

# **Characterising the astrophysical S-factor for $^{12}\text{C} + ^{12}\text{C}$ fusion with wave-packet dynamics**

**ALEXIS DIAZ-TORRES**

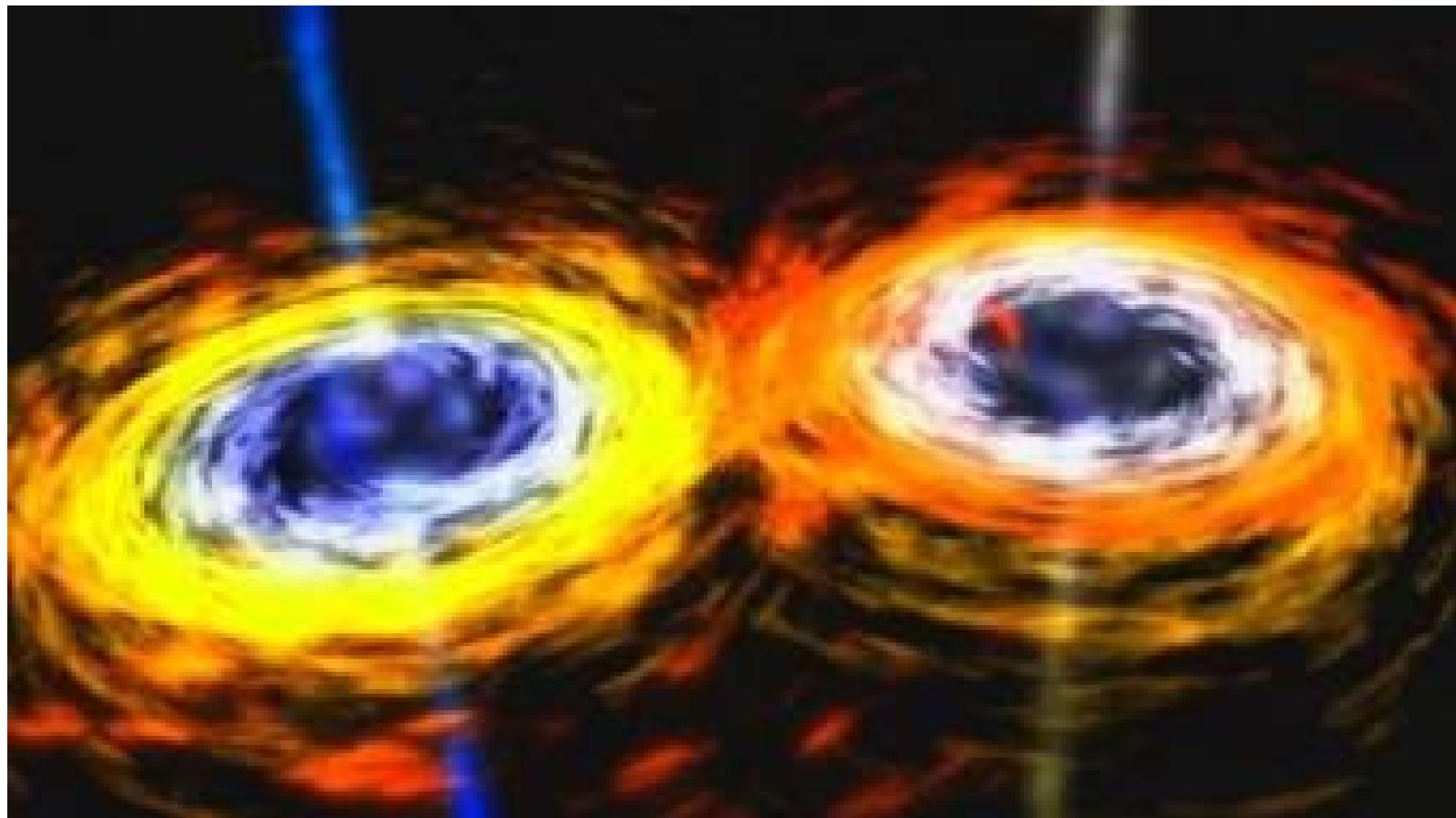


UNIVERSITY OF  
**SURREY**

**In collaboration with Michael Wiescher**



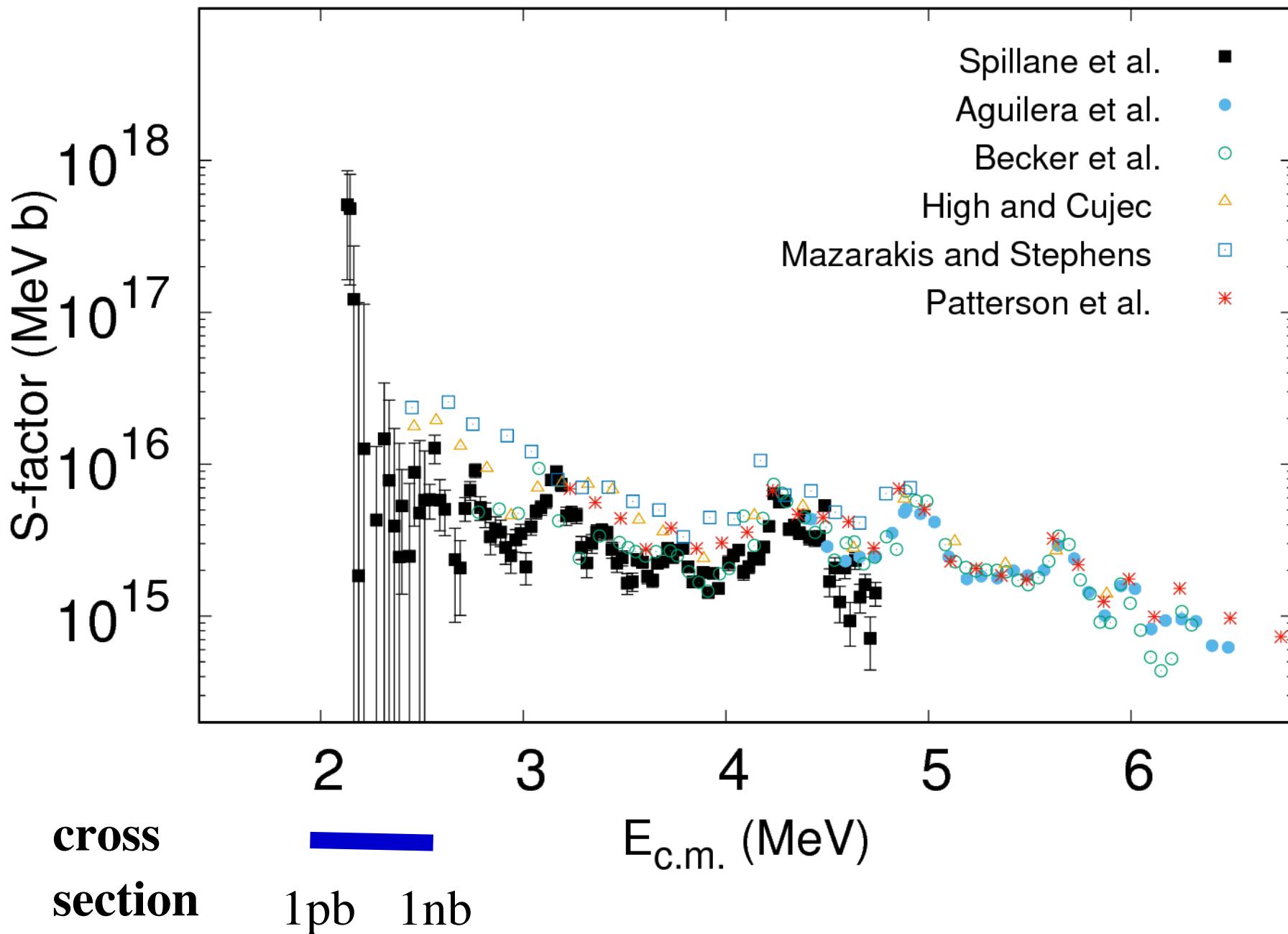
# How do two $^{12}\text{C}$ nuclei fuse at sub-barrier energies?



Picture taken from BBC News

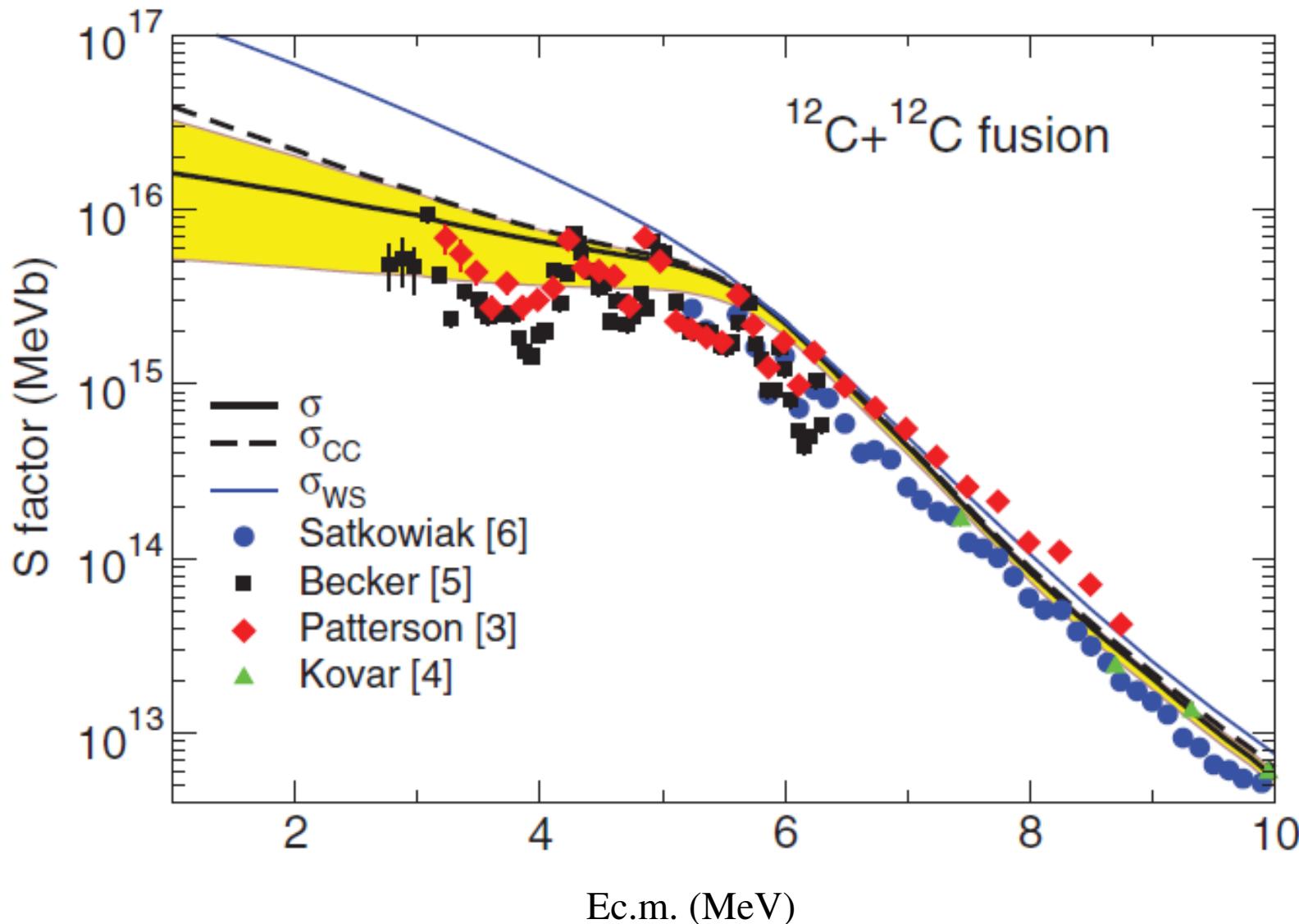
AD-T & Wiescher, Physical Review C 97 (2018) 055802

# Astrophysical S-Factor Excitation Function for $^{12}\text{C} + ^{12}\text{C}$

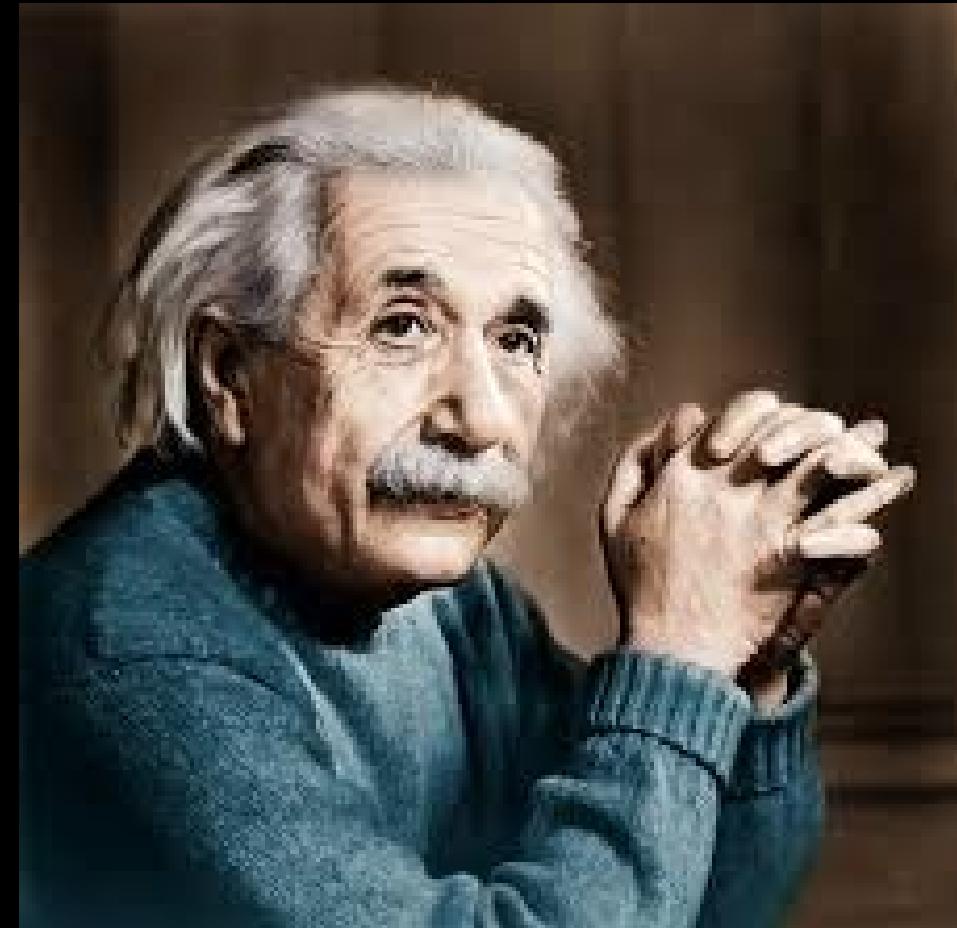


# Coupled-Channels Calculations for $^{12}\text{C} + ^{12}\text{C}$

Jiang, Esbensen et al., PRL 110 (2013) 072701

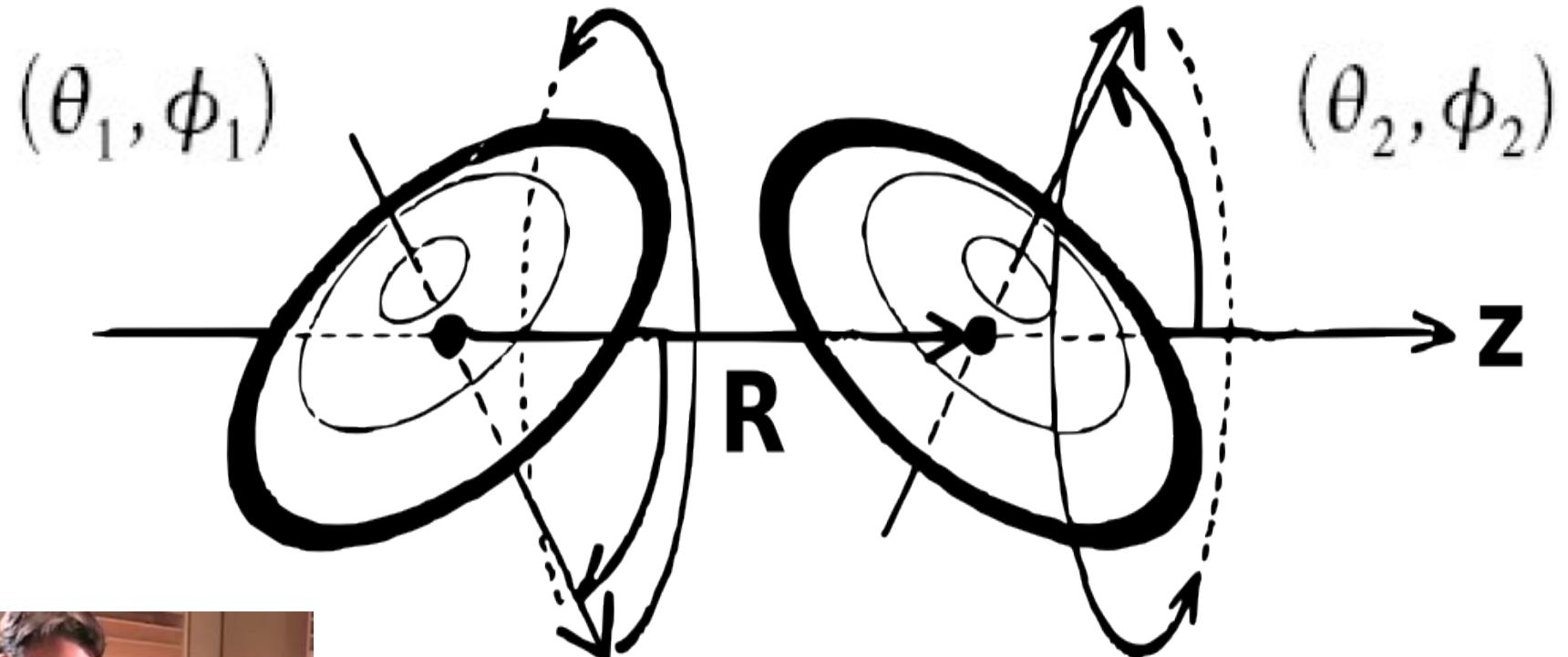


**“WE CANNOT  
SOLVE OUR  
PROBLEMS  
WITH THE SAME  
THINKING WE  
USED WHEN WE  
CREATED THEM”**



# The $^{12}\text{C} + ^{12}\text{C}$ Molecular Structure

Greiner, Park & Scheid, in Nuclear Molecules, World Scientific, 1994



Quadrupole deformation of  $^{12}\text{C}$ :  $\sim -0.5$

**How does this molecular structure affect low-energy fusion?**

# Wave-Packet Dynamics

D.J. Tannor, Quantum Mechanics: a Time-Dependent Perspective, USB, 2007



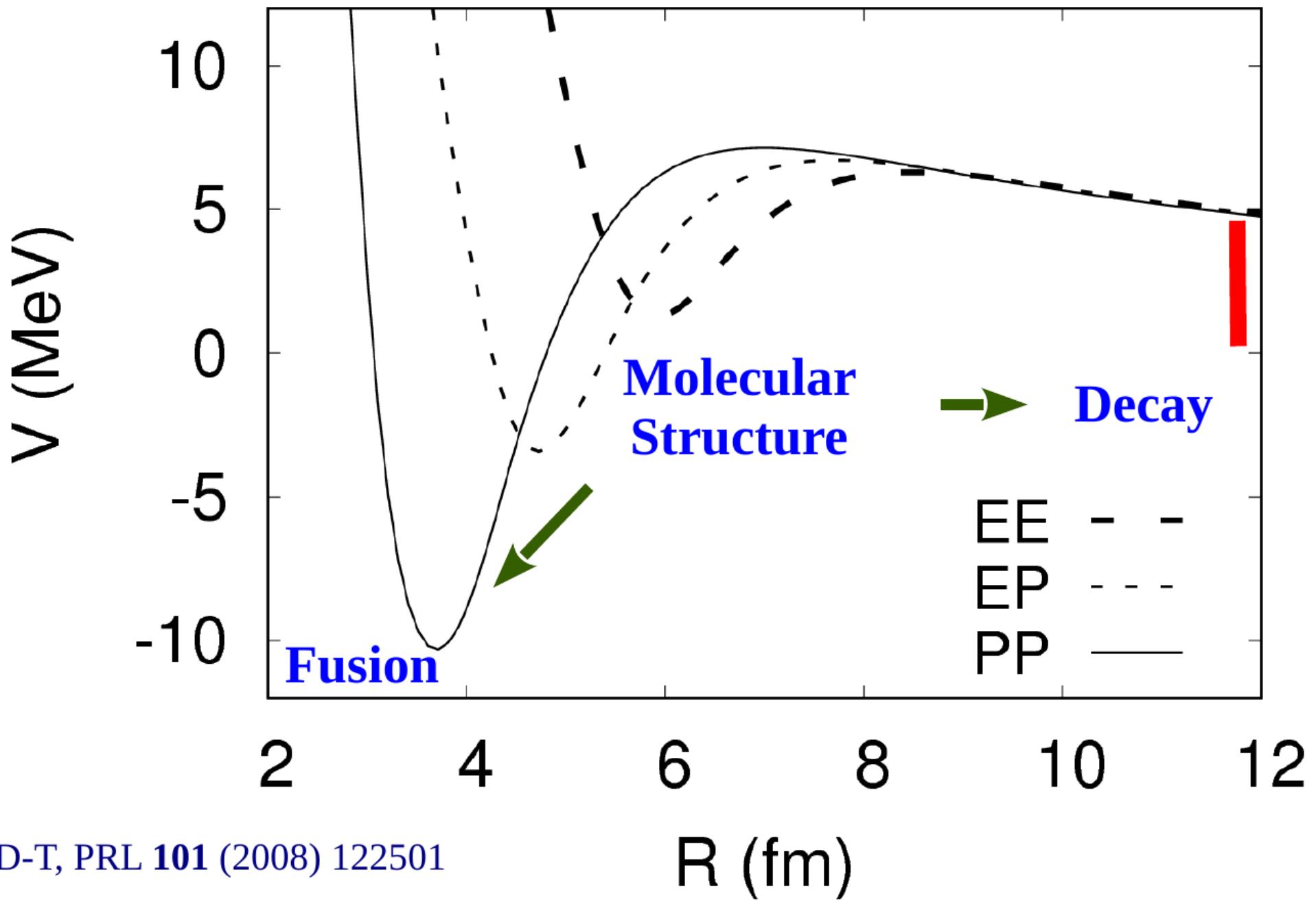
- ◆ **Preparation:** the initial state  $\Psi(t = 0)$

- ◆ **Time propagation:**  $\Psi(0) \rightarrow \Psi(t)$ ,  
guided by the operator,  $\exp(-i \hat{H} t/\hbar)$

$\hat{H}$  is the model Hamiltonian

- ◆ **Analysis:** extraction of probabilities from  
the time-dependent wave function

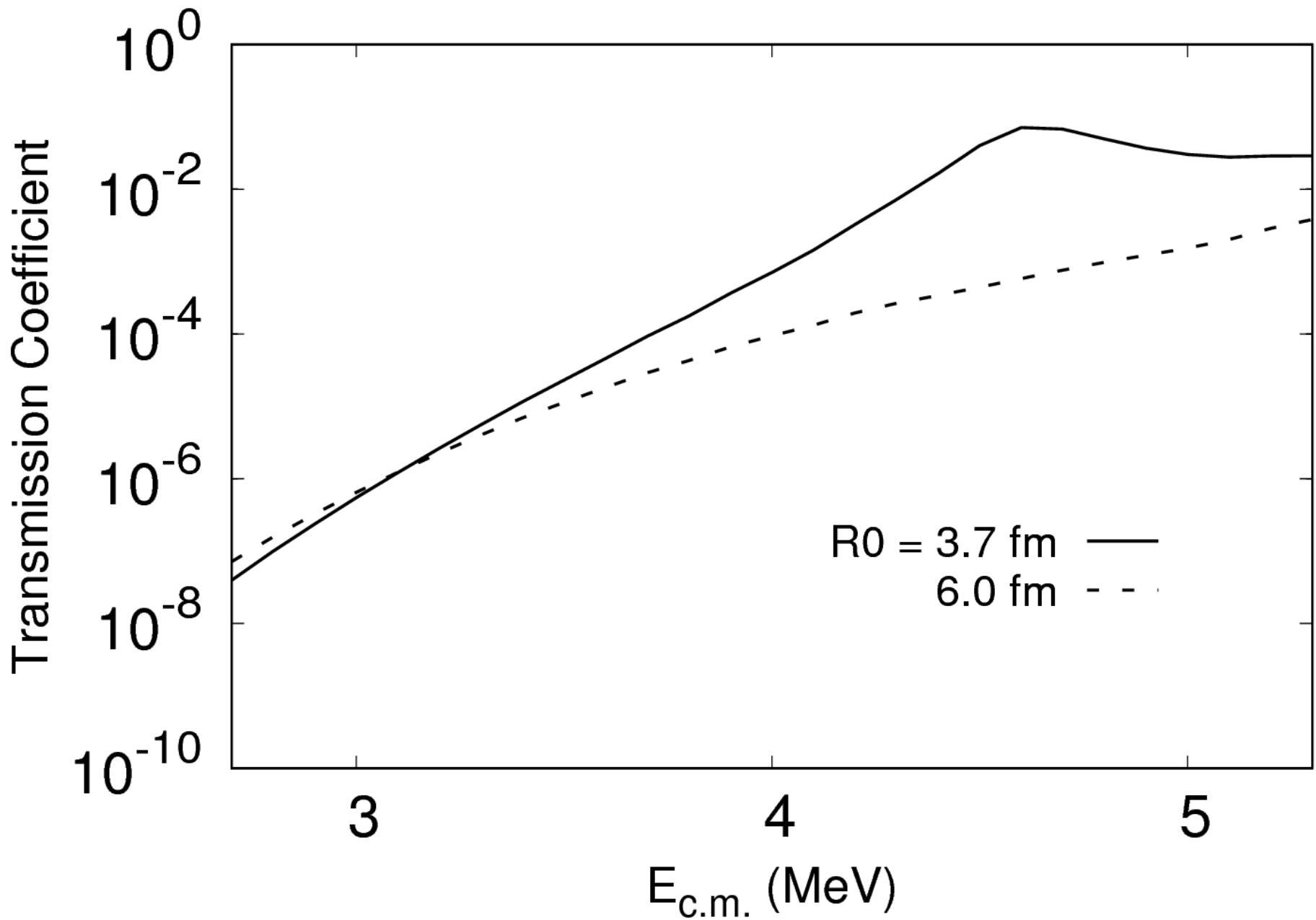
# Collective Potential-Energy Landscape for $^{12}\text{C} + ^{12}\text{C}$



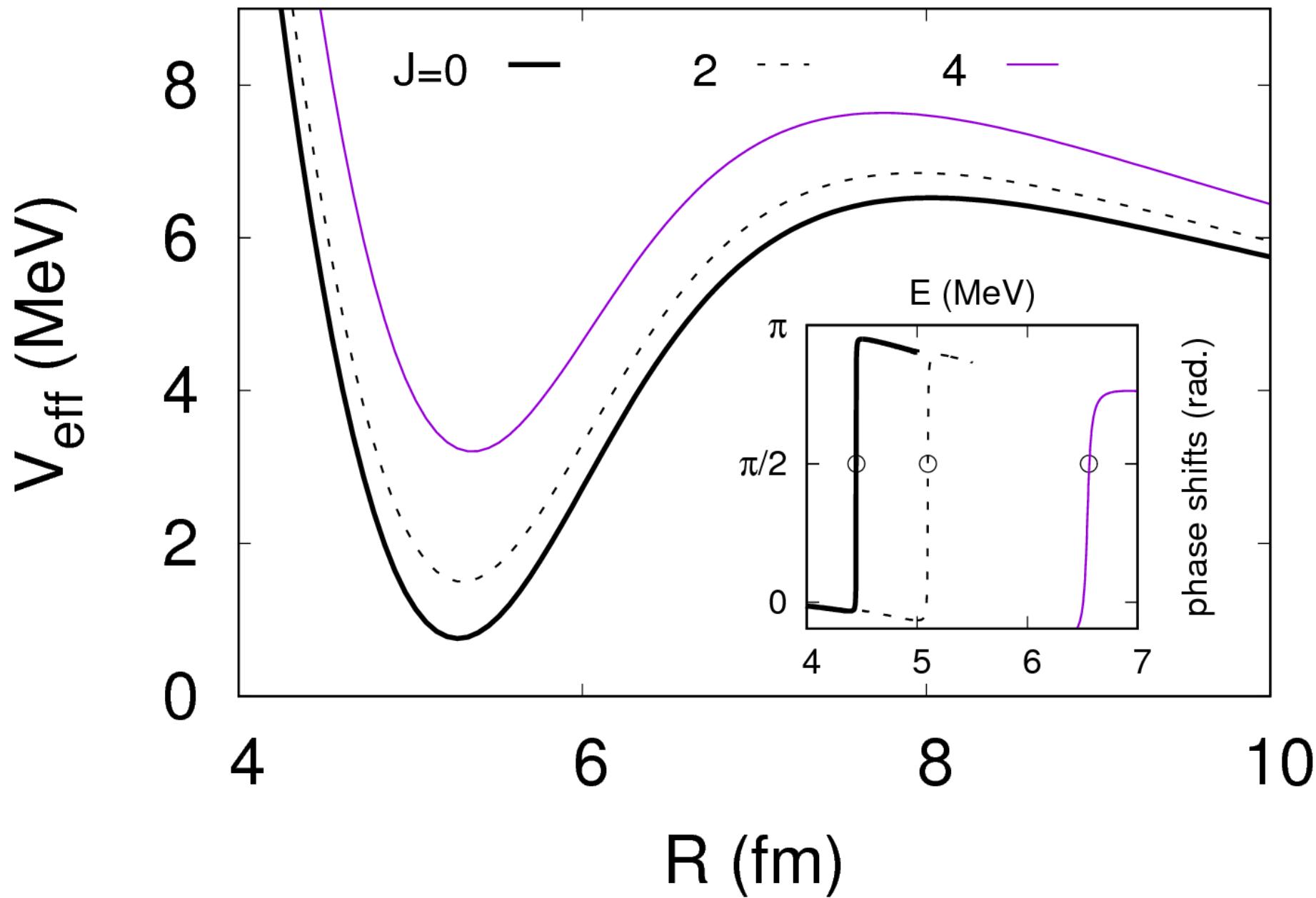
AD-T, PRL 101 (2008) 122501

Moeller & Iwamoto, NPA 575 (1994) 381

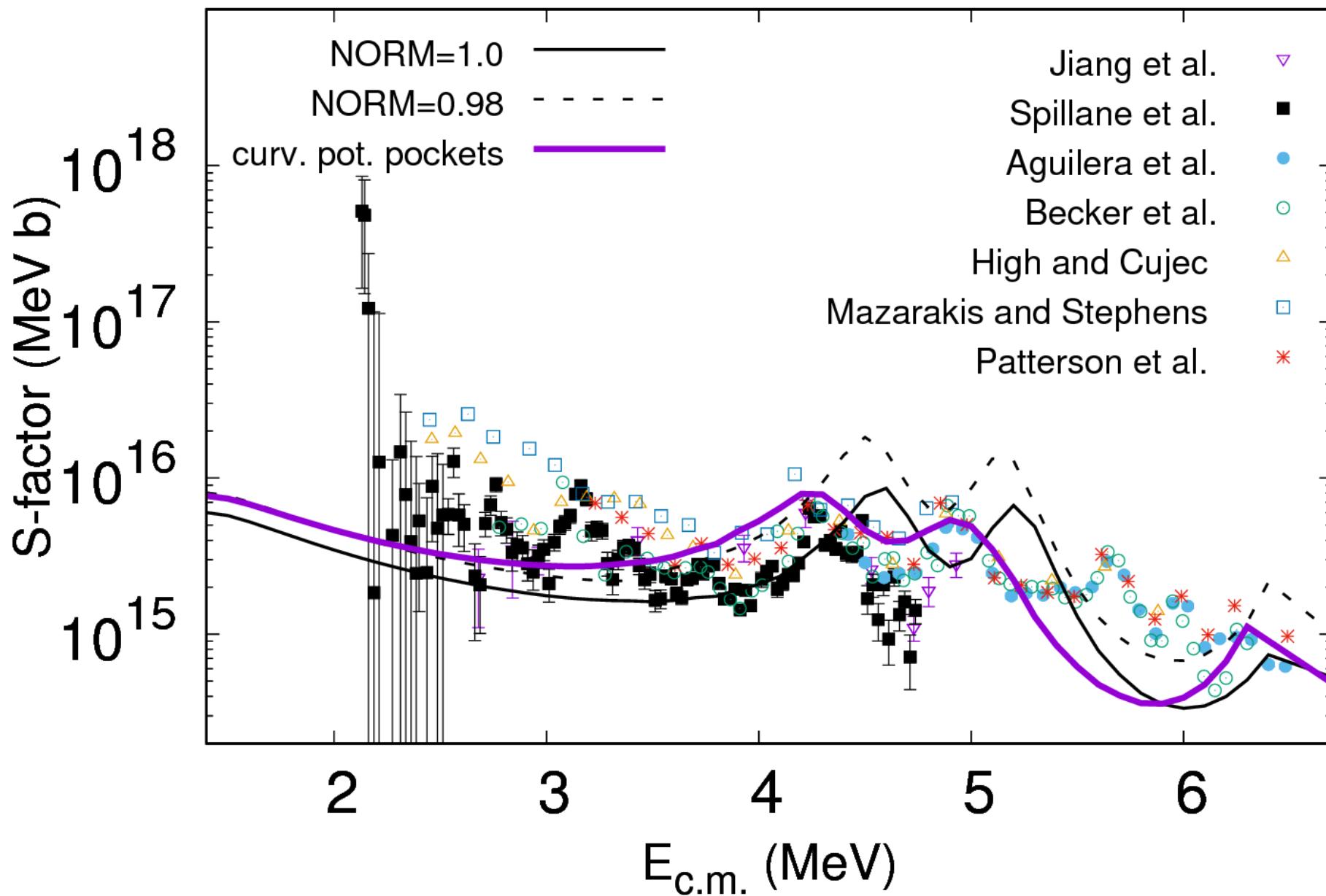
## Role of the imaginary fusion potential in the transmission coefficient



# Phase shift analysis of the effective potentials for molecular configurations



# Astrophysical S-Factor for $^{12}\text{C} + ^{12}\text{C}$ Fusion



# Conclusions

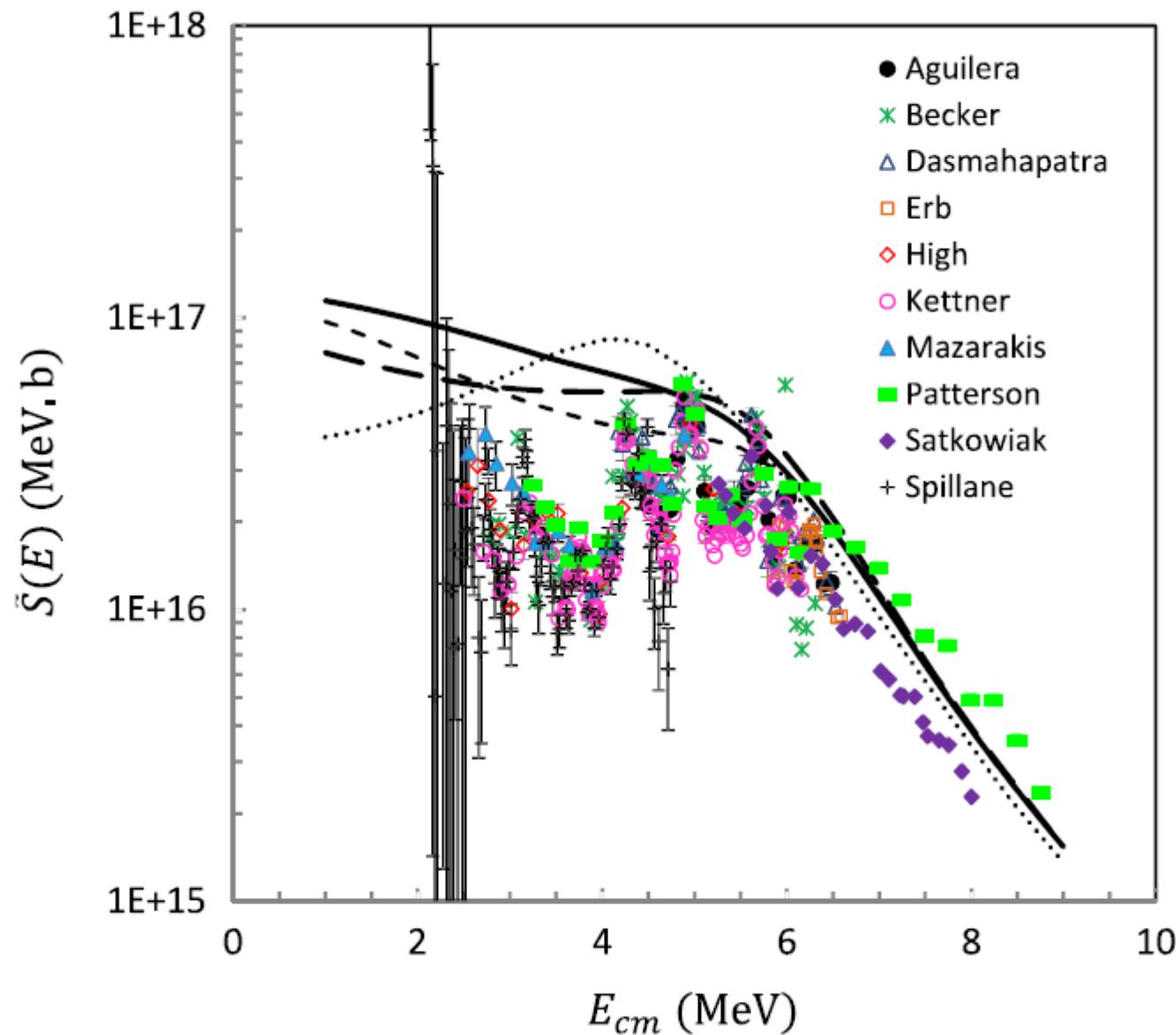
AD-T & Wiescher, Physical Review C **97** (2018) 055802

- The **fusion imaginary potential** for specific dinuclear configurations is crucial for the appearance of resonances.
- **Three resonant structures** are revealed in the calculations, reproducing similar structures in the experimental data.
- Fusion cross sections **monotonically** decline towards stellar energies.
- **Resonant structures** in the experimental data that are not explained may be due to **cluster effects** in the nuclear molecule.

# **EXTRA SLIDES**

# Coupled-Channels Calculations for $^{12}\text{C} + ^{12}\text{C}$

Assuncao & Descouvemont, PLB 723 (2013) 355



# Fusion Cross Section & Astrophysical S-Factor

$$S(E) = \sigma(E) E \exp(2\pi\eta)$$

Structural  
factor

[MeV barn]

Fusion  
cross section

[barn =  $10^{-28} \text{ m}^2$ ]

$$\eta = \left(\frac{\mu}{2}\right)^{1/2} \frac{Z_1 Z_2 e^2}{\hbar E^{1/2}}$$

Sommerfeld  
parameter

$S(E)$  represents the fusion cross section free of Coulomb suppression, which is adequate for extrapolation towards stellar energies

# Kinetic-Energy of Two Deformed Colliding Nuclei

Gatti *et al.*, JCP 123 (2005) 174311

$$\begin{aligned} \frac{2\hat{T}}{\hbar^2} = & -\frac{1}{\mu} \frac{\partial^2}{\partial R^2} + \left( \frac{1}{I_1} + \frac{1}{\mu R^2} \right) \hat{j}_1^2 + \left( \frac{1}{I_2} + \frac{1}{\mu R^2} \right) \hat{j}_2^2 \\ & + \frac{1}{\mu R^2} [\hat{j}_{1,+} \hat{j}_{2,-} + \hat{j}_{1,-} \hat{j}_{2,+} + J(J+1) \\ & - 2k_1^2 - 2k_1 k_2 - 2k_2^2] - \boxed{\frac{C_+(J, K)}{\mu R^2} (\hat{j}_{1,+} + \hat{j}_{2,+})} \\ & - \boxed{-\frac{C_-(J, K)}{\mu R^2} (\hat{j}_{1,-} + \hat{j}_{2,-})} \end{aligned}$$

**Coriolis interaction**

$\mu$  is the reduced mass for the radial motion,

$I_i$  is the  $^{12}\text{C}$  rotational inertia,

$J$  is the total angular momentum with projection  $K = k_1 + k_2$ ,

$C_{\pm}(J, K) = \sqrt{J(J+1) - K(K \pm 1)}$ ,

$\hat{j}_i^2 = -\frac{1}{\sin \theta_i} \frac{\partial}{\partial \theta_i} \sin \theta_i \frac{\partial}{\partial \theta_i} + \frac{k_i^2}{\sin^2 \theta_i}$ ,

$\hat{j}_{i,\pm} = \pm \frac{\partial}{\partial \theta_i} - k_i \cot \theta_i$ , with  $k_i \rightarrow k_i \pm 1$ .

# Initial state $\Psi(t = 0)$ : the $^{12}\text{C}$ nuclei are well separated

$$\Psi_0(R, \theta_1, k_1, \theta_2, k_2) = \chi_0(R) \psi_0(\theta_1, k_1, \theta_2, k_2),$$

↑  
Radial motion

Internal rotational motion

$$\chi_0(R) = (\sqrt{\pi} \sigma)^{-1/2} \exp\left[-\frac{(R - R_0)^2}{2\sigma^2}\right] e^{iP_0(R - R_0)},$$

$$\begin{aligned} \psi_0(\theta_1, k_1, \theta_2, k_2) &= [\zeta_{j_1, m_1}(\theta_1, k_1) \zeta_{j_2, m_2}(\theta_2, k_2) \\ &\quad + (-1)^J \zeta_{j_2, -m_2}(\theta_1, k_1) \zeta_{j_1, -m_1}(\theta_2, k_2)] \\ &\quad / \sqrt{2 + 2 \delta_{j_1, j_2} \delta_{m_1, -m_2}}, \end{aligned}$$

where  $\zeta_{j, m}(\theta, k) = \sqrt{\frac{(2j+1)(j-m)!}{2(j+m)!}} P_j^m(\cos \theta) \delta_{km}$ ,  
and  $P_j^m$  are associated Legendre functions.

# Time Propagation of the Wave Function

$$|\Psi_J(t)\rangle = \underbrace{e^{-i \hat{H} t/\hbar}}_{\text{evolution operator}} |\Psi_J(0)\rangle$$

The evolution operator is represented as a convergent series of modified Chebyshev polynomials

Tannor, Quantum Mechanics from a Time-Dependent Perspective, USB, 2007

## Power Spectrum of the Wave Function

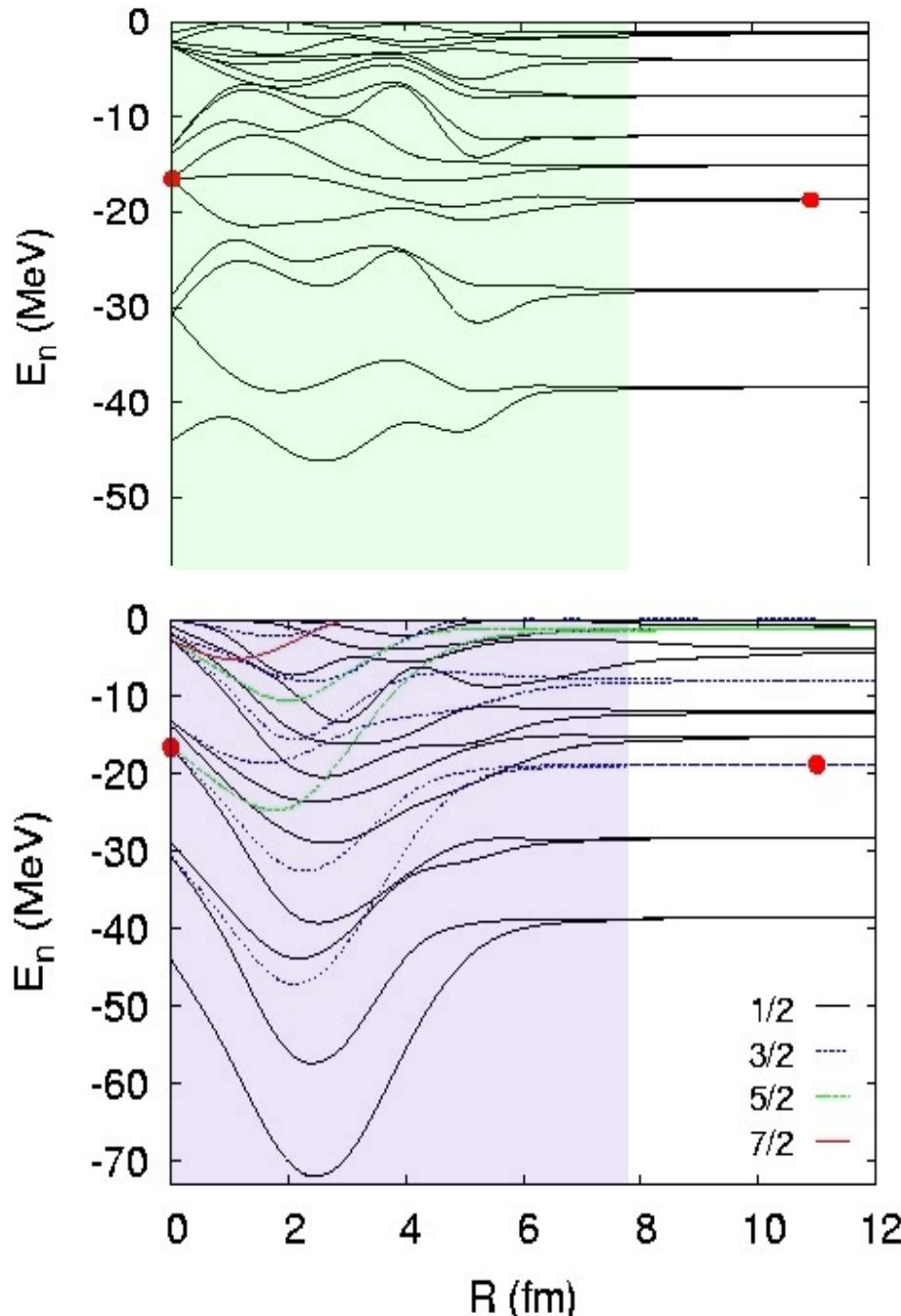
$$\mathcal{P}(E) = \langle \Psi(t) | \underbrace{\delta(E - \hat{H})}_{\text{Energy projector}} | \Psi(t) \rangle$$

## Reflection & Transmission Coefficients

$$\mathcal{R}(E) = \frac{\mathcal{P}^{final}(E)}{\mathcal{P}^{initial}(E)}$$

$$T(E) = 1 - \mathcal{R}(E)$$

# Sensitivity of Molecular Shell Structure to the $^{12}\text{C}$ Alignment

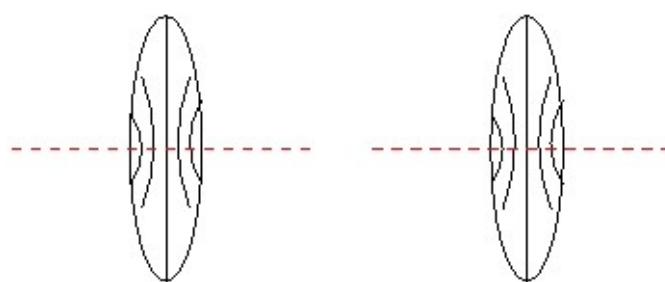


$$V = \sum_{s=1}^2 e^{-i\mathbf{R}_s \hat{k}} \hat{U}(\Omega_s) V_s \hat{U}^{-1}(\Omega_s) e^{i\mathbf{R}_s \hat{k}}$$

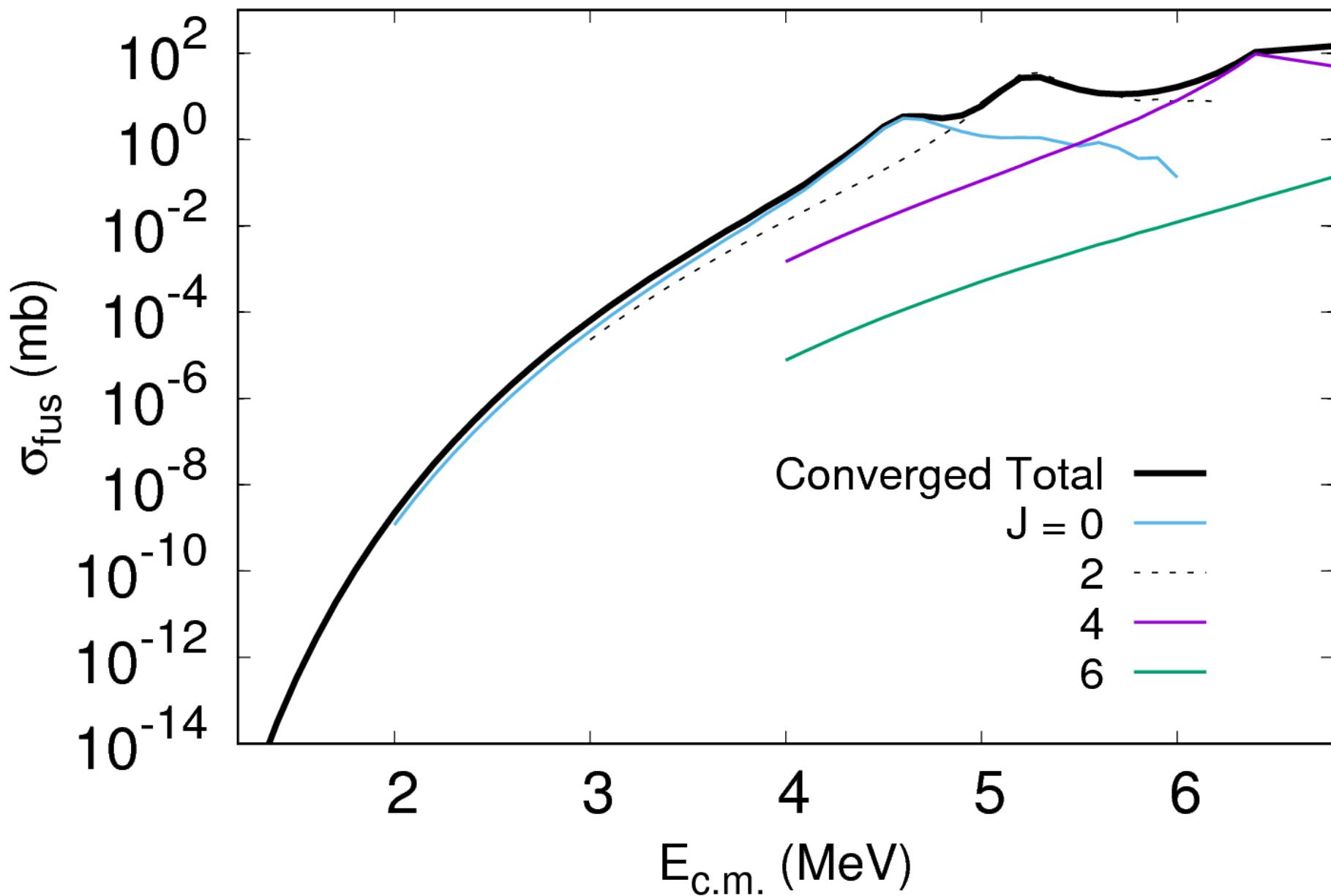
$\uparrow$

$$V_s \approx \sum_{\nu\mu}^N |s\nu\rangle V_{\nu\mu}^s \langle s\mu|$$

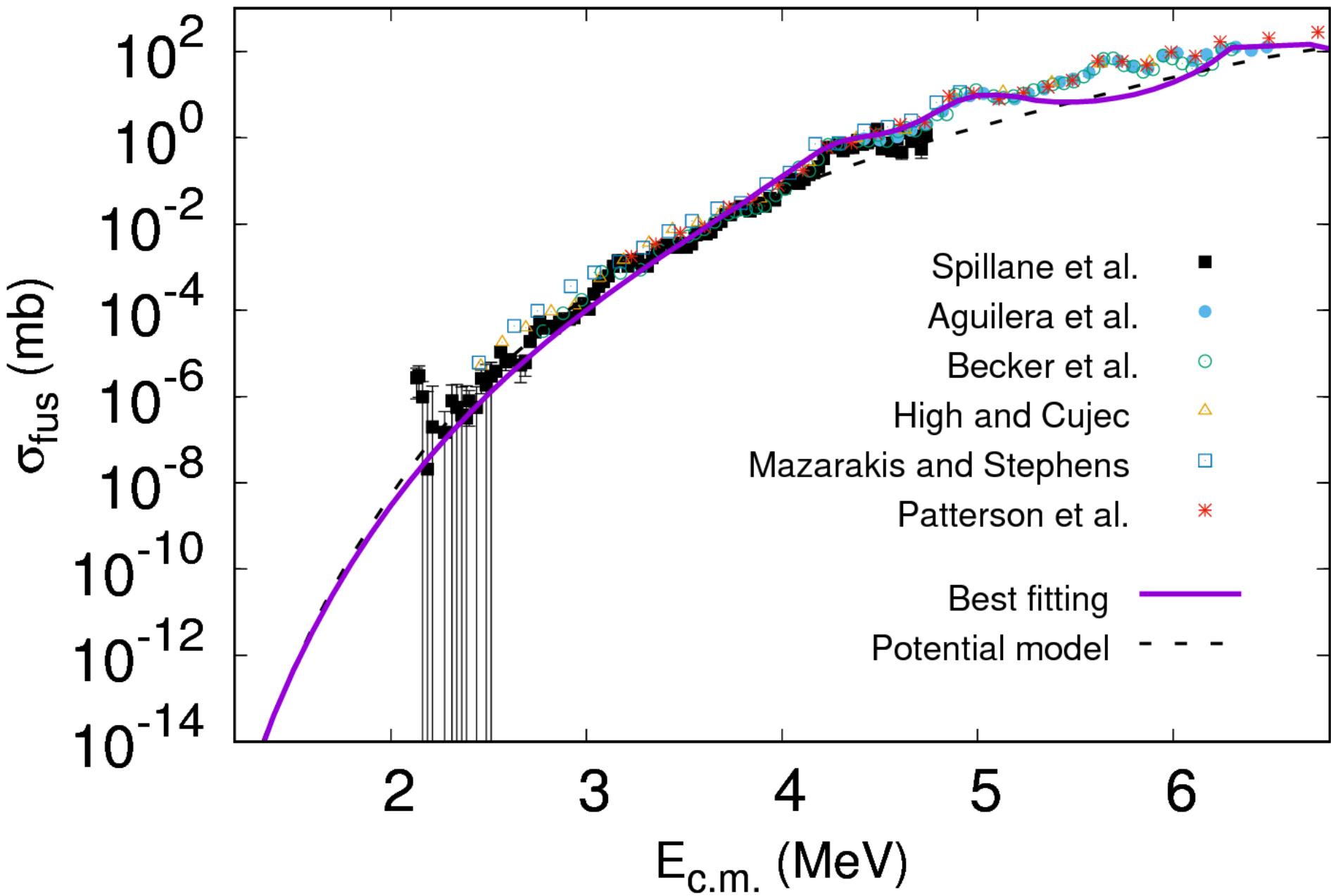
Potential Separable Expansion Method



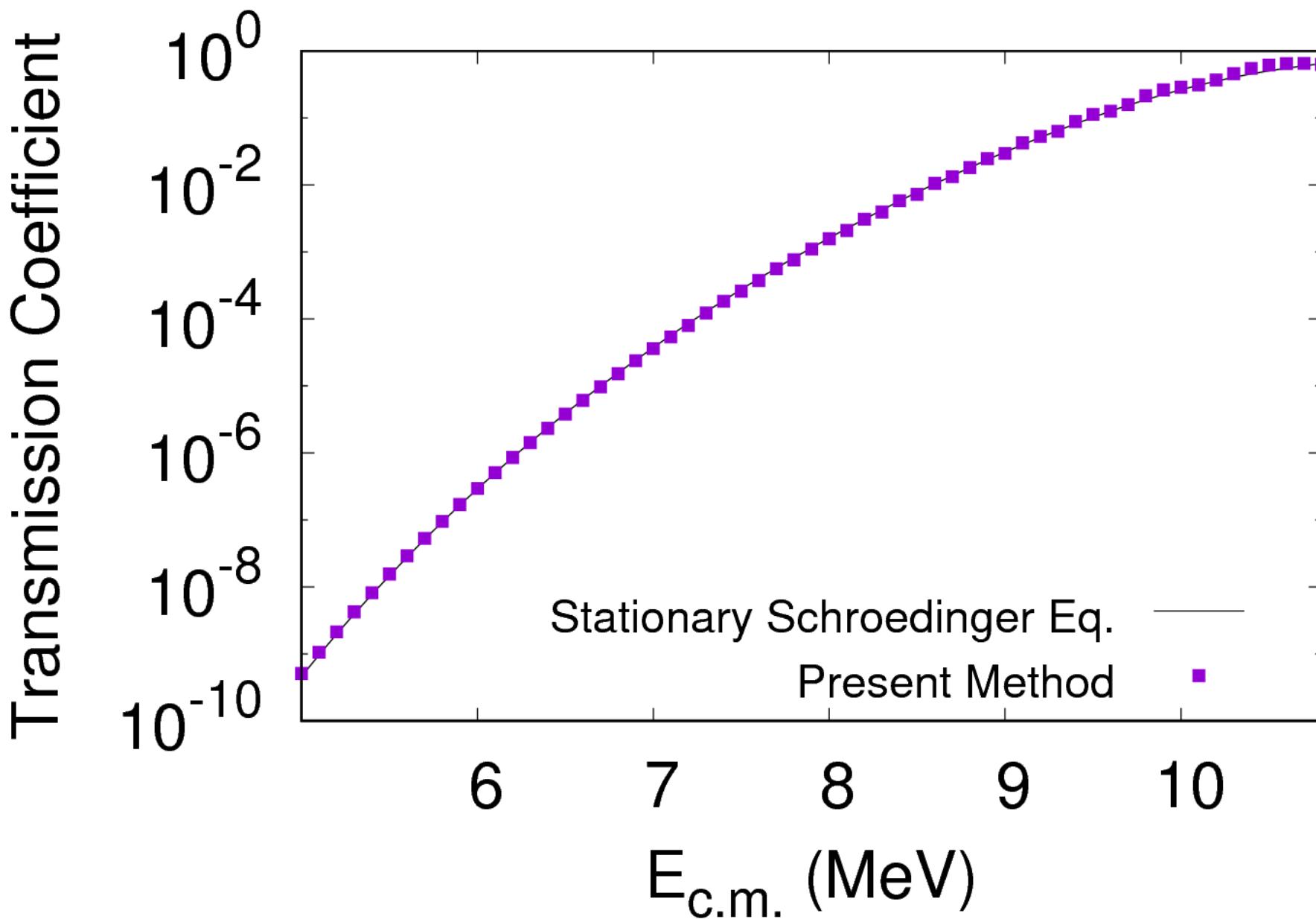
# Fusion Excitation Function for $^{12}\text{C} + ^{12}\text{C}$



# Fusion Excitation Function for $^{12}\text{C} + ^{12}\text{C}$



# Transmission coefficients for $^{16}\text{O} + ^{16}\text{O}$ central collisions



# An increase in the $^{12}\text{C} + ^{12}\text{C}$ fusion rate from resonances at astrophysical energies

A. Tumino<sup>1,2\*</sup>, C. Spitaleri<sup>2,3</sup>, M. La Cognata<sup>2</sup>, S. Cherubini<sup>2,3</sup>, G. L. Guardo<sup>2,4</sup>, M. Gulino<sup>1,2</sup>, S. Hayakawa<sup>2,5</sup>, I. Indelicato<sup>2</sup>, L. Lamia<sup>2,3</sup>, H. Petrascu<sup>4</sup>, R. G. Pizzone<sup>2</sup>, S. M. R. Puglia<sup>2</sup>, G. G. Rapisarda<sup>2</sup>, S. Romano<sup>2,3</sup>, M. L. Sergi<sup>2</sup>, R. Spartá<sup>2</sup> & L. Trache<sup>4</sup>

