

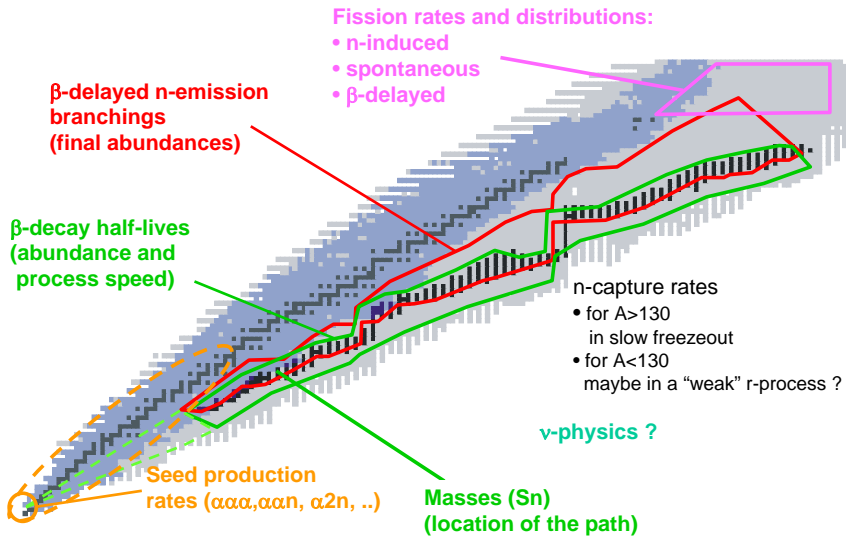
# Beta decay rates of neutron-rich nuclei

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Physics at SPES with non re-accelerated beams  
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# Introduction



Transitions are obtained by solving the pn-RQRPA equations

$$\begin{pmatrix} A & B \\ B^* & A^* \end{pmatrix} \begin{pmatrix} X^\lambda \\ Y^\lambda \end{pmatrix} = E_\lambda \begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix} \begin{pmatrix} X^\lambda \\ Y^\lambda \end{pmatrix}$$

Residual interaction is derived from the Lagrangian density

$$\mathcal{L}_{\rho+\pi} = -g_\rho \bar{\psi} \gamma_\mu \vec{\rho}^\mu \vec{\tau} \psi - \frac{f_\pi}{m_\pi} \bar{\psi} \gamma_5 \gamma^\mu \partial_\mu \vec{\pi} \vec{\tau} \psi$$

Total strength of a particular transition

$$B_{\lambda,J}(GT) = \left| \sum_{pn} \langle p \| \hat{O}_J \| n \rangle \left( X_{pn}^{\lambda,J} u_p v_n - Y_{pn}^{\lambda,J} v_p u_n \right) \right|^2$$

Decay rate:

$$\lambda_i = D \int_1^{W_{0,i}} W \sqrt{W^2 - 1} (W_{0,i} - W)^2 F(Z, W) C(W) dW$$

$$T_{1/2} = \frac{\ln 2}{\lambda}, \quad D = \frac{(G_F V_{ud})^2 (m_e c^2)^5}{2\pi^3 \hbar}$$

Allowed decays shape factor:

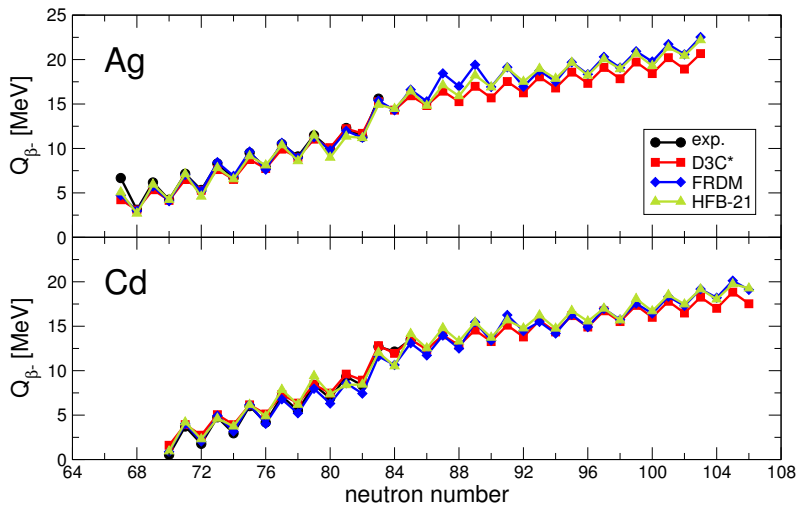
$$C(W) = B(GT)$$

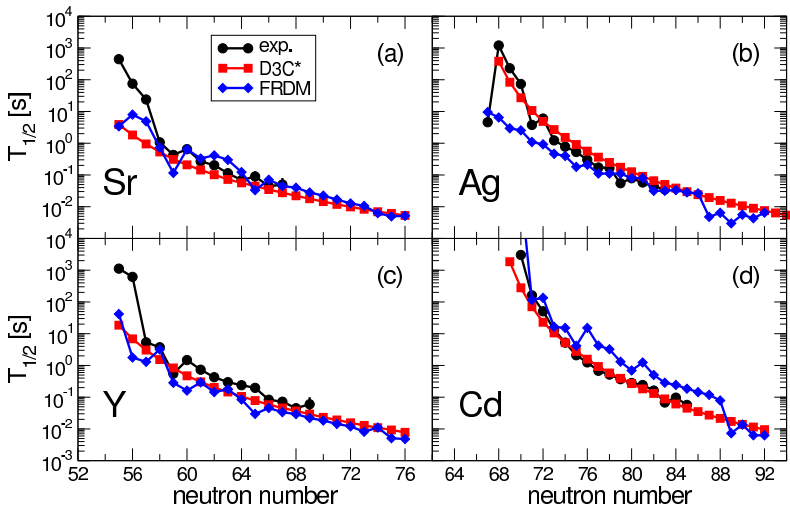
First-forbidden decays shape factor:

$$C(W) = k \left( 1 + aW + bW^{-1} + cW^2 \right)$$

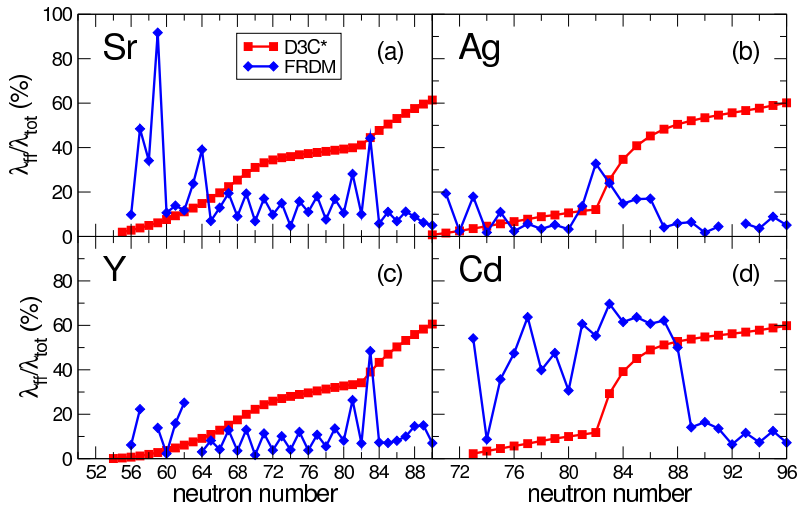
$$\begin{aligned}
k &= \left[ \zeta_0^2 + \frac{1}{9} w^2 \right]_{(0)} + \left[ \zeta_1^2 + \frac{1}{9} (x+u)^2 - \frac{4}{9} \mu_1 \gamma_1 u (x+u) \right. \\
&\quad \left. + \frac{1}{18} W_0^2 (2x+u)^2 - \frac{1}{18} \lambda_2 (2x-u)^2 \right]_{(1)} \\
&\quad + \left[ \frac{1}{12} z^2 (W_0^2 - \lambda_2) \right]_{(2)} \\
ka &= \left[ -\frac{4}{3} uY - \frac{1}{9} W_0 (4x^2 + 5u^2) \right]_{(1)} - \left[ \frac{1}{6} W_0 z^2 \right]_{(2)} \\
kb &= \frac{2}{3} \mu_1 \gamma_1 \left\{ -[\zeta_0 w]_{(0)} + [\zeta_1 (x+u)]_{(1)} \right\} \\
kc &= \frac{1}{18} \left[ 8u^2 + (2x+u)^2 + \lambda_2 (2x-u)^2 \right]_{(1)} \\
&\quad + \frac{1}{12} \left[ (1 + \lambda_2) z^2 \right]_{(2)}
\end{aligned}$$

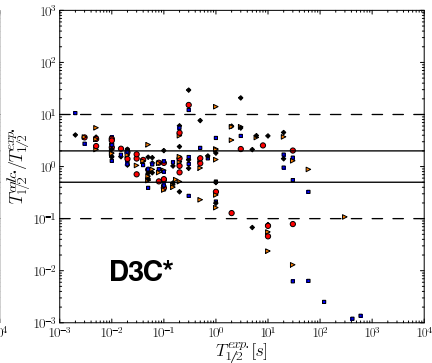
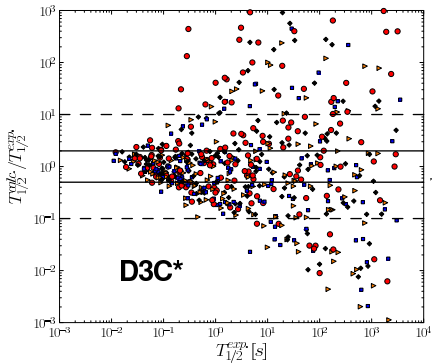
$$\begin{aligned}
w &= -g_A \sqrt{3} \frac{\langle f | \sum_k r_k [\mathbf{c}_1^k \otimes \sigma^k]^0 \mathbf{t}_-^k | i \rangle}{\sqrt{2J_i+1}}, \\
x &= -\frac{\langle f | \sum_k r_k \mathbf{c}_1^k \mathbf{t}_-^k | i \rangle}{\sqrt{2J_i+1}}, \\
u &= -g_A \sqrt{2} \frac{\langle f | \sum_k r_k [\mathbf{c}_1^k \otimes \sigma^k]^1 \mathbf{t}_-^k | i \rangle}{\sqrt{2J_i+1}}, \\
z &= 2g_A \frac{\langle f | \sum_k r_k [\mathbf{c}_1^k \otimes \sigma^k]^2 \mathbf{t}_-^k | i \rangle}{\sqrt{2J_i+1}}, \\
w' &= -g_A \frac{2}{\sqrt{3}} \frac{\langle f | \sum_k r_k l(1,1,1,1,r_k) [\mathbf{c}_1^k \otimes \sigma^k]^0 \mathbf{t}_-^k | i \rangle}{\sqrt{2J_i+1}}, \\
x' &= -\frac{2}{3} \frac{\langle f | \sum_k r_k l(1,1,1,1,r_k) \mathbf{c}_1^k \mathbf{t}_-^k | i \rangle}{\sqrt{2J_i+1}}, \\
u' &= -g_A \frac{2\sqrt{2}}{3} \frac{\langle f | \sum_k r_k l(1,1,1,1,r_k) [\mathbf{c}_1^k \otimes \sigma^k]^1 \mathbf{t}_-^k | i \rangle}{\sqrt{2J_i+1}}.
\end{aligned}$$







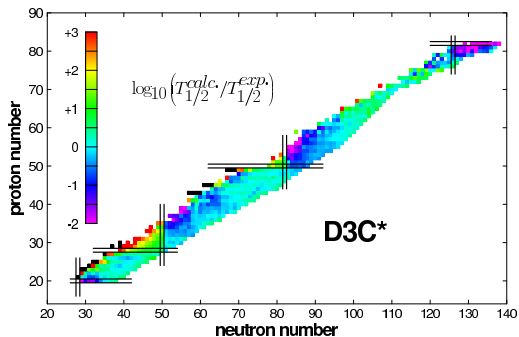




$$\bar{r} = \frac{1}{N} \sum_i \log \frac{T_{\text{th.}}}{T_{\text{exp.}}}$$

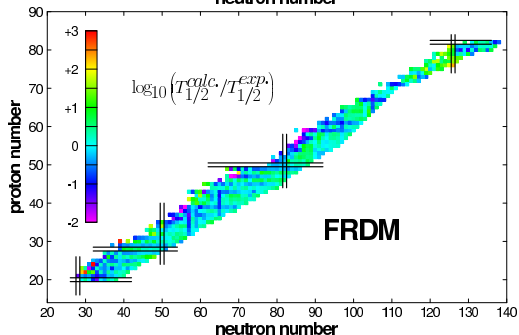
$$\sigma = \left[ \frac{1}{N} \sum_i (r_i - \bar{r})^2 \right]^{1/2}$$

	D3C*		FRDM	
$T_{\text{exp.}} [\text{s}]$	$\bar{r}$	$\sigma$	$\bar{r}$	$\sigma$
< 1000	0.011	0.889	0.021	0.660
< 100	0.057	0.791	0.040	0.580
< 10	0.061	0.645	0.046	0.515
< 1	0.011	0.436	0.019	0.409
< 0.1	0.041	0.195	0.021	0.354



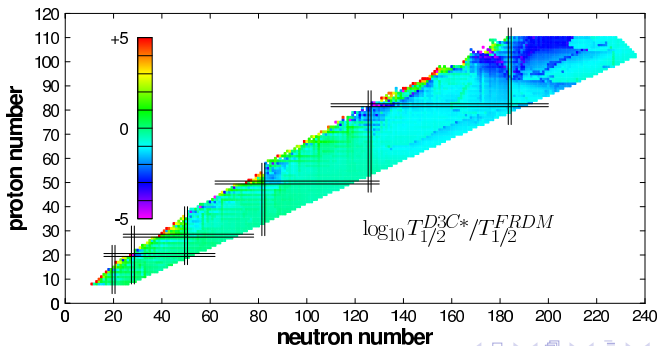
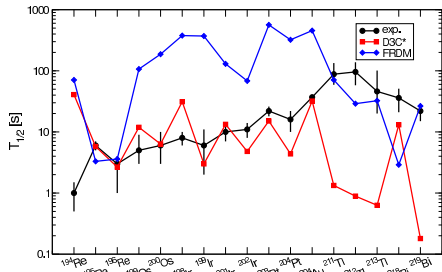
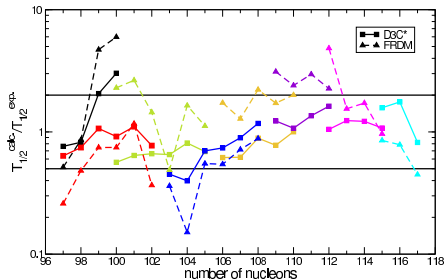
D3C\*

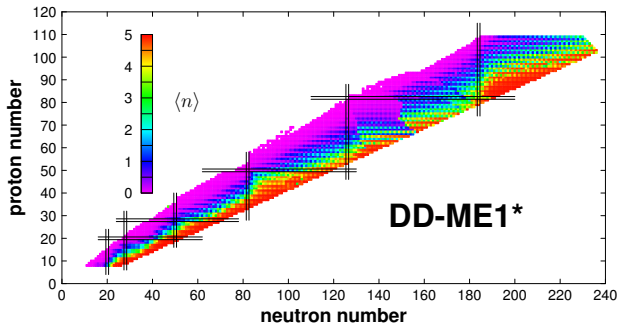
	$\bar{r}$	$\sigma$
even-even	-0.037	0.331
odd-Z	0.054	0.328
odd-N	-0.086	0.387
odd-odd	0.089	0.582
total	0.011	0.436



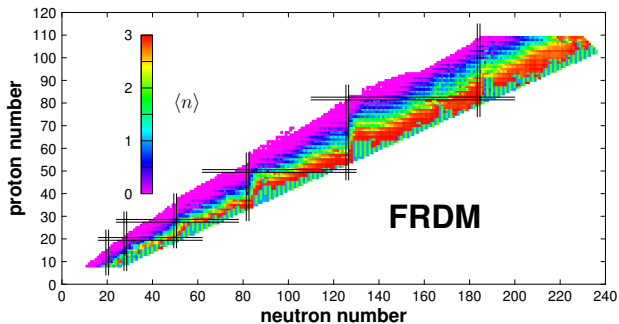
FRDM

	$\bar{r}$	$\sigma$
even-even	0.333	0.226
odd-Z	-0.128	0.288
odd-N	0.124	0.436
odd-odd	-0.179	0.409
total	0.019	0.409

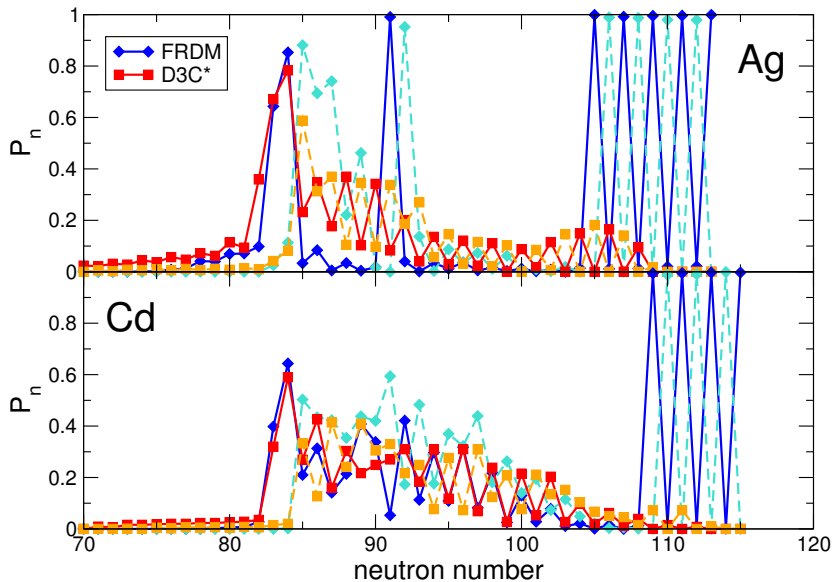




$$P_{xn} = \frac{1}{\lambda_{tot}} \sum_{E_i=S_{xn}}^{S_{(x+1)n}} \lambda_i$$



$$\langle n \rangle = \sum_i iP_{in}$$



# Evaluation of reactor antineutrino spectra

In reactors, 99% of the electrons come from decay of fission products of 4 nuclei.

$$S_{tot}(E) = \sum_{k=^{235}\text{U}, ^{238}\text{U}, ^{239}\text{Pu}, ^{241}\text{Pu}} \alpha_k S_k(E),$$

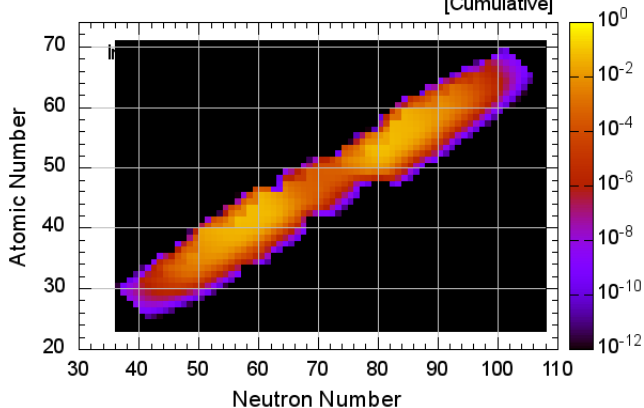
- $\alpha_k$  - number of fissions at considered time
- $S_k(E)$  -  $\beta$  spectrum normalized to one fission
- $E$  - kinetic energy of emitted electrons

Electrons (and antineutrinos) come from the  $\beta$ -decay of resulting fission fragments.

$$S_k(E) = \sum_{f=1}^{N_f} Y_f S_f(E)$$

$$S_f(E) = \sum_{i=1}^{N_t} \frac{\lambda_i}{\lambda_{tot}} S_f^i(Z, A, E_{max}, E).$$

Pu-239 Neutron-induced Fission Yields  
[Cumulative]



JAEA Nuclear Data Center



For allowed transitions the spectrum reads

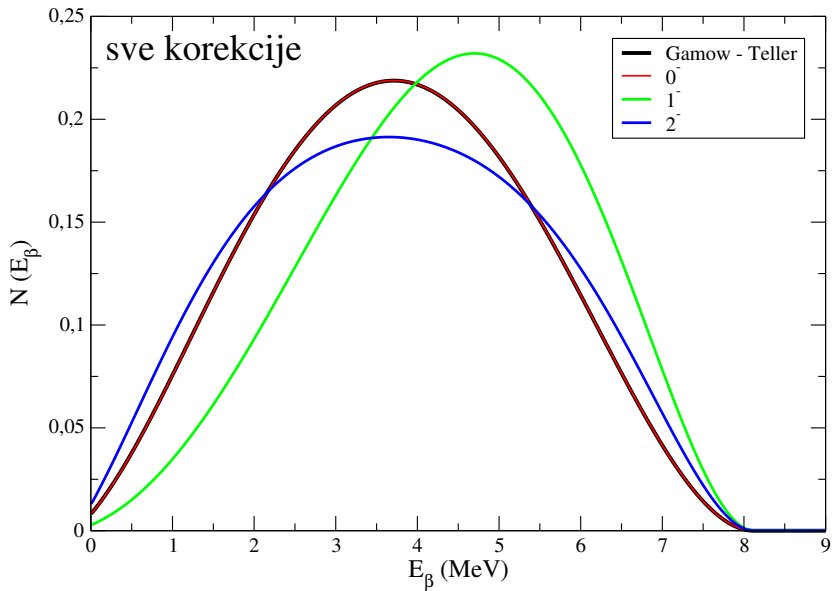
$$S_f^i = F(Z, A, E) \cdot pE(E - E_{max})^2 \cdot L_0(Z, E) \cdot C'(Z, E)$$

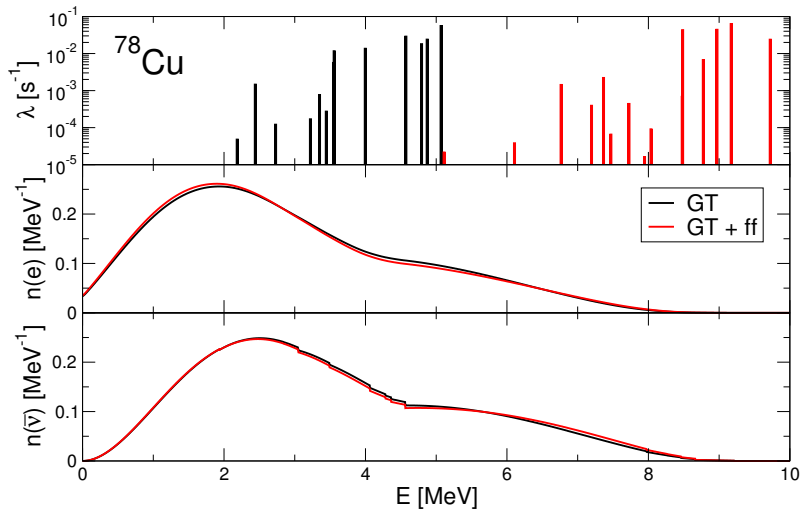
- $F(Z, A, E)$  - Fermi function, correction for the Coulomb field
- $L_0(Z, E)$  - correction for the finite size of the charge distribution
- $C'(Z, E)$  - correction for the nucleon moving within a nuclear potential
- other corrections are neglected

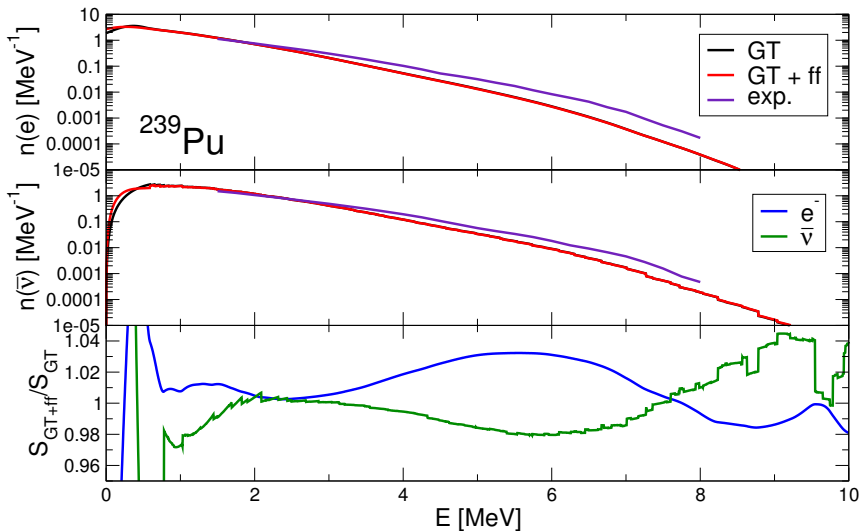
but if we include first-forbidden transitions

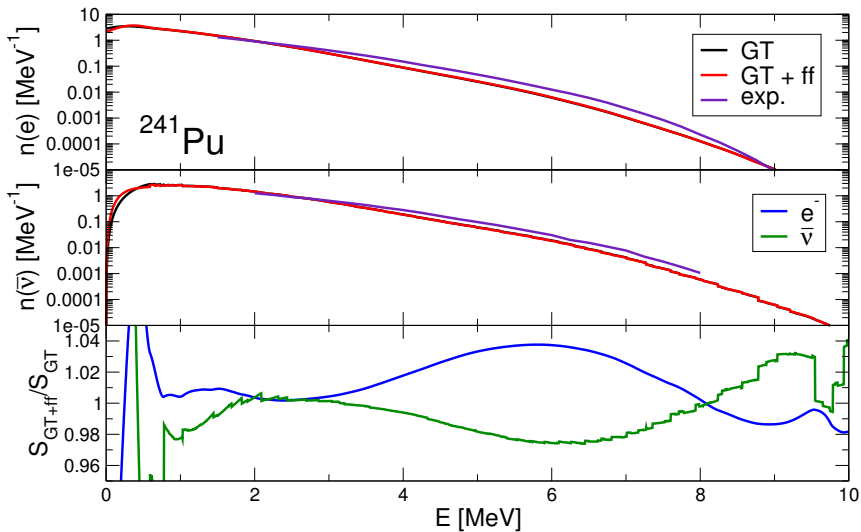
$$S_f^i = F(Z, A, E) \cdot pE(E - E_{max})^2 \cdot C(E) \cdot L_0(Z, E) \cdot C(Z, E)$$

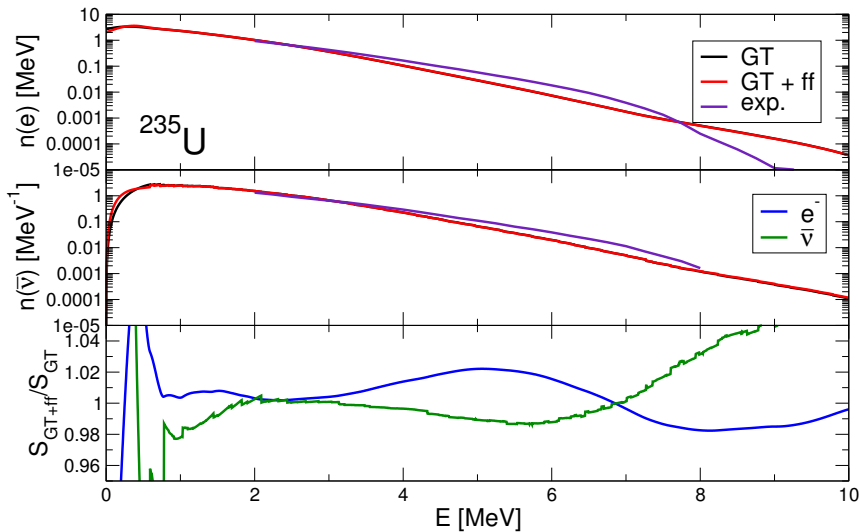
where  $C(E)$  is the *shape factor*

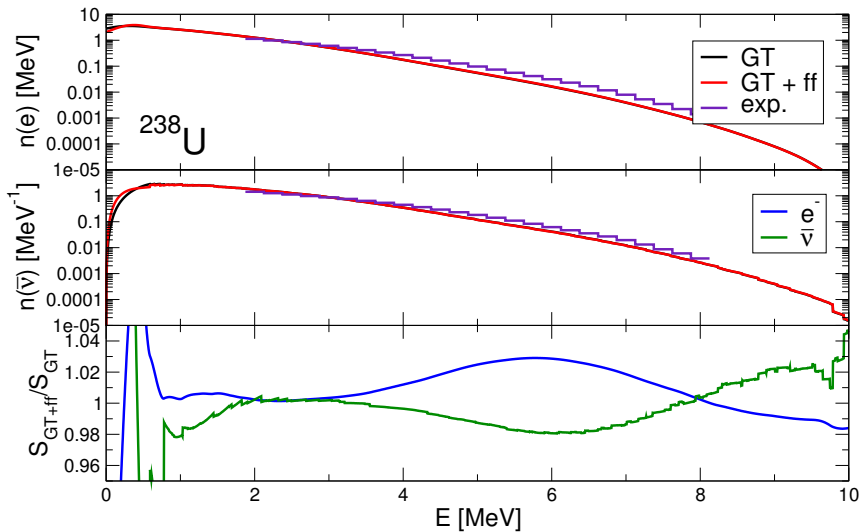


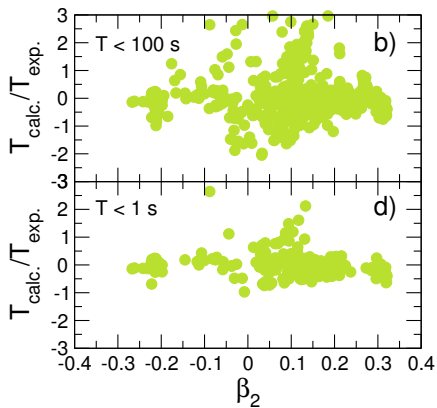
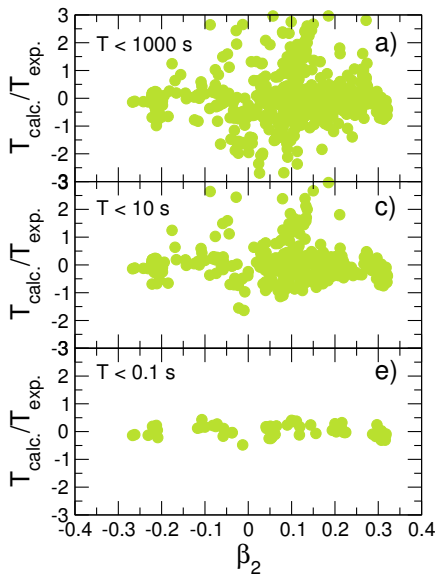




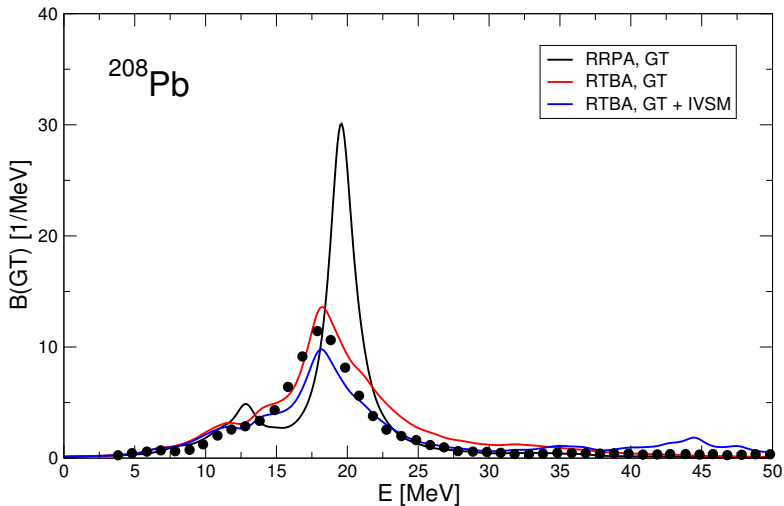












T. Wakasa *et al.*, Phys. Rev. C 85, 064606 (2012)

