

Constraints on the density dependence of the symmetry energy

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IFJ PAN, Kraków

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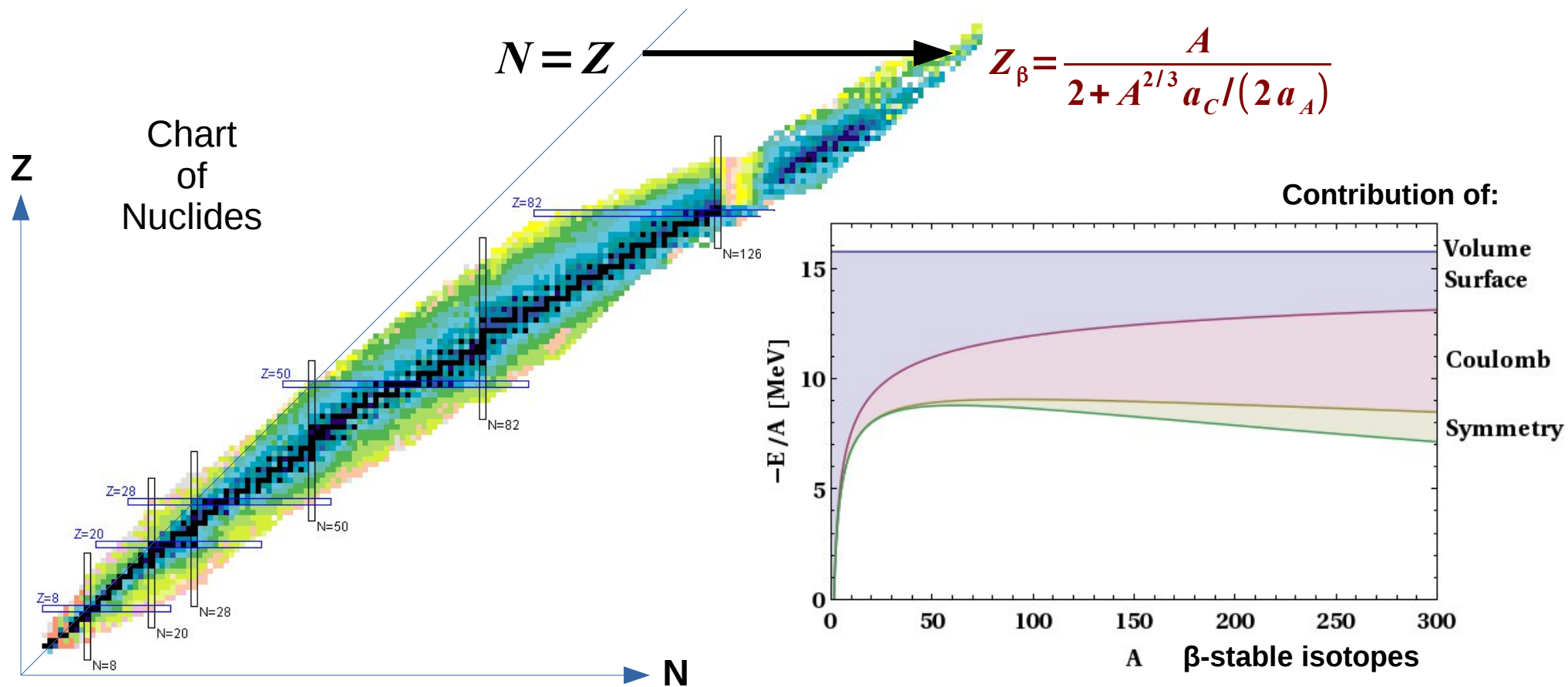
- short history and introduction
- present status
- future

Symmetry Energy (nucleus)

$$\frac{E_B}{A} = \underbrace{-a_V}_{\text{Volume}} + \underbrace{a_S \frac{1}{A^{1/3}}}_{\text{Surface}} + \underbrace{a_C \frac{Z^2}{A^{4/3}}}_{\text{Coulomb}} + \underbrace{a_A \left(\frac{N-Z}{A} \right)^2}_{\text{Symmetry}} \pm \underbrace{a_P \frac{1}{A^{3/2}}}_{\text{Pairing}}$$

Binding Energy/Nucleon
(semiempirical
Bethe-Weizsäcker)
acc. better than 1%

$a_V=15.8$	$a_S=18.0$	$a_C=0.72$	$a_A=23.5$	$a_P=11.5 0$	[MeV]
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Symmetry Energy (nuclear matter)

$$\frac{E_B}{A} = \underbrace{-a_V}_{\text{Volume}} + \underbrace{a_S \frac{1}{A^{1/3}}}_{\text{Surface}} + \underbrace{a_C \frac{Z^2}{A^{4/3}}}_{\text{Coulomb}} + \underbrace{a_A \left(\frac{N-Z}{A} \right)^2}_{\text{Symmetry}} \pm \underbrace{a_P \frac{1}{A^{3/2}}}_{\text{Pairing}}$$

Energy per nucleon in nuclear matter (EoS):

$$E(\rho, \delta) = E(\rho, 0) + E_{sym}(\rho) \delta^2 + o(\delta^4)$$

dominant symmetric matter (N=Z) term (*isoscalar*):

$$E(\rho, 0) \approx -a_V + \frac{K}{18} \left(\frac{\rho - \rho_0}{\rho_0} \right)^2 + \dots$$

symmetry term (*isovector*):

$$E_{sym}(\rho) \approx E_{sym}(\rho_0) + \frac{L}{3} \left(\frac{\rho - \rho_0}{\rho_0} \right) + \frac{K_{sym}}{18} \left(\frac{\rho - \rho_0}{\rho_0} \right)^2 + \dots$$

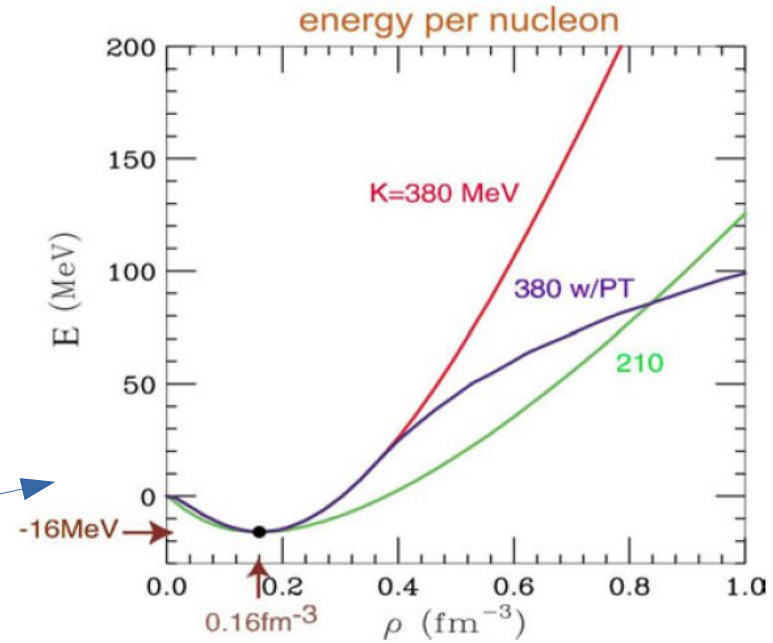
$\rho_n, \rho_p \rightarrow$ neutron, proton densities

$\rho = \rho_n + \rho_p \rightarrow$ nucleon density

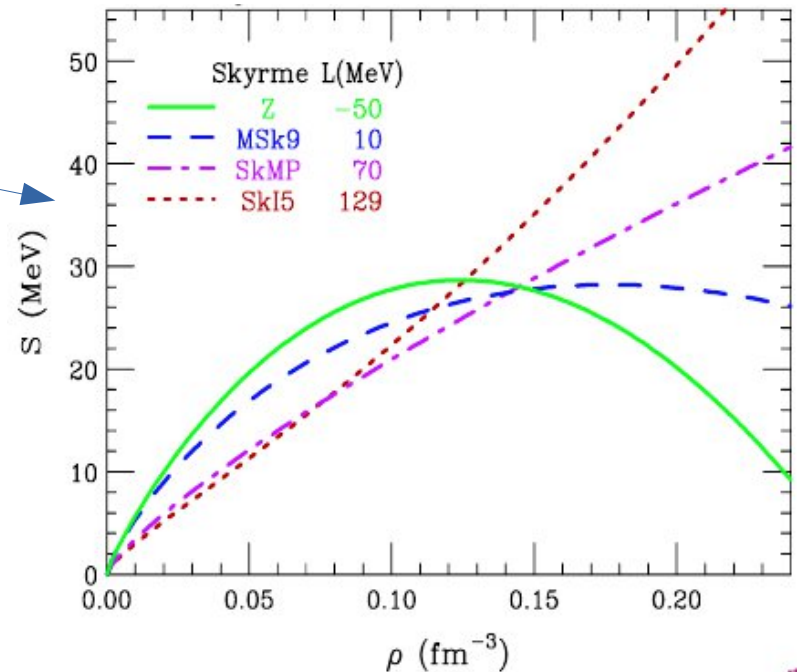
$\delta = \frac{\rho_n - \rho_p}{\rho} \rightarrow$ isospin asymmetry

$L = 3 \rho_0 \left. \frac{\partial E_{sym}}{\partial \rho} \right|_{\rho=\rho_0} \rightarrow \sim$ symmetry pressure

$K = 9 \rho_0^2 \left. \frac{\partial^2 E}{\partial \rho^2} \right|_{\rho=\rho_0} \rightarrow$ compressibility



P. Danielewicz, arXiv:nucl-th/0512009



(*) P. Danielewicz, J. Lee, NPA 818(2009)36

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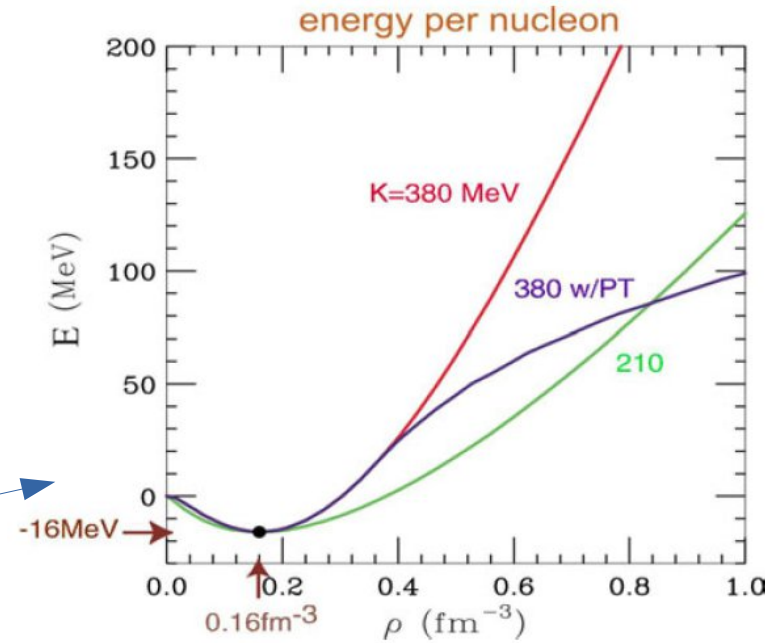
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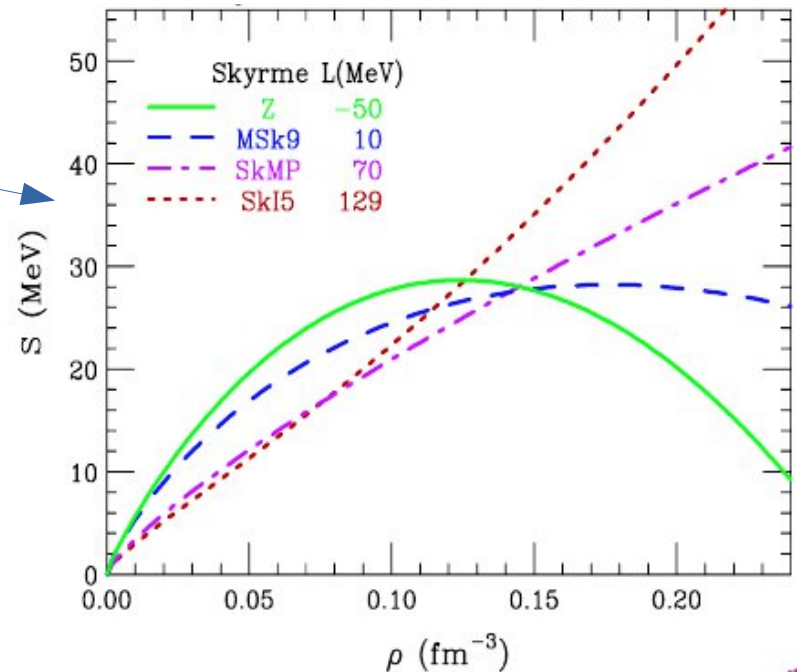
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More exact formula

B.-A. Li, Nucl. Phys. News, 27 (2017) 7

It is well known that the nucleon potential $U_{n/p}(k, \rho, \delta)$ in ANM can be expanded up to the second order in δ as

$$U_{\tau}(k, \rho, \delta) = U_0(k, \rho) + \tau_3 U_{sym,1}(k, \rho) \cdot \delta + U_{sym,2}(k, \rho) \cdot \delta^2 + \mathcal{O}(\delta^3), \quad (2)$$

where $\tau_3 = \pm$ for $\tau = n/p$ and $U_0(k, \rho)$, $U_{sym,1}(k, \rho)$ and $U_{sym,2}(k, \rho)$ are the isoscalar, isovector (symmetry or Lane potential [15]) and the second-order isoscalar potentials, respectively. At the mean-field level, using the Bruckner theory [16, 17] or the Hugenholtz-Van Hove (HVH) theorem [18], the $E_{sym}(\rho)$ and its density slope $L(\rho) \equiv [3\rho(\partial E_{sym}/\partial\rho)]_{\rho}$ at an arbitrary density ρ can be expressed generally as [19-21]

$$E_{sym}(\rho) = \frac{1}{3} \frac{\hbar^2 k_F^2}{2m_0^*} + \frac{1}{2} U_{sym,1}(\rho, k_F), \quad (3)$$

$$L(\rho) = \frac{2}{3} \frac{\hbar^2 k_F^2}{2m_0^*} + \frac{3}{2} U_{sym,1}(\rho, k_F) - \frac{1}{6} \left(\frac{\hbar^2 k^3}{m_0^{*2}} \frac{\partial m_0^*}{\partial k} \right) \Big|_{k_F} + \frac{dU_{sym,1}}{dk} \Big|_{k_F} k_F + 3U_{sym,2}(\rho, k_F), \quad (4)$$

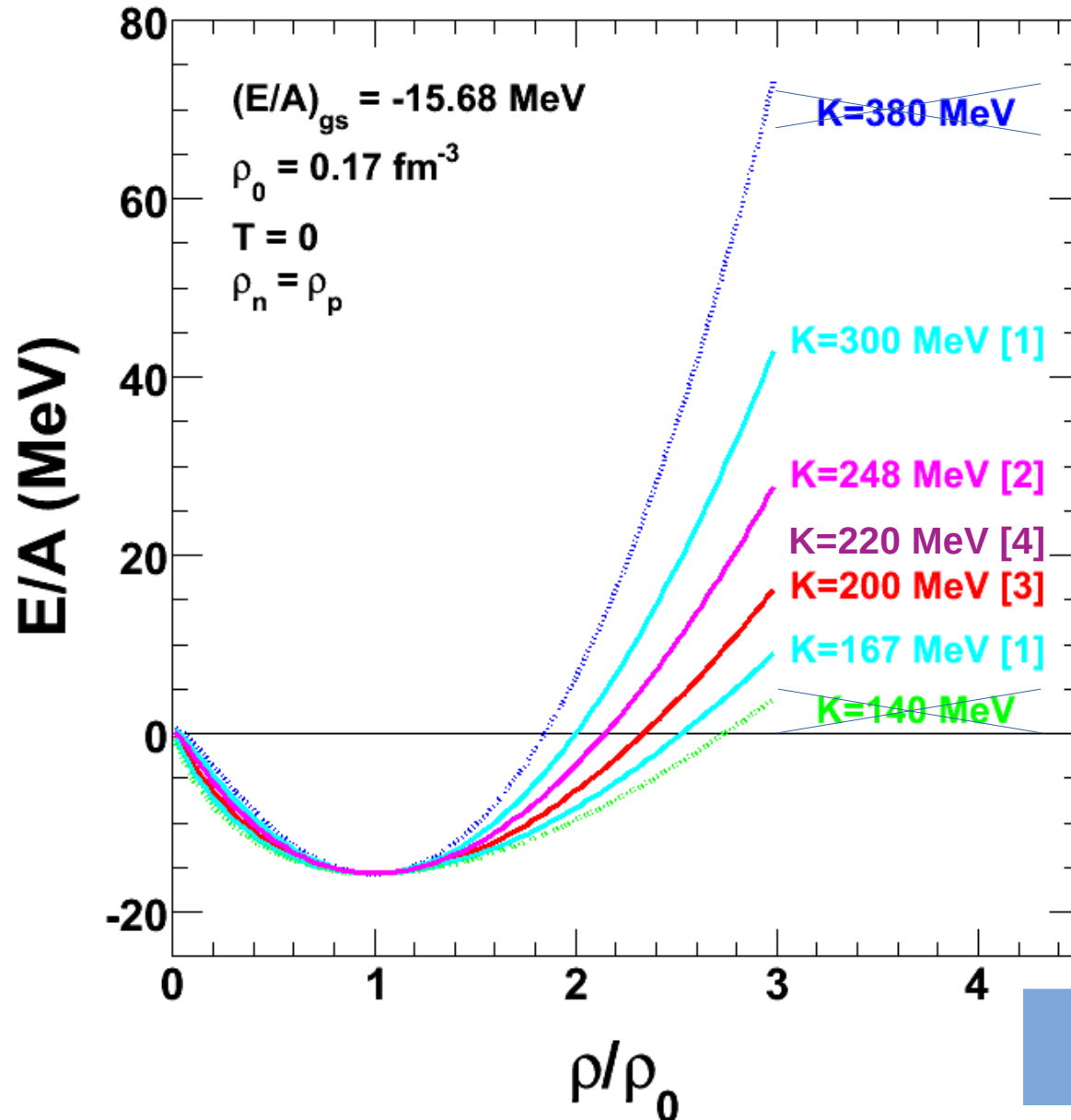
where $k_F = (3\pi^2\rho/2)^{1/3}$ is the nucleon Fermi momentum and $m_0^*/m = (1 + \frac{m}{\hbar^2 k_F} dU_0/dk)^{-1} \Big|_{k_F}$ is the nucleon isoscalar effective mass. The $E_{sym}(\rho)$ consists of two terms, i.e., the kinetic symmetry energy $E_{sym}^1(\rho)$ equivalent to 1/3 the Fermi energy of quasi-nucleons with an isoscalar effective mass m_0^* and the potential symmetry energy $E_{sym}^2(\rho)$ equivalent to 1/2 the isovector potential $U_{sym,1}(\rho, k_F)$ at k_F . The density slope $L(\rho)$ can be cast into five terms

It is important to point out that the $E_{sym}(\rho)$ is closely related to the neutron-proton effective mass splitting $m_{n-p}^* \equiv (m_n^* - m_p^*)/m$, which is a fundamental quantity having broad impacts on many interesting issues in both nuclear physics and astrophysics [23, 24]. In terms of the momentum dependence of the single-nucleon potential or the $E_{sym}(\rho)$ and $L(\rho)$, one has

$$m_{n-p}^* \approx \frac{2\delta m}{\hbar^2 k_F} \left[-\frac{dU_{sym,1}}{dk} - \frac{k_F}{3} \frac{d^2 U_0}{dk^2} + \frac{1}{3} \frac{dU_0}{dk} \right] \left(\frac{m_0^*}{m} \right)^2 \approx \frac{\delta}{E_F(\rho)} \left[3E_{sym}(\rho) - L(\rho) - \frac{1}{3} \frac{m}{m_0^*} E_F(\rho) \right] \left(\frac{m_0^*}{m} \right)^2$$

where $E_F(\rho)$ is the Fermi energy in SNM at density ρ . Therefore, while probing the density dependence of the nuclear symmetry energy, we are also studying the neutron-proton effective mass splitting in neutron-rich nuclear matter [23-25].

K for symmetric matter ($\delta=0$)



[1] Flow:

P. Danielewicz et al., Science 298 (02) 1592
 BME: $K = 167\text{-}300 \text{ MeV}$

[2] ISGMR:

J. Piekarewicz, PRC 69 (04) 041301

RMF: $K=248 \text{ MeV}$

G. Colò et al., PRC 70 (04) 024307

Skyrme HF: $K=230 \text{ MeV}$

S. Shlomo, et al. Eur. Phys. J. A 30, 23 (2006)

$K = 240 \pm 20 \text{ MeV}$

[3] Subthreshold K^+ :

C. Sturm et al., PRL 86 (01) 39

Ch. Hartnack et al., PRL 96 (06) 012302

Flow:

W. Reisdorf (FOPI) arXiv:1307.4210 [nucl-ex]

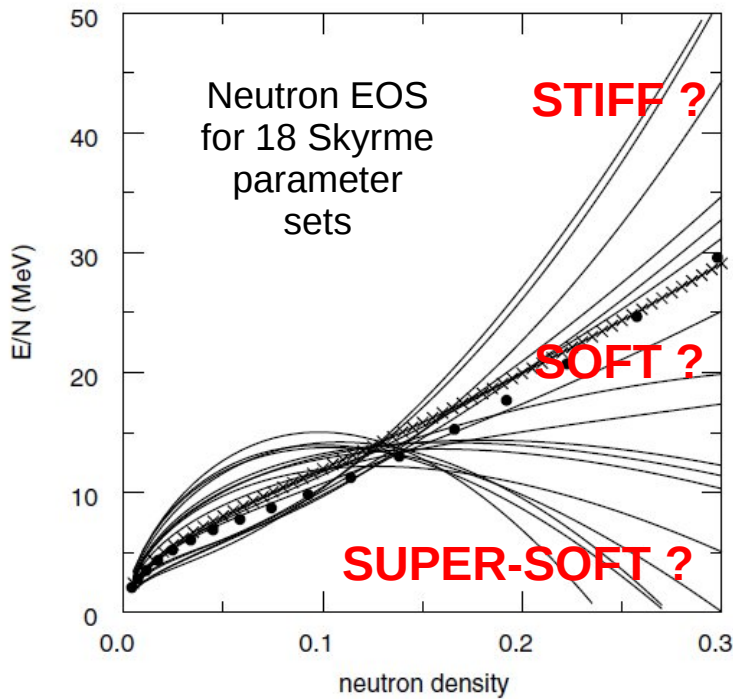
IQMD: $K=200 \text{ MeV}$

[4] Y. Wang et al., PLB 778 (2018) 207

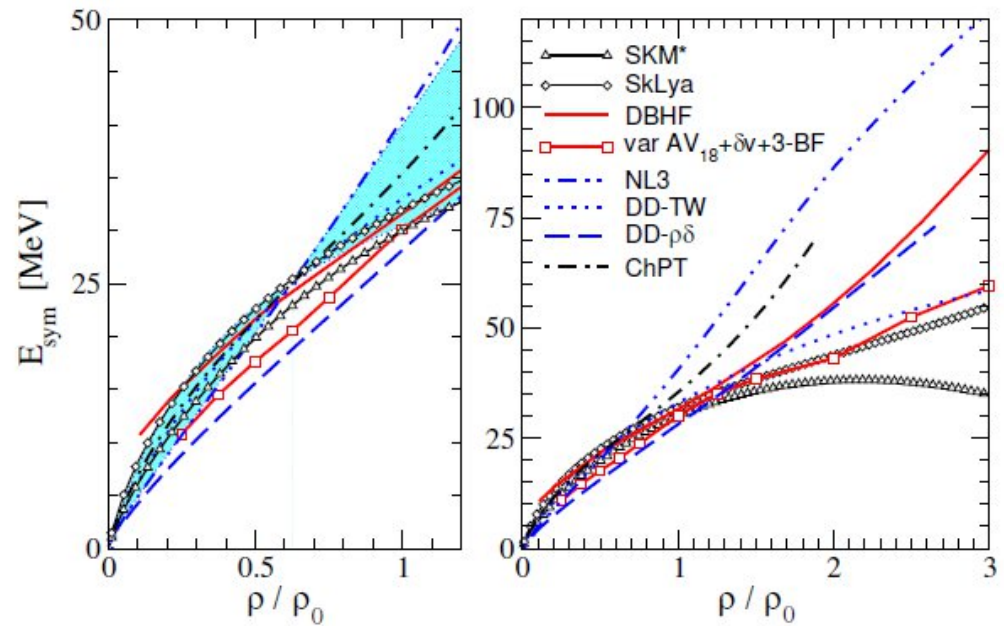
UrQMD + FOPI: $K = 220 \pm 40 \text{ MeV}$

“Generally accepted” value of
 $K = 240 \pm 20 \text{ MeV}$

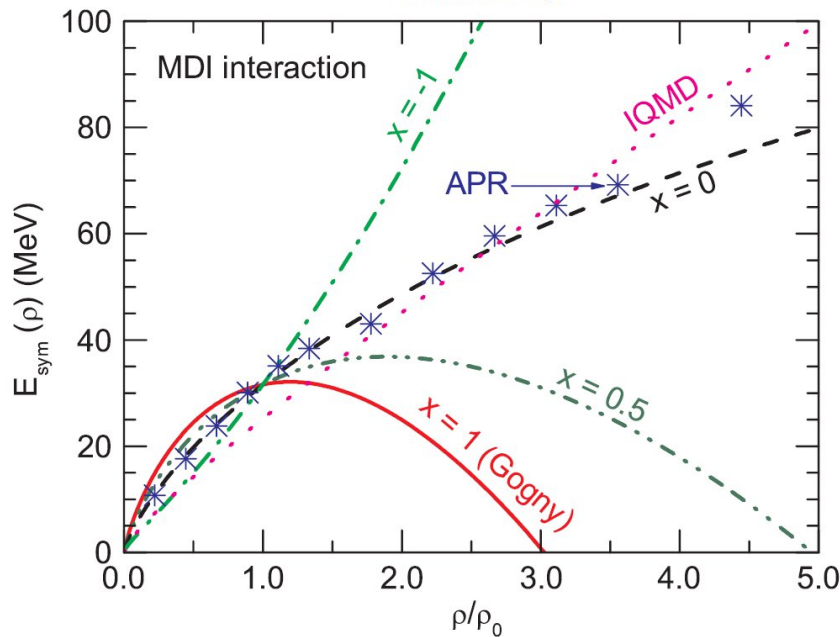
Symmetry term. Why so uncertain?



B. Alex Brown, PRL 85(2000)5296



C. Fuchs and H.H. Wolter, EPJA 30(2006)5



Z. Xiao et al., PRL 102(2009)062502

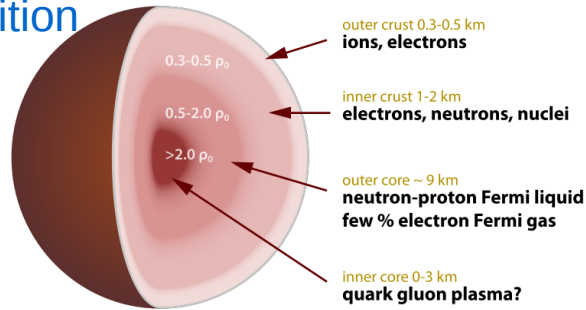
Symmetry energy uncertain at high density and modified by clustering at low density

Phenomenological forces constrained around saturation and for nearly isospin-symmetric matter. Poor knowledge of effective forces in neutron-rich matter. Uncertainties in the nature of the three-neutron force. Uncertain extrapolations above the saturation density.

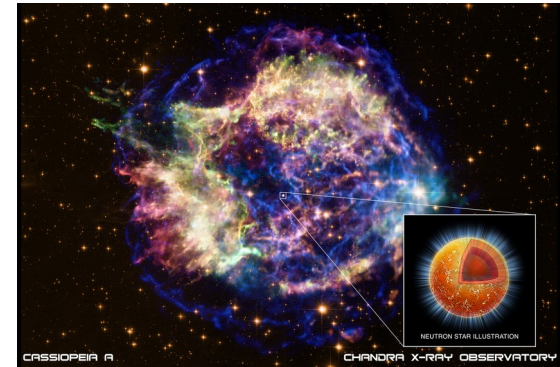
Why so important?

ASTROPHYSICS

- Neutron star structure, composition, size, mass and cooling rate, crust-core transition density
- Pulsars, masses, spin rates
- Supernova explosions
- Stellar nucleosynthesis, r-process



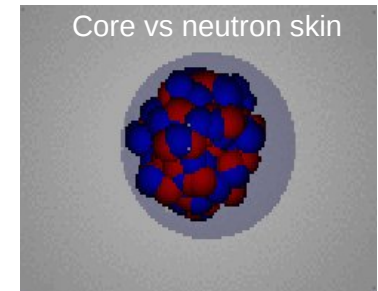
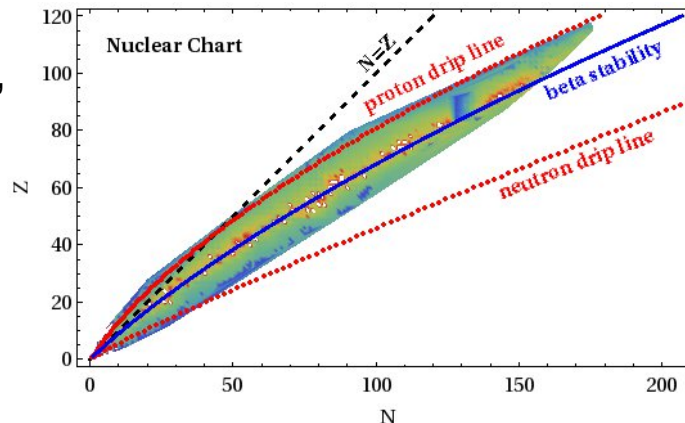
Robert Schulze (Wikipedia)



Cassiopeia A Supernova Remnant
(<http://chandra.harvard.edu/photo/printgallery/2004/>)

STRUCTURE

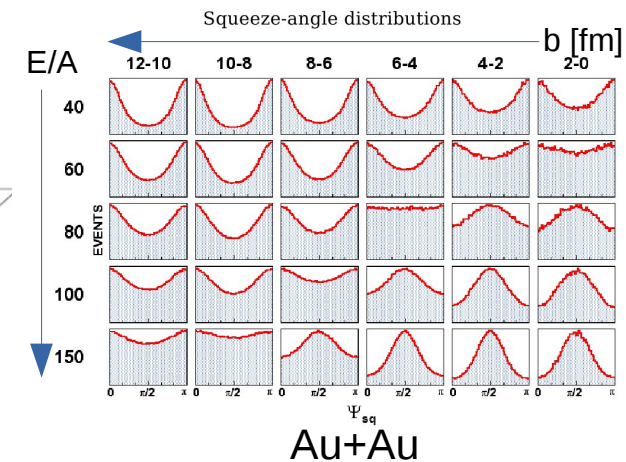
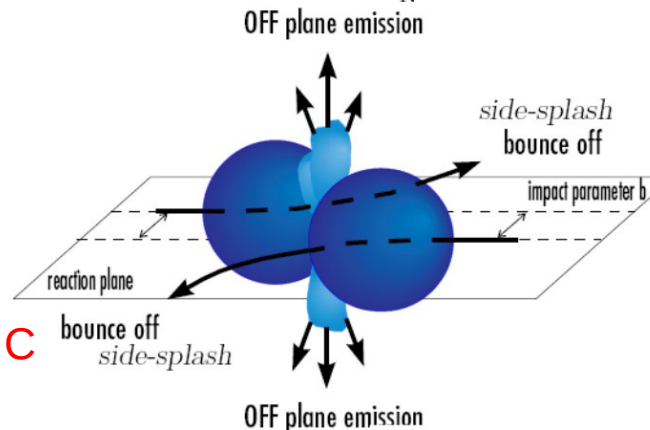
- Structure of exotic nuclei (masses, drip lines)
- Neutron skin thickness
- IvGDR
- Pygmy resonances
- Differences between IAS



PDR

REACTIONS

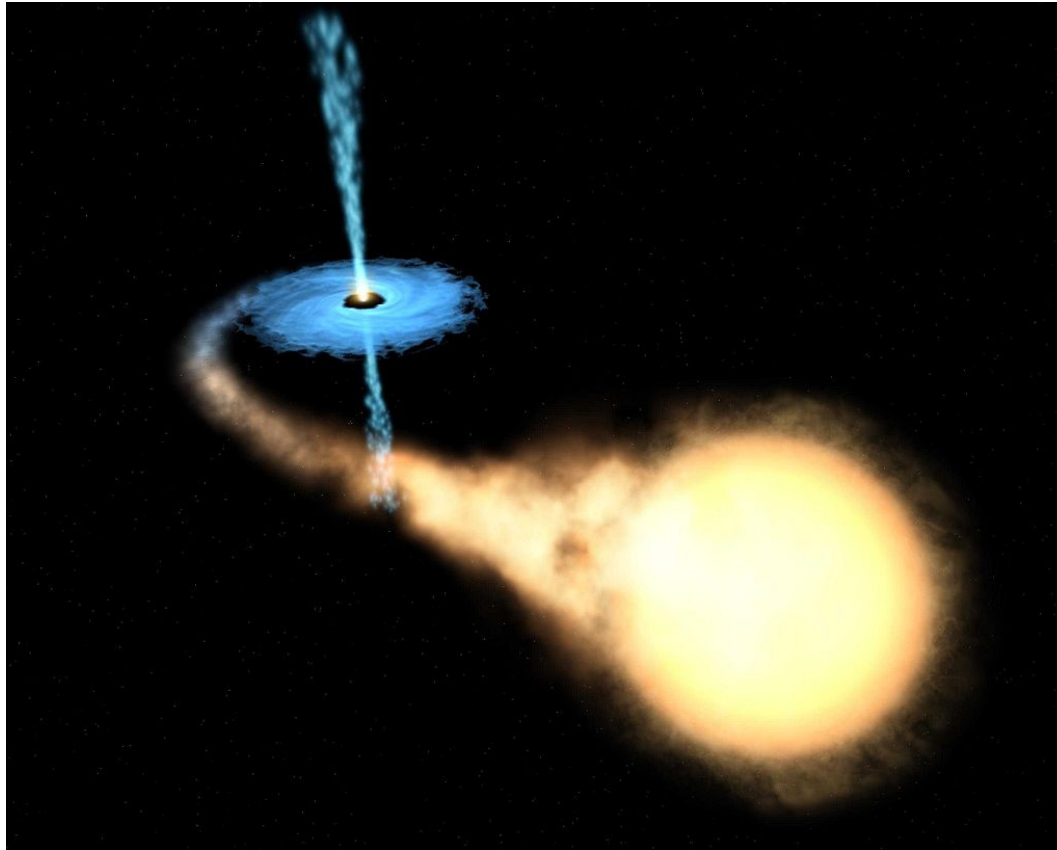
- Flow patterns in HIC
- Multifragmentation, isoscaling, isospin diffusion
- n/p , $t/{}^3\text{He}$, π^-/π^+ , K^+/K^0 ratios in HIC



The Mass & Radius of Neutron Stars

from qLMXB

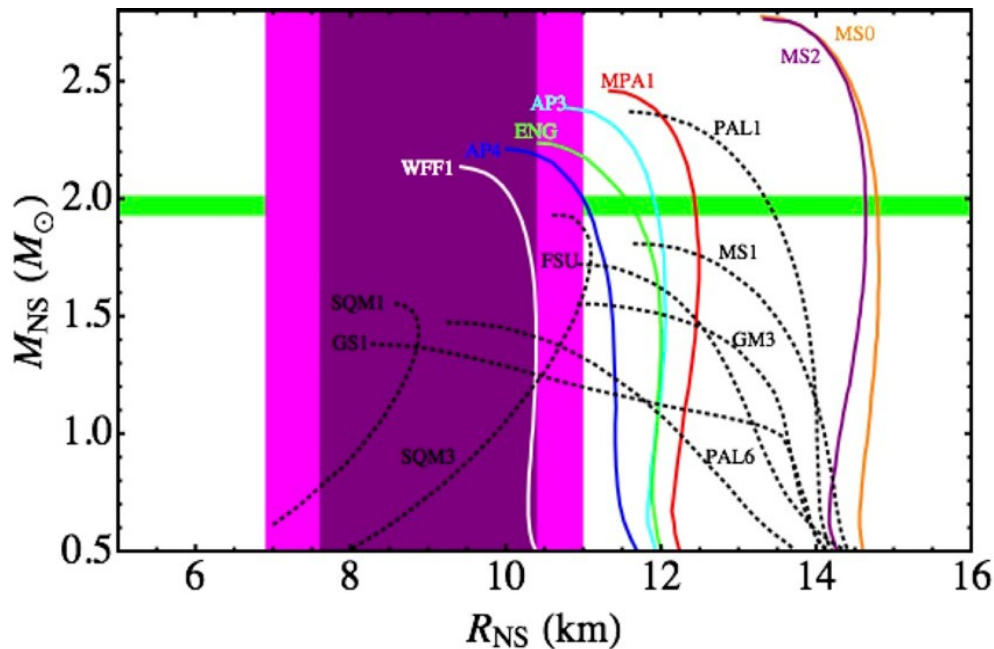
A “nuclear diesel”



The Mass & Radius of Neutron Stars

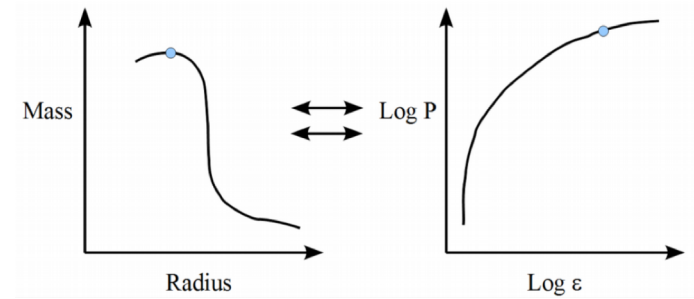
from qLMXB

Guillot et al. (2013)

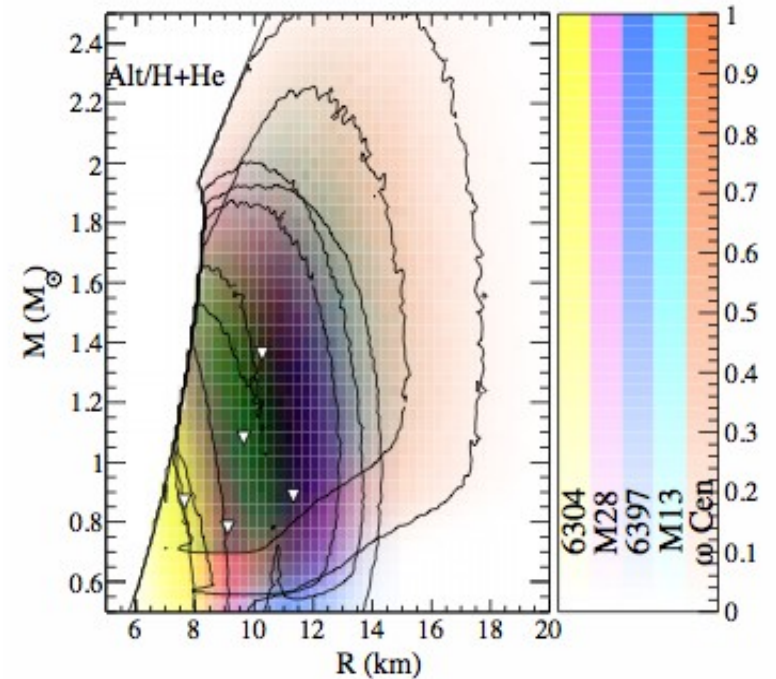


$$R_{\text{NS}} = 9.1_{-1.5}^{+1.3} \text{ km (90\%-confidence)}$$

WFF1: $E_{\text{sym}} \approx 26 \text{ MeV}$, $L \approx 60 \text{ MeV}$



Lattimer & Steiner (2013)



$$10.9 < R_{\text{NS}} < 12.7 \text{ km (2.0 } M_{\odot})$$

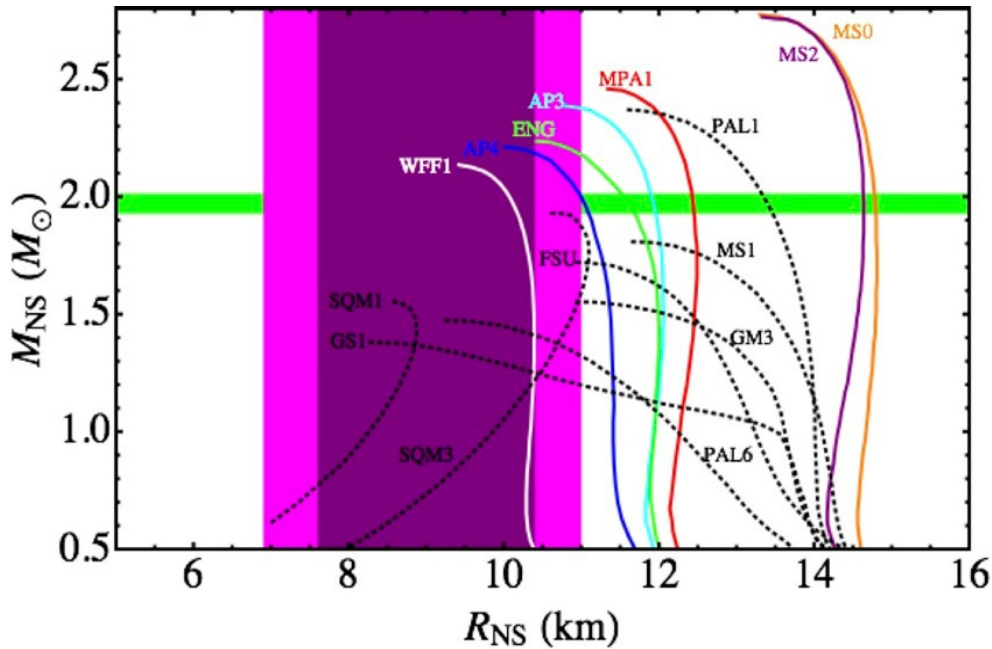
$41 < L < 83 \text{ MeV}$

(analysis of the same data)

The Mass & Radius of Neutron Stars

from qLMXB

Guillot et al., ApJ 772, 7 (2013)



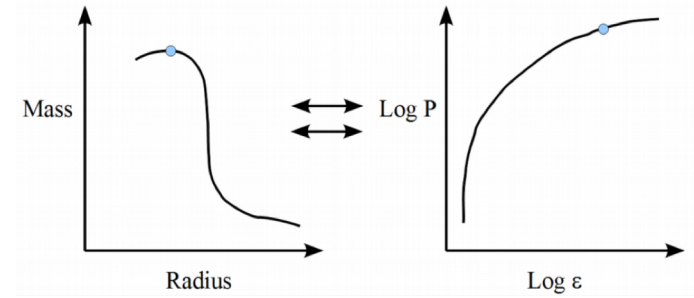
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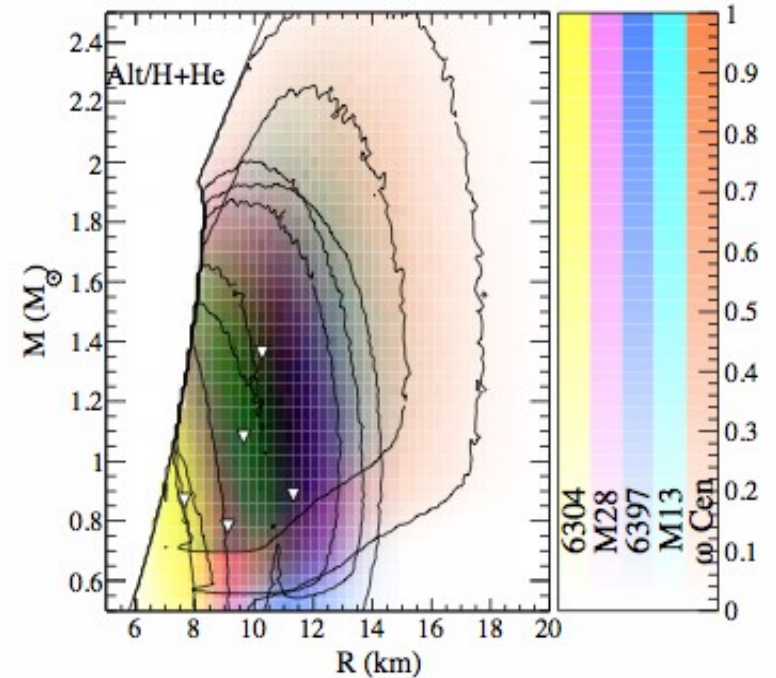
updates

S. Guillot et al. ApJL, 796:L3,2014: $R_{\text{NS}} = 9.4 \pm 1.2 \text{ km}$

F. Özel et al., ApJ, 820:28,2016: $R_{\text{NS}} = 10.1\text{-}11.1 \text{ km (} 1.5 M_{\odot} \text{)}$



Lattimer & Steiner (2013)



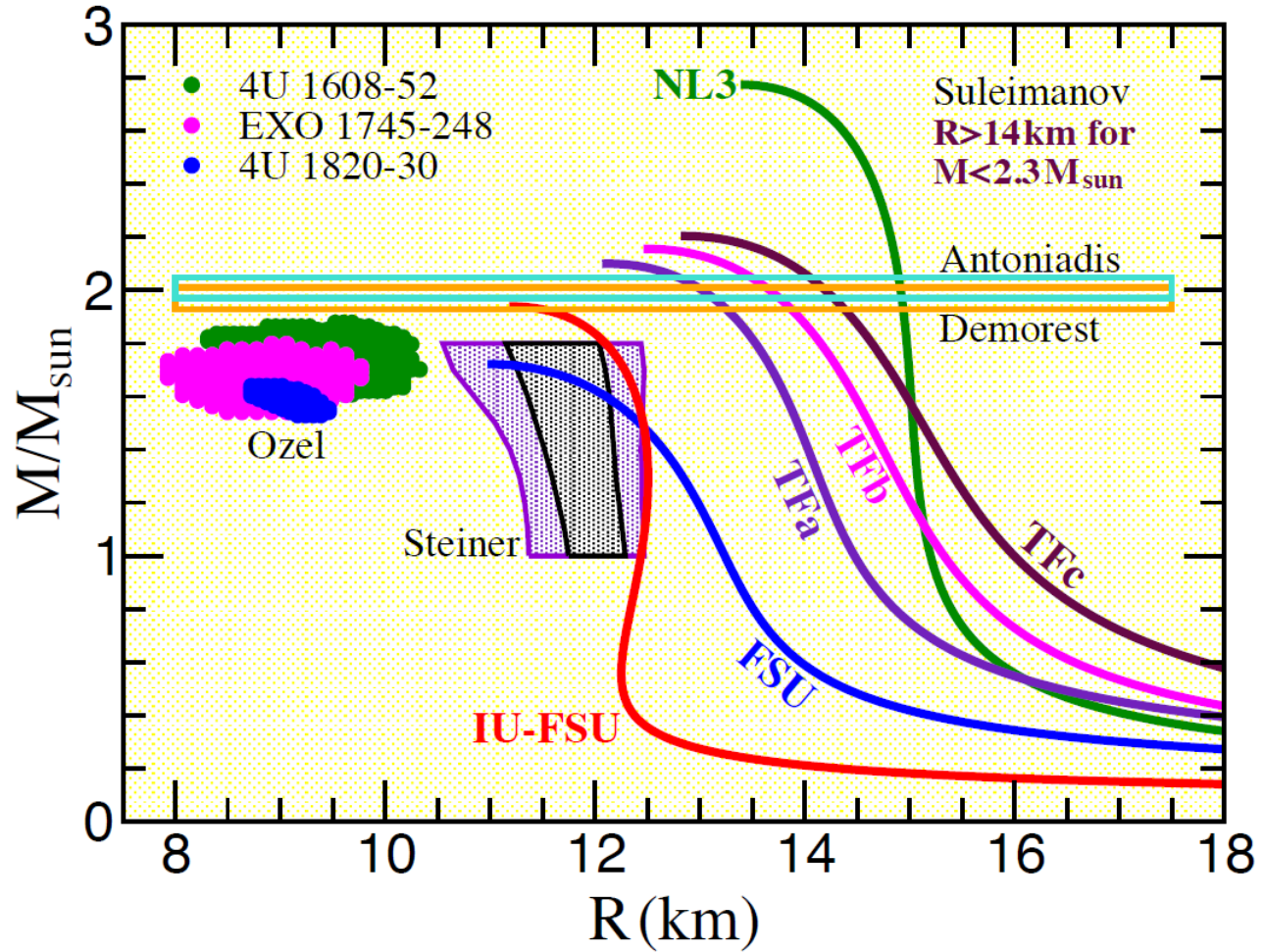
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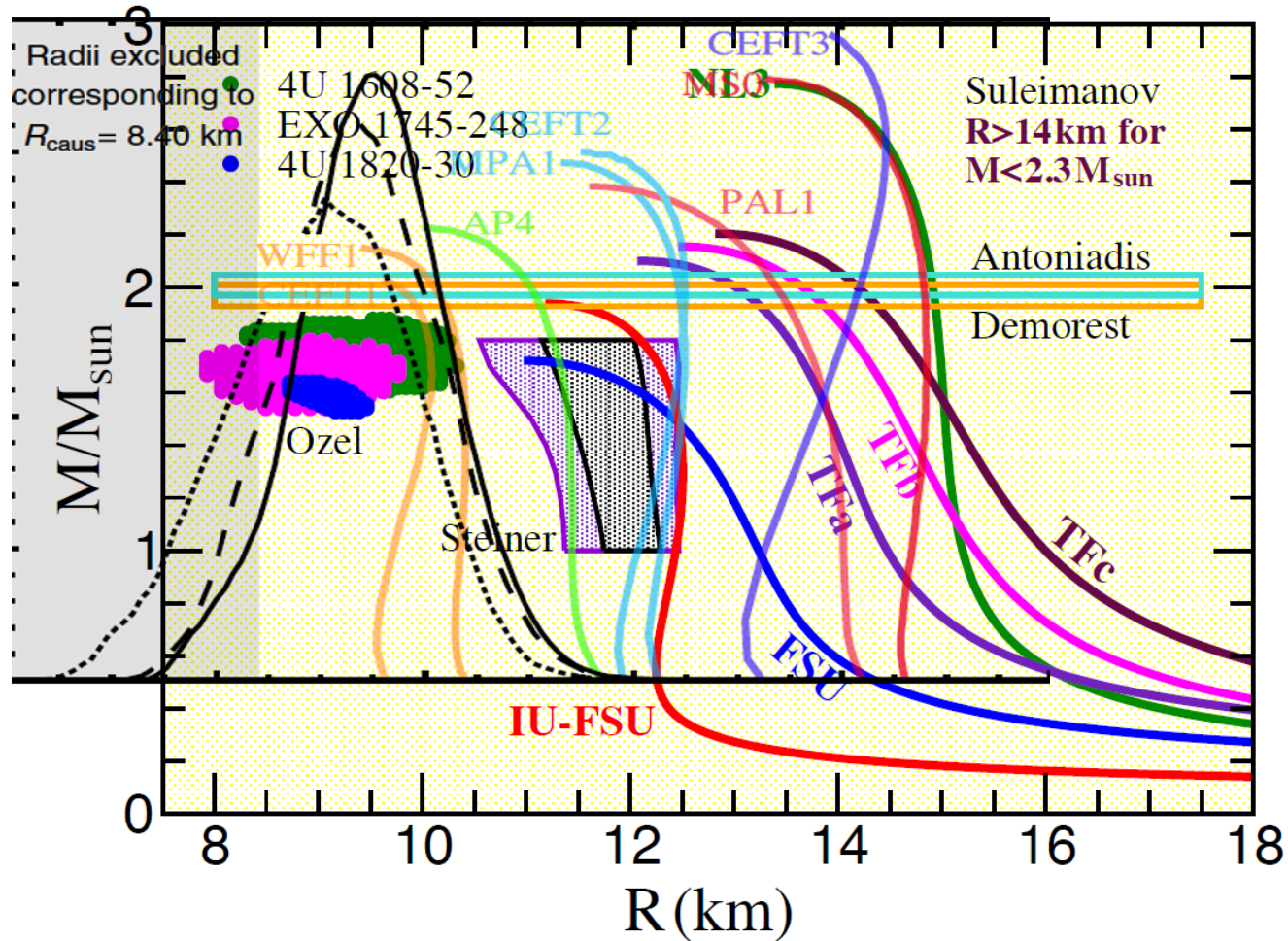
A.W. Shaw et al., arXiv:1803.00029 [astro-ph.HE]

$$R_{\text{NS}} = 12_{-1.7}^{+1.9} \text{ km (} 1.4 M_{\odot} \text{)}$$

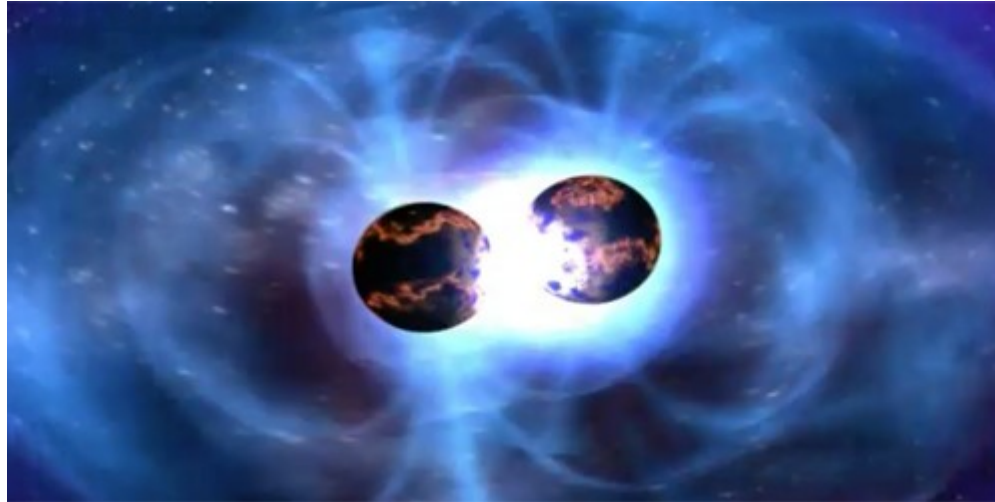
current status



current status

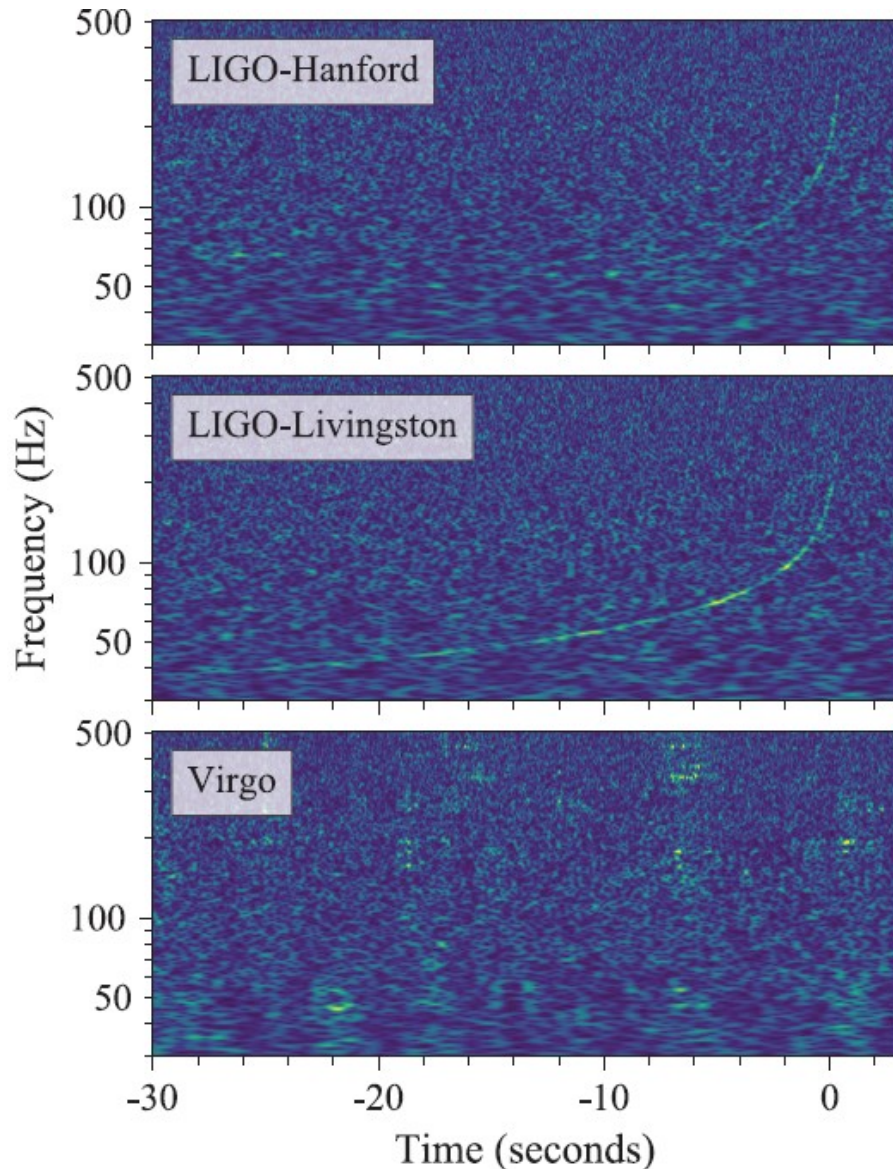


Waltzing Neutron Stars shaking the Universe



NASA/Goddard Space Flight Center

Gravitational Wave Constraints



B. P. Abbott et al. PRL 119, 161101 (2017)
(LIGO and Virgo Collaborations)
GW170817: Observation of Gravitational Waves
from a Binary Neutron Star Inspiral



tidal (quadrupolar) deformability $\Lambda \sim (R/M)^5$
(the stiffness of a star to the stress
caused by its companion's gravity,
measures deviation of a gravitational
field from that of a pointlike star) $\Lambda < 800$

and

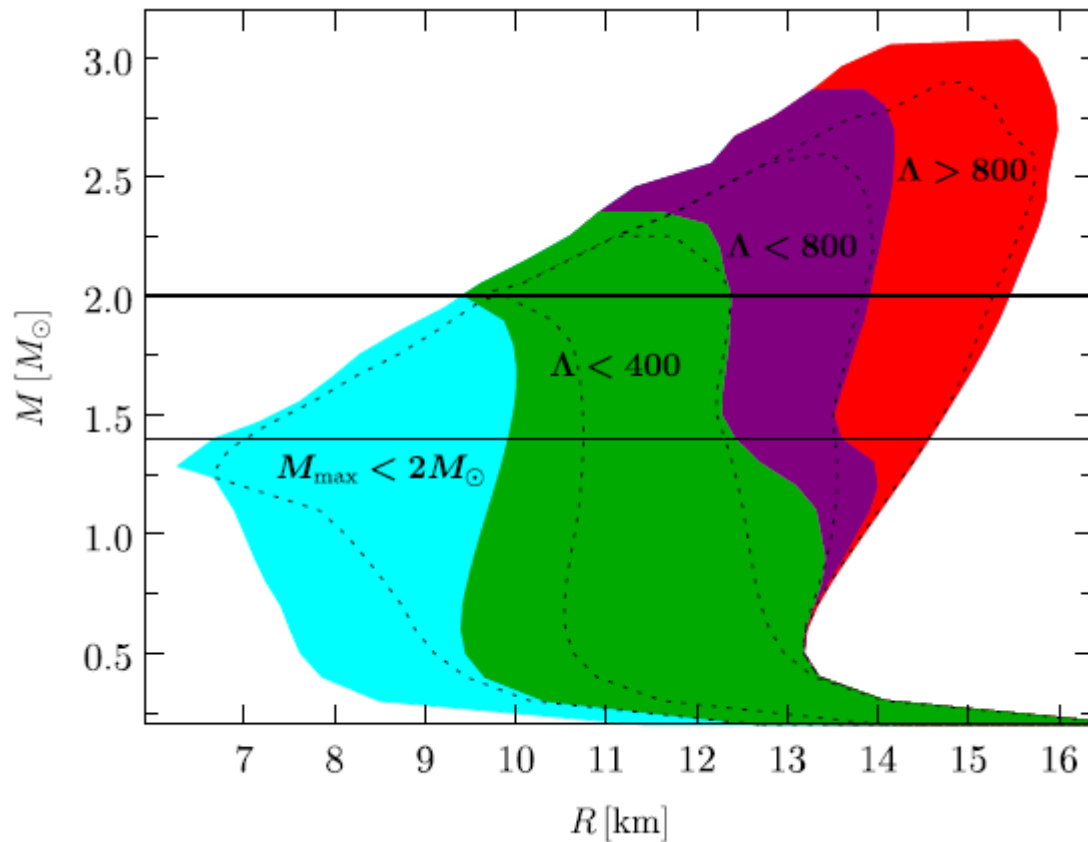
component masses in the range $1.17\text{--}1.60 M_{\odot}$
total mass $2.74 M_{\odot}$

(spin dependent)

were inferred by matching the observed waveform
with a frequency-domain post-Newtonian
waveform model.

GW170817 constraints on the NS Radius and the EOS

E. Annala et al. PRL 120, 172703 (2018)



$$9.9 < R(1.4 M_{\odot}) < 13.6 \text{ km}$$

$$\Lambda(1.4 M_{\odot}) \geq 120$$

J.M. Lattimer, M. Prakash,
Physics Reports 621 (2016) 127:

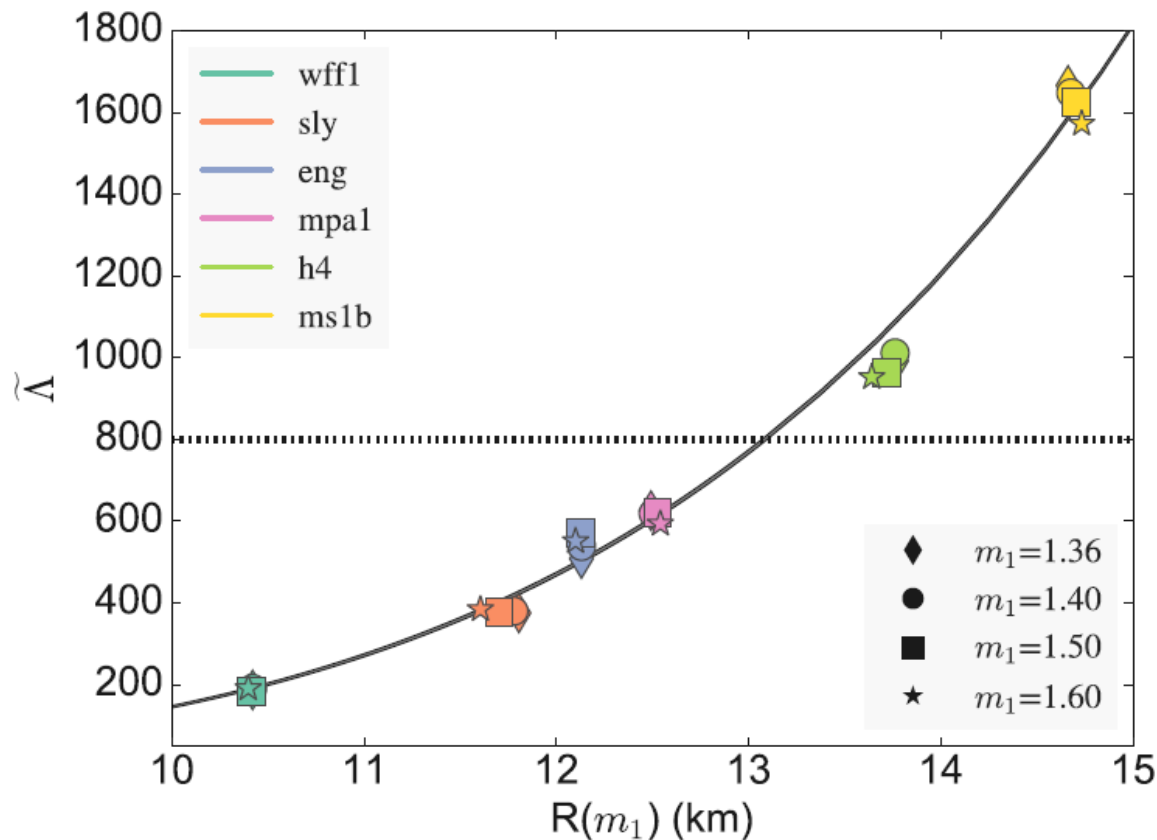
$$\blacktriangleright \sim 26 < L < \sim 84 \text{ MeV}$$

Compatible with larger NS from
qLMBX analyses and thus
supporting the stiffer E_{sym} .

The mass-radius clouds corresponding to our EOSs

GW170817 constraints on the NS Radius and the EOS

C.A. Raithel, F. Özel, & D. Psaltis, AJL 857:L23 (2018)



$R(1.4 M_{\odot}) < 13 \text{ km}$
(no lower limit)

Bayesian Inference:

→ most probable $\Lambda = 400$
→ most likely radius $R_{\text{NS}} = 11.7 \text{ km}$

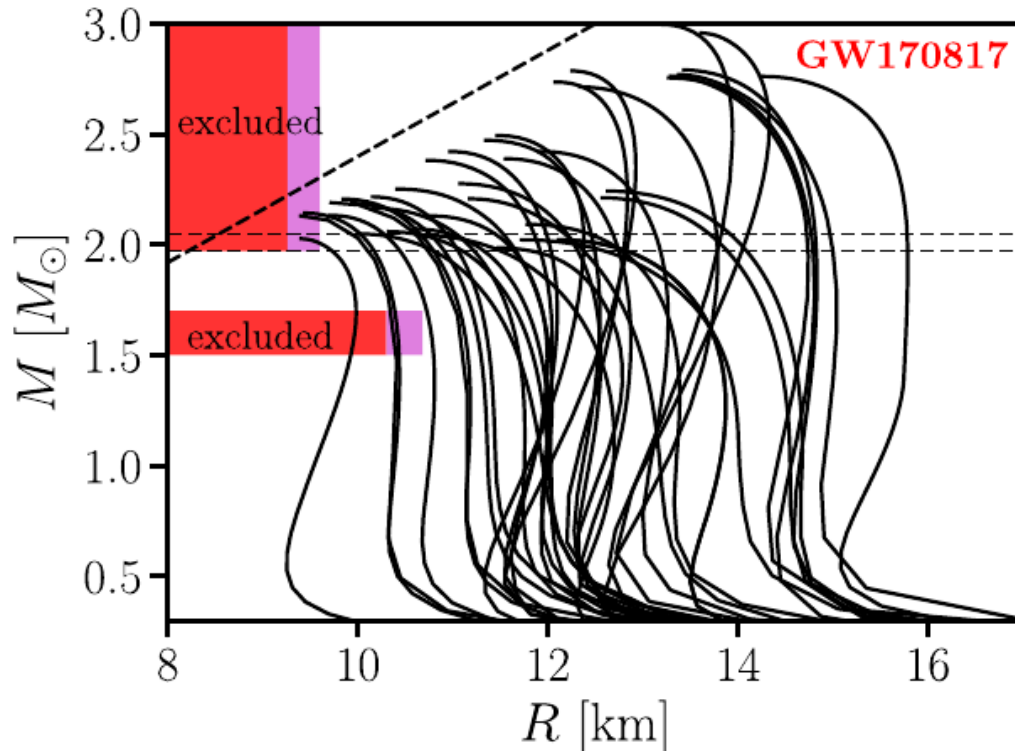
small mass dependence

Supporting rather larger NS

Effective tidal deformability of the binary system as a function of the radius of the primary neutron star, for 6 EoSs.

GW170817 constraints on the NS Radius and the EOS

A. Bauswein et al., Astr. J. Lett., 850:L34, 2017



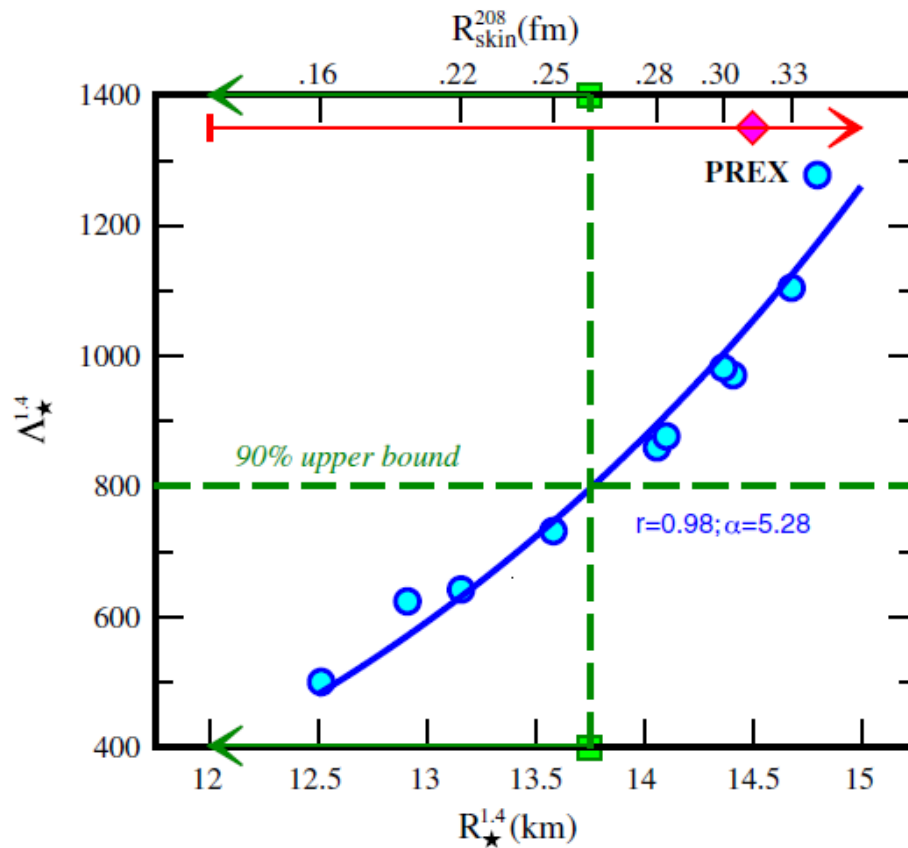
$$R(1.6 M_{\odot}) > 10.68 \text{ km}$$

Lower limit on the stellar radius of a $1.6 M_{\odot}$ neutron star was obtained under the assumption that the merger did not result in a prompt collapse.

Mass–radius relations of different EoSs with very conservative (red area) and “realistic” (cyan area) constraints of this work for $R_{1.6}$ and R_{max} . Horizontal lines display the limit by Antoniadis et al. (2013). The dashed line shows the causality limit.

GW170817 constraints on the NS Radius and the EOS

F. J. Fattoyev, J. Piekarewicz, and C. J. Horowitz, PRL 120, 172702 (2018)



$$R(1.4 M_{\odot}) < 13.76 \text{ km}$$

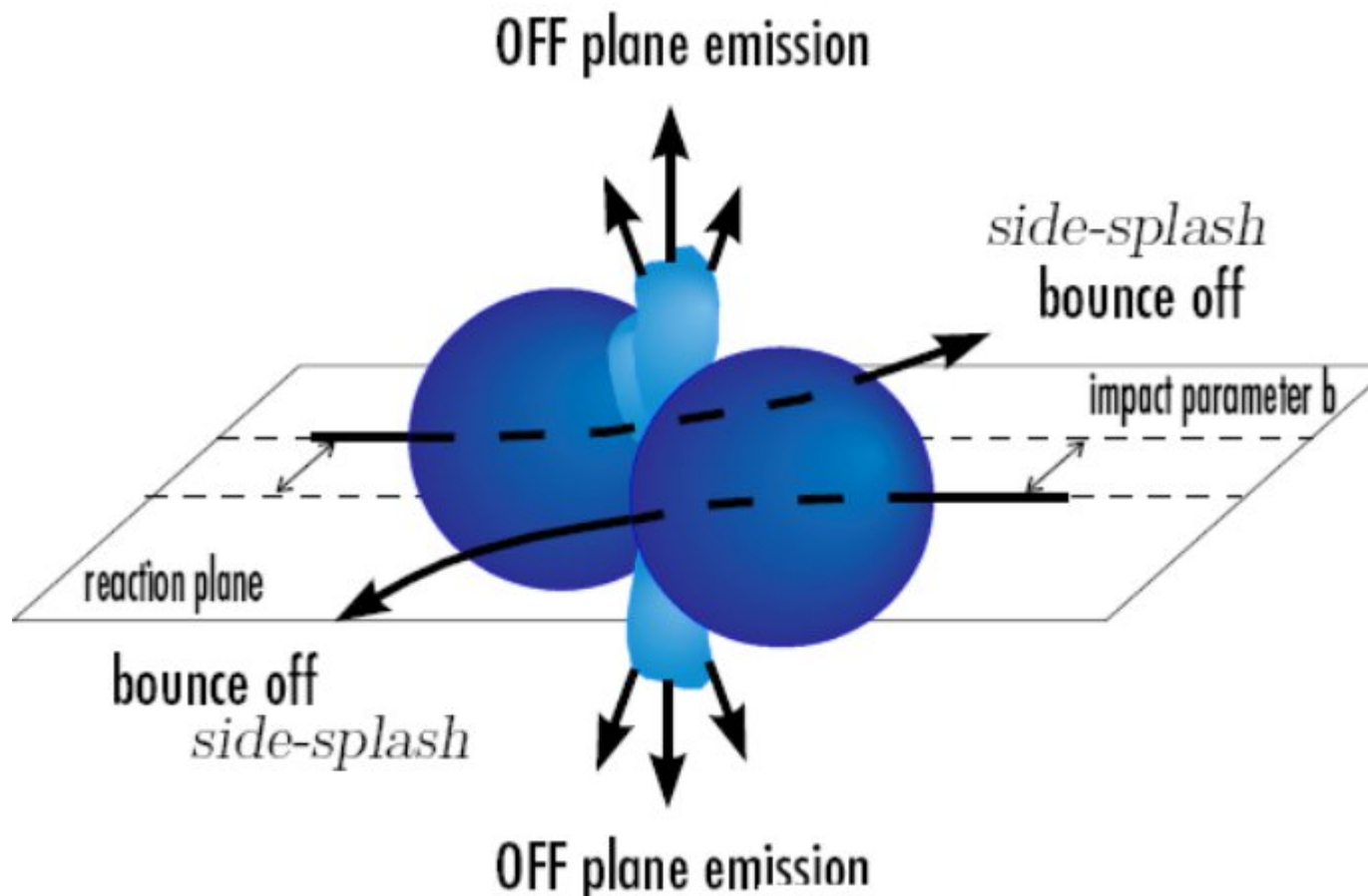
$$\Lambda(1.4 M_{\odot}) \geq 490 \text{ (PREX lower limit)}$$

$$L < 82.5 \text{ MeV (TAMUC-FSUa)}$$

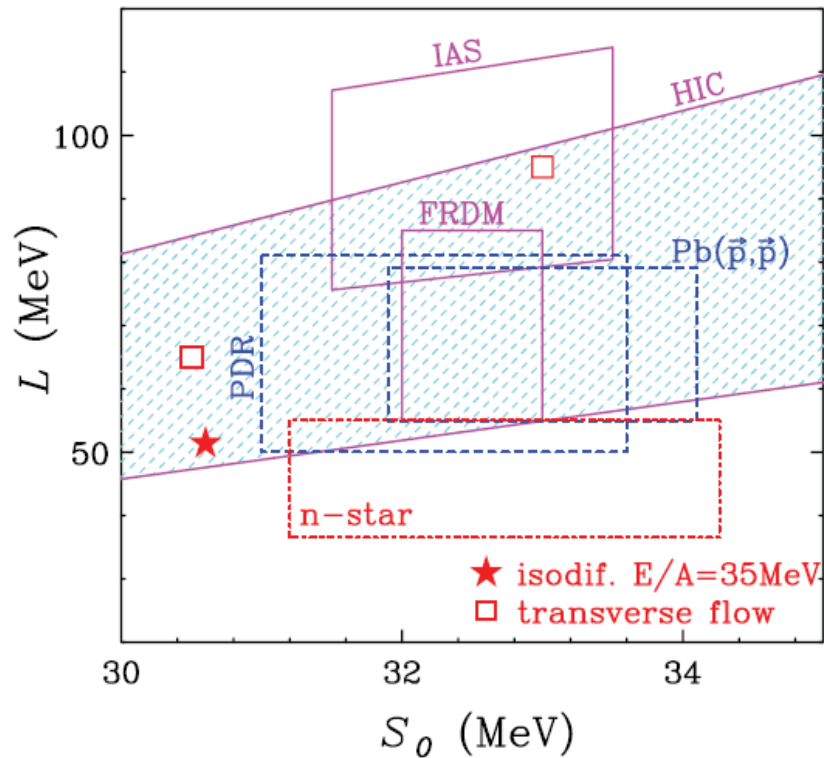
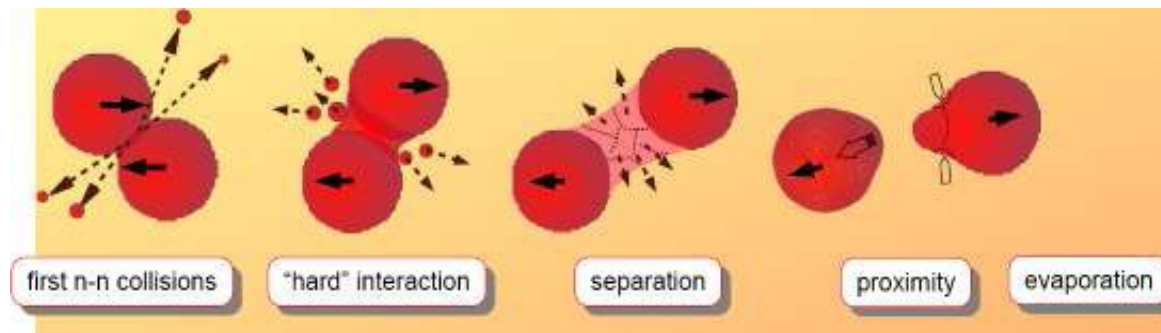
Supporting rather larger NS.
RN and Rskin not quite compatible

The dimensionless tidal polarizability $\Lambda_{1.4}$ of a $1.4M_{\odot}$ neutron star as a function of the neutron-skin thickness of ^{208}Pb (upper abscissa) and the radius of a $1.4M_{\odot}$ neutron star (lower abscissa) as predicted by the 10 RMF models.

Squeezing the Symmetry Energy out of Heavy Ion Collisions



Symmetry energy from low and intermediate energy HI collisions

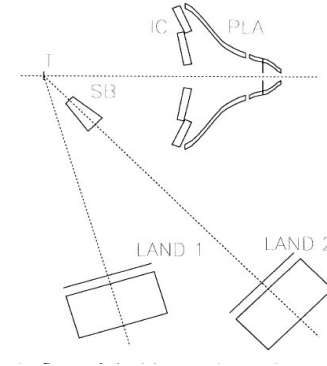


(S_0, L) centered on $\sim (32.5, 70)$ MeV over the range of density between $0.3-1 \rho_0$

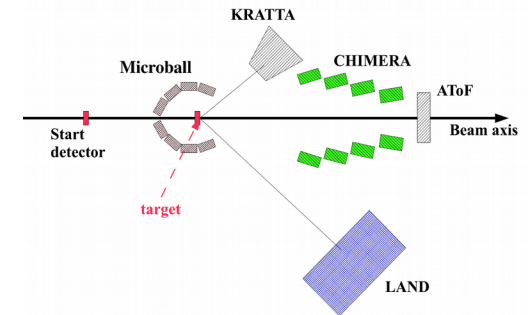
HIC - heavy-ion collisions, isospin diffusion, isoscaling, n/p ratios, Tsang et al., 2009, 2012
IAS - isobaric analog states, Danielewicz & Lee 2009, 2012
PDR - pygmy dipole resonances, Klimkiewicz et al. 2007
FRDM – finite-range droplet model, binding energy fits, Moeller et al. 2012
Pb(p,p) – neutron skin thickness from p elastic scattering, Zenihiro et al. 2010

High energy (density) experiments

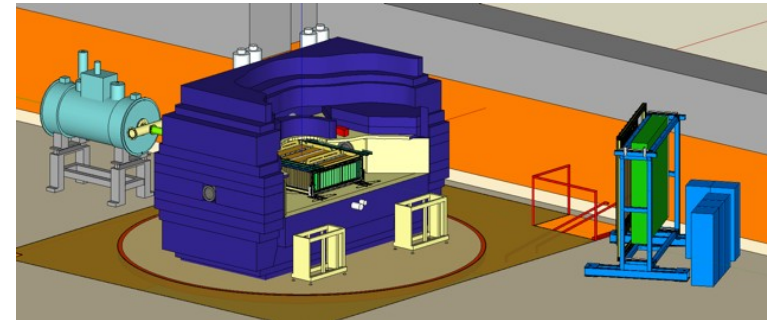
(1) FOPI/LAND data (GSI 1993/2011)



(2) ASY-EOS experiment (GSI 2011)

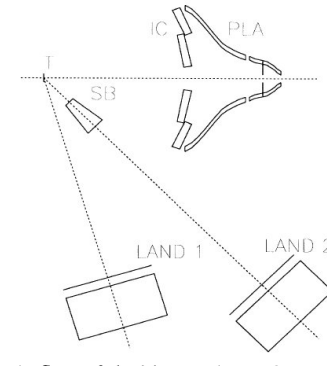


(3) SPiRIT experiment (RIKEN 2016)

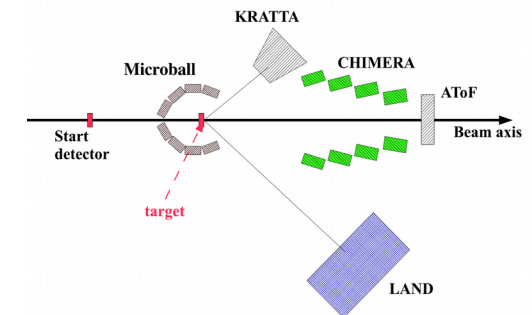


High energy (density) experiments

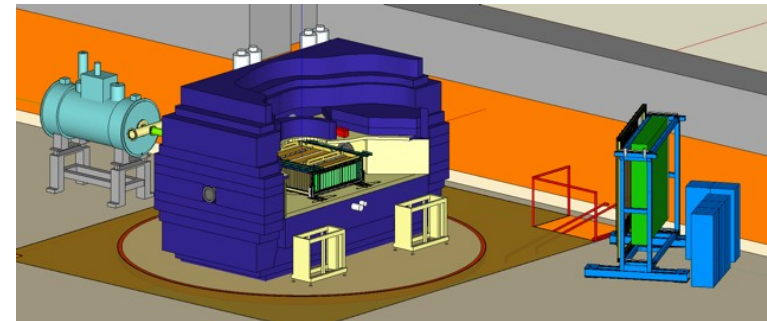
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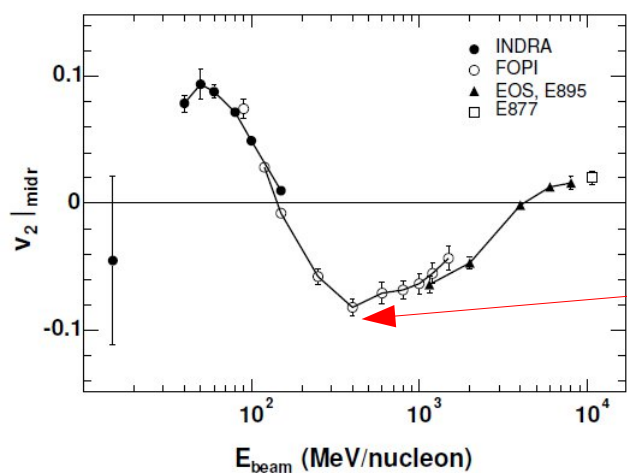
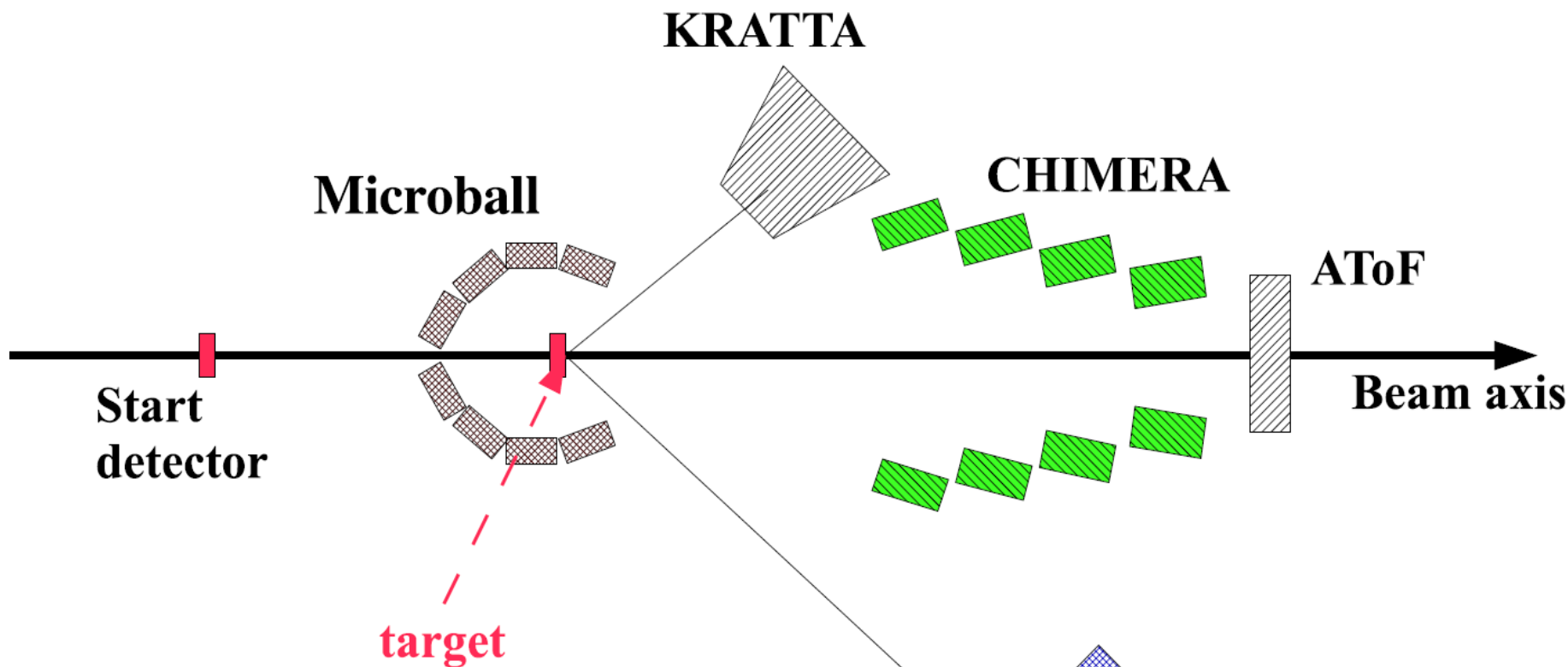
(2) ASY-EOS experiment (GSI 2011)



(3) SPiRIT experiment (RIKEN 2016)

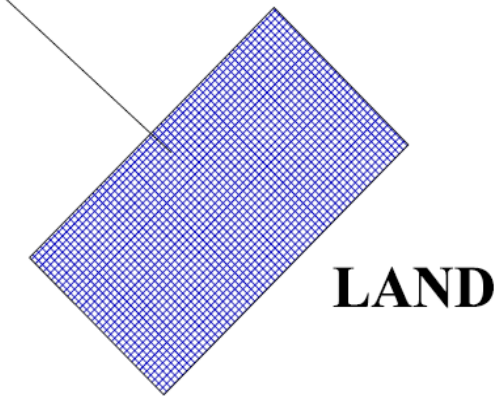


ASY-EOS experimental setup



A. Andronic et al. EPJA 30(2006)31

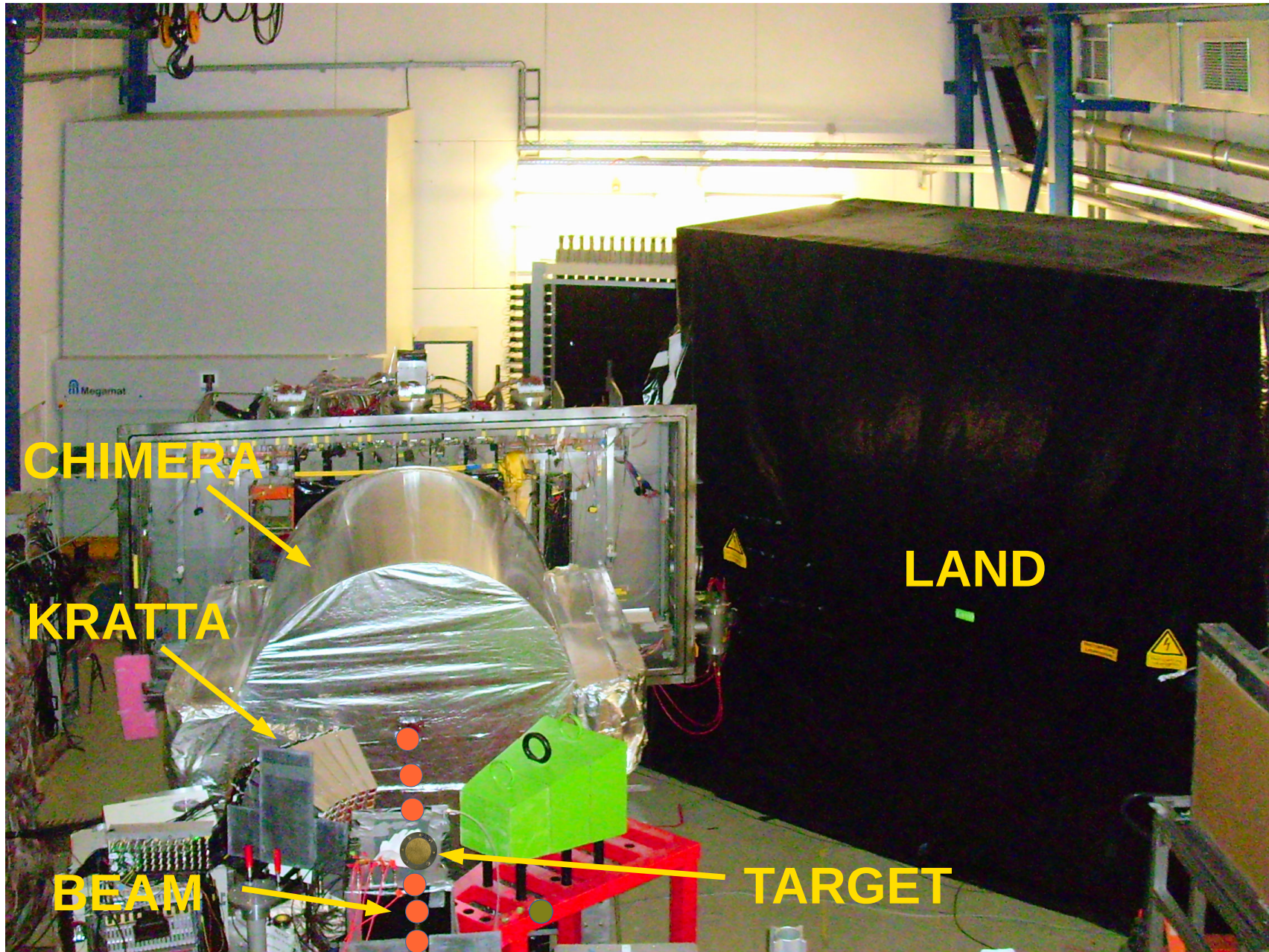
^{197}Au	+	^{197}Au	@ 400 AMeV	$\delta^2 = 0.039$
^{96}Zr	+	^{96}Zr	@ 400 AMeV	$\delta^2 = 0.028$
^{96}Ru	+	^{96}Ru	@ 400 AMeV	$\delta^2 = 0.007$



LAND

KRATTA @ GSI (Au+Au at 400 AMeV, 2011)

ASY-EOS Experimental Setup

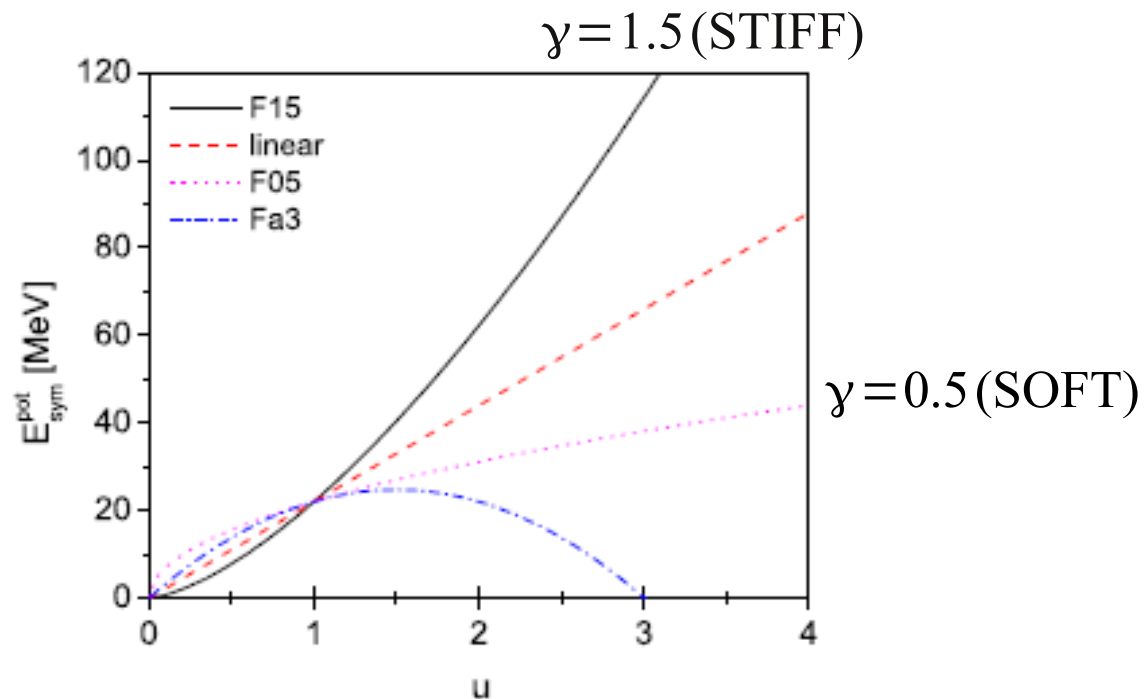


Model assumptions

UrQMD, Q. Li, J.Phys. G 31(2005)1359

„Fermi-gas” parametrization of the symmetry term:

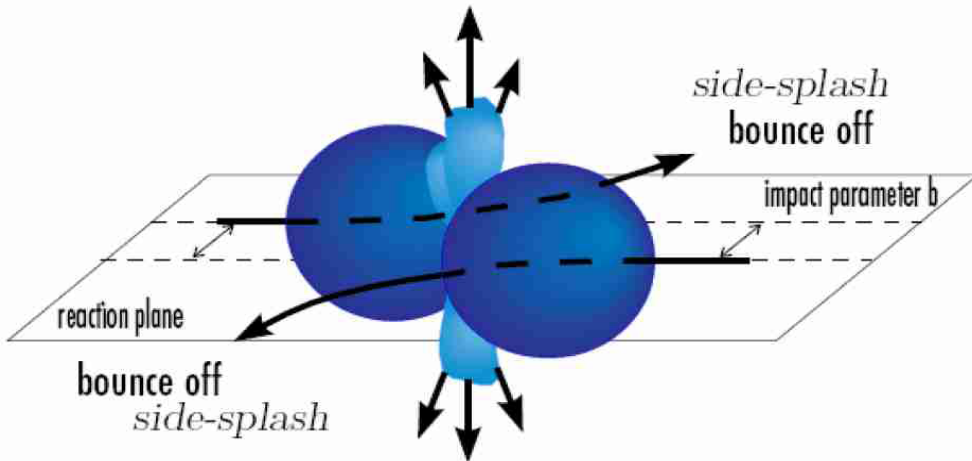
$$E_{sym} = E_{sym}^{pot} + E_{sym}^{kin} = 22 \text{ MeV} \left(\frac{\rho}{\rho_0} \right)^\gamma + 12 \text{ MeV} \left(\frac{\rho}{\rho_0} \right)^{2/3}$$



Flow analysis (v_1, v_2)

elliptic flow

OFF plane emission



directed flow

OFF plane emission

Fourier decomposition of **azimuthal distributions** with respect to the **reaction plane** (ϕ_R):

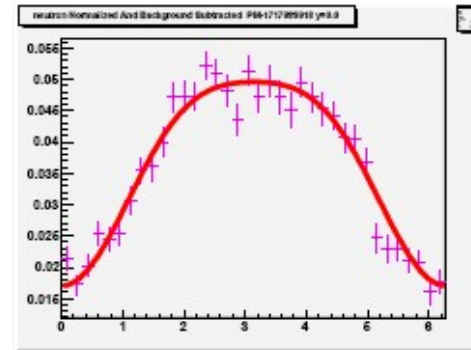
$$\frac{dN}{d(\phi - \phi_R)} = \frac{N_0}{2\pi} \left(1 + 2 \sum_{n \geq 1} v_n \cos n(\phi - \phi_R) \right)$$

directed flow: $v_1 \equiv \langle \cos(\phi - \phi_R) \rangle$

elliptic flow: $v_2 \equiv \langle \cos 2(\phi - \phi_R) \rangle$

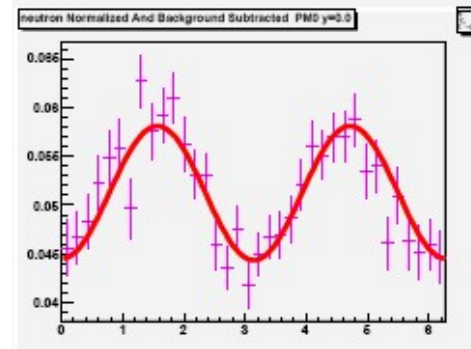
$$v_n = v_n(b, Z, A, y, p^T)$$

W. Trautmann et al., 2011

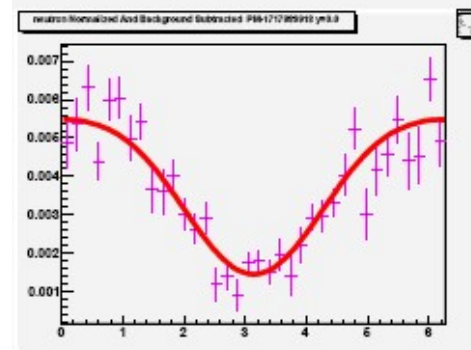


Au+Au @ 400 AMeV

Target rapidity
(strong directed flow)



Mid-rapidity
(strong elliptic flow-squeeze-out)



Projectile rapidity
(strong directed flow)

0 $\Delta\phi$ 2π

FOPI-LAND

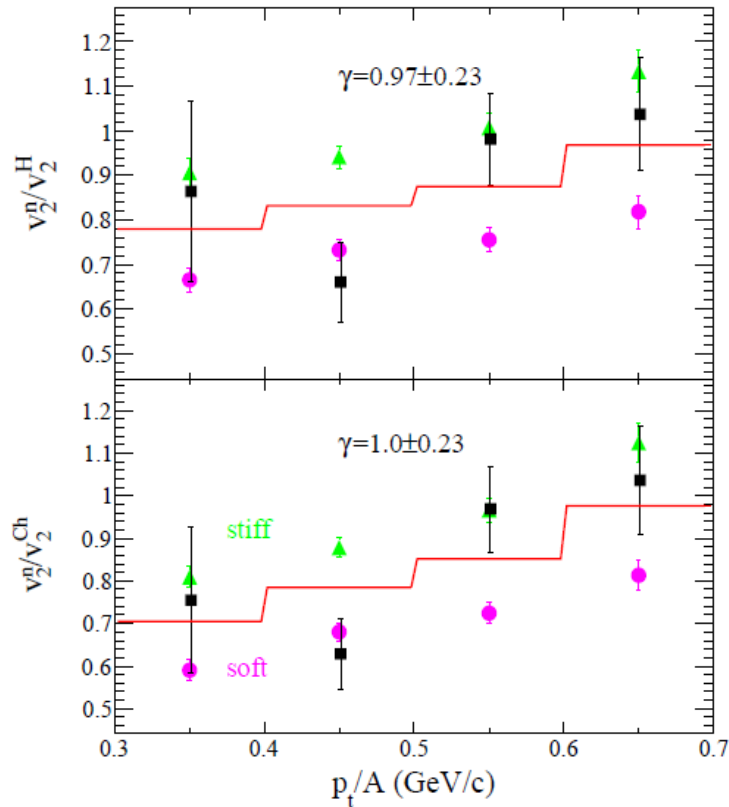


FIG. 14: (Color online) FOPI-LAND data: Elliptic flow ratio of neutrons over all hydrogen isotopes (top) and of neutrons over all charged particles (bottom) for moderately central ($b < 7.5$ fm) collisions of $^{197}\text{Au}+^{197}\text{Au}$ at 400 MeV/nucleon, as a function of the transverse momentum per nucleon p_t/A . The black symbols represent the experimental data. The UrQMD predictions for stiff ($\gamma = 1.5$, green symbols) and soft ($\gamma = 0.5$, purple symbols) are shown. The red line in each panel is the result of a linear interpolation between the predictions; the obtained gamma values and their uncertainties are indicated.

LAND@ASY-EOS

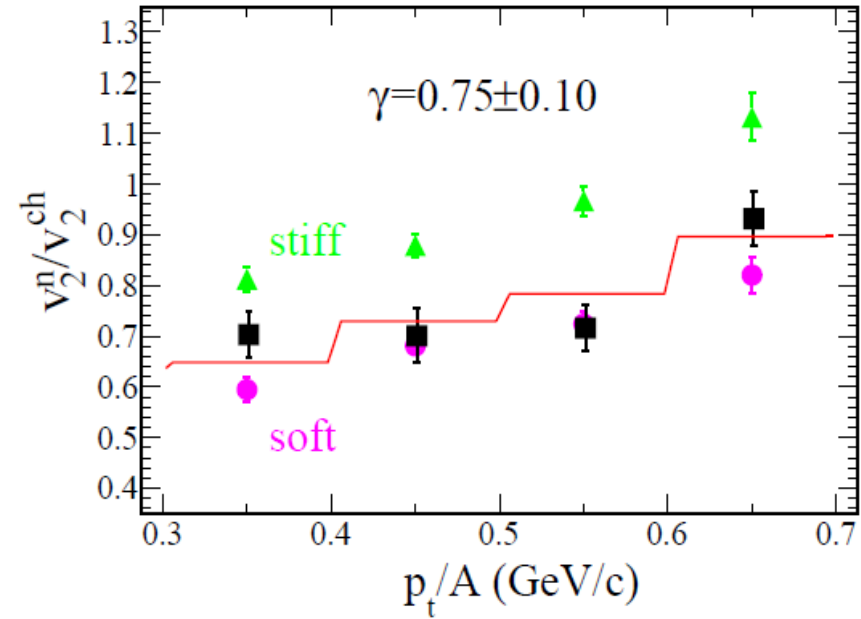
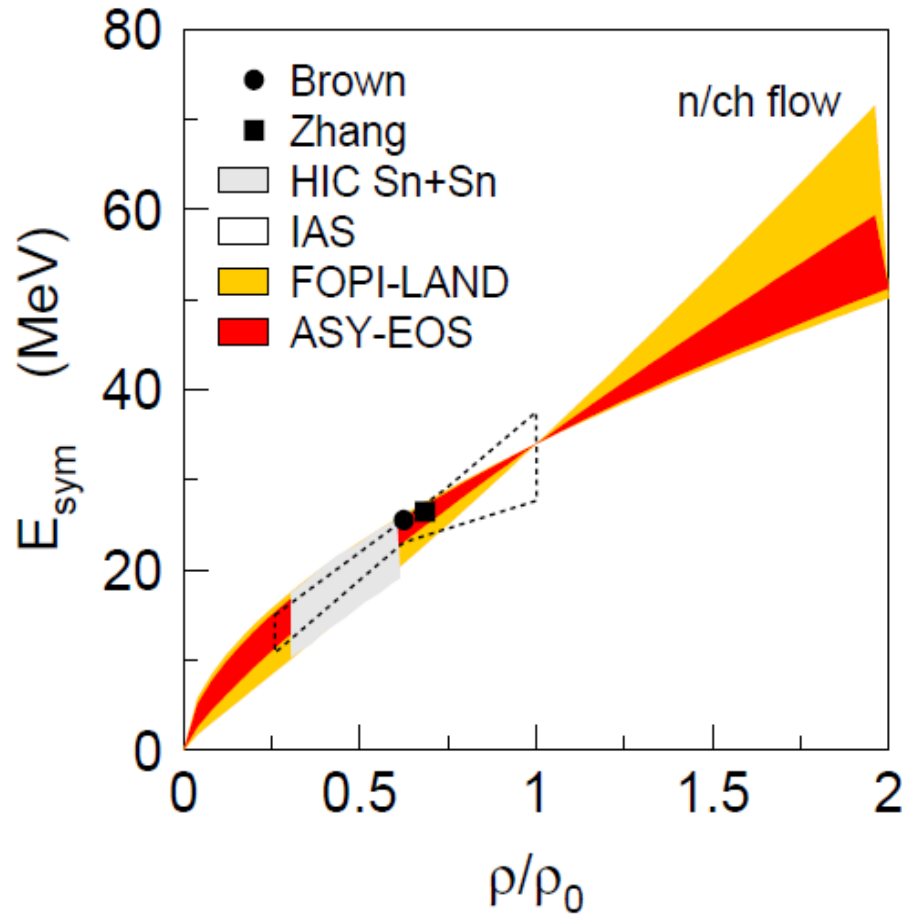


FIG. 15: (Color online) Elliptic flow ratio of neutrons and charged particles for moderately central ($b < 7.5$ fm) collisions of $^{197}\text{Au}+^{197}\text{Au}$ at 400 MeV/nucleon as a function of the transverse momentum per nucleon p_t/A , evaluated with a fraction of 80% for the second step of timing corrections (see Sec. IV A). The full squares represent the experimental data, the triangles and dots represent the UrQMD predictions for stiff ($\gamma = 1.5$) and soft ($\gamma = 0.5$) power-law exponents of the potential term. The full line is the result of a linear interpolation between the predictions, leading to the indicated $\gamma = 0.75 \pm 0.10$.

LAND@ASY-EOS

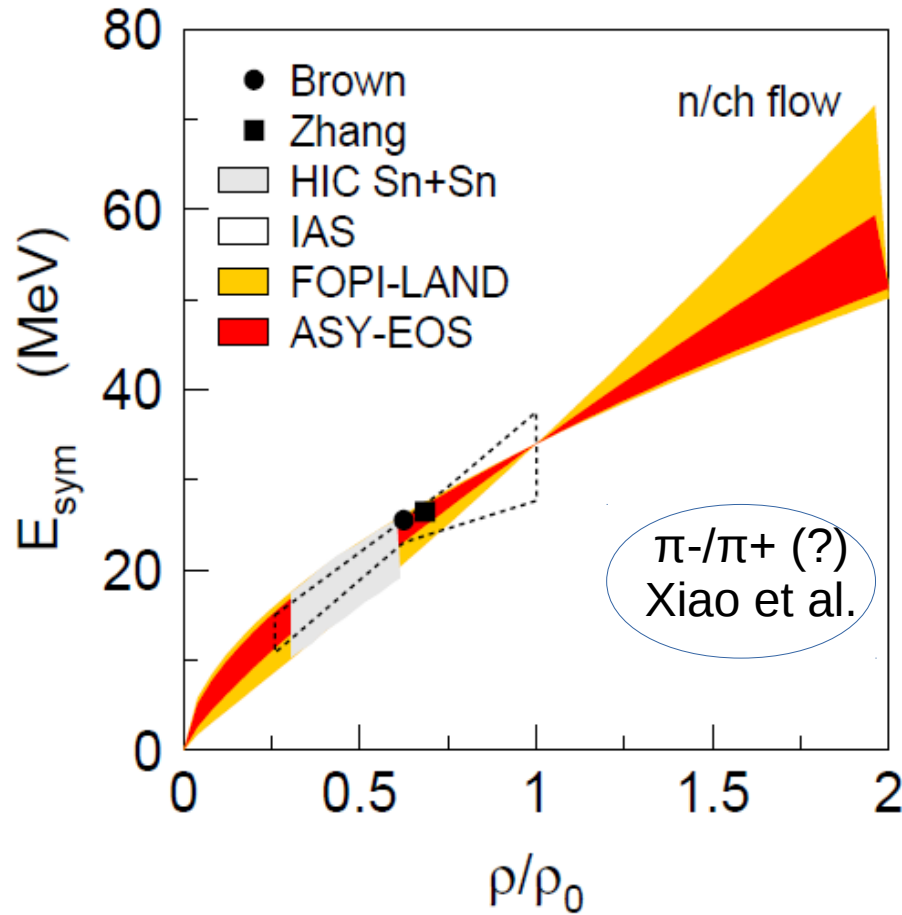


$$\gamma = 0.72 \pm 0.19$$

$$L = 72 \pm 13 \text{ MeV}$$

FIG. 19: (Color online) Constraints deduced for the density dependence of the symmetry energy from the present data in comparison with the FOPI-LAND result of Ref. [5] as a function of the reduced density ρ/ρ_0 . For reference, the low-density results of Refs. [66–69] as presented in Ref. [70] are included.

LAND@ASY-EOS



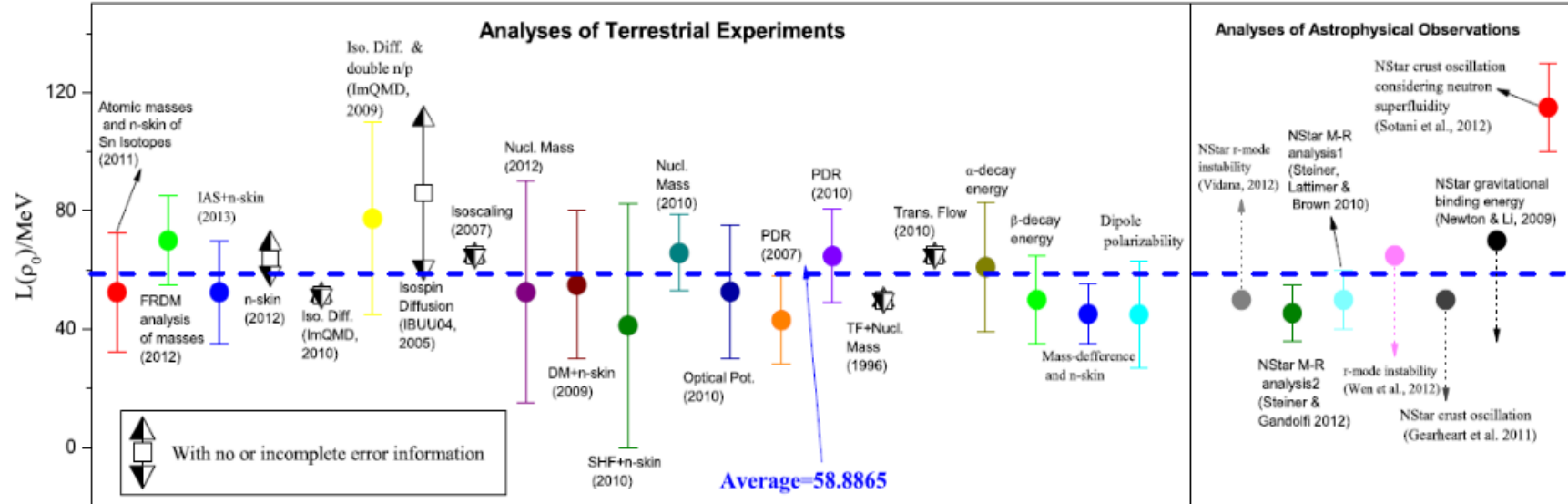
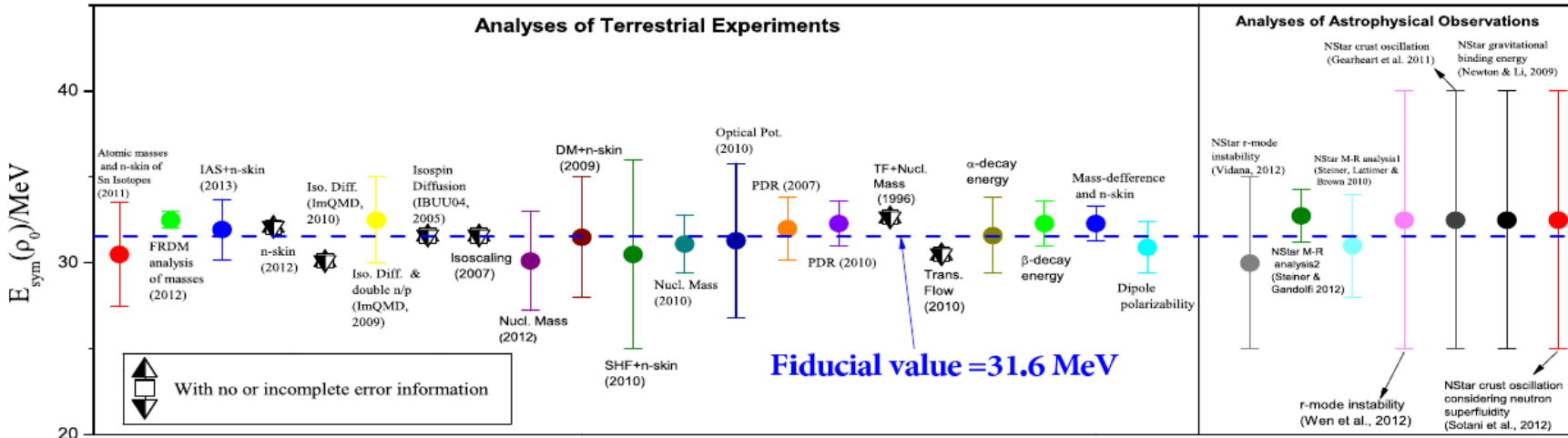
$$\gamma = 0.72 \pm 0.19$$

$$L = 72 \pm 13 \text{ MeV}$$

FIG. 19: (Color online) Constraints deduced for the density dependence of the symmetry energy from the present data in comparison with the FOPI-LAND result of Ref. [5] as a function of the reduced density ρ/ρ_0 . For reference, the low-density results of Refs. [66–69] as presented in Ref. [70] are included.

A well known figure

B.-A. Li, Nucl. Phys. News, 27 (2017) 7



Symmetry energy parameters at ρ_0 from 28 analyses of terrestrial nuclear laboratory experiments and astrophysical observations.

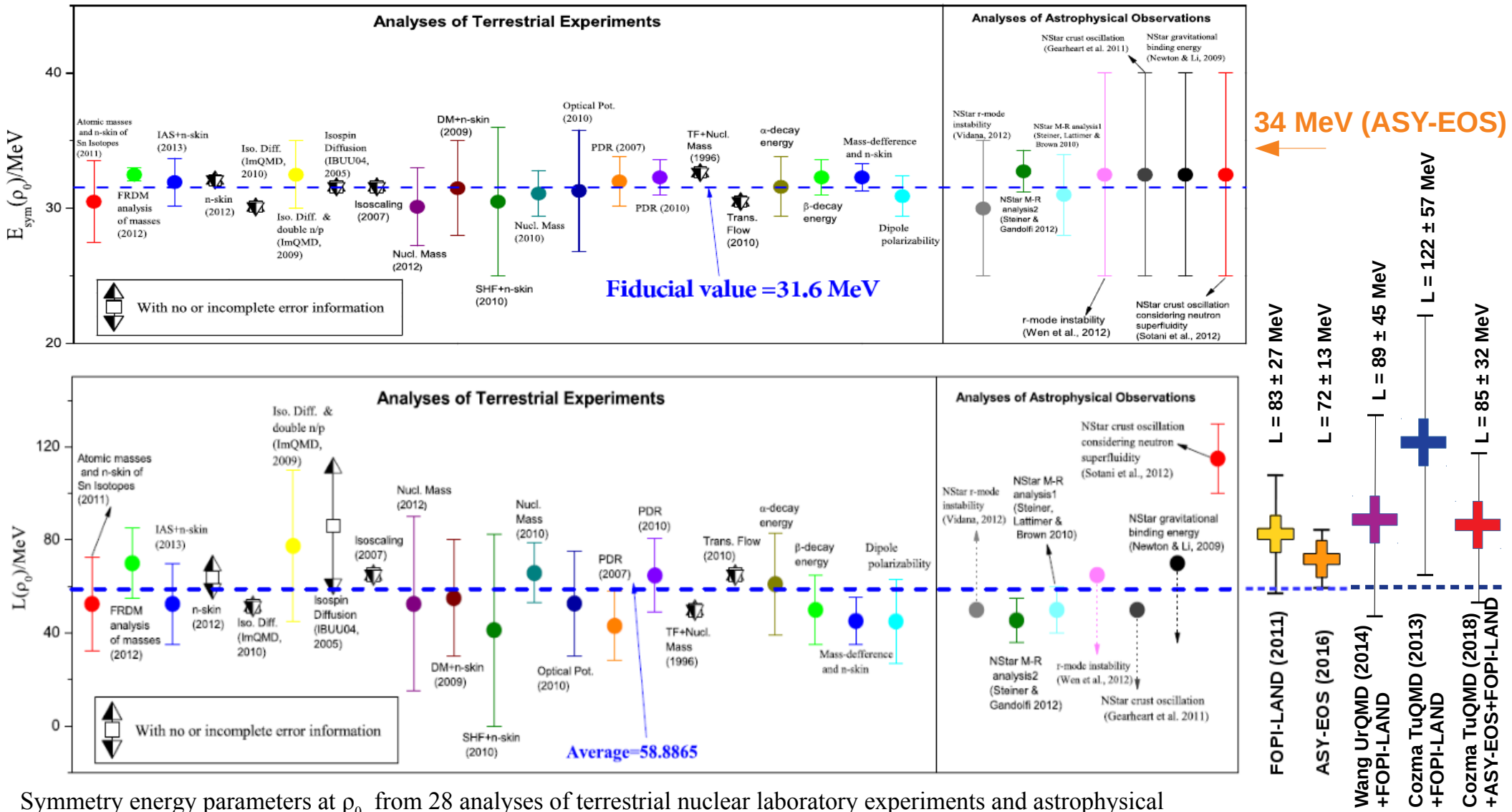
B.-A. Li, Nucl. Phys. News, 27 (2017) 7

P. Russotto et al. Phys. Lett. B 697 (2011) 471; Phys. Rev. C 94 (2016) 034608

Y. Wang et al, Phys. Rev. C 89 (2014) 044603

M.D. Cozma et al., Phys. Rev C 88 (2013) 044912

M.D. Cozma, Eur. Phys. J. A 54 (2018) 40



Symmetry energy parameters at ρ_0 from 28 analyses of terrestrial nuclear laboratory experiments and astrophysical observations. Additional five symbols on the right side represent the results of the FOPI-LAND (yellow) and of the ASY-EOS (orange) high energy experiments, the violet cross shows the Skyrme-UrQMD result for the FOPI-LAND data and the blue and red crosses show the TuQMD result for the FOPI-LAND and the FOPI-LAND and ASY-EOS data, respectively.

B.-A. Li, Nucl. Phys. News, 27 (2017) 7

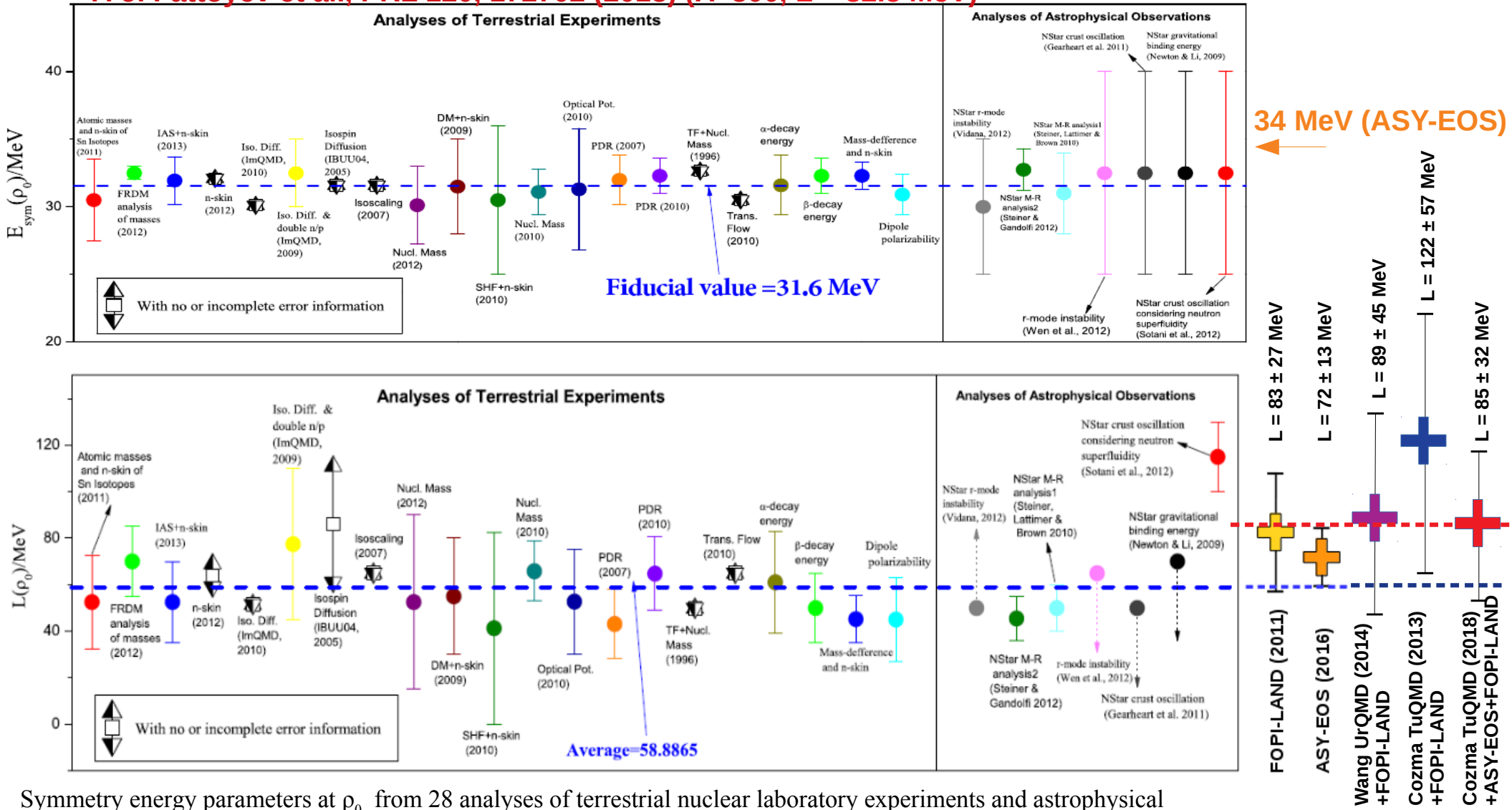
P. Russotto et al. Phys. Lett. B 697 (2011) 471; Phys. Rev. C 94 (2016) 034608

Y. Wang et al, Phys. Rev. C 89 (2014) 044603

M.D. Cozma et al., Phys. Rev C 88 (2013) 044912

M.D. Cozma, Eur. Phys. J. A 54 (2018) 40

F. J. Fattoyev et al., PRL 120, 172702 (2018) ($\Lambda < 800$, $L < 82.5$ MeV) -----



Symmetry energy parameters at ρ_0 from 28 analyses of terrestrial nuclear laboratory experiments and astrophysical observations. Additional five symbols on the right side represent the results of the FOPI-LAND (yellow) and of the ASY-EOS (orange) high energy experiments, the violet cross shows the Skyrme-UrQMD result for the FOPI-LAND data and the blue and red crosses show the TuQMD result for the FOPI-LAND and the FOPI-LAND and ASY-EOS data, respectively.

QMC predictions

S. Gandolfi et al., Phys. Rev. C 85, 032801(R) (2012)

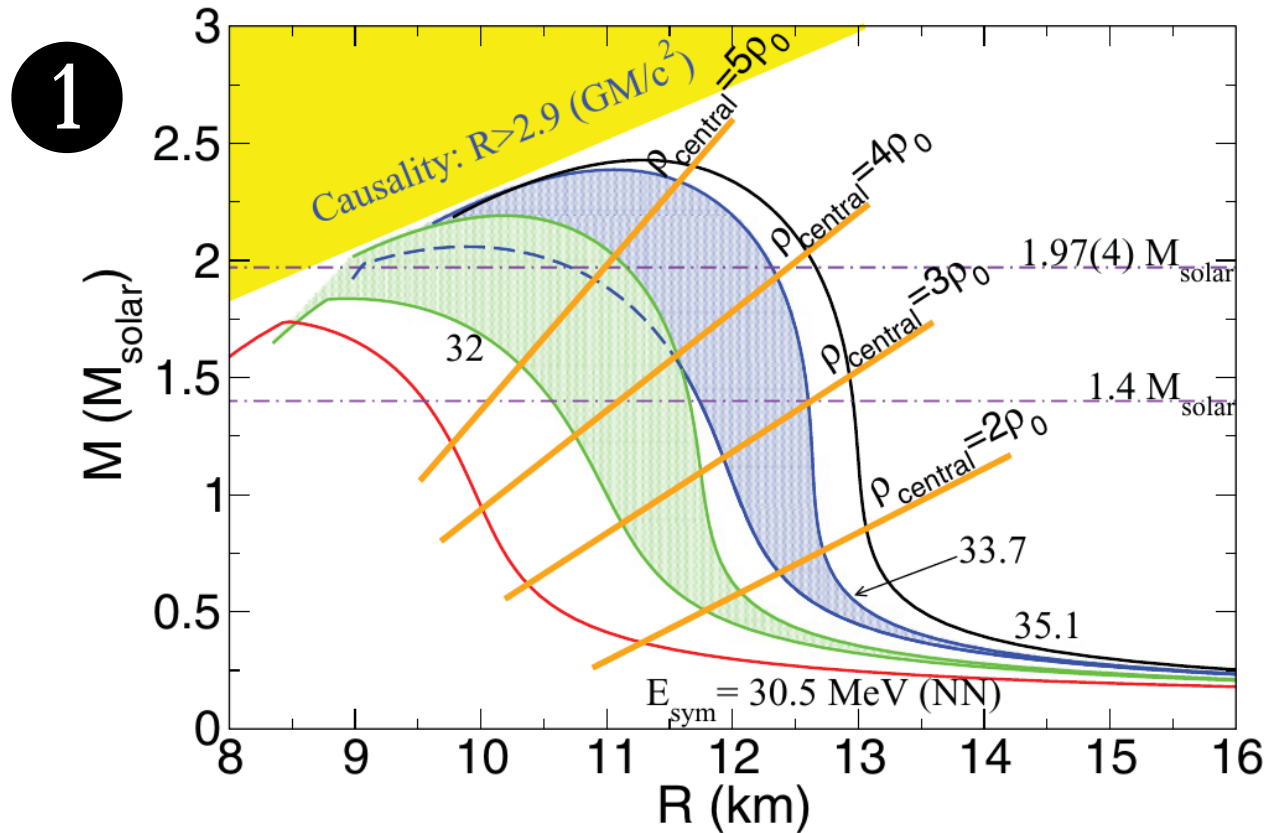


FIG. 2. (Color online) Mass-radius relation for the EoS with three-neutron interactions corresponding to the bands for different E_{sym} shown in Fig. 1. The intersections with the orange lines roughly indicate central densities realized in these stars.

2

“The uncertainty in the measured symmetry energy of ± 2 MeV leads to an uncertainty of about 3 km for the radius.”

1 → measure the symmetry energy at high densities (energies) to probe the EOS in the core of NS

2 → measure the symmetry energy as precise as possible to provide constraints relevant for astrophysics.

densities probed in HIC

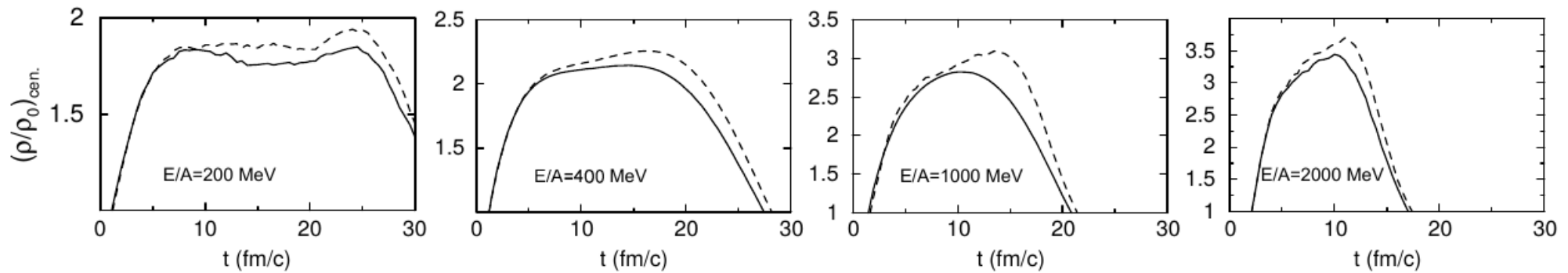


Figure 3: Evolution of the central baryon density in $^{132}\text{Sn}+^{124}\text{Sn}$ collisions at beam energies from 200 to 2000 MeV/nucleon for $b = 1$ fm, as predicted by the hadronic transport model of [31].

adapted from B.-A. Li, Nucl. Phys. A 708 (2002) 365.

UrQMD predictions

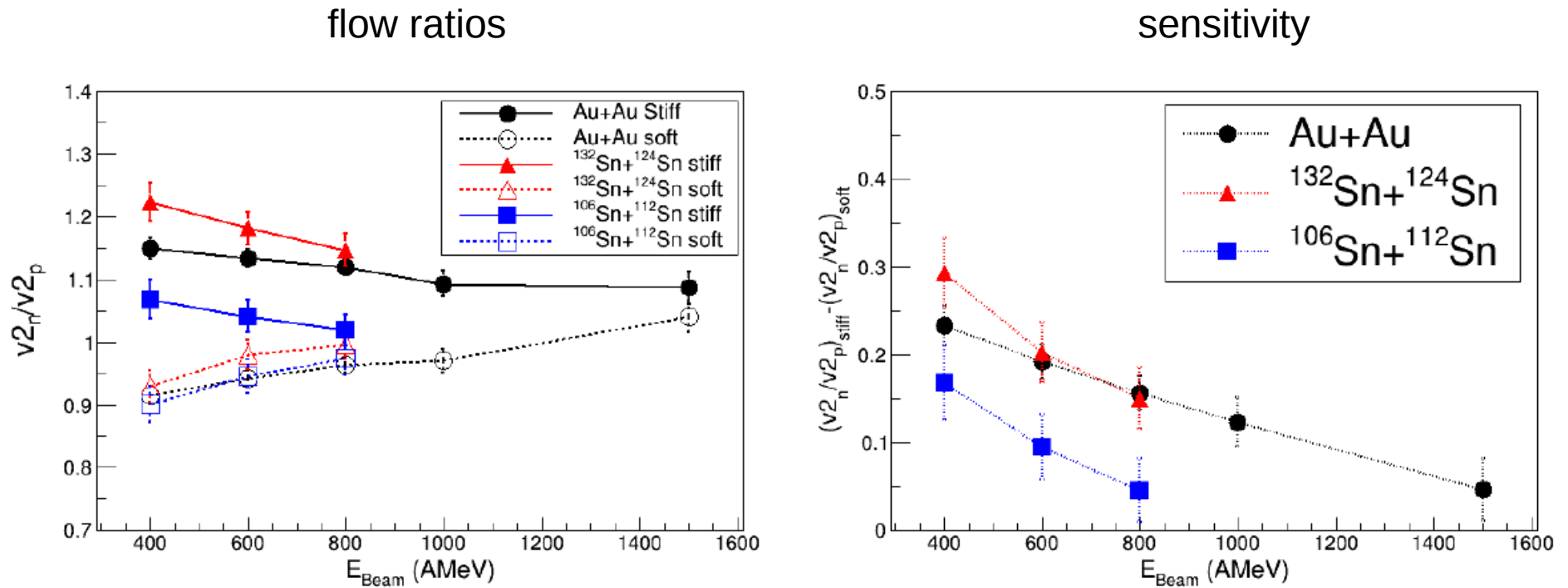


Figure 4: Left panel: Excitation functions of neutron-to-proton elliptic flow ratios, v_{2n}/v_{2p} , at mid-rapidity for semi-central Au+Au, $^{132}\text{Sn}+^{124}\text{Sn}$, and $^{106}\text{Sn}+^{112}\text{Sn}$ collisions, as predicted by the UrQMD model for stiff and soft $E_{\text{sym}}(\rho)$. Right panel: differences between the stiff and soft results.

ASY-EOS II @ FAIR (202?)

Determination of the density dependence of the EOS at supra-saturation densities

Symmetric and asymmetric systems
 ^{108}Sn , ^{132}Sn , ^{197}Au @ 0.4, 1, 1.2 AGeV

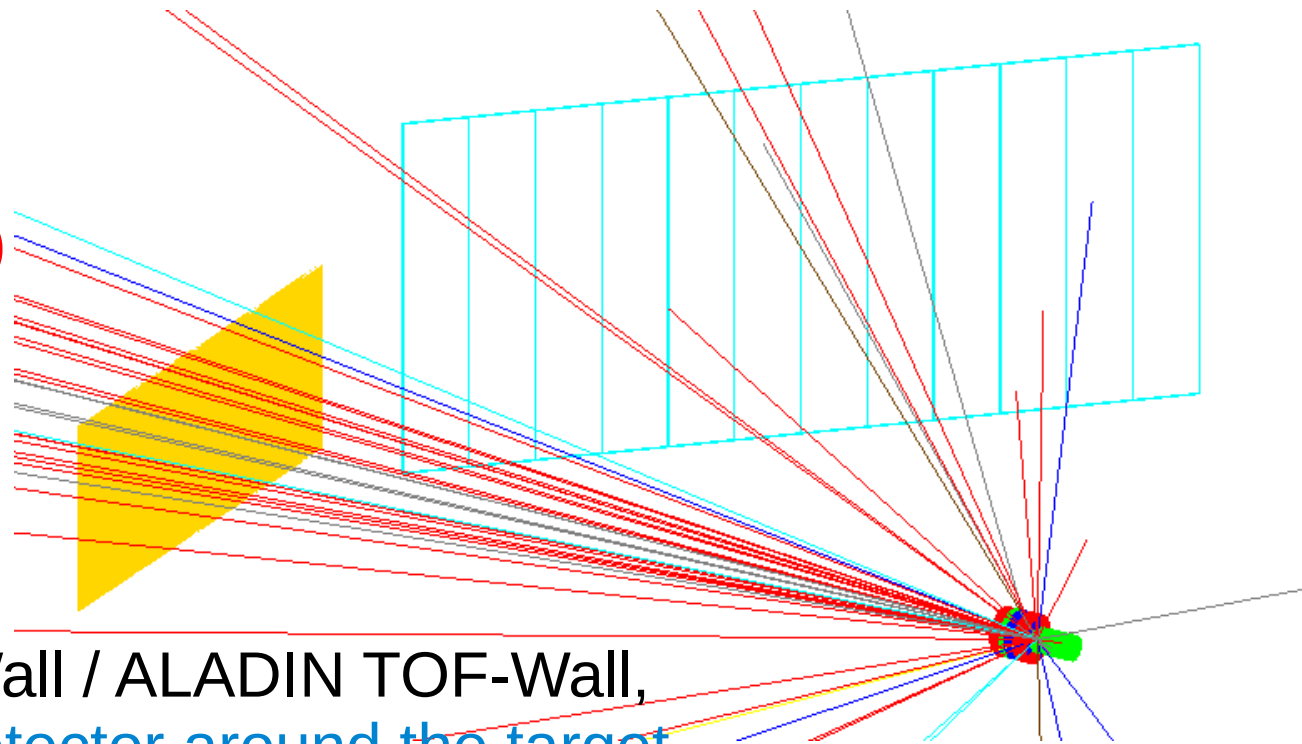
Observables:

ratios: n/p , $t/{}^3\text{He}$, π^-/π^+ (?)

flow: n , p , t , ${}^3\text{He}$

Main detectors:

NeuLAND, FOPI PlasticWall / ALADIN TOF-Wall,
Trigger/Reaction Plane detector around the target



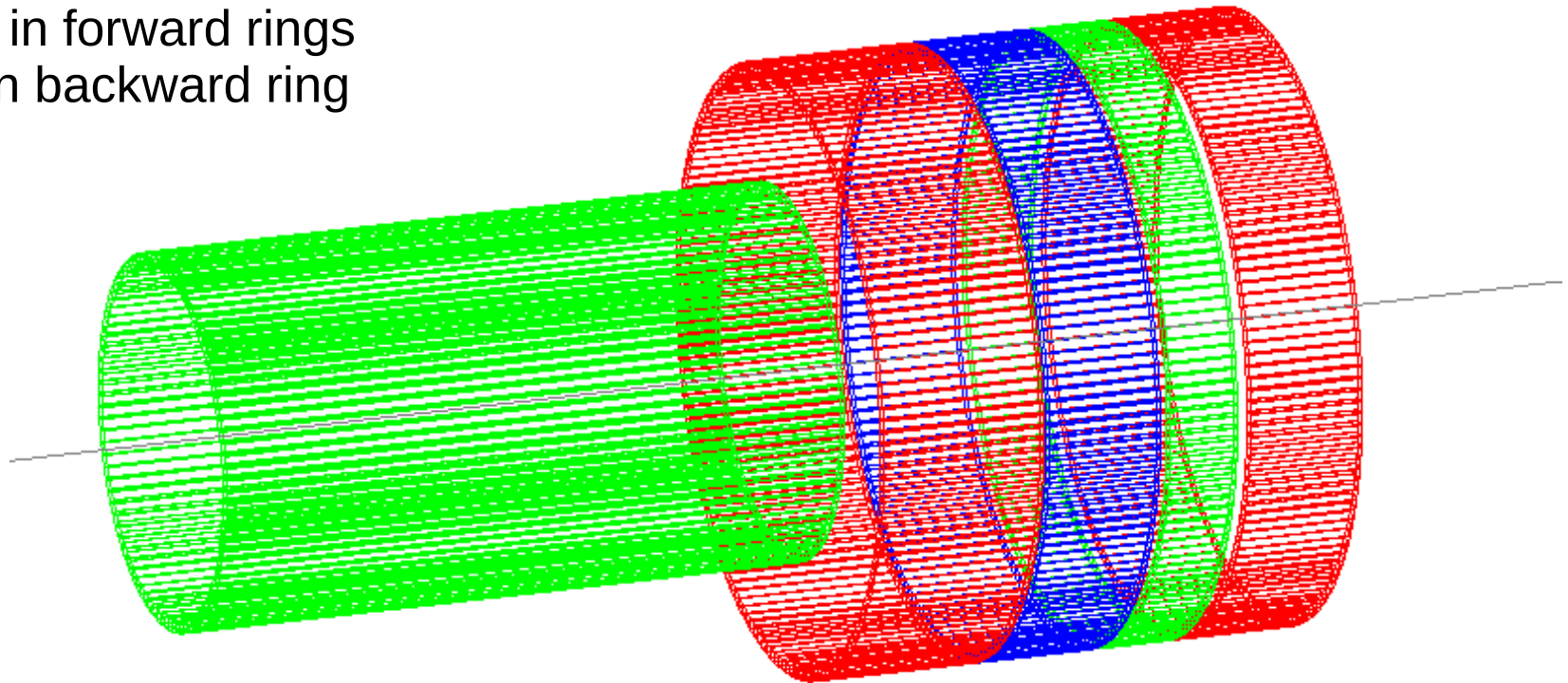
Trigger/Reaction Plane detector around the target

requirements:

- should cover angles $> 30^\circ$,
- high segmentation in azimuthal angle,
- high geometrical efficiency,
- low multihit probability,
- fast timing

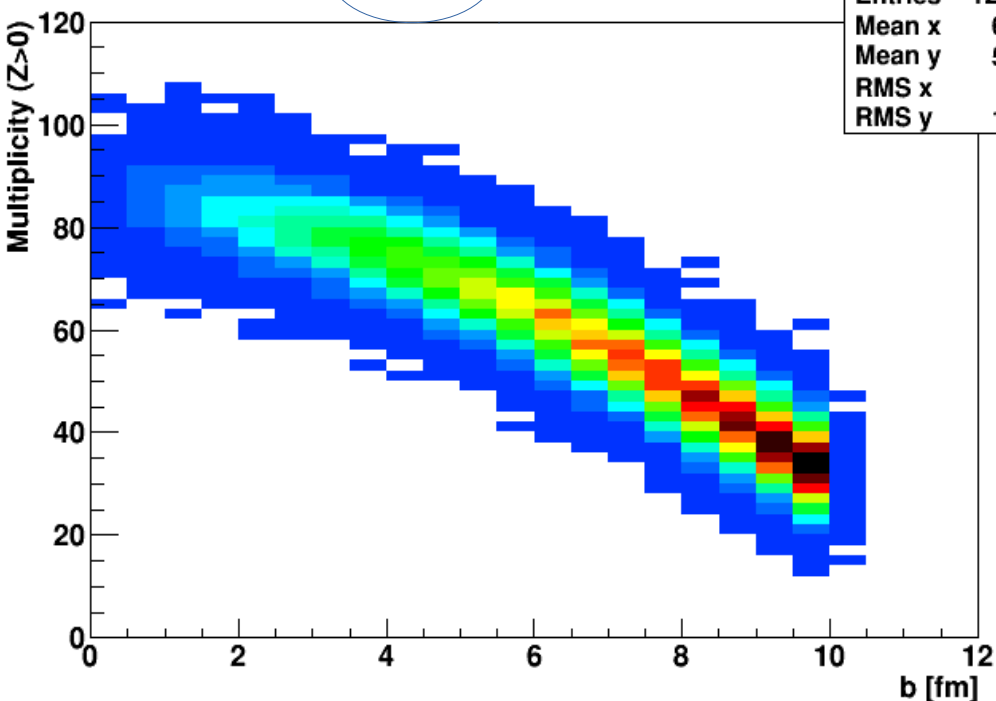
Trigger/Reaction Plane detector around the target:

- 5 rings of 4x4 mm² fast scintillating fibers (e.g. BCF-20) read out by SiPMs
- covers angles from 30° to 165°,
- segmentation assures more or less uniform count rates for Au+Au at 1 AGeV,
- geometrical efficiency ~87%
- ~10% of charged particles involved in multihits,
- ~5% multihit probability
- sufficiently large for radioactive beams
- sufficiently small and lightweight not to disturb neutrons
- min radius - 6 cm,
- max radius - 12 cm
- length 43 cm
- 160 segments in forward rings
- 96 segments in backward ring
- 736 channels



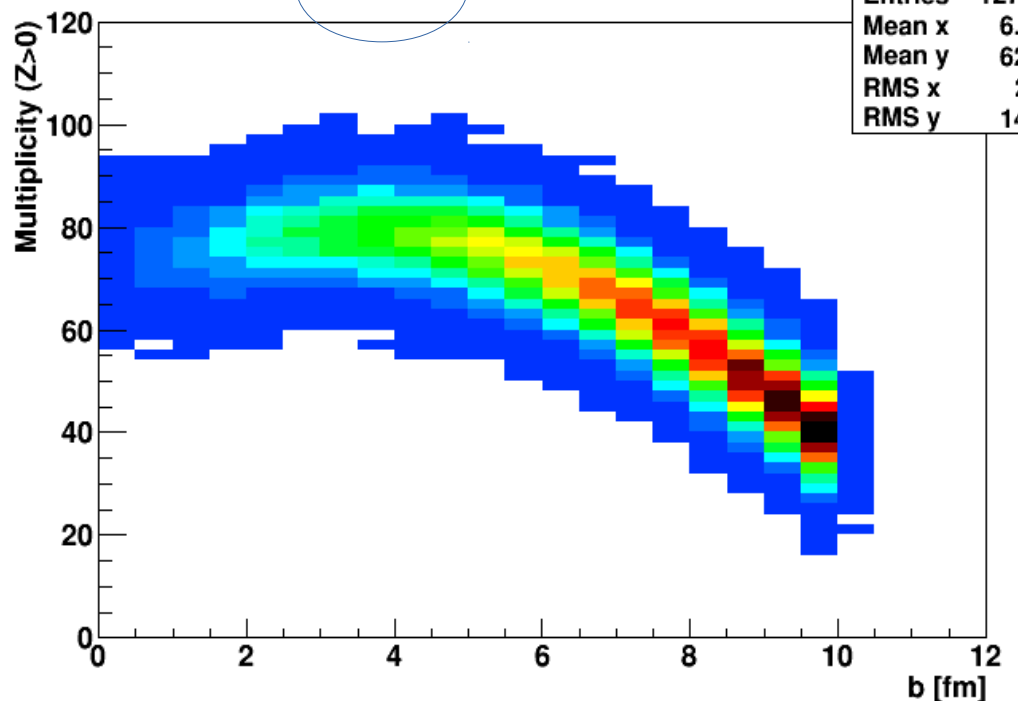
UrQMD + clustering: Au+Au 1000 AMeV, 0-10 fm, 200 fm/c

Sum\$(iz>0&&ia>0&&t>30):b



hbm30	
Entries	12799
Mean x	6.65
Mean y	56.1
RMS x	2.3
RMS y	16.4

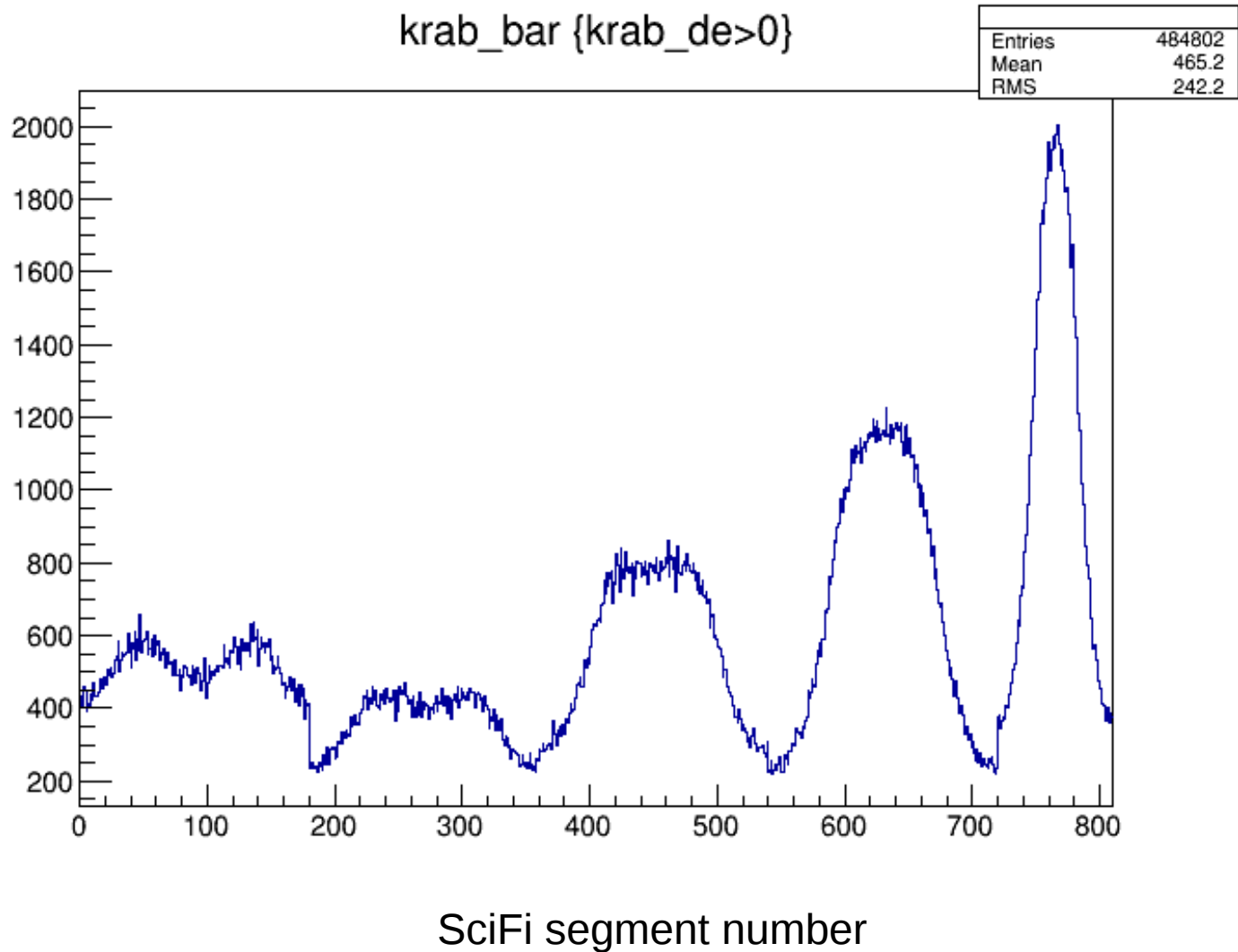
Sum\$(iz>0&&ia>0&&t<30):b



hbm30	
Entries	12799
Mean x	6.65
Mean y	62.8
RMS x	2.3
RMS y	14.3

better correlation

hits/segment



Summary & Conclusions

- ASY-EOS measurements yield a moderately soft to linear density dependence of the symmetry energy:
 $\gamma = 0.72 \pm 0.19$ ($L = 72 \pm 13$ MeV).
- First analyses using the LIGO and VIRGO constraints yield
- $\sim 10 < R_{\text{NS}}(1.4 M_{\odot}) < \sim 13$ km $\rightarrow 26 < L < 84$ MeV, which support the results of Lattimer and Steiner analysis of qLMXB ($10.9 < R_{\text{NS}}(2.0 M_{\odot}) < 12.7$ km) $\rightarrow 41 < L < 83$ MeV, and is compatible with the ASY-EOS measurement.
- The FOPI-LAND and ASY-EOS high energy results for L locate above the average for the low density measurements what seems to indicate a transition towards more stiff symmetry energy above the ρ_0 .
- A precise measurement needed at ~ 1 AGeV to explore densities at 2-3 ρ_0 .
- KRAB detector under construction.

reviews:

Ph. Chomaz, F. Gulminelli, W. Trautmann, S.J. Yennello,
Eur. Phys. J. A 30 (2006).

M. B. Tsang et al.
Phys. Rev. Lett. 102, 122701 (2009)

M.B. Tsang et al.,
Phys. Rev. C 86, 015803 (2012).

W. Trautmann, H.H. Wolter,
Int. J. Mod. Phys. E21, 1230003 (2012).

B.-A. Li, À. Ramos, G. Verde, I. Vidaña,
Eur. Phys. J. A 50 (2014).

C J Horowitz et al.
J. Phys. G: Nucl. Part. Phys. 41 (2014) 093001

M. Baldo, G.F. Burgio, The nuclear symmetry energy,
Prog. Part. Nucl. Phys., 91 (2016) 203

Oertel, M., Hempel, M., Klähn, T., & Typel, S.
Rev. Mod. Phys., 89, 015007 (2017)

B.-A. Li, Nuclear Symmetry Energy Extracted from Laboratory Experiments,
Nucl. Phys. News, 27 (2017) 7

Nai-Bo Zhang, Bao-An Li, Jun Xu
Combined Constraints on the Equation of State of Dense Neutron-Rich Matter from Terrestrial Experiments and Observations of Neutron Stars,
arXiv:1801.06855v1

J. Piekarewicz Nuclear Astrophysics in the New Era of Multimessenger Astronomy
arXiv:1805.04780v1

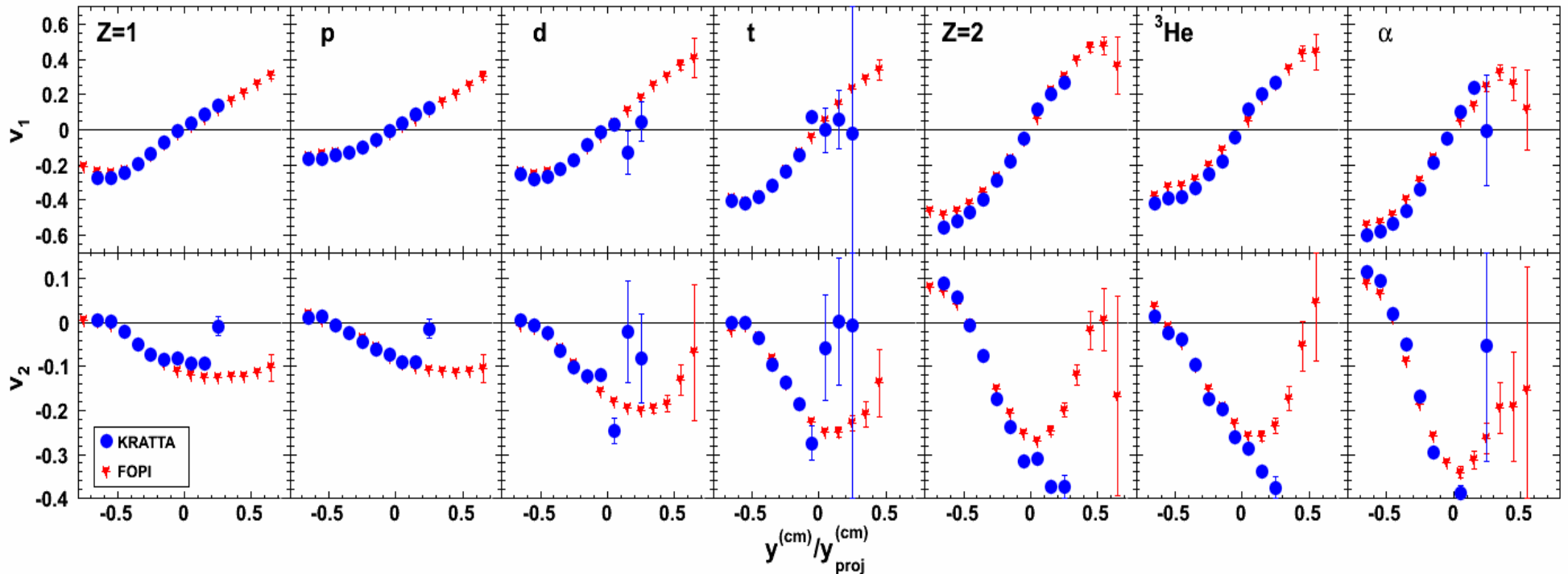
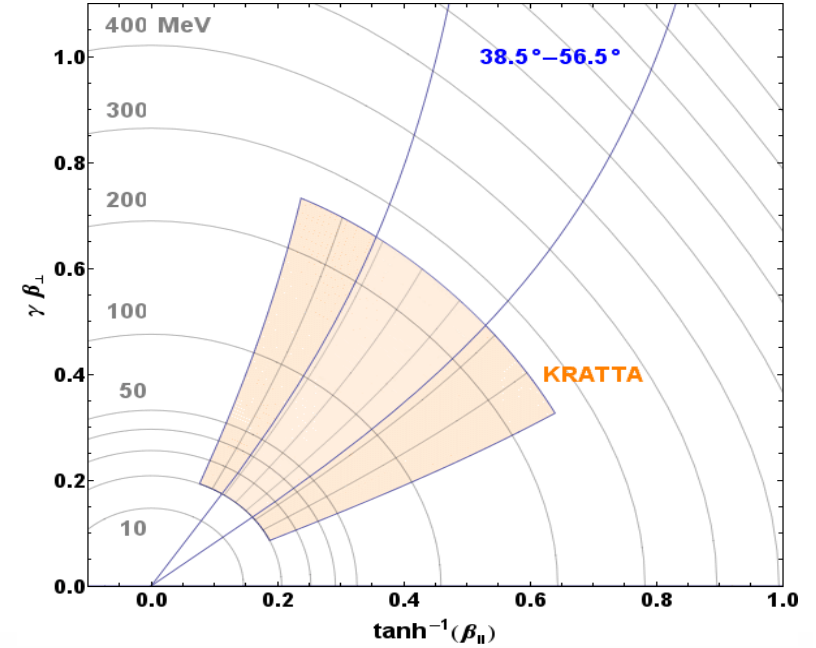
Flows of light charged particles in Au(400 MeV/u) + Au reactions: KRATTA vs FOPI results

Fourier decomposition of the azimuthal distributions with respect to the reaction plane (ϕ_R):

$$\frac{dN}{d(\phi - \phi_R)} = \frac{N_0}{2\pi} \left(1 + 2 \sum_{n \geq 1} v_n \cos n(\phi - \phi_R) \right)$$

$$v_1 \equiv \langle \cos(\phi - \phi_R) \rangle \quad \text{directed flow}$$

$$v_2 \equiv \langle \cos 2(\phi - \phi_R) \rangle \quad \text{elliptic flow}$$

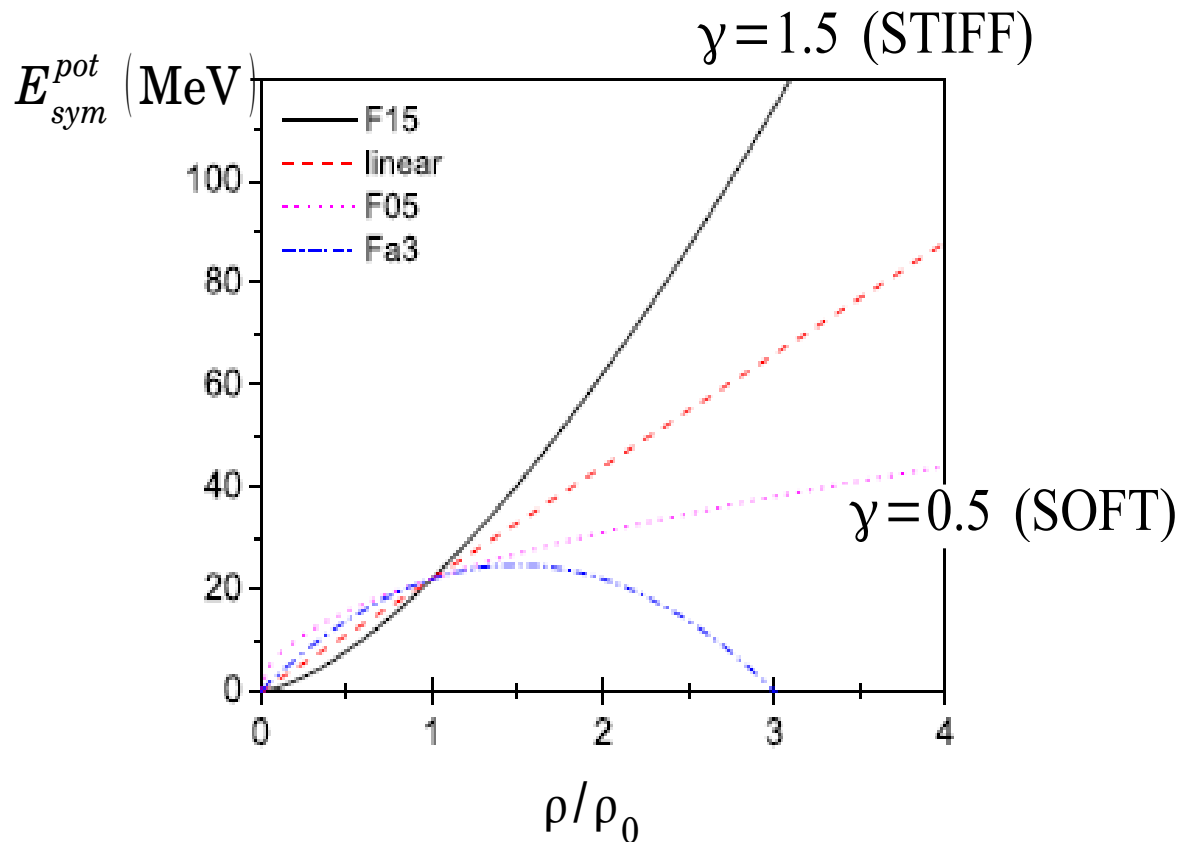


Model simulations

UrQMD Q. Li, J.Phys. G 31(2005)1359

„Fermi-gas” parametrization of the symmetry term:

$$E_{sym} = E_{sym}^{pot} + E_{sym}^{kin} = 22 \text{ MeV} \left(\frac{\rho}{\rho_0} \right)^\gamma + 12 \text{ MeV} \left(\frac{\rho}{\rho_0} \right)^{2/3}$$



Stopping time = 150 fm/c

Nucleons $\rightarrow \{ \vec{r}_i, \vec{p}_i \}$



Clustering procedure

($\Delta r = 2.5$ fm, $\Delta p = 290$ MeV/c)

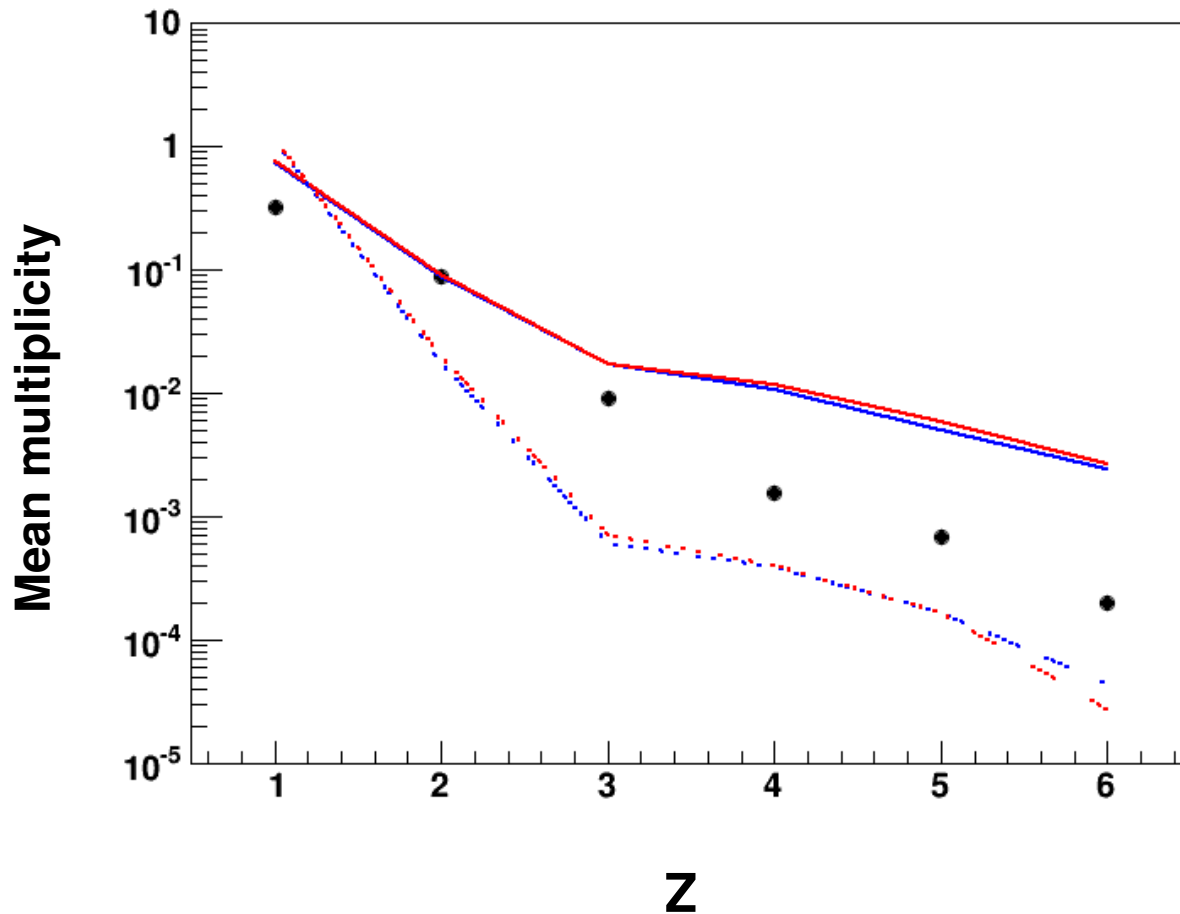
($\Delta r = 3$ fm, $\Delta p = 100$ MeV/c)

⋮

Charge distribution Au(400 MeV/u) + Au

KRATTA data ↔ UrQMD predictions

— ASY - SOFT
— ASY - STIFF



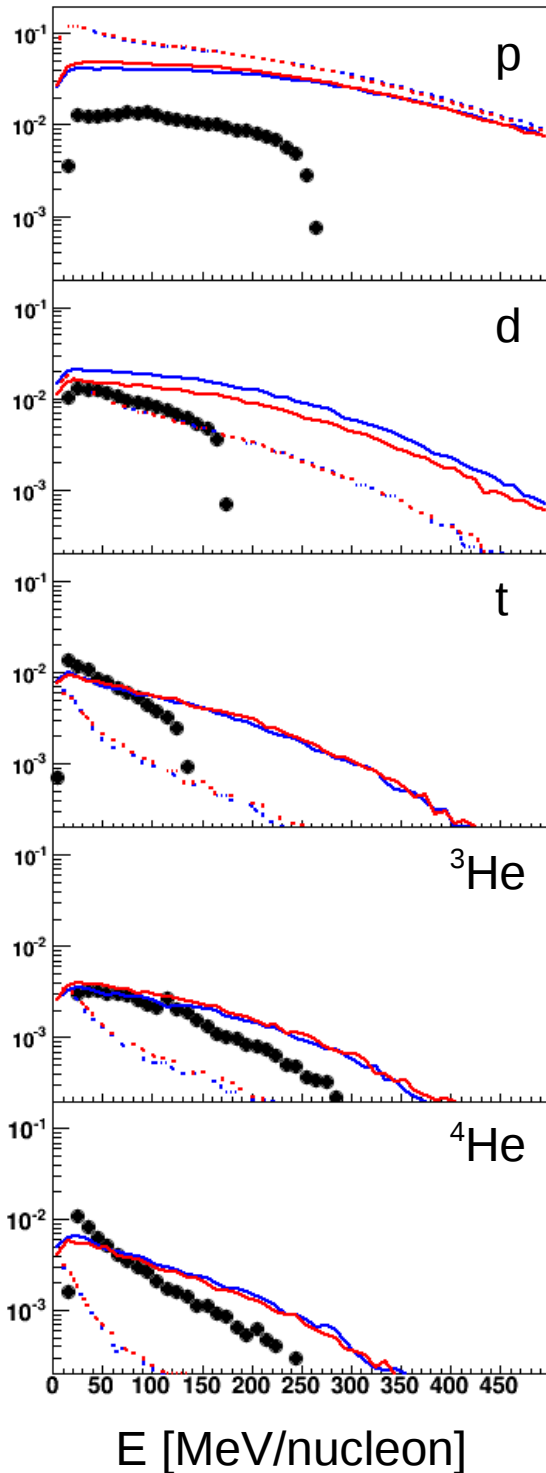
$b < 7.5$ fm
 $24^\circ < \text{lon} < 59.4^\circ$
 $0.7^\circ < \text{lat} < 25.7^\circ$
 $20 < E_{\text{KIN}}/A < 133$ MeV

— $\Delta r=2.5\text{fm } \Delta p=290\text{MeV}/c \ \gamma=0.5$
— $\Delta r=2.5\text{fm } \Delta p=290\text{MeV}/c \ \gamma=1.5$
..... $\Delta r=3.0\text{fm } \Delta p=100\text{MeV}/c \ \gamma=0.5$
..... $\Delta r=3.0\text{fm } \Delta p=100\text{MeV}/c \ \gamma=1.5$
.....●..... Exp.

Energy/nucleon

Au(400 MeV/u) + Au

$d^2\sigma/d\Omega dE$ [arb. units]



- $\Delta r=2.5\text{fm}$ $\Delta p=290\text{MeV}/c$ $\gamma=0.5$
- $\Delta r=2.5\text{fm}$ $\Delta p=290\text{MeV}/c$ $\gamma=1.5$
- ⋯ $\Delta r=3.0\text{fm}$ $\Delta p=100\text{MeV}/c$ $\gamma=0.5$
- ⋯ $\Delta r=3.0\text{fm}$ $\Delta p=100\text{MeV}/c$ $\gamma=1.5$
- ⋯●⋯ Exp.

$b < 7.5$ fm

$24^\circ < \text{lon} < 59.4^\circ$

$0.7^\circ < \text{lat} < 25.7^\circ$

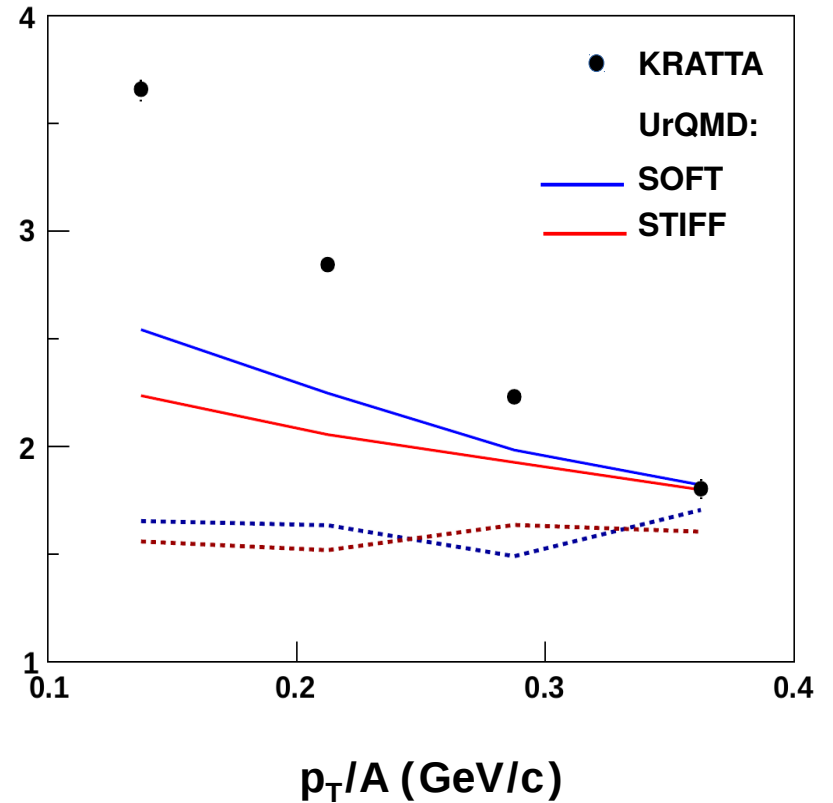
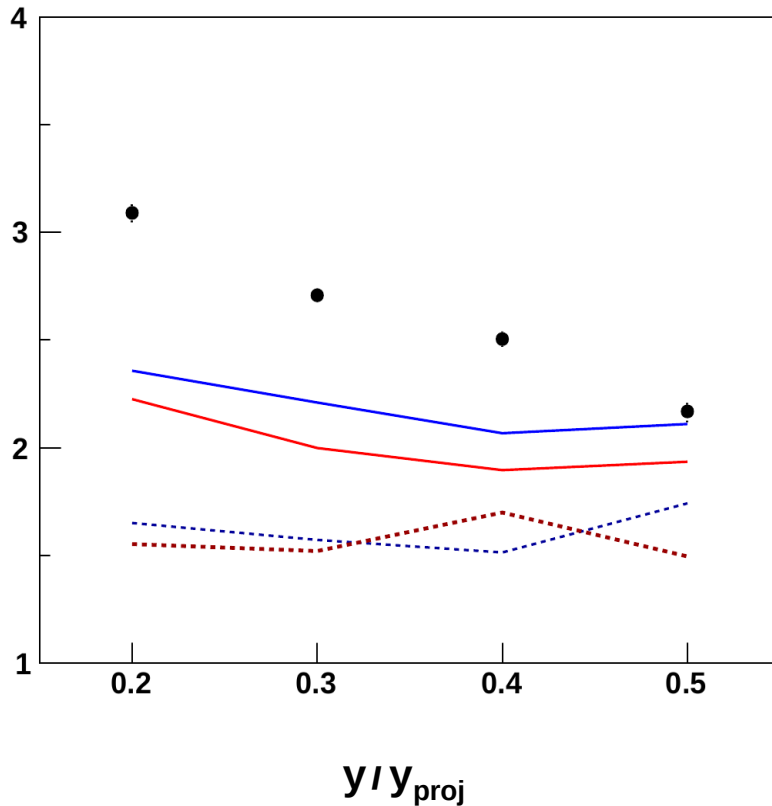
$t/{}^3\text{He}$ isotope ratios ($20 < E_{\text{kin}}/A < 133 \text{ MeV}$)

Au (400 MeV/u) + Au

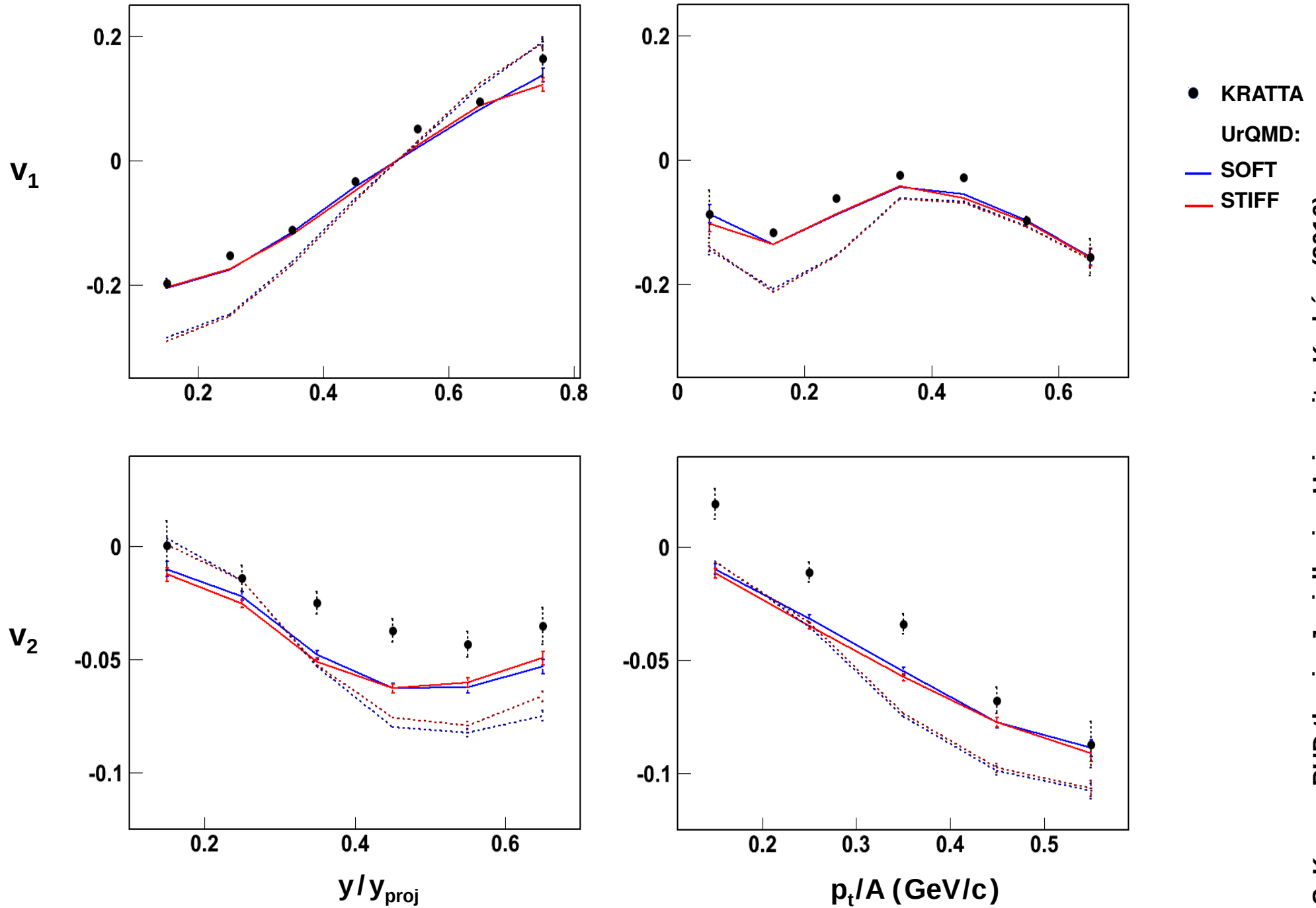
$5.5 < b < 7.5 \text{ fm}$

$24^\circ < \Theta_{\text{LAB}} < 62^\circ$

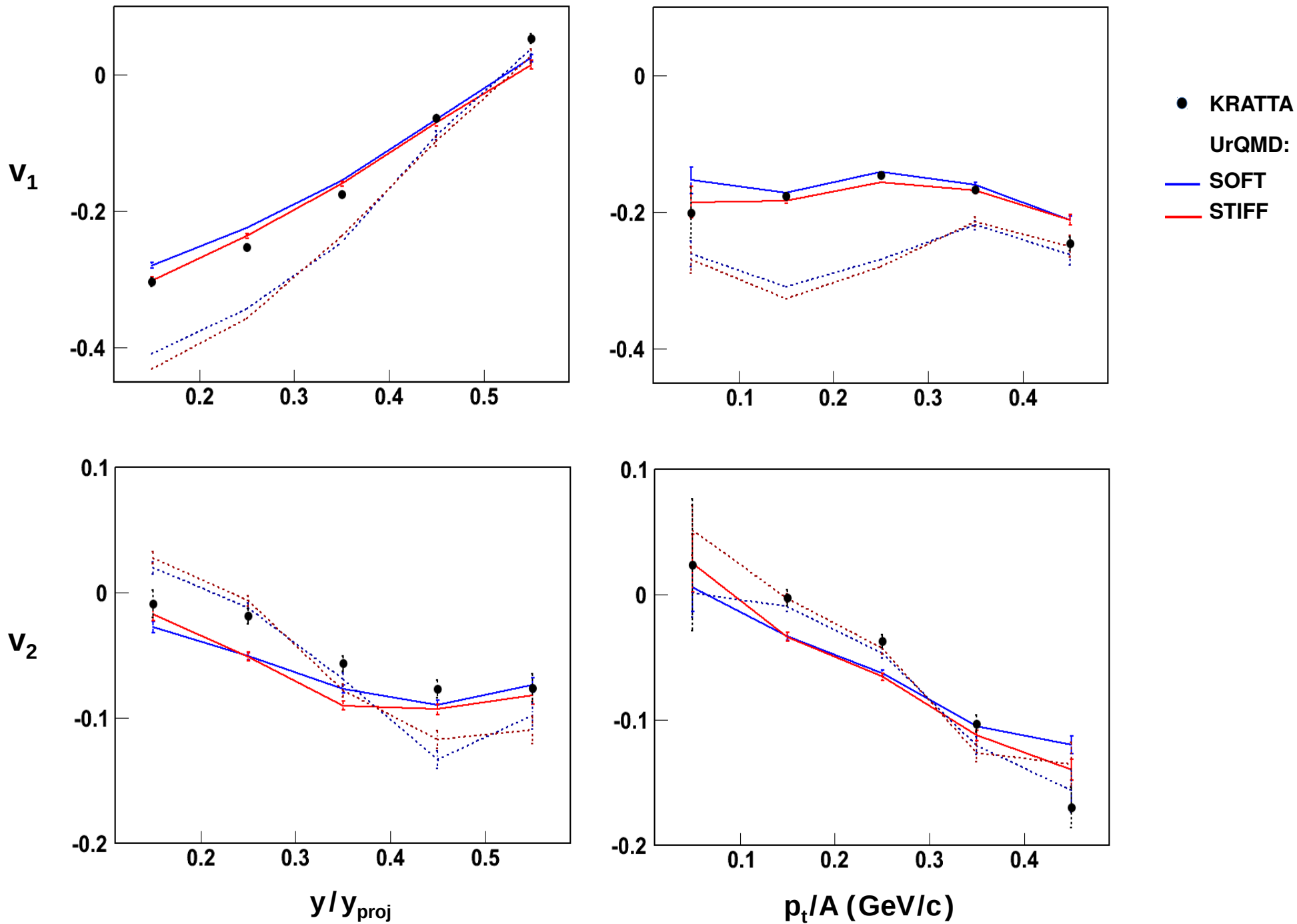
$20 < E_{\text{KIN}}/A < 133 \text{ MeV}$



Proton flow ($20 < E_{\text{kin}} < 250$ MeV)



Deuteron flow ($20 < E_{\text{kin}}/A < 160$ MeV)



current status

