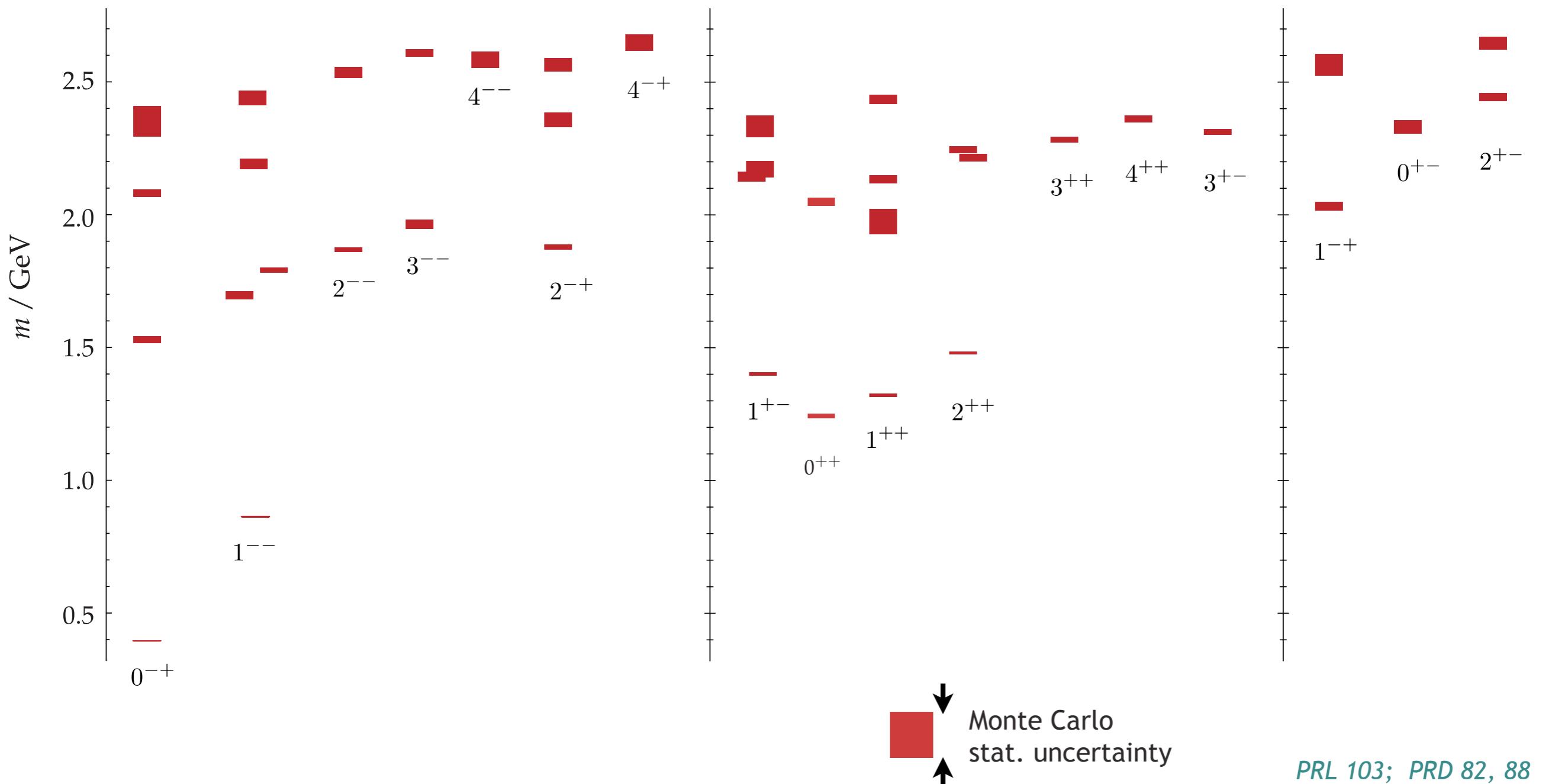
- a large basis can be constructed using covariant derivatives :

$$\mathcal{O} \sim \bar{\psi} \Gamma \overleftrightarrow{D} \dots \overleftrightarrow{D} \psi$$

Glimpse of meson spectrum from lattice QCD

- Appears to be some $q\bar{q}$ -like near-degeneracy patterns - isovectors

$m_\pi \sim 391$ MeV

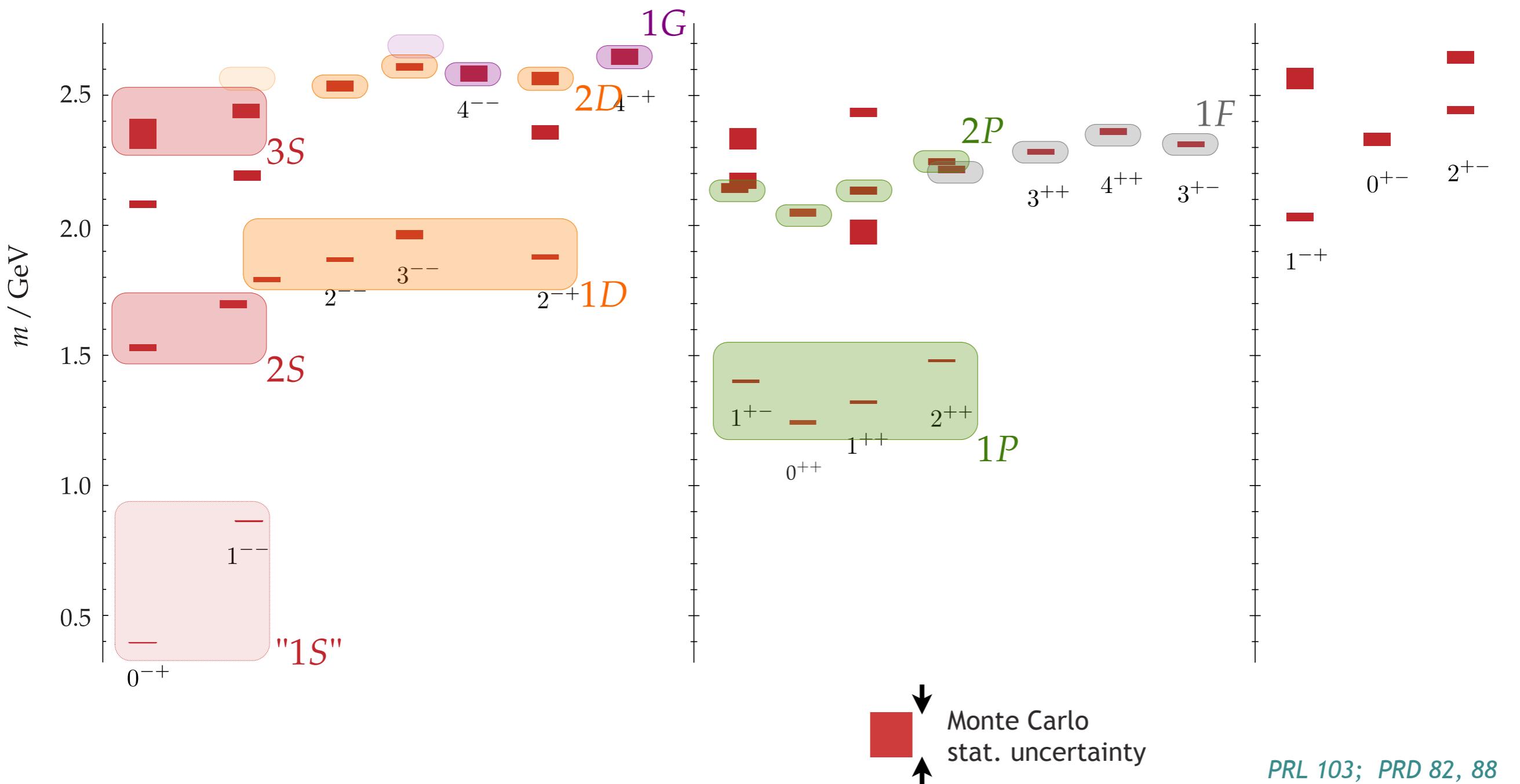


PRL 103; PRD 82, 88

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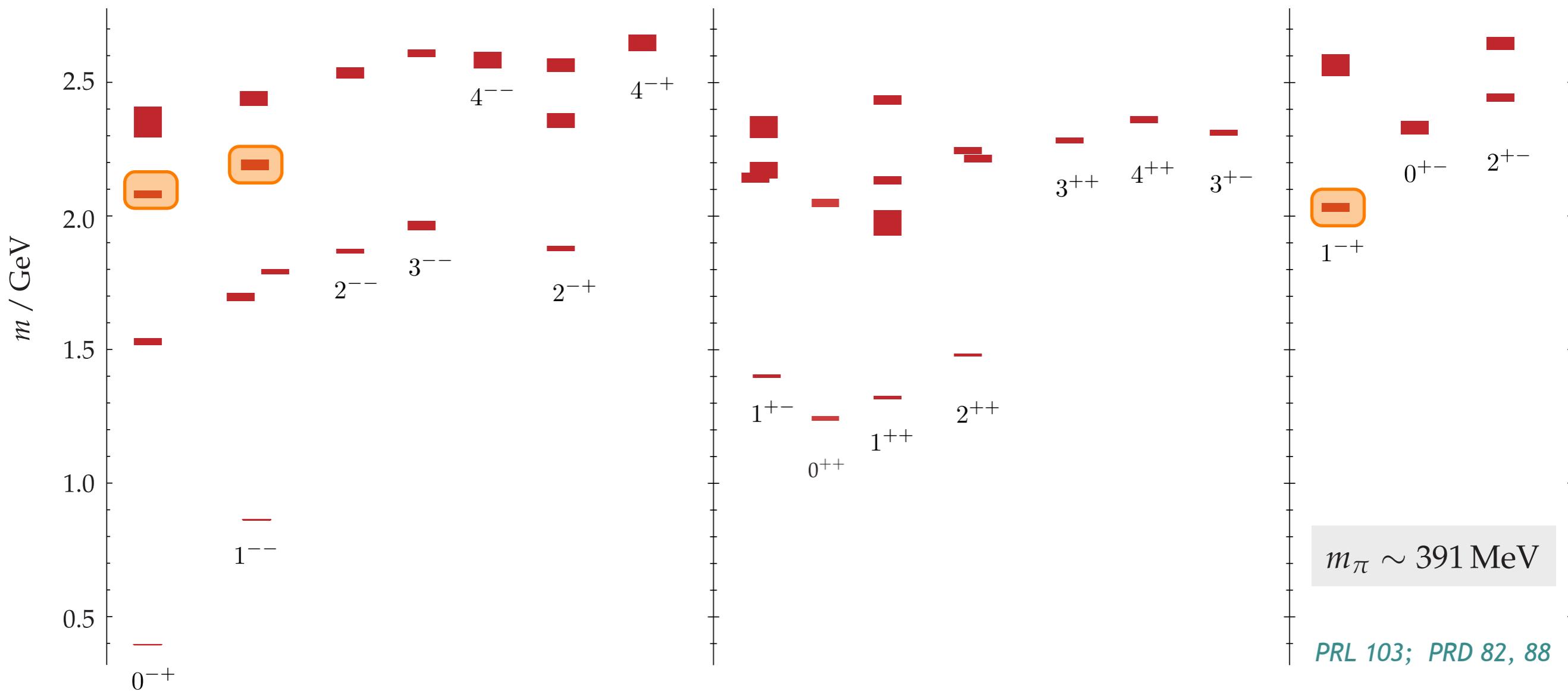


PRL 103; PRD 82, 88

Glimpse of meson spectrum from lattice QCD

- ‘super’-multiplet of **hybrid mesons** roughly 1.2 GeV above the ρ

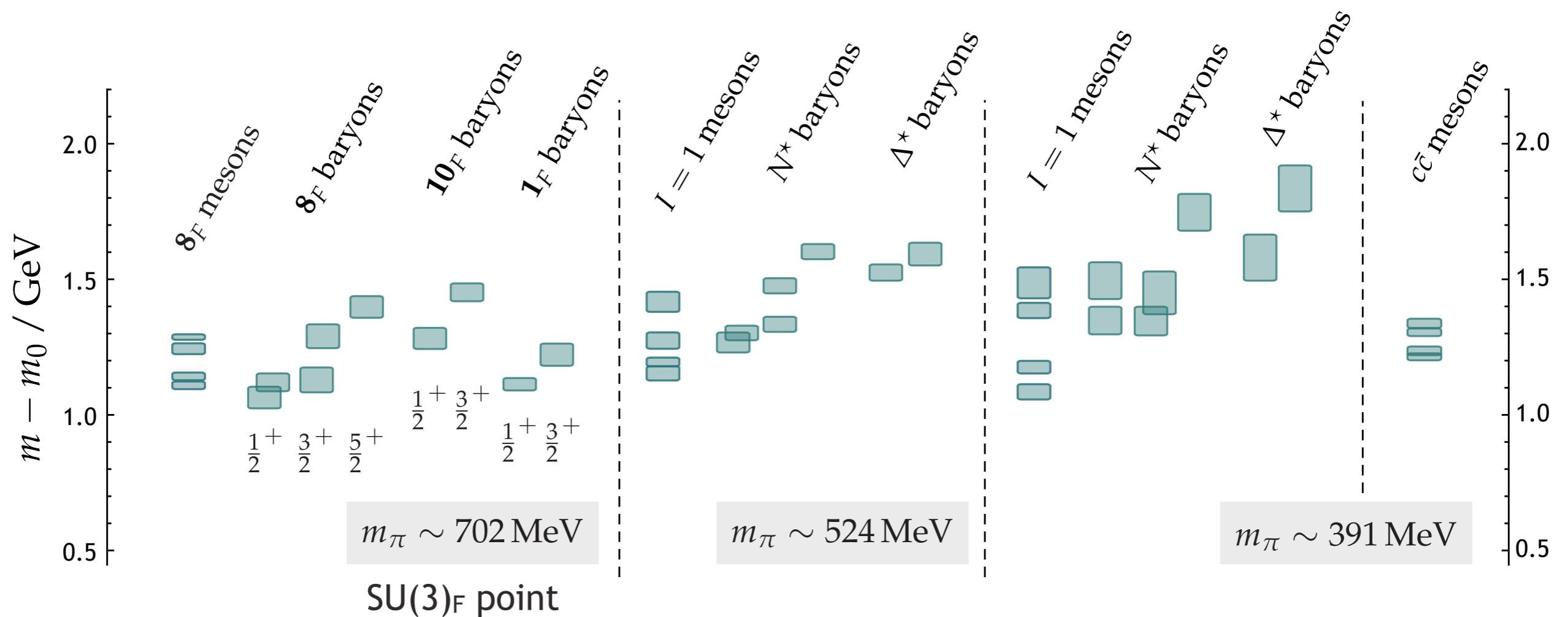
$(0, 1, 2)^{-+}, 1^{--}$



- these states have a dominant overlap onto $\bar{\psi}\Gamma[D, D]\psi \sim [q\bar{q}]_{8_c} \otimes B_{8_c}$

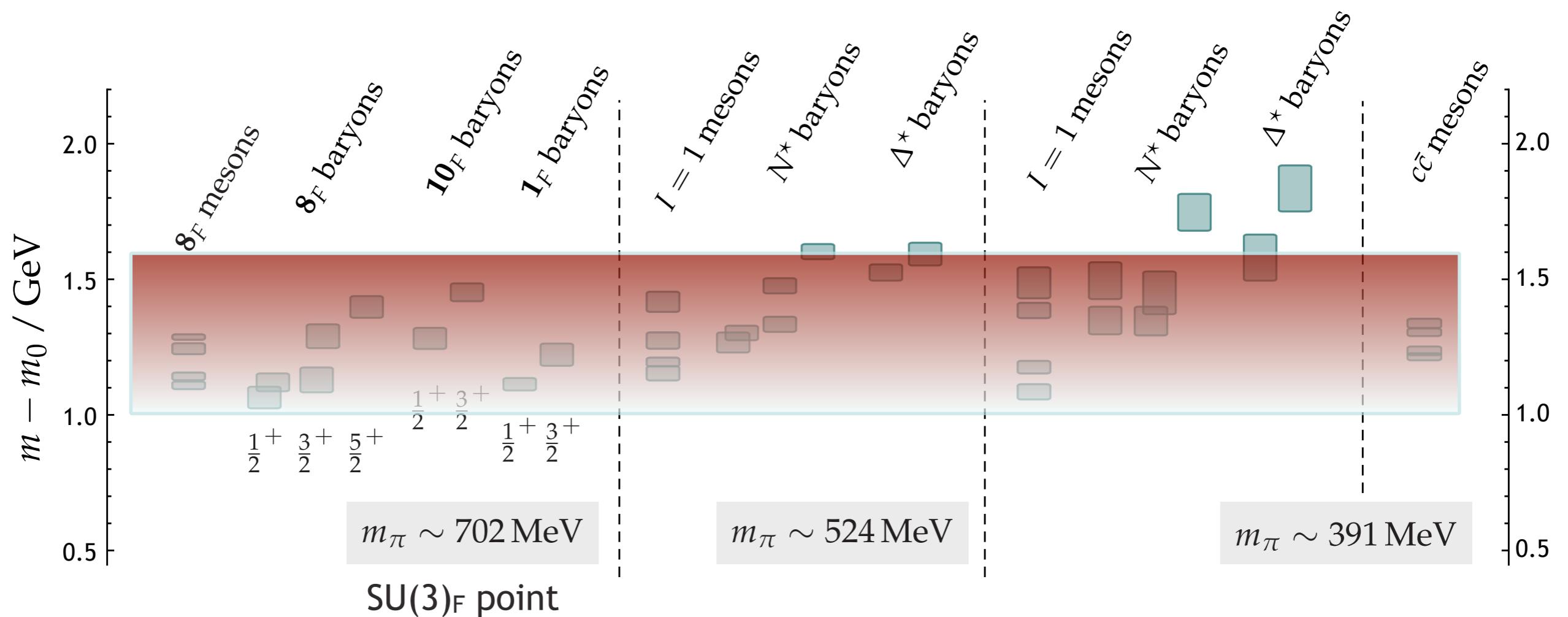
Chromo-magnetic excitation

- Subtract the ‘quark mass’ contribution



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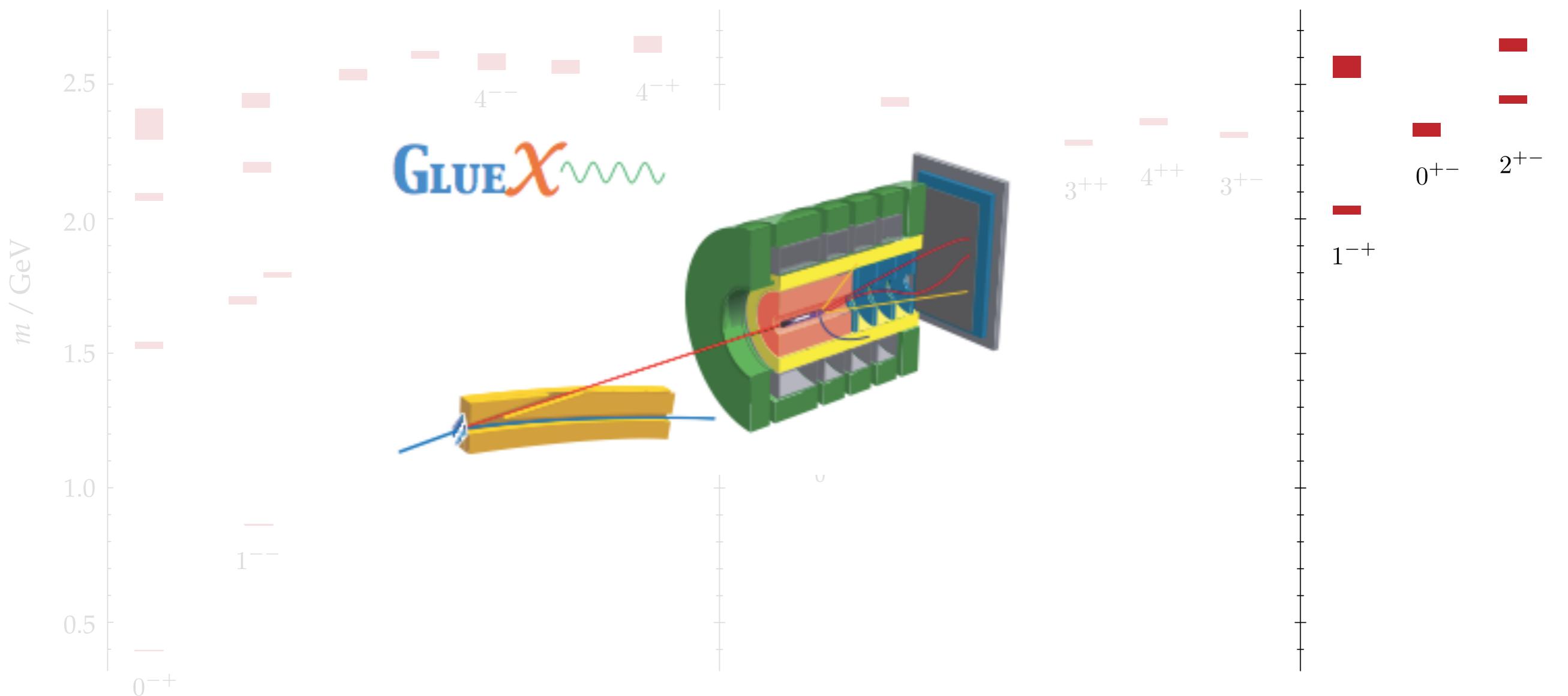
– *Common energy scale of gluonic excitation*

$\sim 1.3 \text{ GeV}$

Glimpse of meson spectrum from lattice QCD

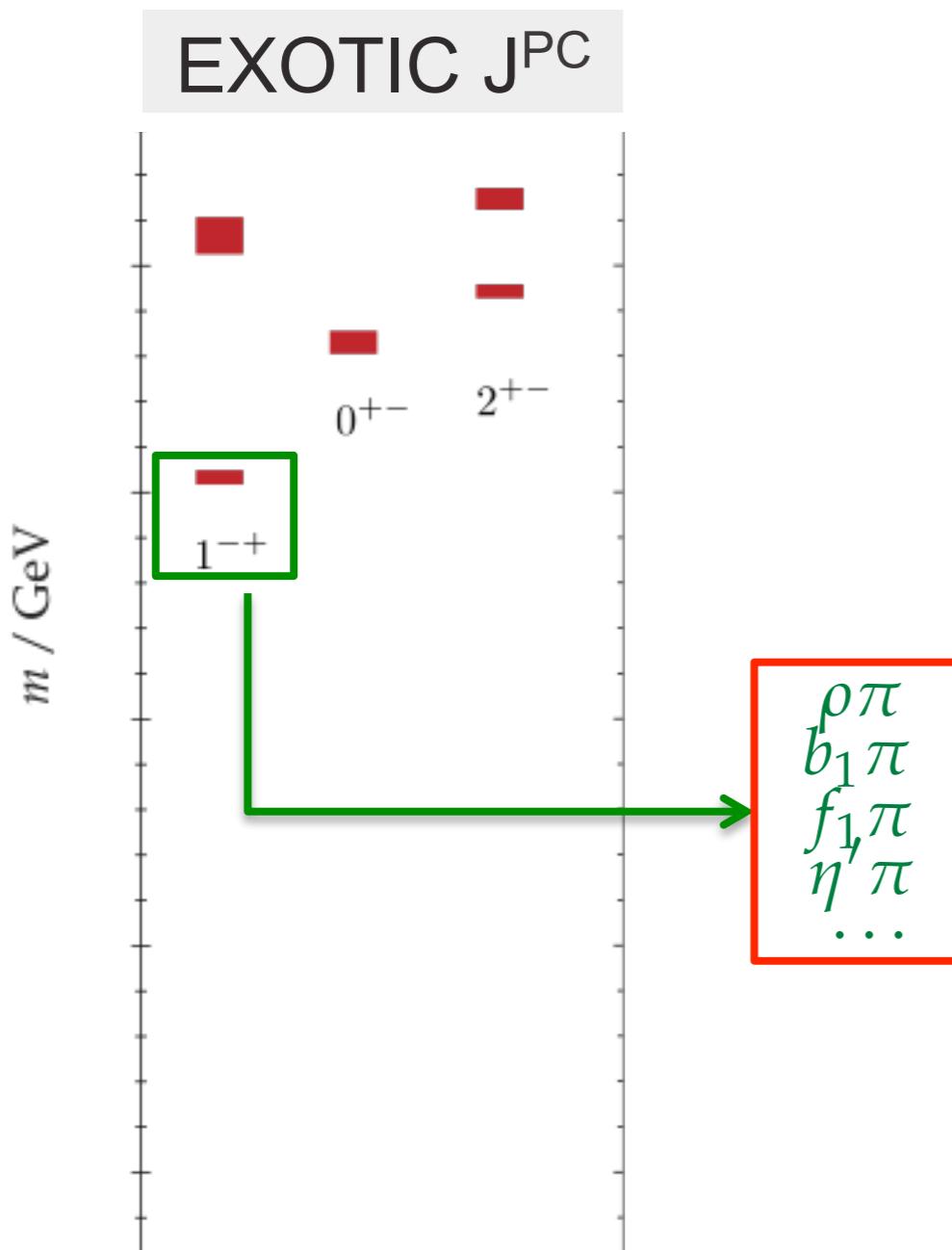
Multiple exotic mesons within range of GlueX

EXOTIC MESONS



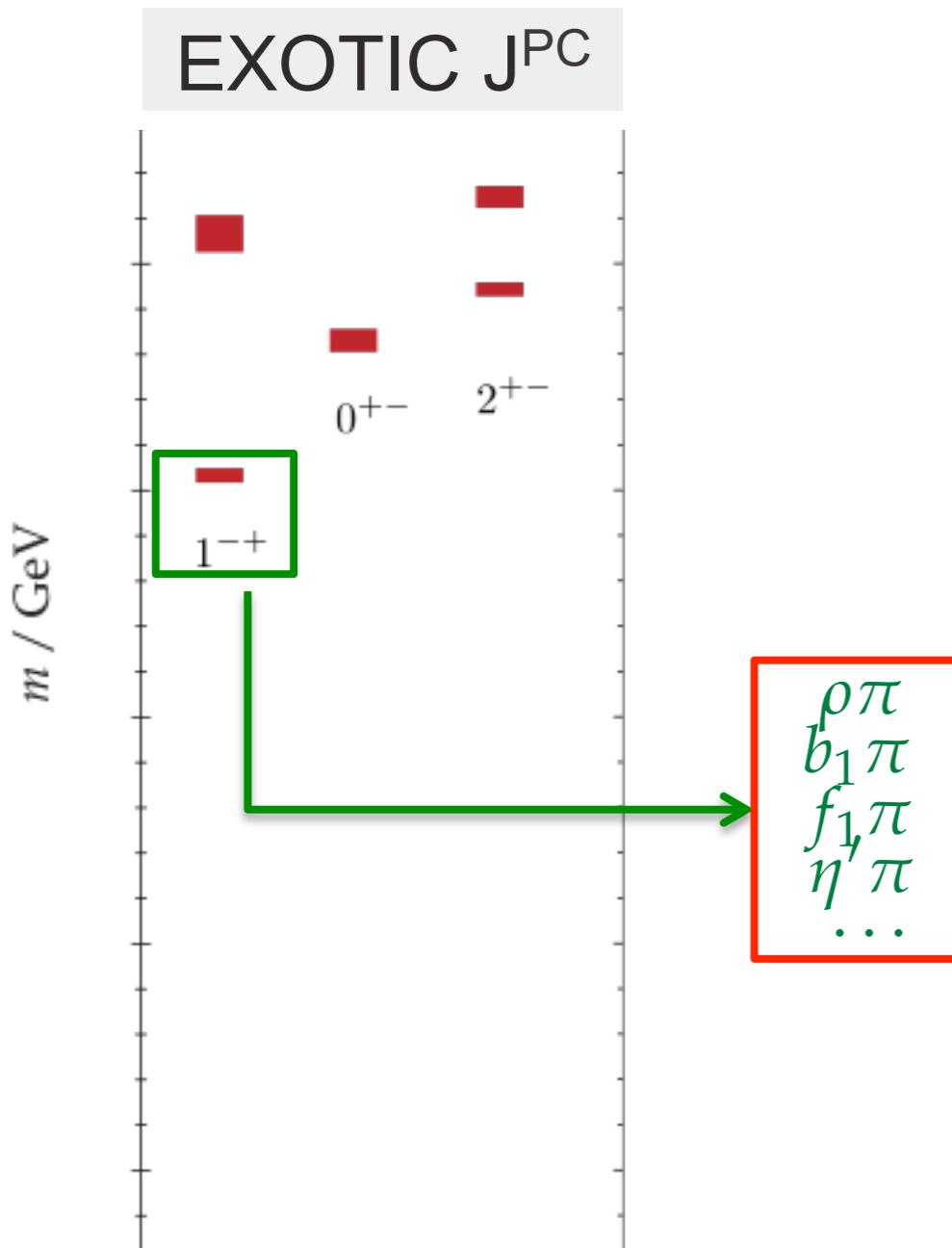
PRL 103; PRD 82, 88

Major objective - exotic meson decays



- LQCD suggests existence of exotic mesons
- Expt. measurement in many decay channels
- Present LQCD calculations missing this info

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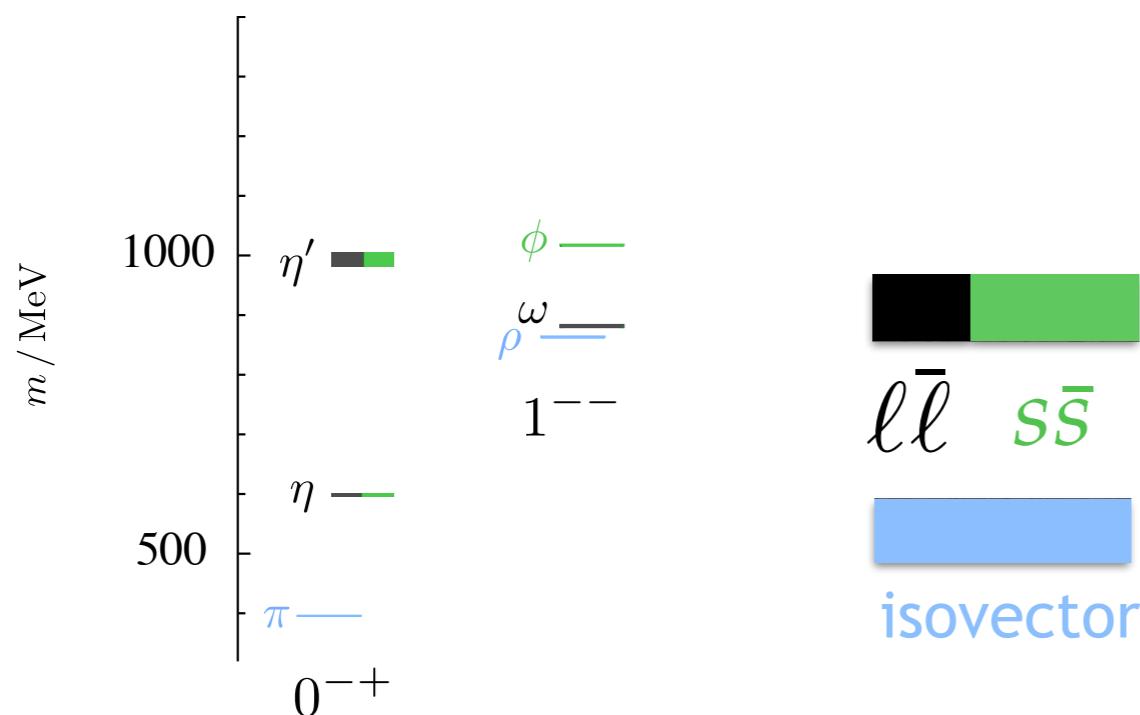


- LQCD suggests existence of exotic mesons
- Expt. measurement in many decay channels
- Present LQCD calculations missing this info
- Objective is to compute them ahead of expt.
 - » Guide expt. analysis

JLab expt. beam has started

Can study many other channels

- Many come as a pre-requisite before tackling exotics
- In particular, low-energy sector of QCD



What pion mass?

- Getting to the physical pion mass **not the most pressing concern here**
- Need to establish feasibility of techniques for resonances

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 - $a_2(1320)$ $M > 9m_\pi$

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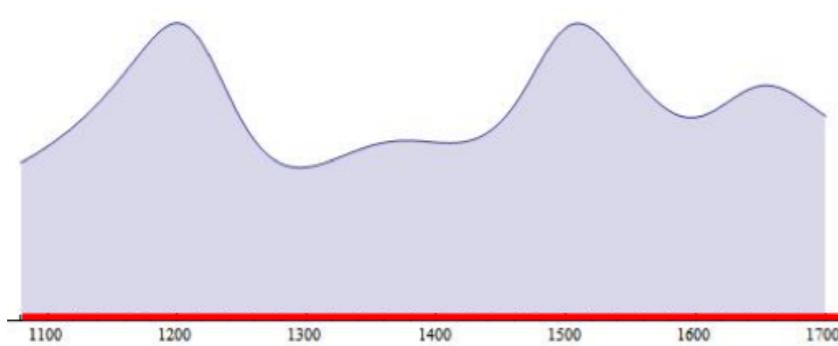
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- Development of three-body formalism required

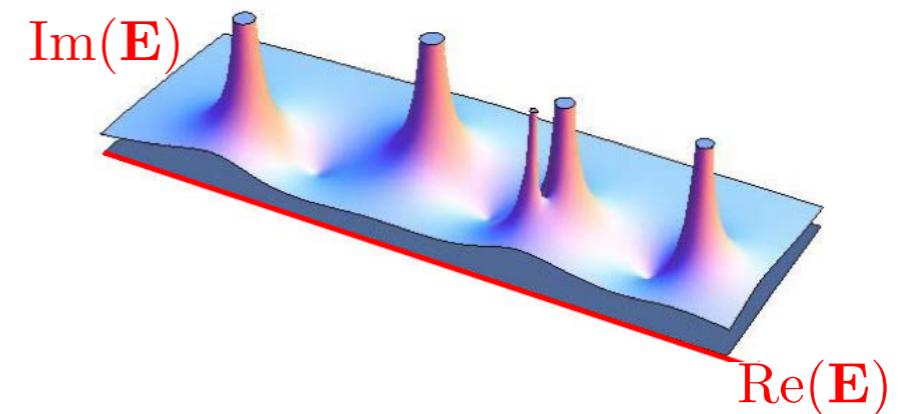
*E.G., HANSEN & SHARPE,
BEANE/BRICENO/RUSETSKY/SAVAGE,... - PROGRESS*

Resonances

- Most hadrons are resonances
 - E.g., $\pi N \pi N$



E (MeV)



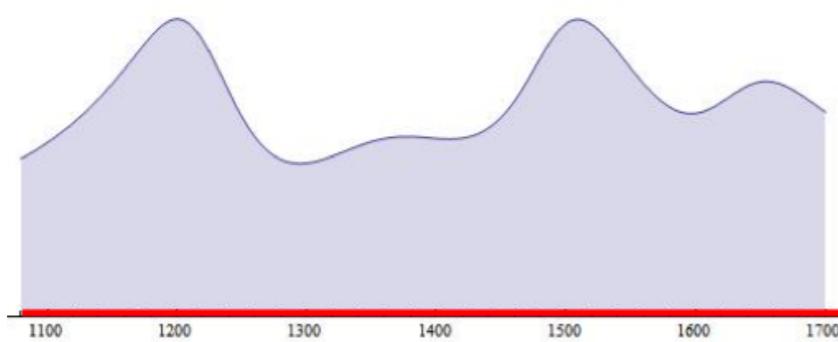
- Formally defined as a pole in a partial-wave scattering amplitude

$$t_l(s) \sim \frac{R}{s_0 - s} + \dots$$

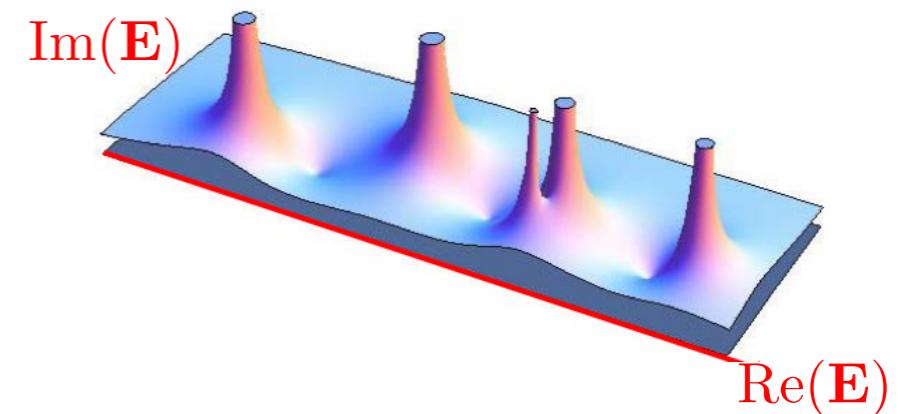
- Different channels should have same pole location
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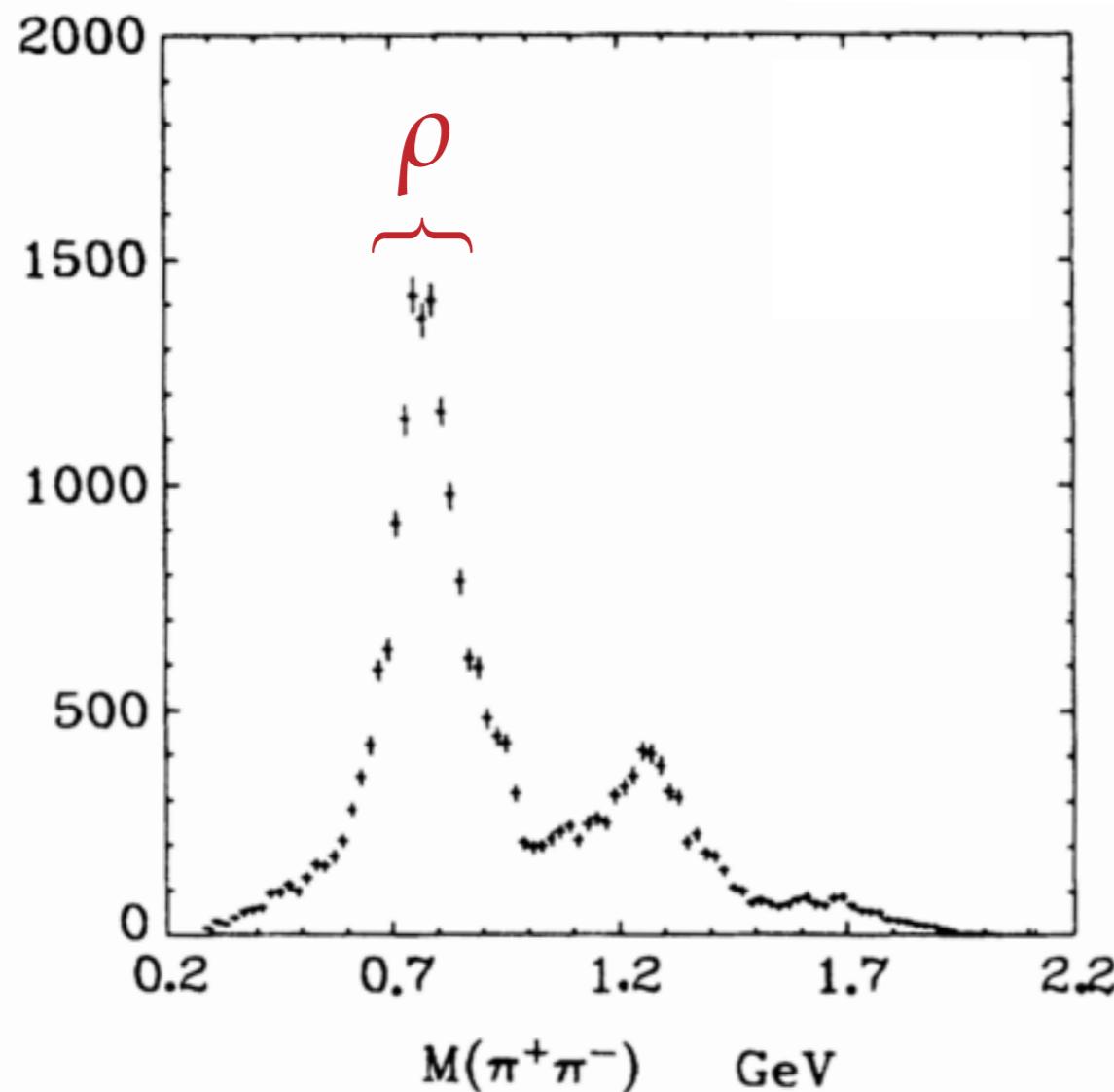
Re(E)

- Formally defined as a pole in a partial-wave scattering amplitude

$$t_l(s) \sim \frac{R}{s_0 - s} + \dots$$

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-
- Can we predict hadron properties from first principles?

Isospin=1 $\pi\pi$ P-wave



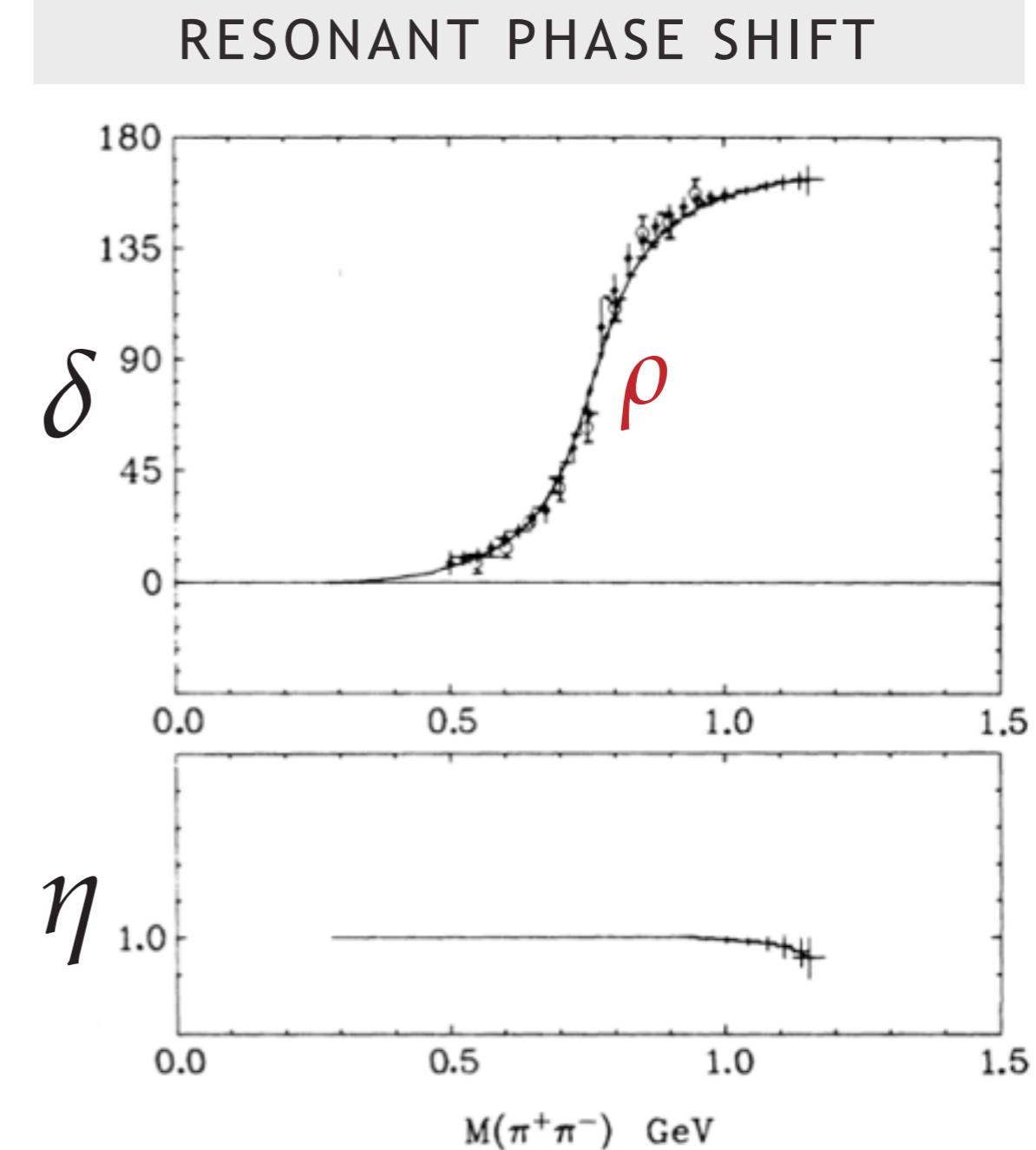
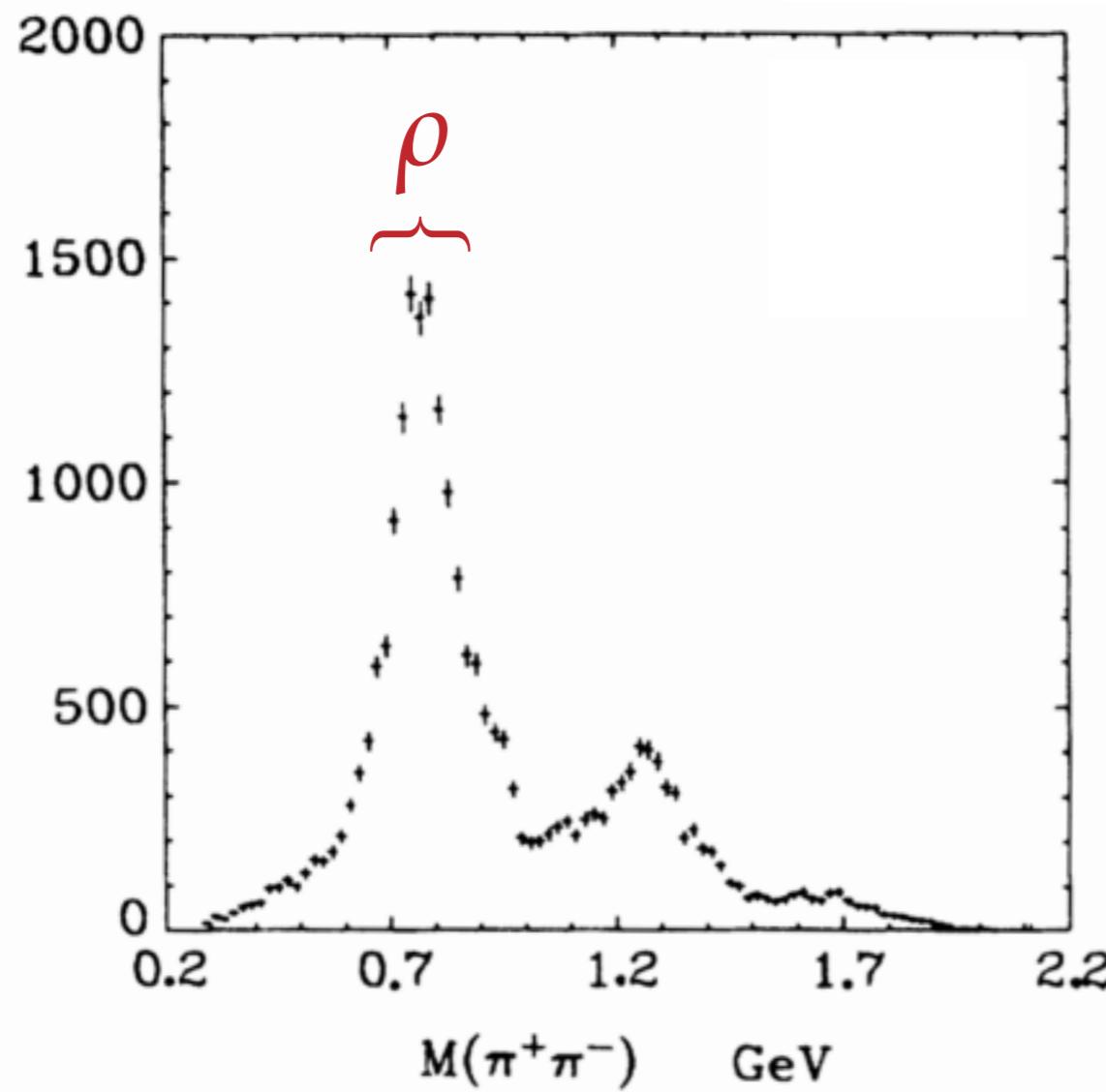
expand angular dependence
in *partial waves*

PARTIAL WAVE AMPLITUDE

$$f_\ell = \frac{1}{2i} (\eta_\ell e^{2i\delta_\ell} - 1)$$

$\eta = 1$ elastic
 $\eta \leq 1$ inelastic

Isospin=1 $\pi\pi$ P-wave



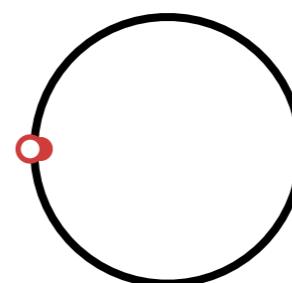
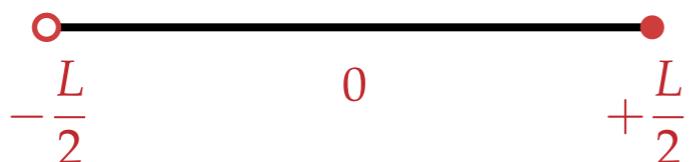
Finite-volume

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 - there isn't one !

Finite-volume

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 - there isn't one !
 - in a finite-volume the spectrum is discrete

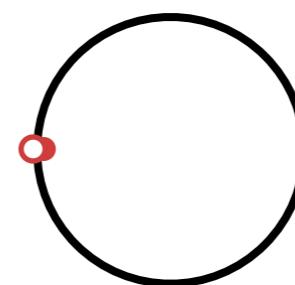
one-dim :



Finite-volume

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one-dim :



e.g. a free particle
 $\psi(x) \sim e^{ipx}$

» periodic boundary condition

$$\begin{aligned}\psi(x) &= \psi(x + L) \\ e^{ipx} &= e^{ip(x+L)} \\ e^{ipL} &= 1\end{aligned}$$

$$p = \frac{2\pi}{L}n$$

discrete
energy
spectrum

Interacting particles in a finite-volume

- Two identical bosons **interacting** through a finite-range potential

$$\psi(z) \sim \cos [p|z| + \delta(p)] \quad \text{outside the range of the potential, } |z| > R$$

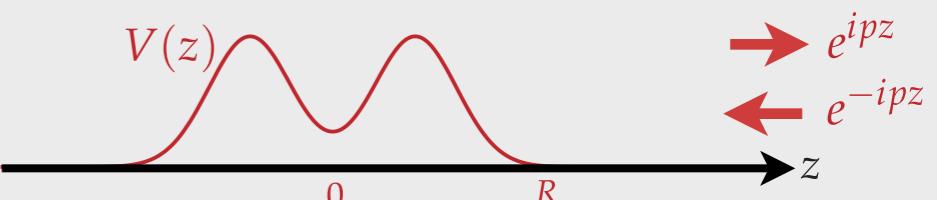
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outside the range
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$$\psi(z > 0) \sim e^{-ipz} + e^{2i\delta(p)} \cdot e^{ipz}$$



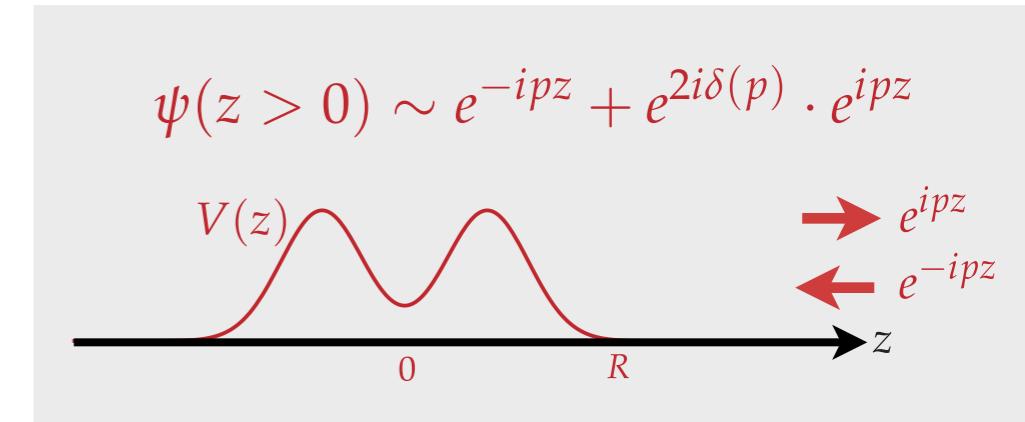
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$$\left. \begin{array}{l} \psi\left(-\frac{L}{2}\right) = \psi\left(\frac{L}{2}\right) \\ \frac{d\psi}{dz}\left(-\frac{L}{2}\right) = \frac{d\psi}{dz}\left(\frac{L}{2}\right) \end{array} \right\} \Rightarrow 0 = \sin \left[\frac{pL}{2} + \delta(p) \right]$$
$$\frac{pL}{2} + \delta(p) = n\pi$$



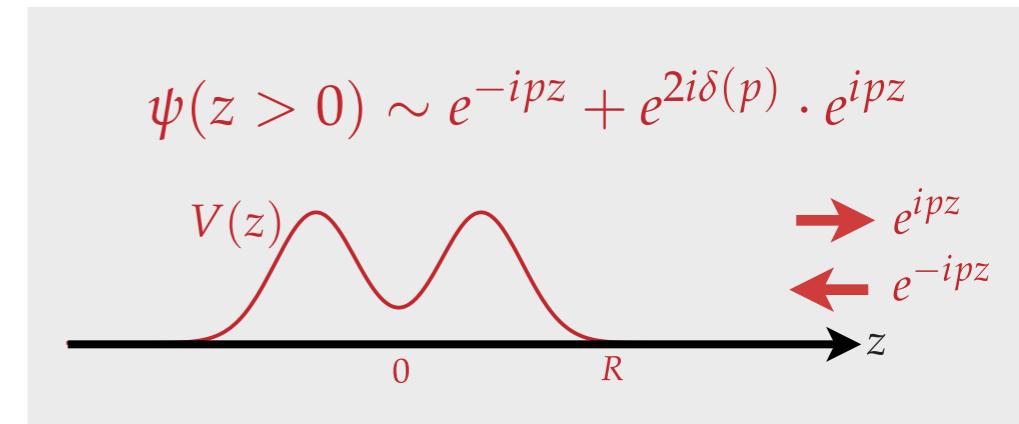
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$$p = \frac{2\pi}{L}n - \frac{2}{L}\delta(p)$$

discrete
energy
spectrum

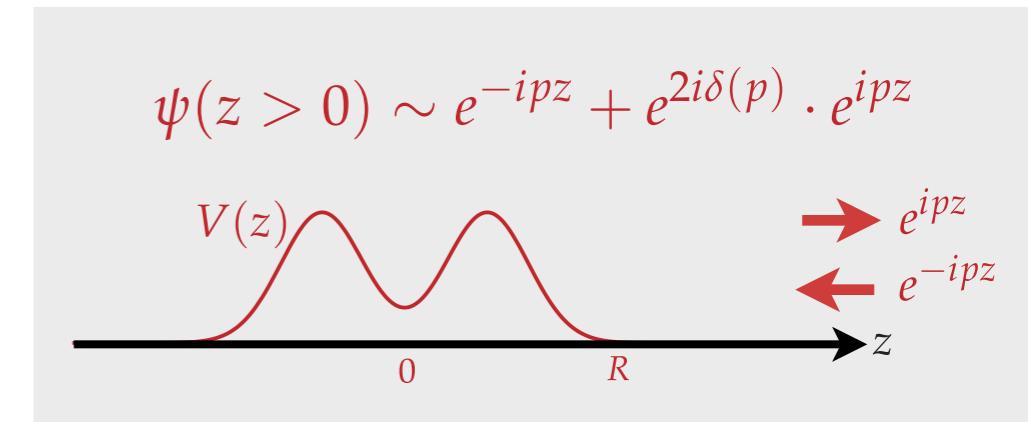
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discrete
energy
spectrum

discrete energy spectrum is determined
by the scattering amplitude

(or vice-versa)

Scattering in a finite cubic volume

- Expect a discrete spectrum in a finite periodic volume $\psi(x + L) = \psi(x)$

e.g. free particle $e^{ip(x+L)} = e^{ipx}$

quantized momentum $p = \frac{2\pi}{L}n$

- For an interacting theory $\cot \delta_\ell(E) = \mathcal{M}_\ell(E, L)$ *LÜSCHER ...*

elastic scattering phase-shift

known function

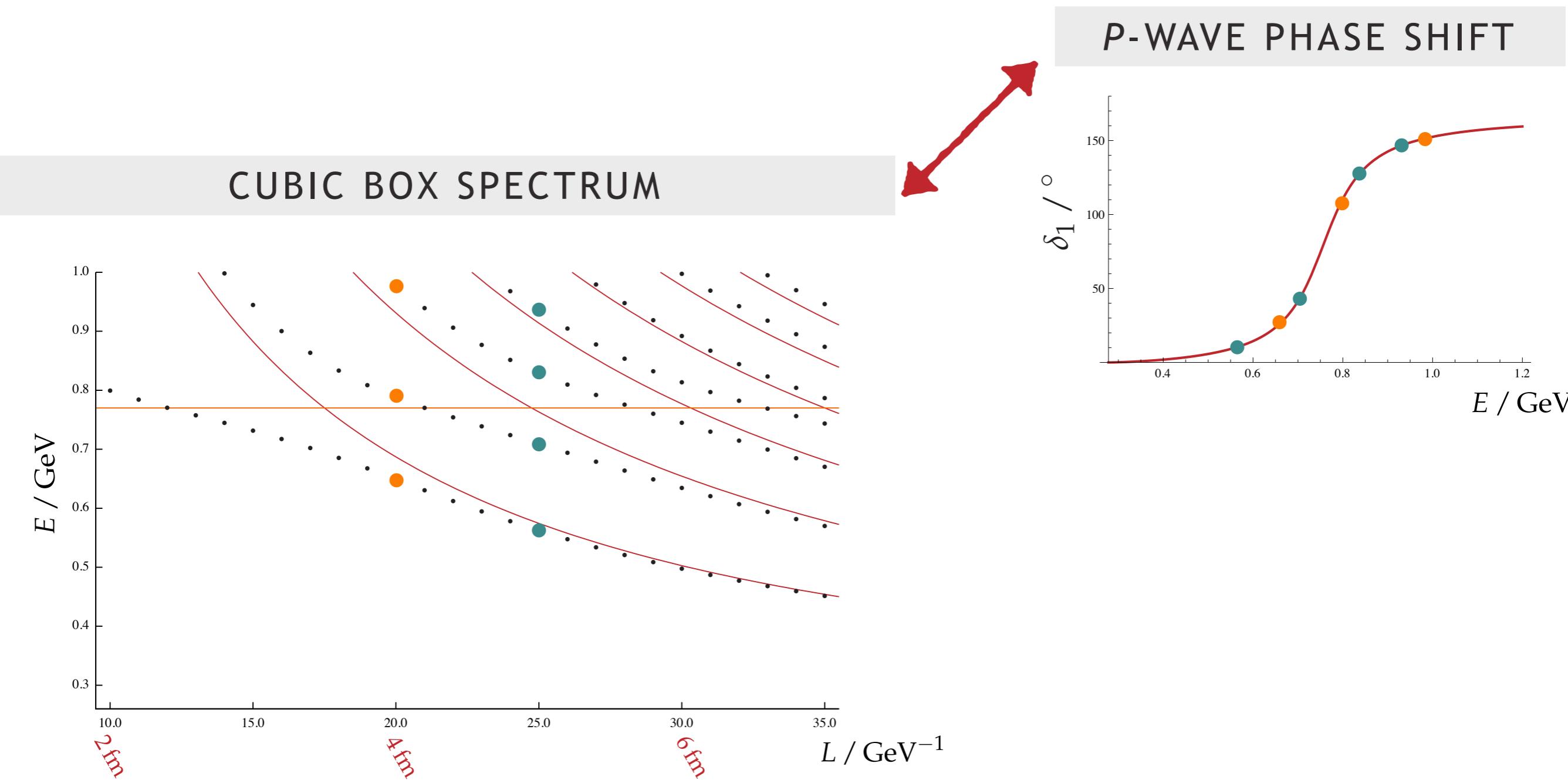
Discrete energies in a finite-volume



Discrete values of the phase-shift

Scattering in a finite cubic volume

- Experimental $\pi\pi l=1$ P -wave scattering amplitude



Coupled-channel scattering

- Finite-volume formalism recently derived (multiple methods)

*HE, JHEP 0507 011
HANSEN, PRD86 016007
BRICENO, PRD88 094507
GUO, PRD88 014051*

$$\det \left[([t^{(\ell)}(E)]_{ij}^{-1} + i\rho_i(E) \delta_{ij}) - \delta_{ij} \mathcal{M}_\ell(p_i(E)L) \right] = 0$$

scattering matrix phase space known functions *matrices in partial-wave space ...*

- However, this is one equation for multiple unknowns (per energy level) $\frac{1}{2}N(N+1)$
for N channels

- parameterize the energy dependence of t
 - try to describe a spectrum globally

“Energy-dependent” analysis

Including multi-meson operators

- Form correlator matrix with both $\bar{\psi}\Gamma\psi$ and $\pi\pi$ -like
- Include operators which resemble a pair of pions and also kaons

$$\bar{d}\Gamma u$$

$$\sum_{\hat{k}_1, \hat{k}_2} C(\Lambda, \vec{P}, \vec{k}_1, \vec{k}_2) \color{red}{\pi^\dagger(\vec{k}_1)\pi^\dagger(\vec{k}_2)}$$

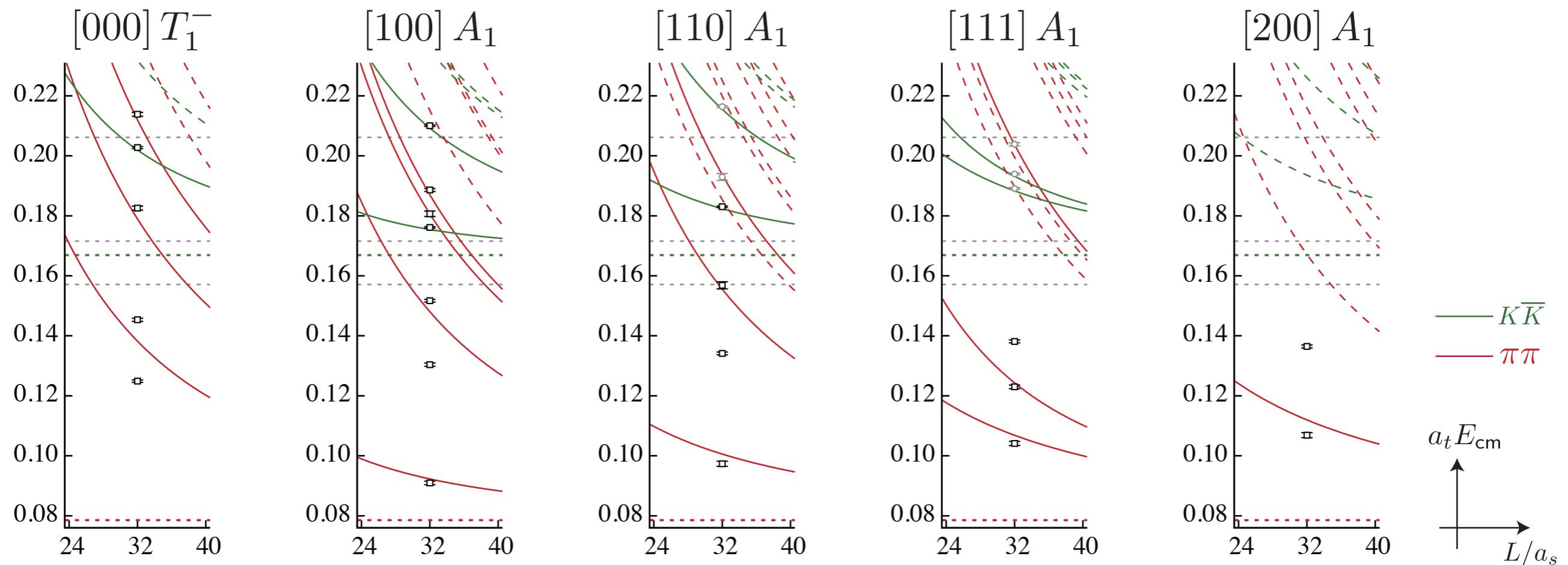
$$\sum_{\hat{k}_1, \hat{k}_2} C(\Lambda, \vec{P}, \vec{k}_1, \vec{k}_2) \color{red}{K^\dagger(\vec{k}_1)\bar{K}^\dagger(\vec{k}_2)}$$

Finite-volume spectrum - moving frames

- Non-interacting thresholds and energies as a function of \vec{k}

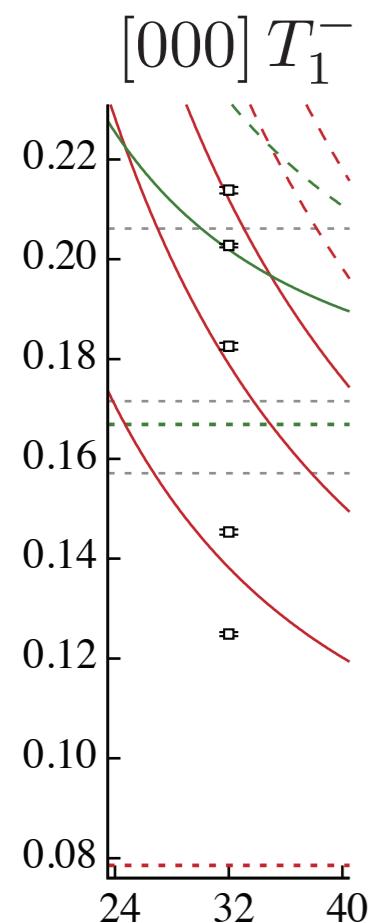
$m_\pi \sim 236 \text{ MeV}$

Momentum & lattice irrep labels: $[\vec{k}] \Lambda$



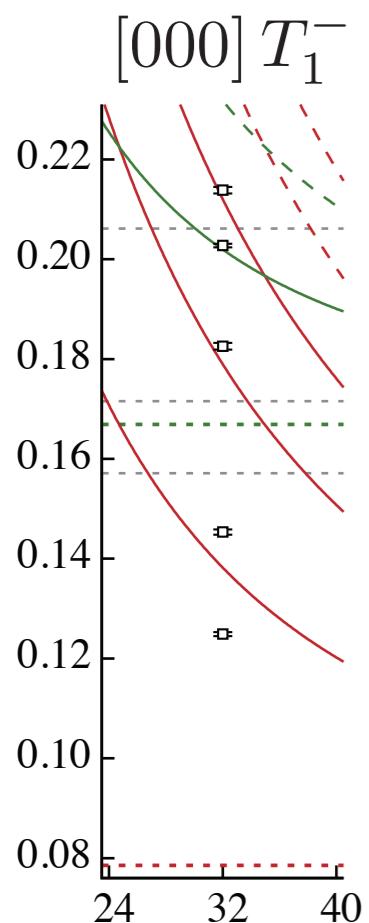
$\pi\pi/KK$ scattering

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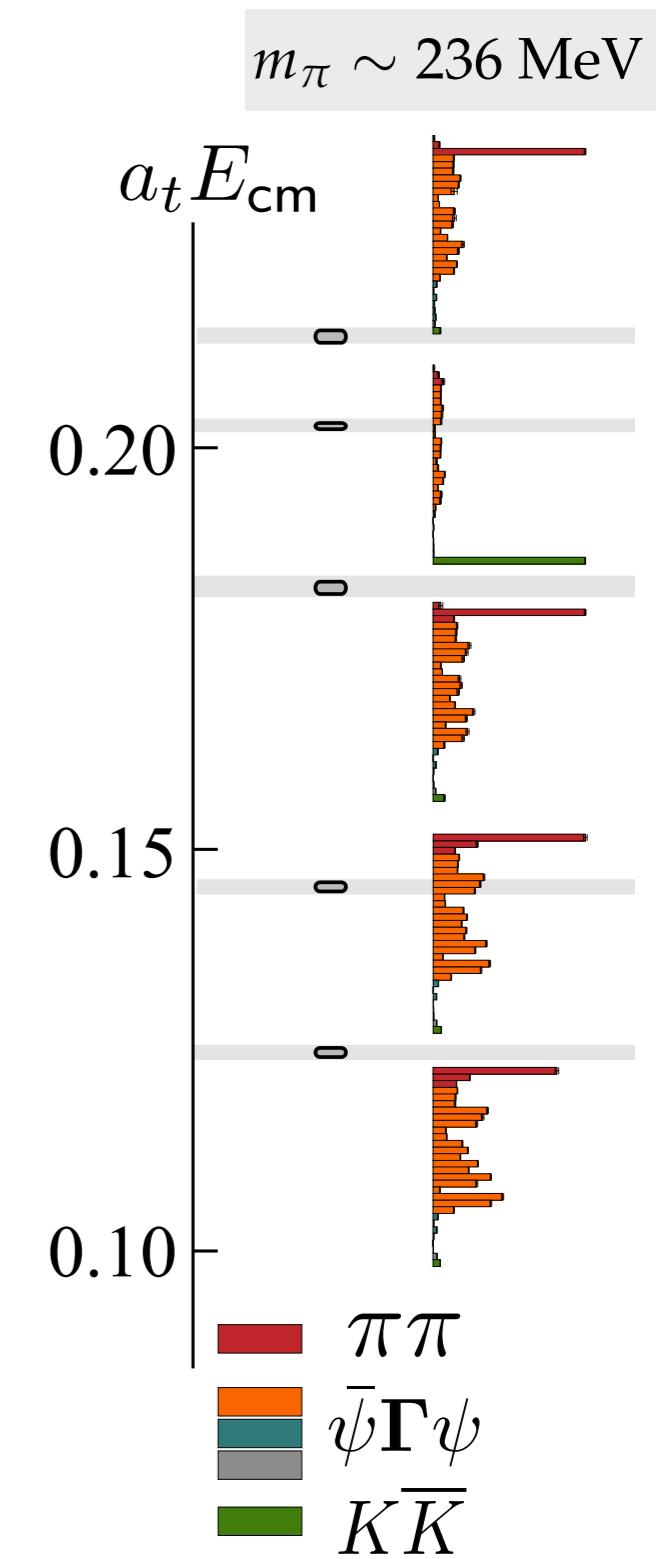
- Consider size of operator overlaps $\langle \mathbf{n} | \mathcal{O}_i^\dagger | \emptyset \rangle$

$\pi\pi/KK$ scattering

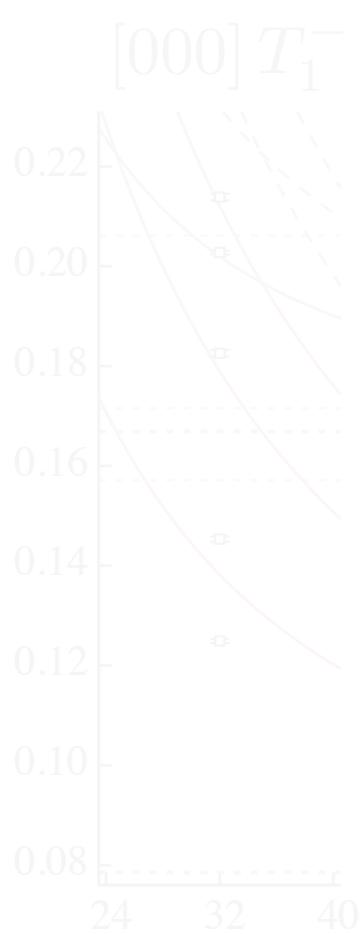


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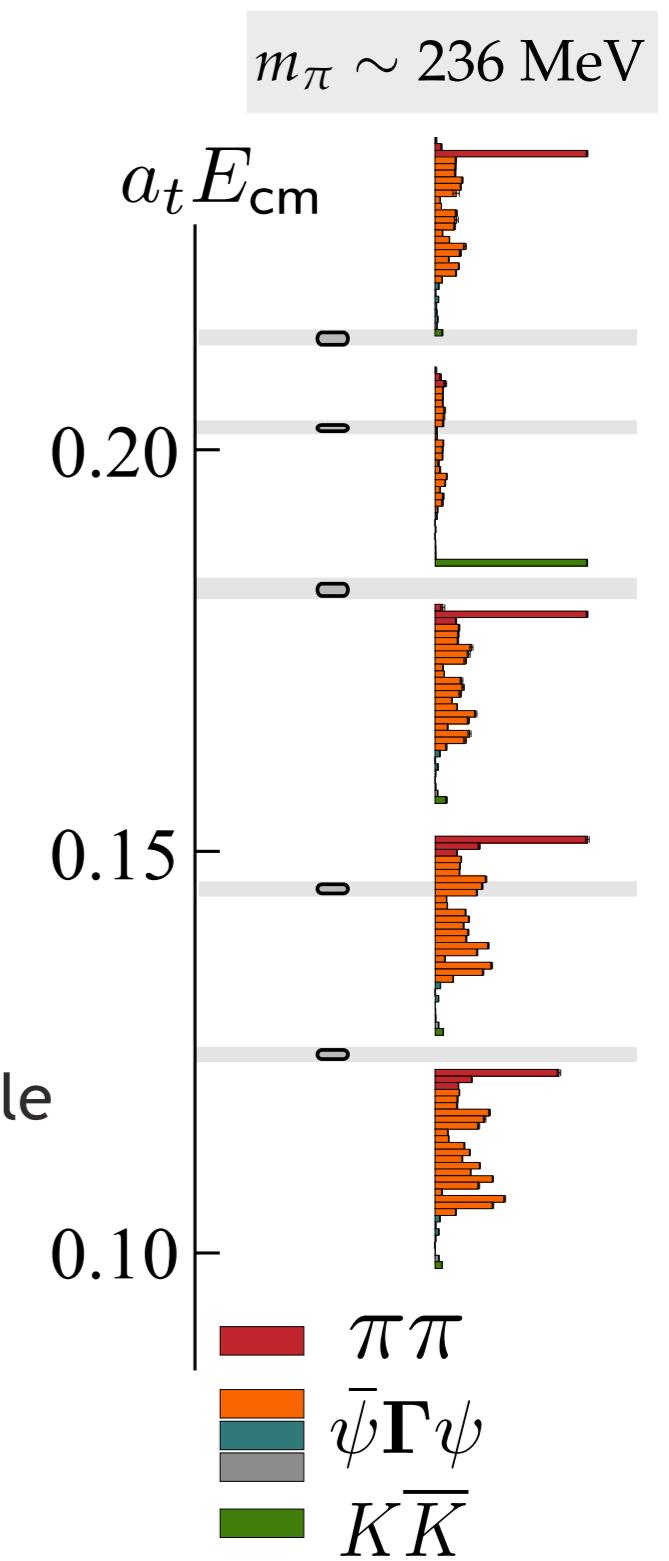
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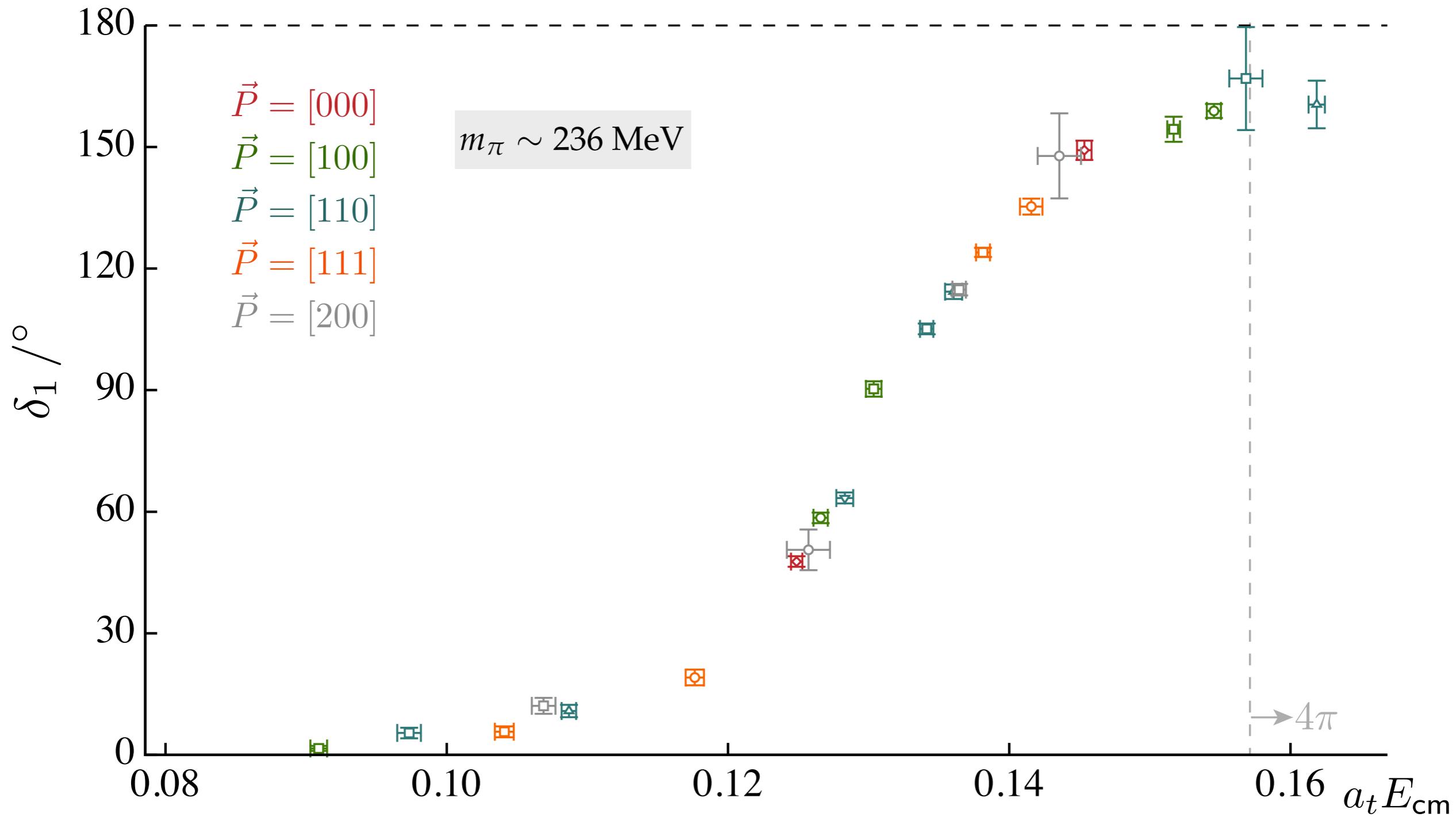
$$\langle n | \mathcal{O}_i^\dagger | \emptyset \rangle$$

» Finite-V: ad-mixture of single & two-particle



$\pi\pi$ P -wave phase-shift

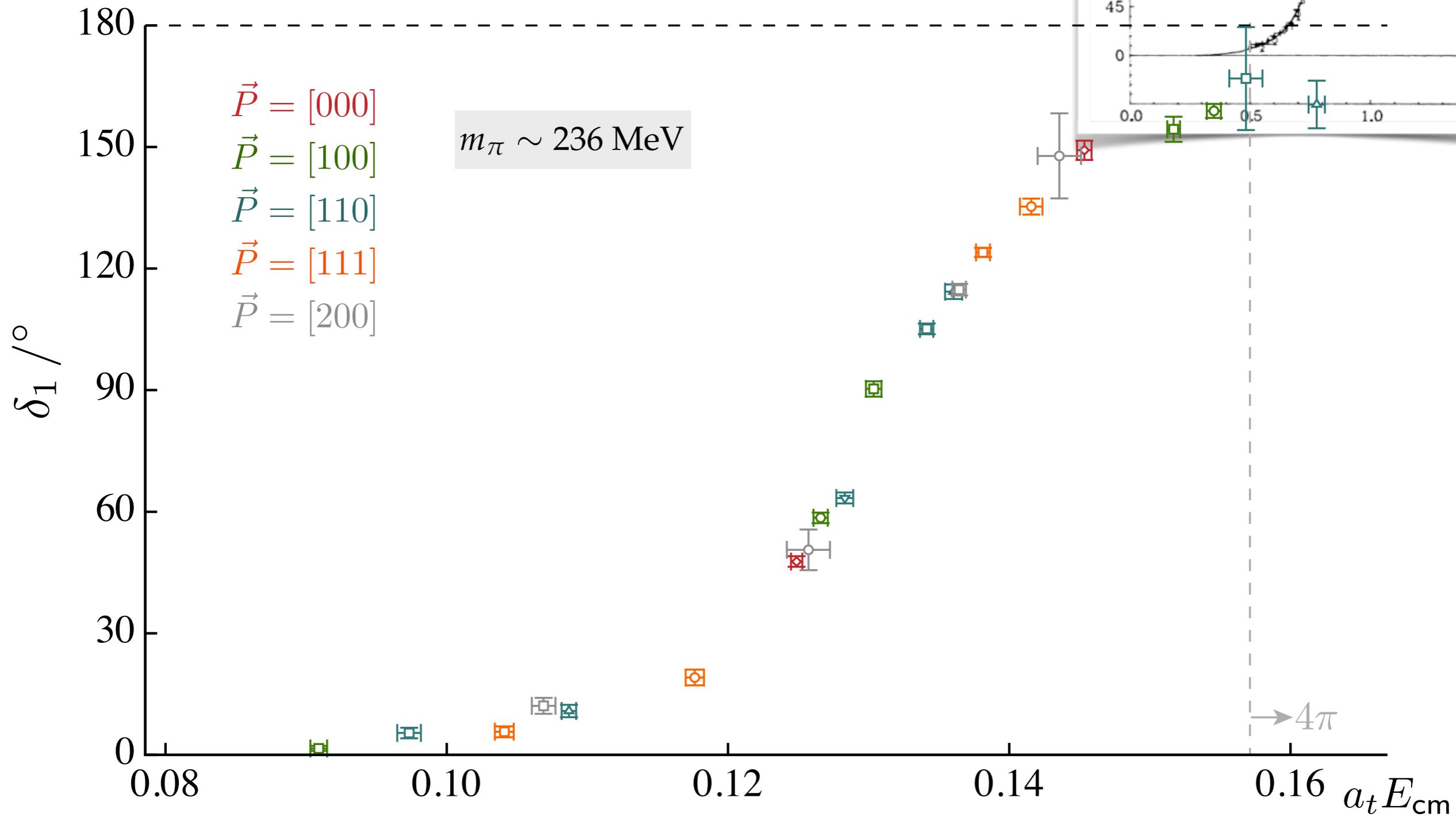
- Restrict to elastic region below $K\bar{K}$ threshold



To appear very soon...

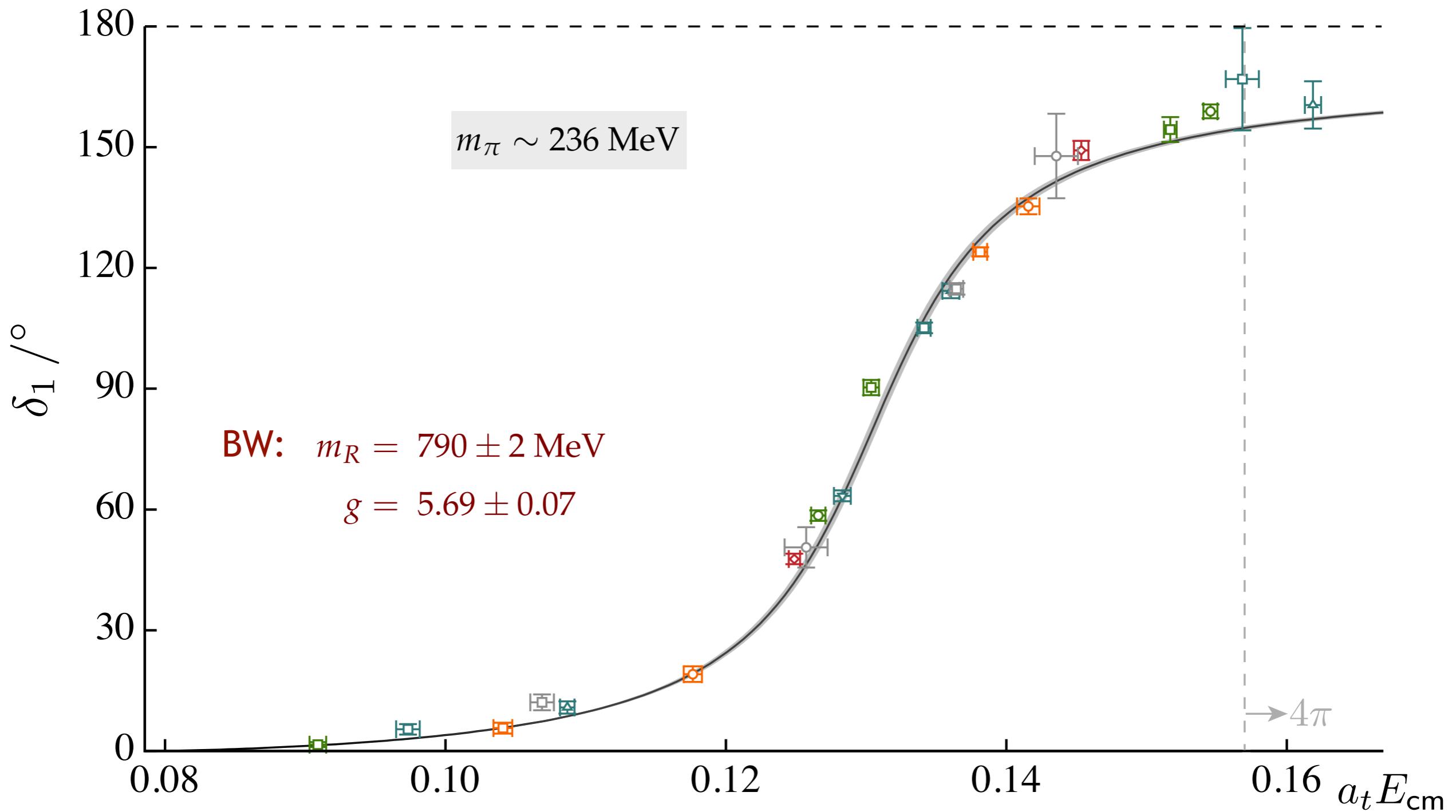
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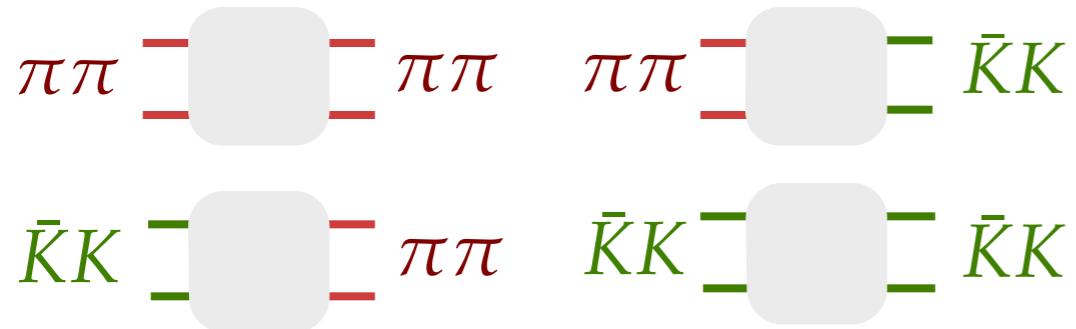
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$\pi\pi$ P -wave phase-shift



ρ resonance as a coupled channel system

- Parameterize the t -matrix in a unitarity conserving way



- Compute finite-volume spectrum

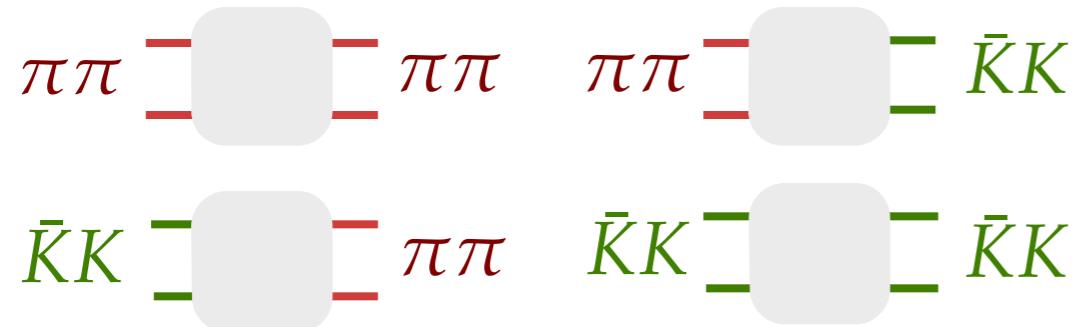
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ρ resonance as a coupled channel system

- Parameterize the t -matrix in a unitarity conserving way



$$t_{ij}^{-1}(E) = K_{ij}^{-1}(E) + \delta_{ij} I_i(E)$$

$$K_{ij}(E) = \frac{g_i g_j}{m^2 - E^2} + \gamma_{ij}$$

- Vary the parameters, solving

$$\det \left[([t^{(\ell)}(E)]_{ij}^{-1} + i\rho_i(E) \delta_{ij}) - \delta_{ij} \mathcal{M}_\ell(E, L) \right] = 0$$

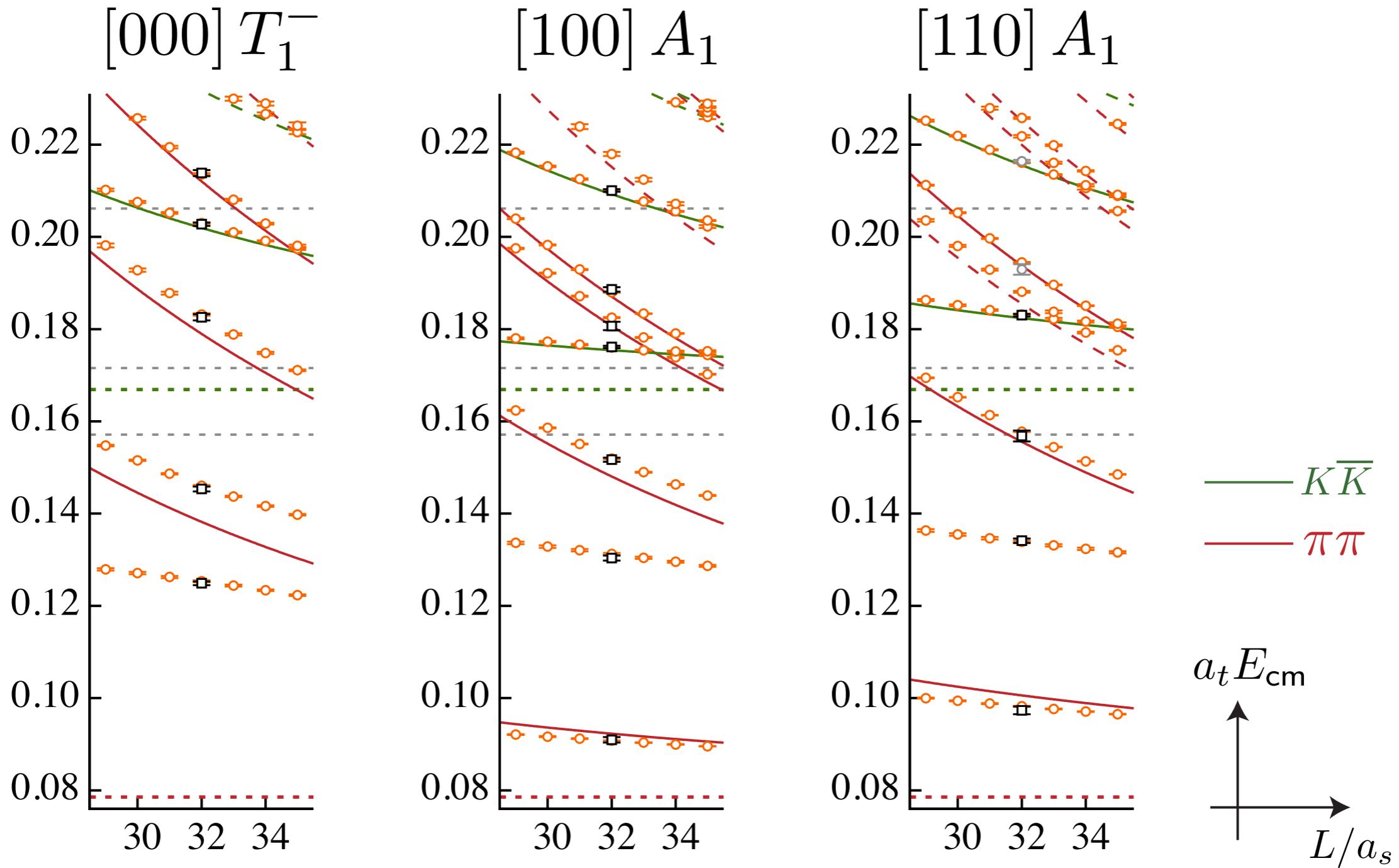
for the spectrum in each irreducible representation & momentum

Want pole mass and couplings of t-matrix

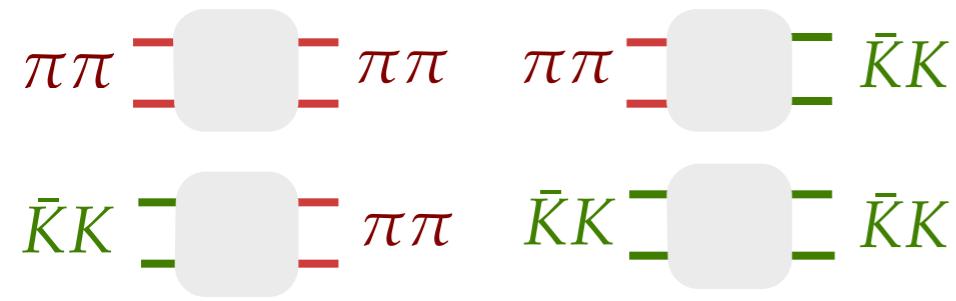
$\pi\pi/KK$ scattering

- Data points (black) compared to parameterization (gold)

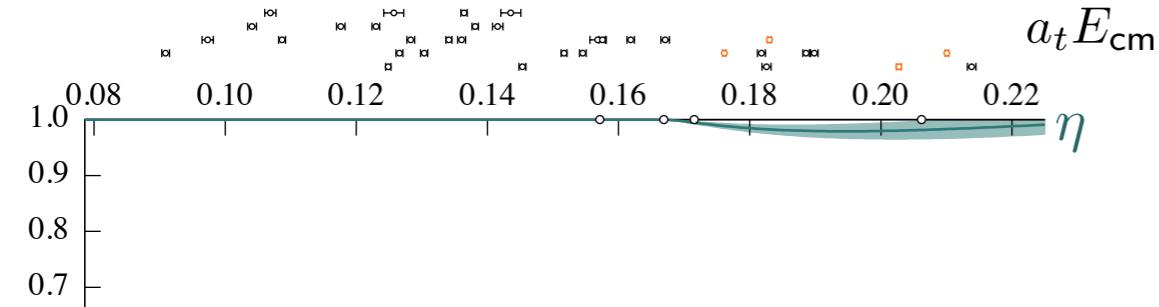
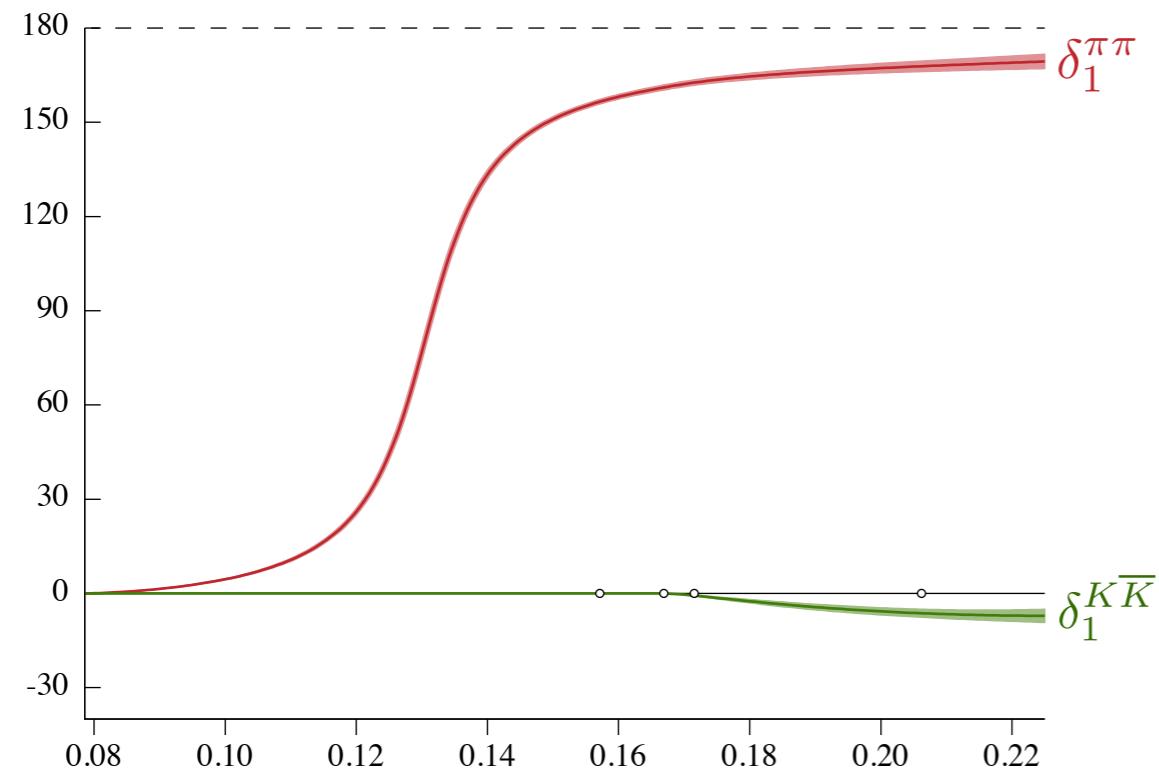
$m_\pi \sim 236$ MeV



ρ resonance as a coupled channel system

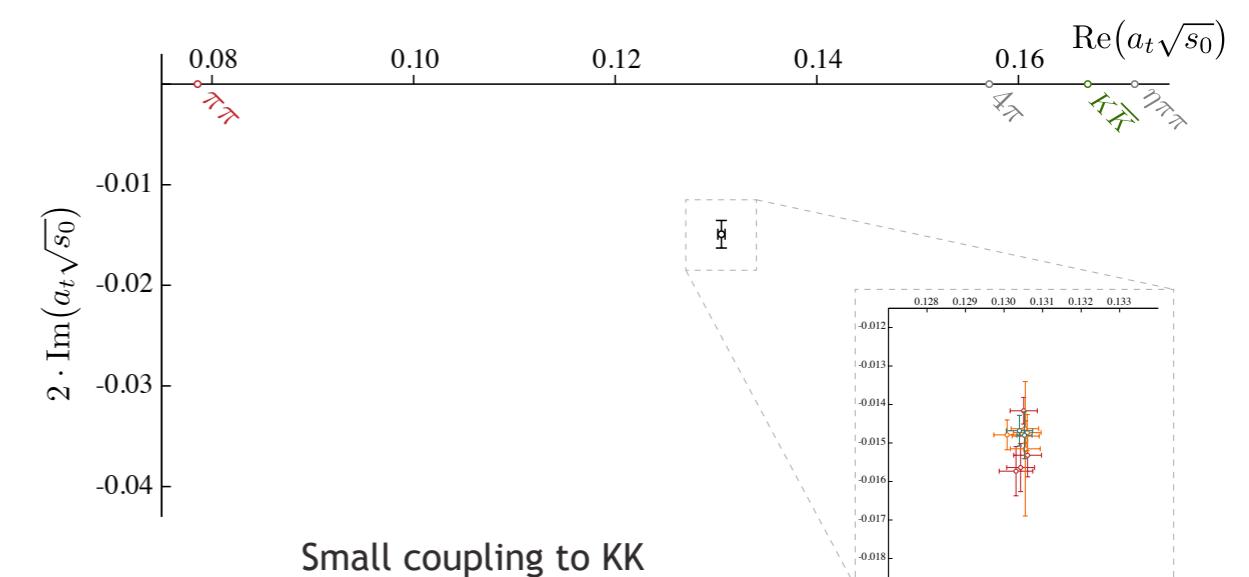


Phase shifts & inelasticity



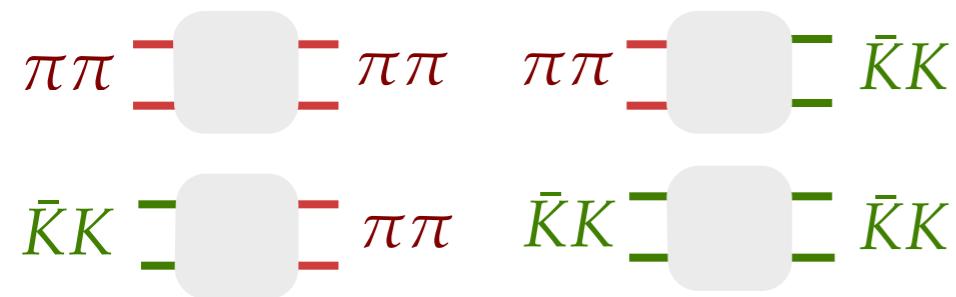
$m_\pi \sim 236 \text{ MeV}$

t-matrix pole location

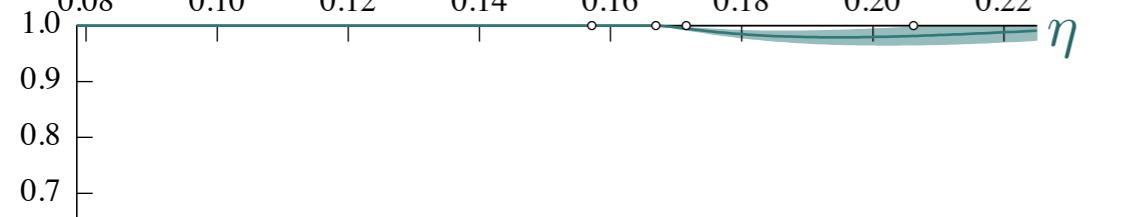
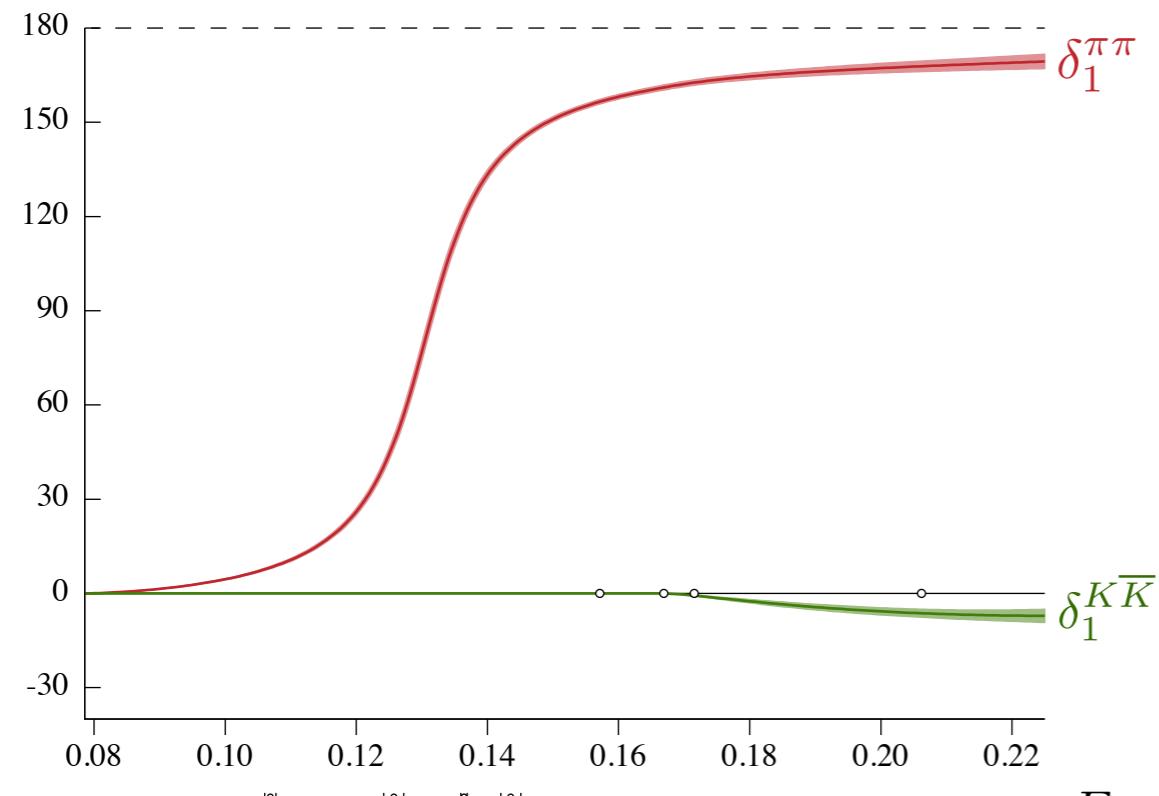


Small coupling to KK

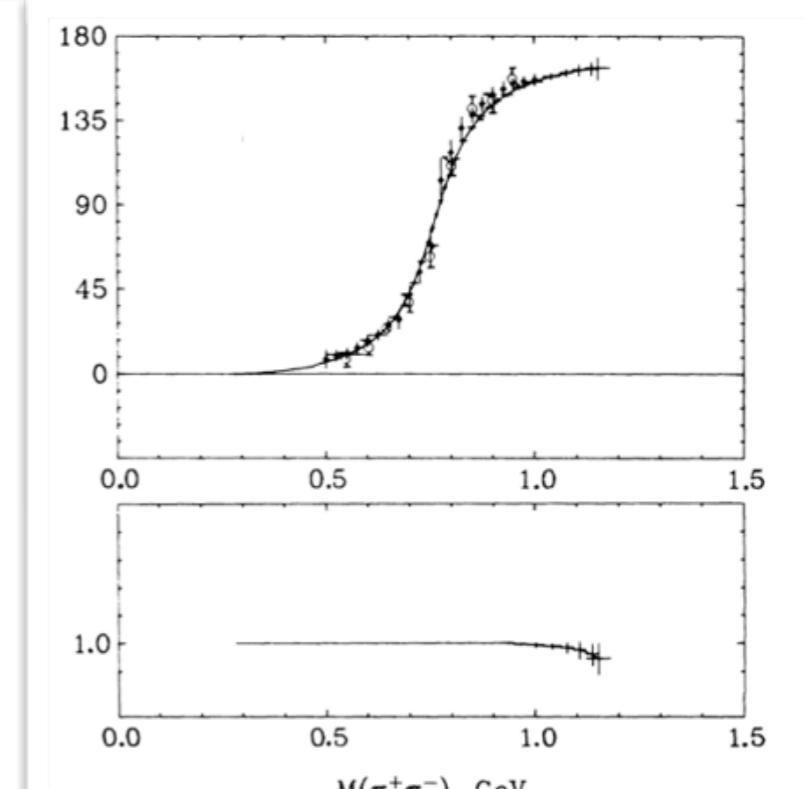
ρ resonance as a coupled channel system



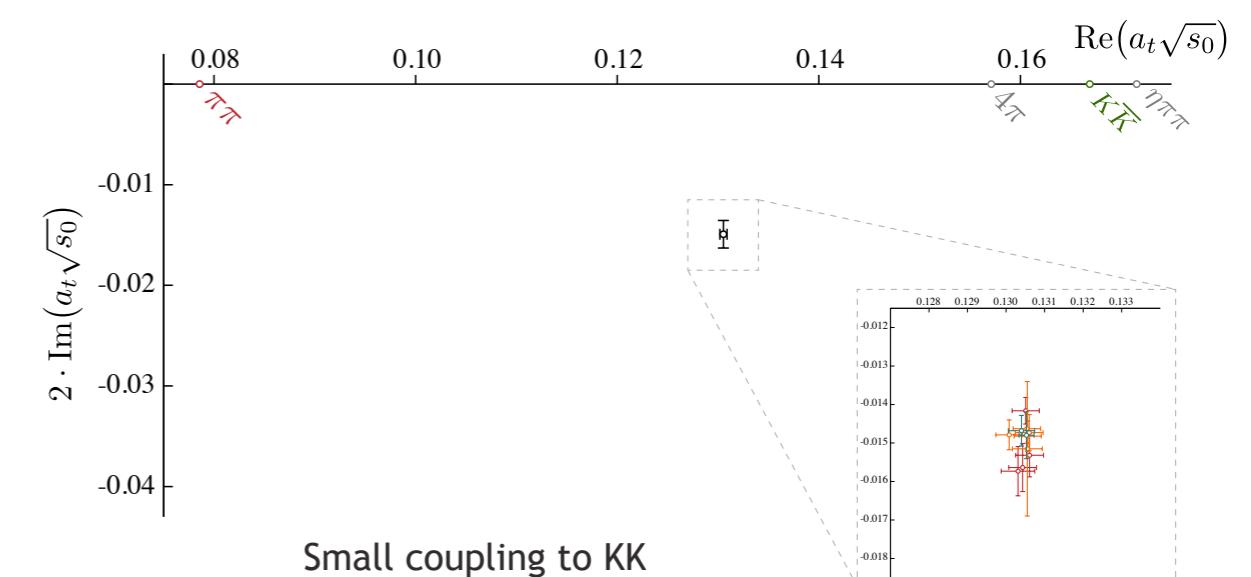
Phase shifts & inelasticity



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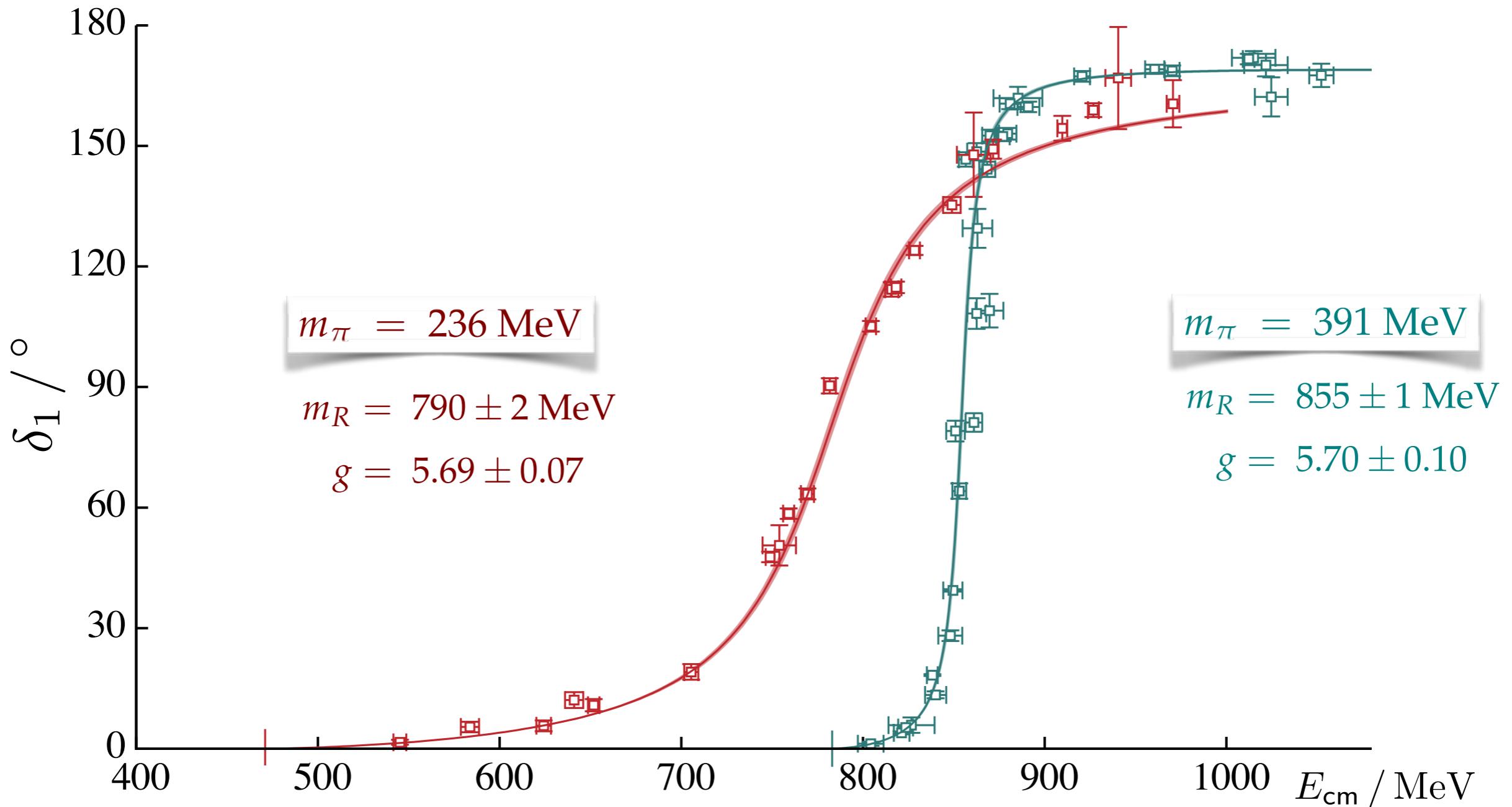
t-matrix pole location



Small coupling to KK

ρ resonance at different pion masses

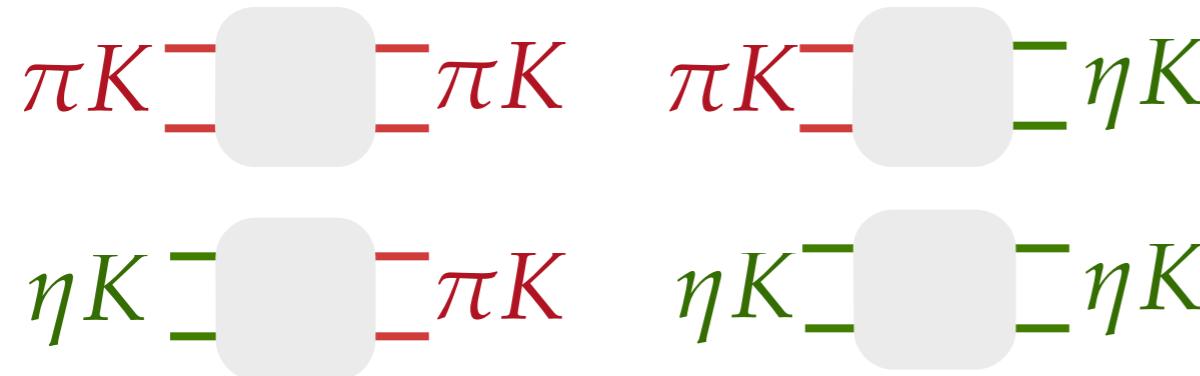
- BW couplings nearly constant in pion mass (will come back to this later...)



PRD 87 034505

$\pi K/\eta K$ scattering & kaon resonances

- Example of coupled-channel scattering



- Compute finite-volume spectrum

$$\bar{u} \Gamma s$$

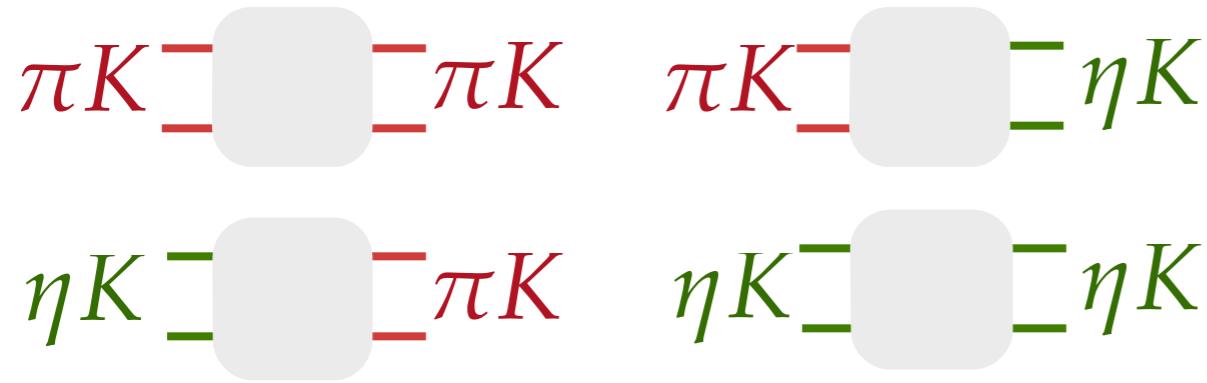
$$\sum_{\hat{k}_1, \hat{k}_2} C(\Lambda, \vec{P}; \vec{k}_1, \vec{k}_2) \pi^\dagger(\vec{k}_1) K^\dagger(\vec{k}_2)$$

$$\sum_{\hat{k}_1, \hat{k}_2} C(\Lambda, \vec{P}; \vec{k}_1, \vec{k}_2) \eta^\dagger(\vec{k}_1) K^\dagger(\vec{k}_2)$$

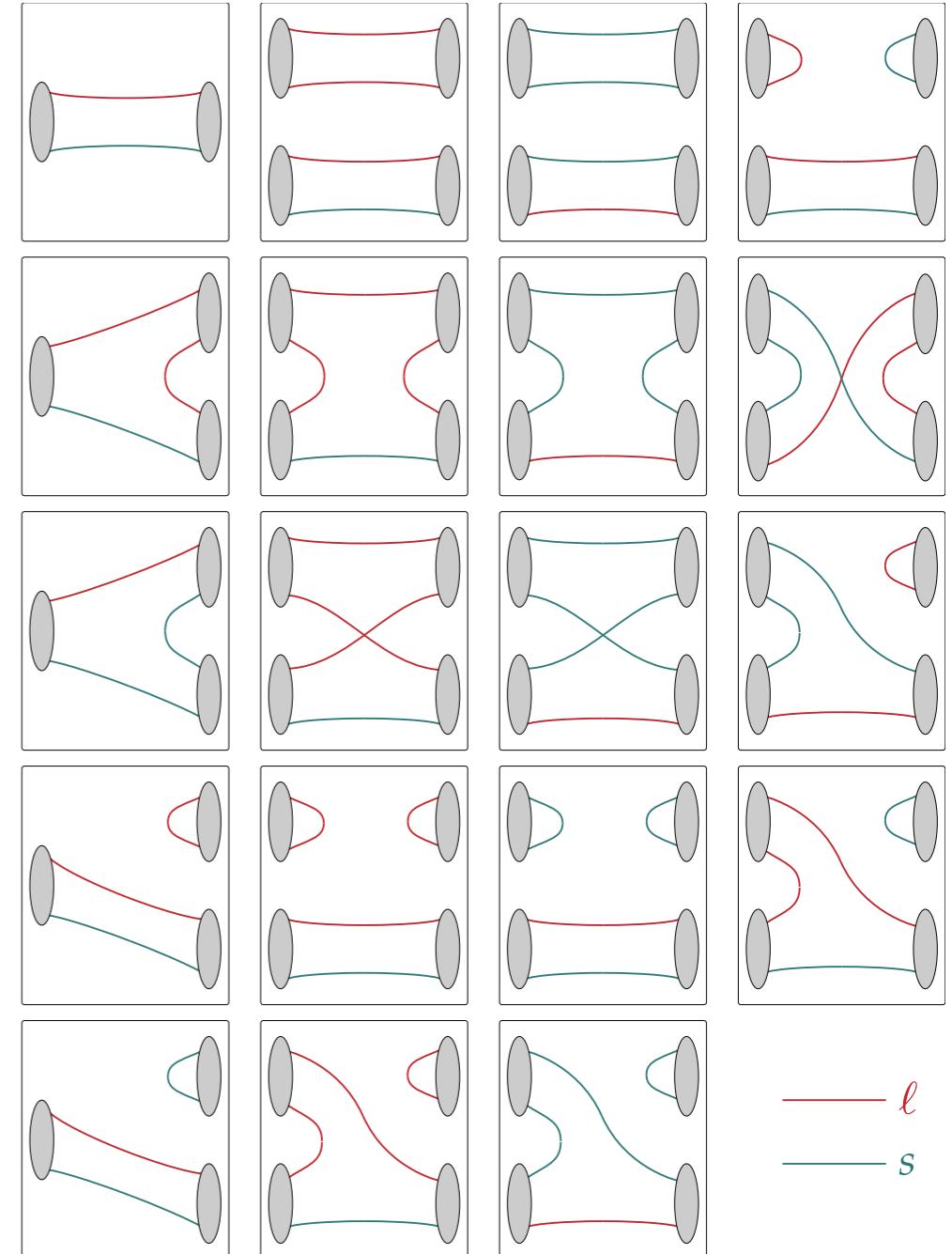
PRL 113 182001
PRD 91 054008

$\pi K/\eta K$ scattering & kaon resonances

- Example of coupled-channel scattering



WICK CONTRACTIONS



- Compute finite-volume spectrum

$$\bar{u} \Gamma s$$

$$\sum_{\hat{k}_1, \hat{k}_2} C(\Lambda, \vec{P}; \vec{k}_1, \vec{k}_2) \pi^\dagger(\vec{k}_1) K^\dagger(\vec{k}_2)$$

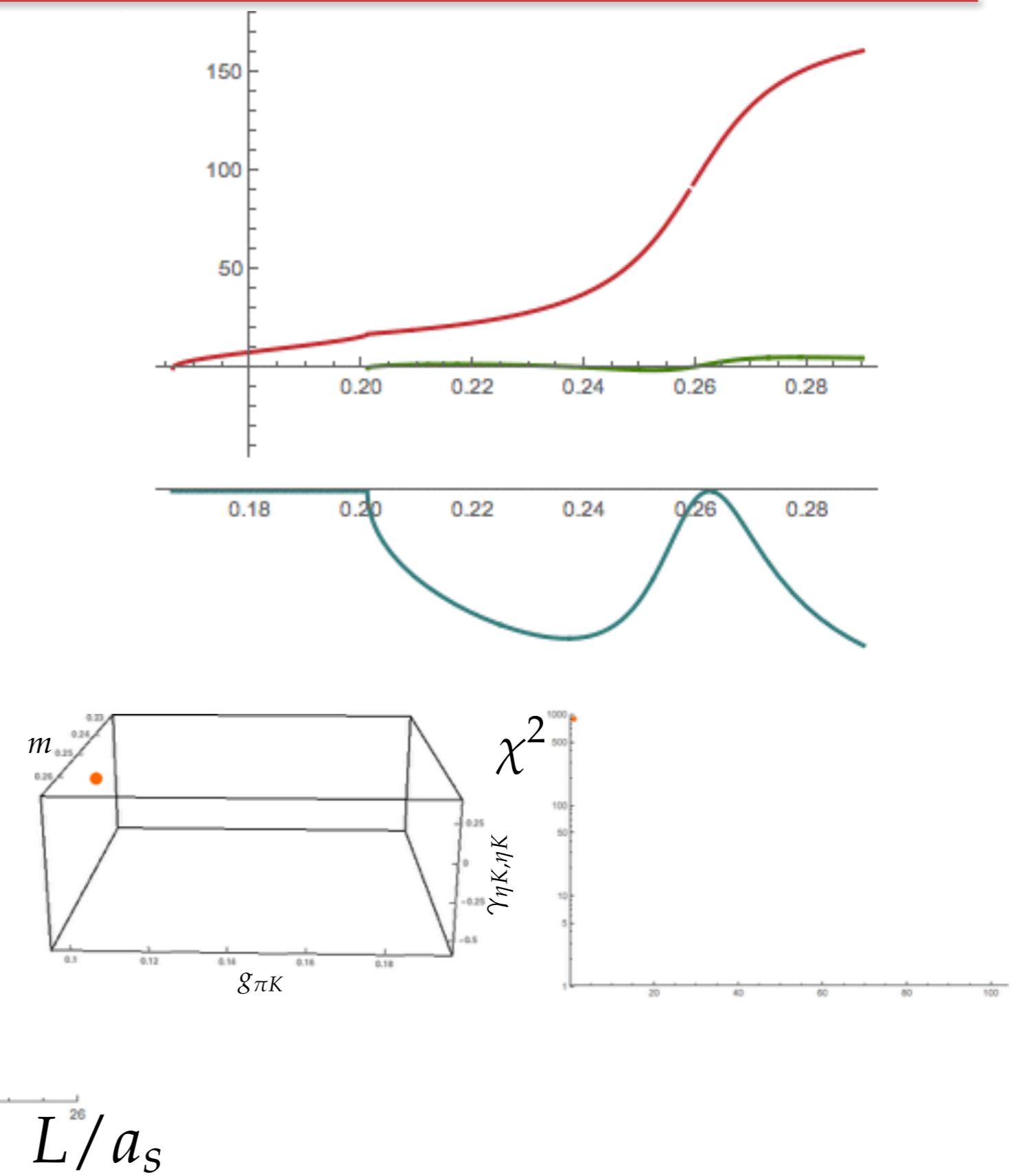
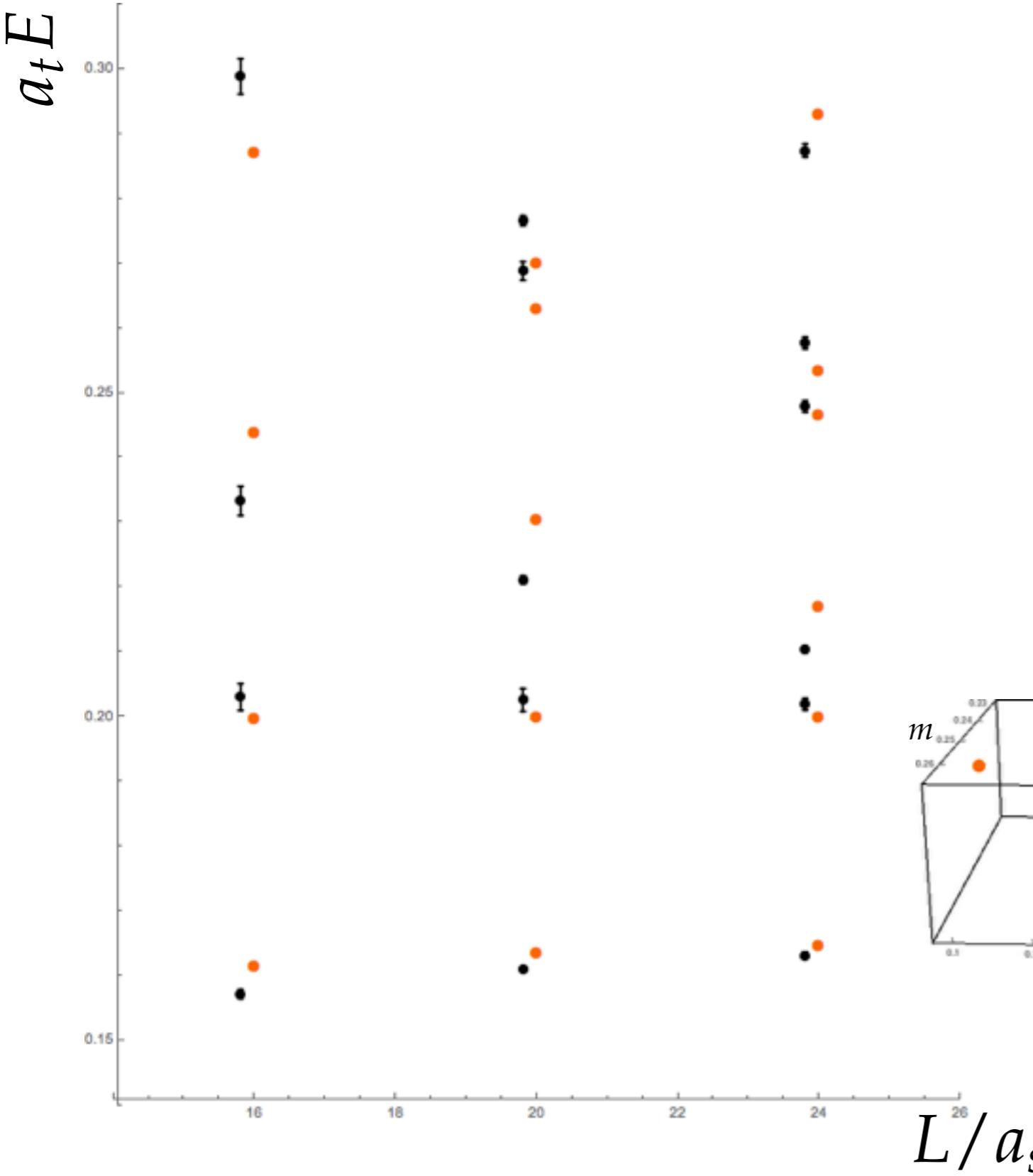
$$\sum_{\hat{k}_1, \hat{k}_2} C(\Lambda, \vec{P}; \vec{k}_1, \vec{k}_2) \eta^\dagger(\vec{k}_1) K^\dagger(\vec{k}_2)$$

PRL 113 182001

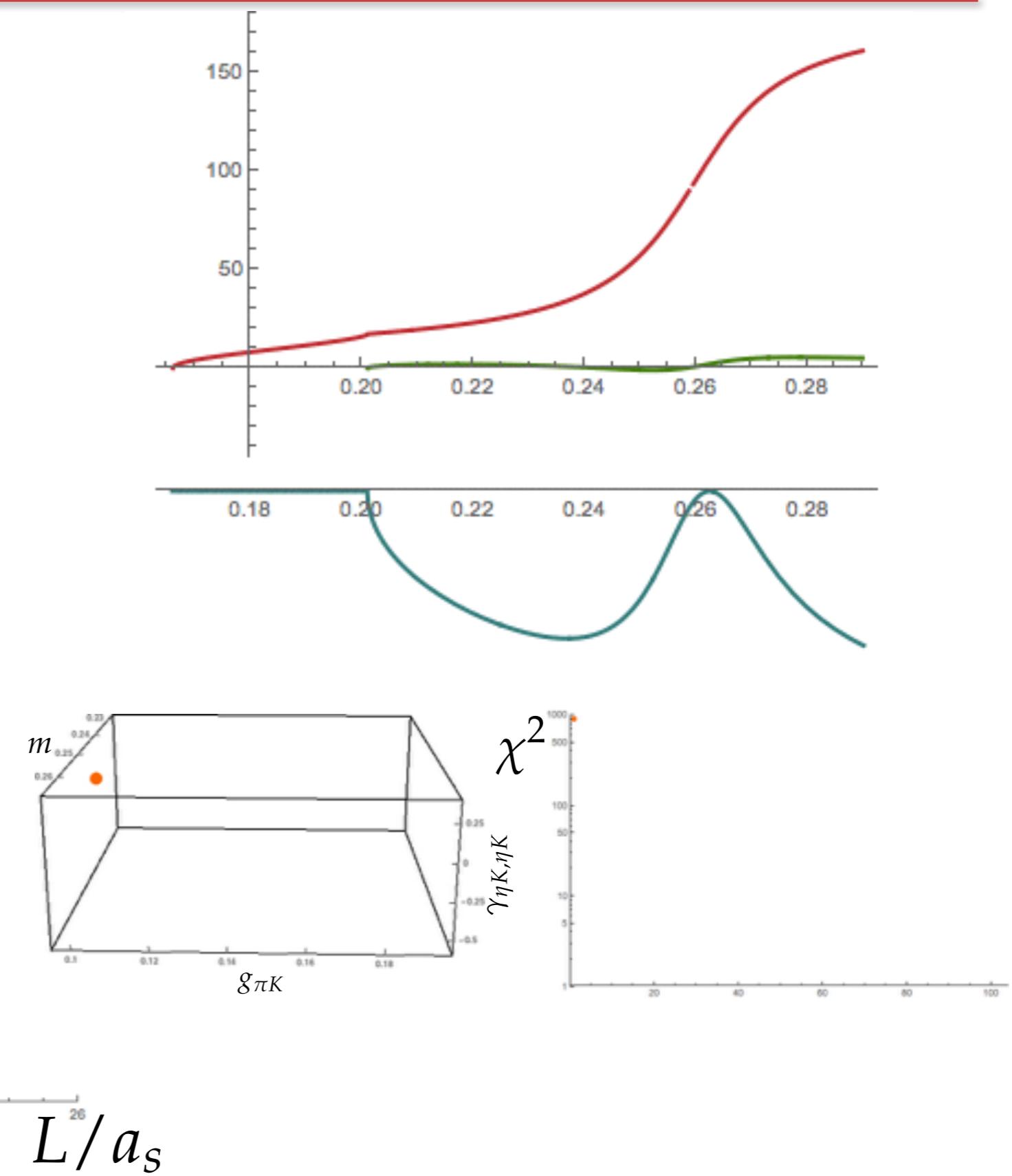
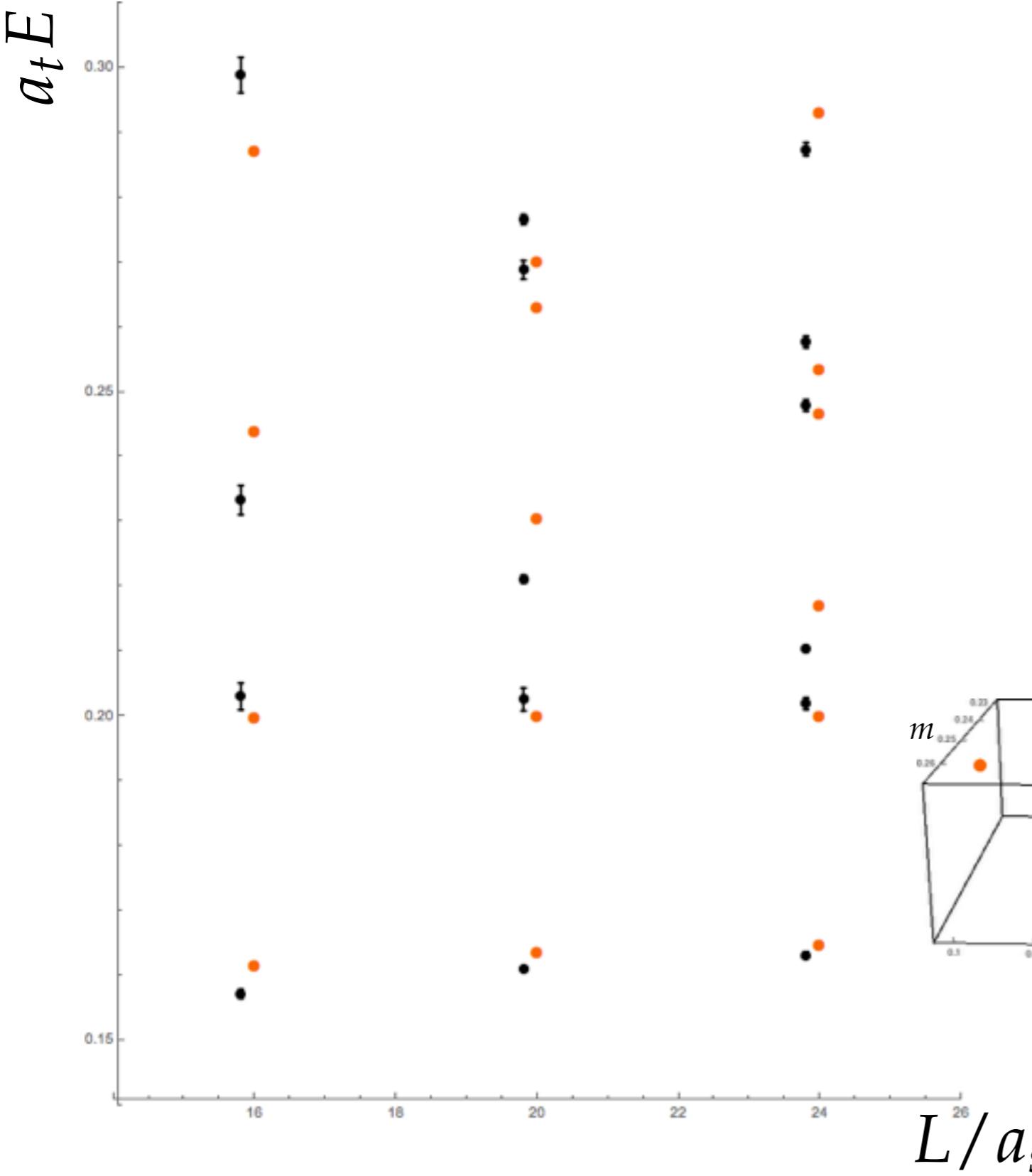
PRD 91 054008

— ℓ
— s

$\pi K/\eta K$ scattering



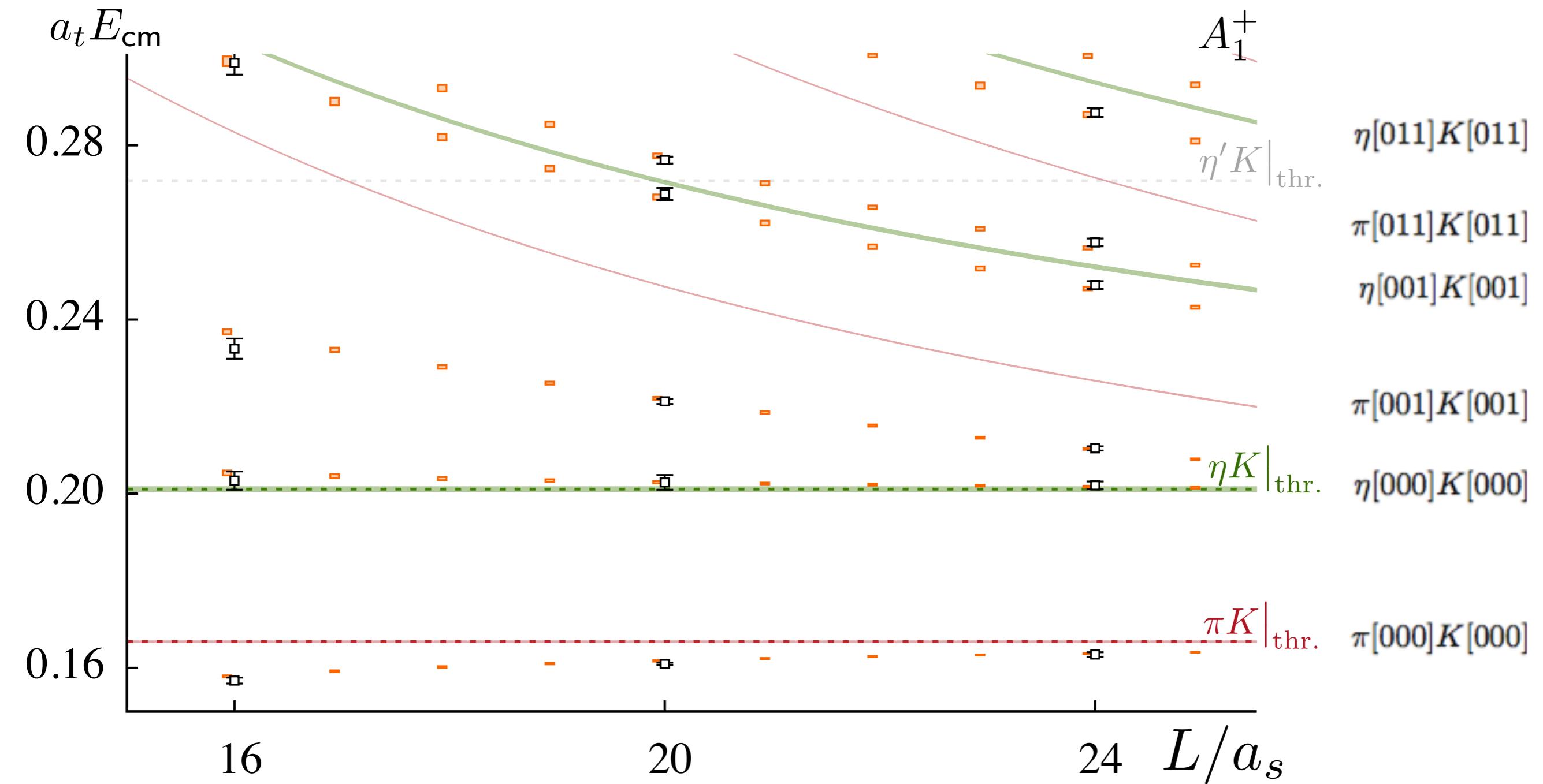
$\pi K/\eta K$ scattering



$\pi K/\eta K$ scattering

$$\chi^2/N_{\text{dof}} = \frac{6.40}{15 - 6} = 0.71$$

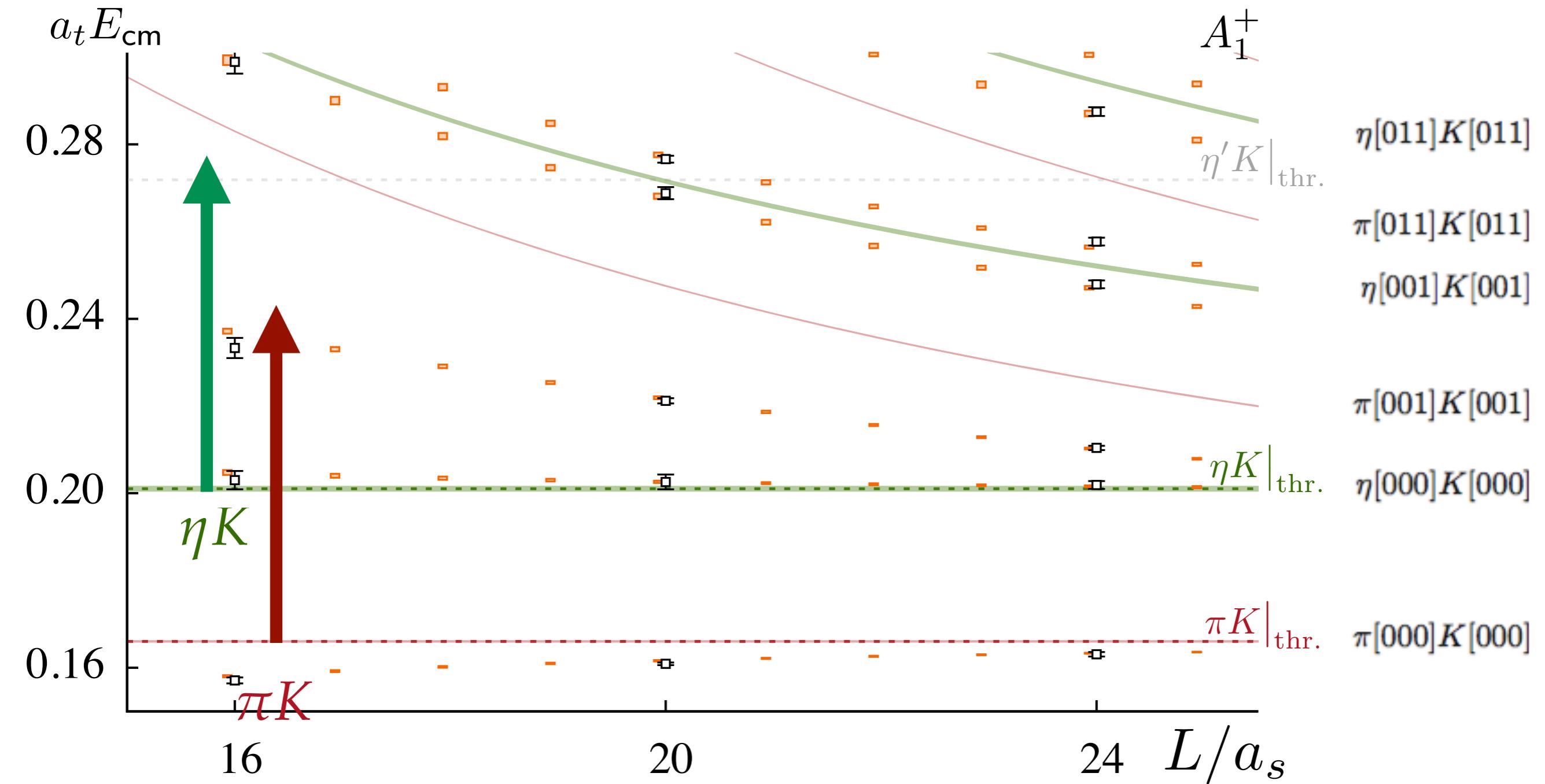
$m_\pi \sim 391 \text{ MeV}$



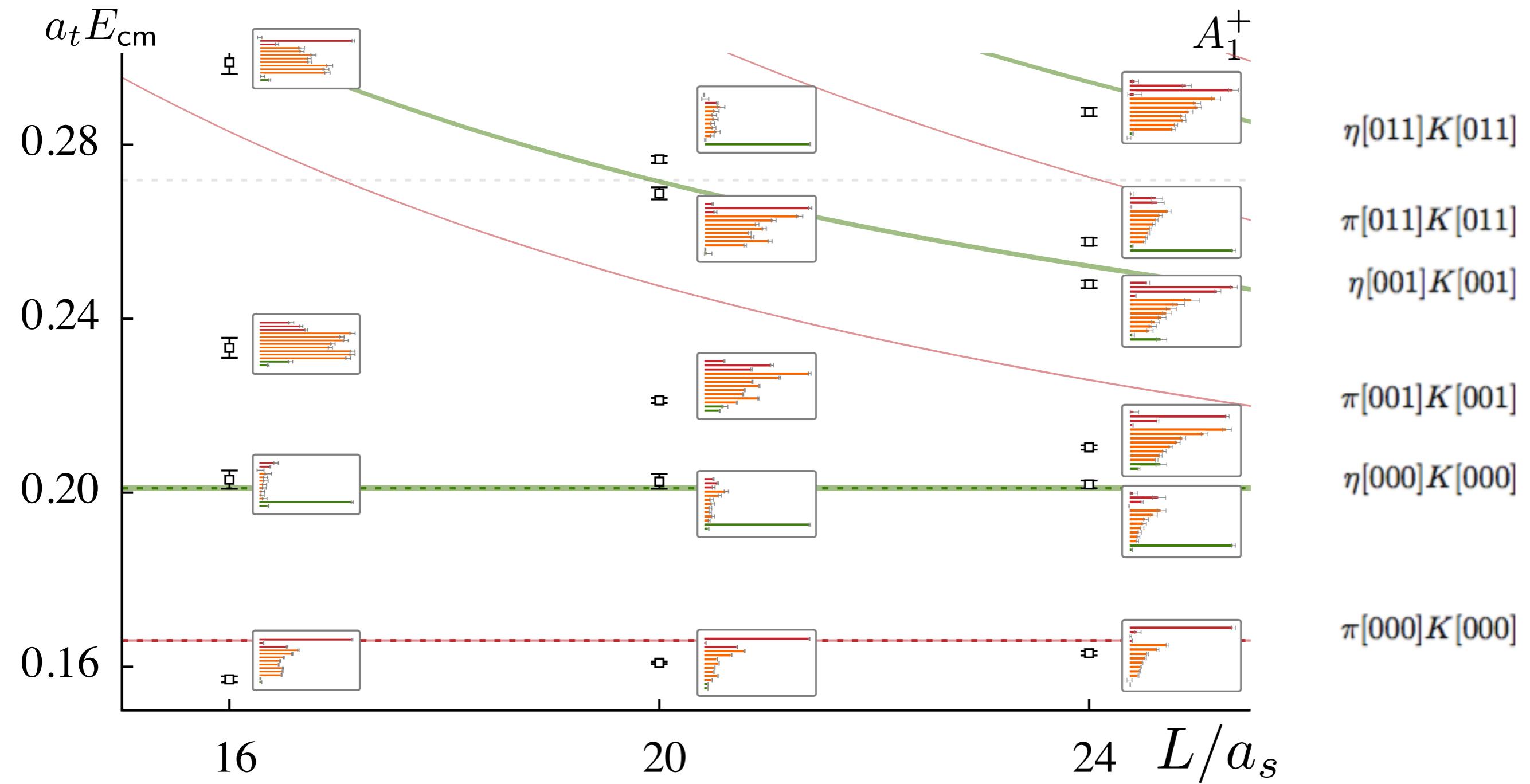
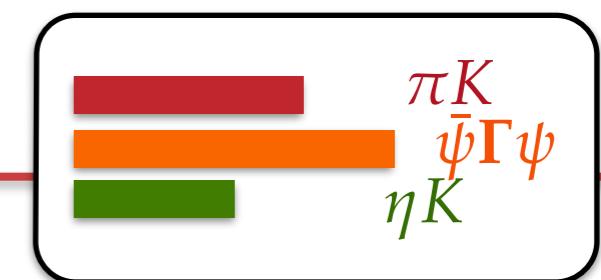
$\pi K/\eta K$ scattering

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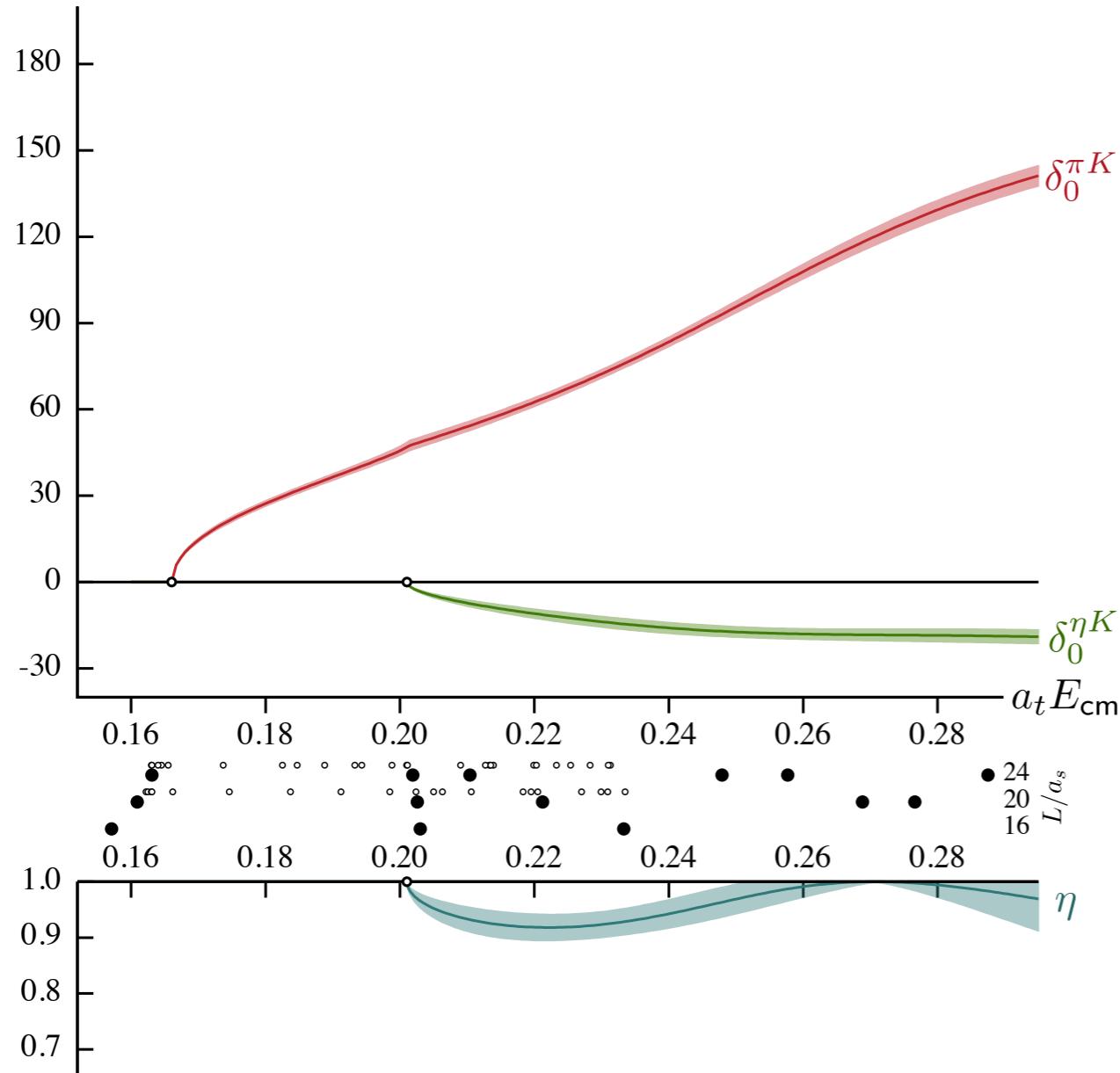
$m_\pi \sim 391$ MeV

- Describe all the finite-volume spectra

$$\chi^2/N_{\text{dof}} = \frac{49.1}{61 - 6} = 0.89$$

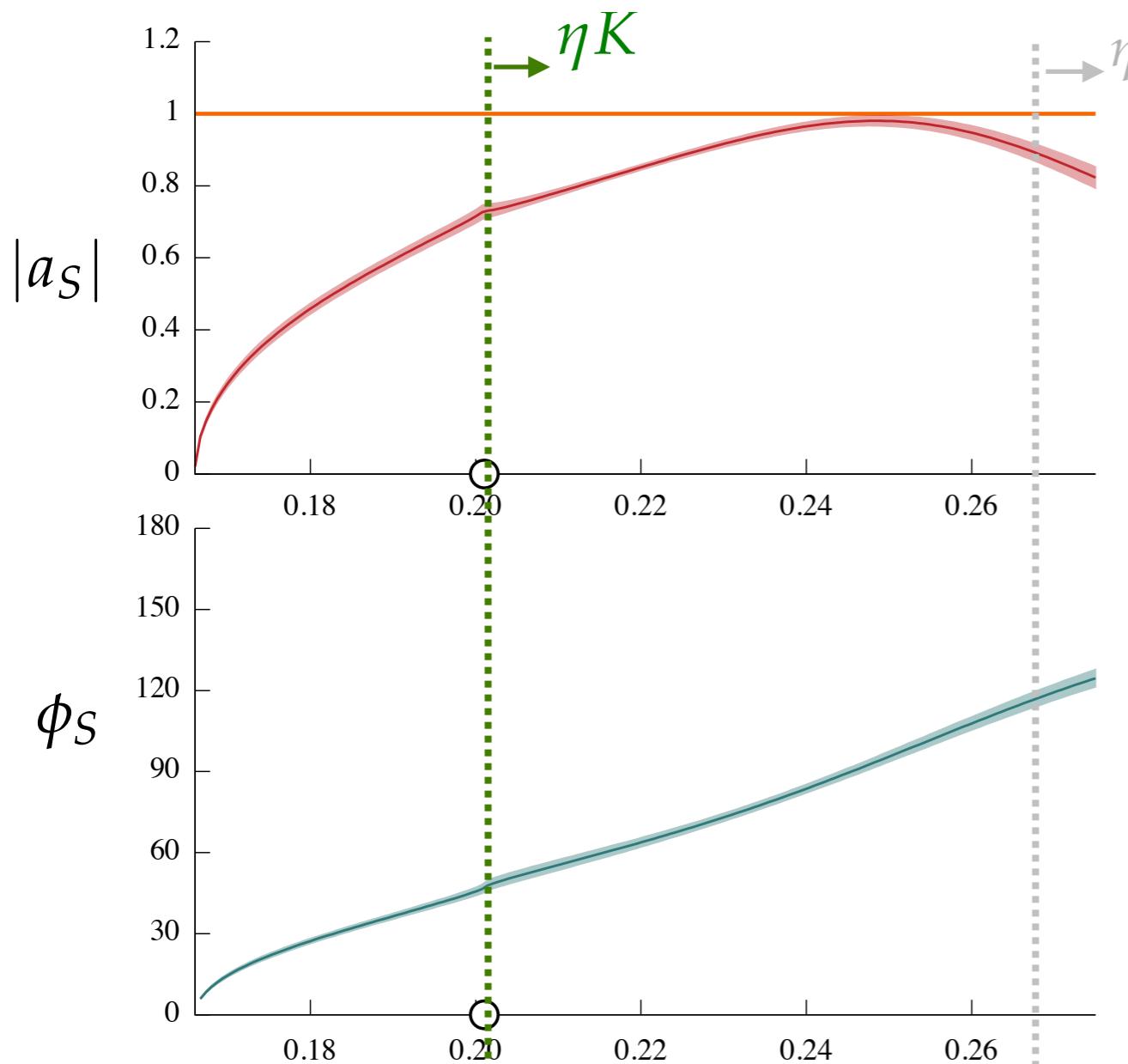
$$S_{\pi K, \pi K} = \eta e^{2i\delta^{\pi K}}$$
$$S_{\eta K, \eta K} = \eta e^{2i\delta^{\eta K}}$$

S-WAVE $\pi K/\eta K$ SCATTERING

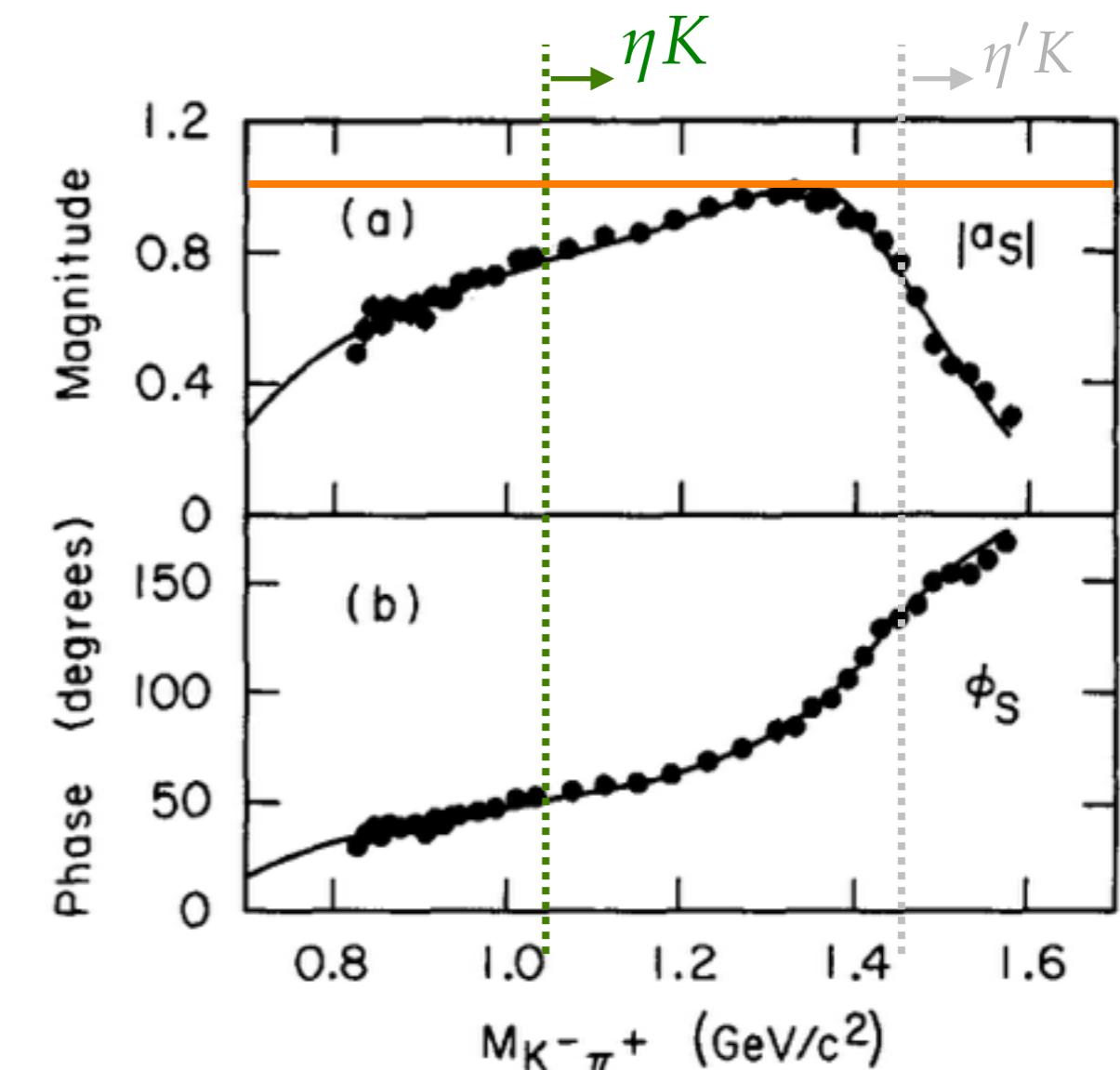


Versus experimental scattering

S-WAVE $\pi K \rightarrow \pi K$ AMPLITUDE



$m_\pi \sim 391 \text{ MeV}$

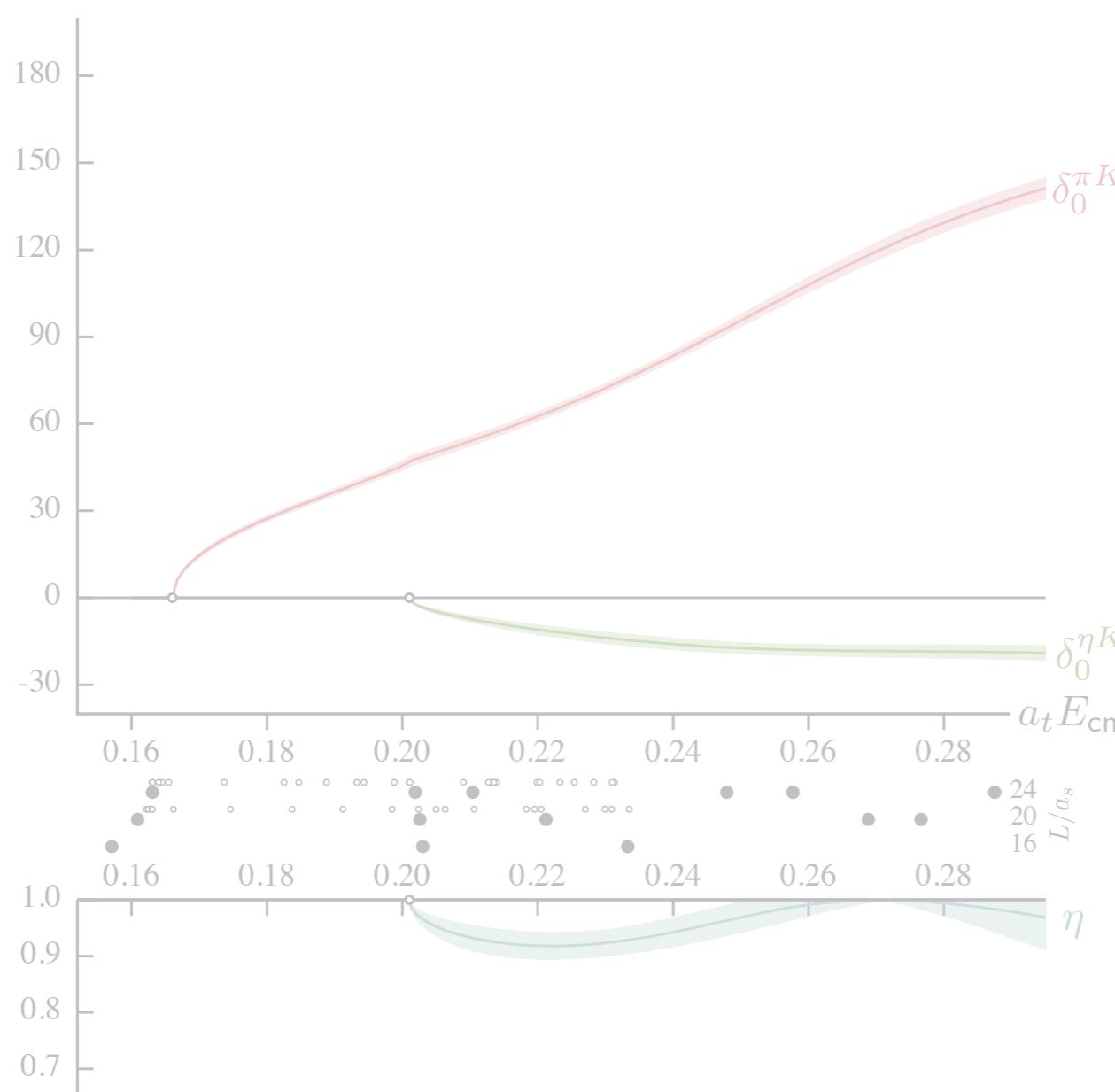


LASS, NPB296 493

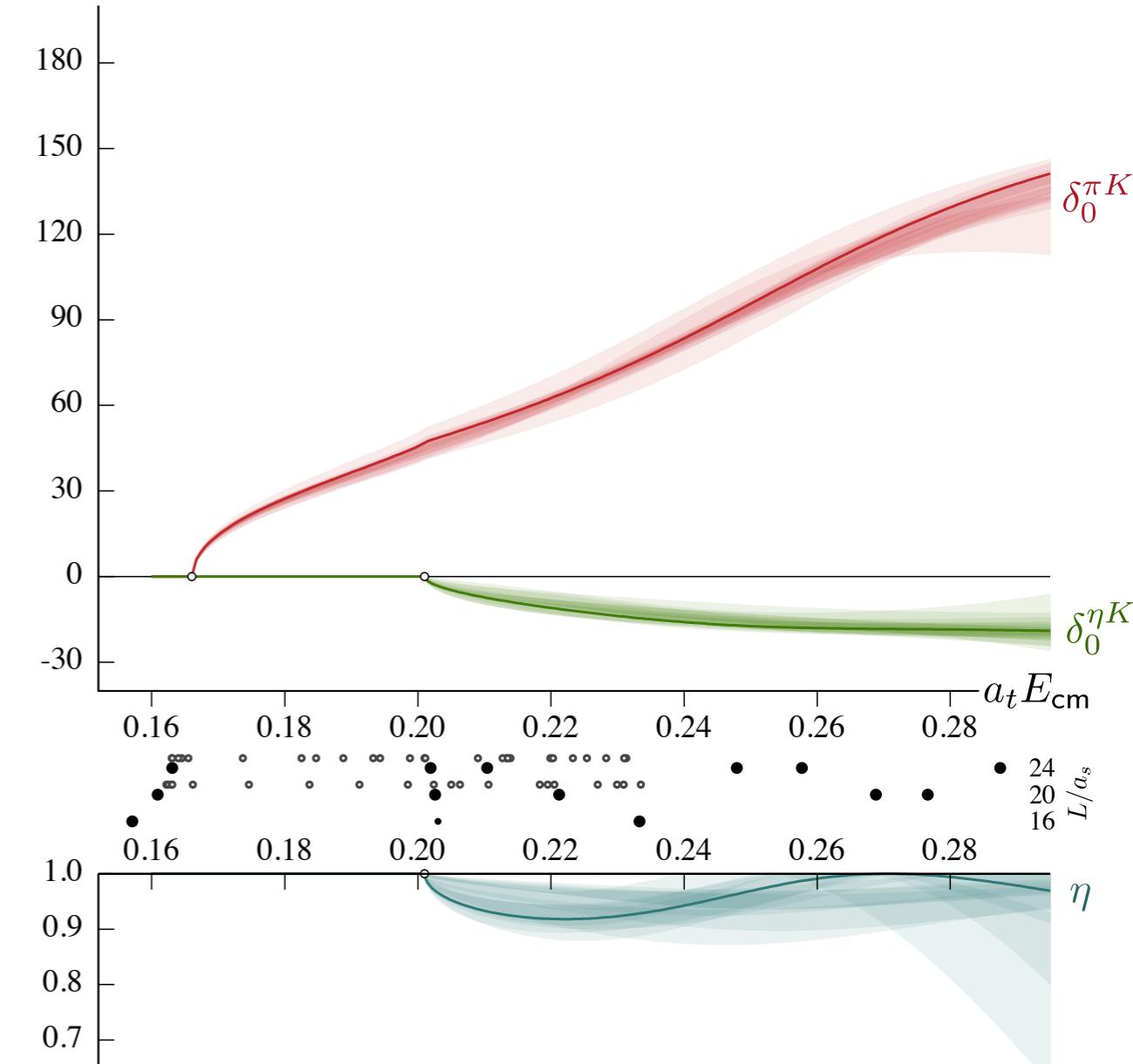
$\pi K/\eta K$ scattering

$m_\pi \sim 391$ MeV

- Are the result parameterization dependent ?
 - Try a range of parameterizations ...



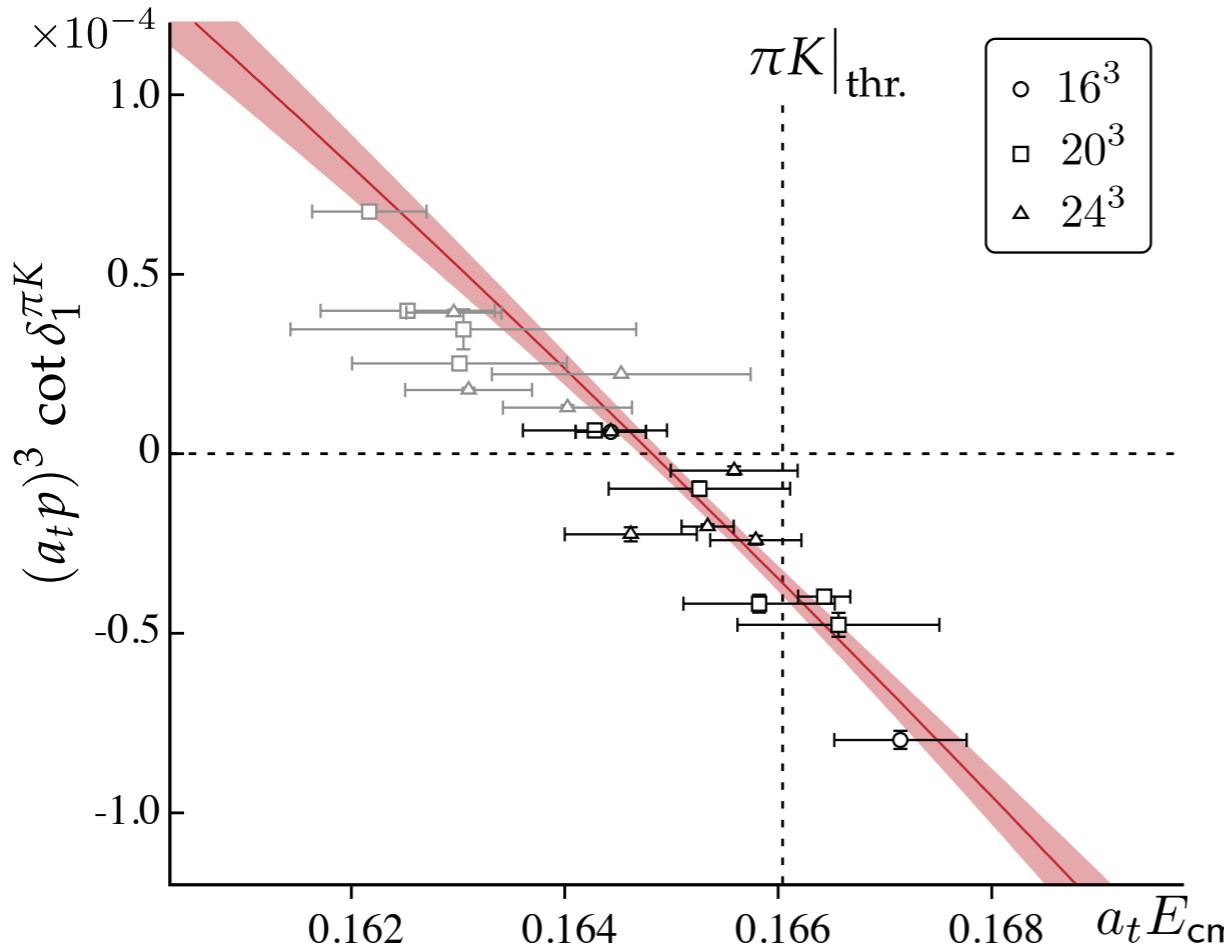
S-WAVE $\pi K/\eta K$ SCATTERING



– gross features are robust

$\pi K/\eta K$ scattering amplitudes

P-WAVE πK SCATTERING



Use a Breit-Wigner with
a subthreshold mass

$$k^3 \cot \delta_1 = (m_R^2 - s) \frac{6\pi\sqrt{s}}{g_R^2}$$

$$\rightarrow g = 5.93(30)$$

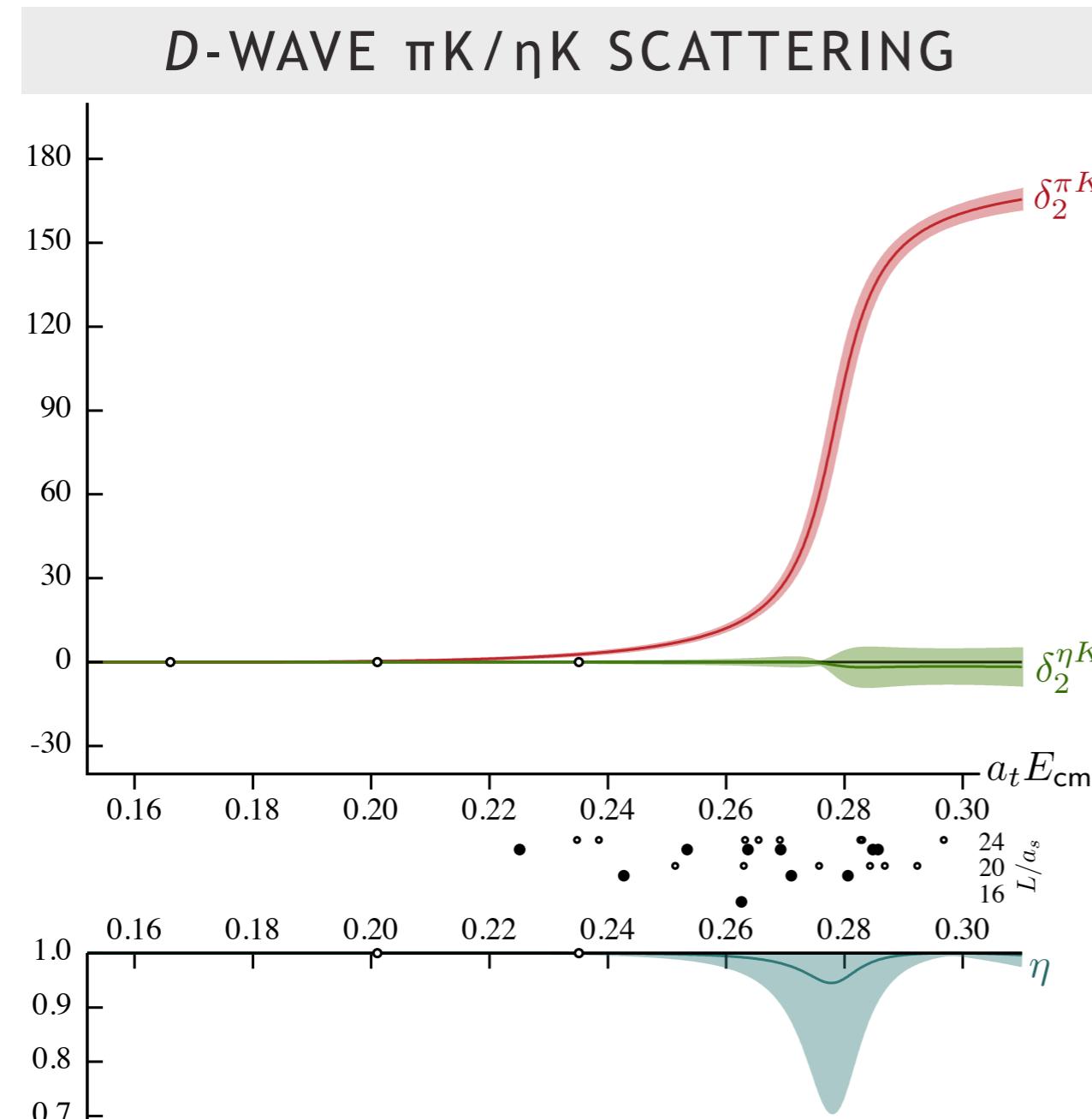
Vector (shallow) bound-state

Quark mass accident that it lies
so close to threshold ...

$g_{\text{phys.}} = 5.5(2)$ PDG

$\pi K/\eta K$ scattering

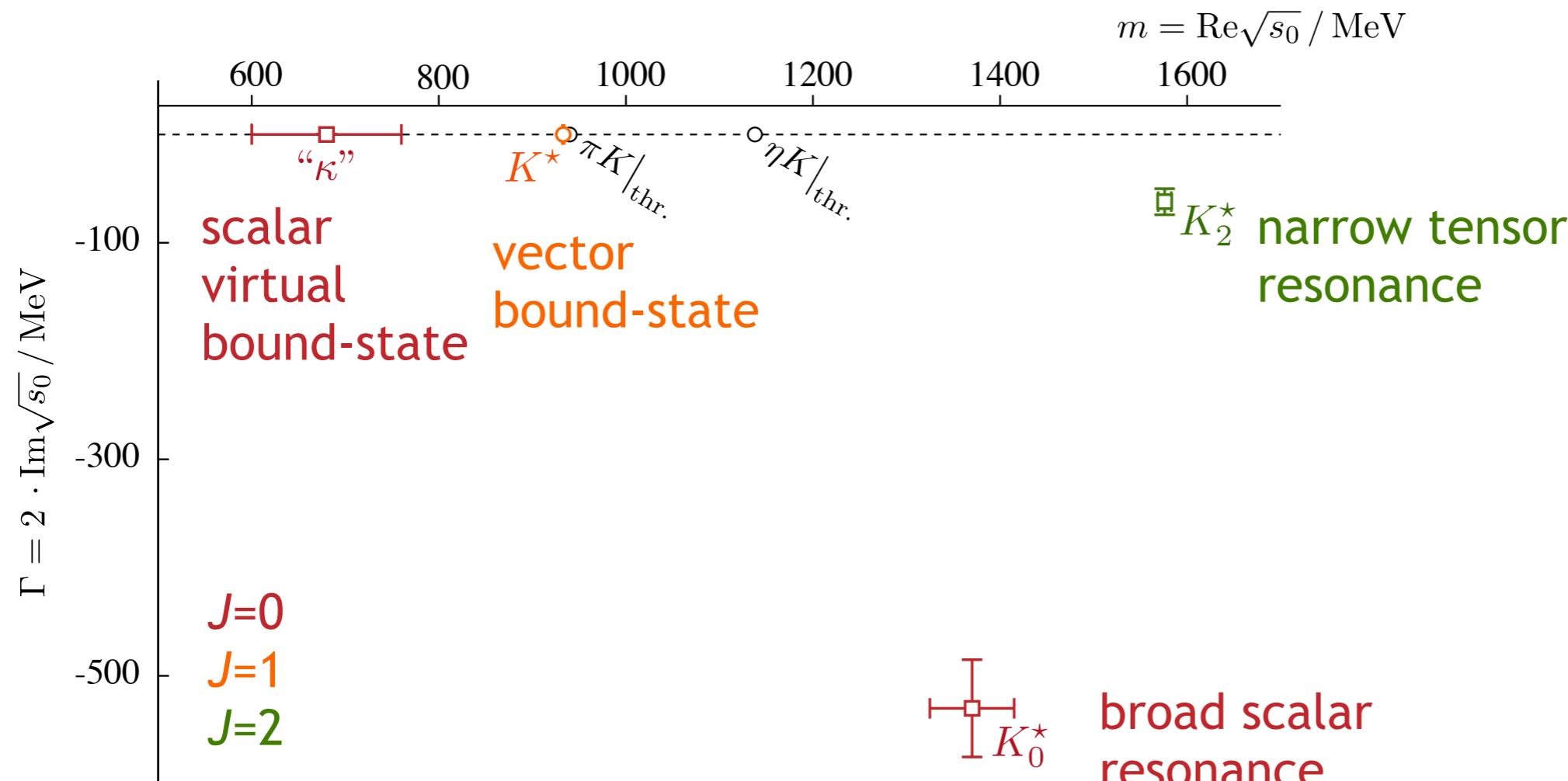
- Clear narrow resonance in D -wave scattering



$m_\pi \sim 391 \text{ MeV}$

Singularity content

- extract t -matrix poles from partial waves

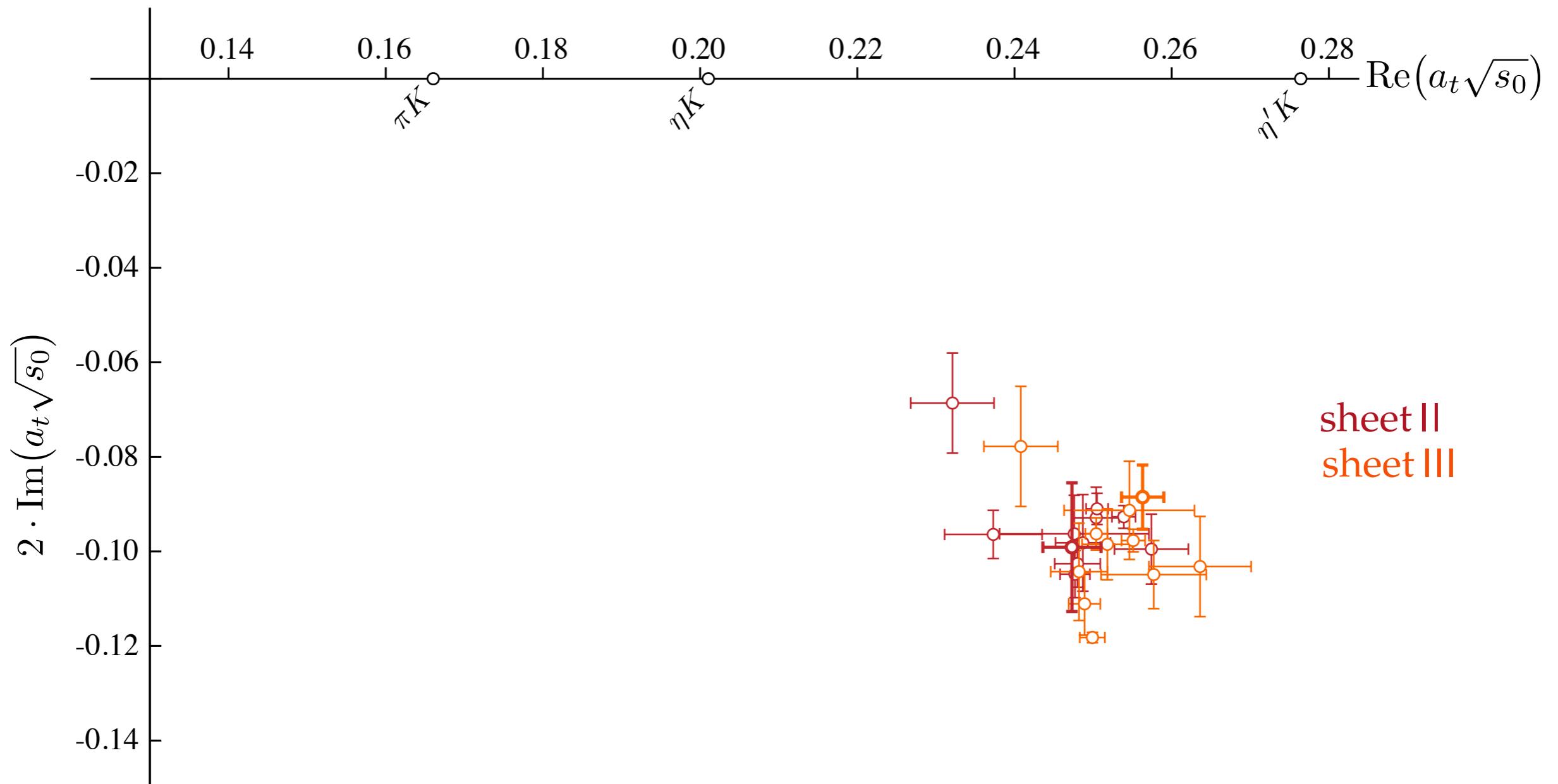


$m_\pi \sim 391 \text{ MeV}$

PRL 113 182001
PRD 91 054008

$\pi K/\eta K$ S-wave resonance

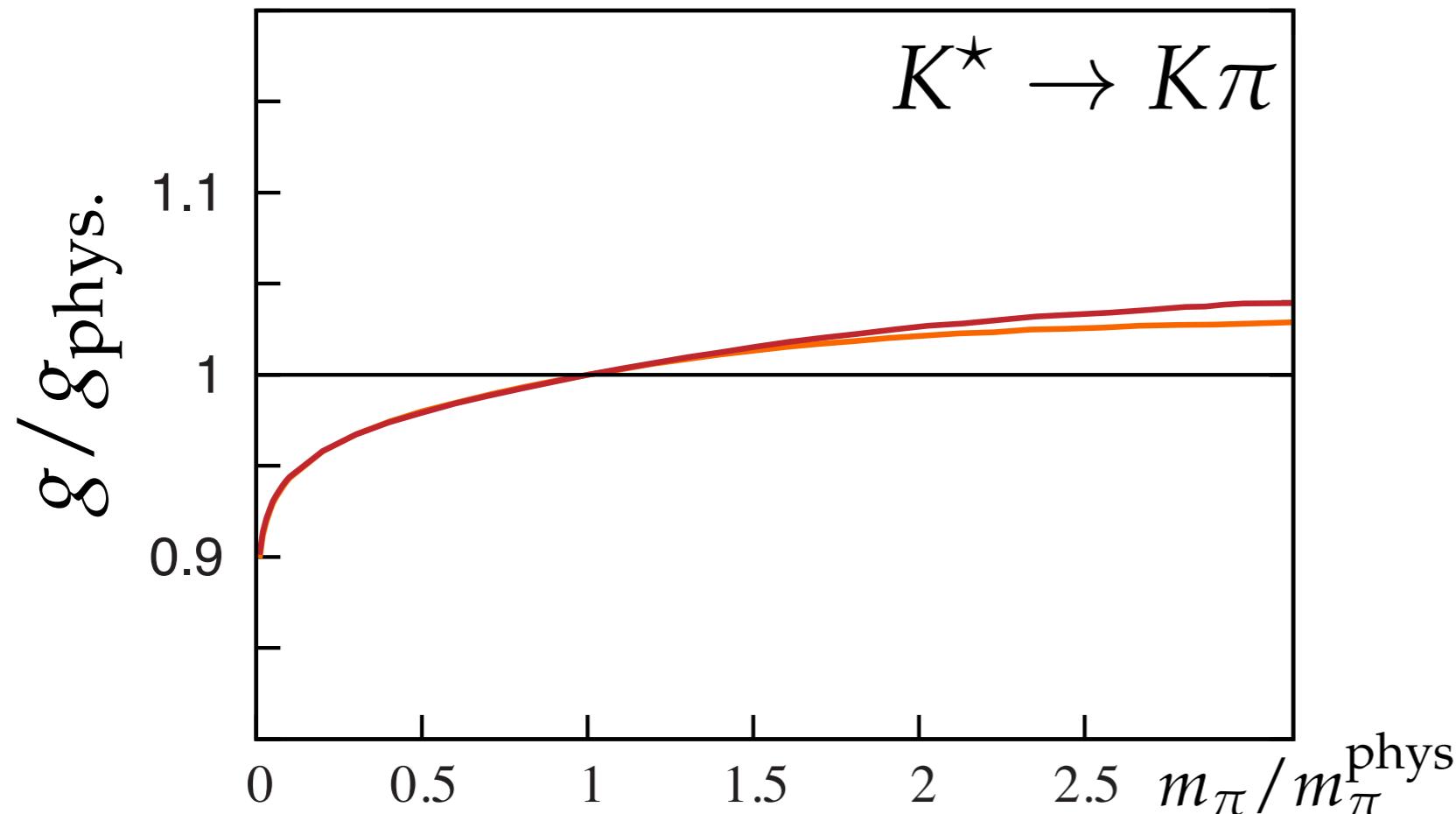
- t -matrix pole position under variation of parameterization



$K^*\pi K$ coupling with changing quark mass

- Unitarized $SU(3)_F$ chiral perturbation theory

NEBREDA & PELAEZ
PRD81 054035 (2010)

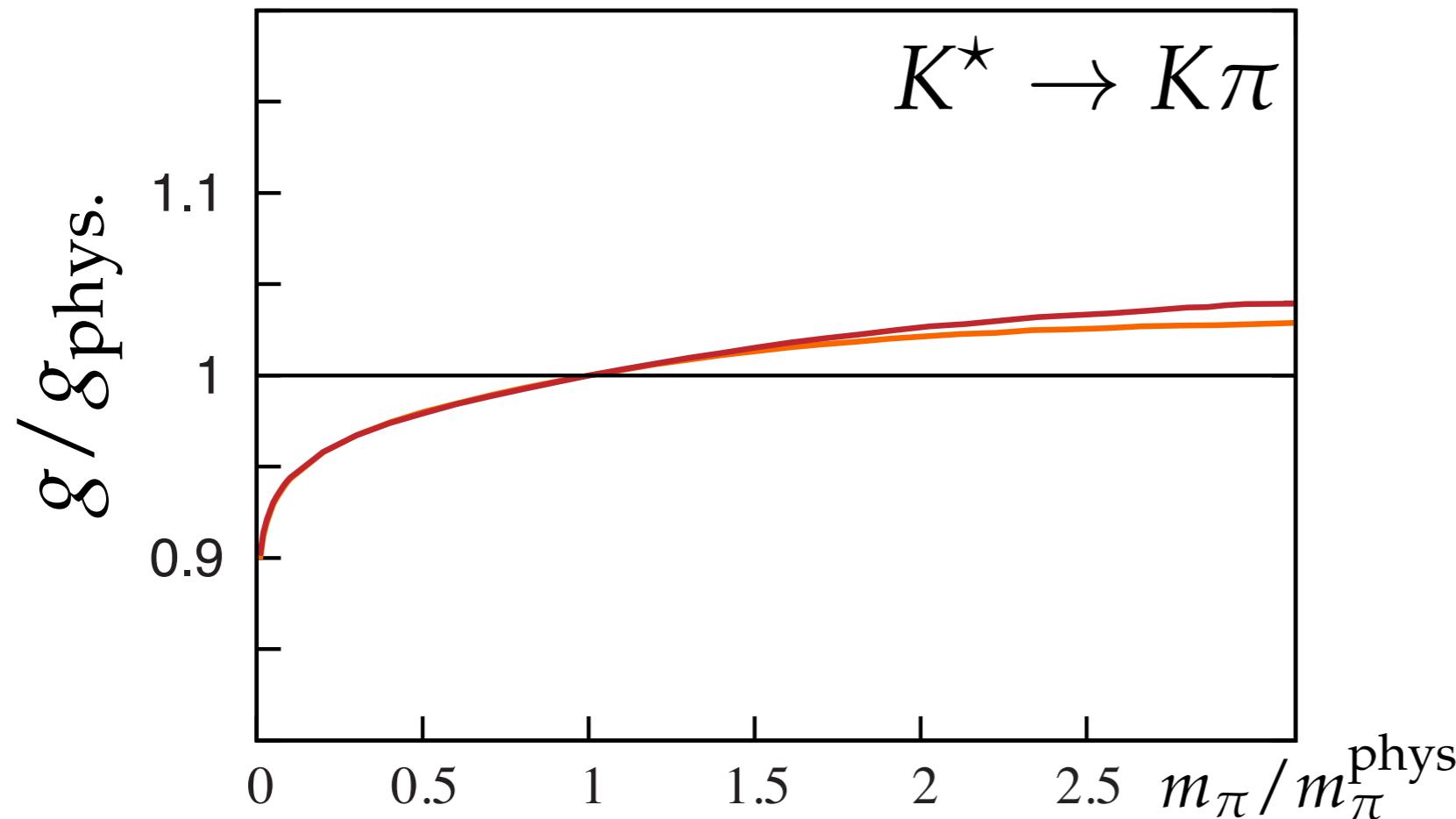


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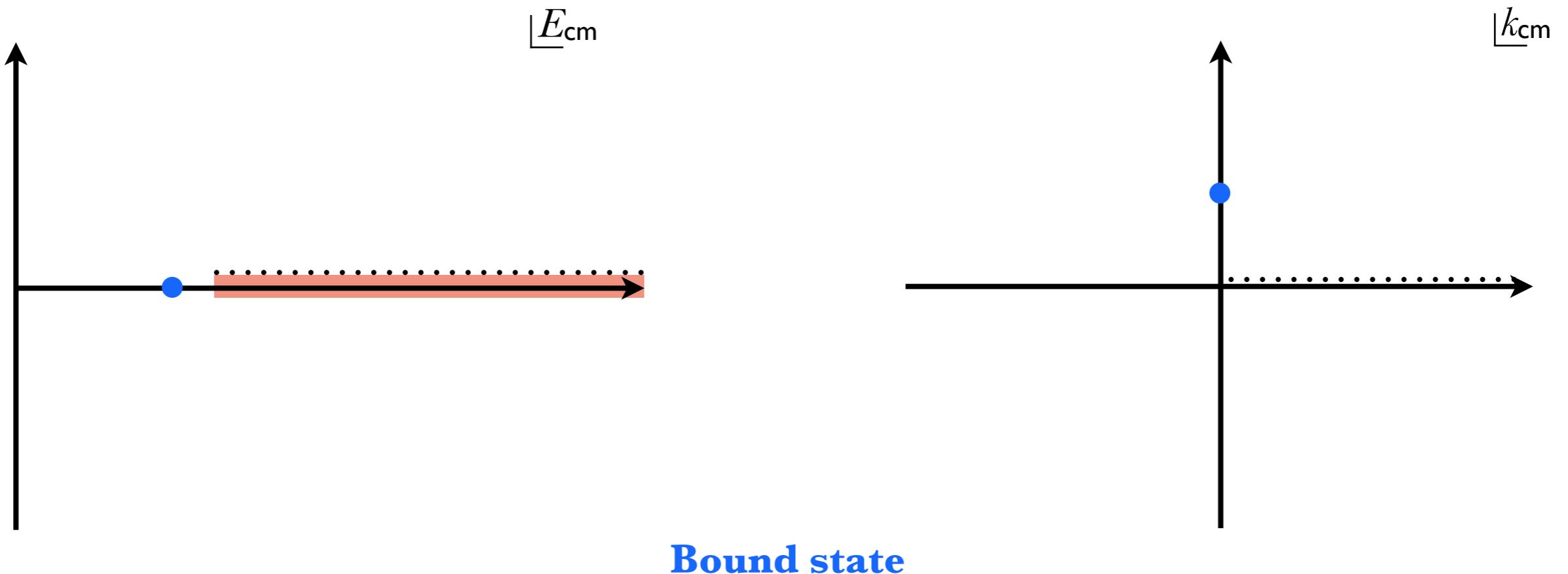
At $m_\pi \sim 391$ MeV we have shallow bound state: $g = 5.93 \pm 0.26$

That's the P-wave - what about S-wave?

Varieties of poles

- Multi-sheet structure around a cut: single channel case

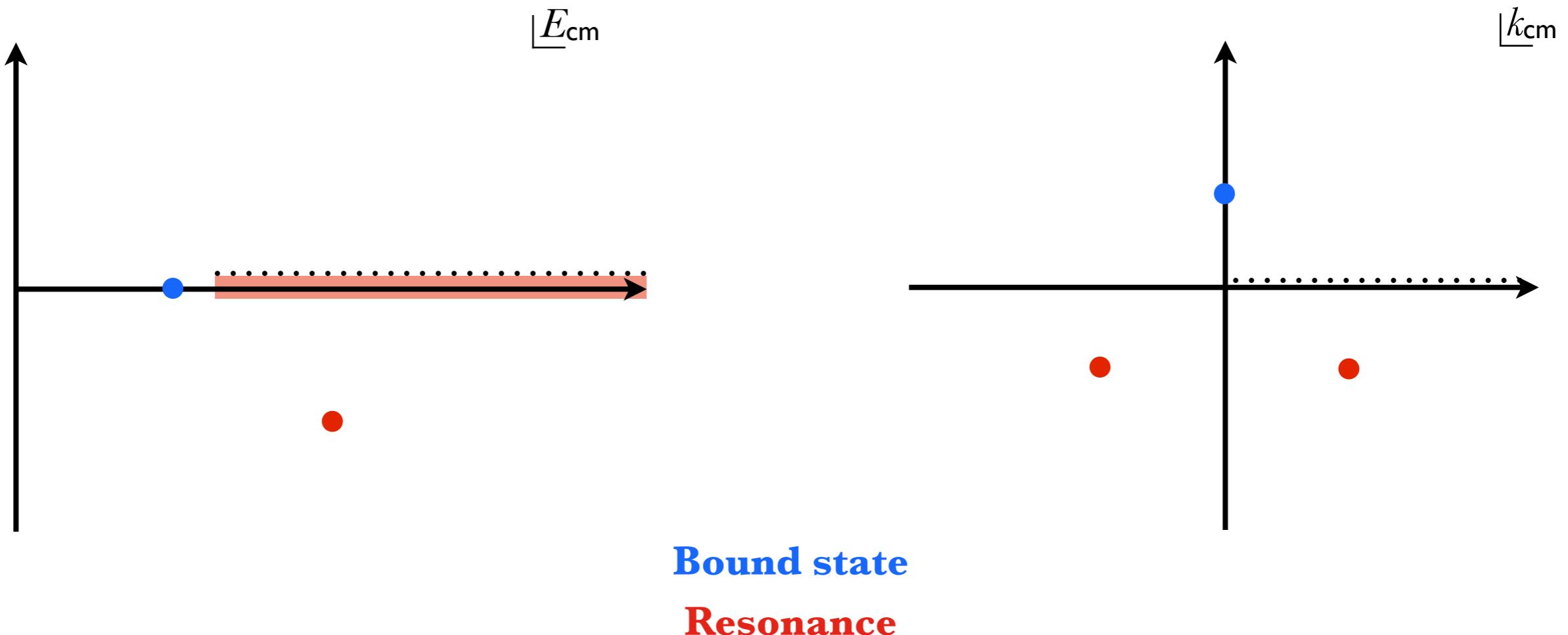
$$k_{cm} = \pm \frac{1}{2} \sqrt{E_{cm}^2 - 4m^2}$$



Varieties of poles

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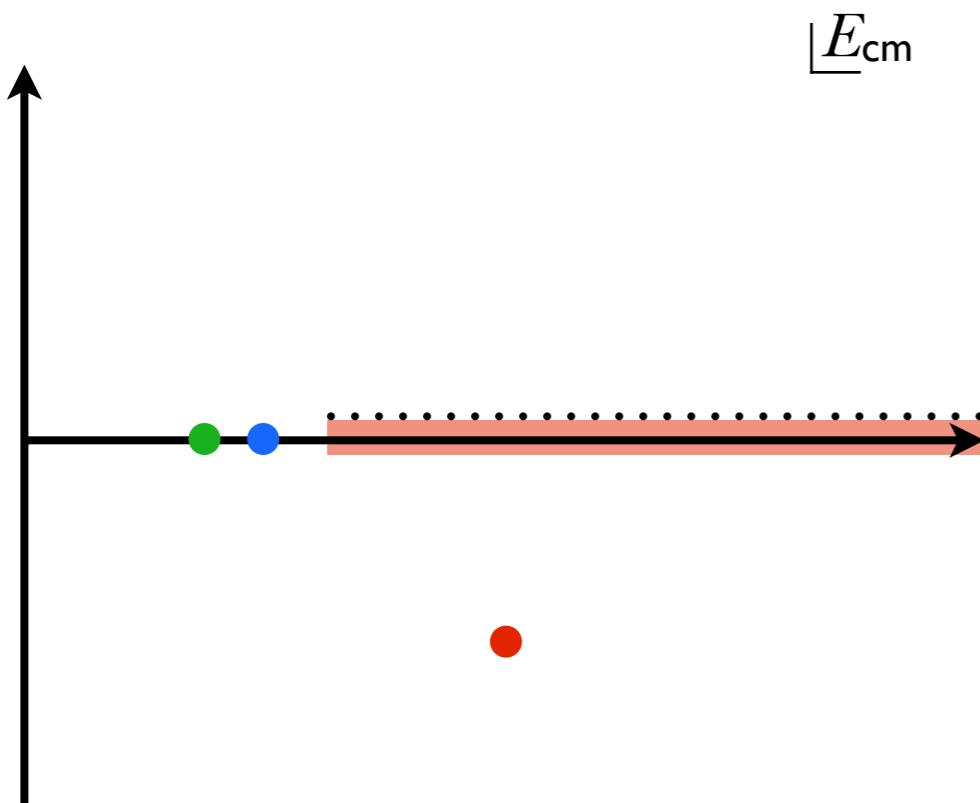
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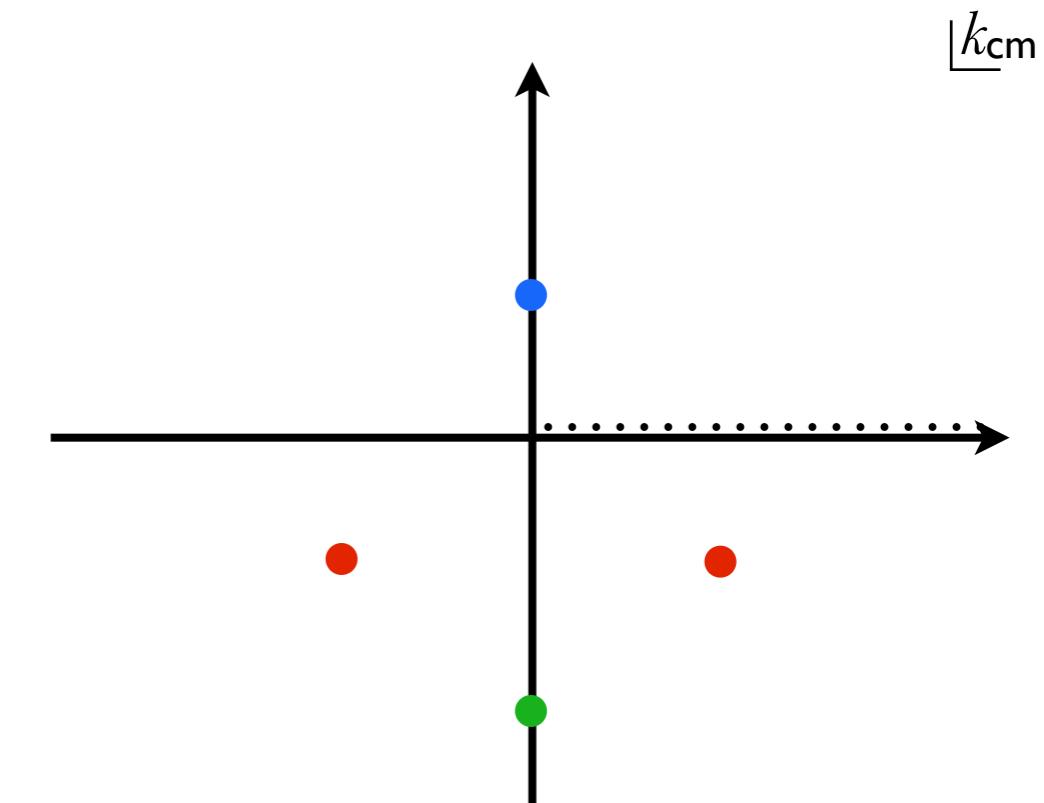
Varieties of poles

- Multi-sheet structure around a cut: single channel case

$$k_{cm} = \pm \frac{1}{2} \sqrt{E_{cm}^2 - 4m^2}$$



Bound state
Resonance
Virtual Bound state



Familiar examples:
N-N 3S_1 deuteron
 ρ
N-N 1S_0

κ (kappa) pole with changing quark mass

- Unitarized SU(3)_F chiral perturbation theory

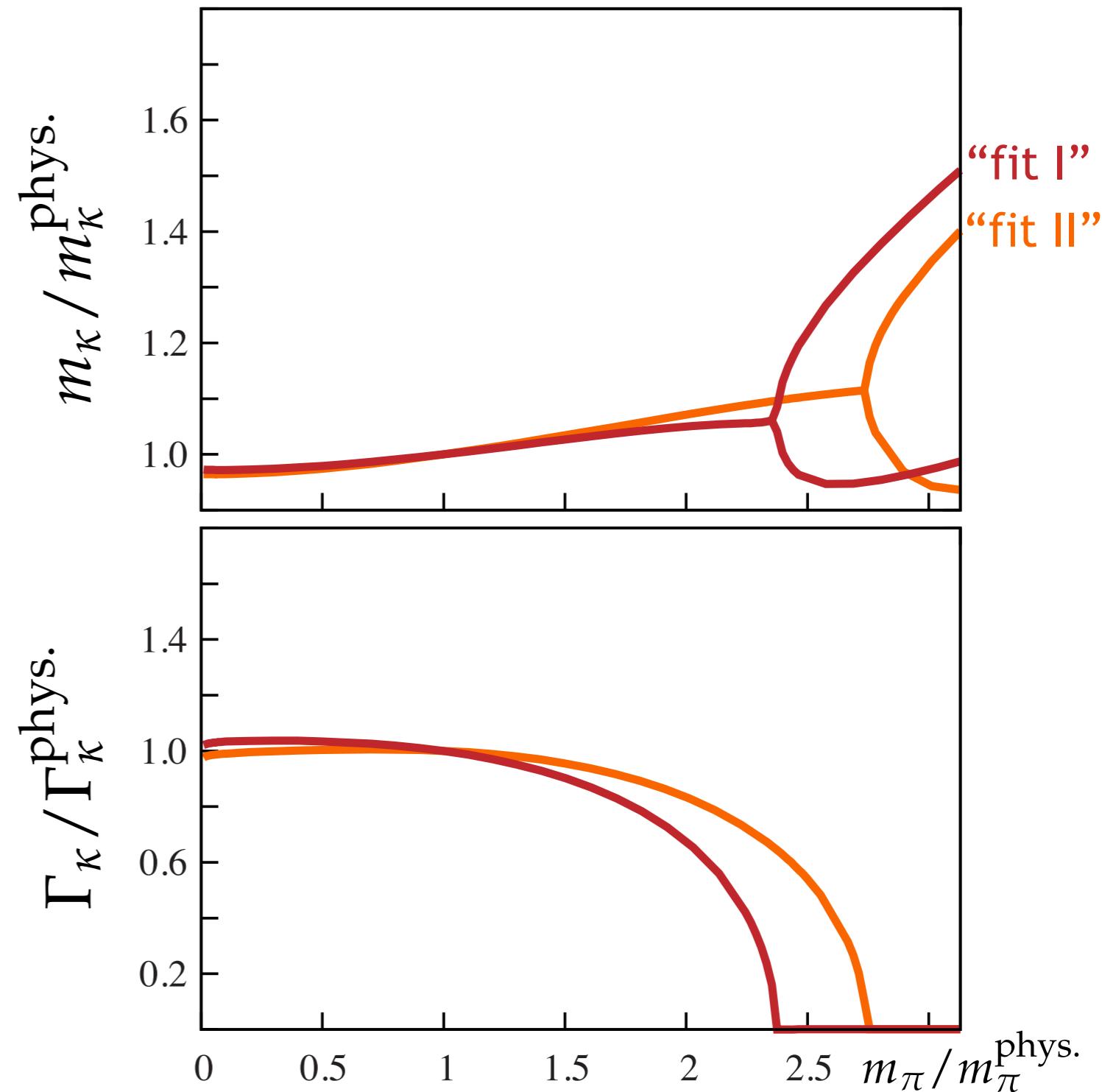
- Resonance poles become virtual bound states somewhere near $m_\pi \sim 2.5 m_\pi^{\text{phys}}$
- At higher pion mass virtual bound-state becomes bound

DESCOTES-GENON

$$\sqrt{s_0} = 660(20) + \frac{i}{2} 550(25) \text{ MeV}$$

NEBREDA & PELAEZ
PRD81 054035 (201)

$$\sqrt{s_0} = m + \frac{i}{2}\Gamma$$



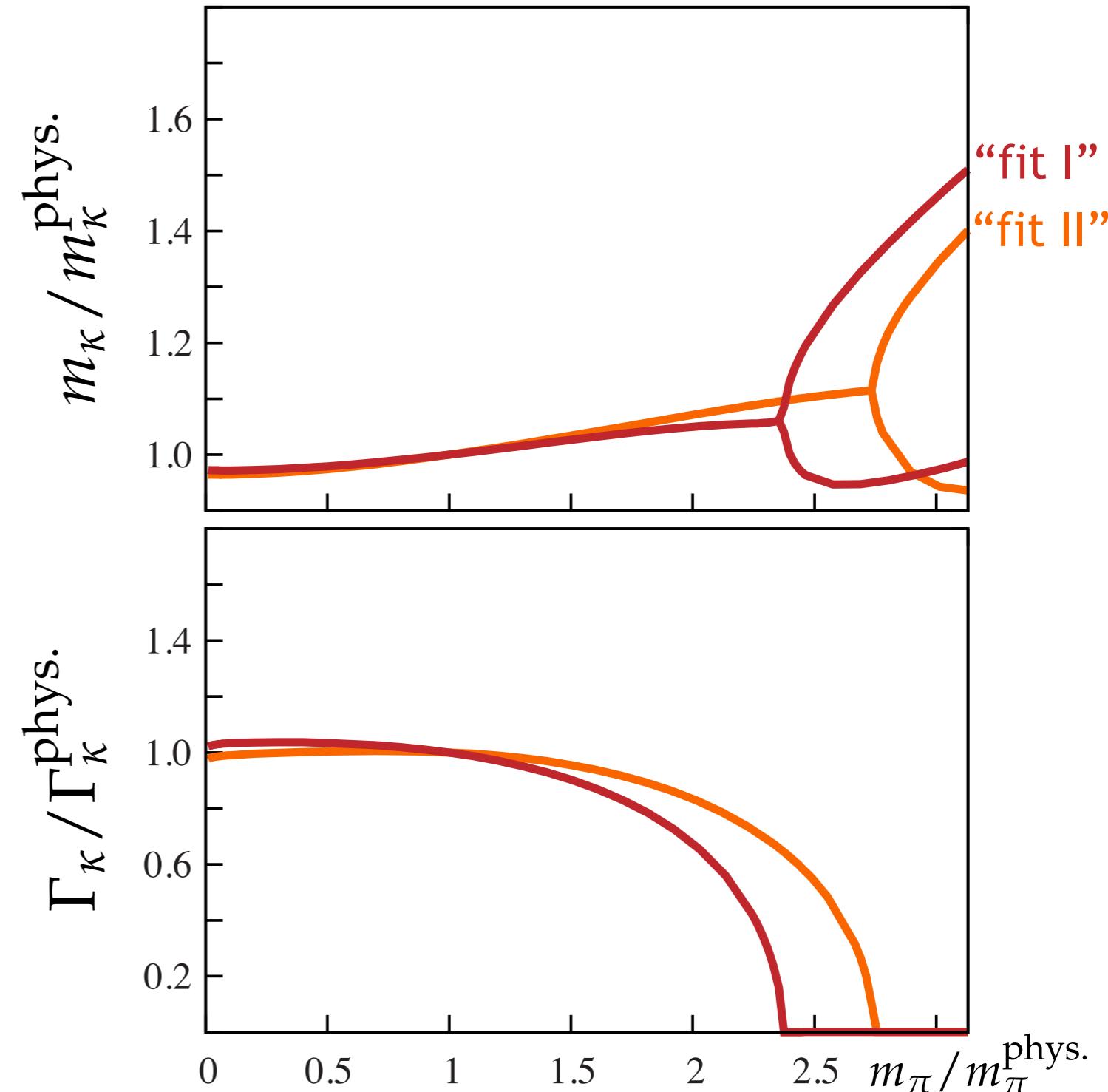
κ (kappa) pole with changing quark mass

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NEBREDA & PELAEZ
PRD81 054035 (201)

$$\sqrt{s_0} = m + \frac{i}{2}\Gamma$$



DESCOTES-GENON

$$\sqrt{s_0} = 660(20) + \frac{i}{2} 550(25) \text{ MeV}$$

At a lower pion mass, do we
see a resonance?
Stay tuned !

Summary

- LQCD spectroscopy program maturing. First phase:
 - With only “single-hadron” operators obtain sketch of hadron spectrum
 - Suggests rich spectrum of mesons & baryons - exotic & non-exotic hybrids

Summary

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 - With only “single-hadron” operators obtain sketch of hadron spectrum
 - Suggests rich spectrum of mesons & baryons - exotic & non-exotic hybrids
- *Goal is to compute resonance information - decays & branching fractions*
 - Including multi-hadron operators leads to richer spectrum
 - Demonstrated viability of finite-volume methods
 - S-matrix formalism increasingly important

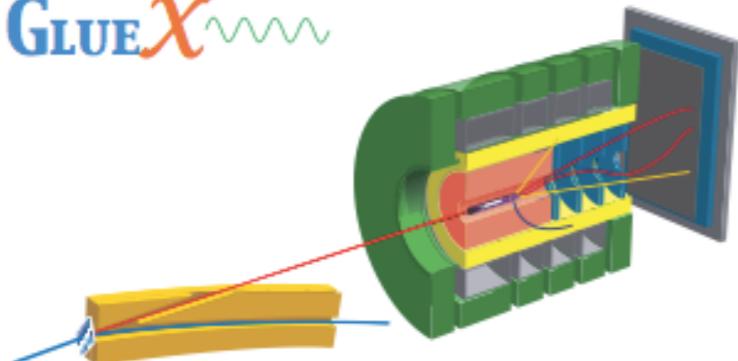
Summary

- LQCD spectroscopy program maturing. First phase:
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- *Goal is to compute resonance information - decays & branching fractions*
 - Including multi-hadron operators leads to richer spectrum
 - Demonstrated viability of finite-volume methods
 - S-matrix formalism increasingly important
- Ultimately, determine underlying structure
 - Besides exotics - disentangle scalar spectrum (in progress)
 - Consider external currents coupled to resonances (in progress)

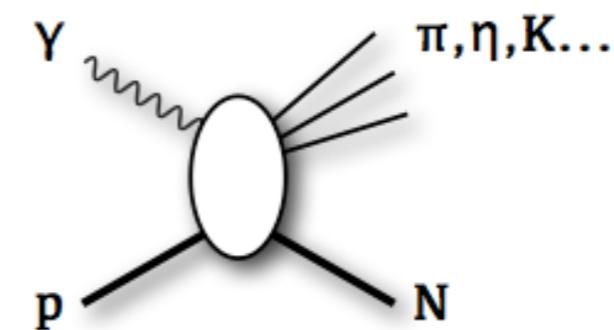
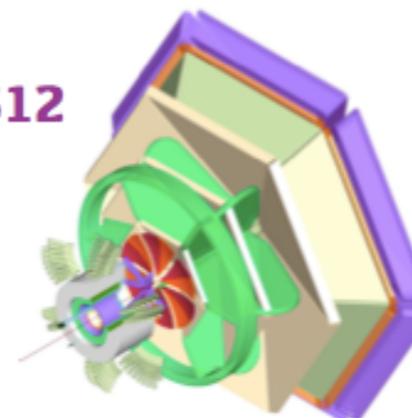
SHULTZ 2015
BRICENO & HANSEN 2015

Spectrum - light meson experiments

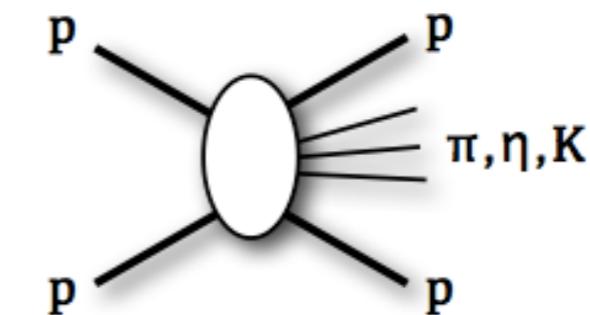
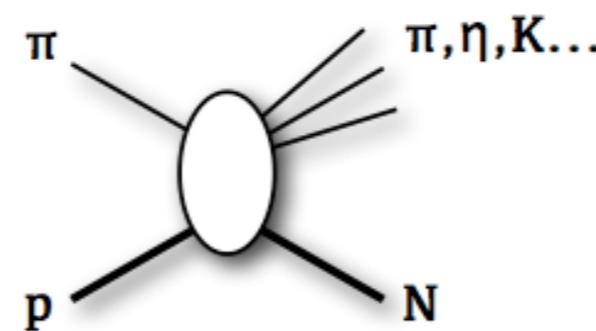
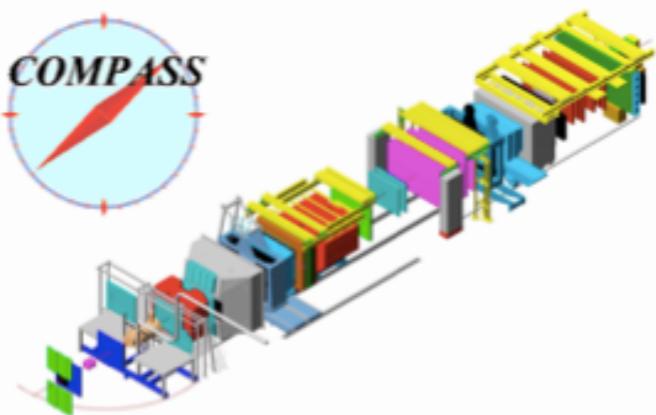
GLUE χ



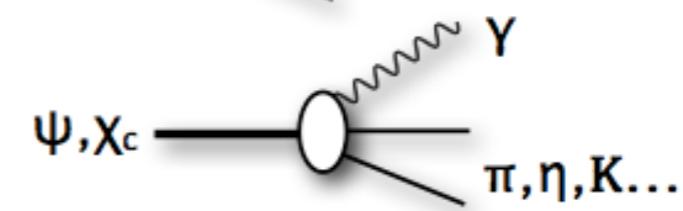
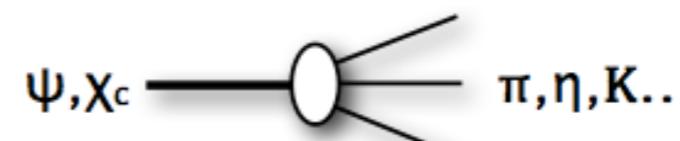
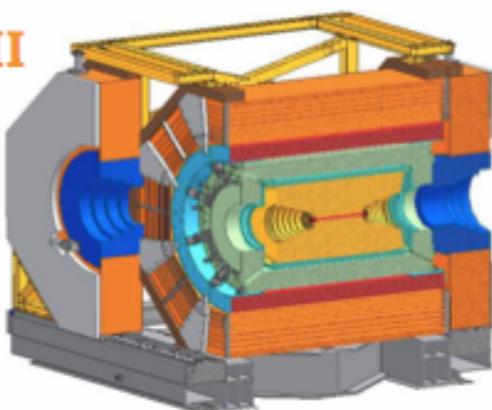
CLAS12



COMPASS



BES III



Scattering with external currents ?

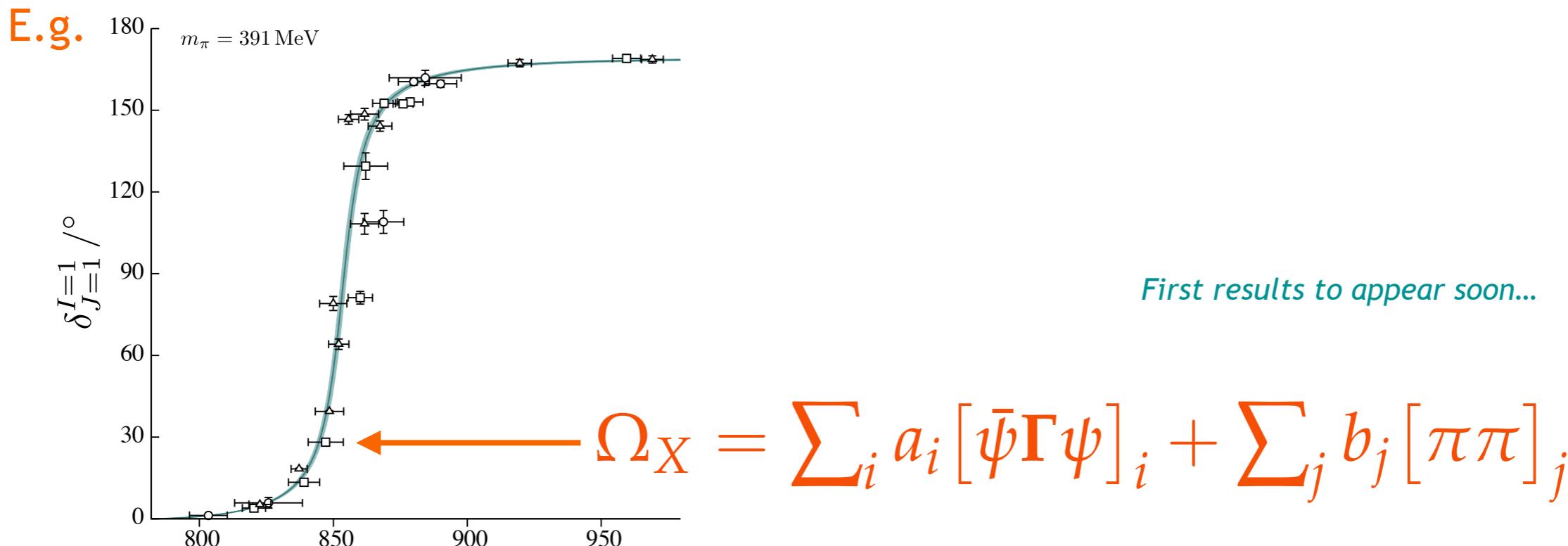
- E.g. $\pi\gamma \rightarrow \pi\pi$ in P -wave : the ρ appears as a resonance
- The observables are the amplitudes $A_\ell(E_{\pi\pi}, Q^2)$
- Can be obtained from correlation functions **in finite-volume**

$$\langle 0 | \Omega_X(t_f) j^\mu(t) \Omega_\pi^\dagger(0) | 0 \rangle$$

PHYSICAL REVIEW D 91, 034501 (2015)

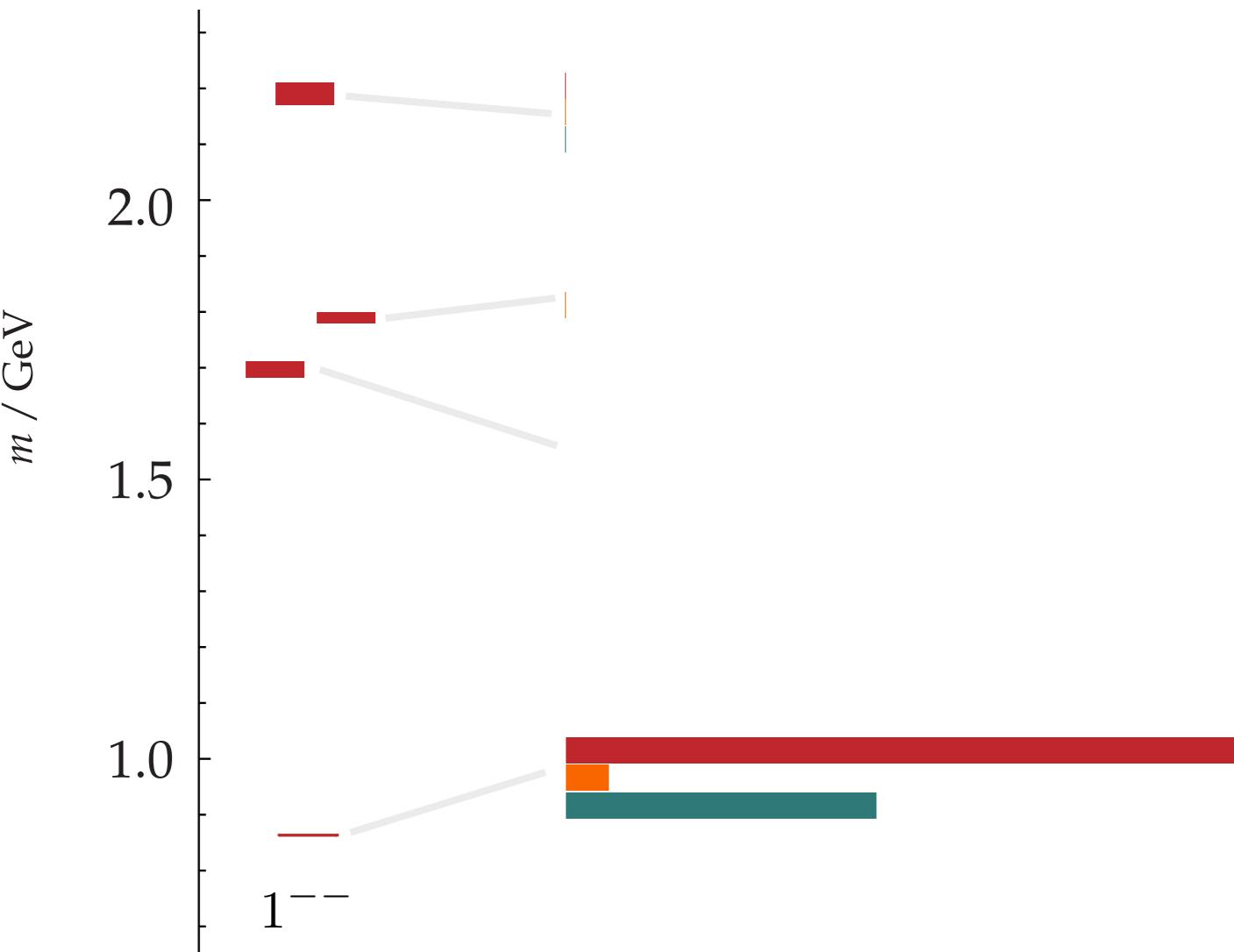
Multichannel $1 \rightarrow 2$ transition amplitudes in a finite volume

Raúl A. Briceño,^{1,*} Maxwell T. Hansen,^{2,†} and André Walker-Loud^{1,3,‡}

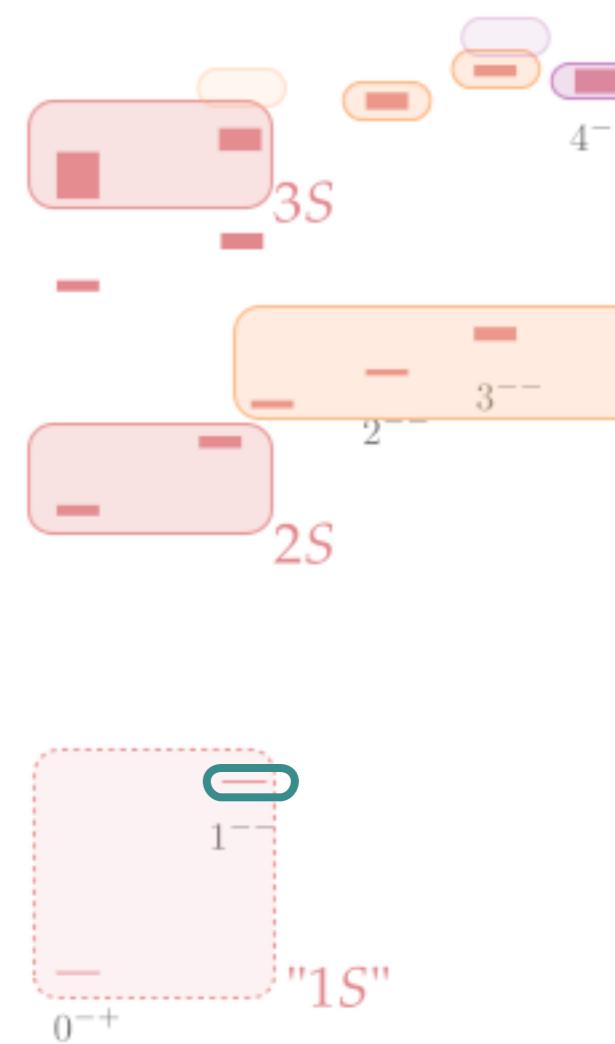


1-- operator overlaps

- Consider the relative size of operator overlaps $\langle \mathbf{n} | \mathcal{O}_i^\dagger | \emptyset \rangle$

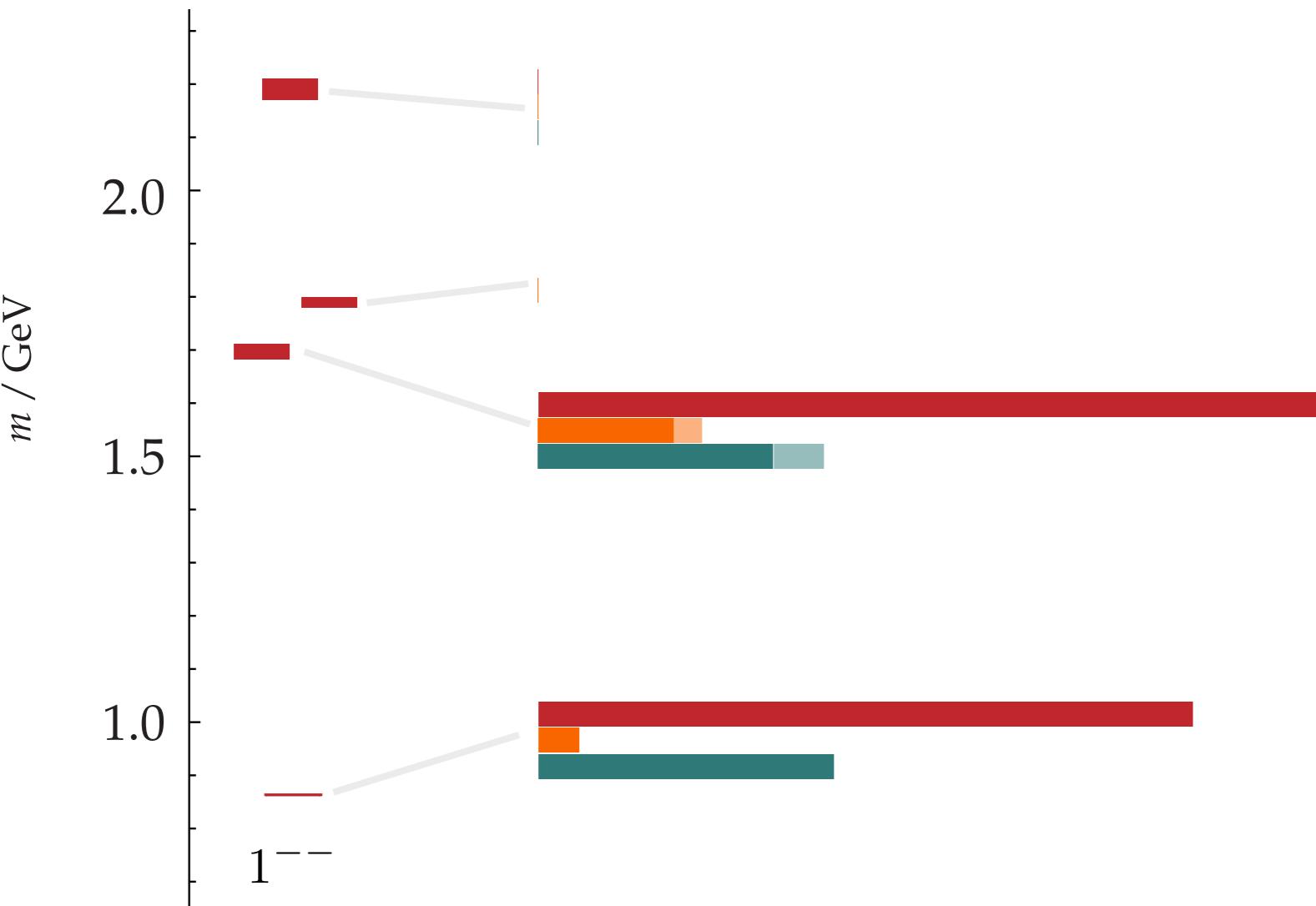


$q\bar{q} [{}^3S_1]$



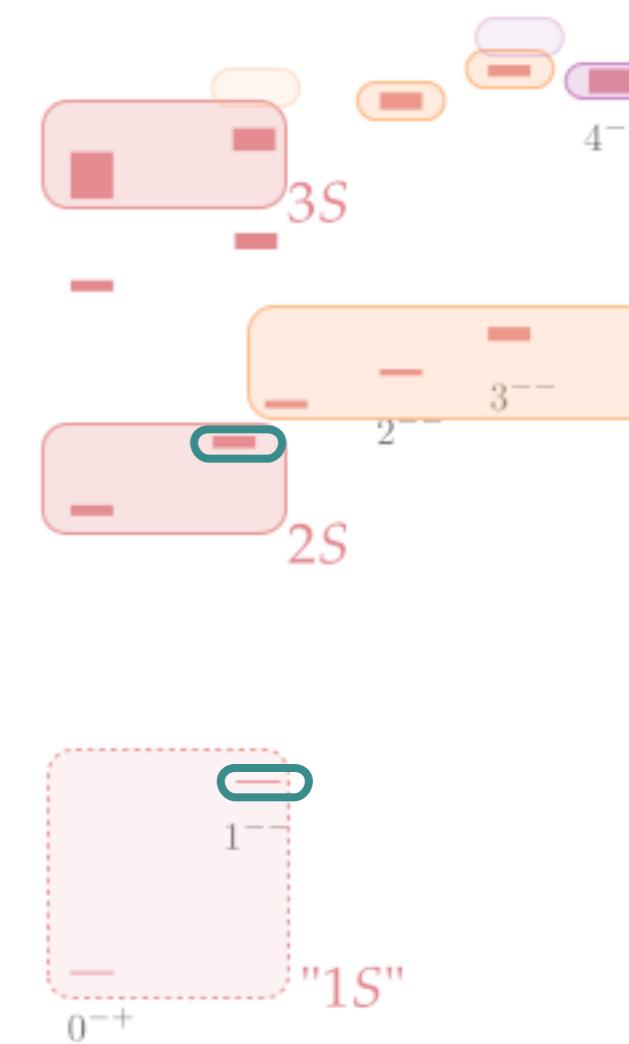
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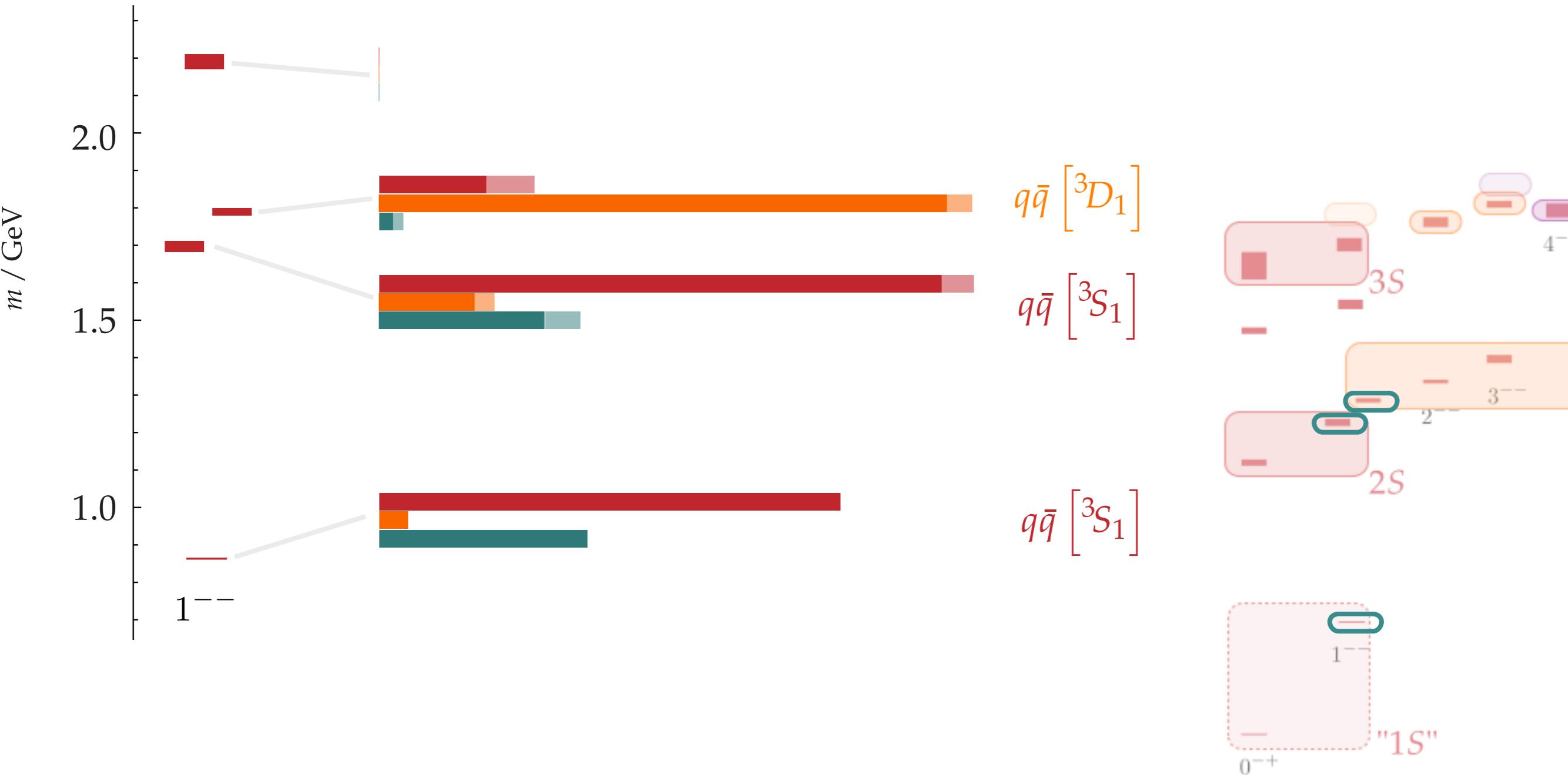
$q\bar{q} [{}^3S_1]$

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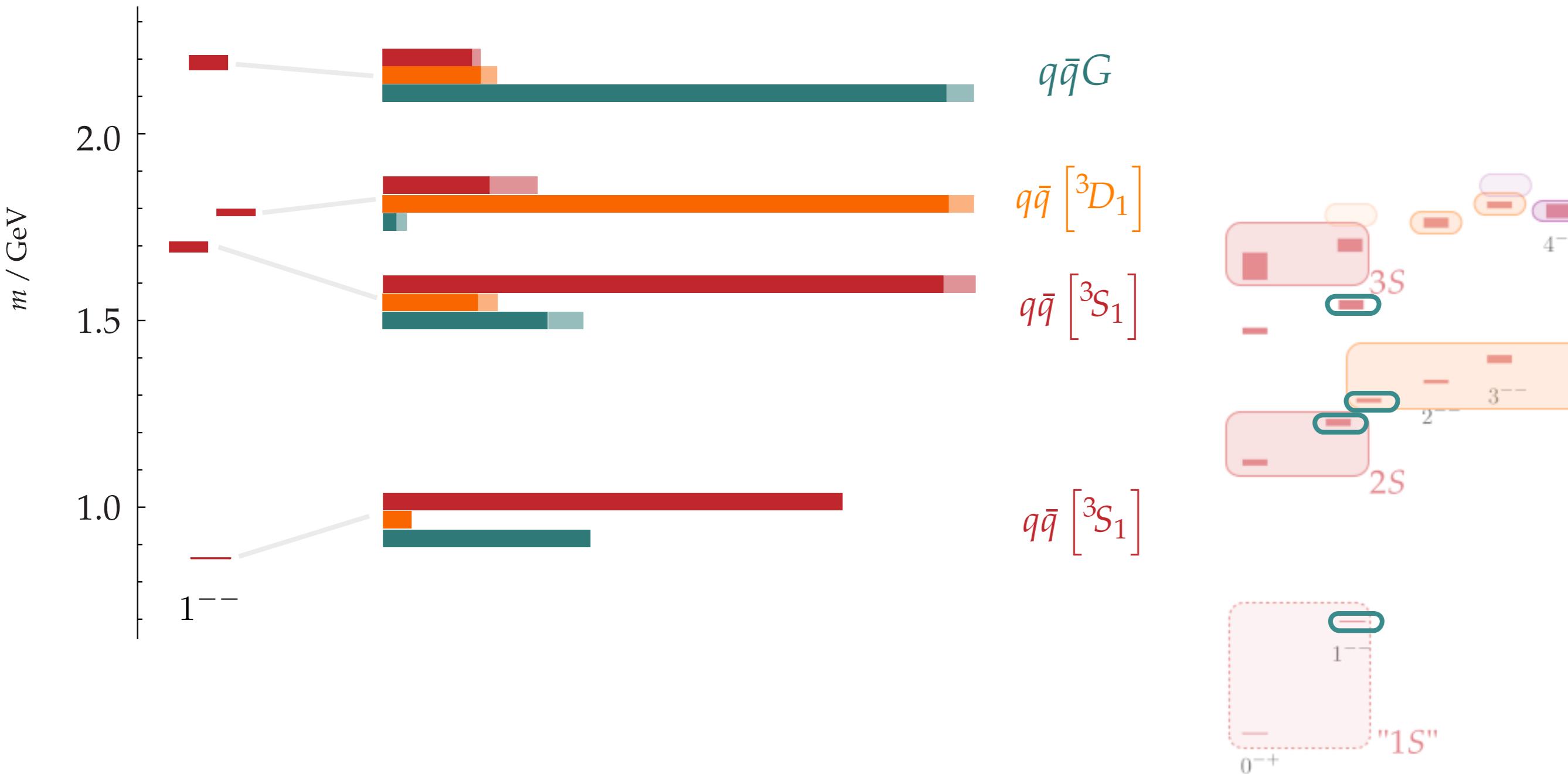
1-- operator overlaps

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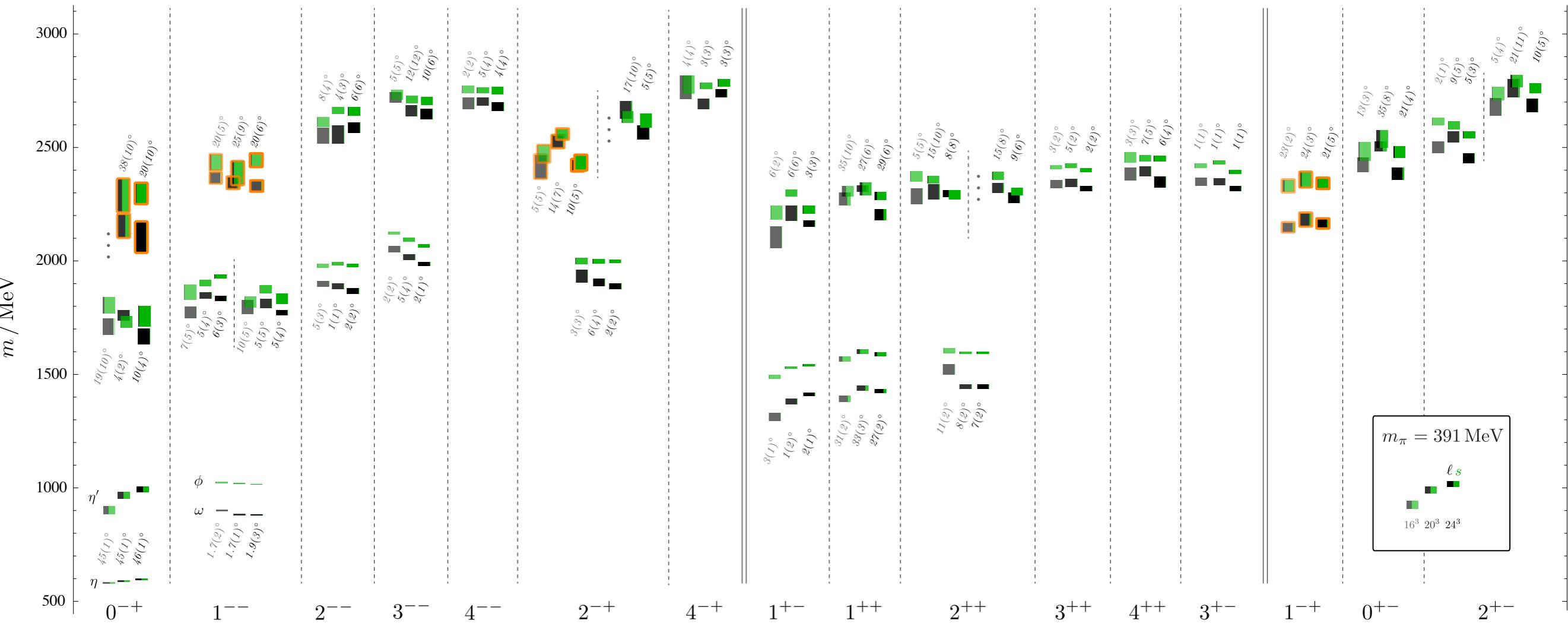
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volume dependence

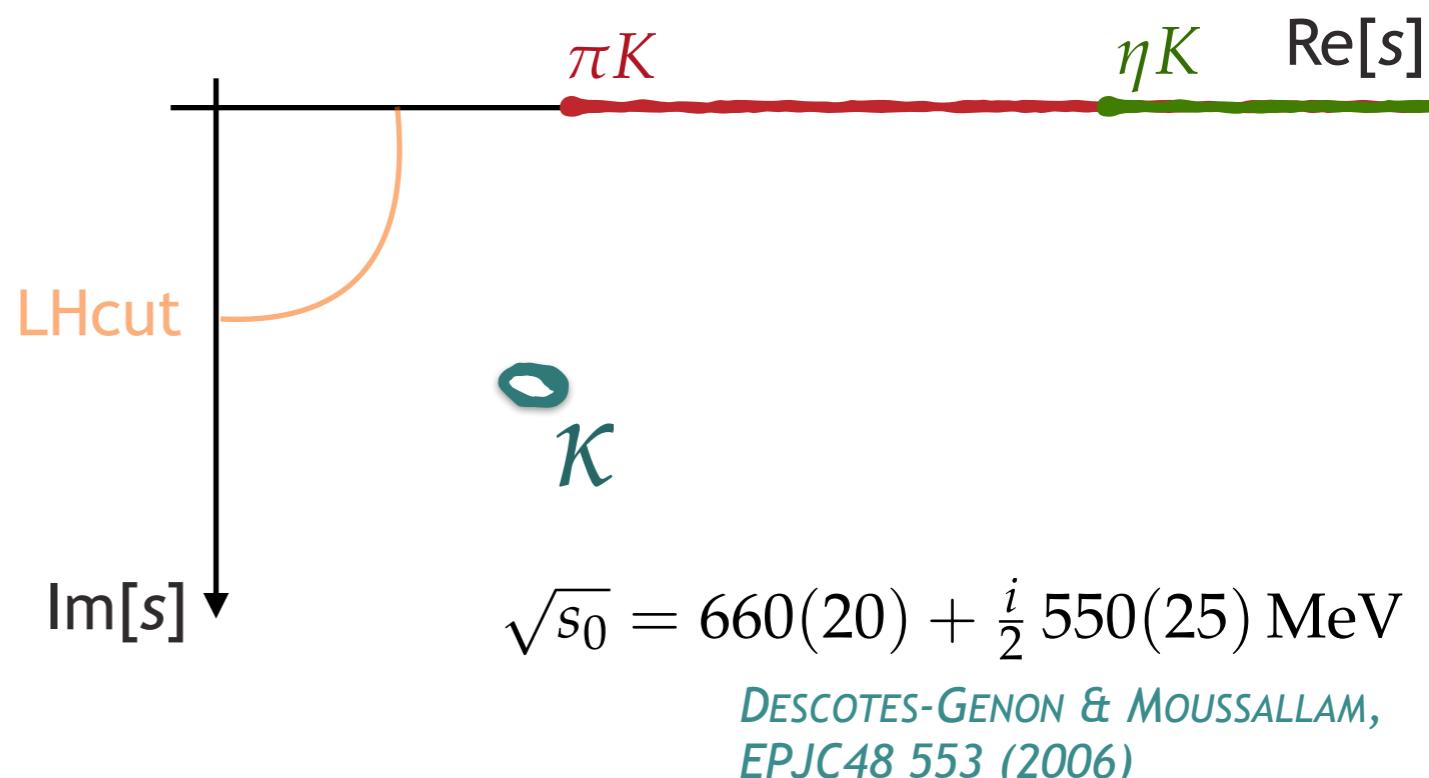
- meson spectrum using large $\bar{\psi}\Gamma D \dots D\psi$ operator basis (no *meson-meson-like*)



little volume dependence - probably picking out narrow resonances

Roy-Steiner & the κ (kappa)

- scattering data: $\pi K \rightarrow \pi K$, $\pi\pi \rightarrow K\bar{K}$
+ analyticity, unitarity, crossing-symmetry
- dispersive representation

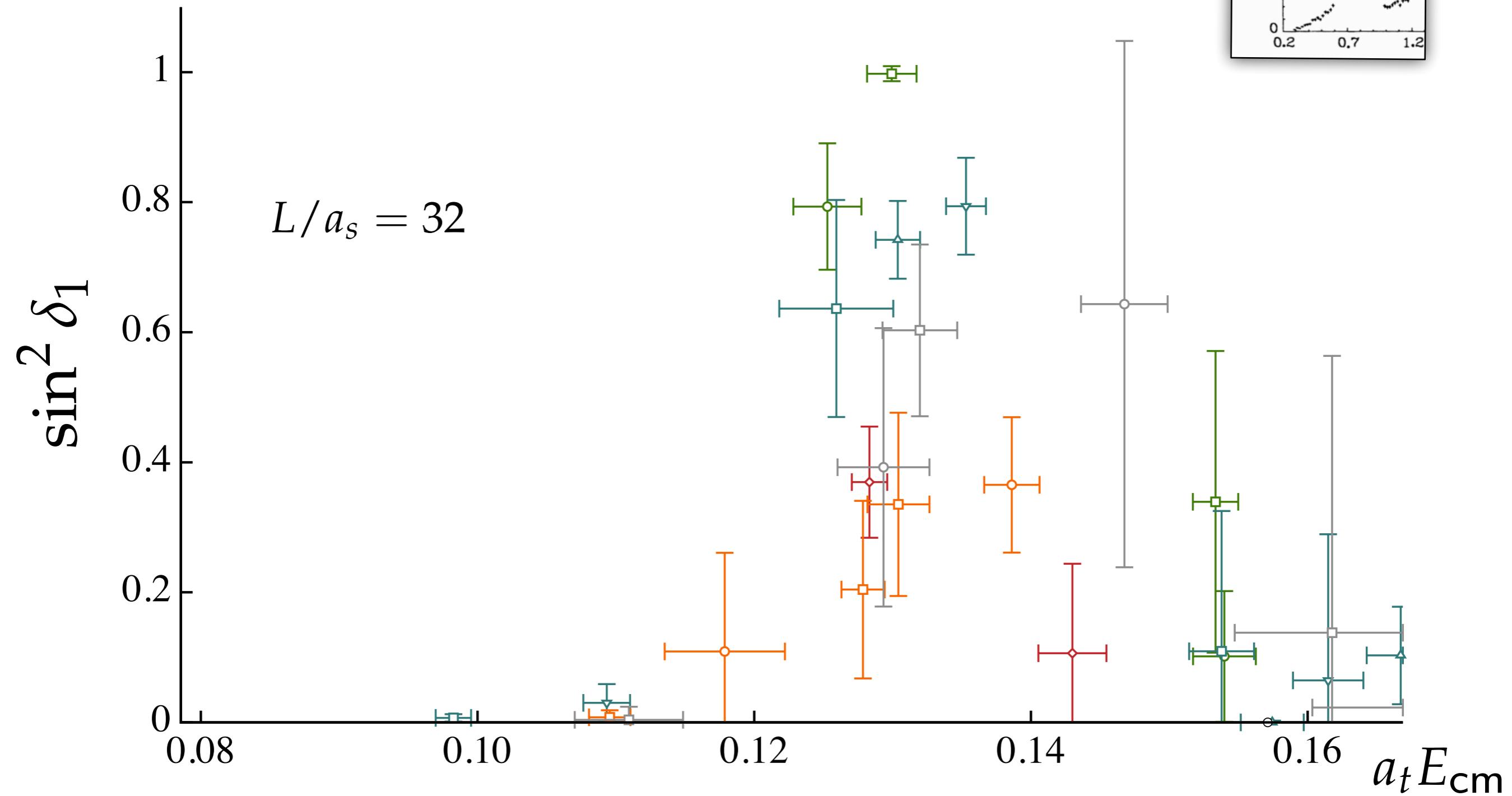
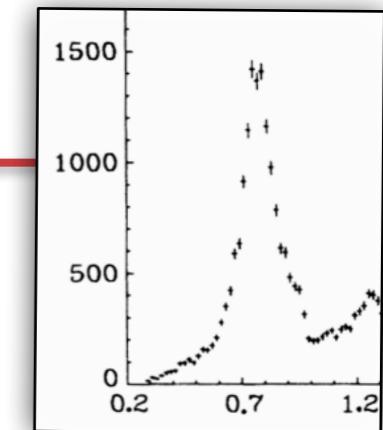


- resonance pole, but deep in the complex plane

ρ resonance - comparison

$m_\pi \sim 230 \text{ MeV}$

Stochastic method

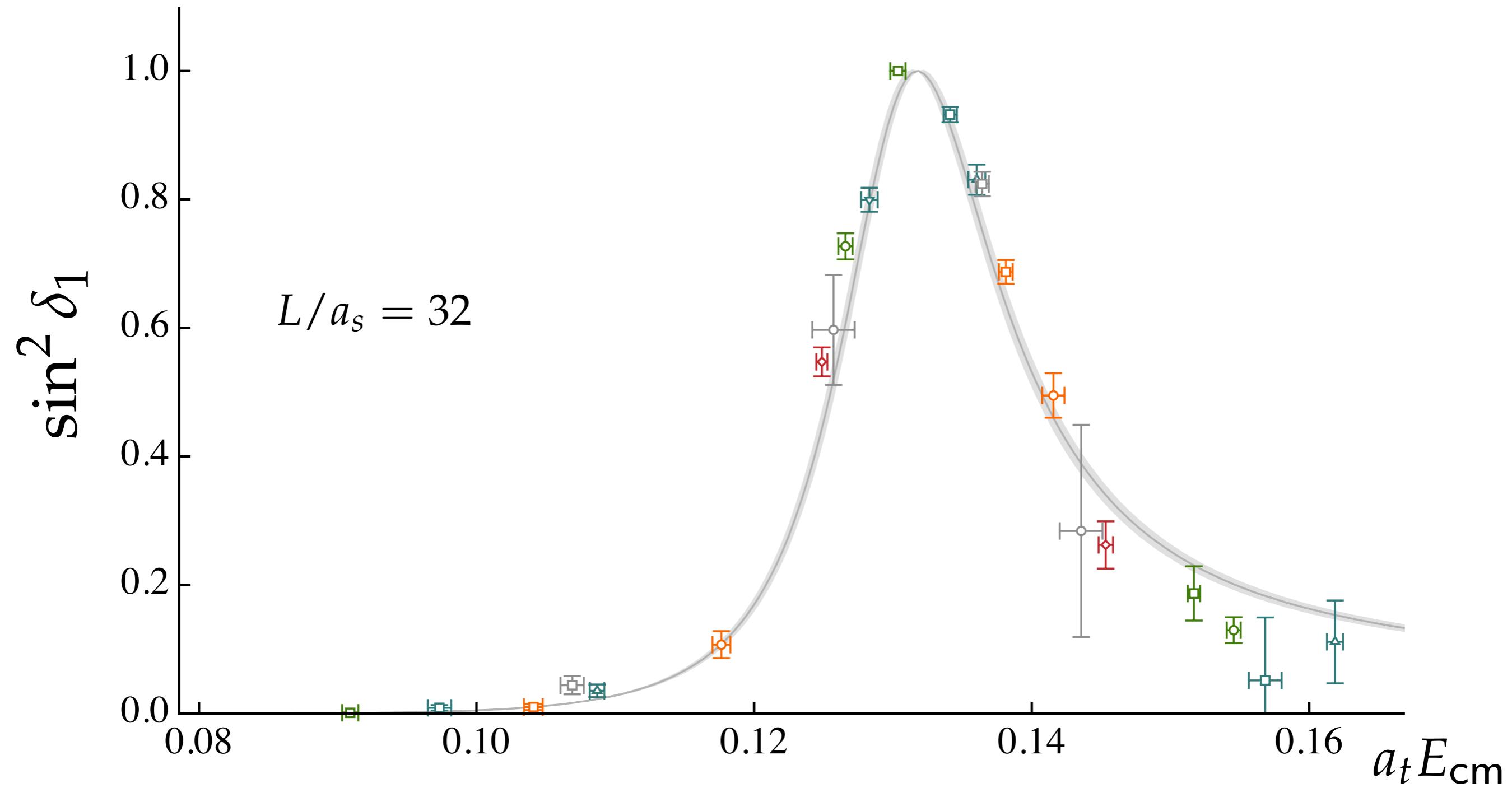


Fahy, et.al. arXiv:1410.8843

ρ resonance - comparison

$m_\pi \sim 230 \text{ MeV}$

Full distillation method



PRELIMINARY ... to appear soon