

# Galactic cosmic rays phenomenology: nuclei, antimatter, dark matter

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NPQCD 1<sup>st</sup> Workshop - Cortona, April 20, 2015



1. What are cosmic rays (CRs)
2. Searches for dark matter in CRs
3. Cross sections for CR phenomenology

# CHARGED (GALACTIC) COSMIC RAYS

are charged particles (nuclei, isotopes, leptons, antiparticles)  
diffusing in the galactic magnetic field

Observed at Earth with  $E \sim 10 \text{ MeV/n} - 10^3 \text{ TeV/n}$

## 1. SOURCES

PRIMARIES: directly produced in their sources

SECONDARIES: produced by spallation reactions of primaries on the interstellar medium (ISM)

## 2. ACCELERATION

SNR are considered the powerhouses for CRs.

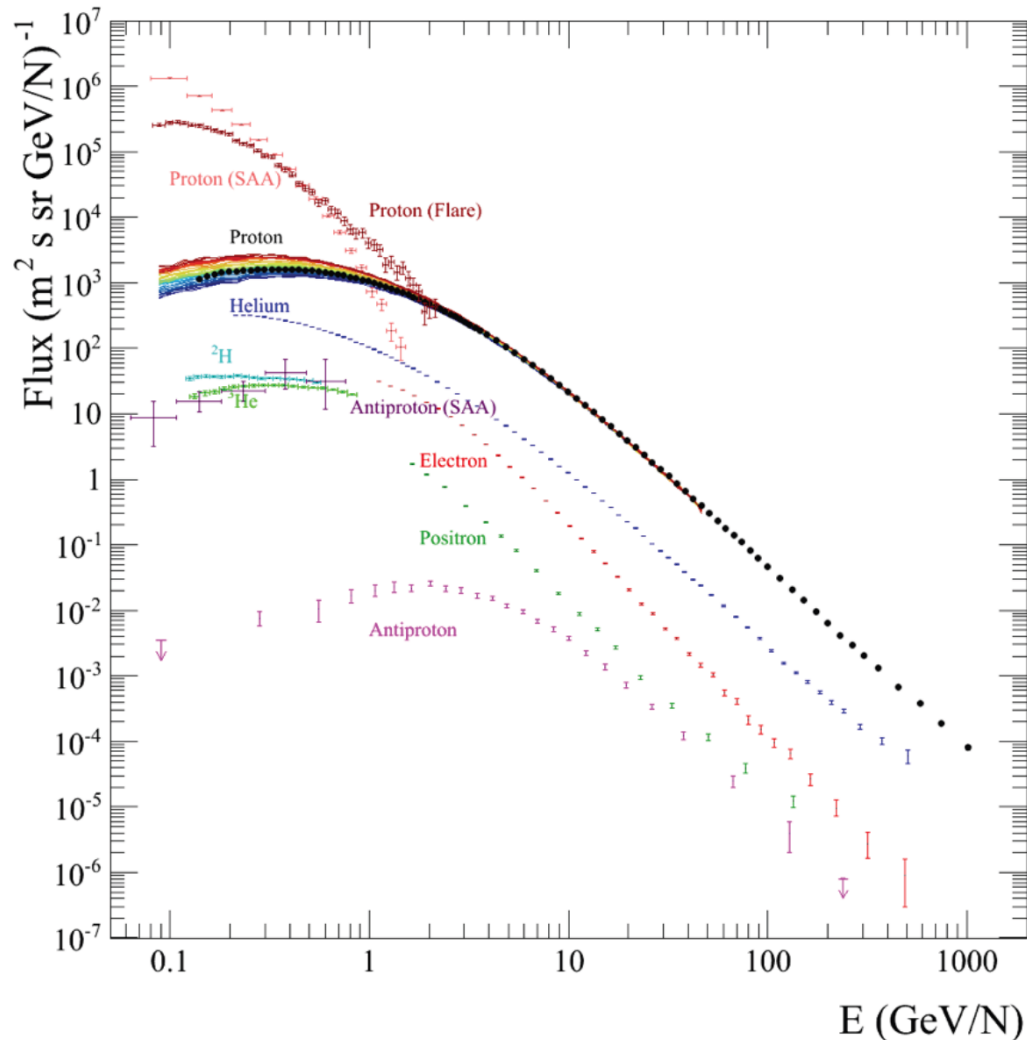
They can accelerate particles at least up to  $10^2 \text{ TeV}$

## 3. PROPAGATION

CRs are diffused in the Galaxy by the inhomogeneities of the galactic magnetic field.

+ loose/gain energy with different mechanisms

# All particle spectra (from Pamela experiment)



Credits: Valerio Formato & Mirko Boezio, Pamela Collaboration, 2013

# Transport equation in diffusion models

$$\begin{aligned}
 & \text{Diffusion} \quad \text{Convection} \\
 & -\vec{\nabla} \left[ K \vec{\nabla} N^j(E) - \vec{V}_c N^j(E) \right] - \Gamma^j N^j \\
 & - \frac{(\vec{\nabla} \cdot \vec{V}_c)}{3} \frac{\partial}{\partial E} \left[ \frac{p^2}{E} N^j(E) \right] = Q^j(E) + \\
 & \frac{\partial}{\partial E} \left[ -b_{\text{tot}}(E) N^j(E) + \beta^2 K_{pp} \frac{\partial N^j(E)}{\partial E} \right] \\
 & \text{Energy losses (EM)} \quad \text{Reacceleration}
 \end{aligned}$$

**Destruction on ISM**

CR sources: primaries,  
**secondaries**  
(spallations)

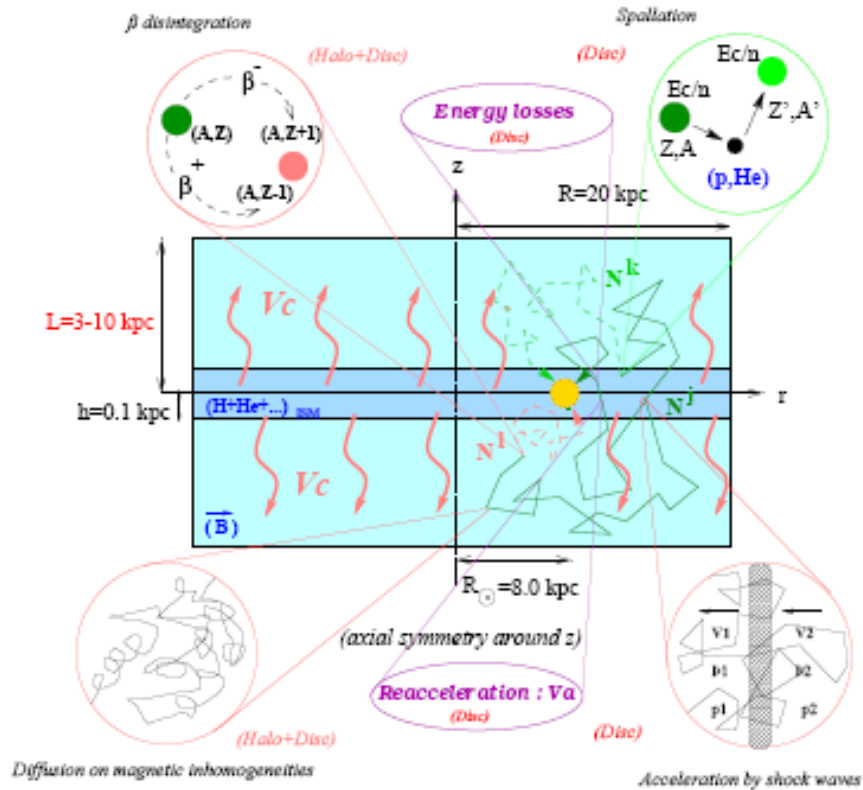
1. **DESTRUCTION:**  $\Gamma = n_{\text{ISM}} v \sigma_R$ ,  $\sigma_R = \sigma^{\text{tot}} - \sigma^{\text{tot}}_{\text{el}}$

2. **SOURCES:**  $\bar{Q}^j \equiv q_0^j Q(E) \hat{q}_i + \sum_k^{m_k > m_j} \tilde{\Gamma}^{kj} N_i^k(0)$

Primary  
production  
in SNR

Secondary production  
by fragmentation  
of heavier nuclei

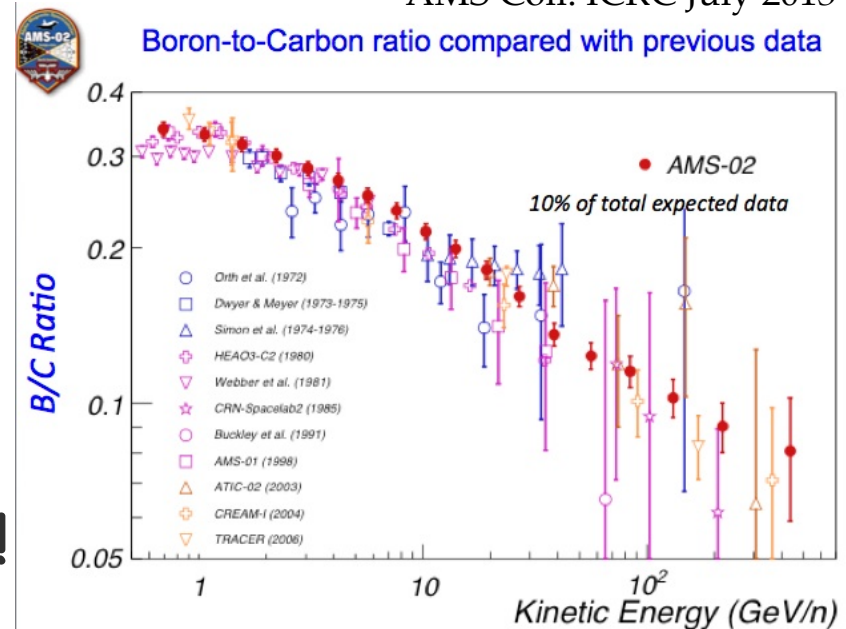
# Cosmic antimatter fluxes from DM annihilation in the Milky Way halo



## Diffusive models for CR propagation in the Galaxy

Jopikii & Parker 1970; Ptuskin & Ginzburg, 1976; Ginzburg, Khazan & Ptuskin 1980; Weber, Lee & Gupta 1992, ....; Maurin, FD, Taillet, Salati 2001; Maurin, Taillet, FD 2002; Putze, Derome, Maurin 2010; Strong & Moskalenko 1998; Moskalenko, Strong, Ormes, Potgieter, 2002; Shibata, Hareyama, Nakazawa, Saito 2004; 2006; Evoli, Gaggero, Grasso, Maccione 2008; Di Bernardo et al. 2010; ...

AMS Coll. ICRC July 2013

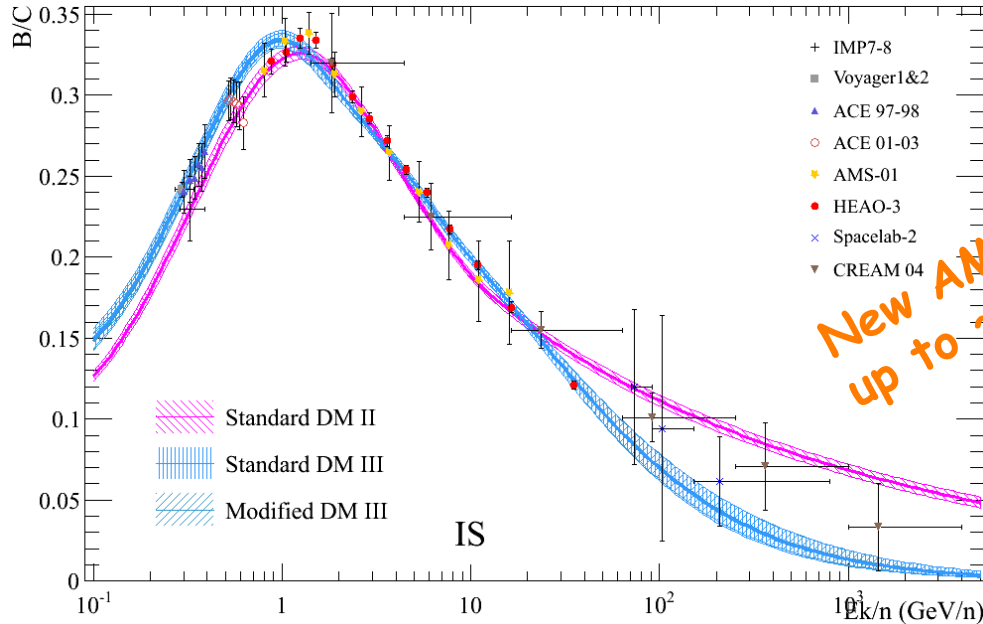


**AMS data on CRs:**  
great step forward for fixing  
Propagation and source models!

# Propagation of CRs: DATA and MODELS

## BORON/CARBON (B/C): A "STANDARD CANDLE"

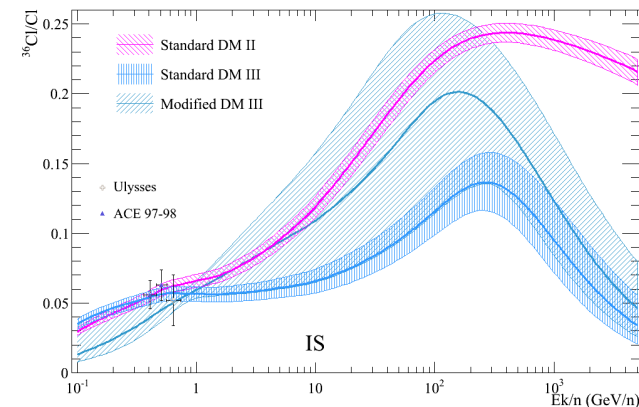
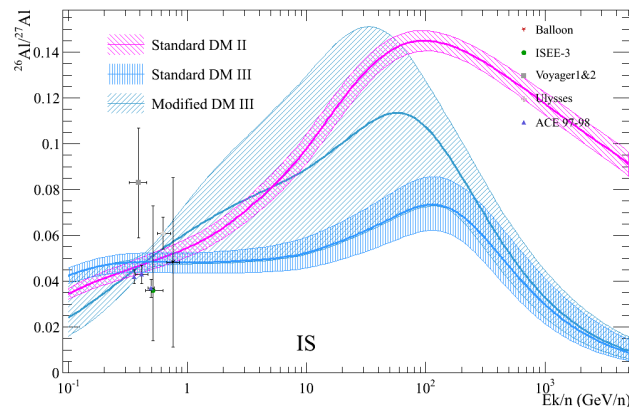
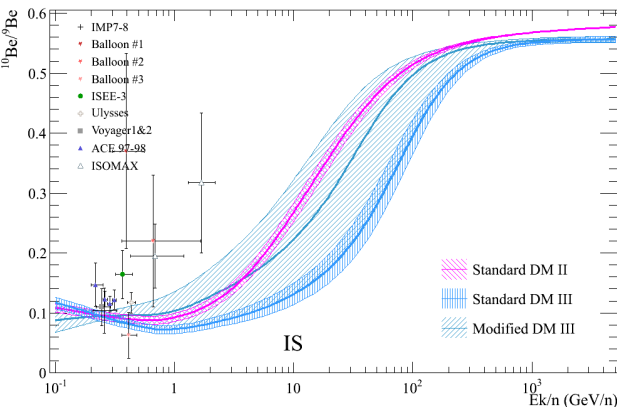
Putze, Derome, Maurin A&A 2010



New AMS-02 data  
up to  $\sim 900$  GeV

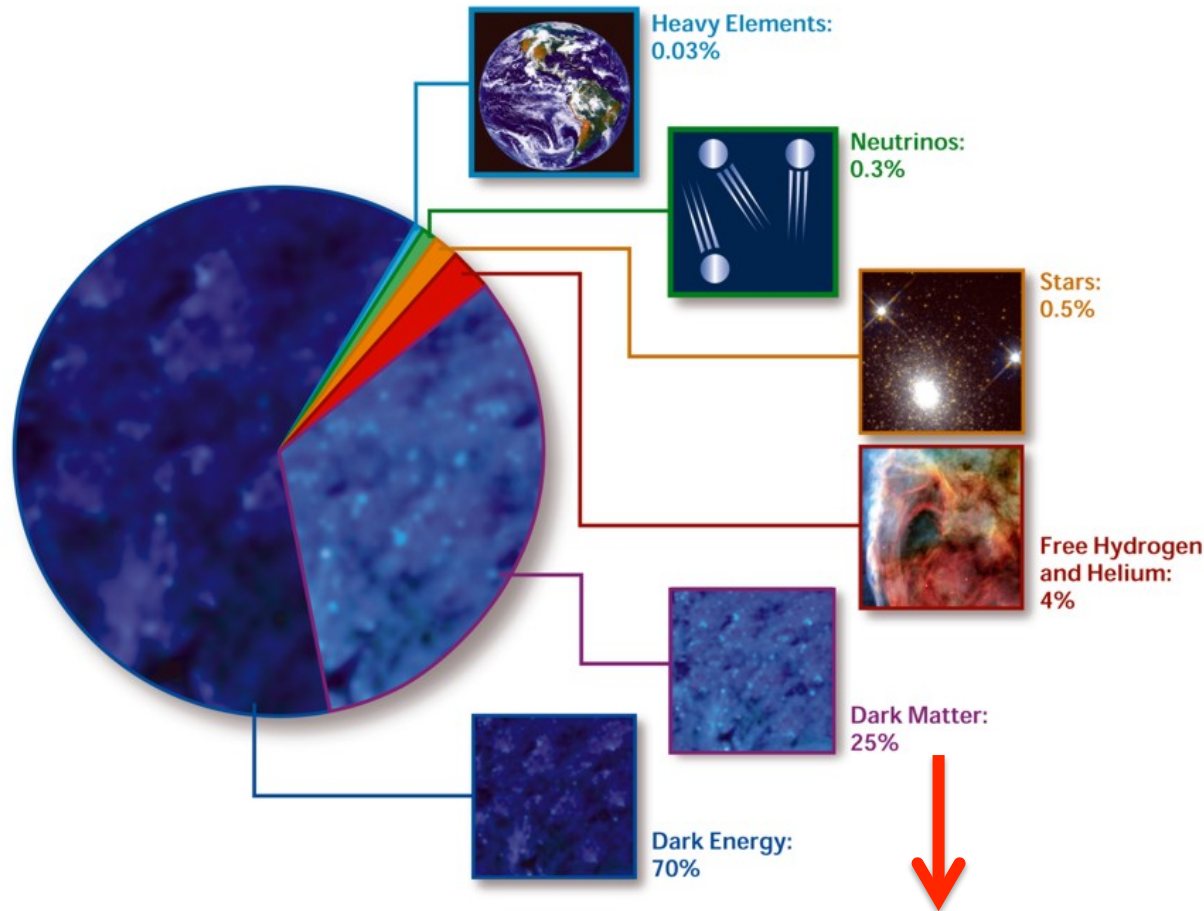
B/C: still high  
degeneracies in the  
propagation models

Radioactive isotopes  
could help break  
degeneracy



# The composition of the Universe

## COMPOSITION OF THE COSMOS

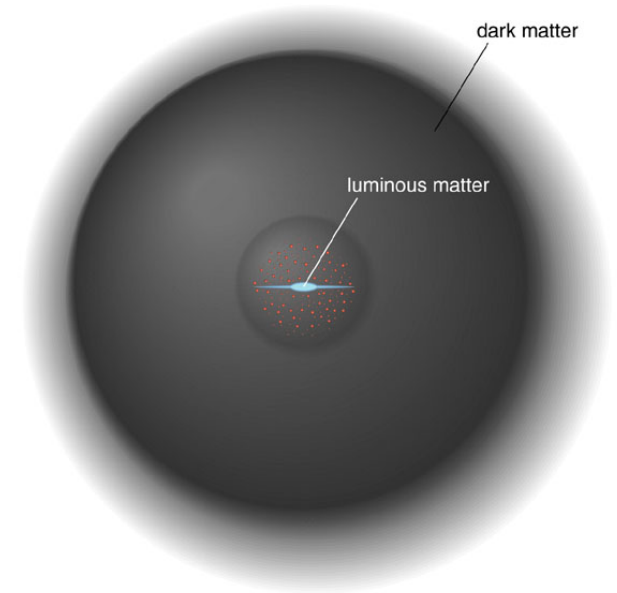
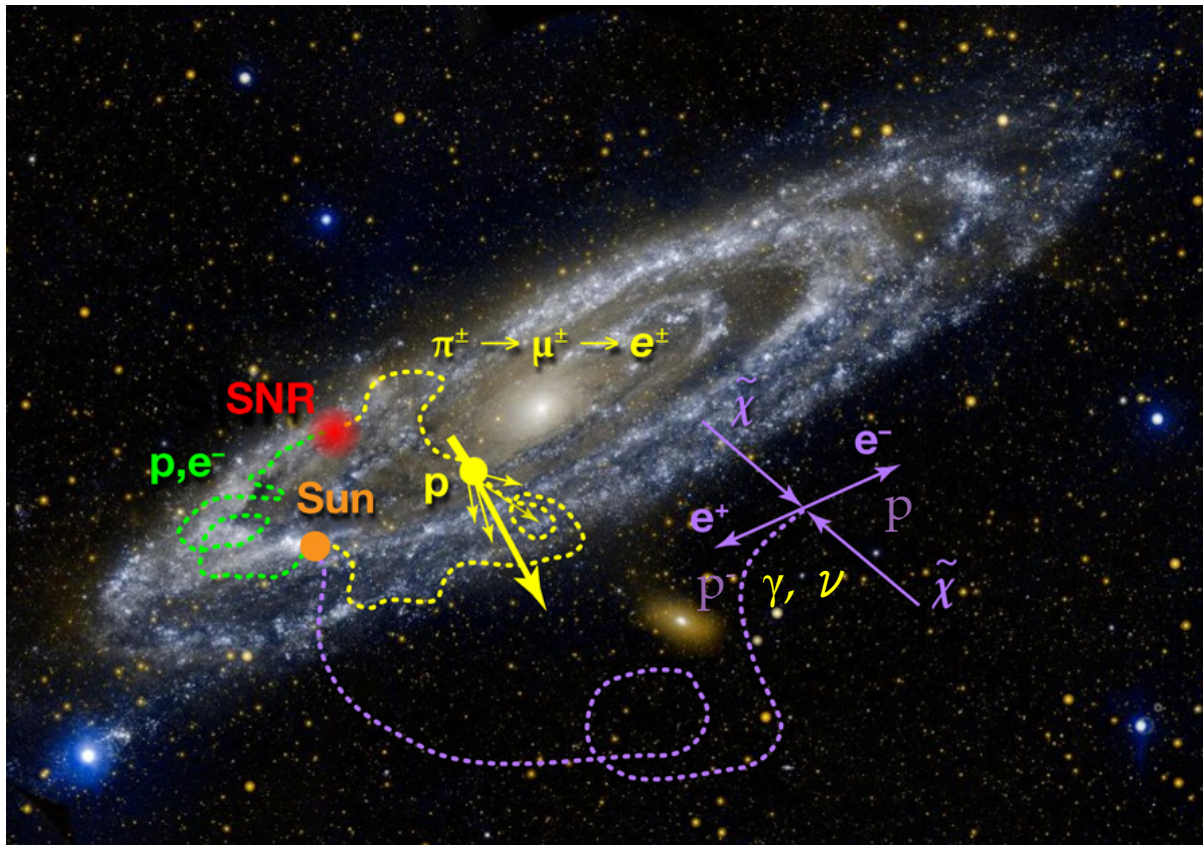


Hypothesis: the solution is a particle,  
a WIMP (weakly interacting massive particle)



# Indirect DARK MATTER searches

Dark matter can annihilate in pairs with standard model final states. Low background expected for ANTIMATTER, and for neutrinos and gamma rays coming from dense DM sites



# WIMP Dark Matter Indirect Searches

Annihilation inside celestial bodies:

→  $\nu$  at  $\nu$  telescopes as up-going  $\mu$ 's

$$\Phi_{\mu}^{(\text{Earth, Sun})} \propto \langle \sigma_{\text{ann}} v \rangle \frac{\rho_{\chi}}{m_{\chi}}$$

Annihilation in the galactic halo(s):

→ Photons ( $\gamma$ -rays, radio,...)

→  $e^+$ ,  $\bar{p}$ ,  $\bar{D}$

$$\Phi(\bar{p}, \bar{D}, e^+, \gamma) \propto \langle \sigma_{\text{ann}} v \rangle \left( \frac{\rho_{\chi}}{m_{\chi}} \right)^2$$

$\nu$  and  $\gamma$  keep directionality

can be detected only if emitted from high  $\chi$  density regions

Clue is **DM** space distribution  $\rho(r)$

Charged particles diffuse in the galactic halo

antimatter searched as rare components in cosmic rays (CRs)

Clue is the **astrophysics of charged cosmic rays**

N.B. New particles are searched at **colliders**  
but we cannot say anything about being  
the solution to the **DM** in the Universe!

The case for  
antiprotons

# Cosmic antiprotons

Antiprotons are produced in the Galaxy by fragmentation of proton and He (and marginally heavier nuclei) on the ISM (**secondary antiprotons**).

These antiprotons would be the background to an exotic component due to **dark matter annihilation** in the galactic halo (**primary antiprotons**).

N. B. Thousands of cosmic antiprotons have already been detected by balloon-borne (Bess, Caprice,...) or satellite experiments (Pamela), and AMS-01, and 290000 (out of 54 billion events) from AMS-02 on the ISS

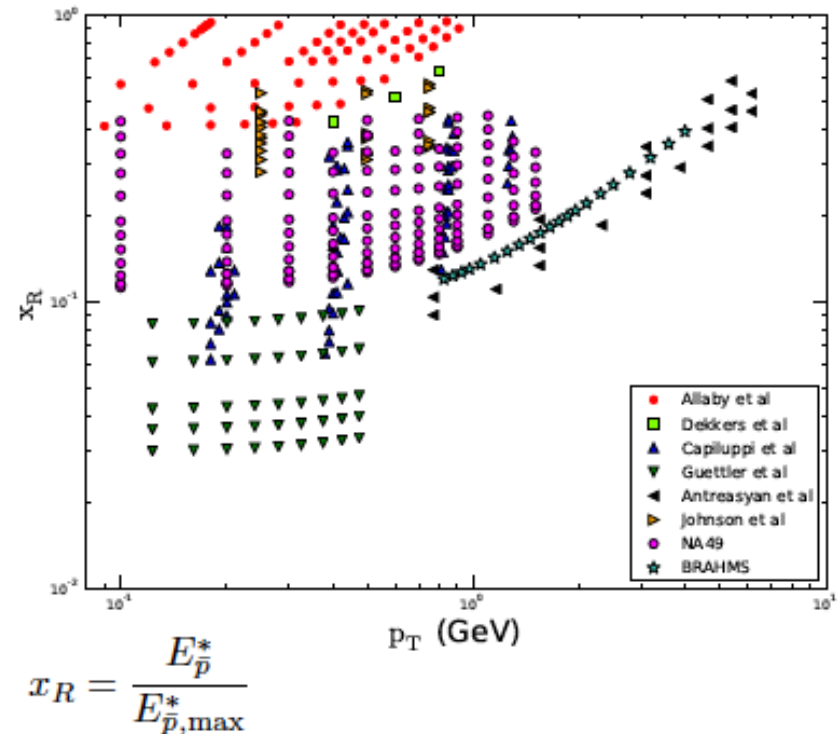
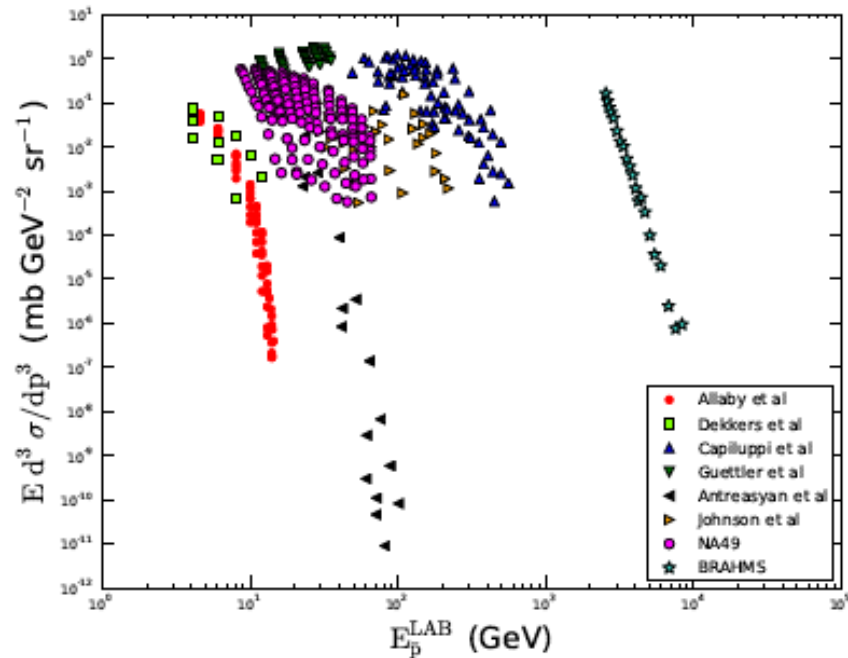


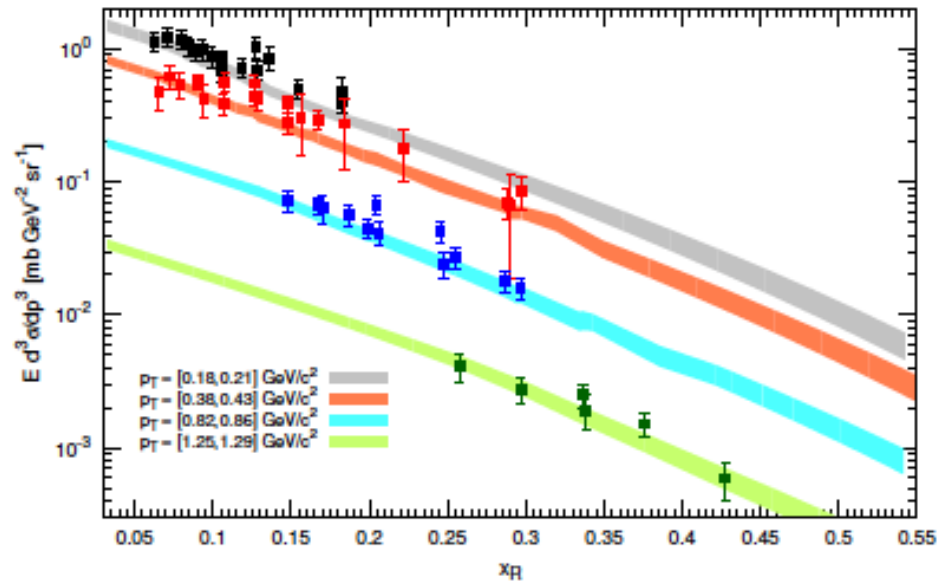
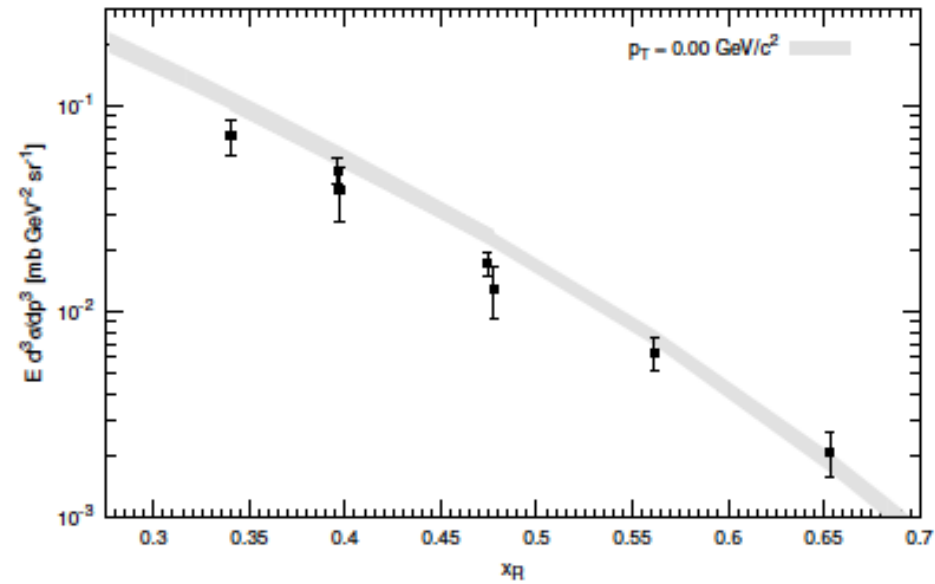
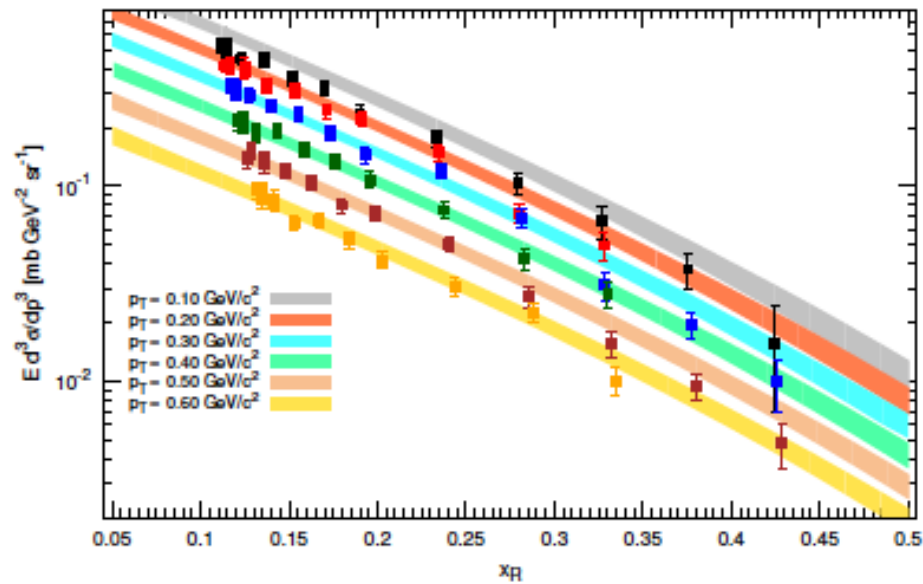
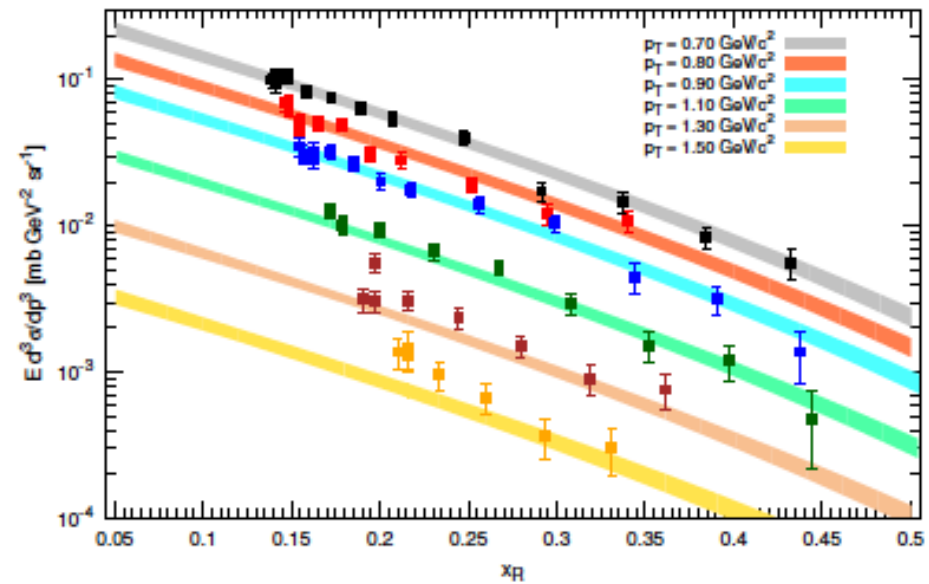
# New analysis of $p\text{-}p \rightarrow p\bar{p}$ data

Di Mauro, FD, Goudelis, Serpico PRD 2014, 1408.0288; Kappl, Winkler 1408.0299, JCAP2014

$$q_{\bar{p}}^{pp}(E_{\bar{p}}) = \int_{E_{th}}^{+\infty} \frac{d\sigma_{pp \rightarrow \bar{p}}}{dE_{\bar{p}}}(E_p, E_{\bar{p}}) n_H(4\pi\Phi_p(E_p)) dE_p$$

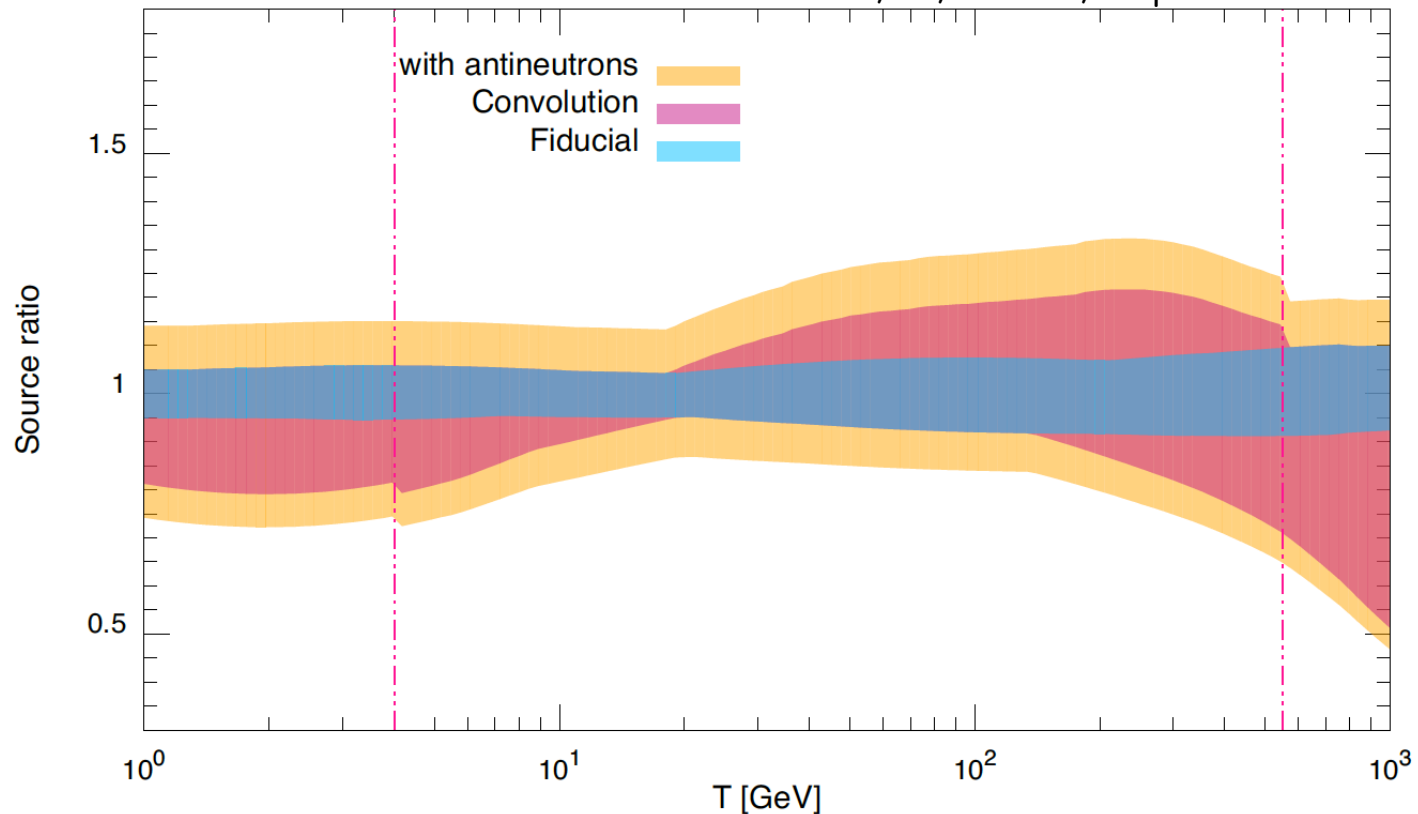
Existing data on the production cross section



Capiluppi et al. 1974  $s^{1/2} = [23.3, 63]$  GeVDekkers et al. 1965  $s^{1/2} = 6.7$  GeVAnticic et al. 2010  $s^{1/2} = 17.3$  GeVAnticic et al. 2010  $s^{1/2} = 17.3$  GeV

# Uncertainties due p-p scattering

Di Mauro, FD, Goudelis, Serpico PRD 2014

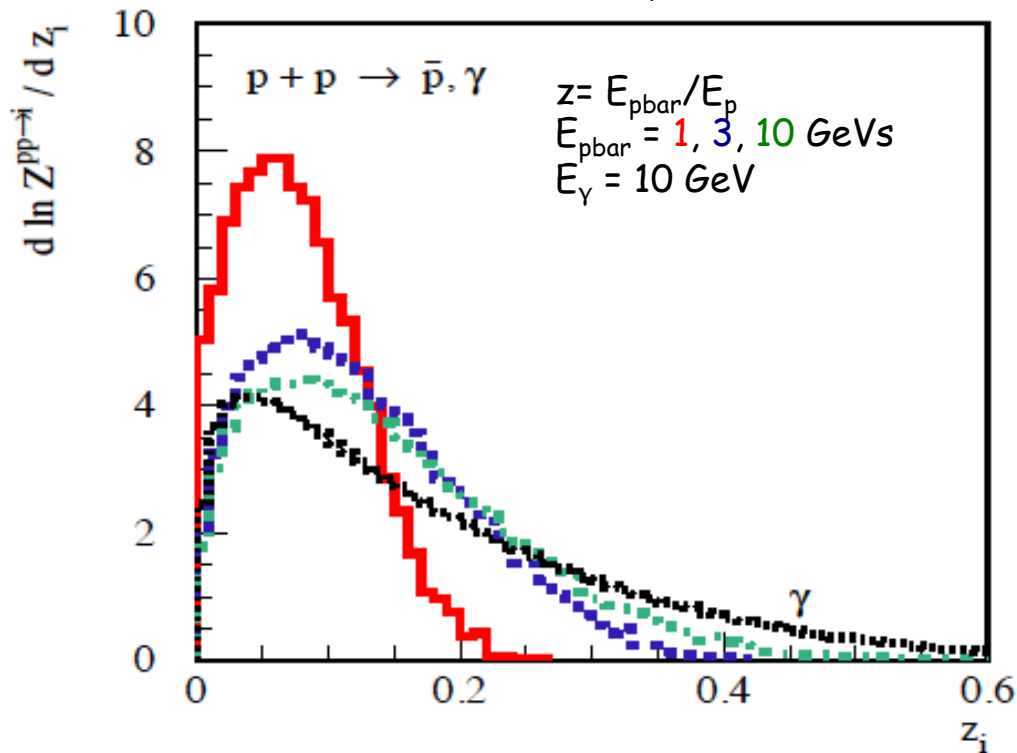


Uncertainties in the  $p\bar{p}$  production spectrum from p-p scattering are at least 10%.

Conservative: 20% at low energies (GeV) up to 50% (TeV)  
(data expected at least up to  $\sim 500$  GeV)

# Protons $\rightarrow$ antiprotons

Kachelriess, Moskalenko, Ostapchenko PRD2015



The bulk of antiprotons  
is produced by **protons**  
with kinetic energy  
10-30 times greater

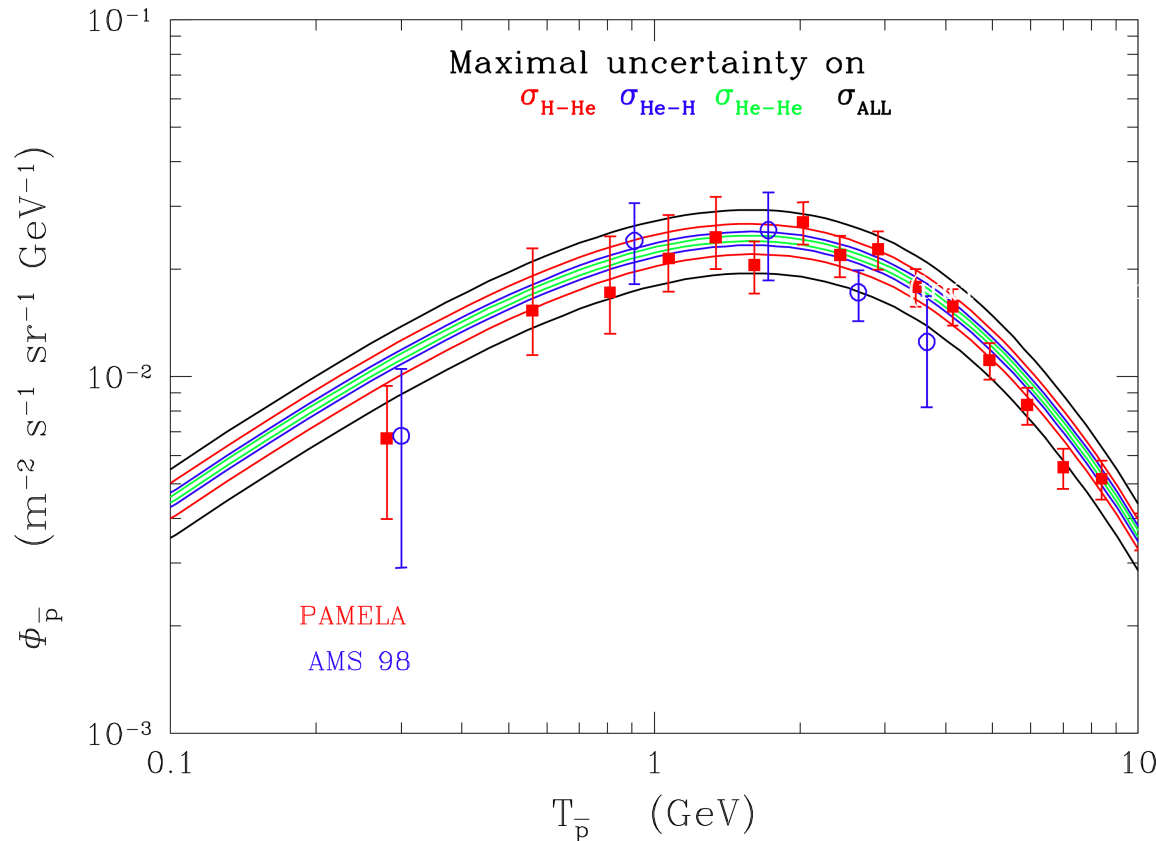
AMS energies  $\sim 1\text{-}500 \text{ GeV}$  in  $E_{\bar{p}}$   $\rightarrow$

Proton energies with  $\sim \text{few GeV} - 10 \text{ TeV}$



# Uncertainties on the antiproton flux from nuclear cross sections

(Donato+ ApJ 2001, PRL 2009)



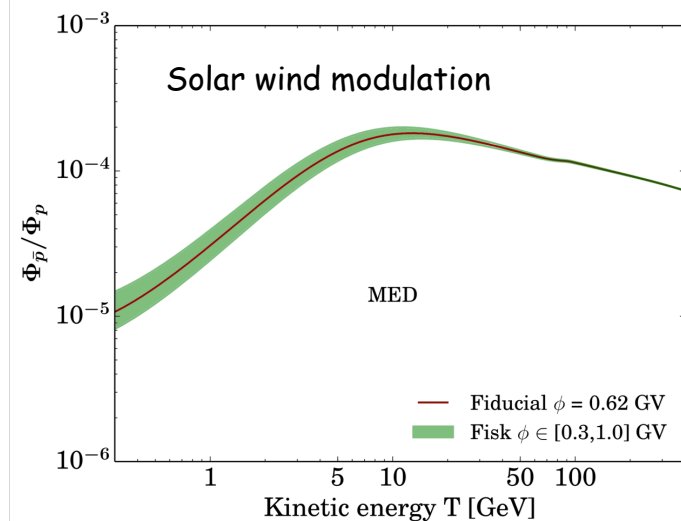
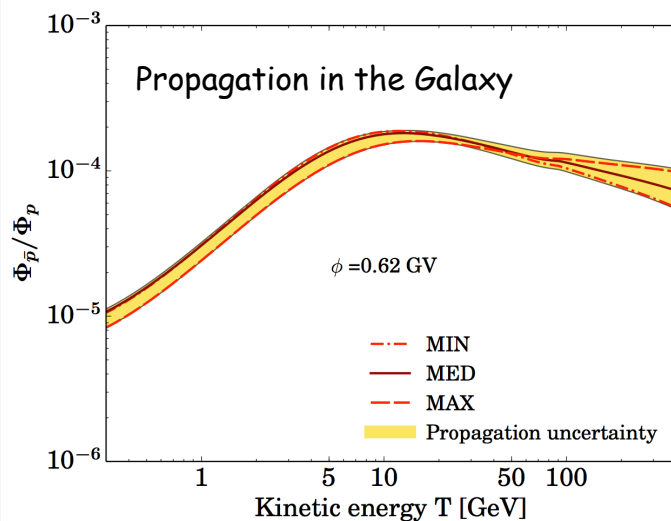
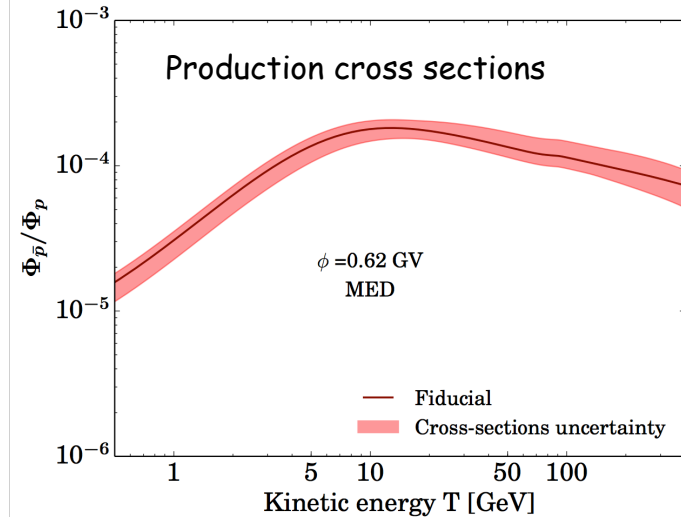
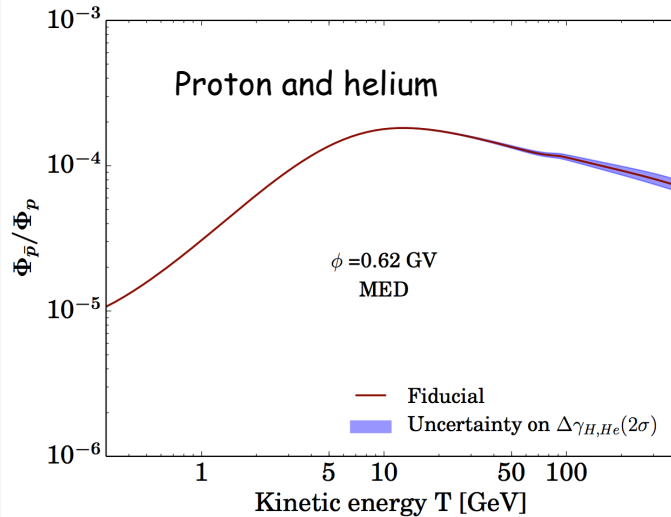
- pp: Tan& Ng
- H-He, He-H, He-He: DTUNUC MC
- Functional form for the cross section derived from other reactions, given **NO DATA!!**

**Maximal uncertainty from p-He cross sections: 20-25%!**

Data from AMS-02 on cosmic antiprotons are at ~ 10% accuracy

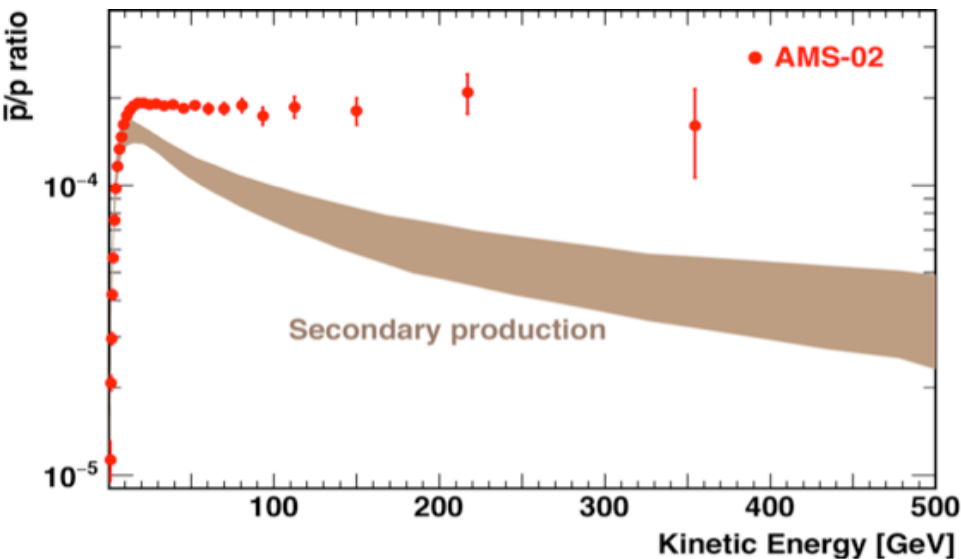
# Secondary antiprotons: theoretical uncertainties

Giesen+ 1504.04276

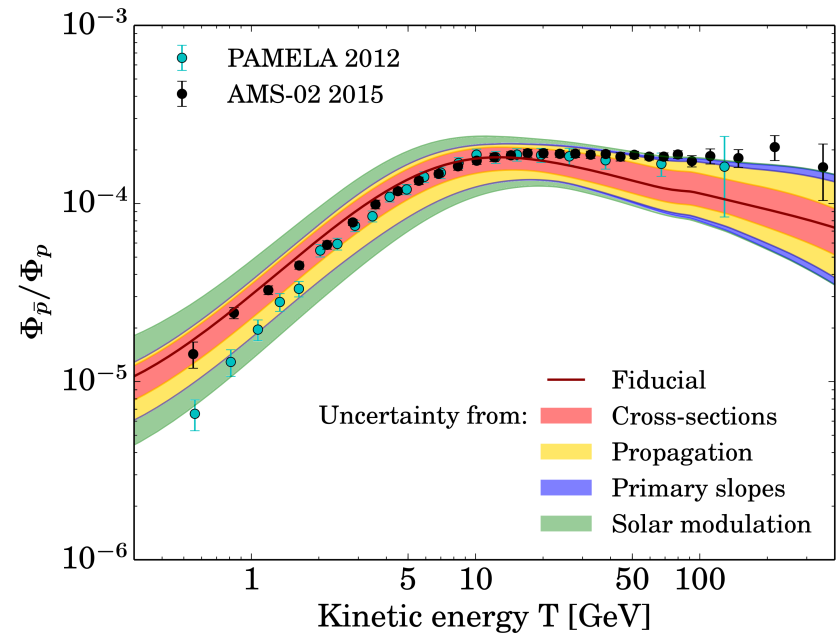


# Prediction and AMS data

AMS Coll., Cern 15.04.2015



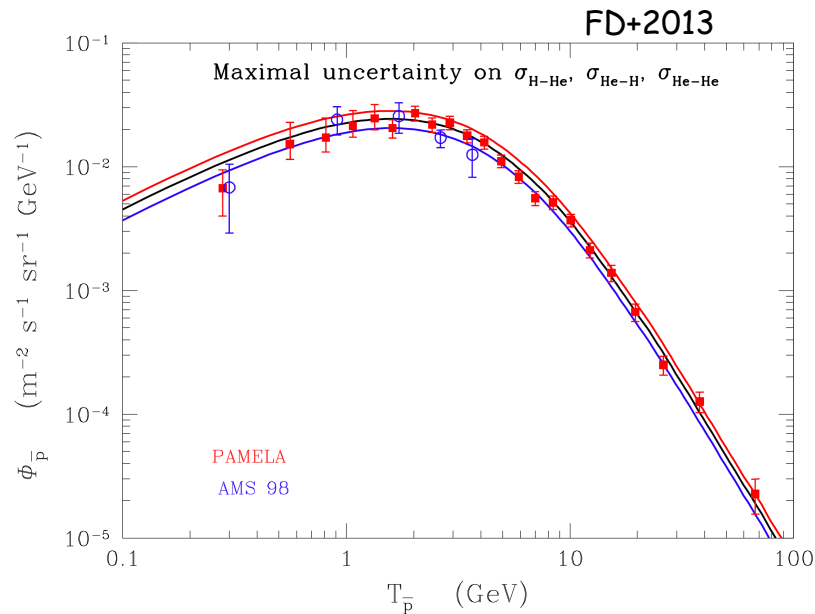
Giesen+ 1504.04276



Very recent AMS-02 results up to 450 GeV  
Can be explained by secondary production in the Milky Way,  
considering several theoretical uncertainties

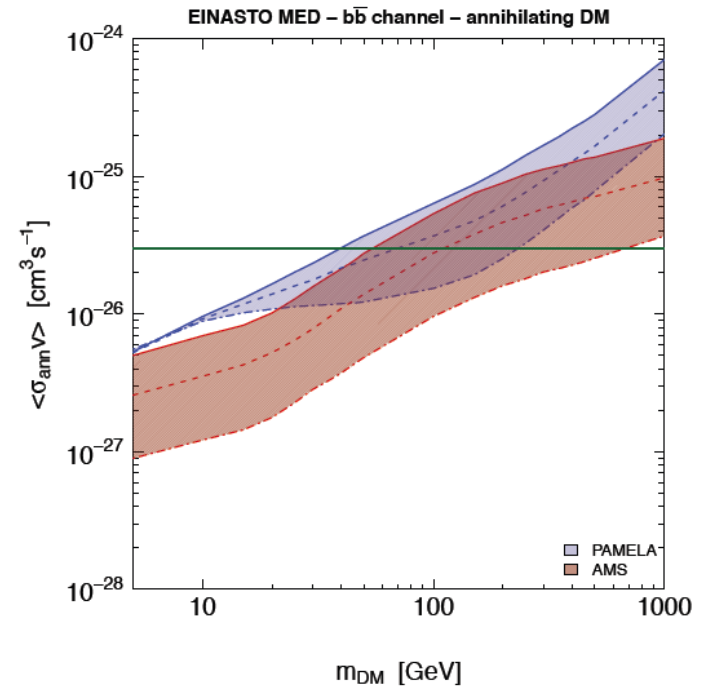
# He cross sections: effect on DM interpretation

Uncertainties due to helium reactions range 40-50% on Secondary CR flux



Effect of cross section uncertainty on DARK MATTER interpretation

Fornengo, Maccione, Vittino JCAP2014



AMS-02 is providing data with much higher precision up to hundreds of GeV!!!  
Their interpretation risks to be seriously limited by nuclear physics



The case for  
positrons

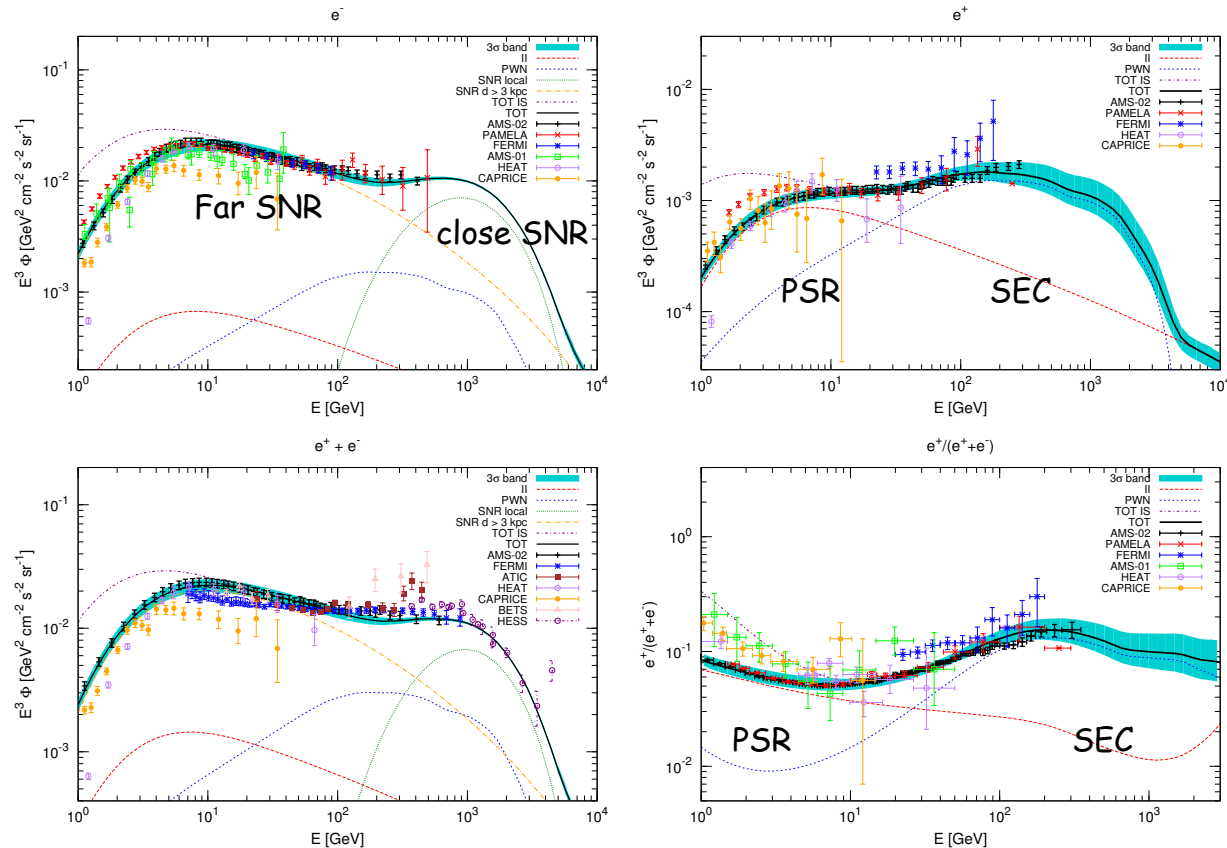
# Sources of positrons in the Milky Way

## Sources of $e^+$ and $e^-$ in the Galaxy:

1. **Secondary  $e^+ e^-$ :** spallation of cosmic p and He on the ISM (H, He)
  - \*  $p+H(He) \rightarrow p+\Delta^+ \rightarrow p+\pi^0 \text{ \& \; } n+\pi^+$  (mainly below 3 GeV)
  - \*  $p+H(He) \rightarrow p+n+\pi^+$
  - \*  $p+H(He) \rightarrow X + K^\pm$
2. **Primary  $e^-$  and  $e^+$  from Pulsars (PSR):**  
pair production in the strong PULSAR magnetosphere
3. **Primary  $e^-$  from SNR:** 1<sup>o</sup> type Fermi acceleration mechanism
4. **Primary  $e^+ e^-$  from exotic sources (DARK MATTER)**

# AMS lepton data: an astrophysical interpretation

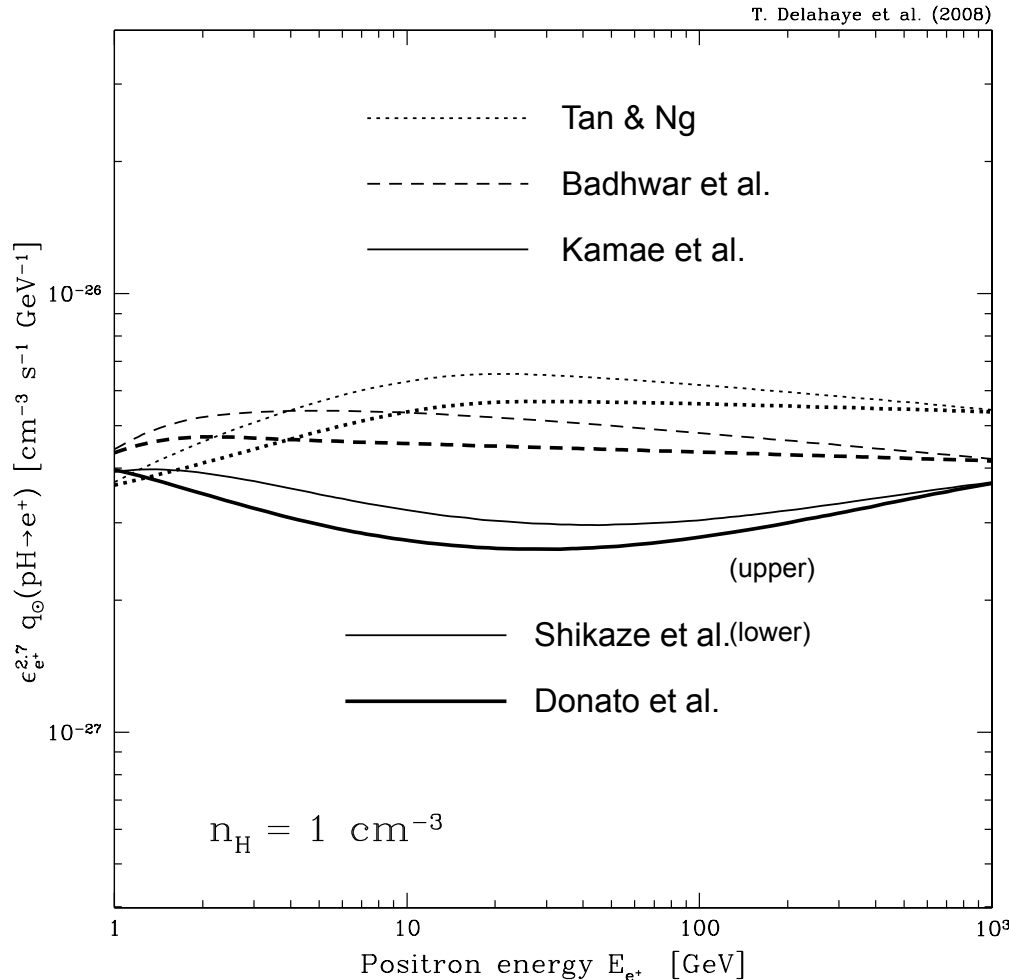
Di Mauro, FD, Fornengo, Vittino JCAP 2014



TH: Secondaries + supernovae + pulsars  
 EXP: Ams data precise + wide range  $\rightarrow$   
 Small features can bring strong information!

# The secondary positron source term: effect of cross sections

Delahaye, FD, Fornengo, Lavallo, Lineros, Salati, Taillet A&A 2009



Effect of proton flux  
determination - negligible

Effect of production  
cross sections is  
not negligible!!

Up to factor 2 in p-p!

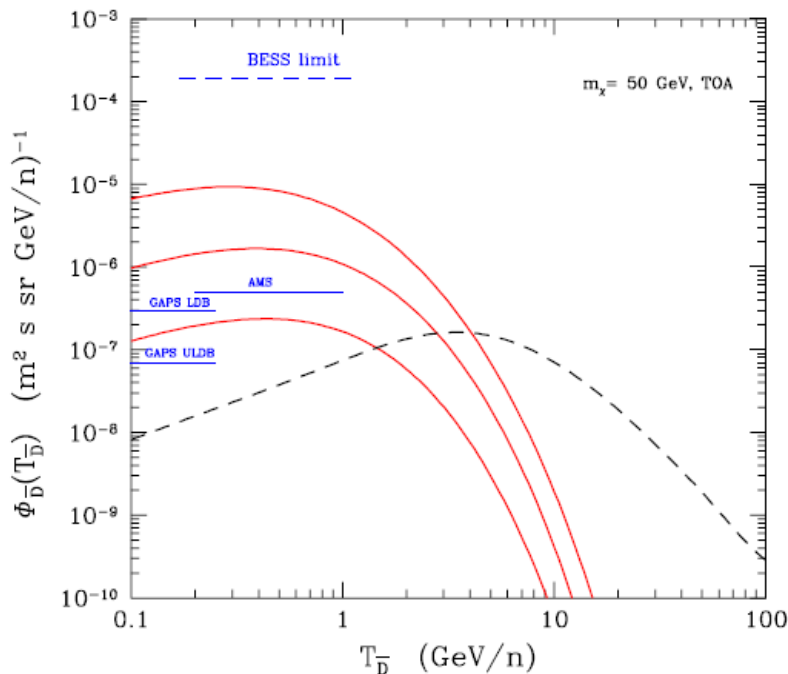
The case for  
antideuteronos



# COSMIC ANTIDEUTERONS

FD, Fornengo, Maurin PRD 2001; 2008; Kadastik, Raidal, Strumia PLB2010; Ibarra, Wild JCAP2013;  
Fornengo, Maccione, Vittino JCAP 2013; ...

In order for fusion to take place, the two antinucleons must have low kinetic energy



Kinematics of **spallation** reactions prevents the formation of very low antiprotons (antineutrons).

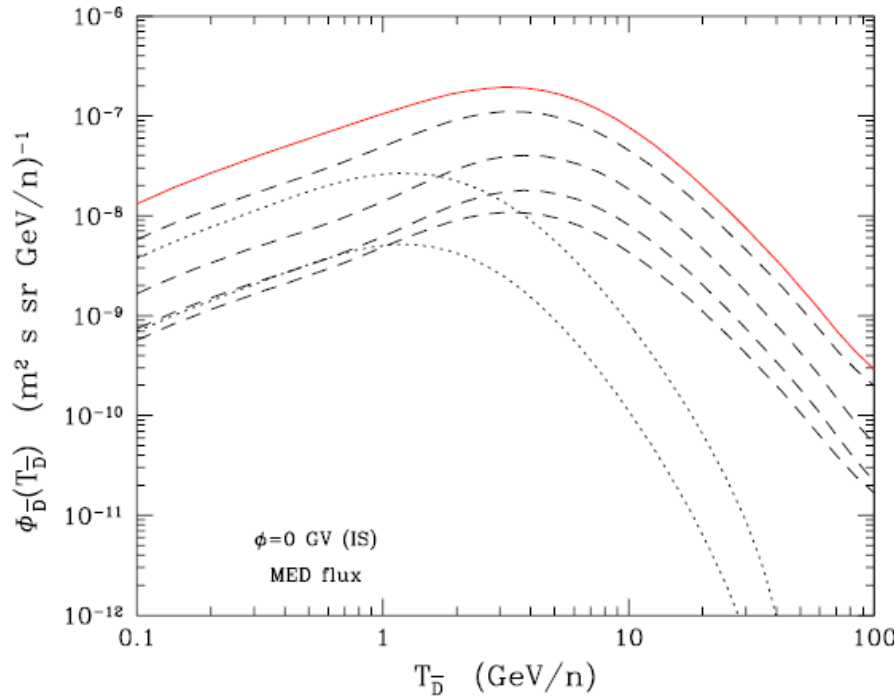
At variance, **dark matter** annihilate almost at rest

N.B: Up to now, NO ANTIDEUTERON has been detected yet.  
Several experiments are on the road: AMS/ISS, BESS-Polar, GAPS ...

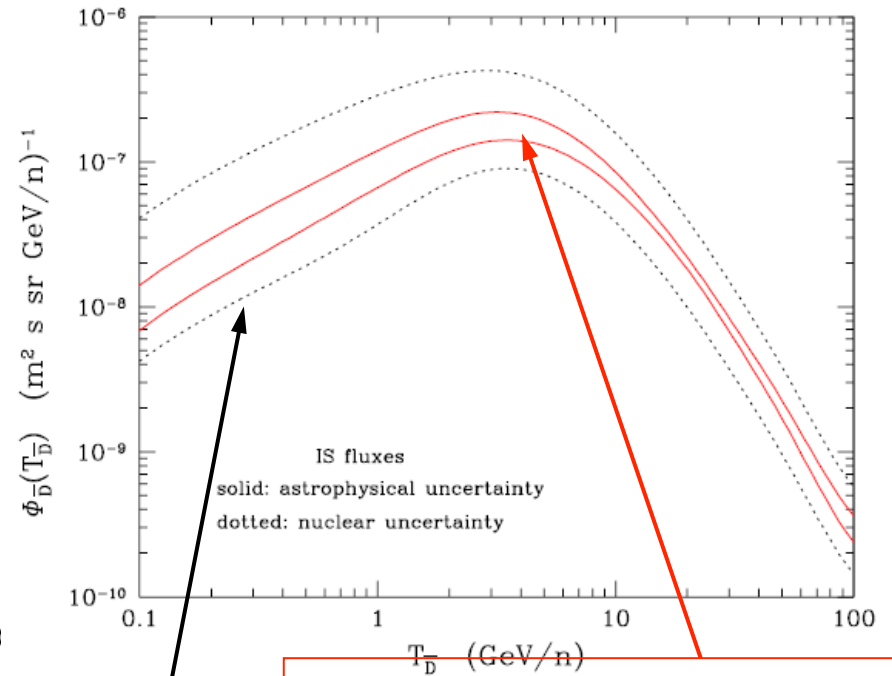
# Secondary antideuteron

FD, Fornengo, Maurin PRD 2008

## Contributions to secondaries



p-p, p-He,  
He-H, He-He  
H- pbar, He-pbar

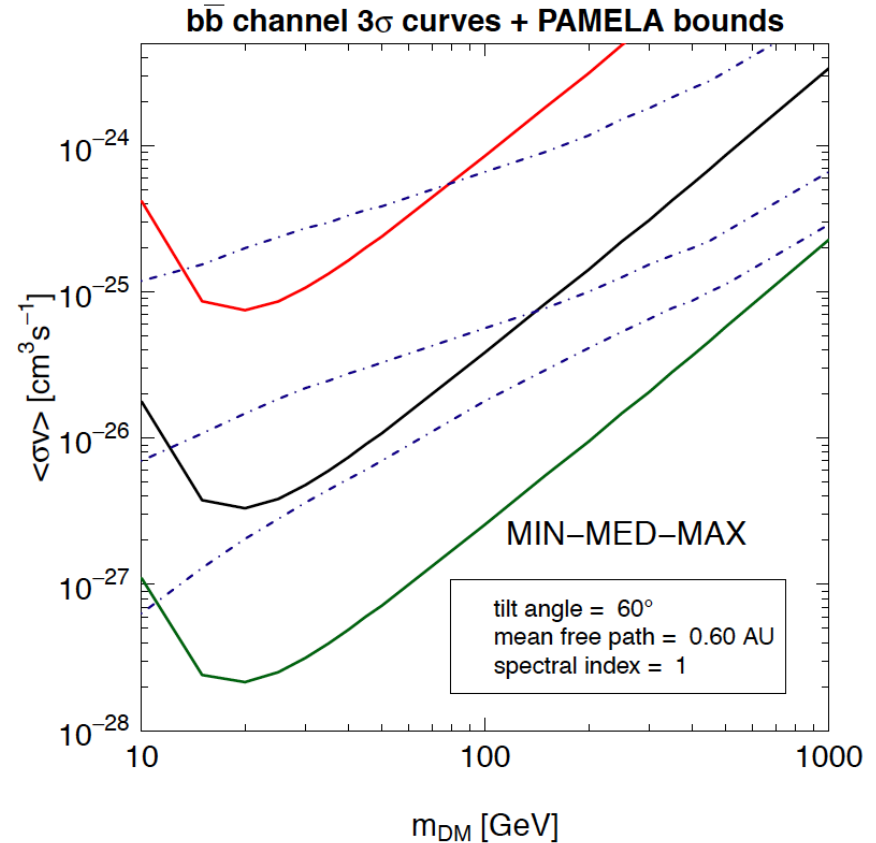
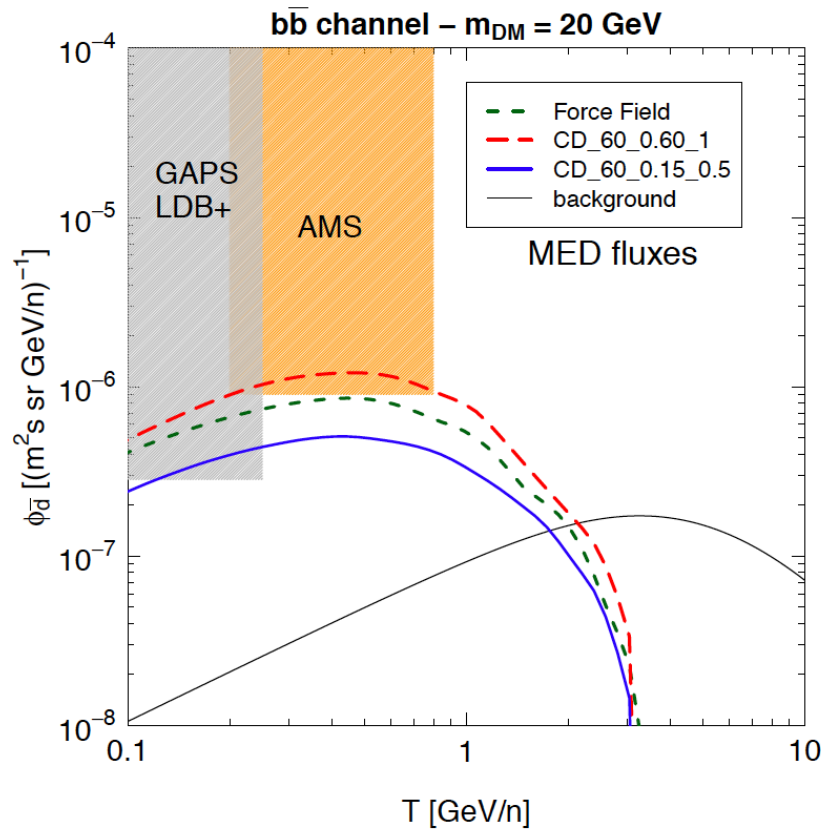


Propagation uncertainties  
Compatibility with B/C

Nuclear uncertainties  
Production cross sections &  $P_{\text{coal}}$   
Production from antiprotons  
Non-annihilating cross sections

# Antideuterons: Dark matter detection perspectives

Fornengo, Maccione, Vittino 1306.4171



$3\sigma$  expected sensitivities

Prospects for  $3\sigma$  detection of antideuteron with GAPS (dotted lines are Pamela bounds from antiprotons)

# Few final remarks

The astrophysics of cosmic rays is entering an era of remarkable **precision** (AMS-02/ISS)

ANTIMATTER is a key element for testing galactic models and searching for DARK MATTER signals

**Propagation uncertainties** confined to less than 20%, close to be  $< 10\%$

The (production, total inelastic, inelastic non-annihilating, ...) cross sections from Fe down, including isotopes and antimatter, rely on very few or (often) NO lab data!!

→ A HUGE EXPERIMENTAL PROGRAM IS REQUIRED FOR A SIGNIFICANT REDUCTION OF UNCERTAINTIES

Proton-helium → antiproton + X looks the most urgent case (contextually, measure of  $\rightarrow e^+, \gamma, D^-$ )

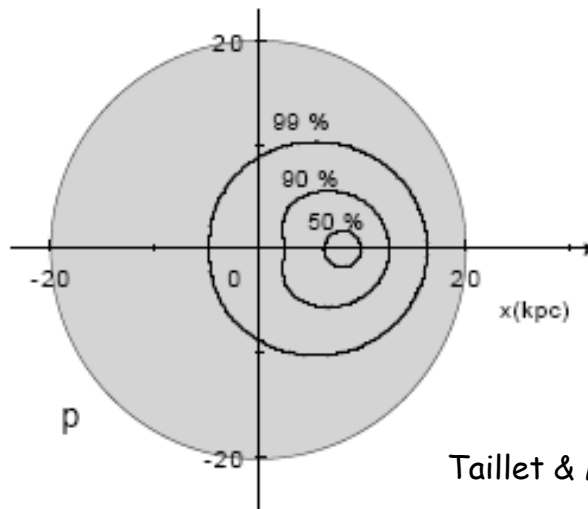
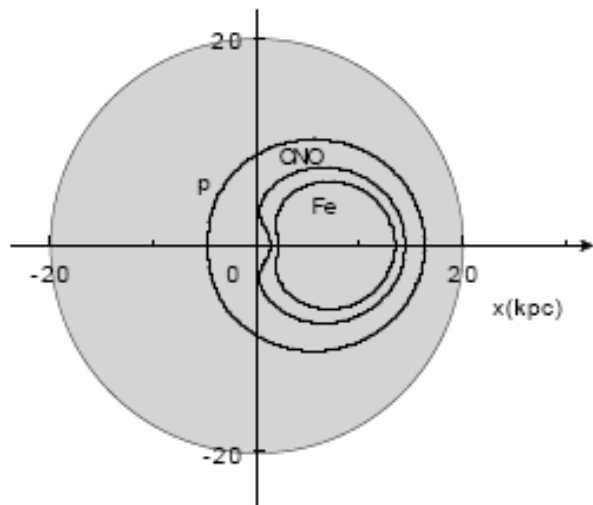
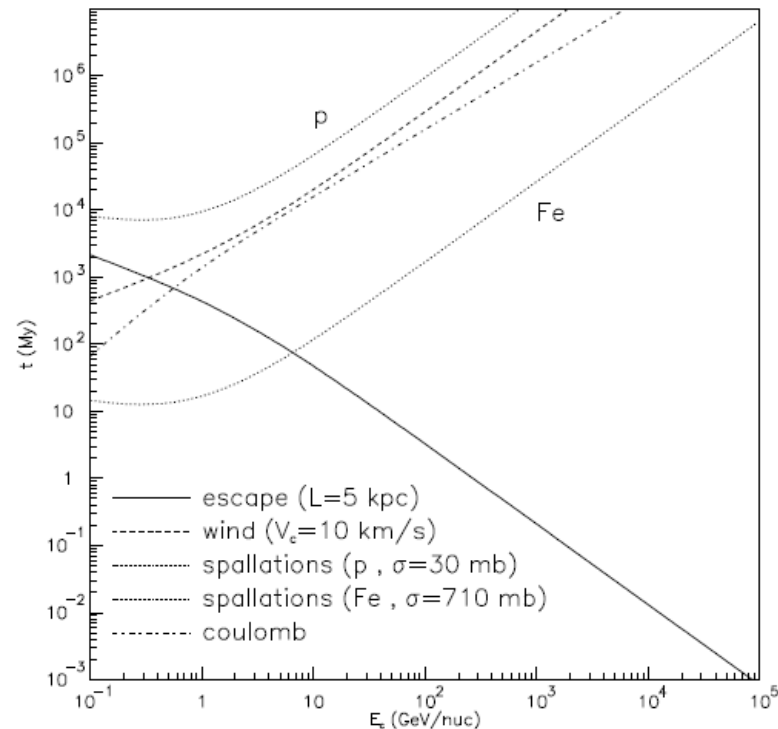
BACKUP SLIDES



# Characteristic times and distances

The smaller the time,  
the most effective the process is

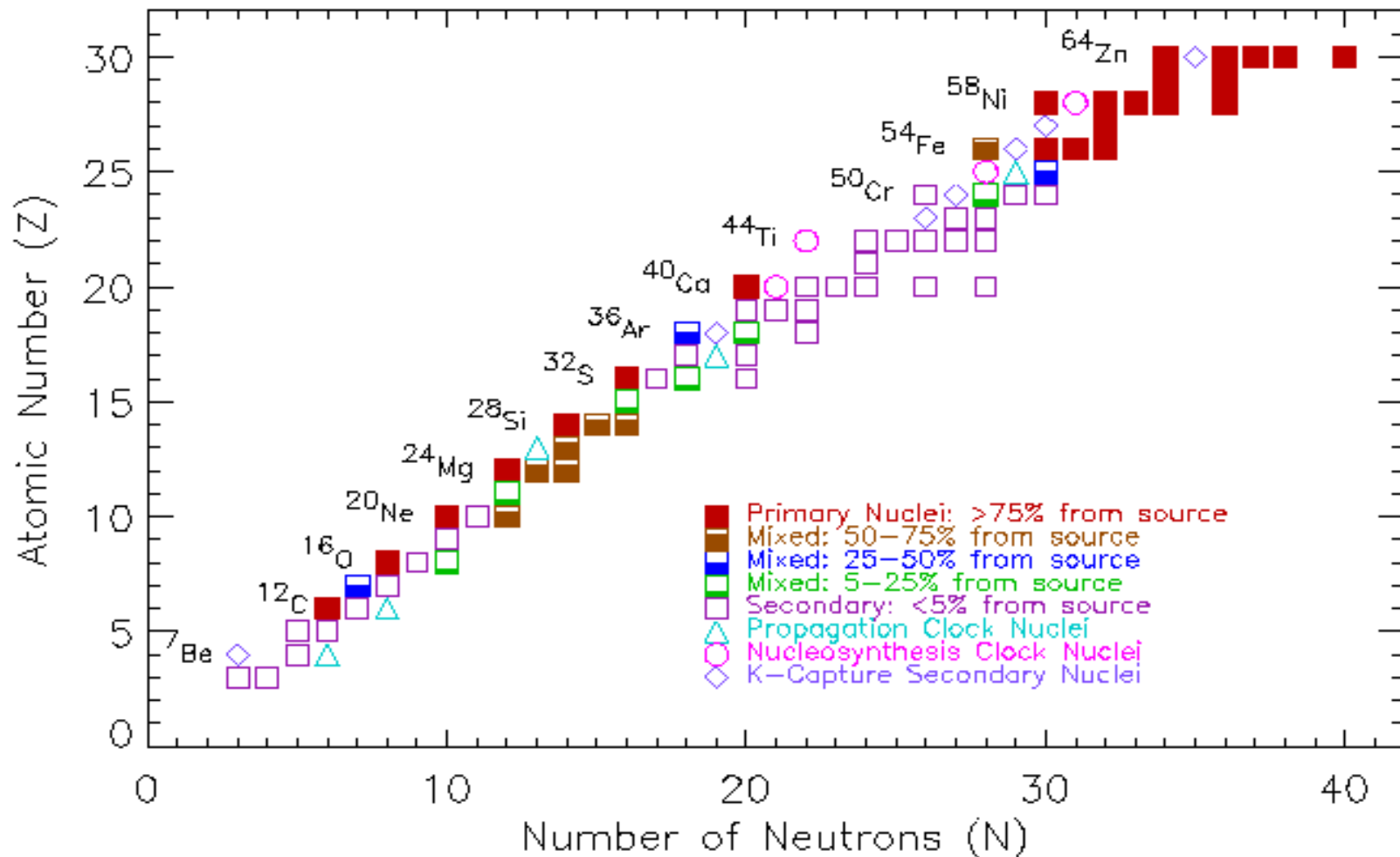
Protons: escape  $E > 1$  GeV  
convection and e.m. losses  $E < 1$  GeV,  
Iron: escape  $E > 1$  GeV  
Spallations  $E < 10$  GeV/n



Fe more local than p  
90% p from 5-6 kpc

**Primaries** = present in sources:  
 Nuclei: H, He, CNO, Fe;  $e^-$ , ( $e^+$ ) in SNR (& pulsars)  
 $e^+$ ,  $p^+$ ,  $d^+$  from Dark Matter annihilation

**Secondaries** = NOT present in sources, thus produced by  
**spallation** of primary CRs (p, He, C, O, Fe) on ISM  
 Nuclei: LiBeB, sub-Fe, ... ;  
 $e^+$ ,  $p^+$ ,  $d^+$ ; ... from inelastic scatterings



# X-sec uncertainties: impact on GCR model parameters

(Slide from D. MAURIN)

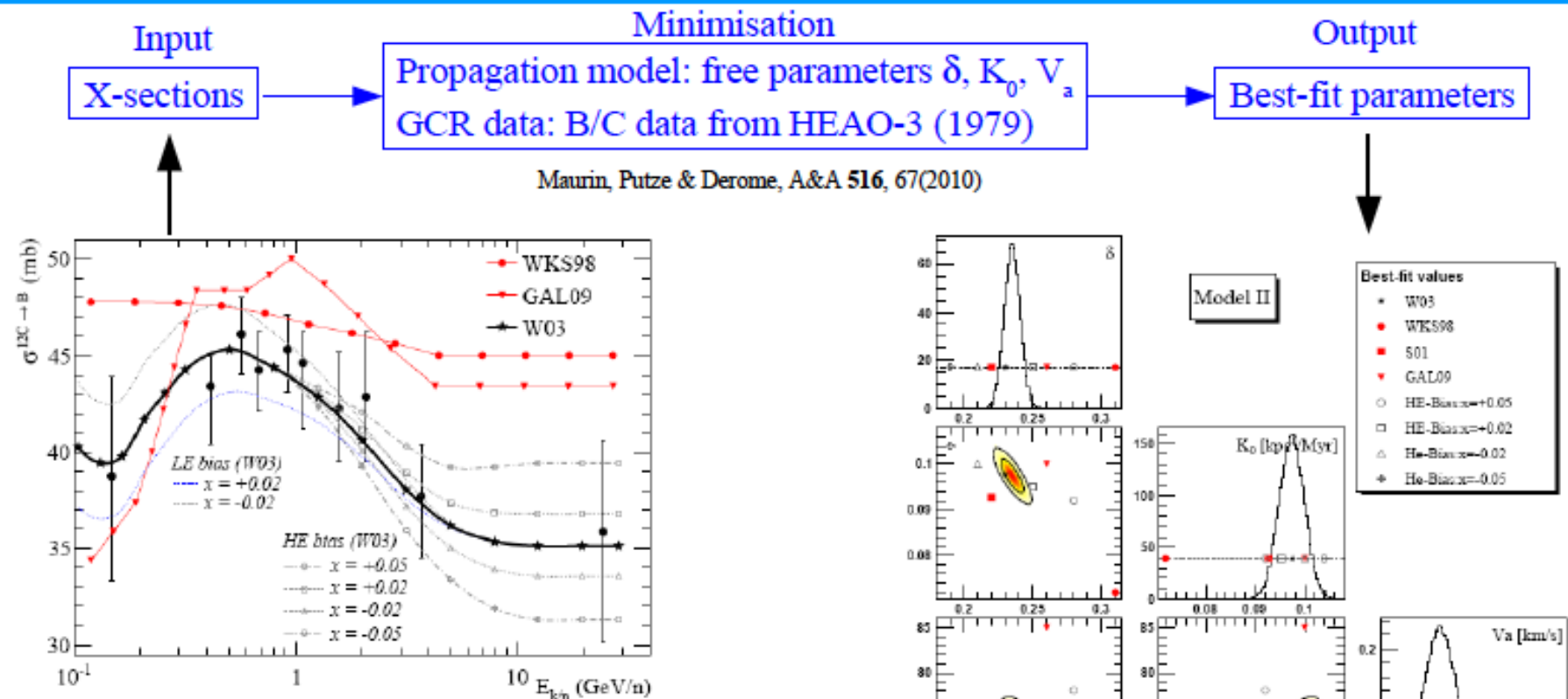


Fig. 3. Production cross-section for  $^{12}\text{C}+\text{H}\rightarrow^{10,11}\text{B}$  (adapted from Webber et al. 2003). The standard sets are shown as solid lines (WKS98: red dots; GAL09: red down triangles; W03: black stars), and the biased sets in dotted ( $|x| = 0.02$ ) and dashed ( $|x| = 0.05$ ) lines.

- W03 and WKS98 are parameterisations of the same 'data' (energy bias)
- GAL09: modern nuclear codes normalised to LANL database [Moskalenko & Mashnik, astro-ph/0306367]

→ Systematics uncertainties (from X-sec) > statistical uncertainties (from GCR data)  
 ... and AMS-02 is at least 100 better than HEAO-3!

III. Uncertainties are too large!

# Inelastic cross sections ( $\sigma_R = \sigma^{tot} - \sigma_{ela}^{tot}$ )

(Slide from D. MAURIN)

- Bradt & Peters (1950)

$$\sigma_R = \pi r_0^2 (A_{proj}^{1/3} + A_{cible}^{1/3} - b_0)^2$$

- Letaw *et al.* (1970-2000): **accuracy <2% for 2<Z<30 and E>300MeV/n**

S.Barshay & al, Phys.Lett **51B**, 5 (1974)

J.R.Letaw & al, ApJSS **51**, 271 (1983)

R.Silberberg & C.H Tsao, Phys.Rep. **191**, 351 (1990)

L.Sihver & al, Phys.Rev.C **47**, 1225 (1993)

H.P.Wellish & D.Axen, Phys.Rev.C **54**, 1329 (1996)

R.Silberberg & al, ApJ **501**, 911 (1998)

C.H.Tsao & al, ApJ **501**, 920 (1998)

$$\mu_R(E_k) = \sigma_R^{HE} [1 - 0.62 \exp(-E_k/200) \sin(10.9 E_k^{-0.28})] \text{ (mb)}$$

$$\sigma_R^{HE} = 45 A^{0.7} [1 + 0.016 \sin(5.3 - 2.63 \ln A)] \text{ (mb)}$$

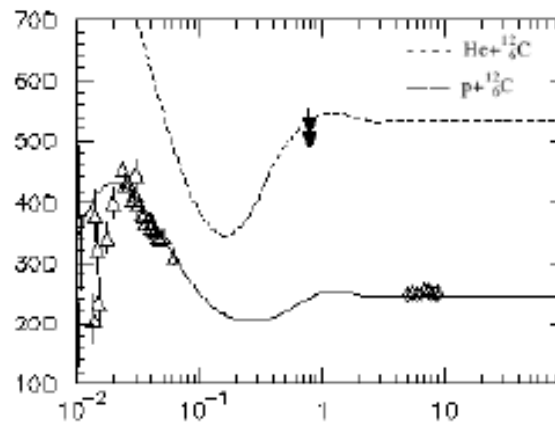
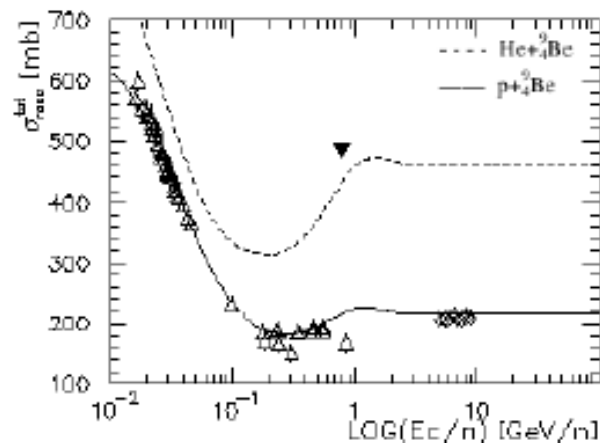
- Tripathi *et al.* (~2000): **~ or better (at low E) than Letaw *et al.*, valid for all N+N reaction!**

R.K.Tripathi & al, NASA, Technical Paper **3621**, (1997)

R.K.Tripathi & al, NASA, Technical Paper **3656**, (1997)

R.K.Tripathi & al, NASA, Technical Paper **209726**, (1999)

$$\sigma_R = \pi r_0^2 \left( A_{proj}^{1/3} + A_{cible}^{1/3} + \delta_E \right)^2 \left( 1 - R_c \frac{B}{E_{cm}} \right) X_m$$



**Data from compilations:**  
Bobchenko 79, Tanihata 85,  
Bauhoff 86, Carlson 96

→ Tripathi *et al.* is the one generally used in the field

# Production cross sections (straight-ahead approx.)

$$\int_0^\infty n_H v' N^k(T') \sigma^{kj}(T, T') dT' = \int_0^\infty n_H v' N^k(T') \sigma^{kj}(T) \delta(T - T') dT' = n_H v N^k(T) \sigma^{kj}(T)$$

## - Semi-empirical approach [Silberberg et Tsao]

- for any Proj. + Targ.  $\rightarrow$  Frag.
- better than Webber if extrapolation ( $Z > 30$ )

## - Empirical approach [Webber et al.]

- for Proj. + H/He  $\rightarrow$  Frag.
- better than Silberberg on 'data' ( $Z < 30$ )

## - More recent empirical codes

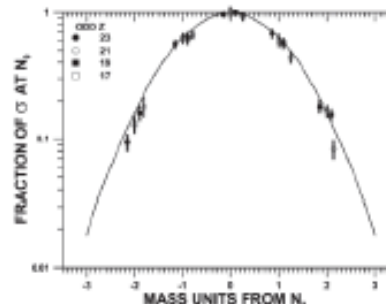
- EPAX2 <http://www-w2k.gsi.de/hellstr/asp/gsi/epaxv21m.asp>
- PHITS [phits.jaea.go.jp](http://phits.jaea.go.jp)

## - Microscopic description

- LAQGSM (Los Alamos Quark Gluon String Model)
- NUCFRG2 (semi-empirical abrasion-ablation model)

R.Silberberg & C.H.Tsao, ApJSS **25**, 315 (1973)  
 R.Silberberg & C.H.Tsao, ApJSS **35**, 129 (1977)  
 R.Silberberg & C.H.Tsao, ApJSS **35**, 137 (1977)  
 R.Silberberg & al, ApJSS **58**, 877 (1985)  
 R.Silberberg & C.H.Tsao, Phys.Rep. **191**, 351 (1990)  
 L.Silver & al, Phys.Rev.C **47**, 1225 (1993)  
 C.H.Tsao & al, Phys.Rev.C **47**, 1257 (1993)  
 C.J.Waddington, ApJ **470**, 1218 (1996)  
 R.Silberberg & al, ApJ **501**, 911 (1998)  
 C.H.Tsao & al, ApJ **501**, 920 (1998)  
 C.H.Tsao & al, ICRC **26**, HE.1.1.04 (1999)

Slide from  
D. MAURIN



P.Ferrando & al, Phys.Rev.C **37**, 1490 (1988)  
 W.R.Webber, J.C.Kish, D.A.Schrier, Phys.Rev.C **41**, 1540 (1990)  
 W.R.Webber & al, ApJ **508**, 940 (1998-a)  
 W.R.Webber & al, ApJ **508**, 949 (1998-b)  
 W.R.Webber & al, Phys.Rev.C **58**, 3539 (1998-c)  
 W.R. Webber et al., ApJSS **144**, 153 (2003)

L.W.Townsend & al, NASA, Technical Paper **3310**, (1993)  
 F.A.Cucinotta, NASA, Technical Paper **3354**, (1993)  
 J.P.Bondorf & al, Phys.Rep. **257**, 133 (1995)  
 J.W.Wilson & al, NASA, Technical Paper **3533**, (1995)  
 F.A.Cucinotta & al, NASA, Technical Paper **3594**, (1996)  
 C.R.Ramsey & al, Phys.Rev.C **57**, 982 (1998)  
 Zeitlin et al., Phys. Rev. C **77**, 034605 (2008)  
 Zeitlin et al., AdSR **46**, 728 (2010)  
 Zeitlin et al., Phys. Rev. C **83**, 034909 (2011)

$\rightarrow$  Webber *et al.* is the one generally used in the field (but for  $Z < 3$  nuclei)  
 with claimed uncertainties  $< 10\%$  (fragments from Li  $\rightarrow$  O) or  $< 20\%$  (from Fe)

NB: it is not straightforward to go from nuclear data/models to X-sec for GCRs



# Transport equation in diffusion models

$$\begin{aligned}
 & \text{Diffusion} \quad \text{Convection} \\
 & -\vec{\nabla} \left[ K \vec{\nabla} N^j(E) - \vec{V}_c N^j(E) \right] - \Gamma^j N^j \\
 & - \frac{(\vec{\nabla} \cdot \vec{V}_c)}{3} \frac{\partial}{\partial E} \left[ \frac{p^2}{E} N^j(E) \right] = Q^j(E) + \\
 & \frac{\partial}{\partial E} \left[ -b_{\text{tot}}(E) N^j(E) + \beta^2 K_{pp} \frac{\partial N^j(E)}{\partial E} \right] \\
 & \text{Energy losses (EM)} \quad \text{Reacceleration}
 \end{aligned}$$

**Destruction on ISM**

CR sources: primaries,  
**secondaries**  
(spallations)

1. **DESTRUCTION:**  $\Gamma = n_{\text{ISM}} v \sigma_R$ ,  $\sigma_R = \sigma^{\text{tot}} - \sigma^{\text{tot}}_{\text{el}}$

2. **SOURCES:**  $\bar{Q}^j \equiv q_0^j Q(E) \hat{q}_i + \sum_k^{m_k > m_j} \tilde{\Gamma}^{kj} N_i^k(0)$

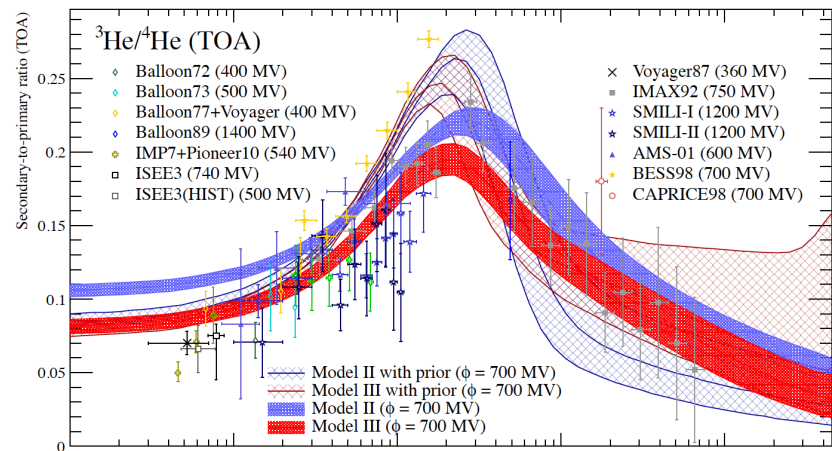
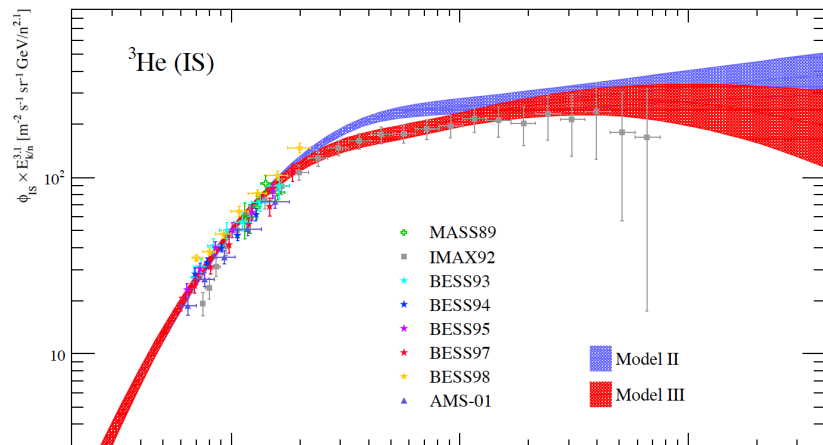
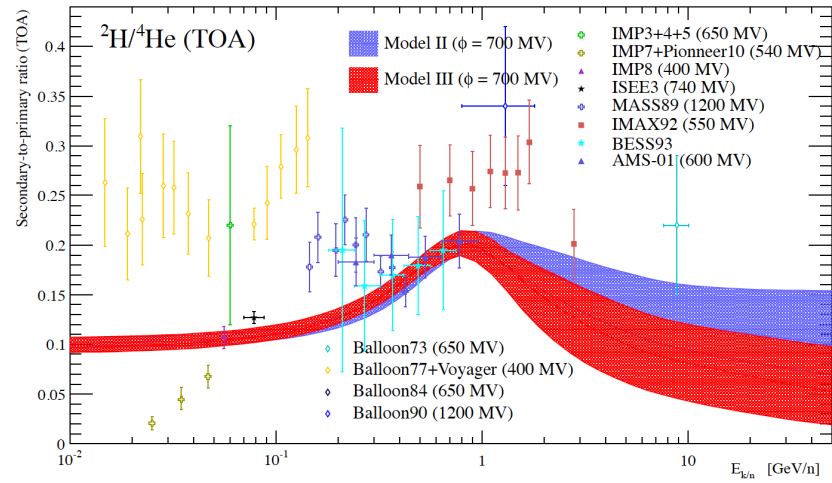
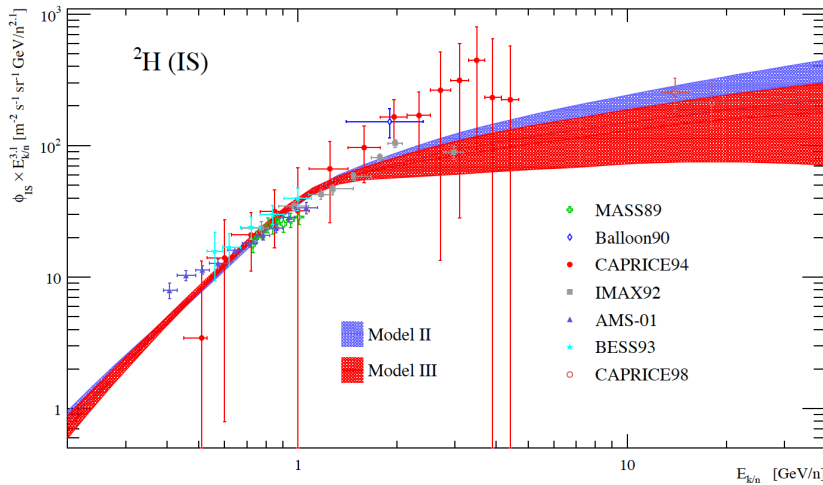
Primary  
production  
in SNR

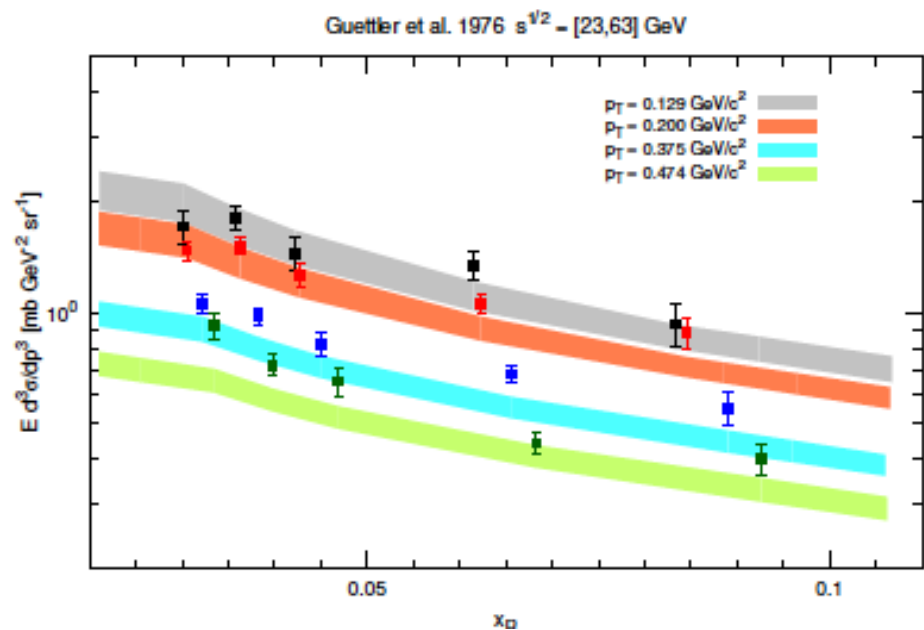
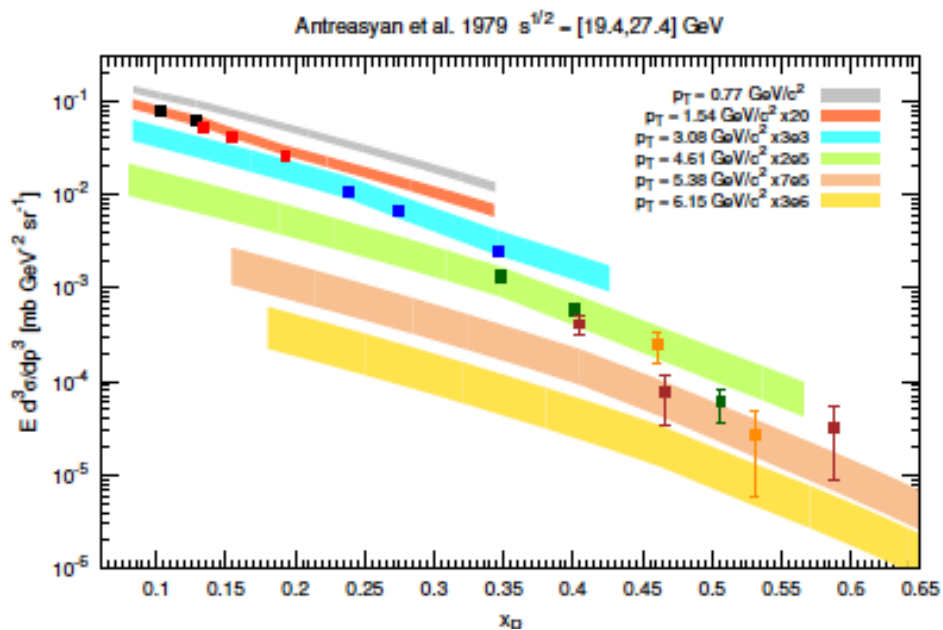
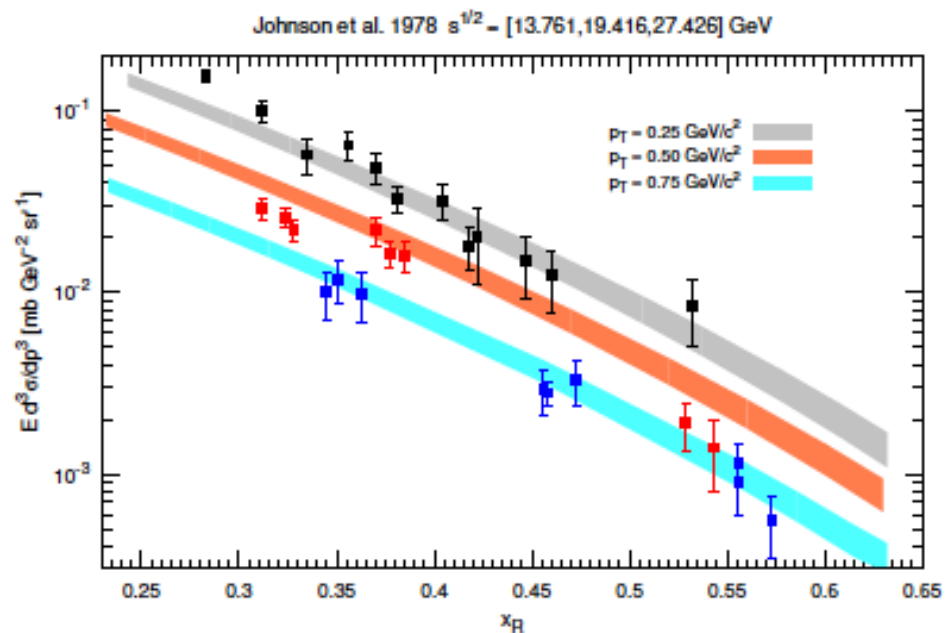
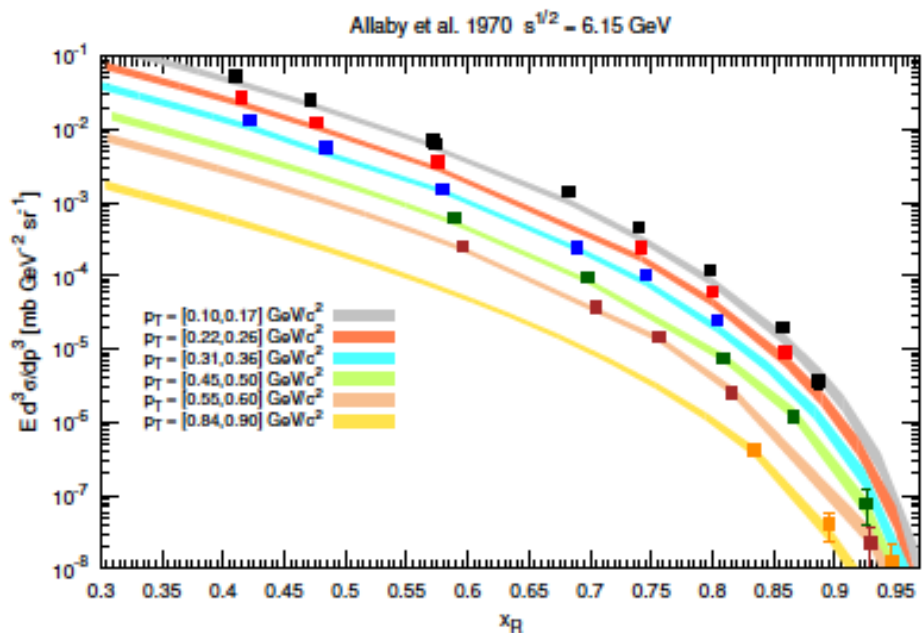
Secondary production  
by fragmentation  
of heavier nuclei

# $Z \leq 2$ Nuclei

Coste, Derome, Maurin, Putze A&A 2012

$^1\text{H}$ ,  $^2\text{H}$ ,  $^3\text{He}$ ,  $^4\text{He}$  almost as powerful as B/C  
Noticeable effort on reliable cross sections





# The role of helium nuclei

Kachelriess, Moskalenko & Ostapchenko 2015

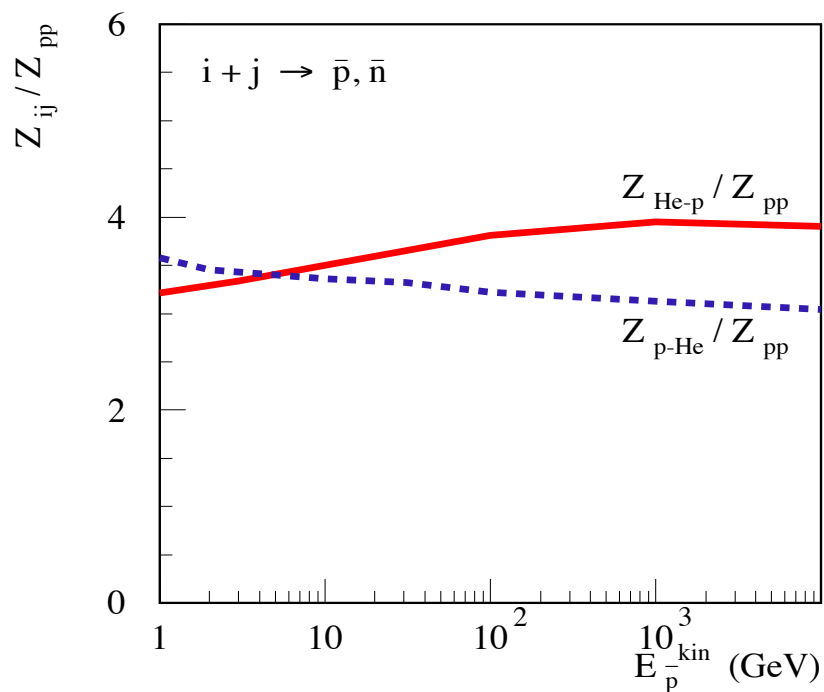
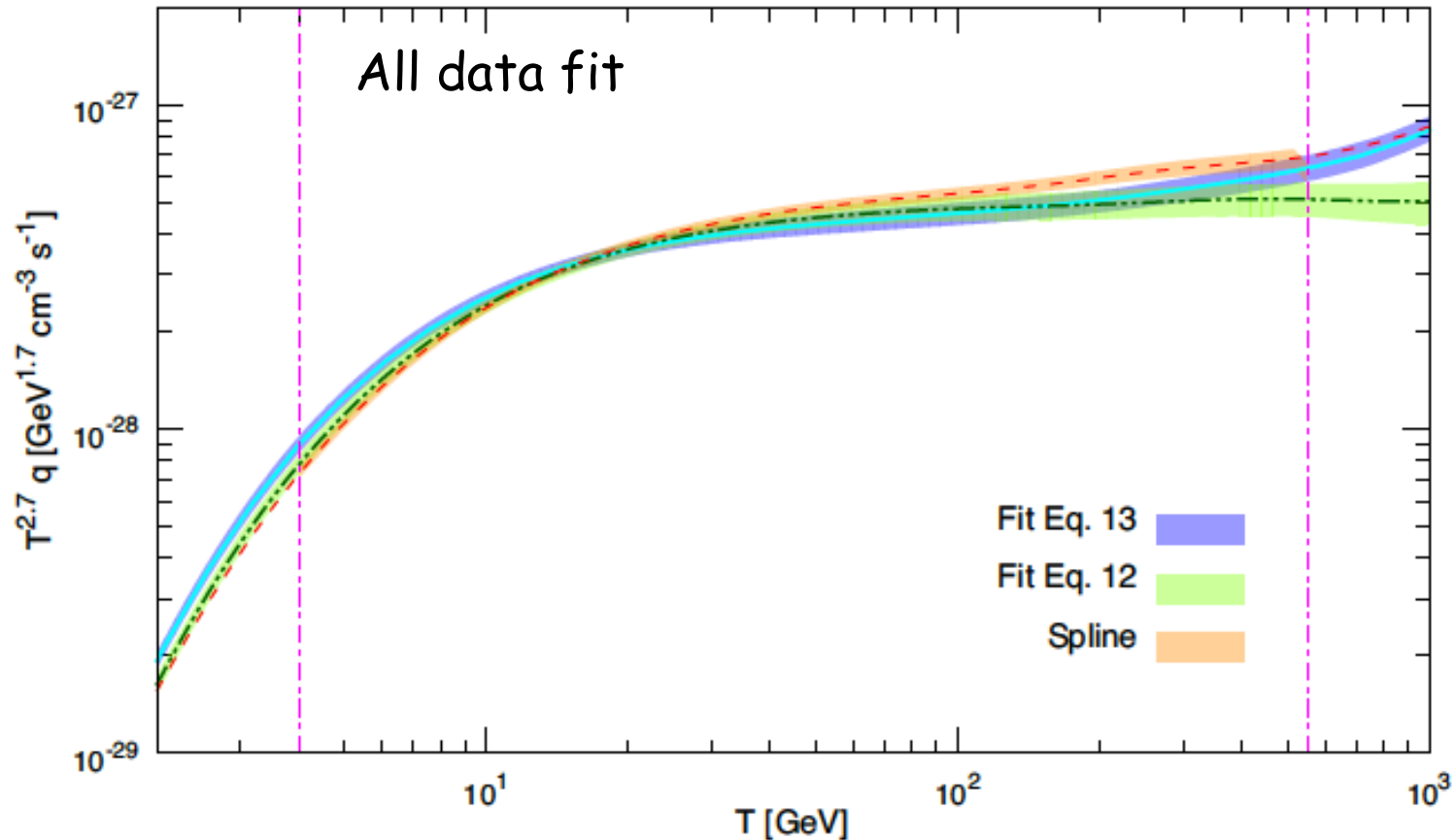


FIG. 7.— Energy dependence of the enhancement of the He  $p$  (solid, red) and  $p$  He (dashed, blue) contributions to the antiproton spectrum, relative to the  $pp$ -case,  $Z_{\bar{p}}^{ij}(E_{\bar{p}}, \alpha)/Z_{\bar{p}}^{pp}(E_{\bar{p}}, \alpha)$  (plotted as a function of  $E_{\bar{p}}^{\text{kin}}$ ), for  $\alpha = 2.6$ .

Note: The cosmic  $p$  flux is  $\sim 10$  times higher than He flux

# Antiproton source spectrum from p-p channels



$$E \frac{d^3\sigma}{dp^3} = \sigma_{\text{in}}(s)(1 - x_R)^{C_1} e^{-C_2 x_R} \quad (12)$$

$$\left[ C_3(\sqrt{s})^{C_4} e^{-C_5 p_T} + C_6(\sqrt{s})^{C_7} e^{-C_8 p_T^2} \right].$$

$$E \frac{d^3\sigma}{dp^3} = \sigma_{\text{in}}(s)(1 - x_R)^{C_1} e^{-C_2 x_R} \left[ C_3(\sqrt{s})^{C_4} e^{-C_5 p_T} + \right.$$

$$\left. C_6(\sqrt{s})^{C_7} e^{-C_8 p_T^2} + C_9(\sqrt{s})^{C_{10}} e^{-C_{11} p_T^3} \right], \quad (13)$$

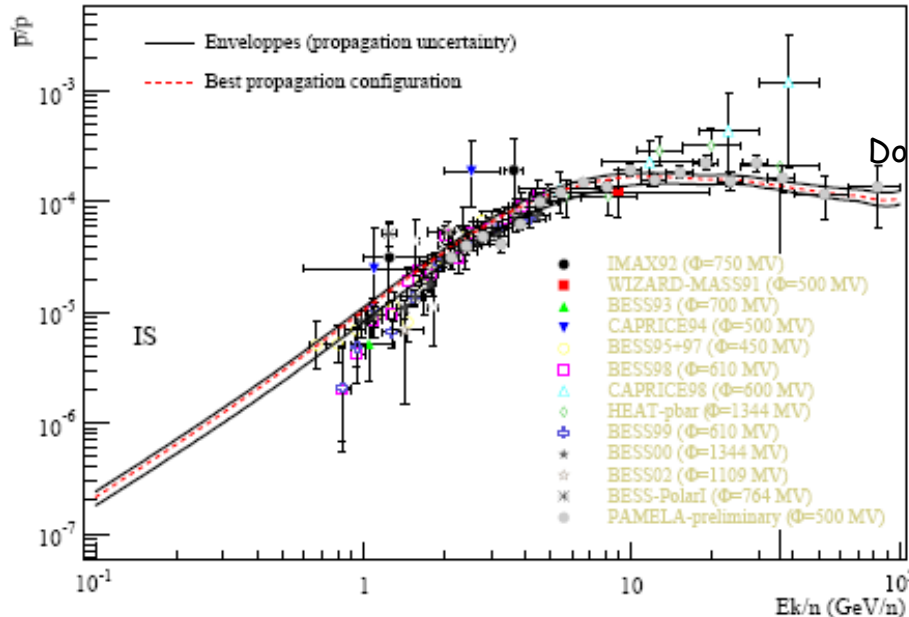
Different analytical functions give similar chi2, but different extrapolation out of validity ranges → uncertainties at low and high energies



# Antiproton in CRs: data and models

Theoretical calculations compatible with CR models

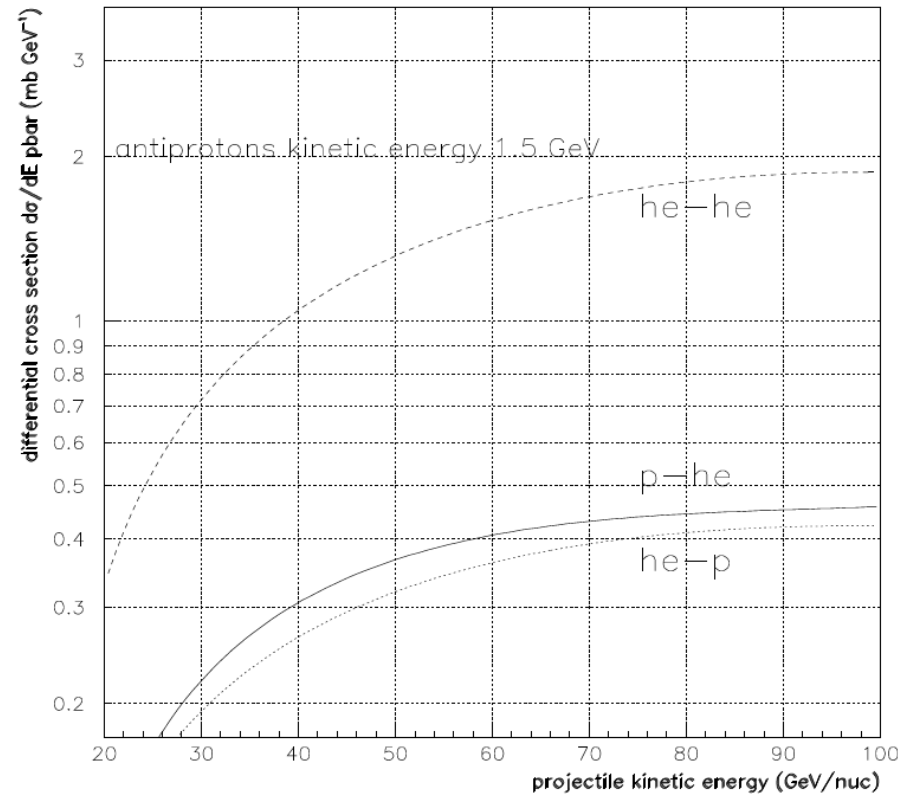
Powerful tool for DARK MATTER models



NO need for new phenomena (astrophysical / particle physics)

**AMS-02 data now at high accuracy 1 - 450 GeV**

# Differential antiproton cross section



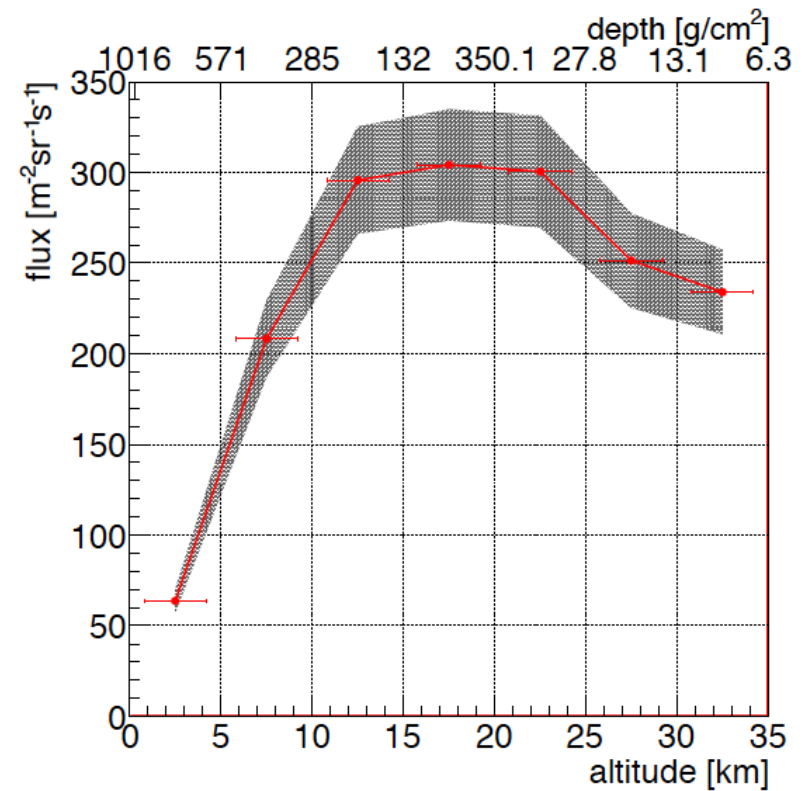
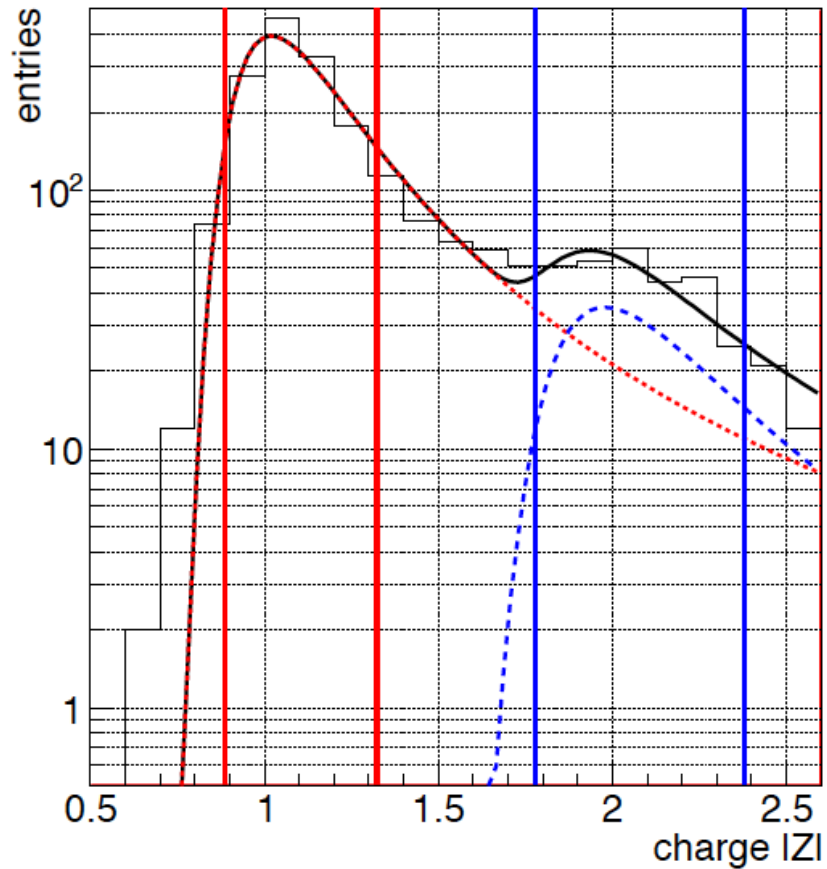
Given that:

1. In the ISM  $n_H=0.9/\text{cm}^3$ ,  $n_{He}=0.1/\text{cm}^3$ ,
2. the cosmic p flux is  $\sim 10$  higher than He

→ the main production channel involving He is  $p_{CR}-He_{ISM}$

# GAPS prototype flight

P. von Doetinchem et al. 1307.3538



# Propagation of $e^+$ & $e^-$

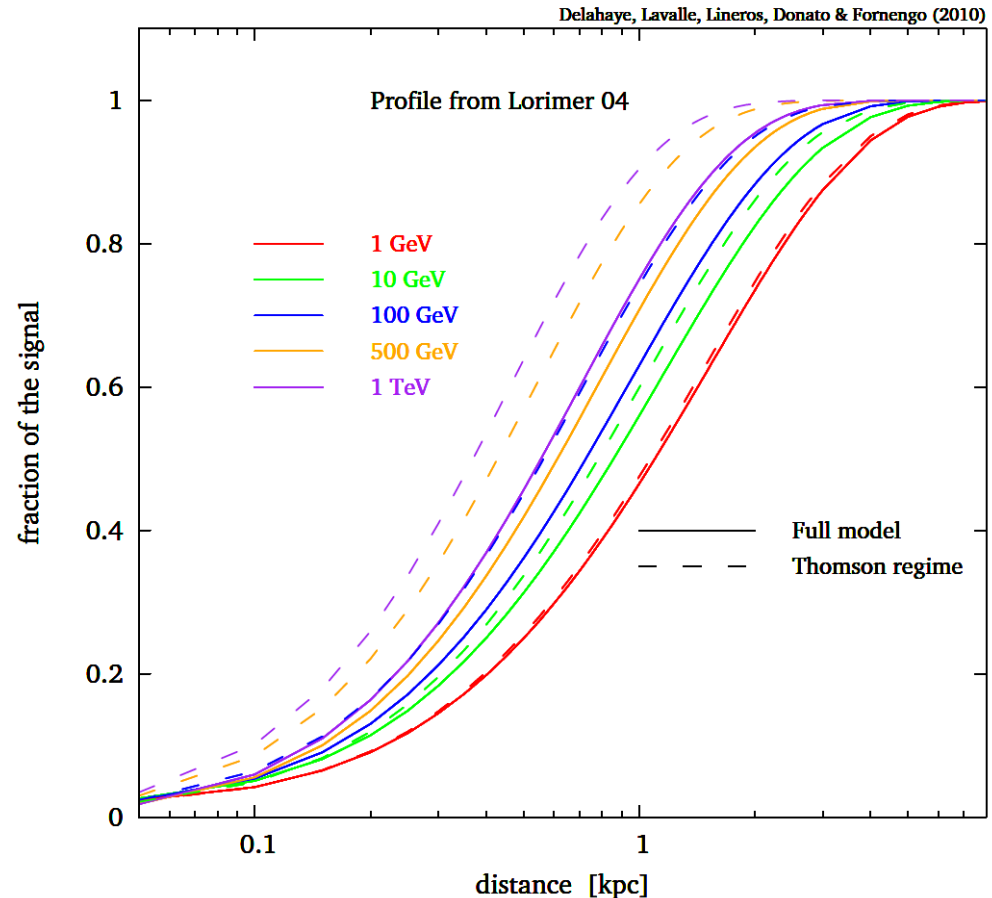
Delahaye, Laval, Lineros, FD, Fornengo, Salati, Taillet A&A 2009

Diffusive semi-analytical model:  
Thin disk and confinement halo  
Free parameters fixed by B/C

Above few GeV: only spatial  
diffusion and energy losses

**Energetic positrons &  
electron are quite local**

**Relativistic (KN) energy  
losses induce a longer  
propagation scale**



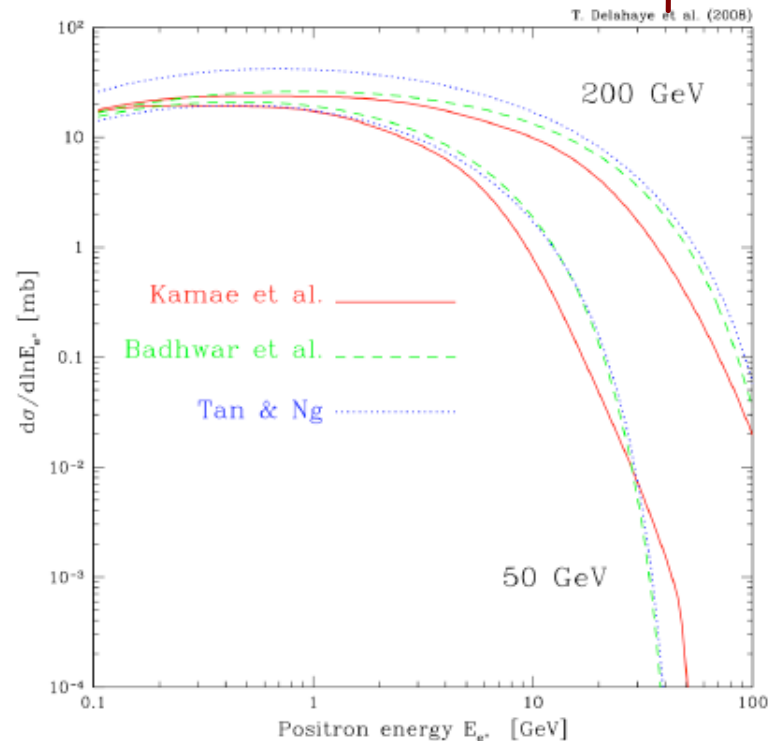
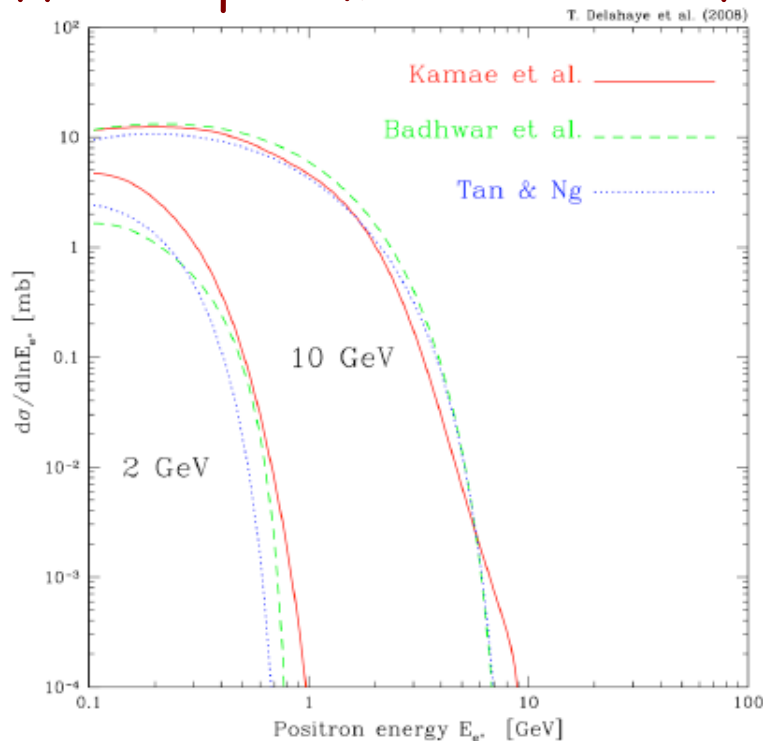
95%(80%) for 2kpc and  $E > 100(10)$  GeV

# Secondary positron production

Spallation of proton and helium nuclei on the ISM (H, He)

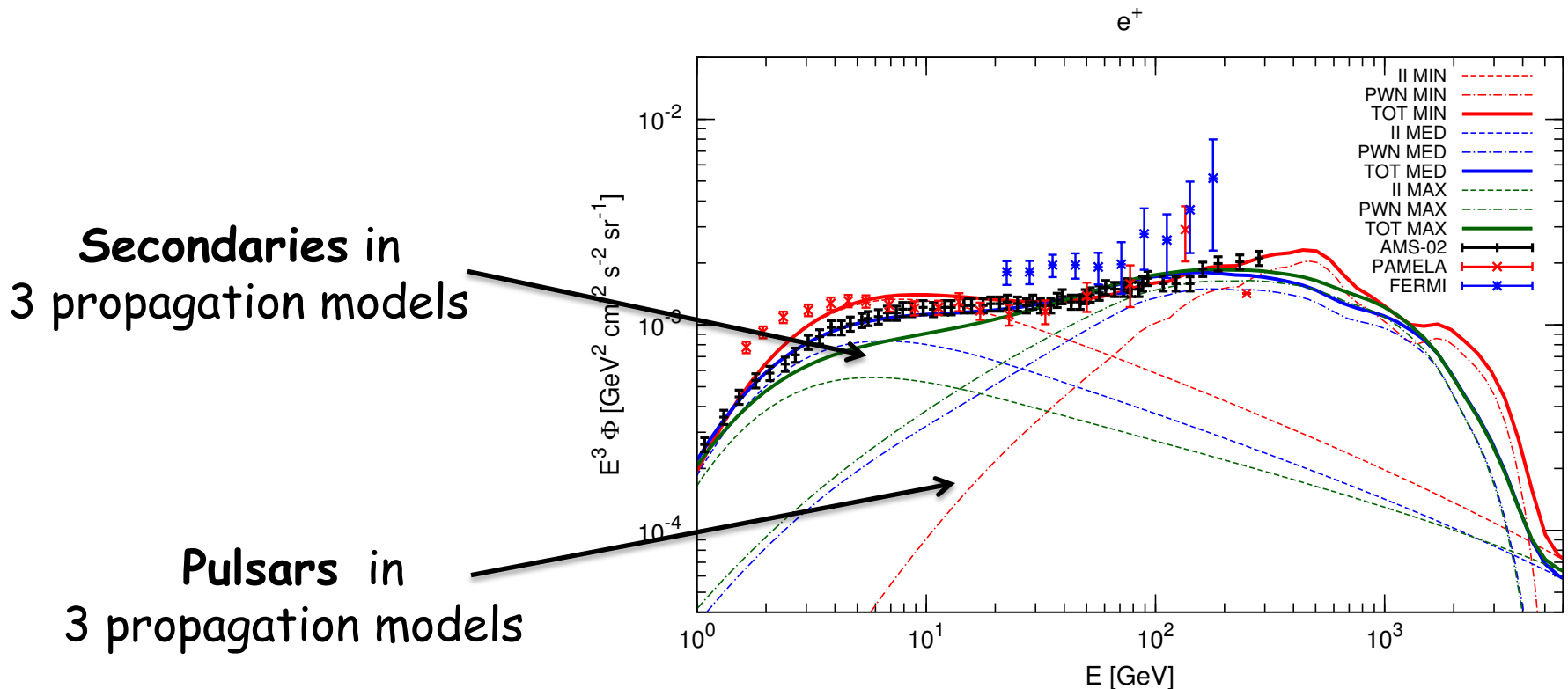
- $p+H \rightarrow p+\Delta^+ \rightarrow p+\pi^0$  &  $n+\pi^+$  (mainly below 3 GeV)
- $p+H \rightarrow p+n+\pi^+$
- $p+H \rightarrow X + K^\pm$

Different parameterizations of cross sections and incident p energy



# Low energy positrons: the power of AMS precision data

Di Mauro, FD, Fornengo, Vittino JCAP 2014; Lavalley, Maurin, Putze PRD 2014



**AMS POSITRON data at low energy are so precise  
to start constraining the Milky Way models (as well as B/C)!**

# The case for gamma rays

(very briefly)



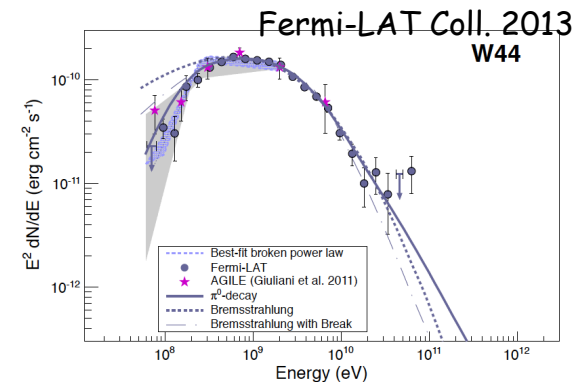
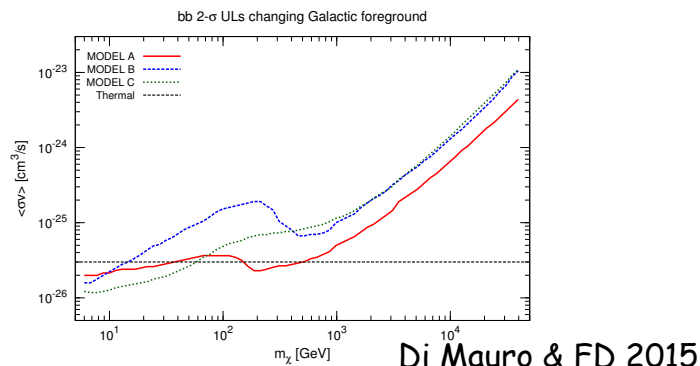
# $p + \text{He} \rightarrow \pi^0 \rightarrow 2\gamma$ : galactic foreground in the Fermi-LAT data

Contextually to  $p + \text{He} \rightarrow$  antiprotons, it would be of very interesting to measure also the photons (gamma rays) coming from the hadronization processes via  $\pi^0$  decay.

This process occurs in the galactic disk and enters the calculation of the galactic emission, crucial for understanding:

1. **DARK MATTER annihilation**  
(Galaxy, dwarf spheroidal galaxies, ..)

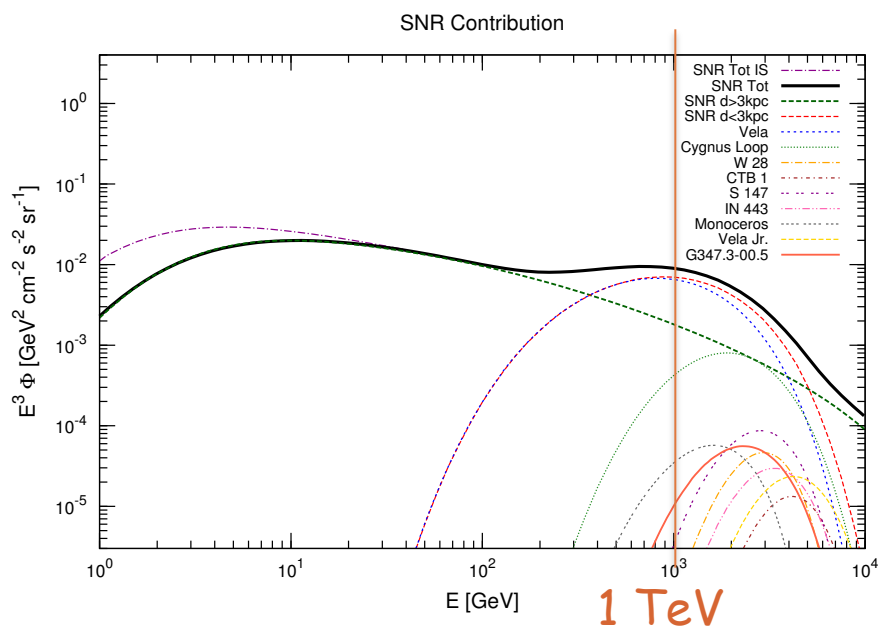
2. **ORIGIN of COSMIC RAYS**  
galaxy emission, SNRs, ..



# Astrophysical sources of $e^+ e^-$

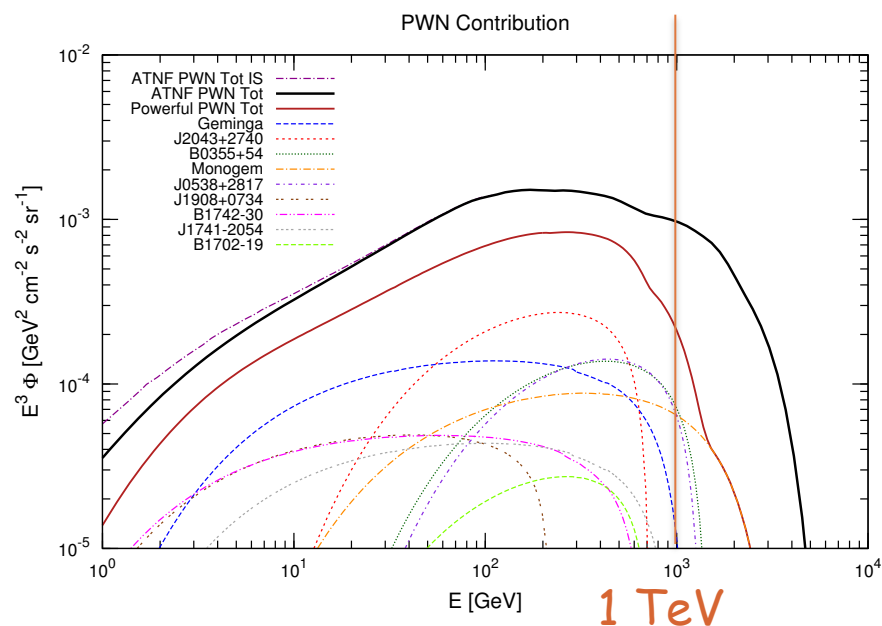
Di Mauro, FD, Fornengo, Vittino JCAP 2014

## Electrons



Supernova remnants

## Positrons



Pulsars

# Partial contributions from CR nuclei

Kachelriß+2015

$$\epsilon_{ij}^{\bar{p}}(E_{\bar{p}}) = \frac{q_{\bar{p}}^{ij}(E_{\bar{p}})}{q_{\bar{p}}^{pp}(E_{\bar{p}})}$$

$$q_{\bar{p}}^{ij}(E_{\bar{p}}) = n_j \int_{E_{\text{thr}}(E_{\bar{p}})}^{\infty} dE \frac{d\sigma^{ij \rightarrow \bar{p}}(E, E_{\bar{p}})}{dE_{\bar{p}}} I_i(E)$$

11

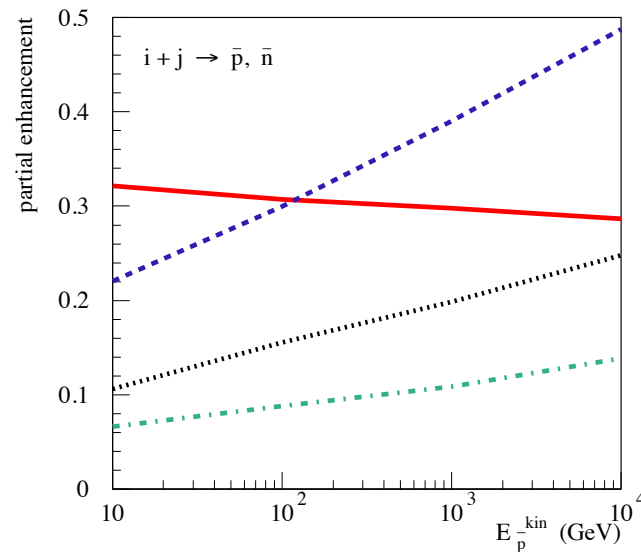
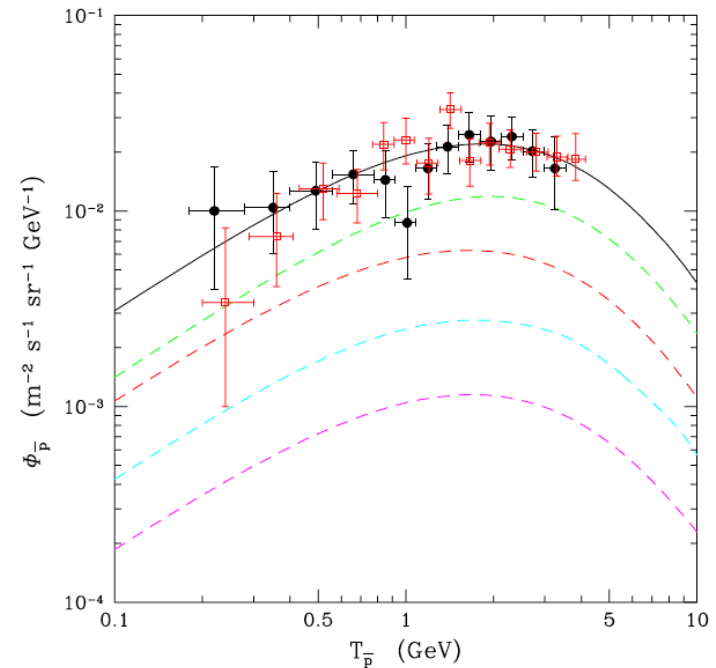
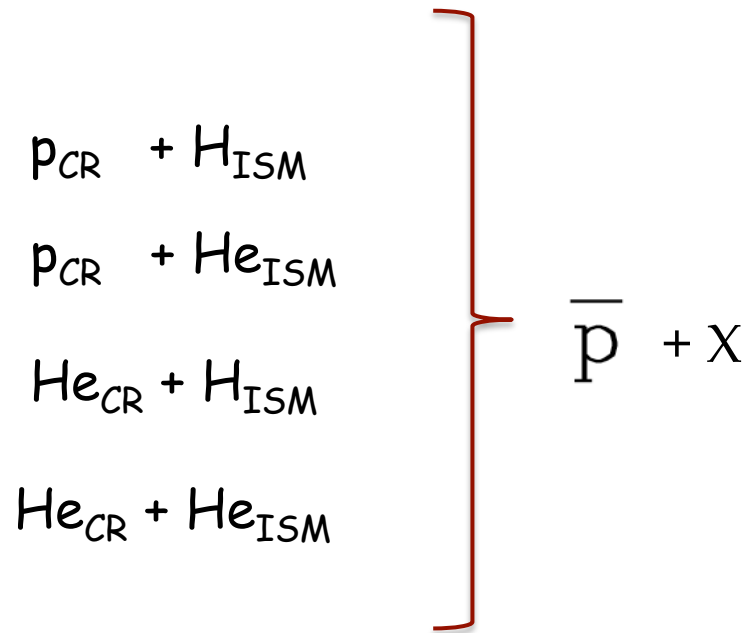


FIG. 8.— Energy dependence of partial contributions  $\epsilon_{ij}^{\bar{p}}(E_{\bar{p}})$  to the nuclear enhancement factor from different interaction channels:  $p$  He (solid, red), He  $p$  (dashed, blue), He He (dot-dashed, green), and all others (dotted, black); the CR composition given in Table 2 is used.

The contribution of He nuclei, in CRs and in ISM, is similar (50%) of the p-p one

# Secondary antiprotons in cosmic rays (CR)

Produced by spallation reactions on the interstellar medium (ISM)

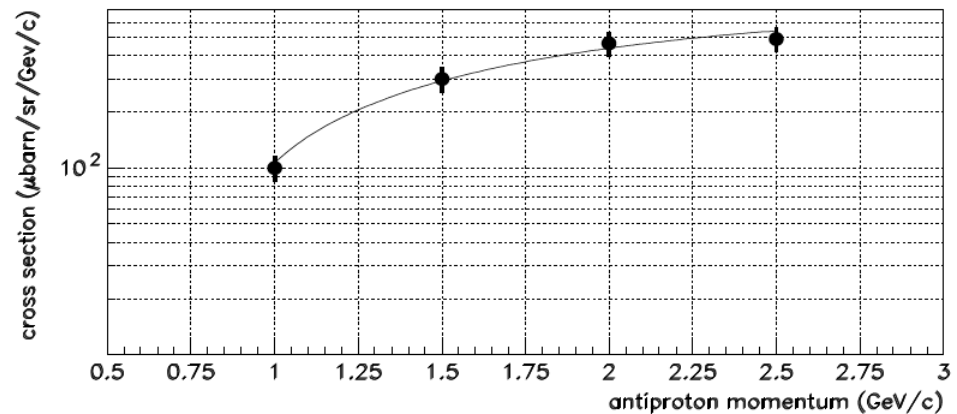
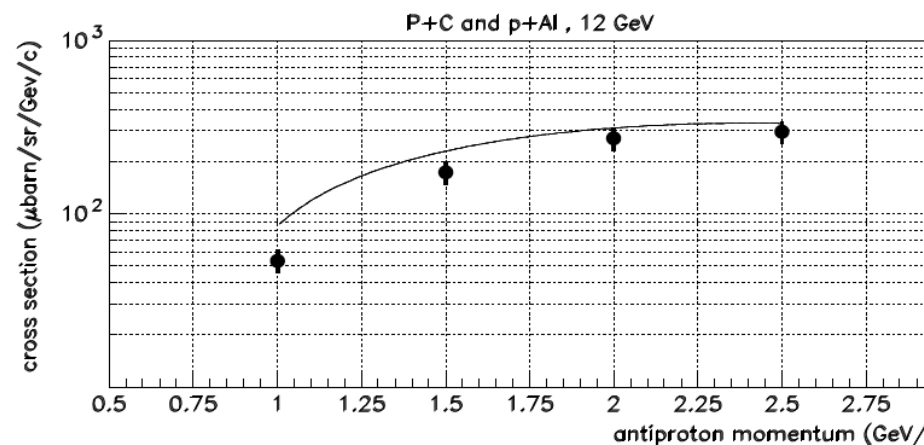


The only measured cross section is  $p\text{-}p \rightarrow \bar{p} + X$

**ALL CROSS SECTIONS  
INVOLVING He (projectile or target)  
ARE DERIVED FROM OTHER DATA!!**

# Antiproton production from p and He

1.  $p+p: \sigma_{p+p \rightarrow \text{antiprotons}}$  NA49 data analysis  
(Di Mauro, FD, Goudelis, Serpico PRD 2014; Kappl, Winkler JCAP 2014)  
or analytical expression (Tan & Ng, PRD26 1982, 1983)
2.  $p + \text{He}, \text{He}+p, \text{He}+\text{He}: \sigma(p + \text{He} \rightarrow \text{antiprotons})$ : NO DATA  
→ derived from MonteCarlo simulations, i.e. DTUNUC  
verified on  $p+C, p+Al$  and heavier nuclei (Duperray et al. 2003, 2005; Donato et al. ApJ 563 (2001)172 )



Data from Sugaya+1998; fit: DTUNUC

Donato+ ApJ2001