

[Renato.campanini@unibo.it](mailto:Renato.campanini@unibo.it)

[campanini@bo.infn.it](mailto:campanini@bo.infn.it)

QGP in pp and pPB ?

Physics Letters B 703 (2011) 237–245



ELSEVIER

Contents lists available at ScienceDirect

Physics Letters B

[www.elsevier.com/locate/physletb](http://www.elsevier.com/locate/physletb)



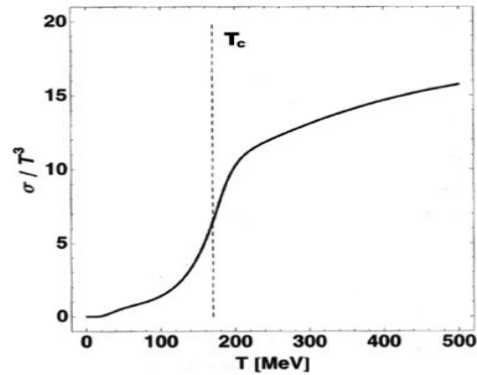
Experimental equation of state in pp and p $\bar{p}$  collisions and phase transition to quark gluon plasma

Renato Campanini<sup>a,b,\*</sup>, Gianluca Ferri<sup>a</sup>

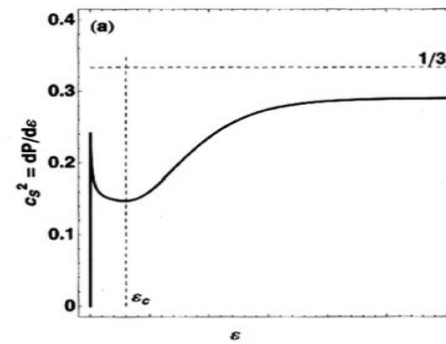
<sup>a</sup> Università di Bologna, Dipartimento di Fisica, viale C. Berti Pichat 6/2, I-40127 Bologna, Italy

<sup>b</sup> INFN, Sezione di Bologna, viale C. Berti Pichat 6/2, I-40127 Bologna, Italy

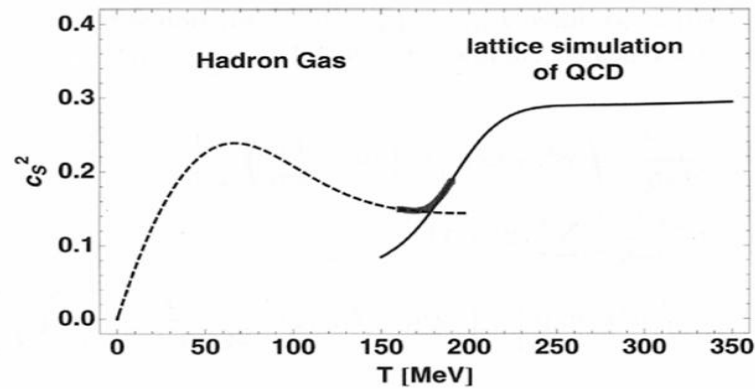
# Entropy, Temperature, Sound velocity



(a)



(b)

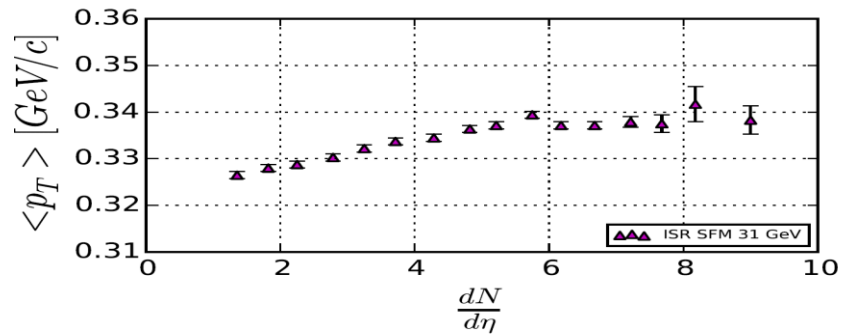
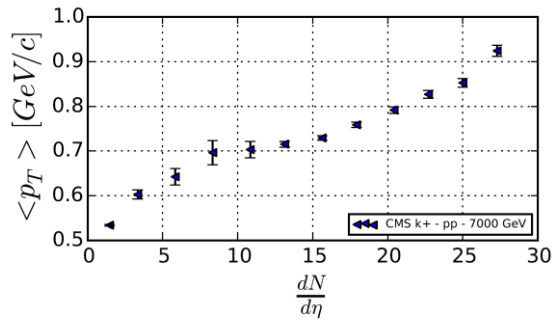
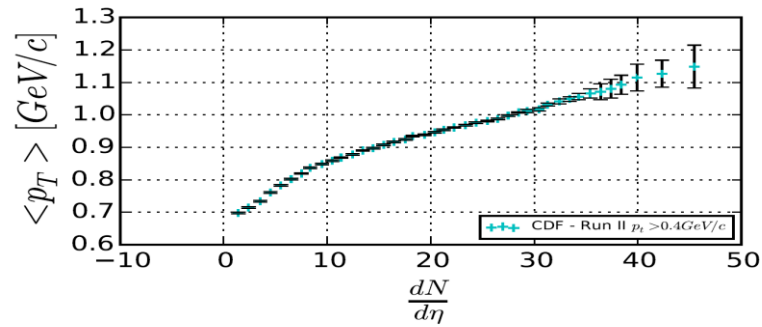
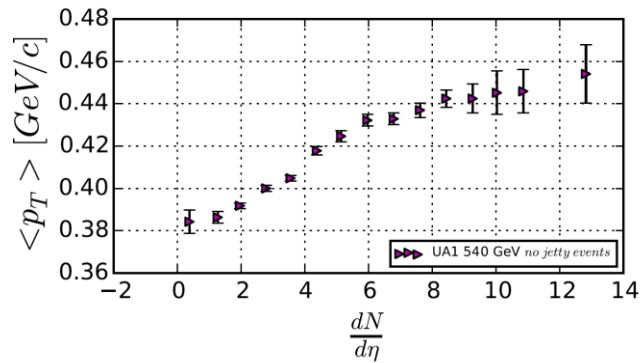


(c)

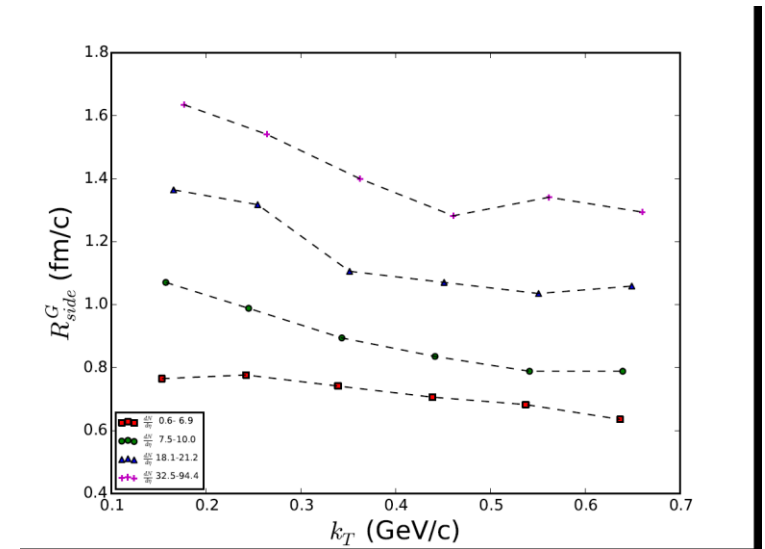
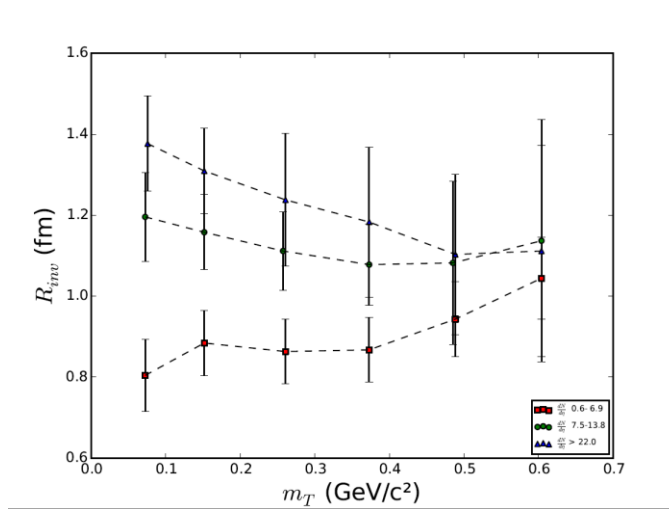
# Results from $\langle Pt \rangle$ vs $dn/d\eta$ and from Transverse Radius vs $dn/d\eta$

- From 31 GeV to 7 TeV there is a slope change in the  $\langle Pt \rangle$  vs  $dn/d\eta$  at  $dn/d\eta$  around 6 (some examples in next slides)
- . From FEMTOSCOPY: the transverse Radius  $R_{side}$  is independent from  $K_T$  (mt) in events with  $dn/d\eta < 7$

# Average Transverse Momentum vs pseudorapidity density, some examples

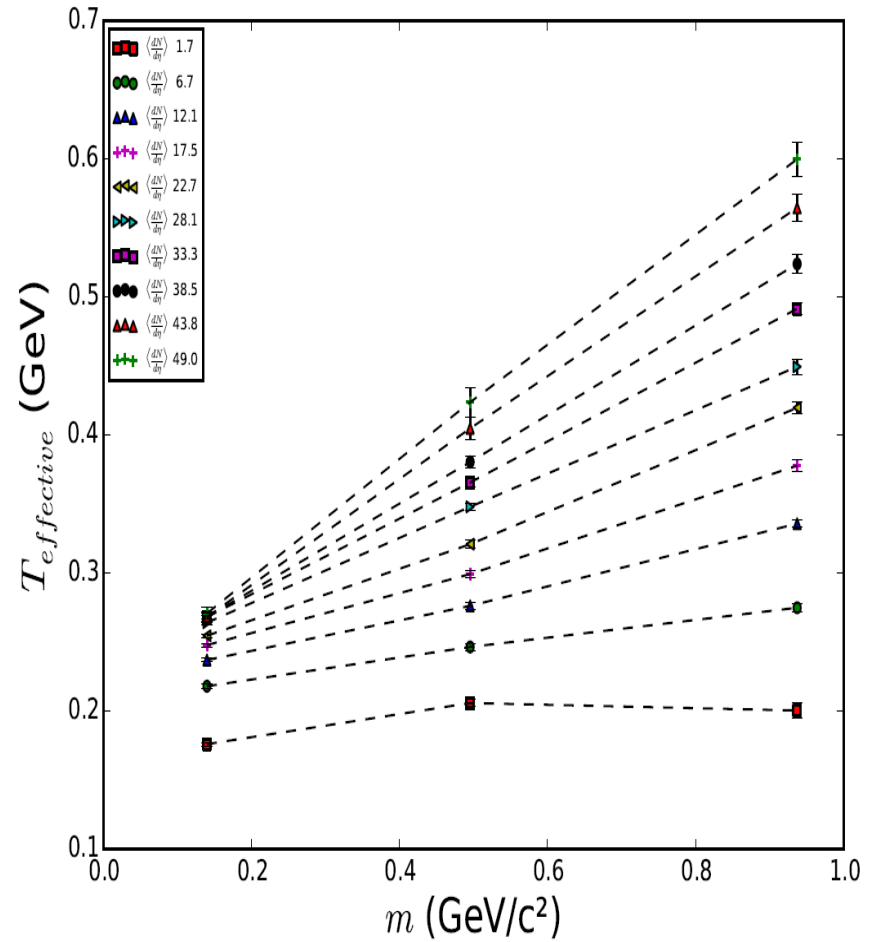
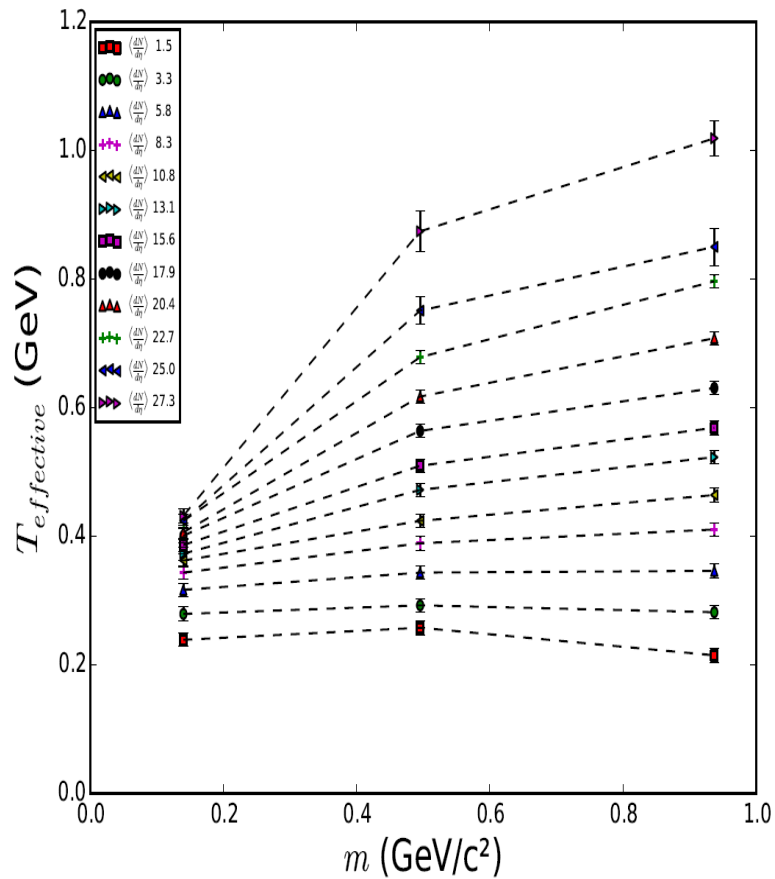


# ALICE pion pion femtoscopy in pp events 7 TeV, $R_{inv}$ and $R_{gside}$ vs $k_T$ in events at different $dn/d\eta$



Teff changes with mass only in events with  $dn/d\eta >$  about 7 (in pp, left and in pPb, right)

TRANSVERSE EXPANSION ONLY IN EVENTS WITH  $dn/d\eta$  about 7



$$\sigma_S / \langle p_T \rangle^3 \quad vs \quad p_T$$

## 2.2. Entropy Density Estimation

The initial energy density in the rest system of a head-on collision has been argued to be [6]:

$$\epsilon \simeq \frac{\frac{dN_{ch}}{d\eta} \cdot \frac{3}{2} \langle p_T \rangle}{V}$$

$V$  denotes the volume into which the energy is deposited. Similarly the initial entropy density is [2]:

$$\sigma \simeq \frac{\frac{dN_{ch}}{d\eta} \cdot \frac{3}{2}}{V}$$

As a result,  $\epsilon$  is equal to  $\sigma \cdot \langle p_T \rangle$ . The volume  $V$  may be estimated as  $V = S \cdot ct$ , where  $S$  is the interaction area and  $ct$  is a longitudinal dimension we can traditionally consider to be about 1 fm long.

In order to study our system, we will use the quantity

$$\sigma_S = \frac{\frac{dN_{ch}}{d\eta} \cdot \frac{3}{2}}{S}$$

as an estimation of entropy density. In models like color glass condensate and percolation, the system physics depends on  $\sigma_S$  [36, 37, 38]. For the estimation

## $\sigma_S / \langle p_T \rangle^3$ vs $p_T$

### 2.4. $\sigma_S / \langle p_T \rangle^3$ vs $p_T$

Using the estimated  $\sigma_S$ , the relation  $\langle p_T \rangle$  vs  $\sigma_S$  can be studied. A slope change in  $\langle p_T \rangle$  vs  $\sigma_S$  plots is found at  $\sigma_S$  between 2.5 and 3 fm<sup>-2</sup>, depending on the method used for the estimation of area  $S$  and corresponds directly to the slope change seen in  $\langle p_T \rangle$  vs  $\frac{dN_{ch}}{d\eta}$  at  $\frac{dN_{ch}}{d\eta} \simeq 6$ .

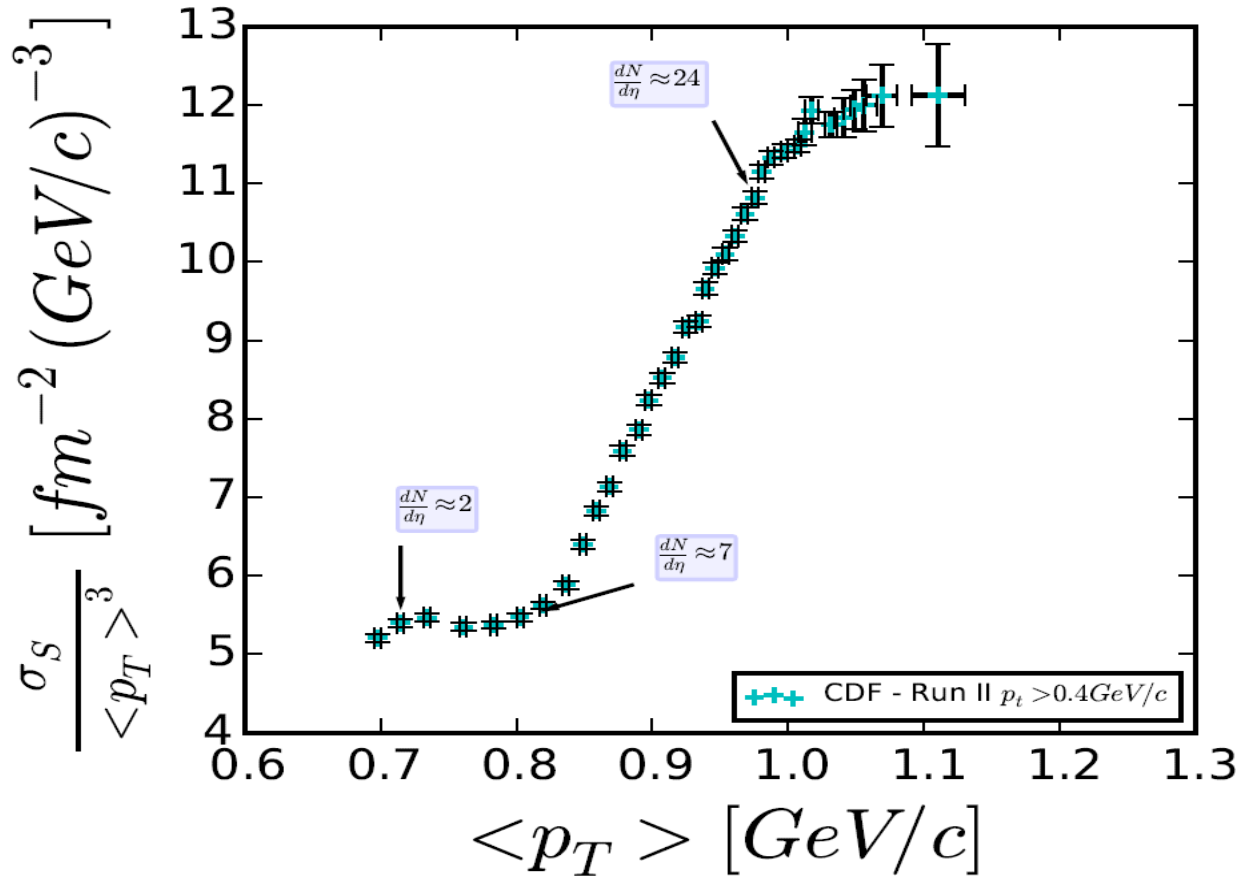
Starting from  $\sigma_S$  and  $\langle p_T \rangle$ , we plotted  $\sigma_S / \langle p_T \rangle^3$  vs  $\langle p_T \rangle$  curves, as an experimental approximation of  $\sigma / T^3$  vs  $T$  curves. See Figs. 3 and 4. We obtained very similar  $\sigma_S / \langle p_T \rangle^3$  vs  $\langle p_T \rangle$  curves from other pp and p $\bar{p}$  experiments [27, 29, 30, 31] (not shown).



$\sigma_S / \langle p_T \rangle^3$  vs  $p_T$

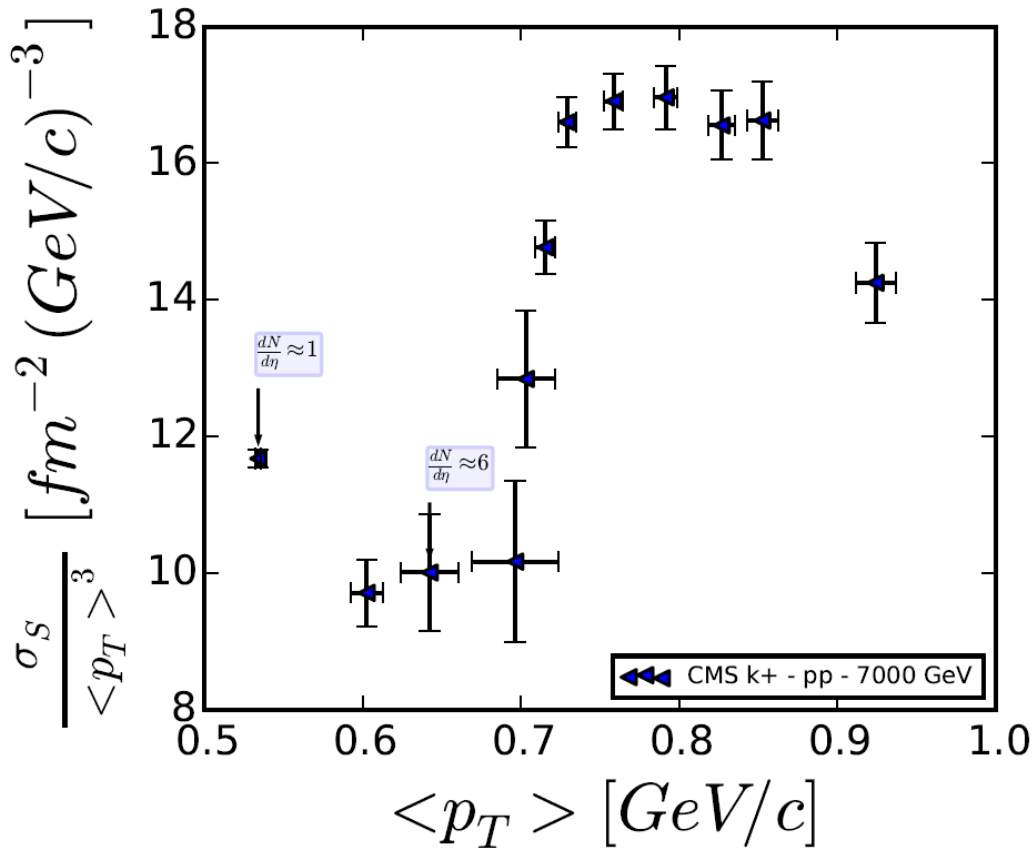
CDF 1.96 TeV

CHARGED



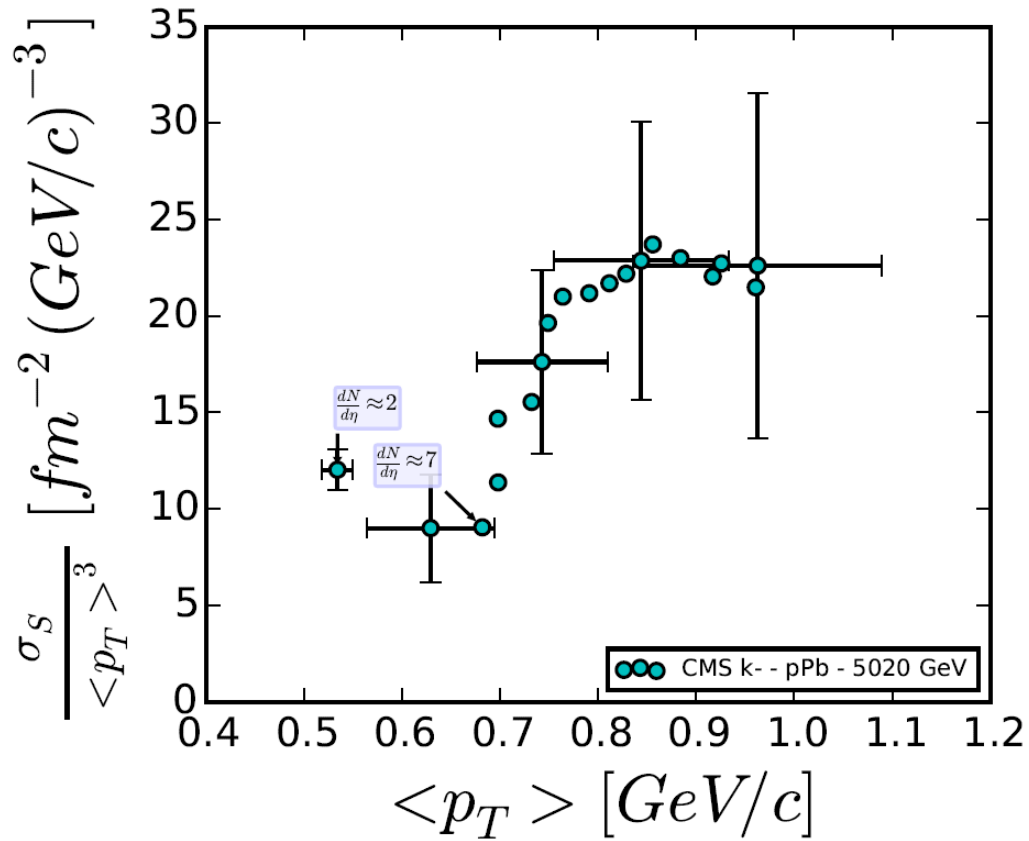
$\sigma_S / \langle p_T \rangle^3$  vs  $p_T$

CMS pp 7 TeV K+



$\sigma_S / \langle p_T \rangle^3$  vs  $p_T$

CMS pPb 5.02 TeV K-



# SOUND VELOCITY

$$c_s^2 = \frac{\sigma}{T} \cdot \frac{dT}{d\sigma}$$

## 2.5. Sound velocity $c_s^2$

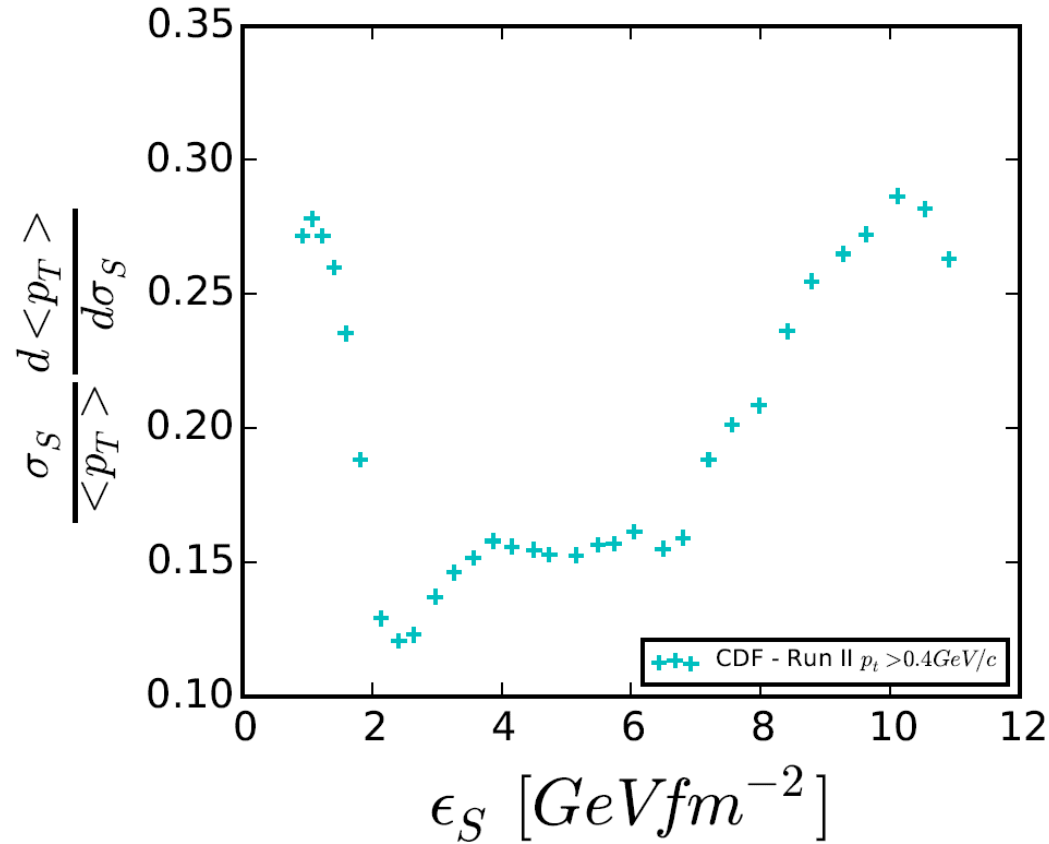
One of the physical quantities used to characterize the state of a system is its squared sound velocity, defined as  $c_s^2 = \frac{\sigma}{T} \cdot \frac{dT}{d\sigma}$ , for constant  $V$  [11]. In our study, we approximate it with  $c_s^2 = \frac{\sigma_S}{\langle p_T \rangle} \cdot \frac{d\langle p_T \rangle}{d\sigma_S}$ . It is really interesting that if  $\langle p_T \rangle$  is proportional to  $T$  and if  $\sigma_S$  is proportional to the entropy density, then the  $c_s^2$  value obtained in this approximation is equal to the right value of  $c_s^2 = \frac{\sigma}{T} \cdot \frac{dT}{d\sigma}$ , because proportionality constants cancel out. In order to obtain our  $c_s^2$  estimation, from  $\langle p_T \rangle$  vs  $\frac{dN_{ch}}{dn}$  curves and from  $\sigma_S$  values, we compute the curve  $p_T$  vs  $\sigma_S$ , to which we apply numerical derivation. We cope with the statistical fluctuation in data points using a combination of Gaussian and Savitzky-Golay filters [43]. Examples of  $c_s^2$  vs  $\langle p_T \rangle$  curves are shown in Figs. 5a and 5c.

# SOUND VELOCITY

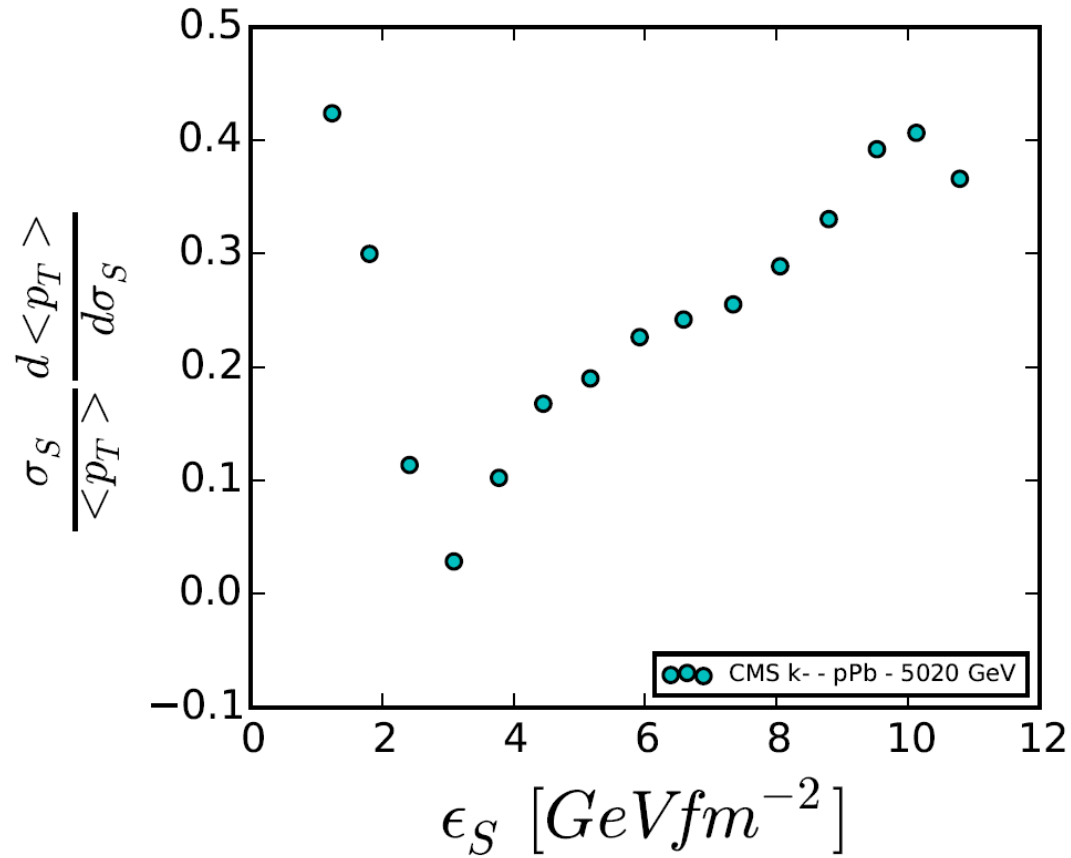
The so obtained  $c_s^2$  estimation resembles the typical shape of a phase transition or a crossover: a descent, a minimum region and a following rise, as it's also obtained analytically from EOSs which present a phase transition or a crossover. The minimum value reached by the estimation of  $c_s^2$  in the different experimental curves varies from 0.08 to 0.18 and could correspond to what it's called the EOS softest point [17].

Recently Refs. [44, 45, 46] estimate  $c_s^2$  minimum value for realistic EOS to be around 0.14. From  $\epsilon_S \simeq \langle p_T \rangle \cdot \sigma_S$  we compute  $c_s^2$  vs  $\epsilon_S$  curves, that are approximations of  $c_s^2$  vs energy density. We report these curves in Figs. 5b and 5d.

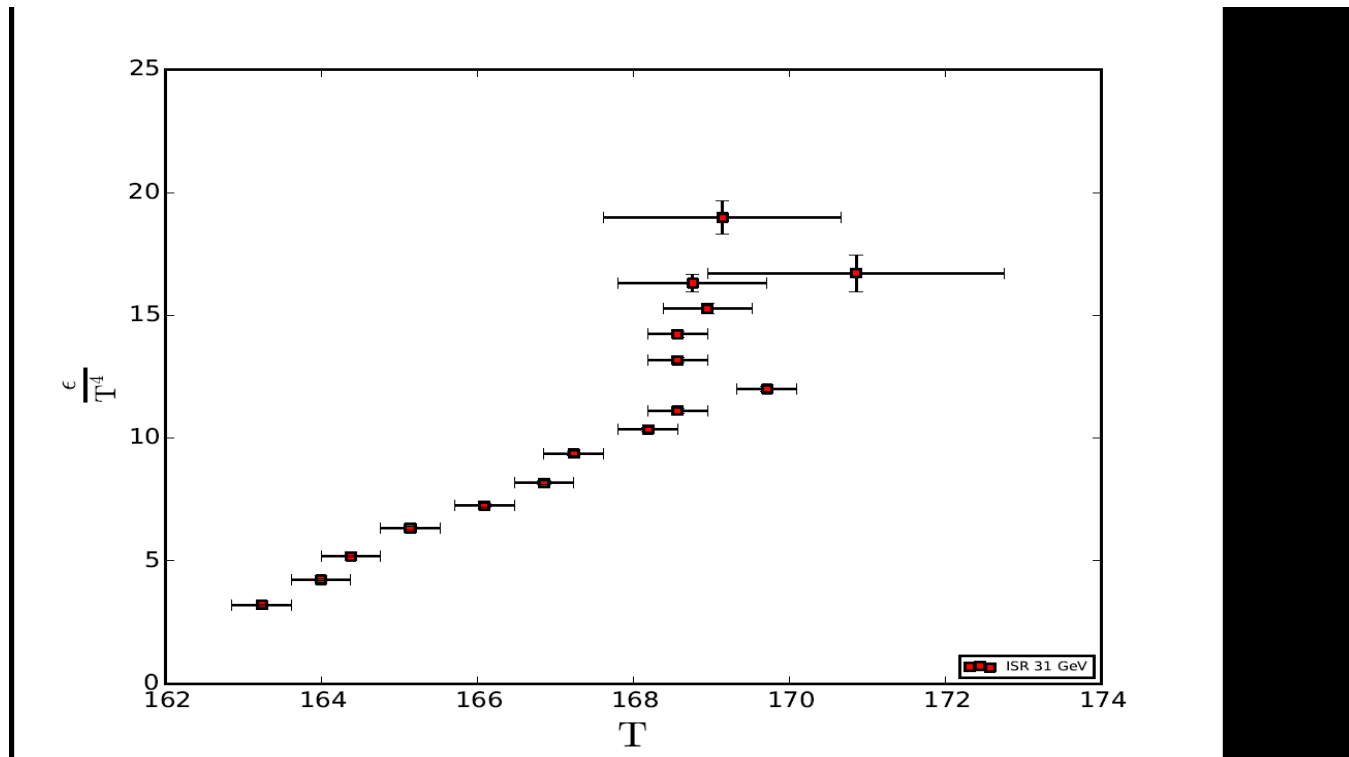
# SOUND VELOCITY CDF pp 1.96 TeV



# SOUND VELOCITY CMS K- pPb 5.02 TeV

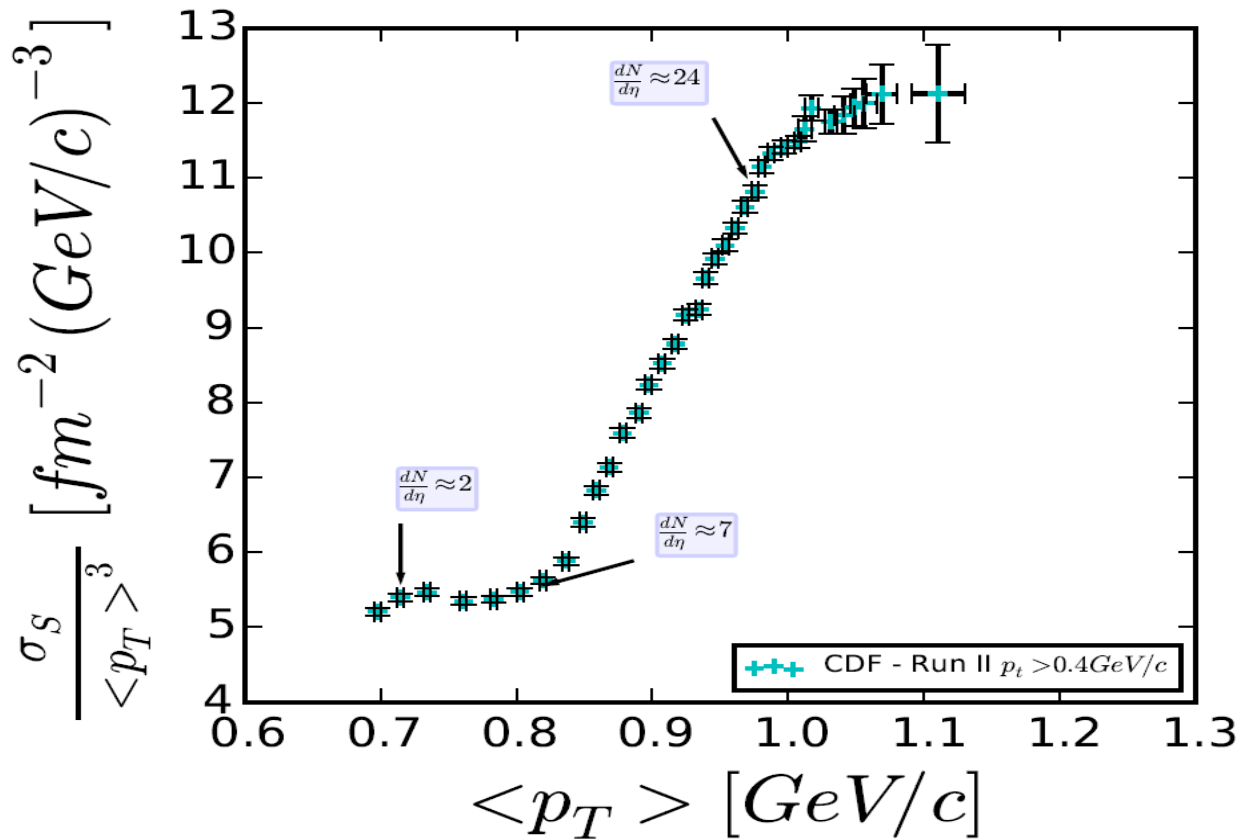


- AT LOW ENERGY (31 GeV) IN EVENTS WITH LOW : NO JETS, NO TRANSVERSE EXPANSION → TRANSVERSE MOMENTUM MEASURES TEMPERATURE! (exponential or Hagedorn Model)





# CROSSOVER FROM $dn/d\eta$ 6 to about 24?



# WHAT HAPPENS IN EVENTS with $dn/d\eta$ in the range 6 - 24?

- Ridge
- Multiparticle correlations
- Particle yield
- $Y(3s)/Y(1s)$  ;  $y(2s)/y(1s)$  ratio

# Ridge, multiparticle correlations

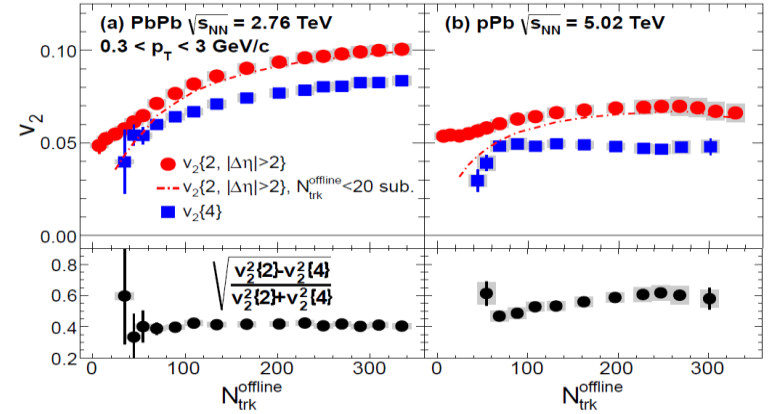
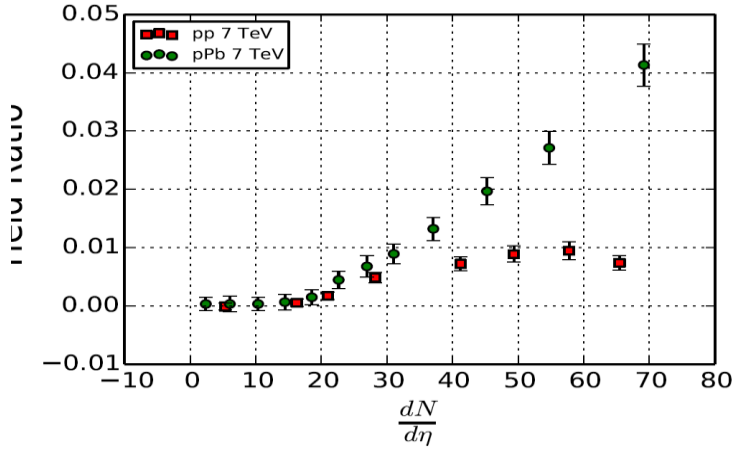


Figure 9: Top: the  $v_2\{2, |\Delta\eta| > 2\}$  (circles) and  $v_2\{4\}$  (squares) values as a function of  $N_{trk}^{offline}$  for  $0.3 < p_T < 3$  GeV/c, in 2.76 TeV PbPb collisions (left) and 5.02 TeV pPb collisions (right). Bottom: upper limits on the relative  $v_2$  fluctuations estimated from  $v_2\{2\}$  and  $v_2\{4\}$  in 2.76 TeV PbPb collisions (left) and 5.02 TeV pPb collisions (right). The error bars correspond to statistical uncertainties, while the shaded areas denote the systematic uncertainties. Results after subtracting the low-multiplicity data ( $N_{trk}^{offline} < 20$ ) are also shown (curves).

Multi-particle azimuthal correlations in p-Pb and Pb-Pb collisions ALICE Collaboration

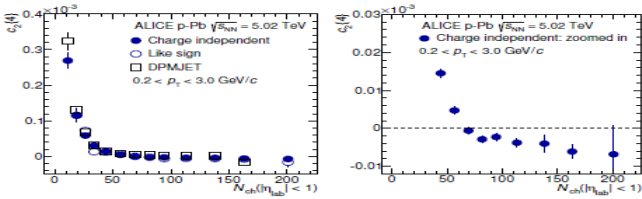


Fig. 4: Mid-rapidity ( $|\eta| < 1$ ) measurements of  $c_2\{4\}$  as a function of multiplicity for p-Pb collisions. Only statistical errors are shown as these dominate the uncertainty. See table 1 for systematic uncertainties. The right panel shows a zoomed in version of the solid points in the left panel.

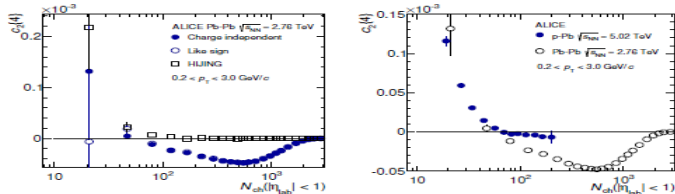


Fig. 5: Left Panel: Mid-rapidity ( $|\eta| < 1$ ) measurements of  $c_2\{4\}$  as a function of multiplicity for Pb-Pb collisions. Right Panel: Comparison of  $c_2\{4\}$  for p-Pb and Pb-Pb collisions. Only statistical errors are shown as these dominate the uncertainty. See table 1 for systematic uncertainties.

(proton + antiproton) / pions in pp and pPb v, D meson production in pp,  $\Upsilon$  (ns) /  $\Upsilon$  (ms) vs  $dN/d\eta$

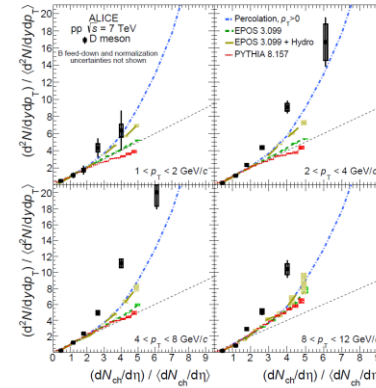
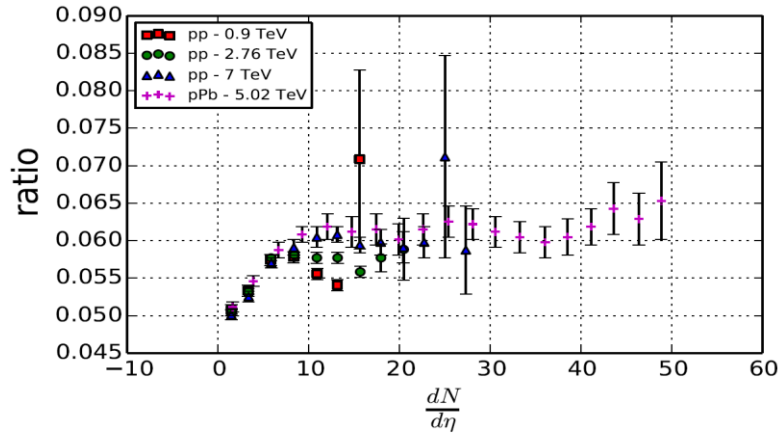
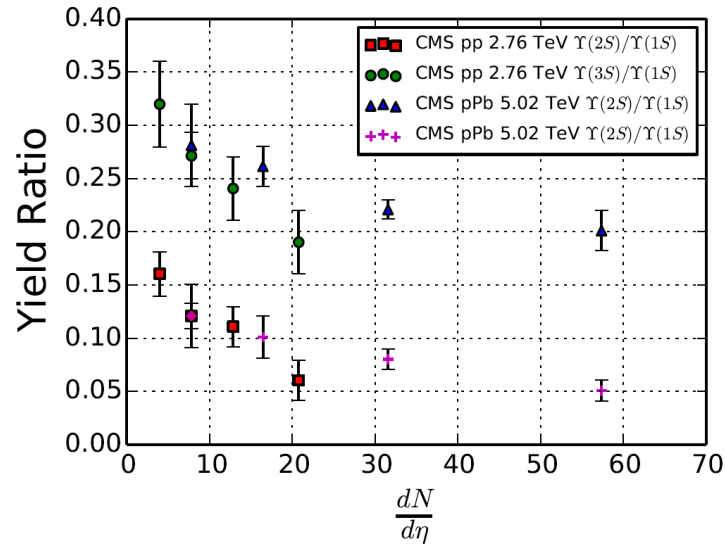


Figure 10: Average D-meson relative yield as a function of the relative charged-particle multiplicity at central rapidity in different  $p_T$  intervals. The systematic uncertainties on the data normalisation (+6%/-3%), on the  $(dN_{ch}/d\eta)/(dN_{ch}/d\eta)$  values ( $\pm 6\%$ ), and on the feed-down contribution are not shown in this figure. Different calculations are presented: PYTHIA 8.157 [30, 31], EPOS 3 with and without hydro [71, 72] and a  $p_T$ -integrated percolation model [41, 73]. The coloured lines represent the calculation curves, whereas the shaded bands represent their statistical uncertainties at given values of  $(dN_{ch}/d\eta)/(dN_{ch}/d\eta)$ . The diagonal (dashed) line is shown to guide the eye.



# Possible saturation of particle ratios and $Y(ns)/Y(ms)$ ratio at high $dn/d\eta$ . Volume effect? Does the volume saturate as well?



CERN-PH-EP-2014-264  
Submitted to: Eur. Phys. J. C

1 [hep-ex] 27 Feb 2015

Two-particle Bose–Einstein correlations in  $pp$  collisions at  $\sqrt{s} = 0.9$  and 7 TeV measured with the ATLAS detector

The ATLAS Collaboration

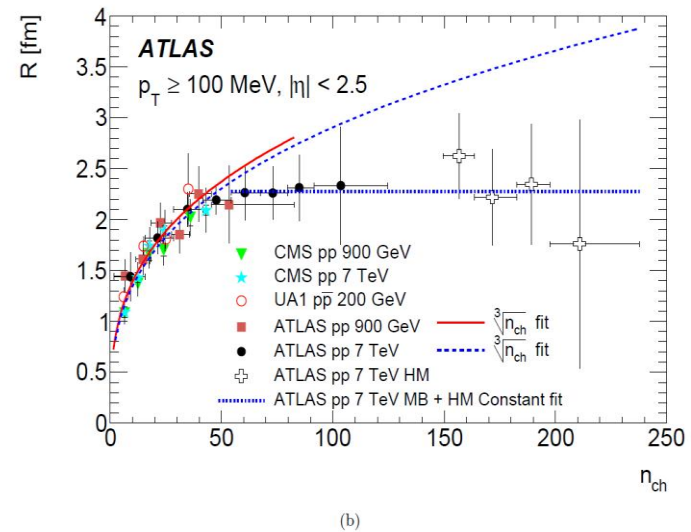


Fig. 3. Multiplicity,  $n_{ch}$ , dependence of the parameters (a)  $\lambda$  and (b)  $R$  obtained from the exponential fit to the two-particle double-ratio correlation functions  $R_2(Q)$  at  $\sqrt{s} = 0.9$  and 7 TeV, compared to the equivalent measurements of the CMS [38, 39] and UA1 [67] experiments. The solid and dashed curves are the results of (a) the exponential and (b)  $\sqrt[3]{n_{ch}}$  for  $n_{ch} < 55$  fits. The dotted line in (b) is a result of a constant fit to minimum-bias and high-multiplicity events data at 7 TeV for  $n_{ch} \geq 55$ . The error bars represent the quadratic sum of the statistical and systematic uncertainties.