

Neutrino-Flavoured Sneutrino Dark Matter

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Motivation

- SUSY with R-parity has many nice features
 - Solves the hierarchy problem
 - Precision gauge coupling unification
 - Natural dark matter candidate
- Is the MSSM enough?
 - No explanation for neutrino masses
 - Can there be interesting structure in the dark sector (c.f. our sector)?
 - Is the neutralino the only good dark matter candidate?
- Can simple extensions to the MSSM answer some of these questions and still give a natural dark matter candidate?

Extending the MSSM...

hep-ph/0403067 March-Russell & West

- Add 3 right-handed neutrinos N and their super-partners $\tilde{n}$

$$\Delta\mathcal{L} = \int d^2\theta \left(\lambda_{ij} L_i N_j H_u + \frac{1}{2} M_N N_i N_i \right) \\ + m_{\tilde{n}}^2 |\tilde{n}_i|^2 + A \tilde{L}_i \tilde{n}_i h_u + \frac{1}{2} \lambda_{ij} M_B^2 \tilde{n}_i \tilde{n}_j + h.c.$$

- Important features:

- Right-handed neutrinos have **weak scale masses**: $M_N \sim \text{TeV}$
- Yukawa coupling **suppressed**: $\lambda_{ij} \sim 10^{-7} - 10^{-8}$
- A-term **weak scale** and **flavour diagonal**: $A \sim \text{TeV}$
- Small right-handed sneutrino lepton-number violating B-term **with flavour structure**: $\lambda_{ij} M_B^2 \sim (100\text{MeV})^2$

- These terms all arise from higher dimensional operators and are justified with an R-symmetry

The Sneutrino Spectrum

- Initially turn off A and B-terms
- 3 complex right and 3 complex left handed sneutrinos

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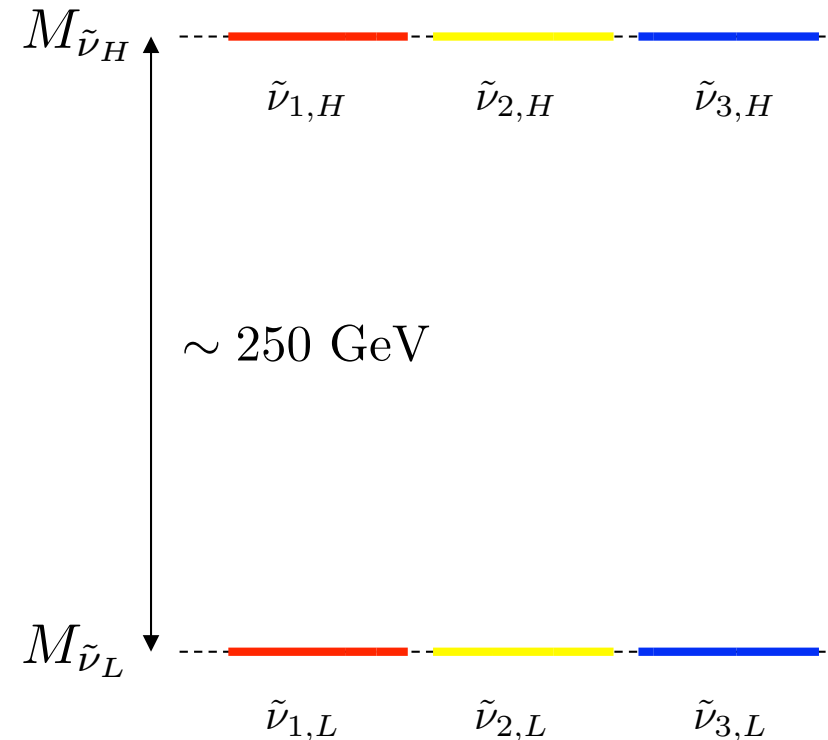
$$\tilde{\nu}_{i,H} = \cos \theta \tilde{\nu} + \sin \theta \tilde{n}$$

$$\tilde{\nu}_{i,L} = -\sin \theta \tilde{\nu} + \cos \theta \tilde{n}$$

$$M_{\tilde{\nu}_H}^2 = \frac{M_L^2 \cos^2 \theta - M_R^2 \sin^2 \theta}{\cos 2\theta}$$

$$M_{\tilde{\nu}_L}^2 = \frac{M_R^2 \cos^2 \theta - M_L^2 \sin^2 \theta}{\cos 2\theta}$$

$$\tan 2\theta = \frac{2Av \sin \beta}{M_R^2 - M_L^2}$$

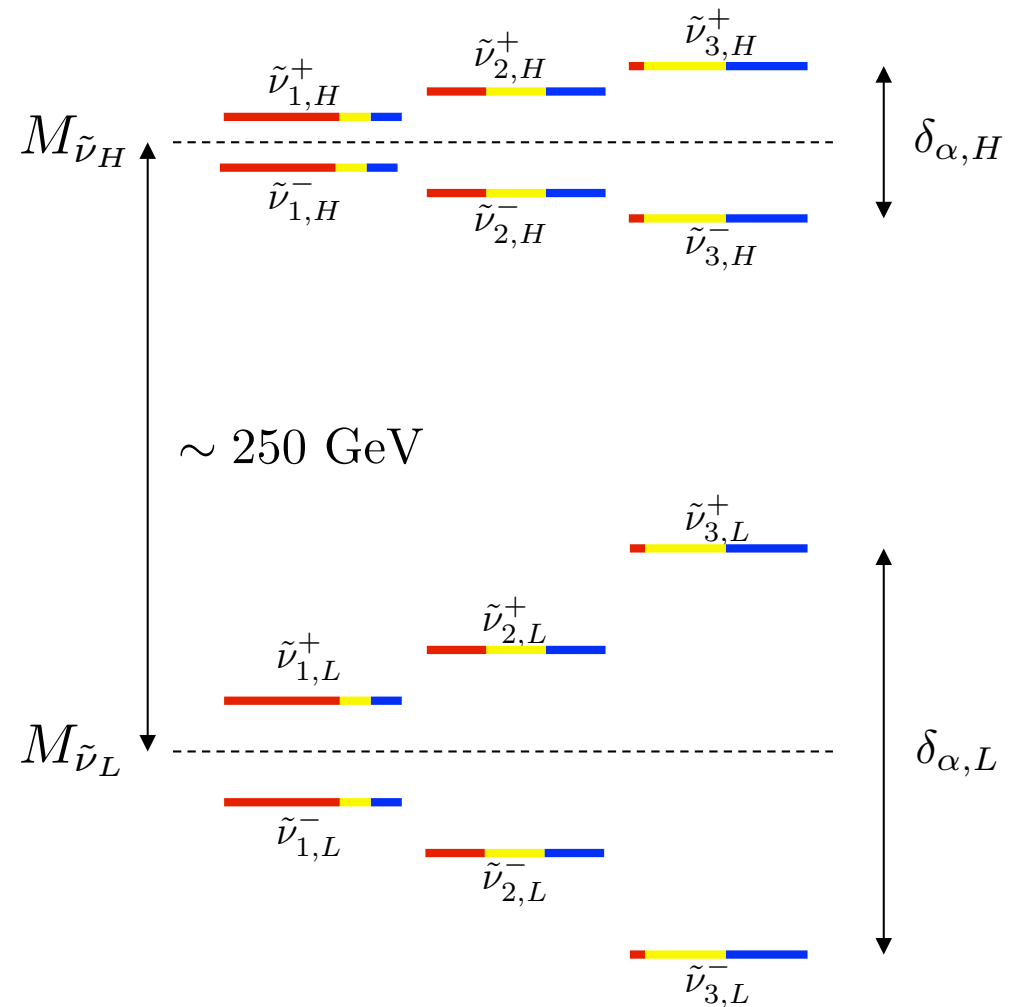


The Sneutrino Spectrum

- Initially turn off A and B-terms
- 3 complex right and 3 complex left handed sneutrinos
- Turn on A
- Mixes left and right sneutrinos
- Turn on B
- Splits spectrum into CP even and CP odd states
- Mixes **flavours**. Structure determined by: λ_{ij}

$$\delta_{\alpha,H} \sim 10 \text{ keV}$$

$$\delta_{\alpha,L} \sim 100 \text{ keV}$$



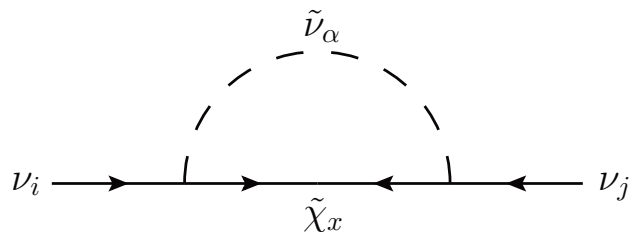
Generating Neutrino Properties...

- Neutrino masses will be generated in two ways
 - See-saw mechanism gives a **small** contribution:

$$M_{\nu_{ij}} = \frac{v^2 \sin^2 \beta}{M_N} \lambda_{ik} \lambda_{kj} \sim 10^{-5} \text{ eV}$$

- Recall $\lambda_{ij} \sim 10^{-7} - 10^{-8}$

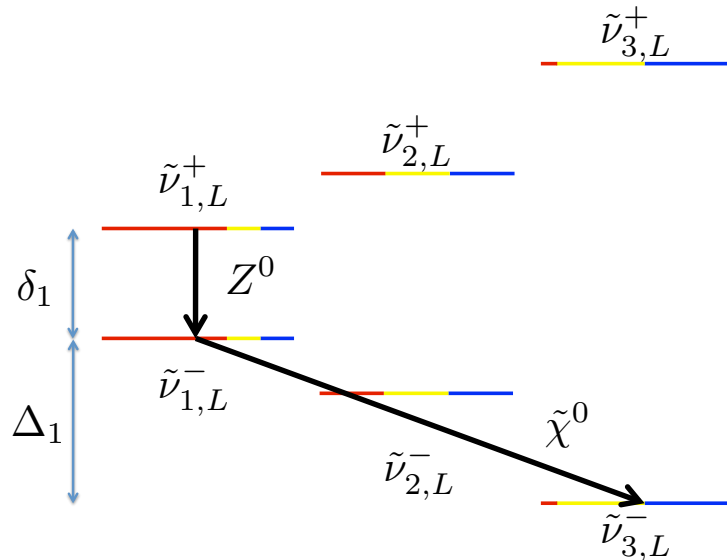
- **Dominate** contribution from sneutrino-neutralino loop:



$$M_{\nu_{ij}} \sim \lambda_{ij} \times f(M_{\tilde{\chi}_x^0}, M_{\tilde{\nu}_L}, M_{\tilde{\nu}_L}, \theta)$$

- Notice that flavour structure determined by λ_{ij} . BUT recall that the sneutrino flavour structure is also determined by λ_{ij}
 - ✧ We have **neutrino-flavoured sneutrinos**!

Lifetimes of the lightest states



- Scalars can decay to pseudoscalars via a Z boson with lifetime:

$$\Gamma_Z^{\tilde{\nu}_{1,L}^+} \sim (9 \times 10^4 \text{ yrs})^{-1} \left(\frac{\sin \theta}{0.1} \right)^4 \left(\frac{\delta_1}{100 \text{ keV}} \right)^5$$

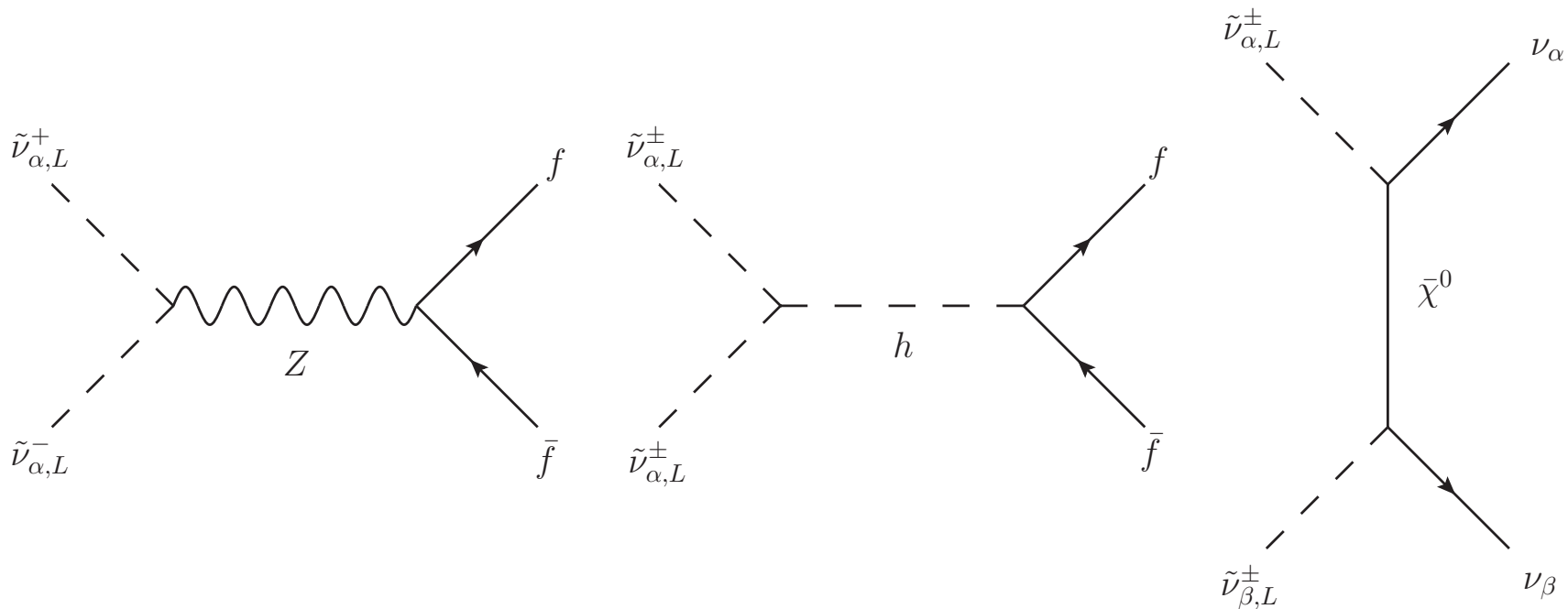
- Pseudoscalars can decay by neutralino exchange with lifetime:

$$\Gamma_{\tilde{\chi}}^{\tilde{\nu}_{1,L}^-} \sim (10^{10} \text{ yrs})^{-1} \left(\frac{\chi}{0.22} \right) \left(\frac{\sin \theta}{0.1} \right)^4 \left(\frac{\Delta_1}{25 \text{ keV}} \right)^5 \left(\frac{100 \text{ GeV}}{m_{\tilde{\nu}}} \right)^2 \left(\frac{200 \text{ GeV}}{m_{\chi_1}} \right)^2$$

- Depending on parameters, we can have a dark matter cocktail containing the three lightest pseudoscalars or just the lightest
- For this talk, we'll assume all have decayed to the lightest state

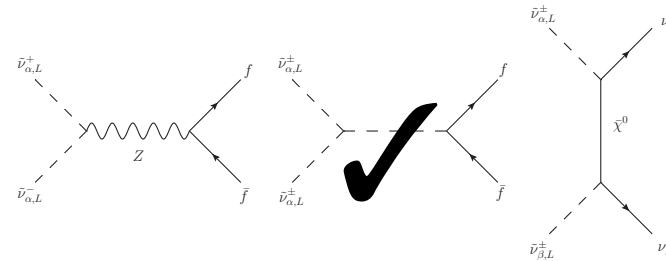
Can Sneutrinos Be The Dark Matter?

- Thermal Relic Density?
 - Main diagrams:

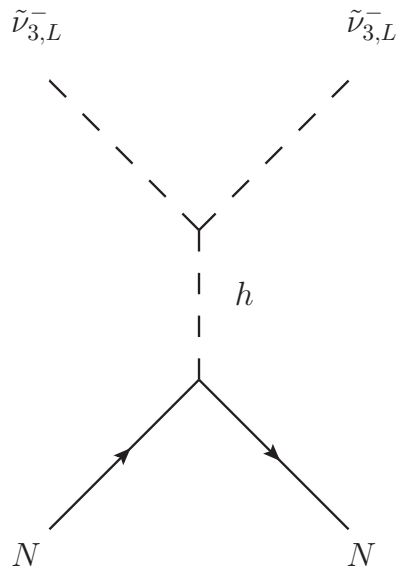


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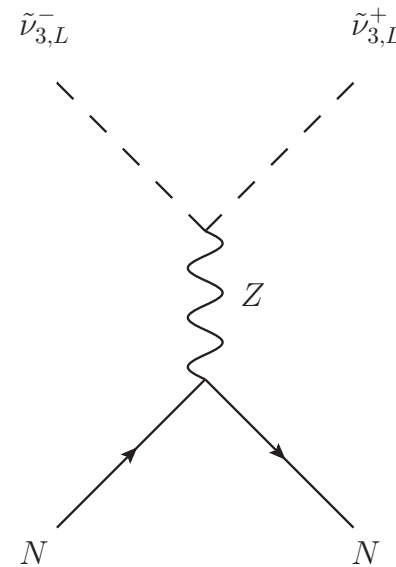
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- Direct detection limits?
 - Elastic scattering



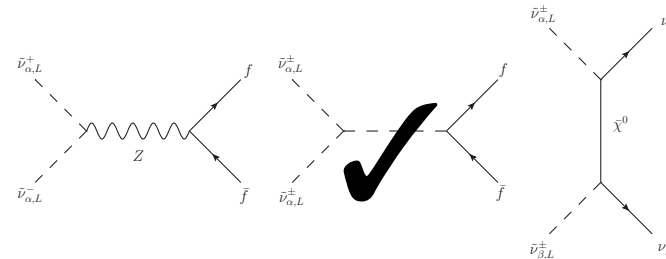
- Inelastic scattering



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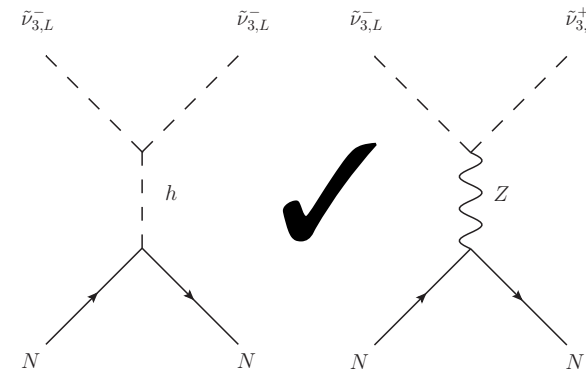
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- Indirect detection limits?

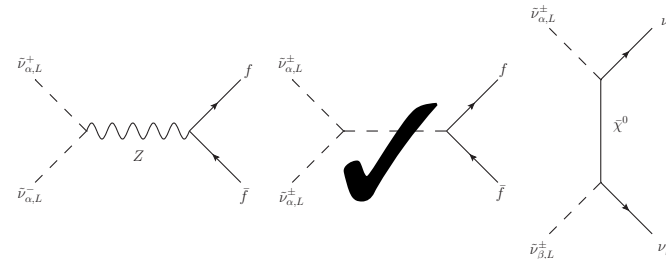
- Dark matter capture in the Sun
 - Limits from the measured neutrino flux at Super-Kamiokande

arXiv/0905.1333 Nussinov et al
arXiv/0905.1847 Menon et al

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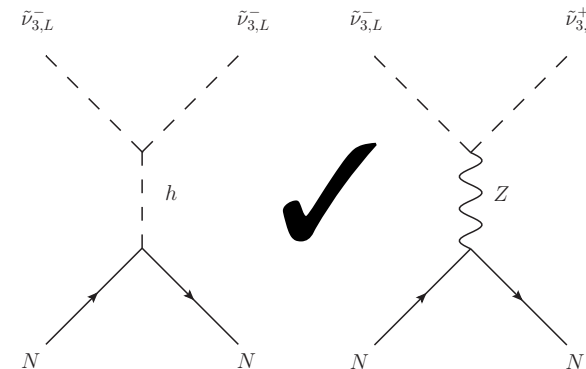
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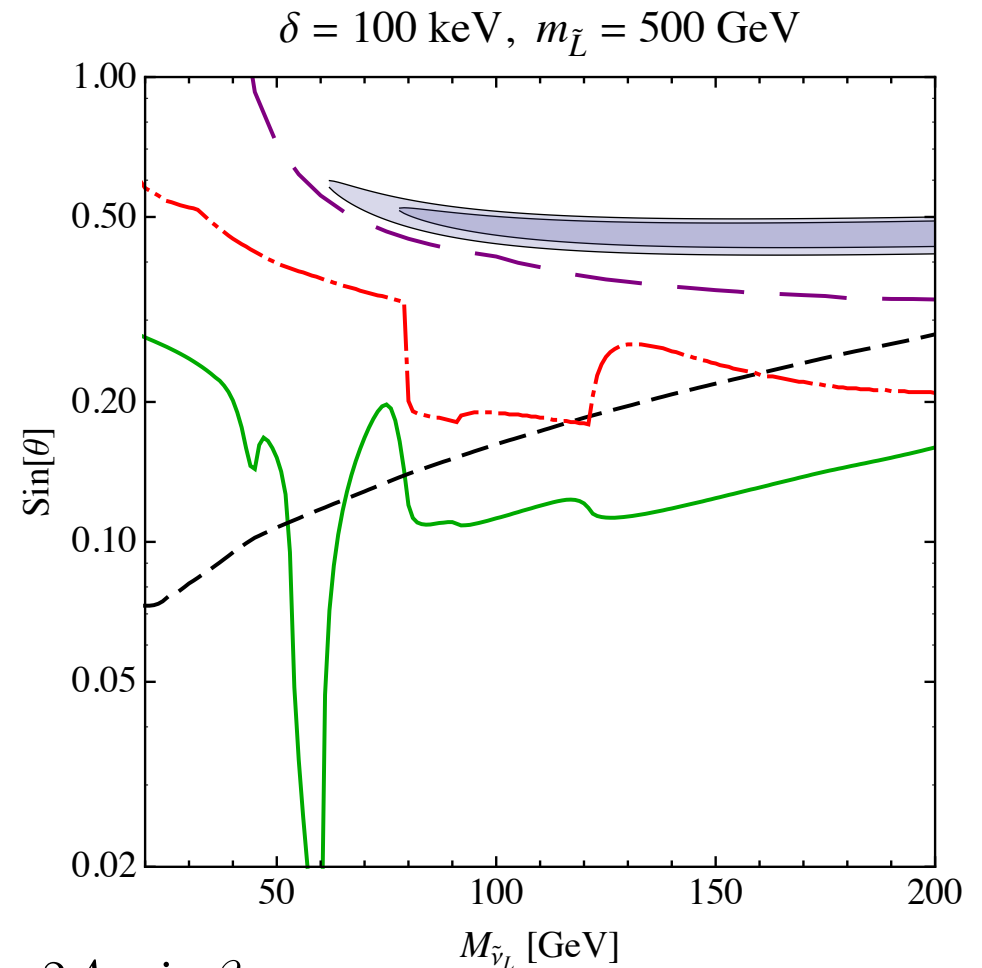
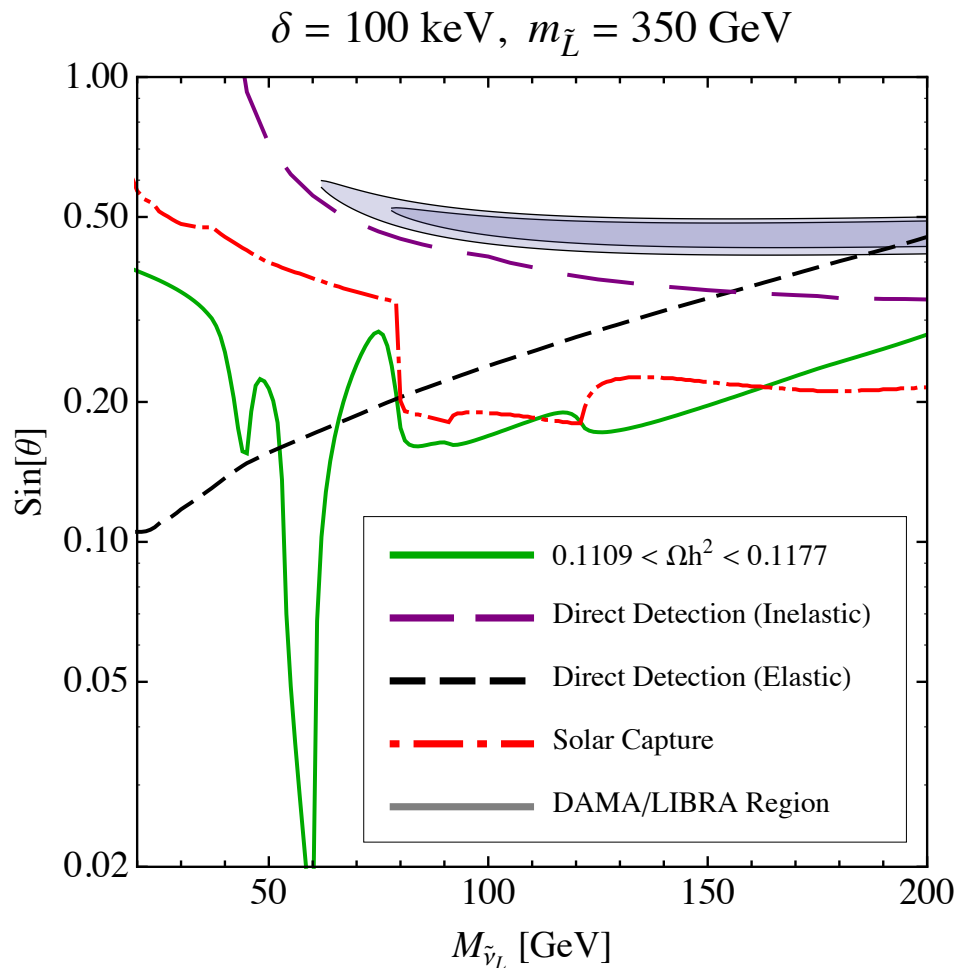


- Indirect detection limits?

- Dark matter capture in the Sun
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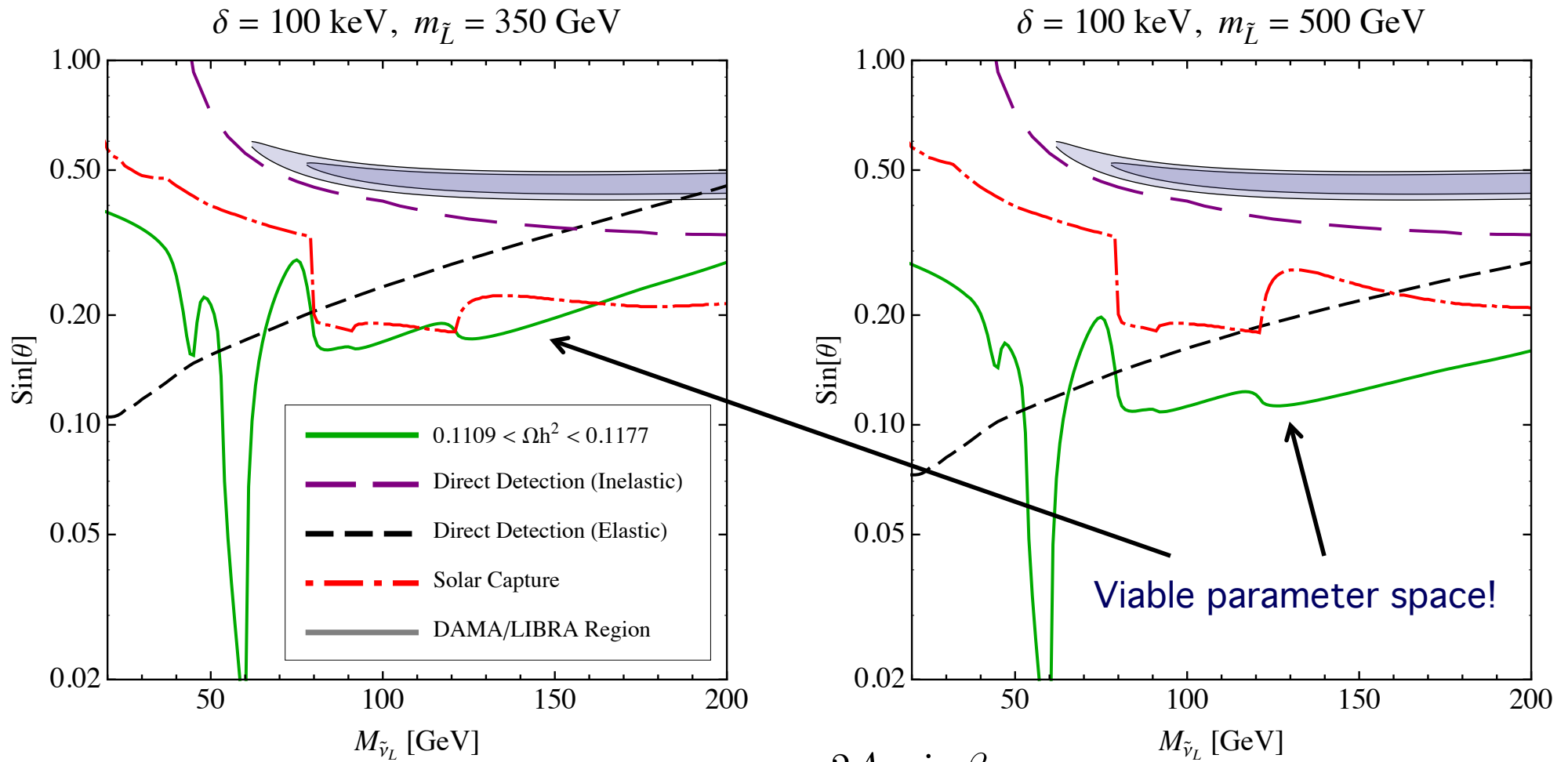
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Examples of parameter space I



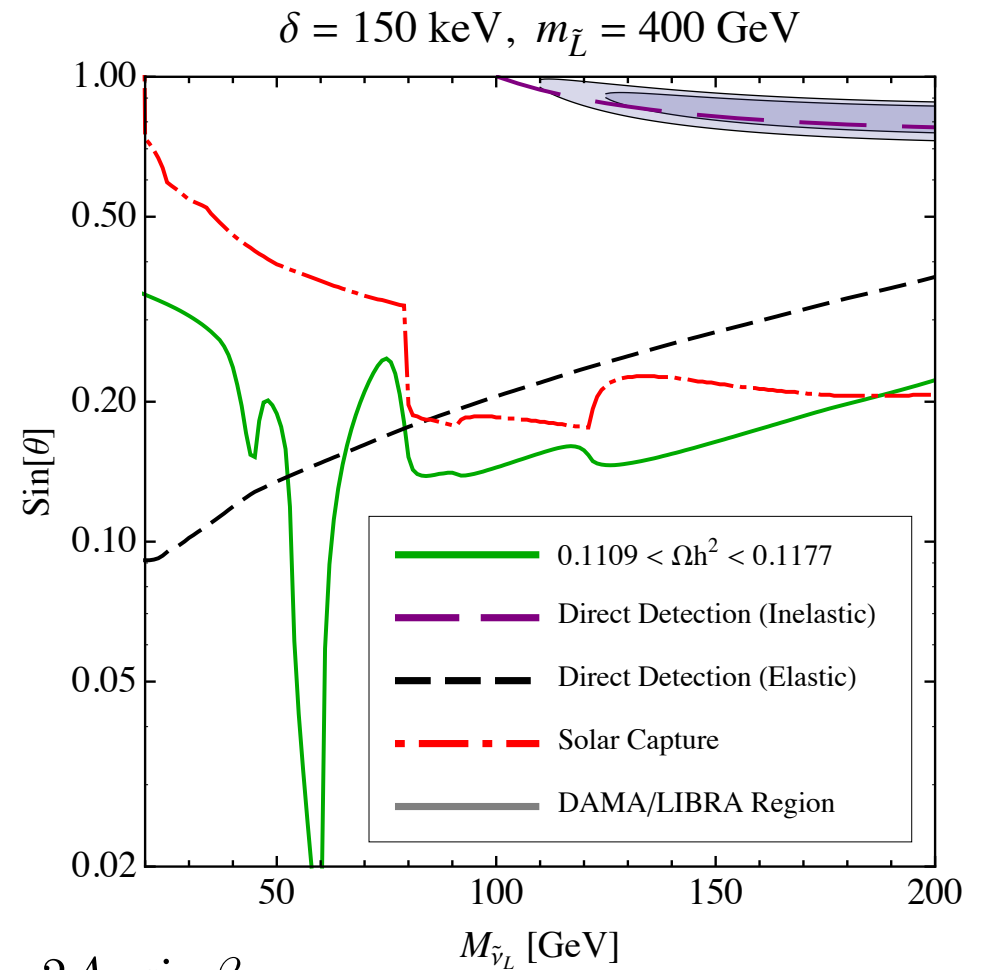
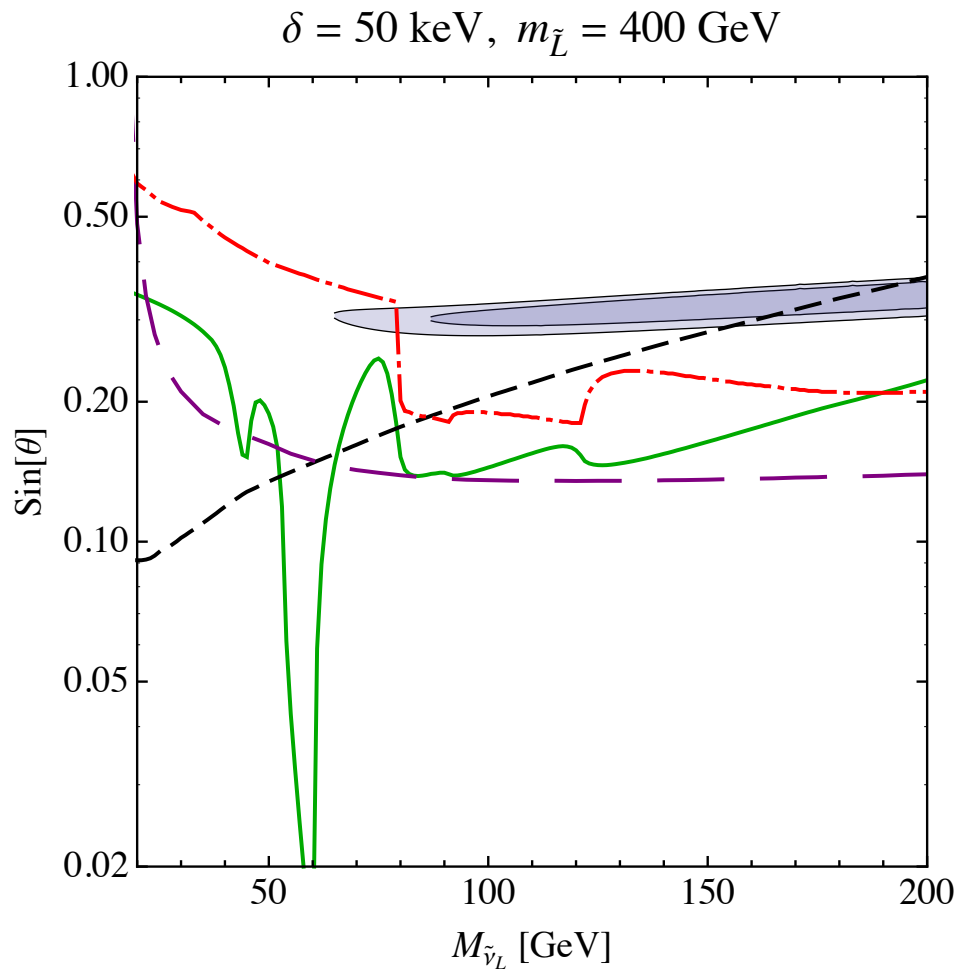
Recall: $\tan 2\theta = \frac{2Av \sin \beta}{M_R^2 - M_L^2}$

Examples of parameter space I



Recall:
$$\tan 2\theta = \frac{2Av \sin \beta}{M_R^2 - M_L^2}$$

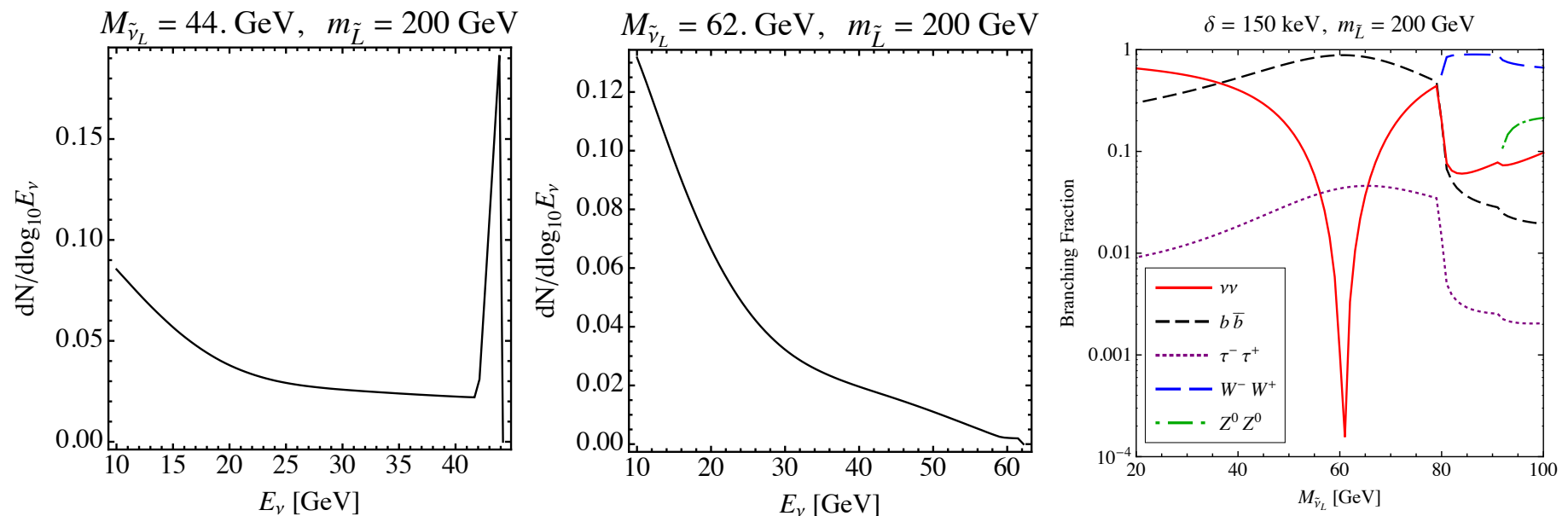
Examples of parameter space II



Recall:
$$\tan 2\theta = \frac{2A v \sin \beta}{M_R^2 - M_L^2}$$

Tests of the model/Future signals?

- Possible spectrum at neutrino telescopes



- Possible peak from direct annihilation into neutrinos?

Conclusions

- A simple extension to the MSSM can lead to a much richer structure in the dark matter sector
 - A link between dark matter and neutrino physics
 - Interesting direct detection signals?
 - Elastic scattering
 - Inelastic scattering
 - Interesting indirect signals at neutrino telescopes?