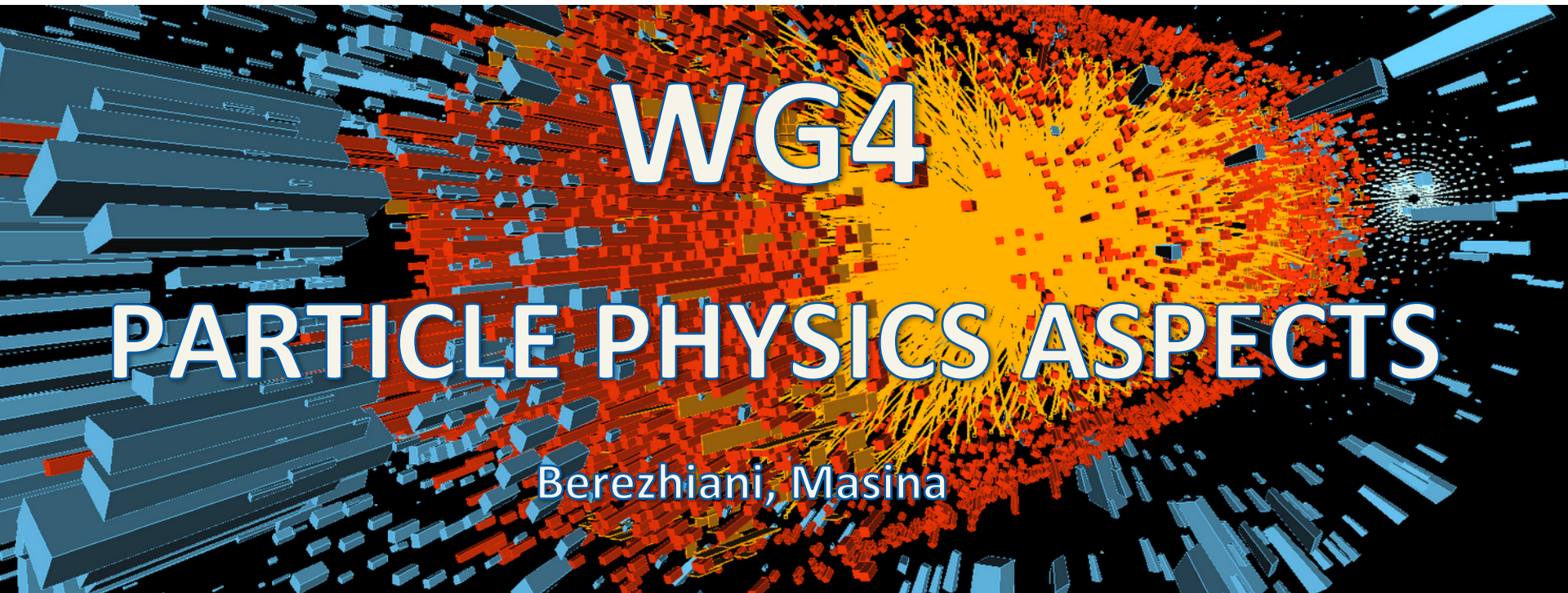


PRIN 2012 Project
THEORETICAL ASTROPARTICLE PHYSICS



PRIN 2012 MidTerm Workshop, Torino, 9-10/07/14

PRIN 2012 Project
THEORETICAL ASTROPARTICLE PHYSICS

WG4

“...where the intimate connection between cosmological/astrophysical analyses is coordinated with particle physics.”

PRIN 2012 Project
THEORETICAL ASTROPARTICLE PHYSICS

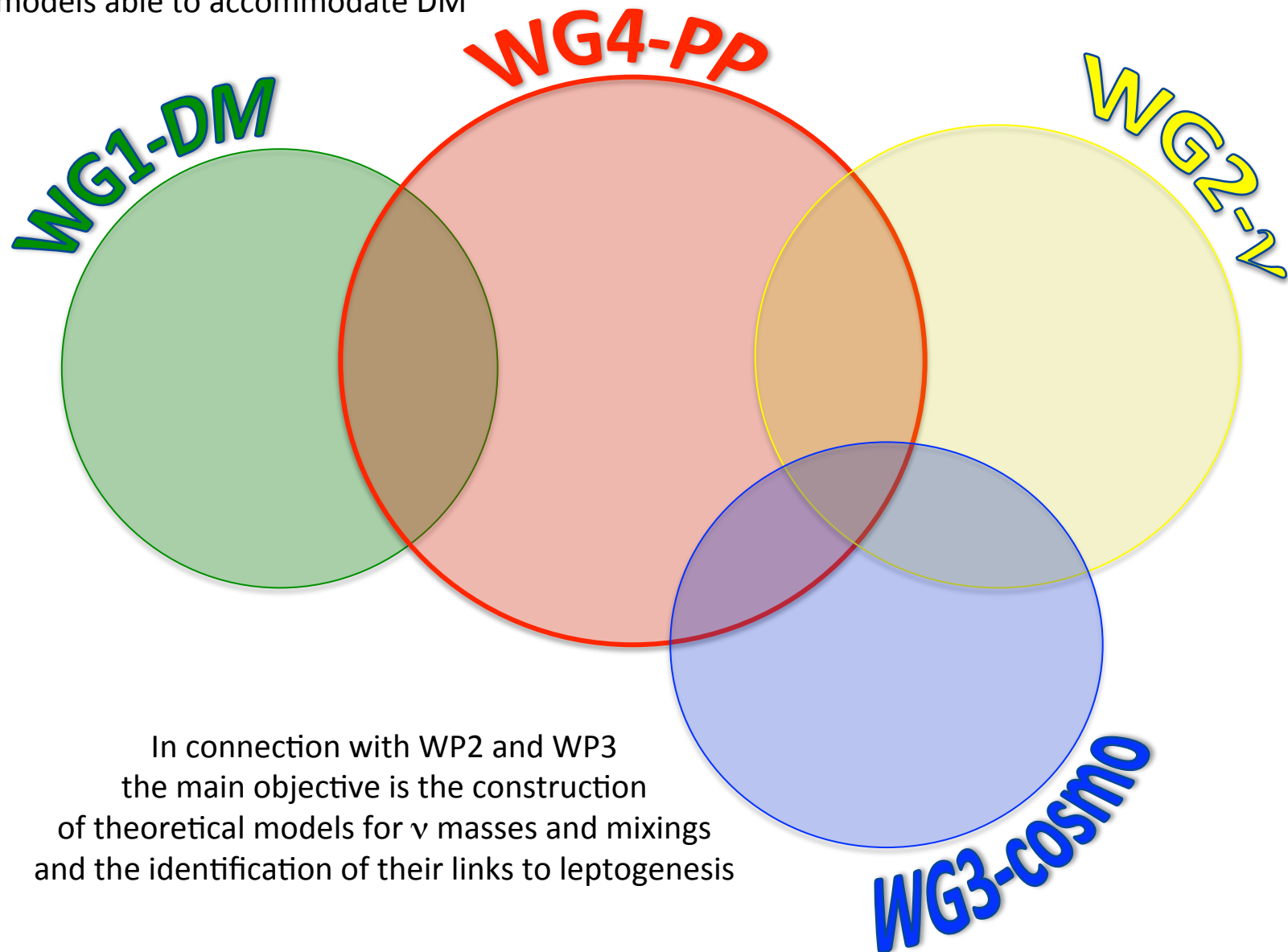
WG4

“...where the intimate connection between cosmological/astrophysical analyses is coordinated with particle physics.”

Goal:

- 1) To identify a path toward the development of particle physics frameworks beyond the SM able to solve the astroparticle physics puzzles studied in WP1-WP3
- 2) To attempt, if possible, to find common solutions.

In connection with WP1
the main objective is to identify
NP models able to accommodate DM



In connection with WP2 and WP3
the main objective is the construction
of theoretical models for ν masses and mixings
and the identification of their links to leptogenesis

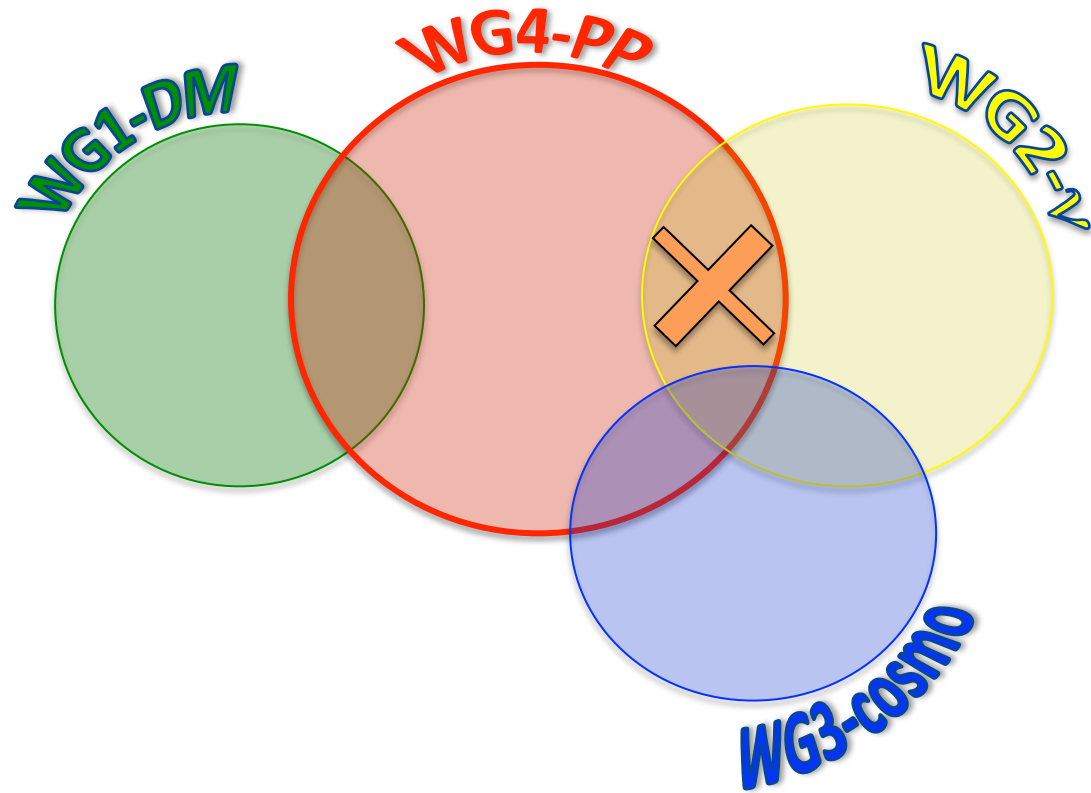
Total n. of papers in area WG4 since PRIN start: about

30

TOPICS:

- | | |
|----------------------------------|--------------------|
| 1) Lepton flavor and CP (8) | SISSA/LNF-RM/NA-SA |
| 2) Leptogenesis (2) | SISSA/LNF-RM |
| 3) Solar physics constraints (2) | FE/AQ |
| 4) Higgs and inflation (1) | FE |
| 5) Susy and LHC (4) | PD/AQ/NA-SA/FE |
| 6) Quantum theories (7) | TO/NA-SA/AQ |
| 7) Gravity (1) | BA |
| 8) Neutrons and Exotica (5) | AQ |

1) WG2/WG4: LEPTON FLAVOR AND CP (SISSA/LNF-RM/NA-SA)



1) WG2/WG4: LEPTON FLAVOR AND CP (SISSA/LNF-RM/NA-SA)

1) S. M. Boucenna, R. M. Fonseca, F. Gonzalez-Canales, J. W. F. Valle
Small neutrino masses and gauge coupling unification
arXiv:1411.0566 [hep-ph], Phys.Rev. D91 (2015) 031702

2) Sofiane M. Boucenna, Jose W. F. Valle, Avelino Vicente
Predicting charged lepton flavor violation from gauge symmetry
arXiv:1502.07546 [hep-ph]

3) Sofiane M. Boucenna, Jose W. F. Valle, Avelino Vicente
Are the B decay anomalies related to neutrino oscillations?
arXiv:1503.07099 [hep-ph]

Connection with
gauge symmetries

4) L.L. Everett, T.Garon and A.J. Stuart
A Bottom-Up Approach to Lepton Flavor and CP Symmetries
arXiv:1501.04336 [hep-ph], JHEP 1504 (2015) 069

Talk by Stuart

Bottom-up
approach

1) WG2/WG4: LEPTON FLAVOR AND CP (SISSA/LNF-RM/NA-SA)

5) S.T. Petcov

Predicting the values of the leptonic CP violation phases in theories with discrete flavour symmetries

arXiv:1405.6006 [hep-ph], Nucl. Phys. B 892 (2015) 400

6) Cesar Bonilla, Stefano Morisi, Eduardo Peinado, Jose W. F. Valle

Relating quarks and leptons with the T7 flavour group

arXiv:1411.4883[hep-ph], Phys.Lett. B742 (2015) 99

7) Claudia Hagedorn, Aurora Meroni, Emiliano Molinaro


Lepton mixing from $D(3n^2)$ and $D(6n^2)$ and CP

arXiv:1408.7118 [hep-ph], Nucl.Phys. B891 (2015) 499

8) I. Girardi, S.T. Petcov, A.V. Titov

Determining the Dirac CP Violation Phase in the Neutrino Mixing Matrix from Sum Rules

arXiv:1410.8056 [hep-ph], Nucl. Phys. B 894 (2015) 733



Discrete
Flavor
symmetries

I. Girardi, S.T. Petcov, A.V. Titov

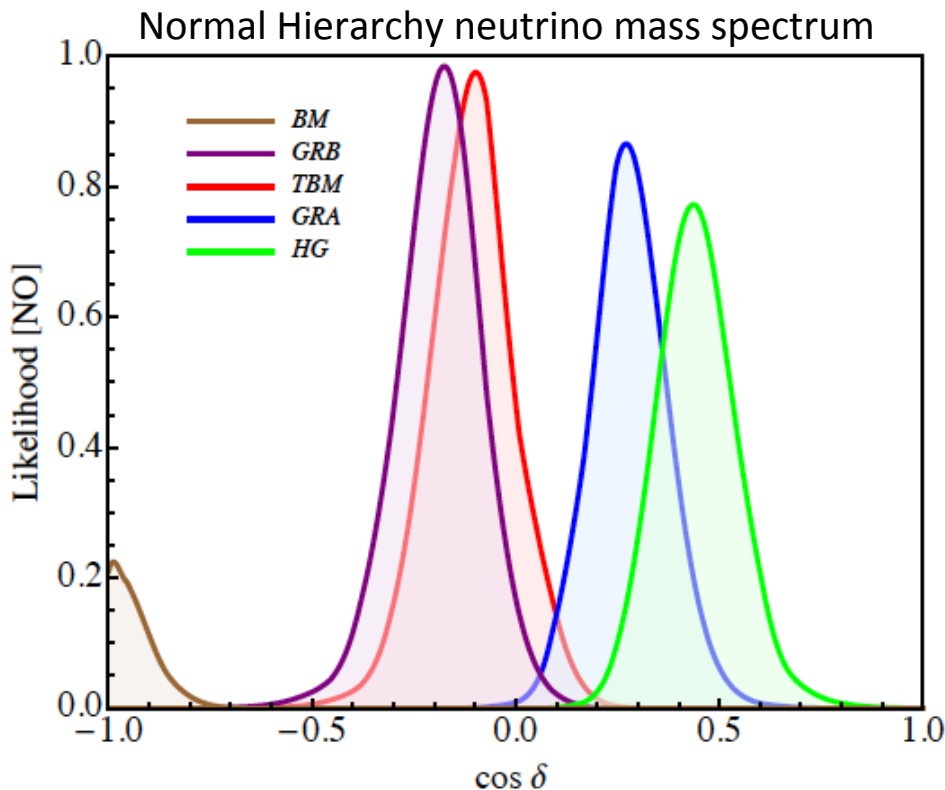
Determining the Dirac CP Violation Phase in the Neutrino Mixing Matrix from Sum Rules

arXiv:1410.8056 [hep-ph], Nucl. Phys. B 894 (2015) 733

$$U = U_e^\dagger \Psi_1 \tilde{U}_\nu Q_0$$

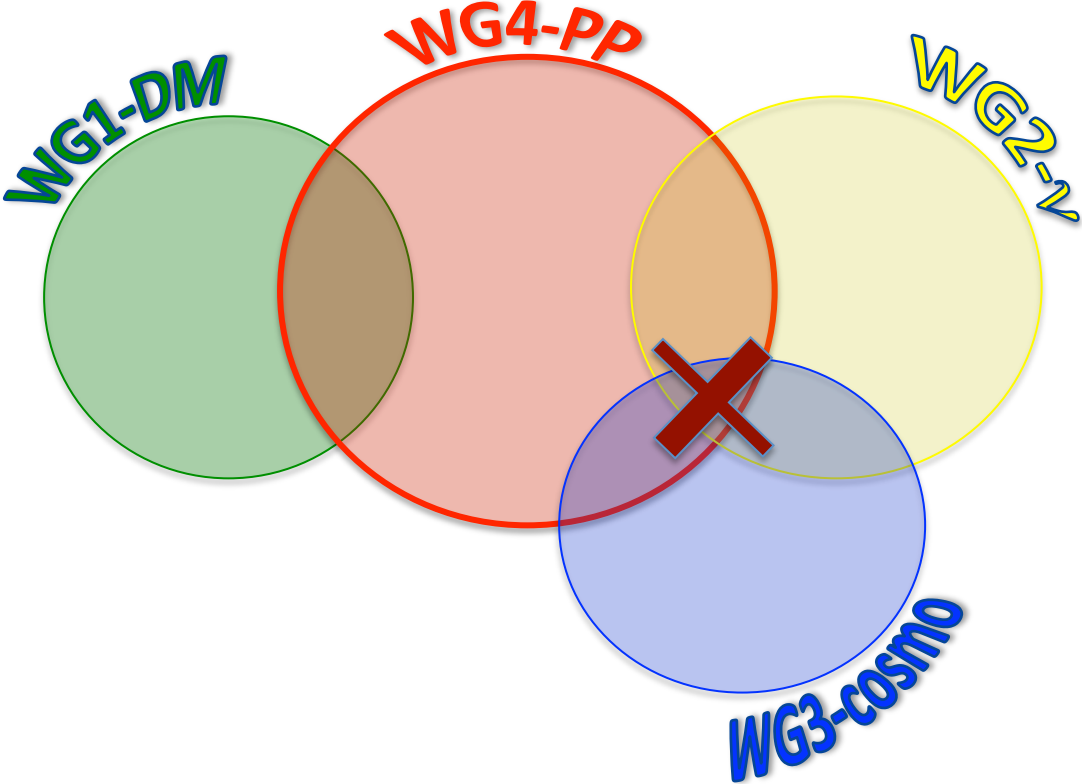
← phases

← BM, TBM, GRA, GRB, HG...



Tension for bimixing
Possible in future to discriminate

2) WG2/WG3/WG4: LEPTOGENESIS (SISSA/LNF-RM)

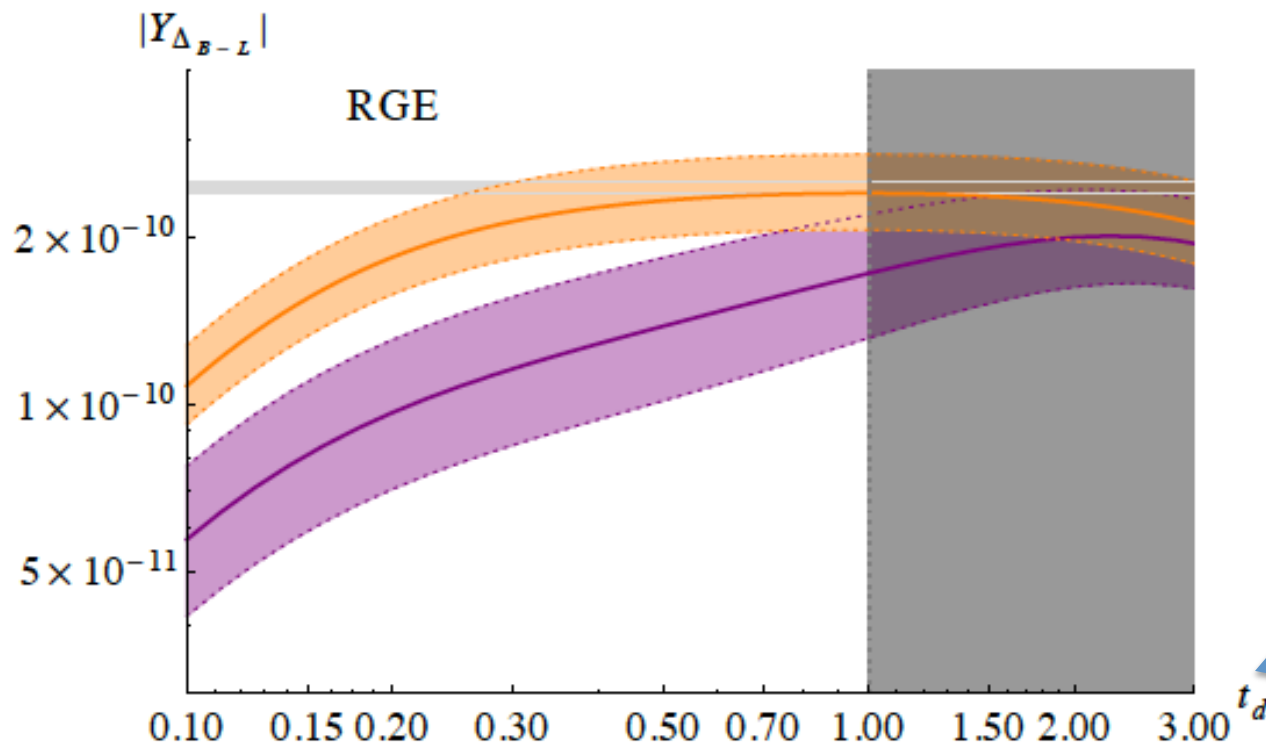


2) WG2/WG3/WG4: LEPTOGENESIS (SISSA/LNF-RM)

1) J. Gehrlein, S.T. Petcov, M. Spinrath, X. Zhang
Leptogenesis in an $SU(5) \times A5$ Golden Ratio Flavour Model
arXiv:1502.00110 [hep-ph], Nucl. Phys. B 896 (2015) 311

2) Chee Sheng Fong, Davide Meloni, Aurora Meroni, Enrico Nardi
Leptogenesis in $SO(10)$
arXiv:1412.4776 [astro-ph.HE], JHEP 1501 (2015) 111

In this work we have considered leptogenesis in a non-supersymmetric SO(10) GUT with fermion masses from the $\mathbf{10} \oplus \overline{\mathbf{126}}$ Higgs representations, which can (i) fit well all the low energy data, (ii) successfully account for unification of the gauge couplings, and (iii) allow for a sufficiently long lifetime of the proton. We have shown that, once the model parameters are fixed in terms of measured low energy observables, the requirement of successful leptogenesis can fix the only one remaining high energy parameter. We have

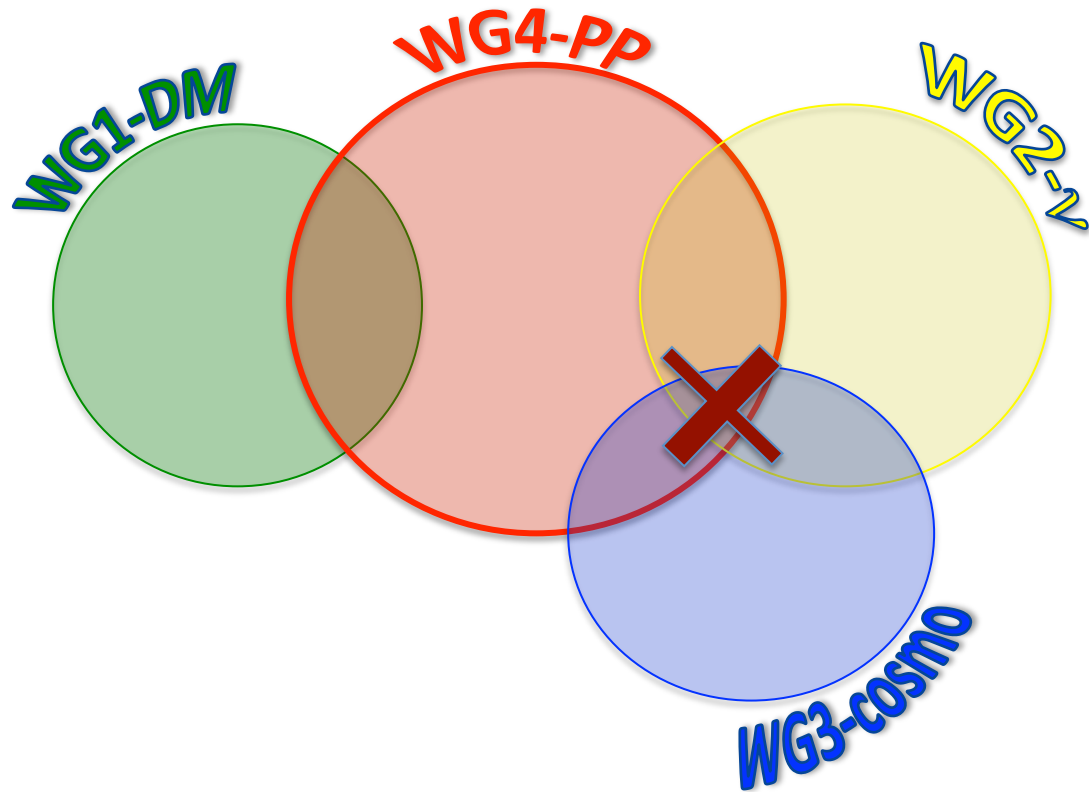


Two conditions

$$(i) M_{\Delta} \gtrsim M_{N_2} \simeq 2.0 \times 10^{11} \text{ GeV}$$

$$(ii) 0.3 \lesssim \left[\frac{\langle \Sigma_d^\dagger \Sigma_d \rangle}{\langle H_d^\dagger H_d \rangle} \right]^{1/2} \lesssim 1$$

3) WG2/WG3/WG4: SOLAR DATA CONSTRAINTS (FE/AQ)

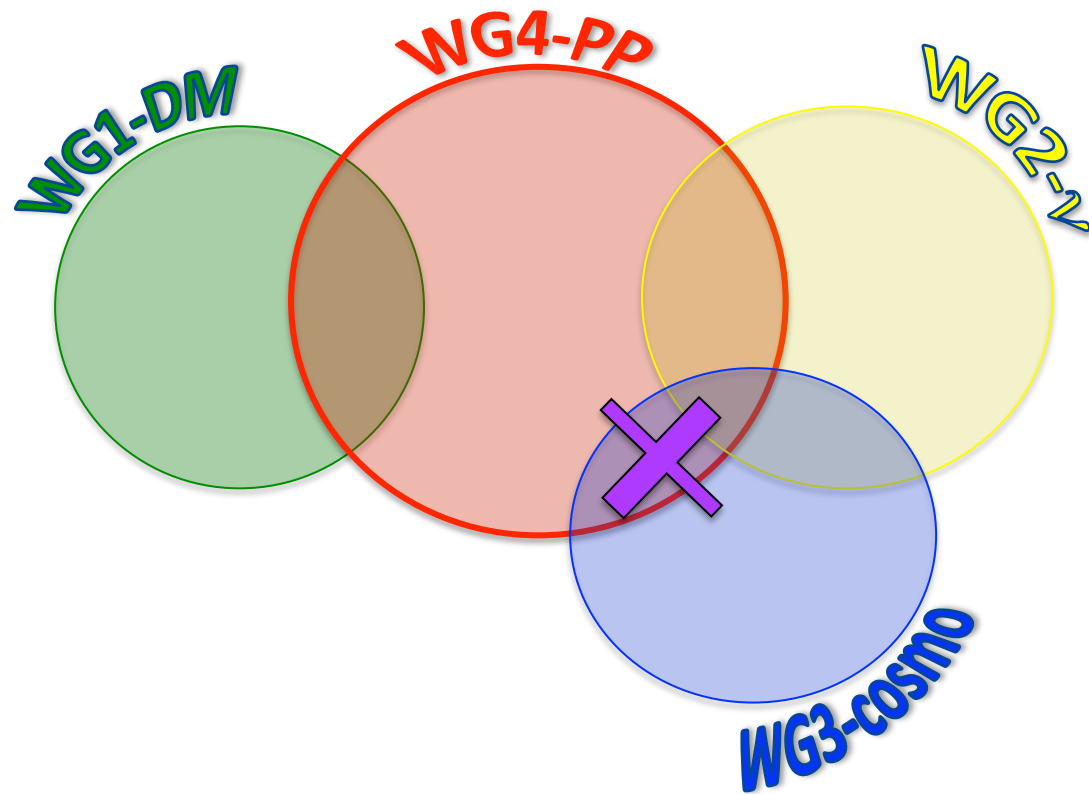


3) WG2/WG3/WG4: SOLAR DATA CONSTRAINTS (FE/AQ)

1) N. Vinyoles, A. Serenelli, F.L. Villante, S. Basu, J. Redondo, J. Isern
New axion and hidden photon constraints from a solar data global fit
arXiv:1501.01639 [astro-ph.SR]

2) Andrea Zanzi , Barbara Ricci
Chameleon fields and solar physics
arXiv:1405.1581 [hep-ph], Mod.Phys.Lett. A30 (2015) 10, 1550053

4) WG2/WG4: HIGGS AND INFLATION (FE)



4) WG2/WG4: HIGGS AND INFLATION (FE)

1) I. Masina

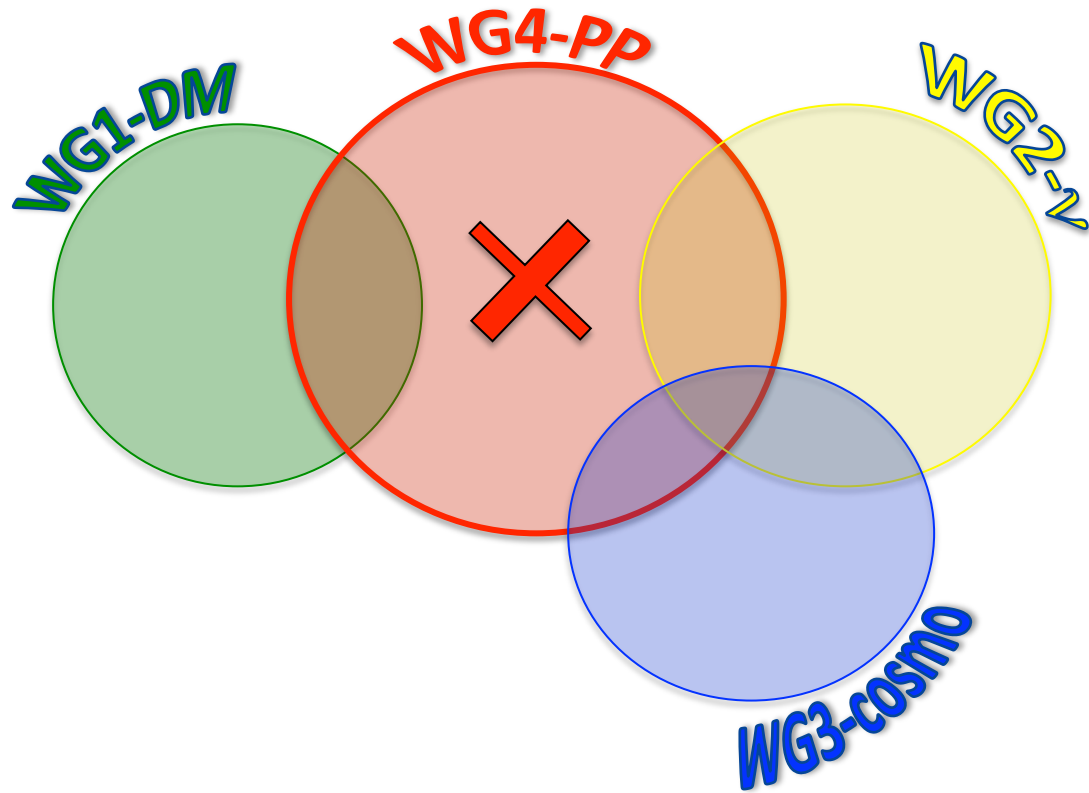
Gravitational wave background and Higgs false vacuum inflation

arXiv:1403.5244 [hep-ph], Phys. Rev. D 89 (2014) 12, 123505

... work in progress on non-minimal coupling →

Talk by Iacobellis

5) WG4: SUSY and LHC (PD/FE/AQ/NA-SA)



5) WG4: SUSY and LHC (PD/FE/AQ/NA-SA)

1) L. Calibbi, I. Galon, A. Masiero, P. Paradisi, Y. Shadmi

Talk by Paradisi

Charged Slepton Flavor post the 8 TeV LHC: A Simplified Model Analysis of Low-Energy Constraints and LHC SUSY Searches
e-Print: arXiv:1502.07753 [hep-ph]

2) Z. Berezhiani, M. Chianese, G. Miele, S. Morisi

Talk by Miele

Chances for SUSY-GUT in the LHC Epoch
arXiv:1505.04950 [hep-ph]

3) P.P. Giardino, K. Kannike, I. Masina, M. Raidal, A. Strumia

The universal Higgs fit

arXiv:1303.3570 [hep-ph], JHEP 1405 (2014) 046

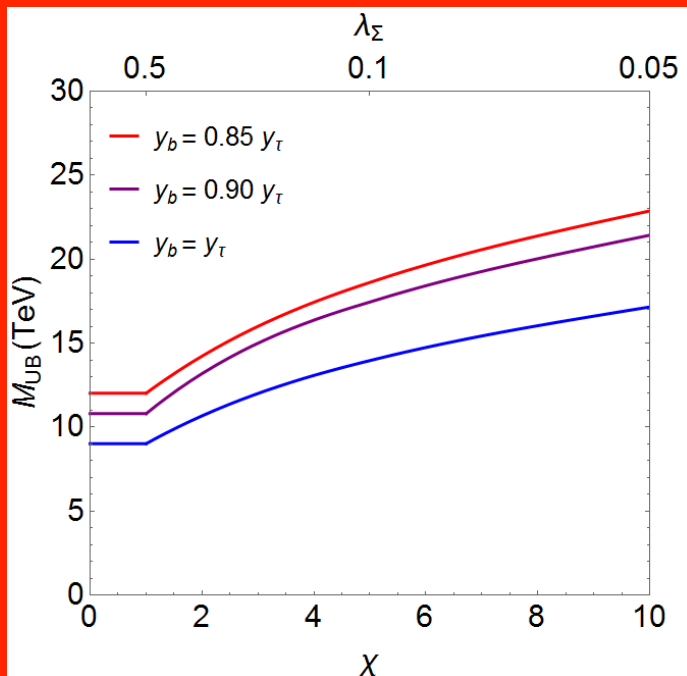
4) I. Masina, G. Nardini and M. Quiros,

Electroweak vacuum stability and finite quadratic radiative corrections

arXiv:1502.06525 [hep-ph]

Z. Berezhiani, M. Chianese, G. Miele, S. Morisi
Chances for SUSY-GUT in the LHC Epoch
arXiv:1505.04950 [hep-ph]

- After LHC run I (8 TeV), for planning new colliders it is of interest:
 - to reanalyze the room still remaining for SUSY-GUT inspired models;
 - to determine the **upper bound** for the energy below which SUSY signatures have to show up.
- Assuming one step unification (SU(5) bottleneck), under natural assumptions we have obtained general bounds on SUSY mass spectrum.



- We claim that if a SUSY-GUT model is the proper way to describe physics beyond the SM, the highest gluino or higgsino cannot have a mass larger than

20 TeV

5) WG4: SUSY and LHC (PD/FE/AQ/NA-SA)

1) L. Calibbi, I. Galon, A. Masiero, P. Paradisi, Y. Shadmi

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I.Masina, G.Nardini and M.Quiros,

Electroweak vacuum stability and finite quadratic radiative corrections

arXiv:1502.06525 [hep-ph]

Assuming the MSSM is the SM UV completion, we calculated the quadratic corrections to the (SM) Higgs mass: they turn out to be equivalent to finite threshold contributions

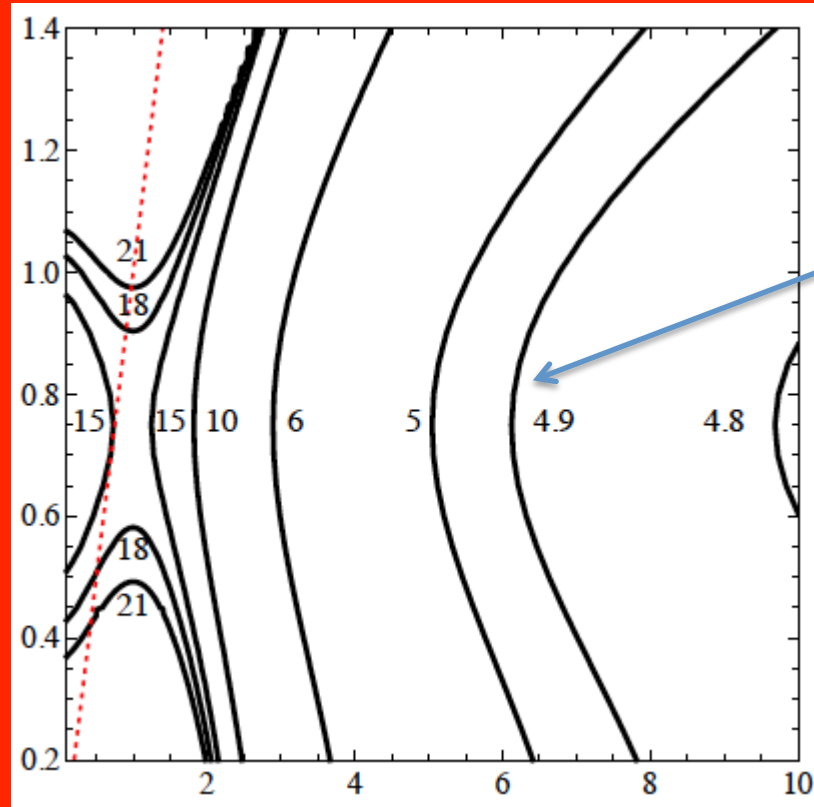
$$\begin{aligned} -\Delta m^2 = & \frac{1}{32\pi^2} \left\{ 6y_t^2 [G(m_Q^2) + G(m_U^2)] + 6y_b^2 [G(m_Q^2) + G(m_D^2)] + 2y_\tau^2 [G(m_L^2) + G(m_E^2)] \right. \\ & + 6y_t^2 X_t^2 \frac{G(m_Q^2) - G(m_U^2)}{m_Q^2 - m_U^2} + 6y_b^2 X_b^2 \frac{G(m_Q^2) - G(m_D^2)}{m_Q^2 - m_D^2} + 2y_\tau^2 X_\tau^2 \frac{G(m_L^2) - G(m_E^2)}{m_L^2 - m_E^2} \\ & - g_Y^2 (G(m_Q^2) - 2G(m_U^2) + G(m_D^2) - G(m_L^2) + G(m_E^2)) \cos 2\beta \\ & - 6g_2^2 \frac{M_2^2 G(M_2^2) - \mu^2 G(\mu^2)}{M_2^2 - \mu^2} - 2g_Y^2 \frac{M_1^2 G(M_1^2) - \mu^2 G(\mu^2)}{M_1^2 - \mu^2} \\ & - \left(12g_2^2 M_2 \mu \frac{G(M_2^2) - G(\mu^2)}{M_2^2 - \mu^2} + 4g_Y^2 M_1 \mu \frac{G(M_1^2) - G(\mu^2)}{M_1^2 - \mu^2} \right) \sin \beta \cos \beta \\ & \left. + G(m_H^2) (-6g_Z^2 \cos^2 2\beta + 2g_Z^2 + g_2^2) \right\}, \end{aligned} \quad (46)$$

Depending on the values of the different susy mass scales, these contributions can cancel even in the multi-TeV range

(unlike the case of a single susy mass scale, where the Veltman result is obtained)

I.Masina, G.Nardini and M.Quiros,
Electroweak vacuum stability and finite quadratic radiative corrections
arXiv:1502.06525 [hep-ph]

$$M_{1/2}/M_0$$

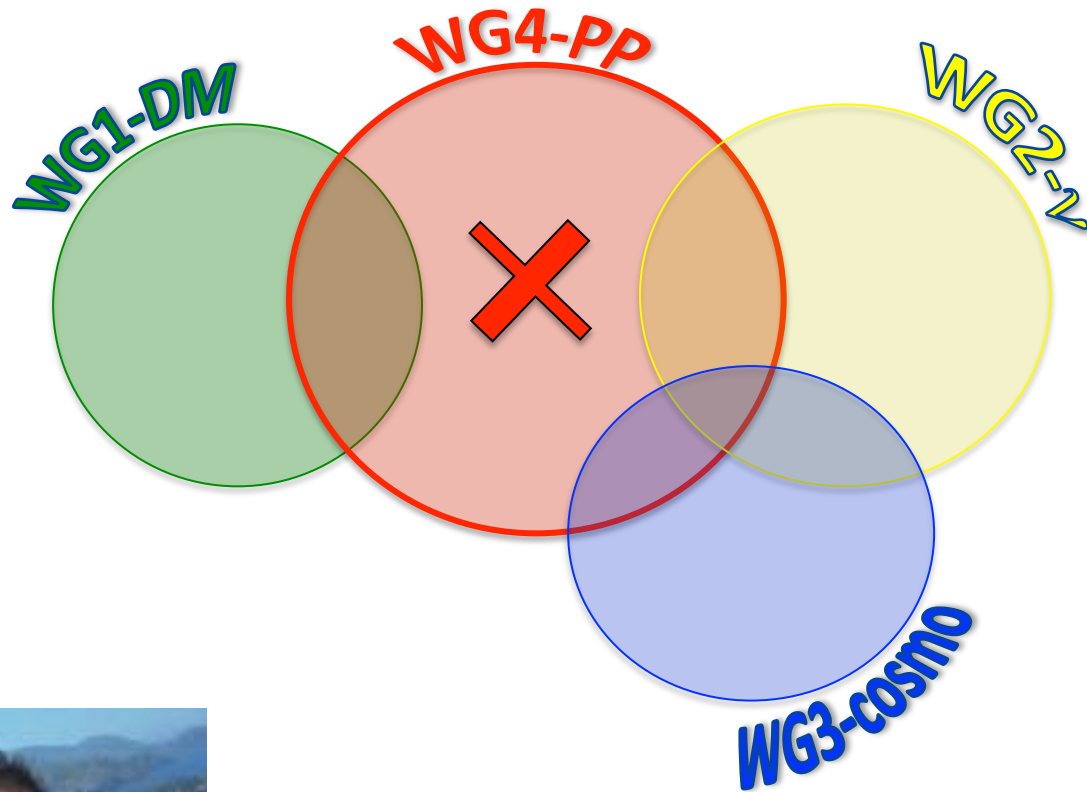


$\text{Log}_{10} M_0 / \text{GeV}$

$$M_H/M_0$$

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6) WG4: QUANTUM THEORIES (TO/NA-SA/AQ)



6) WG4: QUANTUM THEORIES (TO/NA-SA/AQ)

1) S.A. Alavi, C. Giunti

Which is the Quantum Decay Law of Relativistic Particles?

arXiv:1412.3346 [quant-ph], Europhys.Lett. 109 (2015) 6, 60001

2) G. Esposito, G. Marmo, G. Miele, G. Sudarshan

Advanced concepts in quantum mechanics

Published by Cambridge University Press , 2015

3) A. Addazi, G. Esposito

Nonlocal quantum field theory without acausality and nonunitarity at quantum level: is SUSY the key?

arXiv:1502.01471 [hep-th], Int.J.Mod.Phys. A30 (2015) 1550103

4) A. Iorio, G. Lambiase

Quantum Field Theory in curved graphene spacetime, Locachevsky geometry, Weyl symmetry, Hawking effect, and all that

arXiv:1308.0265 [hep-th], Phys. Rev. D 90 (2014) 025006

5) Luigi Rosa, Lucia Trozzo

Casimir energy in a spherical surface within surface impedance approach: the Drude model

e-Print: arXiv:1408.2366 [hep-th]

6) F.E. Canfora, D. Dudal , I.F. Justo, P. Pais, L. Rosa, D. Vercauteren

Effect of the Gribov horizon on the Polyakov loop and vice versa

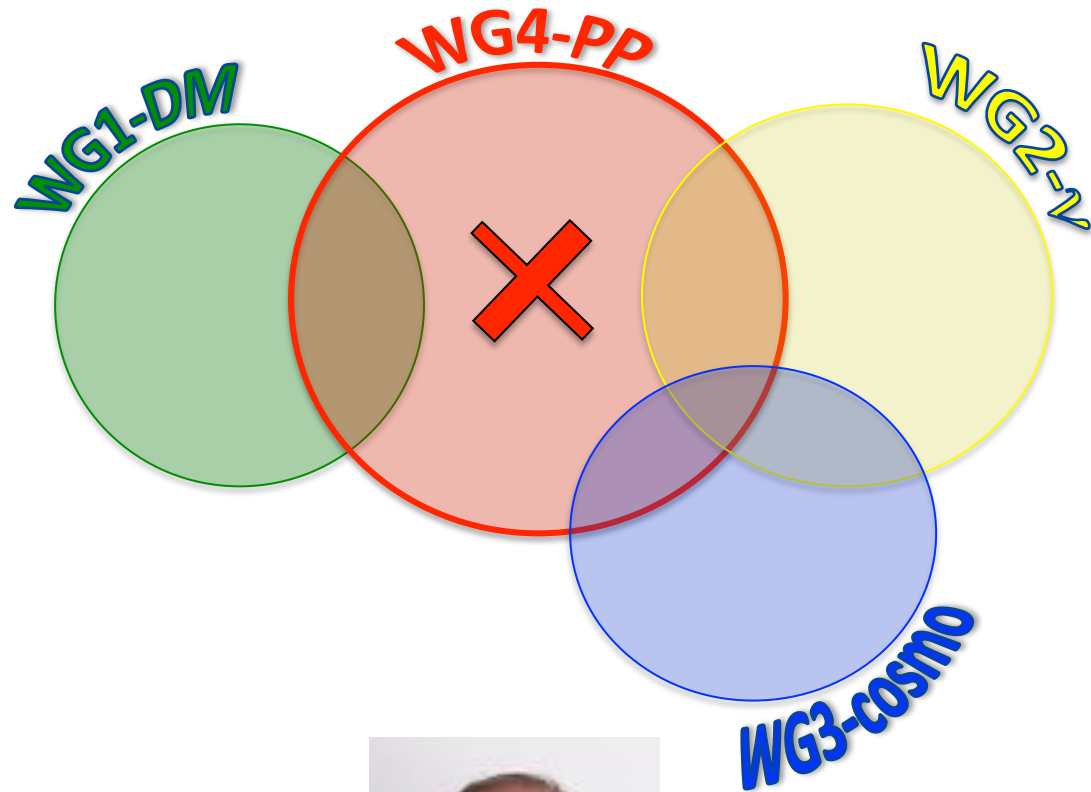
e-Print: arXiv:1505.02287 [hep-th]

7) F. Canfora, M. Kurkov, L. Rosa, P. Vital

The Gribov problem in Noncommutative QED

e-Print: arXiv:1505.06342 [hep-th]

7) WG2/WG4: GRAVITY (BA)



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1) M. Gasperini

The twin paradox in the presence of gravity

arXiv:1409.1818 [gr-qc], Mod. Phys. Lett. A29 (2014) 1450149

7) WG2/WG4: GRAVITY (BA)

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The twin paradox in the presence of gravity

arXiv:1409.1818 [gr-qc], Mod. Phys. Lett. A29 (2014) 1450149

Conventional wisdom, based on kinematic (flat-space) intuition, tell us that a static twin is aging faster than his traveling twin brother. However, such a situation could be exactly inverted if the two twins are embedded in an external gravitational field, and if the (dynamical) distortion of the space-time geometry, due to gravity, is strong enough to compensate the kinematic effect of the relative twin motion.

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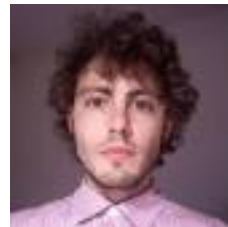
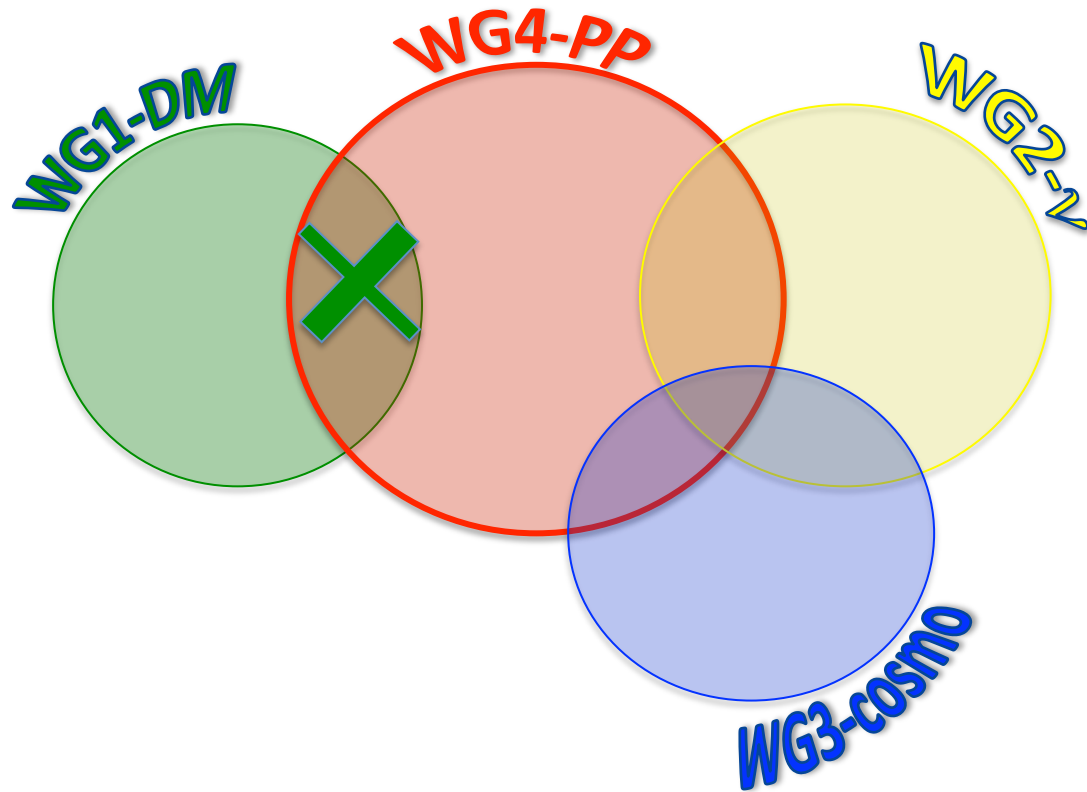
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arXiv:1409.1818 [gr-qc], Mod. Phys. Lett. A29 (2014) 1450149

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Let us hope that future technology will be able to apply this exceptional virtue of gravity to keep us young, as long as possible!

8) WG4: NEUTRON-ANTINEUTRON (AQ)



8) WG4: NEUTRON-ANTINEUTRON (AQ)

1) D.G. Phillips, W.M. Snow, K. Babu, S. Banerjee, D.V. Baxter, Z. Berezhiani et al.
Neutron-Antineutron Oscillations: Theoretical Status and Experimental Prospects
arXiv:1410.1100 [hep-ex], Accepted in Phys. Rept.

2) Z. Berezhiani, A. Vainshtein
Neutron-Antineutron Oscillation as a Signal of CP Violation
arXiv:1506.05096 [hep-ph]

Talk by Berezhiani

Neutron-antineutron oscillations phenomenology

3) A. Addazi, M. Bianchi
Neutron Majorana mass from exotic instantons
arXiv:1407.2897 [hep-ph], JHEP 1412 (2014) 089

4) A. Addazi, M. Bianchi
Neutron Majorana mass from Exotic Instantons in a Pati-Salam model
arXiv:1502.08041 [astro-ph.CO], JHEP 1506 (2015) 012

5) A. Addazi
Exotic vector-like pair of color-triplet scalars
arXiv:1501.04660 [hep-ph], JHEP 1504 (2015) 153

DM and Neutron oscillations and/or exotica

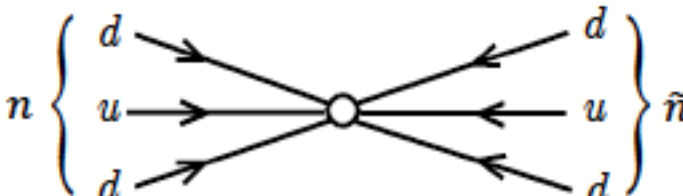
Z. Berezhiani, A. Vainshtein
 Neutron-Antineutron Oscillation as a Signal of CP Violation
 arXiv:1506.05096 [hep-ph]

Assuming the Lorentz and CPT invariances we show that neutron-antineutron oscillation implies breaking of CP along with baryon number violation. Presence of external magnetic field does not add any new operator to mixing of neutron and antineutron.

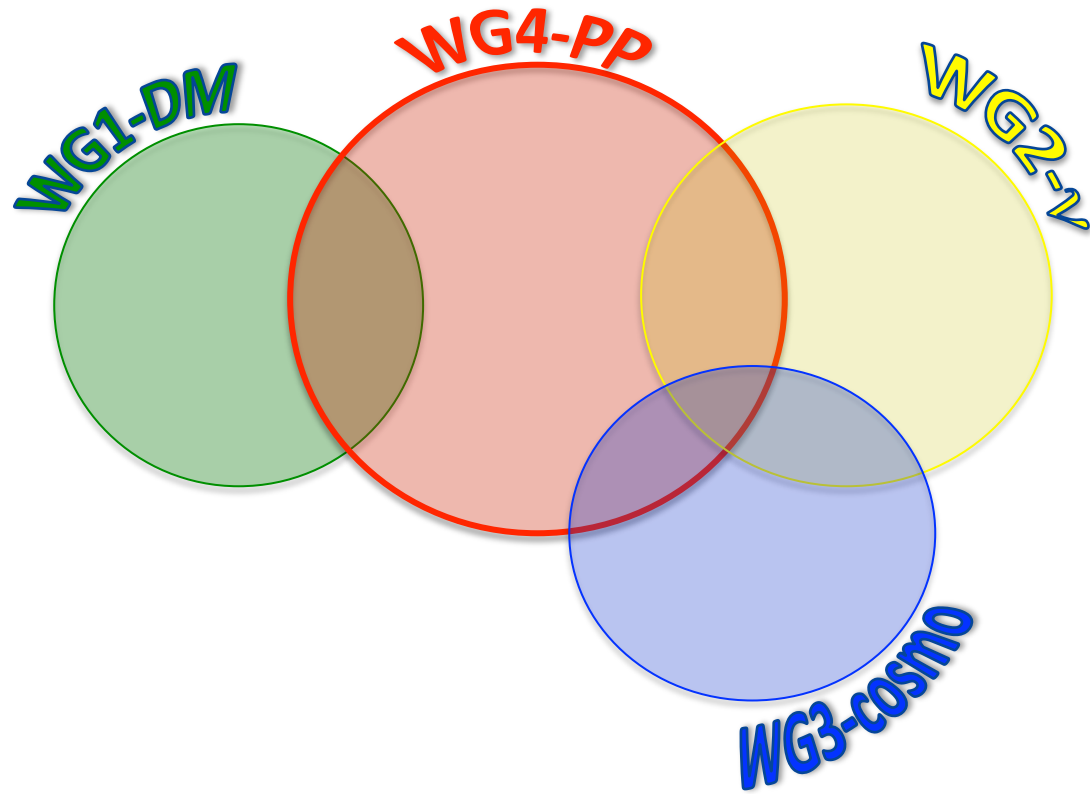
B can be broken by two units via Majorana terms $\Delta\mathcal{L}_B = -\frac{1}{2} \epsilon [n^T C n + \bar{n} C \bar{n}^T]$

$\Delta\mathcal{L}_B$ relevant for $n - \bar{n}$ oscillation (generates mass splitting): *conserves* C but *breaks* P and CP

Neutron antineutron oscillation emerges in BSM physics via six-quark effective operators

$$\mathcal{O} = \frac{1}{M^5} u d d u d d$$


automatically satisfy two (in fact all three) of Sakharov's conditions.



... buon lavoro!

Beyond Standard Model

Several phenomena or open problems suggest the presence of physics beyond Standard Model (SM):

- hint of gauge unification;
 - structure of fermion masses;
 - hierarchy problem;
 - baryogenesis.
- The Supersymmetric Grand Unified Theories (SUSY-GUTs) are able to address part of these problems.
 - After the 8 TeV LHC run I, we try:
 - to reanalyze the room still remaining for SUSY-GUT inspired models;
 - to determine the **upper bound** M_{UB} for the energy below which SUSY signatures have to show up.

Assumptions

To perform our study, we require:

- One step gauge unification at a single energy scale M_{GUT} , without intermediate symmetry scales;



SU(5) Bottleneck

- Consistency of third family fermion masses and possible Yukawa b- τ unification;

$$y_b(M_{GUT}) = y_\tau(M_{GUT}) \left(1 + \mathcal{O} \left(\frac{y_\mu(M_{GUT})}{y_\tau(M_{GUT})} \right) \right)$$

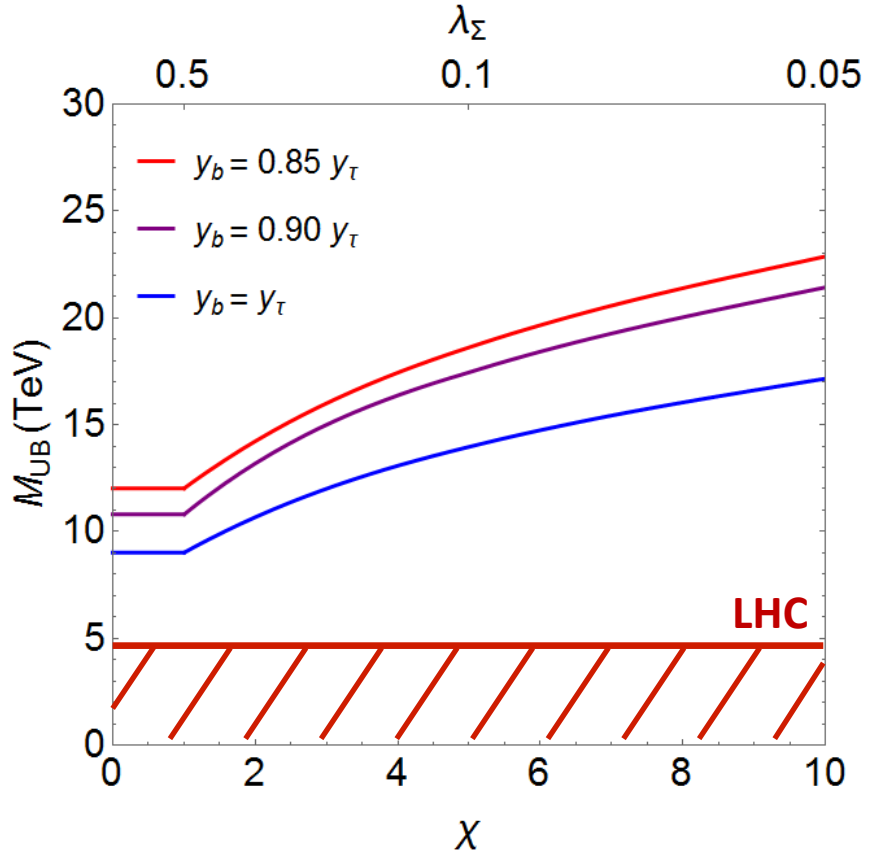
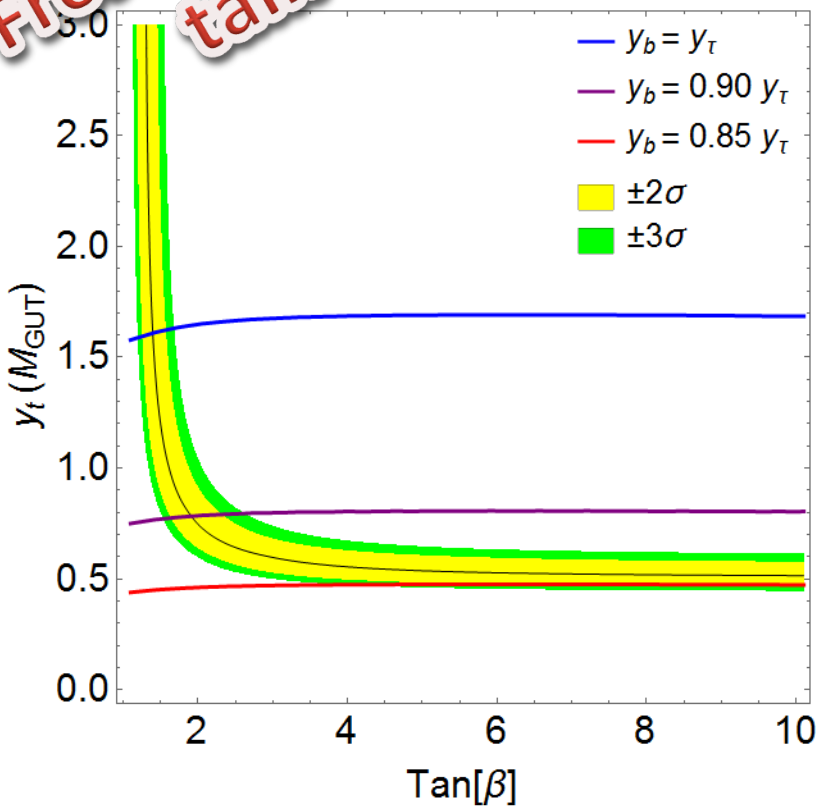
- Consistency with the experimental limit on proton decay;
- Absence of special fine tunings among the parameters in the GUT scenario, implying couplings $\mathcal{O}(1)$ above M_{GUT} .



Naturalness

Upper bound for SUSY physics

From Miele's talk



• Larger values for GUT threshold χ are not allowed since:

- naturalness requirement implies $\lambda \downarrow \Sigma \sim \mathcal{O}(1)$;
- $M \downarrow GUT$ unnaturally approaches $M \downarrow Plack$.



$M_{UB} \sim 20$ TeV

Conclusions

After LHC run I (8 TeV), for planning new colliders it is of interest:

- to reanalyze the room still remaining for SUSY-GUT inspired models;
- to determine the **upper bound** M_{UB} for the energy below which SUSY signatures have to show up.
- Assuming one step unification (SU(5) bottleneck), under natural assumptions we have obtained general bounds on SUSY mass spectrum.
- We claim that if a SUSY-GUT model is the proper way to describe physics beyond the SM, the lightest gluino or higgsino cannot have a mass larger than

$$M_{UB} \sim 20 \text{ TeV}$$