

Pseudoscalar transition form factors: $(g-2)_\mu$, $P \rightarrow e^+e^-$, η - η' mixing

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Work done in collaboration with
Pablo Sanchez-Puertas



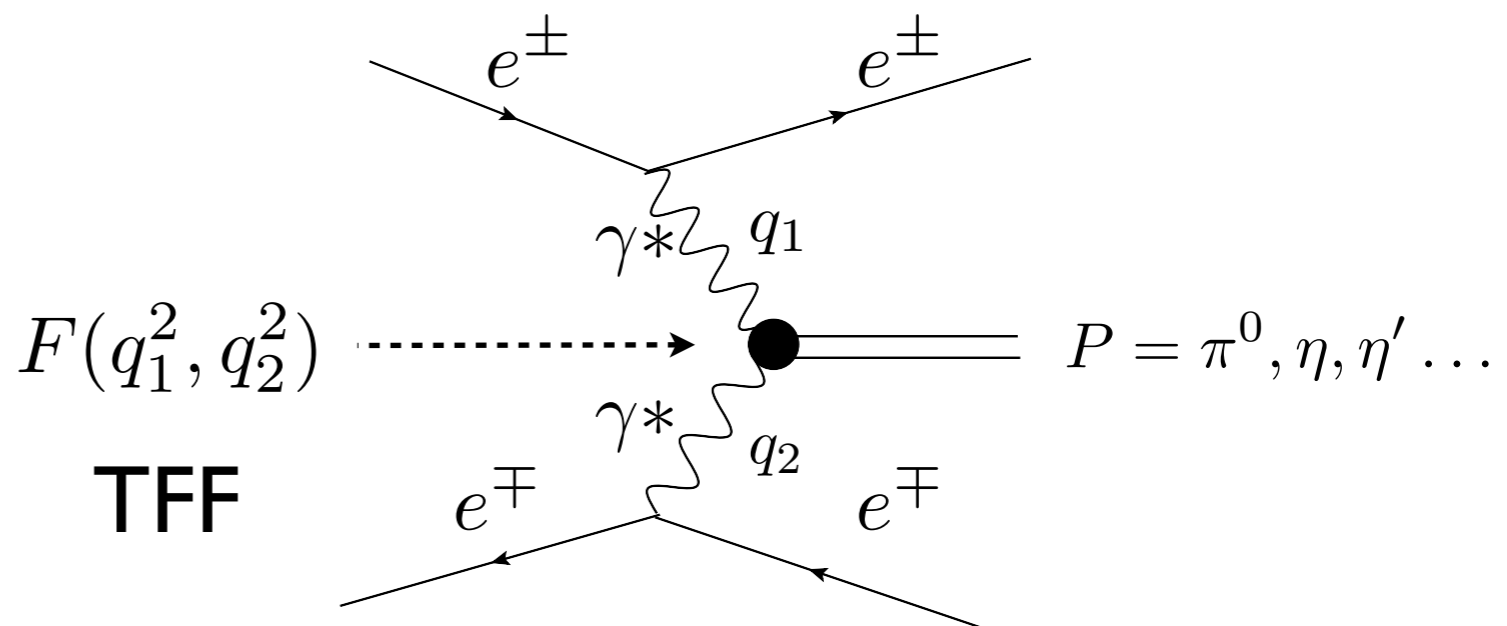
FCCP2015, Anacapri, 11th Sept.

Outline

- Pseudoscalar Transition Form Factors
- How to use data for dressing the TFFs
- Applications
 - $(g-2)_\mu$, $P \rightarrow e^+e^-$, η - η' mixing
- Conclusions and outlook

Pseudoscalar Transition Form Factors

- Study of $ee \rightarrow ee\gamma^*\gamma^*$
with $\gamma^*\gamma^* \rightarrow \pi, \eta, \eta'$
but also $P \rightarrow ee\gamma, 4e, 2e$



- Meson Structure
 - Transition Form Factors (TFF) give access to Meson Distribution Amplitudes
- Precision Tests of the Standard Model
 - Relation to mixing parameters, rare decays, and muon anomalous $(g-2)_\mu$

How do we do that?

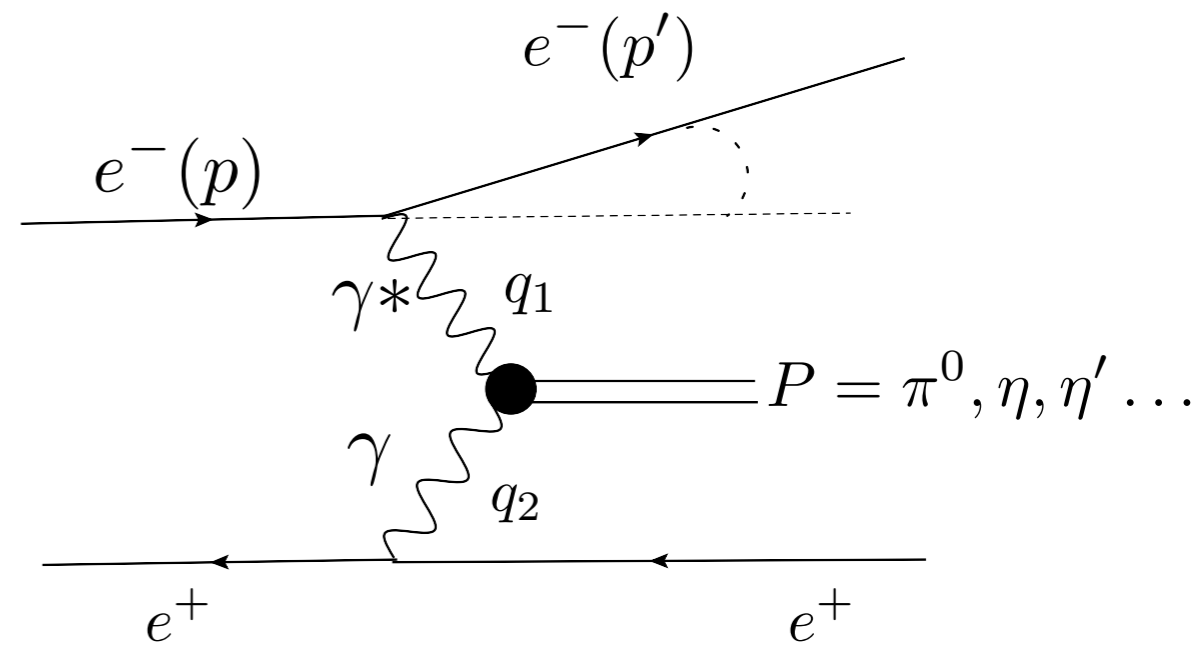
- Single Tag Method can access the Meson Transition Form Factor

Selection criteria

- 1 e^- detected
- 1 e^+ along beam axis
- Meson full reconstructed

Momentum transfer

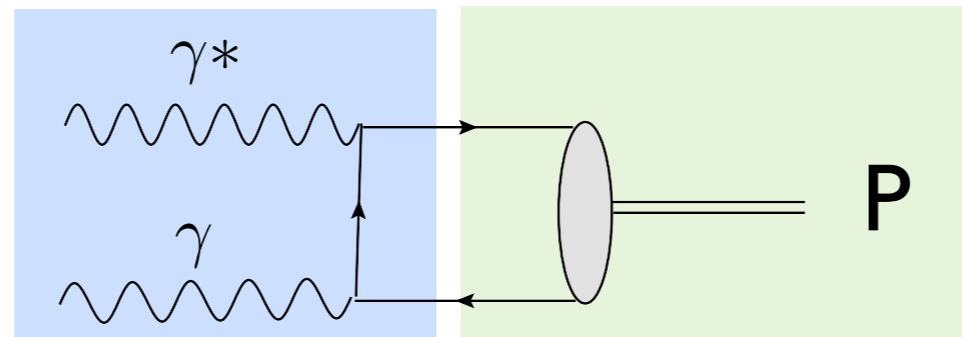
- tagged: $Q^2 = -q_1^2 = -(p - p')^2$
⇒ highly virtual photon
- untagged: $q^2 = -q_2^2 \sim 0 \text{ GeV}^2$
⇒ quasi-real photon



How do we do that?

Cross section for P production depends only on $F(q_1^2, q_2^2)$

With the Single Tag Method: $F(q_1^2, q_2^2) \rightarrow F(Q^2)$



$$F(Q^2) = \int T_H(x, Q^2) \Phi_P(x, \mu_F) dx$$

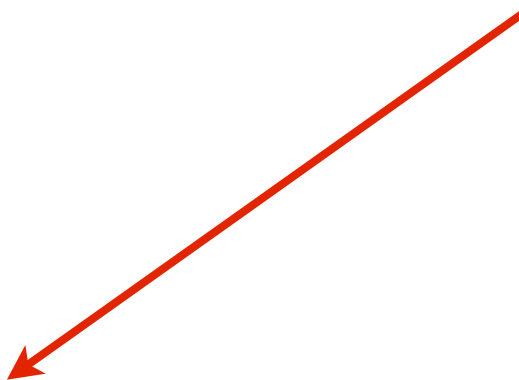
$$T_H(\gamma^* \gamma \rightarrow q\bar{q}) \quad \Phi_P(q\bar{q} \rightarrow P)$$

- μ_F is scale between soft and hard
- x-dependence of $\Phi_P(x, Q^2)$ not known but models
- Experimental data on $F(Q^2)$ is needed

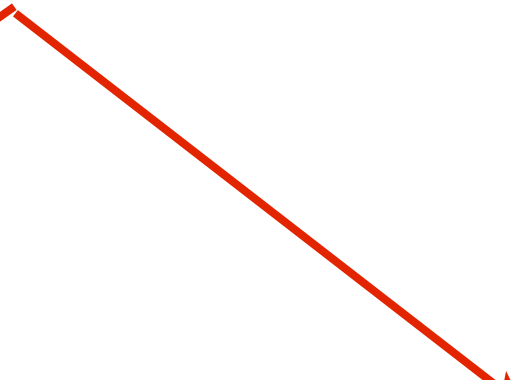
convolution of perturbative and non-perturbative regimes

The role of experimental data

$$F_{P^* \gamma^* \gamma^*}(q_3^2, q_1^2, q_2^2)$$



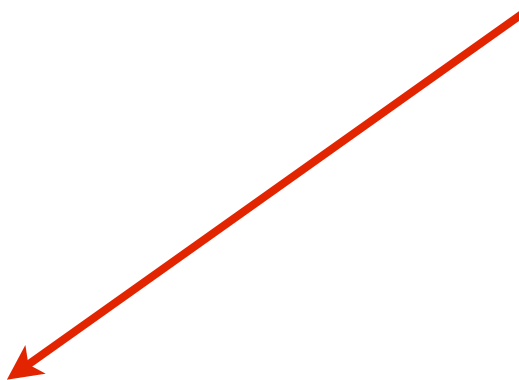
Use hadronic models
constrained with
chiral and large- N_c
arguments



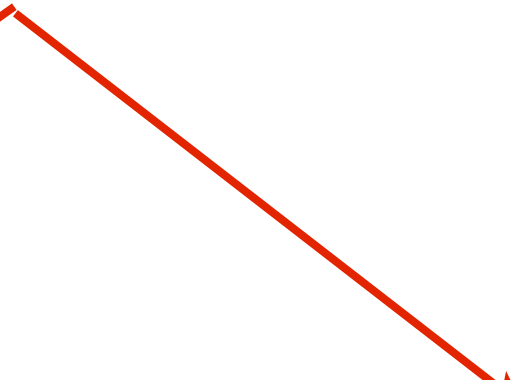
Use data from
the Transition Form Factor
for input calculations

The role of experimental data

$$F_{P^* \gamma^* \gamma^*}(q_3^2, q_1^2, q_2^2)$$



Use hadronic models
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Use data from
the Transition Form Factor
for input calculations

The role of experimental data

$$F_{P^* \gamma^* \gamma^*} (q_3^2, q_1^2, q_2^2)$$

- We want a method, not a *model*: a TOOLKIT
- based on analyticity and unitarity
 - heading towards 10% errors
- Simple (not black box)
- Approaches yes (improvable), assumptions no
- Systematic:
 - easy to update with new data
 - error from approach
- Predictive (checkable)

The role of experimental data

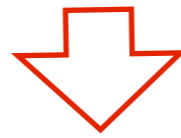
Use data from
the Transition Form Factor
for numerical integral

$$F_{P^* \gamma^* \gamma^*}(q_3^2, q_1^2, q_2^2)$$

The role of experimental data

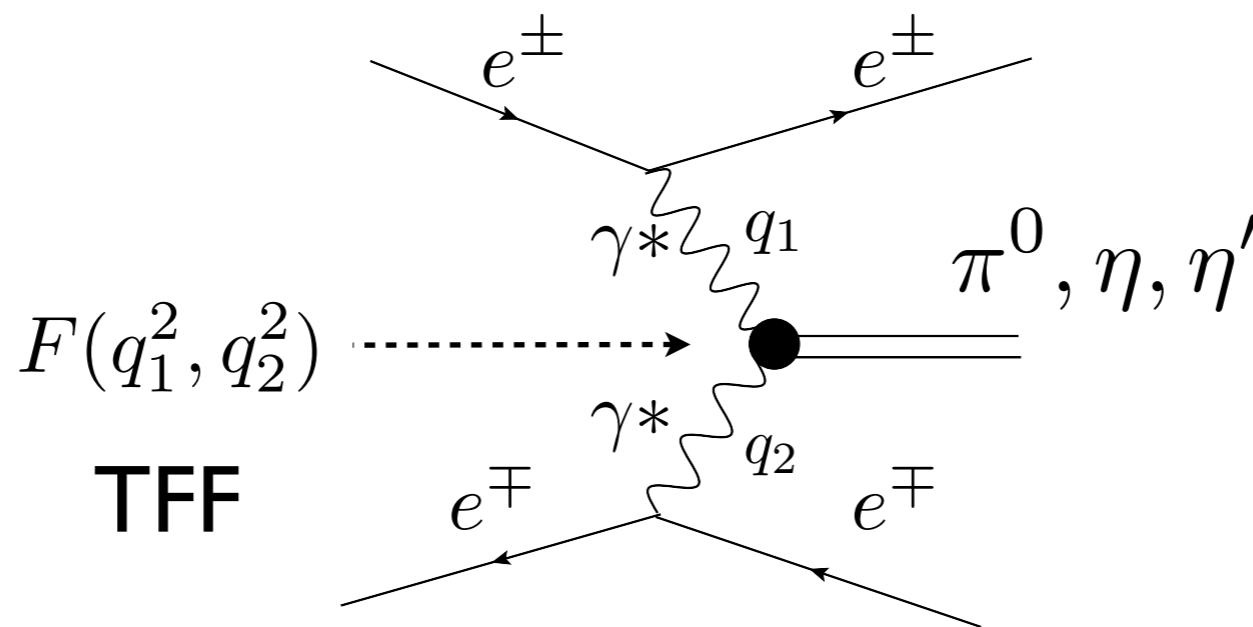
~~Use data from
the Transition Form Factor
for numerical integral~~

~~$$F_{P^* \gamma^* \gamma^*}(q_3^2, q_1^2, q_2^2)$$~~



$$F_{P \gamma^* \gamma^*}(m_P^2, q_1^2, q_2^2)$$

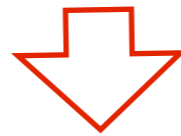
double-tag method



The role of experimental data

~~Use data from the Transition Form Factor for numerical integral~~

~~$$F_{P^* \gamma^* \gamma^*}(q_3^2, q_1^2, q_2^2)$$~~



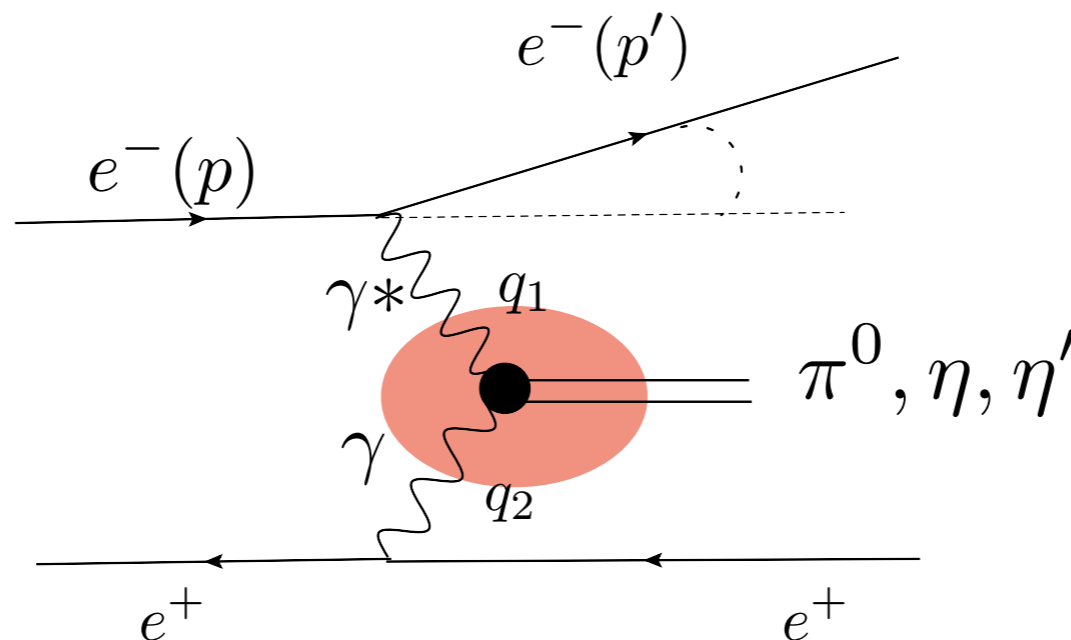
~~$$F_{P \gamma^* \gamma^*}(m_P^2, q_1^2, q_2^2)$$~~



$$F_{P \gamma^* \gamma}(m_P^2, q_1^2, 0)$$

single-tag method

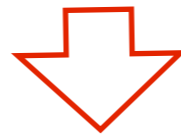
Use data from the Transition Form Factor to constrain your hadronic model



The role of experimental data

~~Use data from
the Transition Form Factor
for numerical integral~~

$$\del F_{P^* \gamma^* \gamma^*}(q_3^2, q_1^2, q_2^2)$$



$$\del F_{P \gamma^* \gamma^*}(m_P^2, q_1^2, q_2^2)$$



Use data from
the Transition Form Factor
to constrain your
hadronic model

$$F_{P \gamma^* \gamma}(m_P^2, q_1^2, 0)$$

How??

Nice synergy between experiment and theory

Simple, easy, systematic, user-friendly method

Our proposal: use Padé Approximants

[P.M.'12; P.M., M.Vanderhaeghen'12; R. Escribano, P.M., P. Sanchez-Puertas, '13, '15]

We need low-energy region (data driven) + high-energy tail
we don't want a model rather a method providing systematics

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$$F_{P\gamma^*\gamma}(Q^2, 0) = a_0^P \left(1 + b_P \frac{Q^2}{m_P^2} + c_P \frac{Q^4}{m_P^4} + \dots \right)$$

$\Gamma_{P \rightarrow \gamma\gamma}$ slope curvature

We have published space-like data for $Q^2 F_{P\gamma^*\gamma}(Q^2, 0)$

$$Q^2 F_{P\gamma^*\gamma}(Q^2, 0) = a_0 Q^2 + a_1 Q^4 + a_2 Q^6 + \dots$$

$$P_M^N(Q^2) = \frac{T_N(Q^2)}{R_M(Q^2)} = a_0 Q^2 + a_1 Q^4 + a_2 Q^6 + \dots + \mathcal{O}((Q^2)^{N+M+1})$$

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$$Q^2 F_{P\gamma^*\gamma}(Q^2, 0) = a_0 Q^2 + a_1 Q^4 + a_2 Q^6 + \dots$$

$$P_1^1(Q^2) = \frac{a_0 Q^2}{1 - a_1 Q^2} \longrightarrow \begin{aligned} P_1^N(Q^2) &= P_1^1(Q^2), P_1^2(Q^2), P_1^3(Q^2), \dots \\ P_N^N(Q^2) &= P_1^1(Q^2), P_2^2(Q^2), P_3^3(Q^2), \dots \end{aligned}$$

sequence of approximations, i.e., theoretical error

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Convergence (making use of analytical properties):

$$\lim_{N \rightarrow \infty} P_1^N(Q^2) = F_{P\gamma^*\gamma}(Q^2, 0) \quad \text{Montessus Theorem}$$

Conv. from pole at $-Q^2$ to Q^{*2} : good at LE, bad at HE. Fantastic for LEPs and cheap

Our proposal: use Padé Approximants

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Convergence (making use of analytical properties):

$$\lim_{N \rightarrow \infty} P_N^N(Q^2) = F_{P\gamma^*\gamma}(Q^2, 0) \quad \text{Pommerenke Theorem}$$

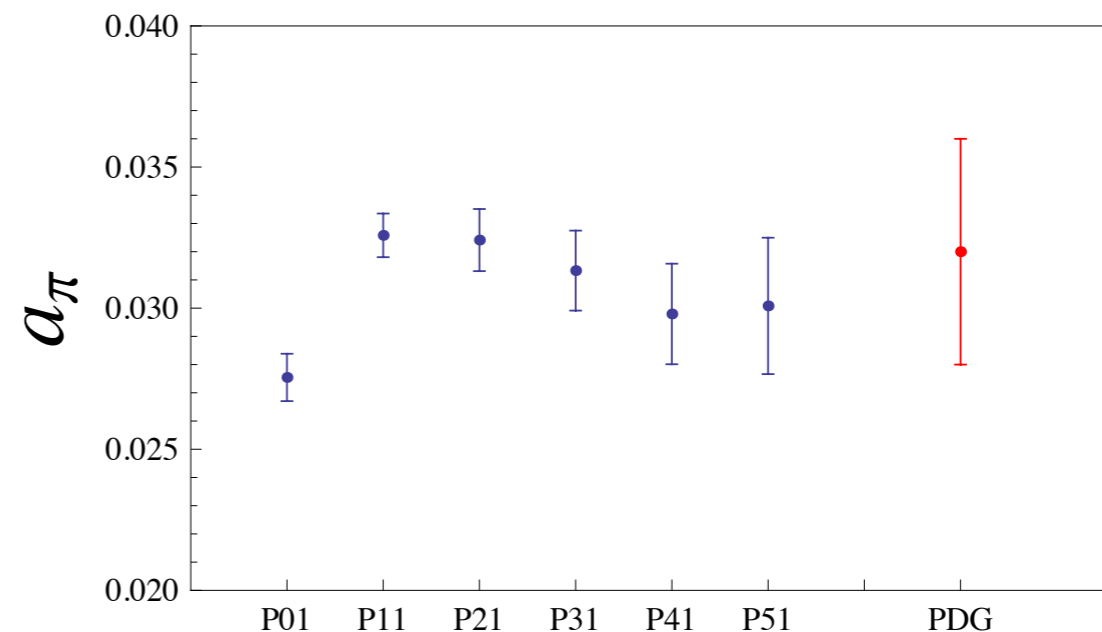
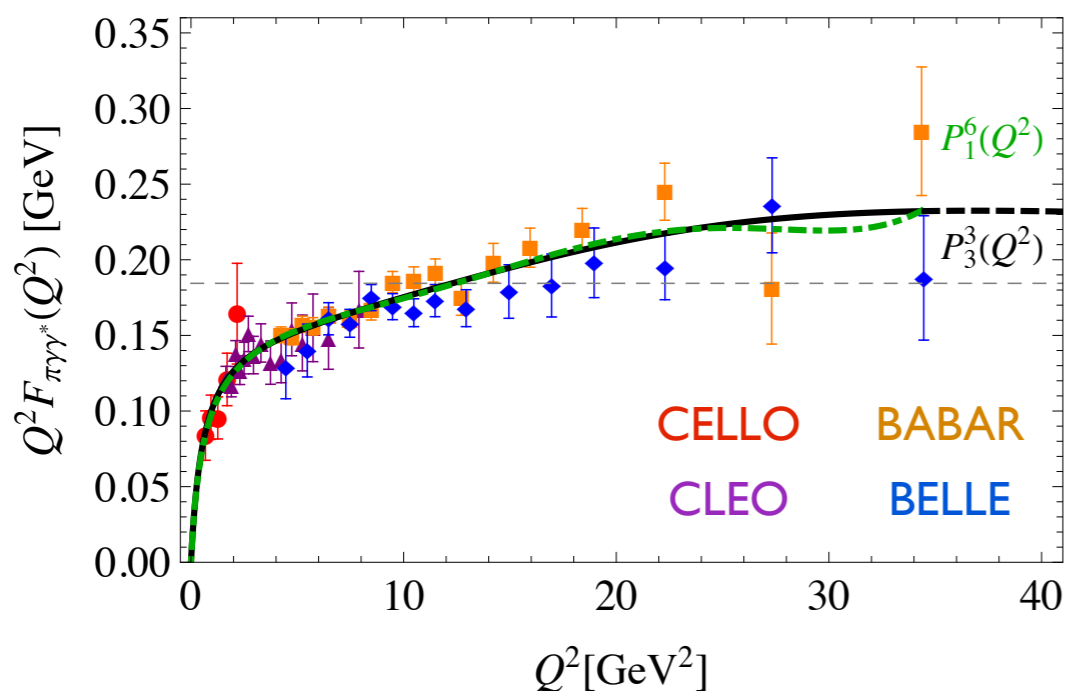
Conv. from cut at $-Q^2$ to ∞ : good at LE and HE. Good for LEPs and no cheap

Our proposal: use Padé Approximants

[P.M.'12; P.M., M.Vanderhaeghen'12; R. Escribano, P.M., P. Sanchez-Puertas, '13, '15]

Fit to Space-like data: CELLO'91, CLEO'98, BABAR'09 and Belle'12

$$P_1^N(Q^2) \quad \text{up to } N=5 \quad [\text{P.M.}, '12]$$



$$P_N^N(Q^2) \quad \text{up to } N=3$$

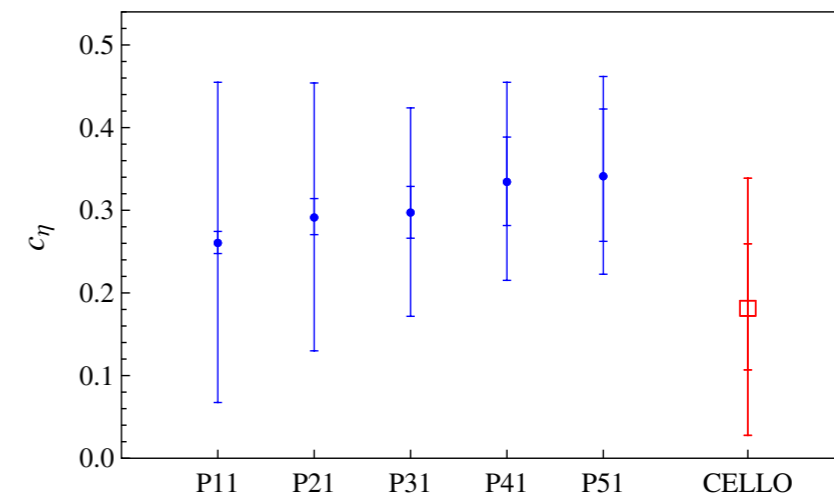
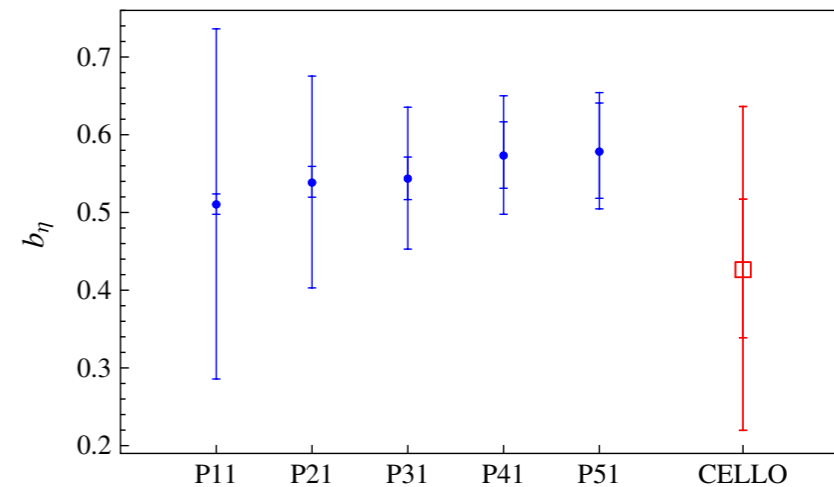
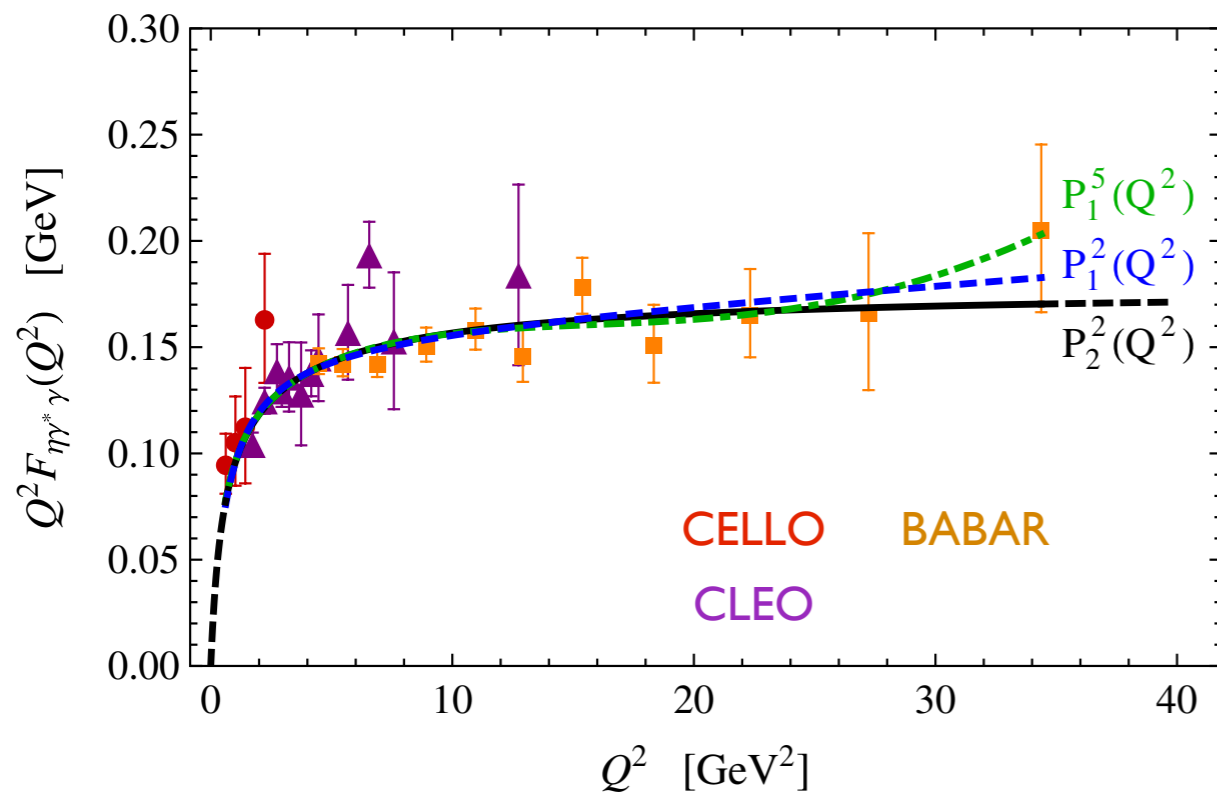
Accurate description of the low-energy region making full use of available experimental data

η -TFF

Fit to Space-like data: CELLO'91, CLEO'98, BABAR'11 + $\Gamma_{\eta \rightarrow \gamma\gamma}$

[R.Escribano, P.M., P. Sanchez-Puertas, '13]

$P_1^N(Q^2)$ up to N=4



$P_N^N(Q^2)$ up to N=2

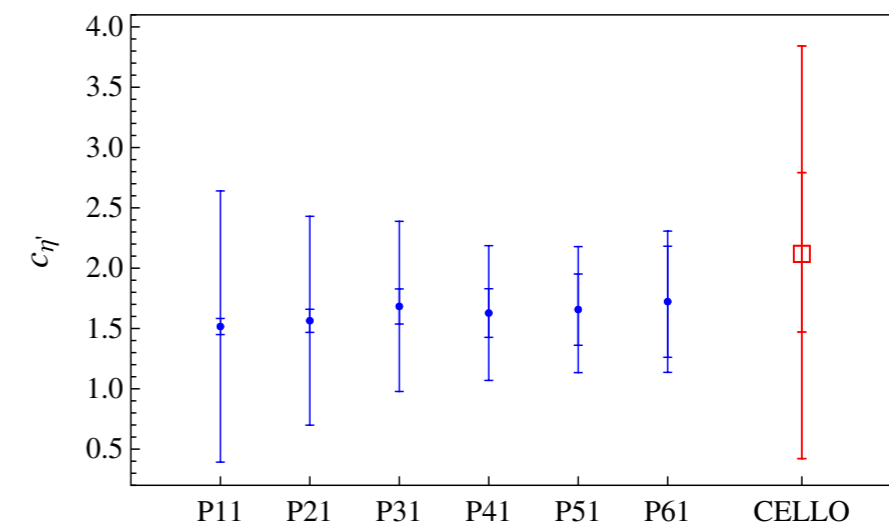
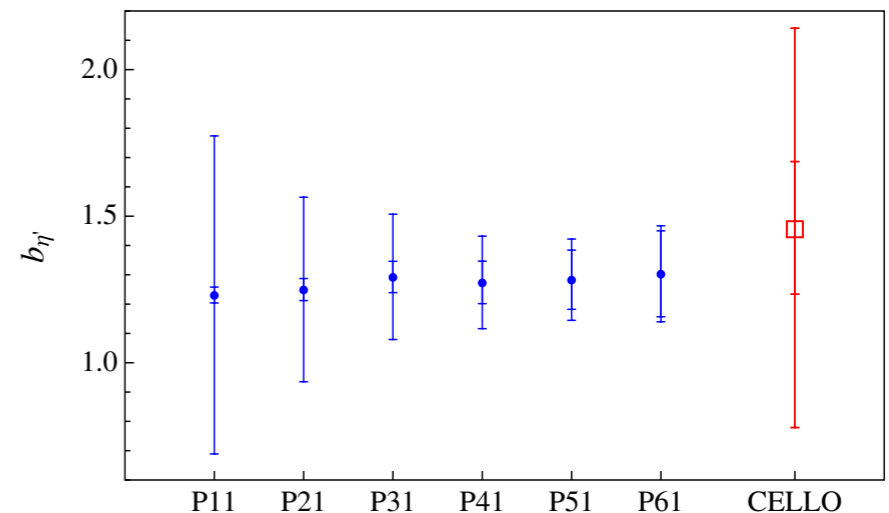
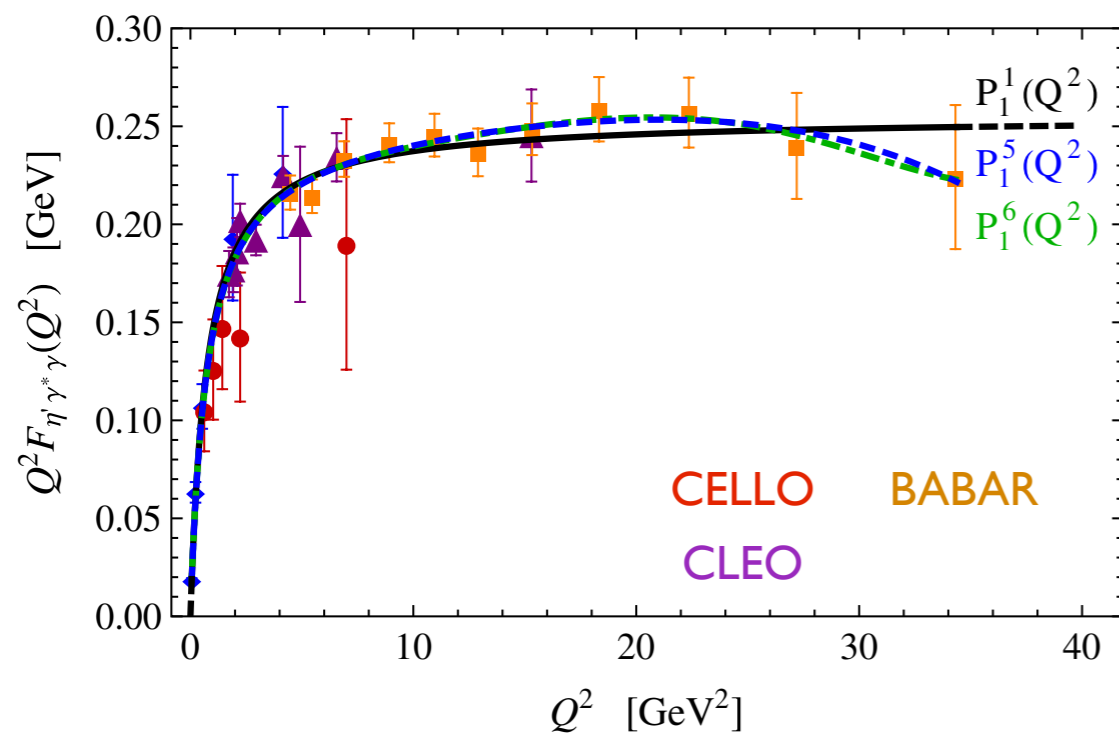
$$\lim_{Q^2 \rightarrow \infty} Q^2 F_{\eta\gamma^*\gamma}(Q^2, 0) = 0.164(2) \text{ GeV}$$

η' -TFF

Fit to Space-like data: CELLO'91, CLEO'98, L3'98, BABAR'11 + $\Gamma_{\eta' \rightarrow \gamma\gamma}$

[R.Escribano, P.M., P. Sanchez-Puertas, '13]

$P_1^N(Q^2)$ up to N=5



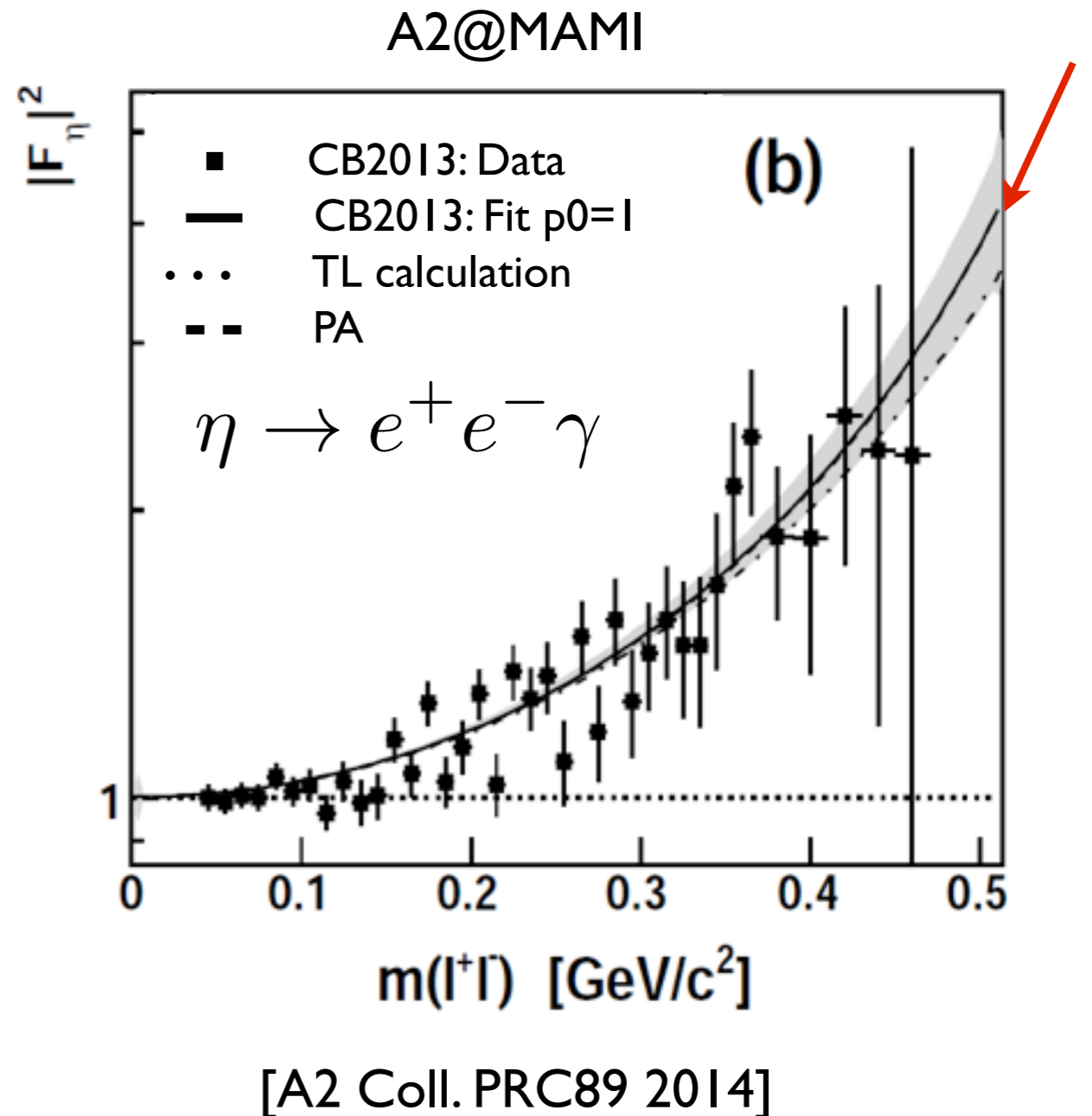
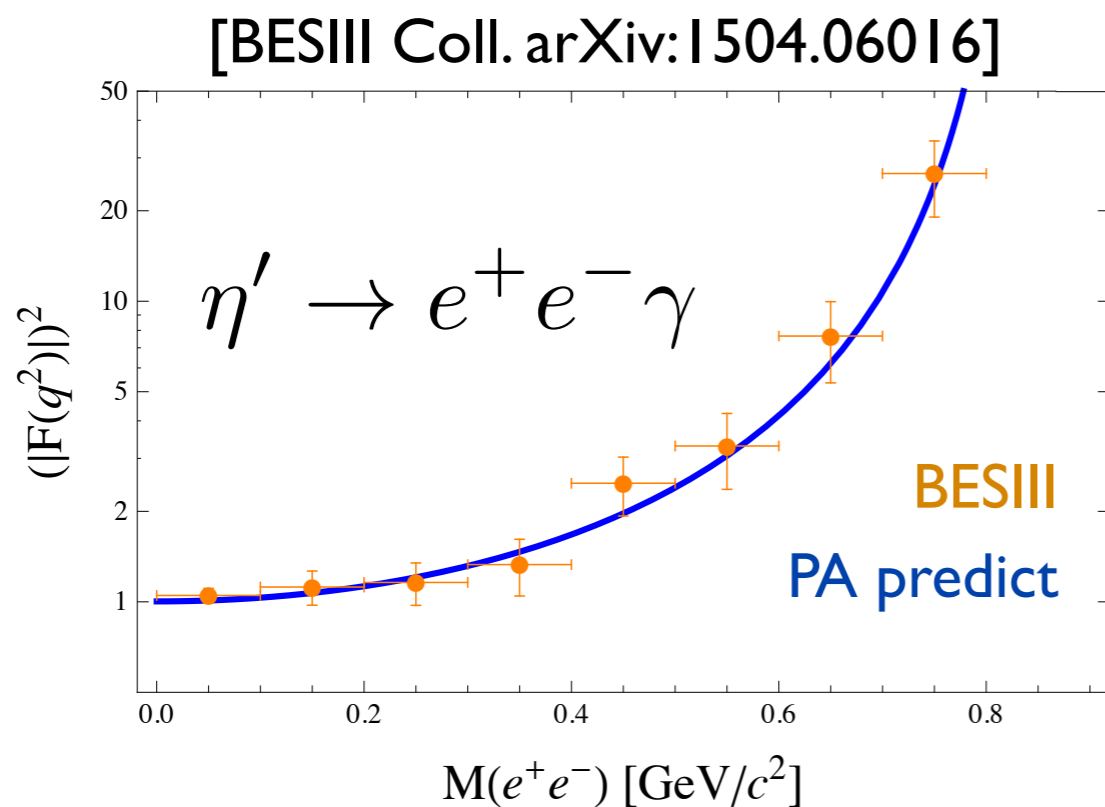
$P_N^N(Q^2)$ up to N=1

$$\lim_{Q^2 \rightarrow \infty} Q^2 F_{\eta' \gamma^* \gamma}(Q^2, 0) = 0.254(4) \text{ GeV}$$

time-like TFF

Predictive method!

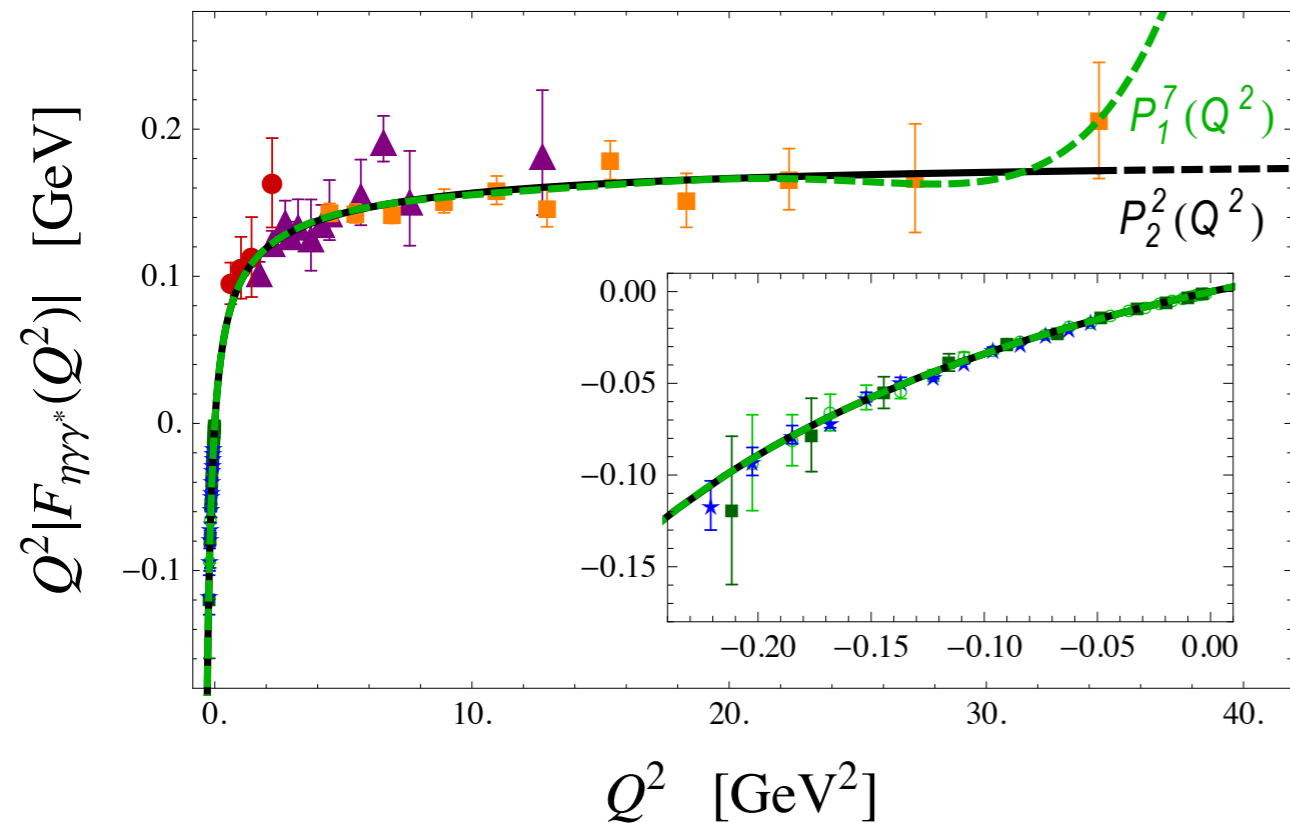
- Study Dalitz decays
 $\eta(\prime) \rightarrow \gamma^* \gamma \rightarrow e^+ e^- \gamma$
- Prediction of the time-like
from space-like data



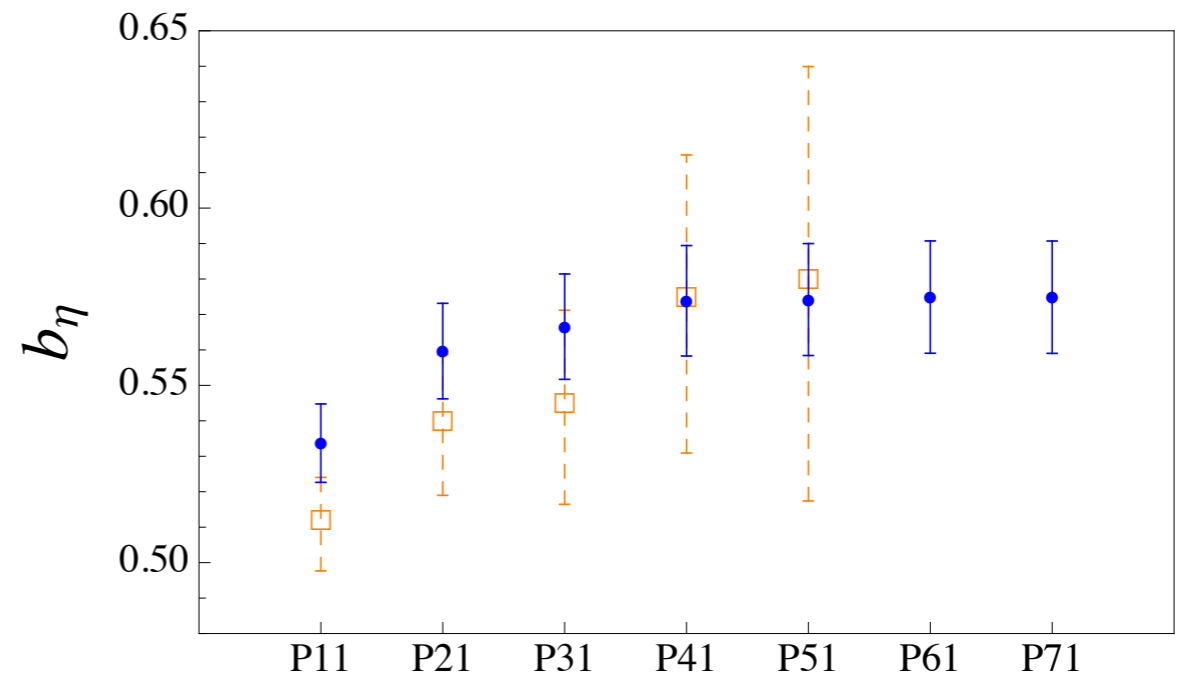
η -TFF

Fit to Space-like data [CELLO'91, CLEO'98, BABAR'11] + $\Gamma_{\eta \rightarrow \gamma\gamma}$
 + Time-like data [NA60'09, A2'11, A2'13]

[R.Escribano, P.M., P. Sanchez-Puertas, '15]



$P_1^N(Q^2)$ up to N=7



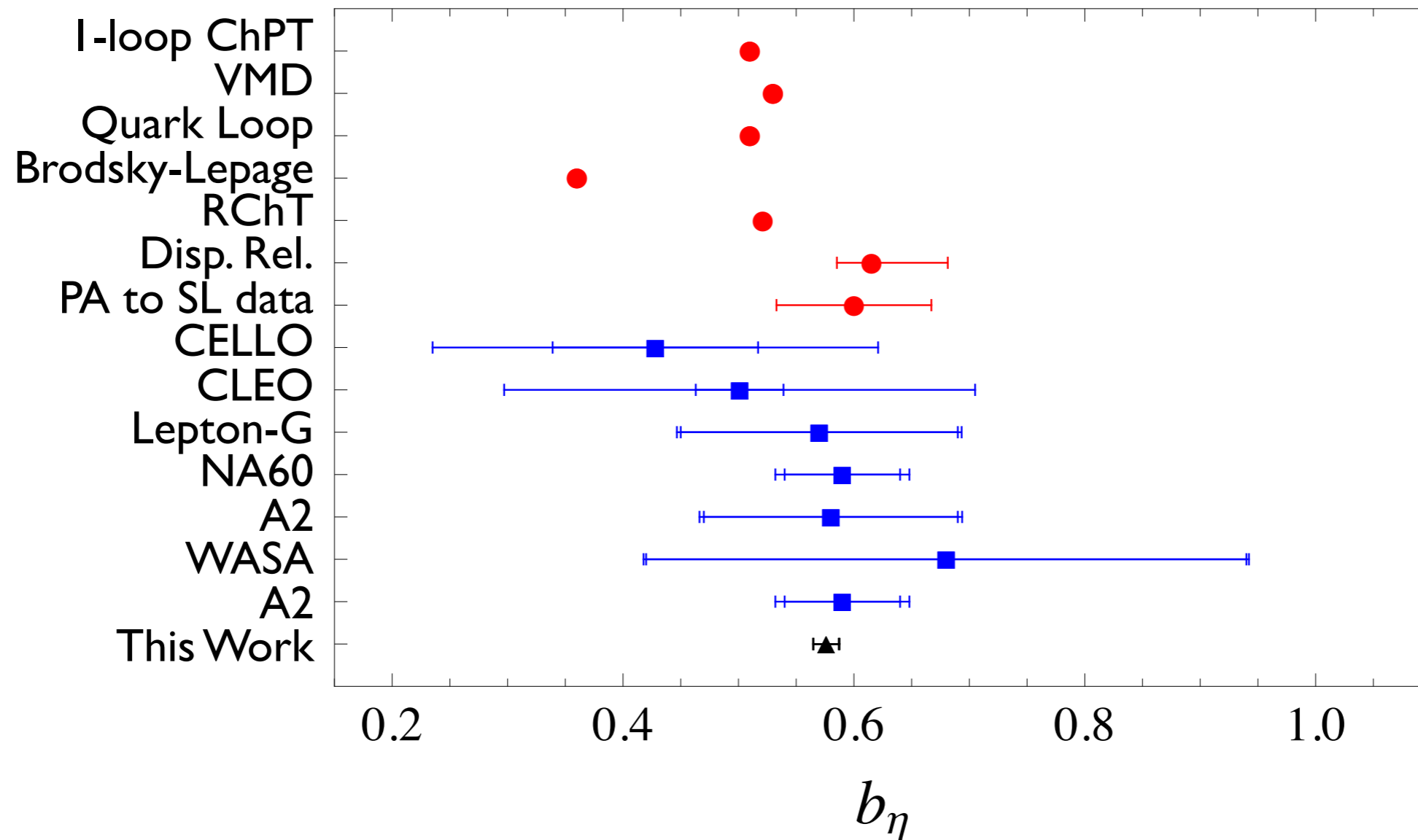
$P_N^N(Q^2)$ up to N=2

$$\lim_{Q^2 \rightarrow \infty} Q^2 F_{\eta\gamma^*\gamma}(Q^2, 0) = 0.177(15) \text{ GeV}$$

η -TFF

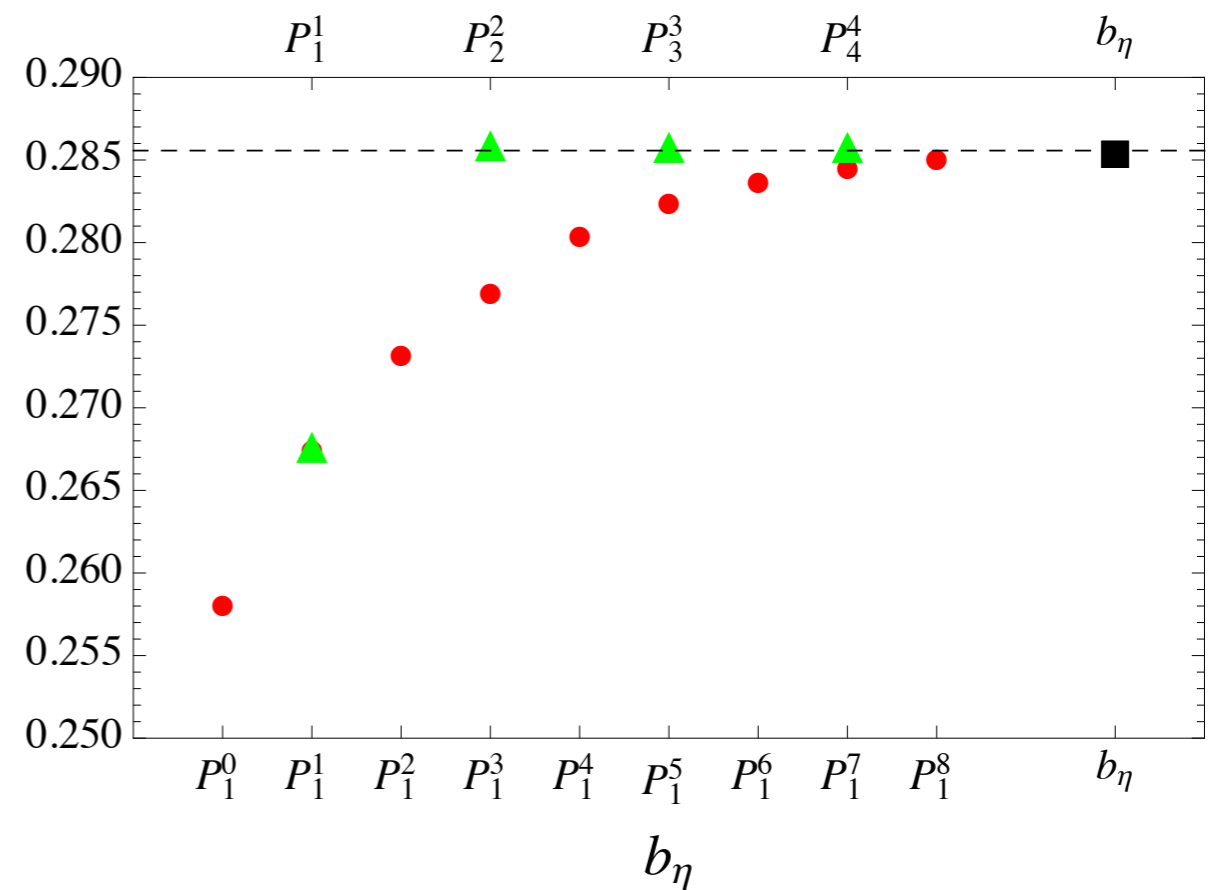
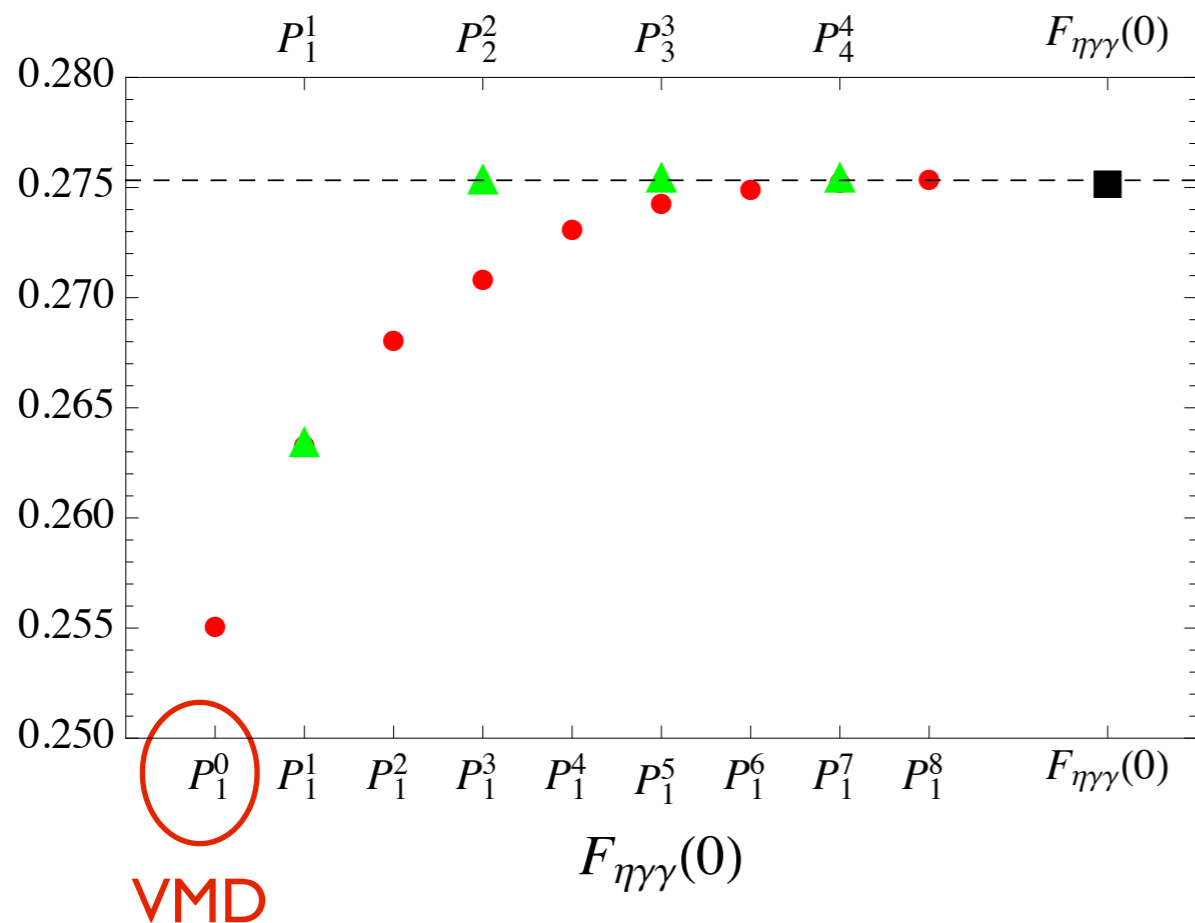
Fit to Space-like data [CELLO'91, CLEO'98, BABAR'11] + $\Gamma_{\eta \rightarrow \gamma\gamma}$
+ Time-like data [NA60'09, A2'11, A2'13]

[R.Escribano, P.M., P. Sanchez-Puertas, '15]



A word on systematics

- Consider a model for η TFF
- Generate a pseudodata set emulating the physical situation (SL+TL)
- Build up your PA sequence
- Fit and compare

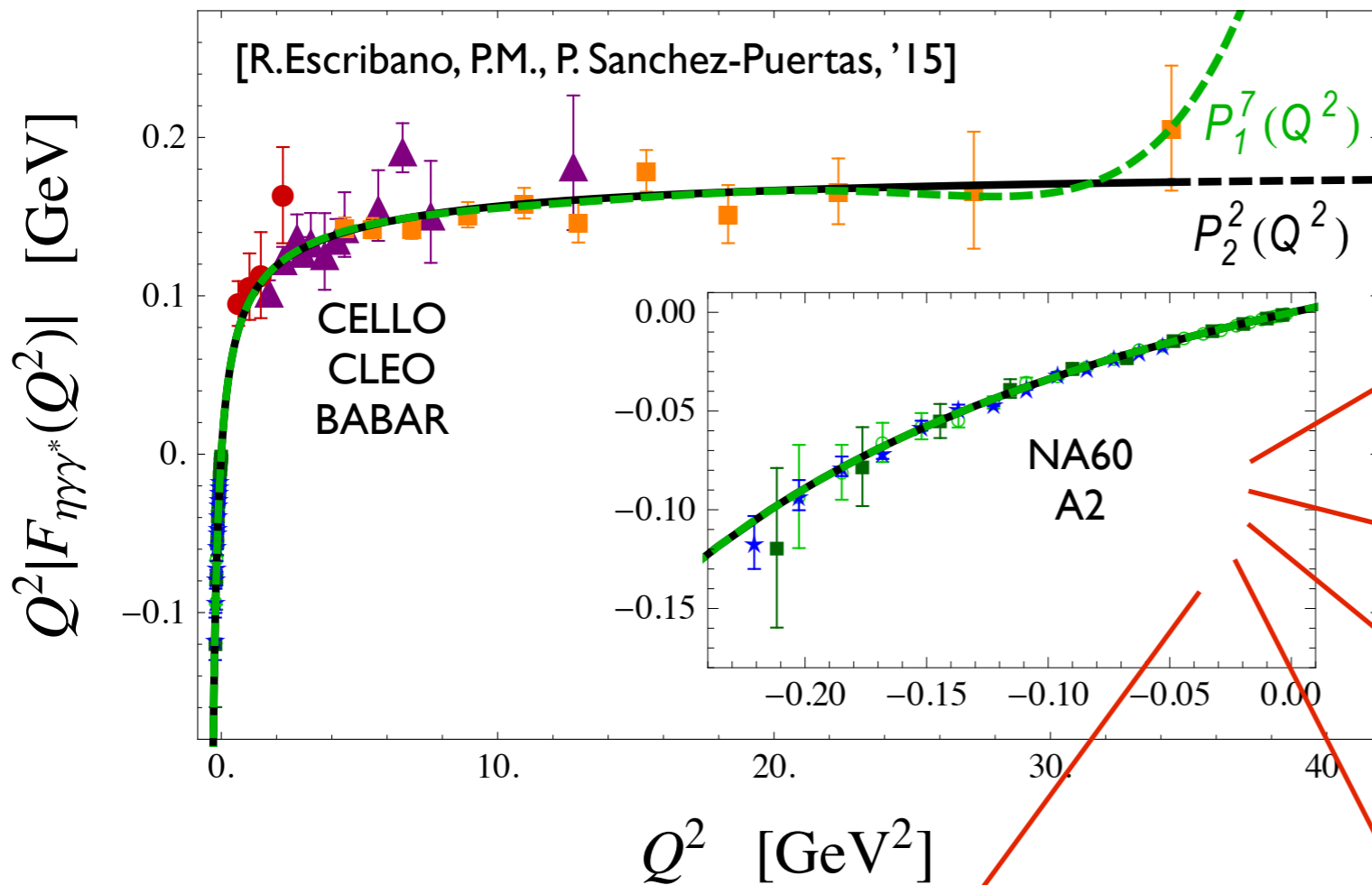


Why PAs?

- To THEORY:
 - our problem is “the general rational Hermite interpolation problem”: solution is PAs
 - in many cases we need $\int ds f(s)$ not $f(s)$
 - VMD, LMD, LMD+V, HQCD... are already a certain kind of PAs (PTs)
 - it is proven: PTs slower convergence than PAs
- To DISPERSION Theory:
 - excellent interpolation tool: bring DRs (TL!) to ∞ (link with pQCD)
 - even more: PA can be the function to be used for the DR
 - or DR can provide LECs to feed in the PA
- To LATTICE:
 - used recently for lattice fitting of HVP: suggestion use it correctly!
- To EXPERIMENT:
 - what energy and precision to measure: use PAs to identify!
(high energies are as well important vs DRs)

PS-TFF

space-like and time-like data



Low-energy parameters
up to the third derivative!

$$(g-2)_\mu$$

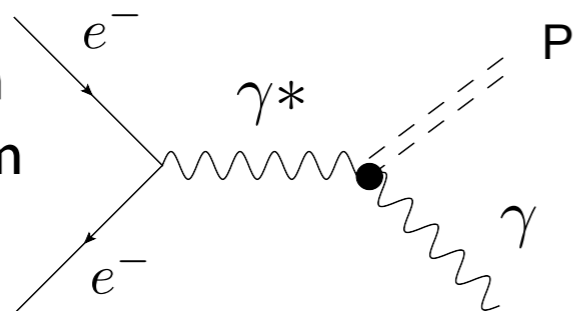
mixing parameters
of the η - η' system

$$f_q, f_s, \phi$$

Rare decays

$$\Gamma_{P \rightarrow l^+ l^-} \quad \Gamma_{P \rightarrow 4l}$$

PS continuum
on quarkonium
region



beautiful synergy experiment - theory

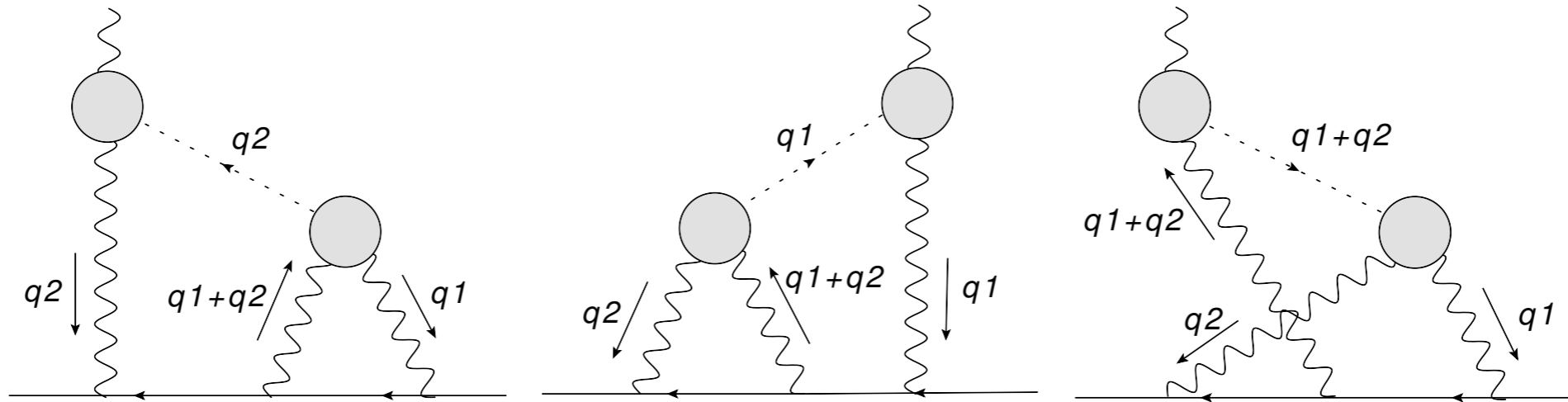
Applications

1. Hadronic Light-by-Light contribution to muon ($g-2$)
2. PS decays into lepton pairs ($\pi^0 \rightarrow e^+e^-$)
3. η - η' mixing

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Dissection of the HLbL contribution

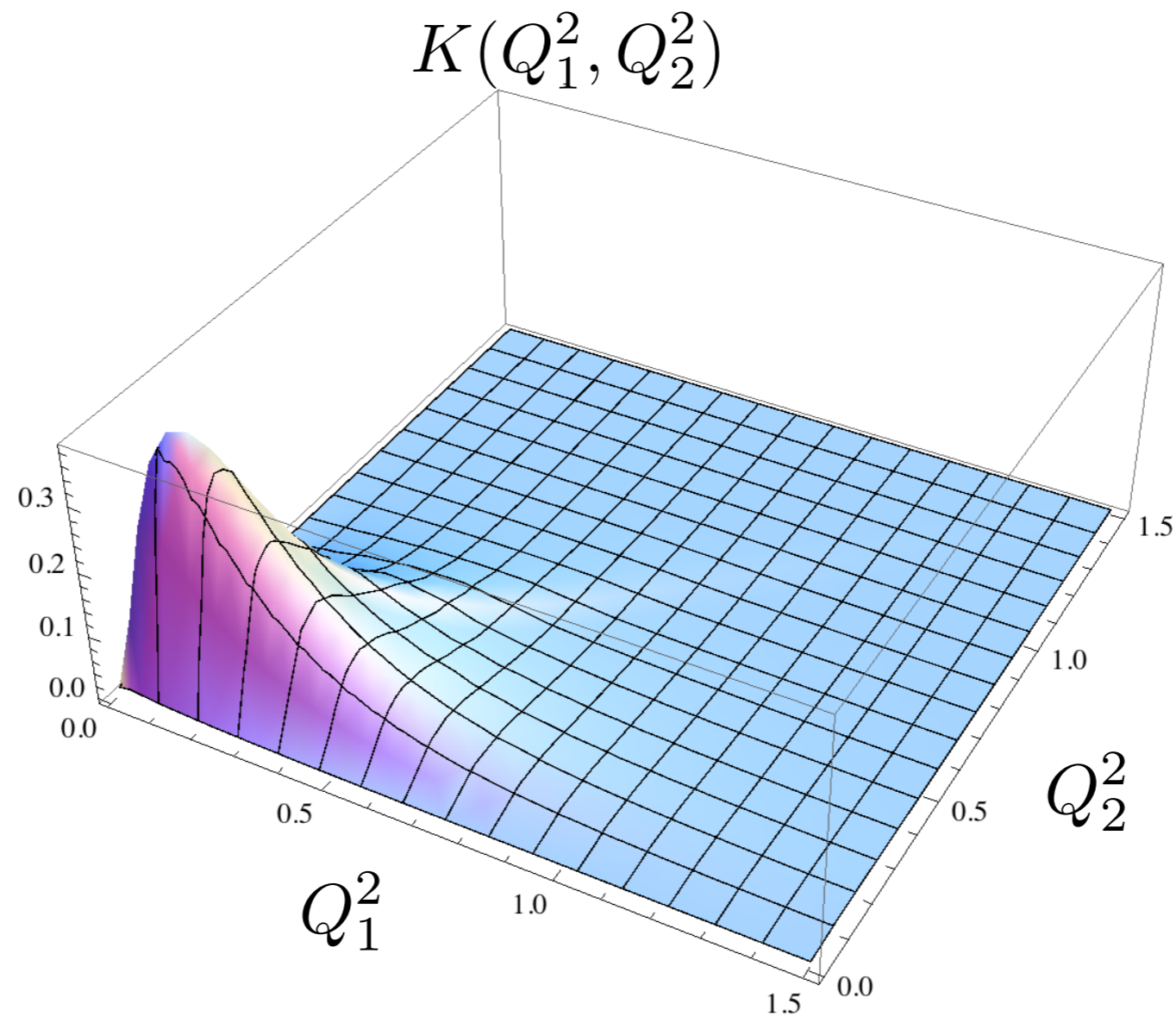


$$\begin{aligned}
 a_{\mu}^{LbL;P} = & -e^6 \int \frac{d^4 q_1}{(2\pi)^4} \int \frac{d^4 q_2}{(2\pi)^4} \frac{1}{q_1^2 q_2^2 (q_1 + q_2)^2 [(p + q_1)^2 - m^2][(p - q_2)^2 - m^2]} \\
 & \times \left(\frac{F_{P^* \gamma^* \gamma^*}(q_2^2, q_1^2, (q_1 + q_2)^2) F_{P^* \gamma^* \gamma^*}(q_2^2, q_2^2, 0)}{q_2^2 - M_P^2} T_1(q_1, q_2; p) \right. \\
 & \left. + \frac{F_{P^* \gamma^* \gamma^*}((q_1 + q_2)^2, q_1^2, q_2^2) F_{P^* \gamma^* \gamma^*}((q_1 + q_2)^2, (q_1 + q_2)^2, 0)}{(q_1 + q_2)^2 - M_P^2} T_2(q_1, q_2; p) \right)
 \end{aligned}$$

Dissection of the HLBL contribution

$$a_{\mu}^{LbyL;\pi^0} = e^6 \int \frac{d^4 Q_1}{(2\pi)^4} \int \frac{d^4 Q_2}{(2\pi)^4} K(Q_1^2, Q_2^2)$$

Using $F_{\pi^0\gamma^*\gamma^*}(Q_1^2, Q_2^2) \sim VMD(Q_1^2, Q_2^2)$



Dissection of the HLBL contribution

a la Melnikov-Vainshtein

Central value:

[Knecht-Nyffeler '01]

$$F_{\pi^0 \gamma^* \gamma^*}^{LMD+V}(q_1^2, q_2^2) = \frac{f_\pi}{3} \frac{q_1^2 q_2^2 (q_1^2 + q_2^2) + h_1 (q_1^2 + q_2^2)^2 + h_2 q_1^2 q_2^2 + h_5 (q_1^2 + q_2^2) + h_7}{(q_1^2 - M_{V_1}^2)(q_1^2 - M_{V_2}^2)(q_2^2 - M_{V_1}^2)(q_2^2 - M_{V_2}^2)}$$

Publication:

$$F_\pi = 92.4 \text{ MeV}$$

$$m_\rho = 769 \text{ MeV}$$

$$m_{\rho'} = 1465 \text{ MeV}$$

$$h_1 = 0 \text{ (BL limit)}$$

$$h_5 = 6.93 \text{ GeV}^4$$

$$h_2 = -10 \text{ GeV}^2$$

$$a_\mu^{\text{HLBL}, \pi} = 7.7 \times 10^{-10}$$

Dissection of the HLBL contribution

a la Melnikov-Vainshtein

Central value:

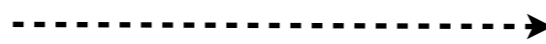
[Knecht-Nyffeler '01]

$$F_{\pi^0 \gamma^* \gamma^*}^{LMD+V}(q_1^2, q_2^2) = \frac{f_\pi}{3} \frac{q_1^2 q_2^2 (q_1^2 + q_2^2) + h_1 (q_1^2 + q_2^2)^2 + h_2 q_1^2 q_2^2 + h_5 (q_1^2 + q_2^2) + h_7}{(q_1^2 - M_{V_1}^2)(q_1^2 - M_{V_2}^2)(q_2^2 - M_{V_1}^2)(q_2^2 - M_{V_2}^2)}$$

Publication:

Preliminary, using exp data

$$F_\pi = 92.4 \text{ MeV}$$



$$\Gamma_{\pi^0 \rightarrow \gamma\gamma}$$

$$m_\rho = 769 \text{ MeV}$$

$$m_{\rho'} = 1465 \text{ MeV}$$



$$m_\rho = 775 \text{ MeV}$$

$$h_1 = 0 \text{ (BL limit)}$$

$$h_5 = 6.93 \text{ GeV}^4$$



curvature

$$h_1 = 0 \text{ (BL limit)}$$

$$h_2 = -10 \text{ GeV}^2$$

slope

$$h_2 = -10 \text{ GeV}^2$$

$$a_\mu^{\text{HLBL}, \pi} = 7.7 \times 10^{-10}$$



$$a_\mu^{\text{HLBL}, \pi} = 9.8 \times 10^{-10}$$

Dissection of the HLBL contribution

a la Knecht-Nyffeler '01

Central value:

$$F_{\pi^0 \gamma^* \gamma^*}^{LMD+V}(q_1^2, q_2^2) = \frac{f_\pi}{3} \frac{q_1^2 q_2^2 (q_1^2 + q_2^2) + h_1 (q_1^2 + q_2^2)^2 + h_2 q_1^2 q_2^2 + h_5 (q_1^2 + q_2^2) + h_7}{(q_1^2 - M_{V_1}^2)(q_1^2 - M_{V_2}^2)(q_2^2 - M_{V_1}^2)(q_2^2 - M_{V_2}^2)}$$

Publication:

$$F_\pi = 92.4 \text{ MeV}$$

$$m_\rho = 769 \text{ MeV}$$

$$m_{\rho'} = 1465 \text{ MeV}$$

$$h_1 = 0 \text{ (BL limit)}$$

$$h_5 = 6.93 \text{ GeV}^4$$

$$h_2 = -10 \text{ GeV}^2$$

$$a_\mu^{\text{HLBL}, \pi} = 6.3 \times 10^{-10}$$

Dissection of the HLBL contribution

a la Knecht-Nyffeler '01

Central value:

$$F_{\pi^0 \gamma^* \gamma^*}^{LMD+V}(q_1^2, q_2^2) = \frac{f_\pi q_1^2 q_2^2 (q_1^2 + q_2^2) + h_1 (q_1^2 + q_2^2)^2 + h_2 q_1^2 q_2^2 + h_5 (q_1^2 + q_2^2) + h_7}{3 (q_1^2 - M_{V_1}^2)(q_1^2 - M_{V_2}^2)(q_2^2 - M_{V_1}^2)(q_2^2 - M_{V_2}^2)}$$

Publication:

$$F_\pi = 92.4 \text{ MeV}$$

$$m_\rho = 769 \text{ MeV}$$

$$m_{\rho'} = 1465 \text{ MeV}$$

$$h_1 = 0 \text{ (BL limit)}$$

$$h_5 = 6.93 \text{ GeV}^4$$

$$h_2 = -10 \text{ GeV}^2$$

Preliminary, using exp data:

$$\Gamma_{\pi^0 \rightarrow \gamma\gamma}$$

$$m_\rho = 775 \text{ MeV}$$

curvature

$$h_1 = 0 \text{ (BL limit)}$$

slope

$$h_2 = -10 \text{ GeV}^2$$

$$a_\mu^{\text{HLBL},\pi} = 6.3 \times 10^{-10}$$

$$a_\mu^{\text{HLBL},\pi} = 7.5 \times 10^{-10}$$

Dissection of the HLBL contribution

a la Padé

P.M., S. Peris, 07 P.M.'12

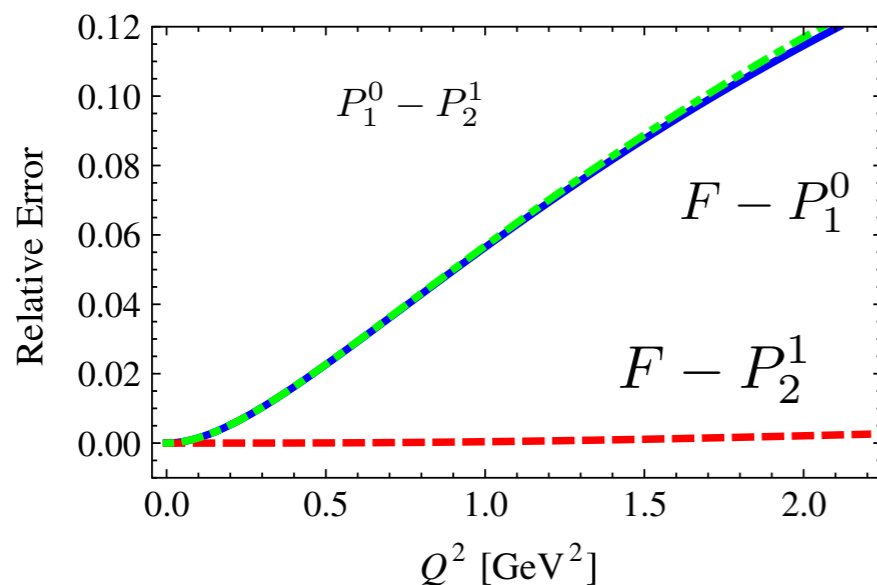
P.M., Vanderhaeghen'12

R. Escribano, P.M., P. Sanchez-Puertas, '13

$$F_{P^*\gamma^*\gamma^*}^{P01}(P_P^2, Q_1^2, Q_2^2) = a \frac{b}{Q_1^2 + b} \frac{b}{Q_2^2 + b} (1 + cP_P^2)$$

$$F_{P^*\gamma^*\gamma^*}^{P12}(P_P^2, Q_1^2, Q_2^2) = \frac{a + bQ_1^2}{(Q_1^2 + c)(Q_1^2 + d)} \frac{a + bQ_2^2}{(Q_2^2 + c)(Q_2^2 + d)} (1 + cP_P^2)$$

	b_P	c_P	$\lim_{Q^2 \rightarrow \infty} Q^2 F_{P\gamma^*\gamma}(Q^2)$	$a_\mu^{\text{HLBL;P}}$
π^0	0.0324(22)	$1.06(27) \cdot 10^{-3}$	$2f_\pi$	$6.49(56) \cdot 10^{-10}$
η	0.60(7)	0.37(12)	0.160(24)GeV	$1.25(15) \cdot 10^{-10}$
η'	1.30(17)	1.72(58)	0.255(4)GeV	$1.27(19) \cdot 10^{-10}$



Systematic error from approach:

$$P_1^0(Q_1^2, Q_2^2) \text{ vs } P_2^1(Q_1^2, Q_2^2) \longrightarrow \boxed{5\%}$$

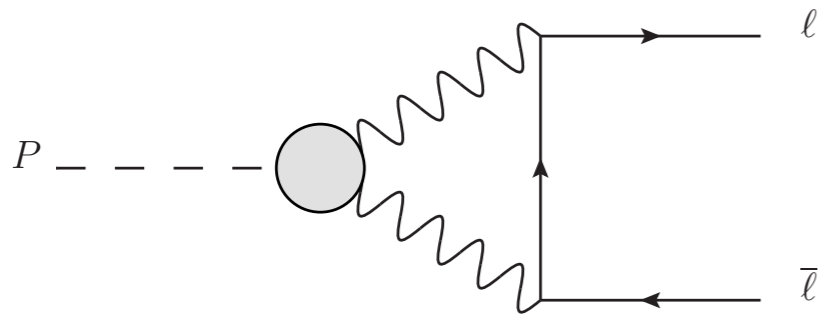
[P.M.,Peris,'07]

Applications

1. Hadronic Light-by-Light contribution to muon ($g-2$)
- 2. P decays into lepton pairs ($\pi^0 \rightarrow e^+ e^-$)**
3. η - η' mixing

Introduction and Motivation

Experiment



$$\frac{BR(P \rightarrow \bar{l}l)}{BR(P \rightarrow \gamma\gamma)} = 2 \left(\frac{\alpha m_\ell}{\pi m_P} \right)^2 \beta_\ell(m_P^2) |\mathcal{A}(m_P^2)|^2$$

$$\sim 1.5 \cdot 10^{-10}$$

KTeV '07:

$$BR(\pi^0 \rightarrow e^+e^-(\gamma), x > 0.95) = (6.44 \pm 0.25 \pm 0.22) \times 10^{-8}$$

Extrapolation to $x=1$ + radiative correction + Dalitz decay background

$$BR_{\text{KTeV}}^{w/o rad}(\pi^0 \rightarrow e^+e^-) = (7.48 \pm 0.29 \pm 0.25) \times 10^{-8}$$

(dominates de PDG)

Prague contribution: Radiative corrections

Vasko, Novotny '11 + Husek, Kampf, Novotny'14

$$\frac{\text{BR}(\pi^0 \rightarrow e^+e^-(\gamma), x > 0.95)}{\text{BR}(\pi^0 \rightarrow \gamma\gamma)} = \frac{\Gamma(\pi^0 \rightarrow e^+e^-)}{\Gamma(\pi^0 \rightarrow \gamma\gamma)} [1 + \delta^{(2)}(0.95) + \Delta^{BS}(0.95) + \delta^D(0.95)]$$

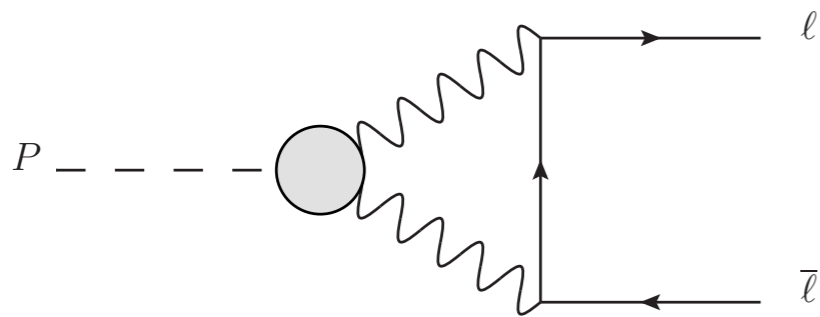
$$\delta^{(2)}(0.95) \equiv \delta^{\text{virt.}} + \delta_{\text{soft}}^{\text{BS}}(0.95) = (-5.8 \pm 0.2) \% \quad \text{vs} \quad \sim -13\%$$

$$\Delta^{\text{BS}}(0.95) = (0.30 \pm 0.01) \% \quad \delta^D(0.95) = \frac{1.75 \times 10^{-15}}{[\Gamma^{\text{LO}}(\pi^0 \rightarrow e^+e^-)/\text{MeV}]}$$

$$BR_{\text{"KTeV"}}^{w/o rad}(\pi^0 \rightarrow e^+e^-) = (6.87 \pm 0.36) \times 10^{-8}$$

Introduction and Motivation

Theory



$$\frac{BR(P \rightarrow \bar{l}l)}{BR(P \rightarrow \gamma\gamma)} = 2 \left(\frac{\alpha m_\ell}{\pi m_P} \right)^2 \beta_\ell(m_P^2) |\mathcal{A}(m_P^2)|^2$$

The only unknown $\mathcal{A}(m_P^2)$ from loop calculation where the TFF enters.

$$\mathcal{A}(q^2) = \frac{2i}{\pi^2} \int d^4k \frac{q^2 k^2 - (k \cdot q)^2}{k^2 (k - q)^2 ((p - k) - m_\ell^2)} \frac{F_{P\gamma^*\gamma^*}(k^2, (q - k)^2)}{F_{P\gamma\gamma}(0, 0)}$$

Dissection of $\pi^0 \rightarrow e^+ e^-$

As model independent as possible:

Cutcosky rules provides the imaginary part

$$\text{Im}\mathcal{A}(q^2) = \frac{\pi}{2\beta_I(q^2)} \ln \left(\frac{1 - \beta_I(q^2)}{1 + \beta_I(q^2)} \right); \quad \beta_I(q^2) = \sqrt{1 - \frac{4m_l^2}{q^2}}$$

$q^2 = m_P^2$

Assuming $|\mathcal{A}|^2 \geq (\text{Im}\mathcal{A})^2$

$$B(\pi^0 \rightarrow e^+ e^-) \geq B^{\text{unitary}}(\pi^0 \rightarrow e^+ e^-) = 4.69 \cdot 10^{-8}$$

(doesn't depend on TFF)

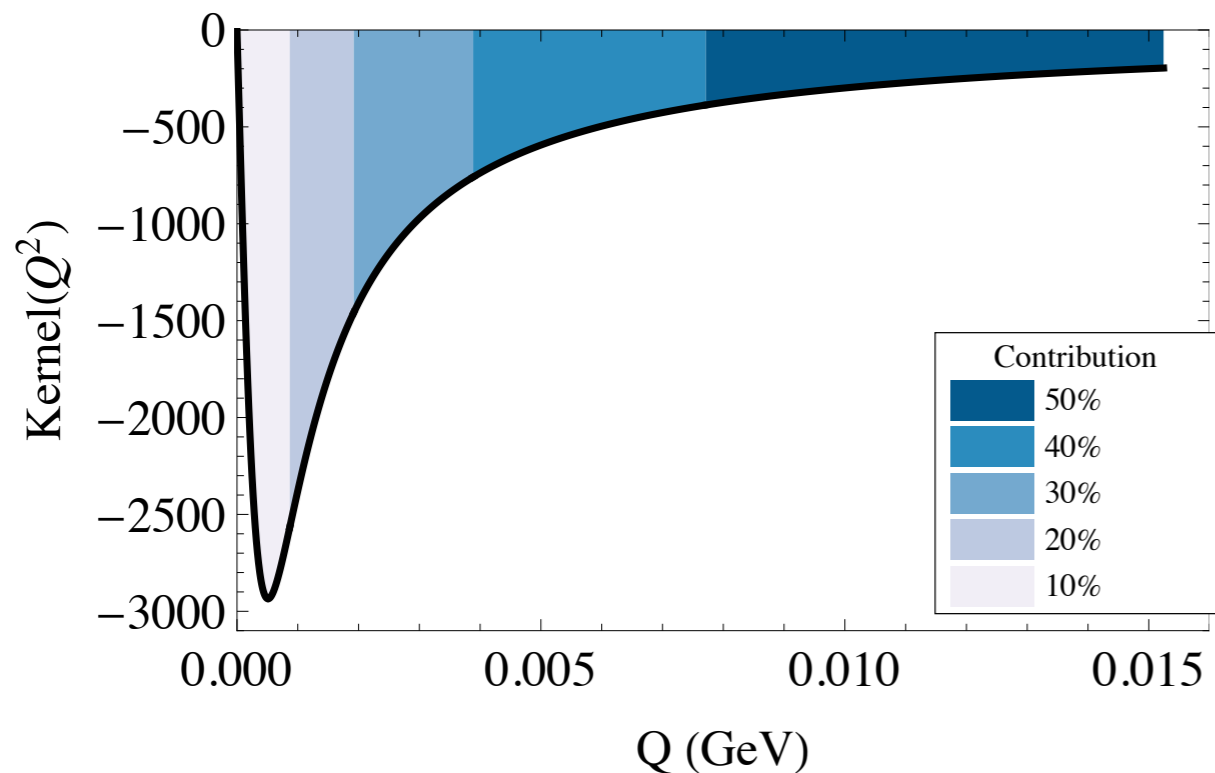
Dissection of $\pi^0 \rightarrow e^+e^-$

$$\text{Re}(\mathcal{A}(m_P^2)) = \left(-\frac{5}{4} + \int_0^\infty dQ^2 \text{Kernel}(Q^2) \right) + \frac{\pi^2}{12} + \ln^2 \left(\frac{m_l}{m_P} \right)$$

$$\text{Re}(\mathcal{A}(m_P^2)) = \int_0^\infty dQ^2 \text{Kernel}(Q^2) + 30.7$$

Dissection of $\pi^0 \rightarrow e^+e^-$

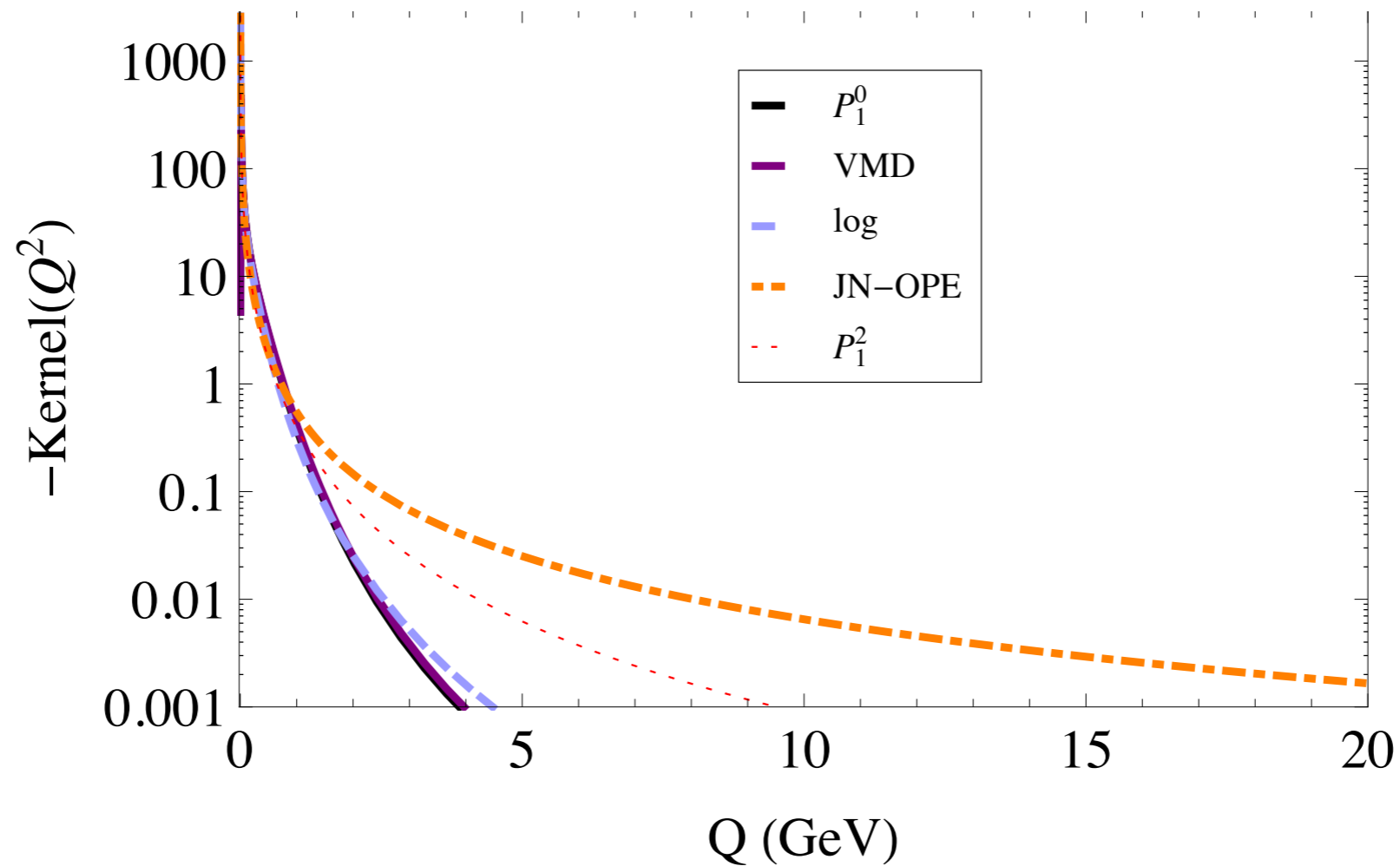
$$\text{Re}(\mathcal{A}(m_P^2)) = \int_0^\infty dQ^2 \text{Kernel}(Q^2) + 30.7$$



- Its contribution is negative: lowers the BR.
- Peaks at $\sim 2m_e$ and $\langle Q \rangle = 0.09$ GeV.
- Low energies relevant only: slope is enough.

Dissection of $\pi^0 \rightarrow e^+e^-$

$$\text{Re}(\mathcal{A}(m_P^2)) = \int_0^\infty dQ^2 \text{Kernel}(Q^2) + 30.7$$

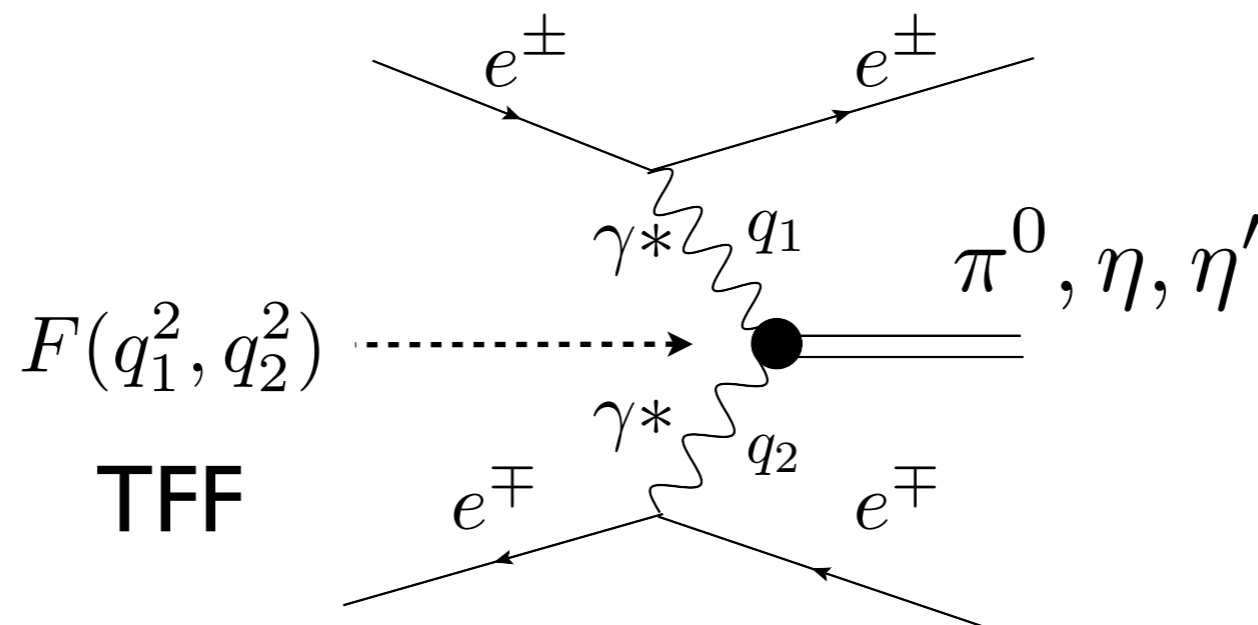


Mainz contribution: TFF parameterization

Use data from
the Transition Form Factor
for numerical integral

$$F_{P\gamma^*\gamma^*}(m_P^2, q_1^2, q_2^2)$$

double-tag method



Remember: only low-energy region is needed

Doubly virtual π^0 -TFF

[P.M., P. Sanchez-Puertas, '15]

For $BR_{SM}(\pi^0 \rightarrow e^+ e^-)$ we need $F_{\pi^0 \gamma^* \gamma^*}(Q^2, Q^2)$

Proposal: Canterbury PA

Chisholm '73

$$P_M^N(Q_1^2, Q_2^2) = \frac{T_N(Q_1^2, Q_2^2)}{R_M(Q_1^2, Q_2^2)} = a_0 + a_1(Q_1^2 + Q_2^2) + a_{1,1}Q_1^2Q_2^2 + a_2(Q_1^4 + Q_2^4) + \dots$$

$$P_1^0(Q_1^2, Q_2^2) = \frac{a_0}{1 + a_1(Q_1^2 + Q_2^2) + (2a_1^2 - a_{1,1})Q_1^2Q_2^2}$$

Doubly virtual π^0 -TFF

Proposal: Canterbury PA

Chisholm '73

$$P_1^0(Q_1^2, Q_2^2) = \frac{a_0}{1 + a_1(Q_1^2 + Q_2^2) + (2a_1^2 - a_{1,1})Q_1^2Q_2^2}$$

a_1 from accurate study of space-like data

$a_{1,1}$ from a systematic fit to doubly virtual SL data

Doubly virtual π^0 -TFF

Proposal: Canterbury PA

Chisholm '73

$$P_1^0(Q_1^2, Q_2^2) = \frac{a_0}{1 + a_1(Q_1^2 + Q_2^2) + (2a_1^2 - a_{1,1})Q_1^2Q_2^2}$$

a_1 from accurate study of space-like data

$a_{1,1}$ from a systematic fit to doubly virtual SL data

OPE indicates: $\lim_{Q^2 \rightarrow \infty} P_1^0(Q^2, Q^2) \sim Q^{-2}$ i.e., $a_{1,1} = 2a_1^2$

Doubly virtual π^0 -TFF

Proposal: Canterbury PA

Chisholm '73

$$P_1^0(Q_1^2, Q_2^2) = \frac{a_0}{1 + a_1(Q_1^2 + Q_2^2) + (2a_1^2 - a_{1,1})Q_1^2 Q_2^2}$$

a_1 from accurate study of space-like data

$$0 \leq a_{1,1} \leq 2a_1^2$$

[P.M, P. Sanchez-Puertas, I5]

$$BR_{SM}^{PA}(\pi^0 \rightarrow e^+ e^-) = (6.20 - 6.41)(4) \times 10^{-8}$$

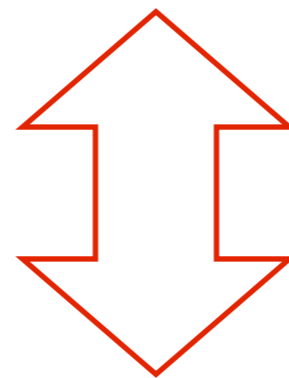
statistics+theoretical error 

method checked for different models
+ to shrink the window: data (data-driven approach)

Doubly virtual π^0 -TFF

[P.M, P. Sanchez-Puertas, 15]

$$BR_{\text{KTeV}}^{w/o rad}(\pi^0 \rightarrow e^+e^-) = (7.48 \pm 0.29 \pm 0.25) \times 10^{-8}$$



$$\sim (3.3 - 2.8)\sigma$$

$$BR_{SM}^{PA}(\pi^0 \rightarrow e^+e^-) = (6.20 - 6.41)(4) \times 10^{-8}$$

[with new RC $\sim 2\sigma$]

Impact of $\pi^0 \rightarrow e^+e^-$ on HLBL

[P.M, P. Sanchez-Puertas, 15]

	Model	Published model		Modified model	
		$\pi^0 \rightarrow e^+e^-$ ($\times 10^8$)	<i>HLBL</i> ($\times 10^{10}$)	$\pi^0 \rightarrow e^+e^-$ ($\times 10^8$)	<i>HLBL</i> ($\times 10^{10}$)
Jegerlehner and Nyffeler '09	LMD+V	6.33	6.29	6.47	5.22
Dorokhov et al '09	VMD	6.34	5.64	6.87	2.44
Our proposal '14	PA	6.36	5.53	6.87	2.85

$$\Delta a_{\mu}^{SM} \sim 6 \times 10^{-10}$$

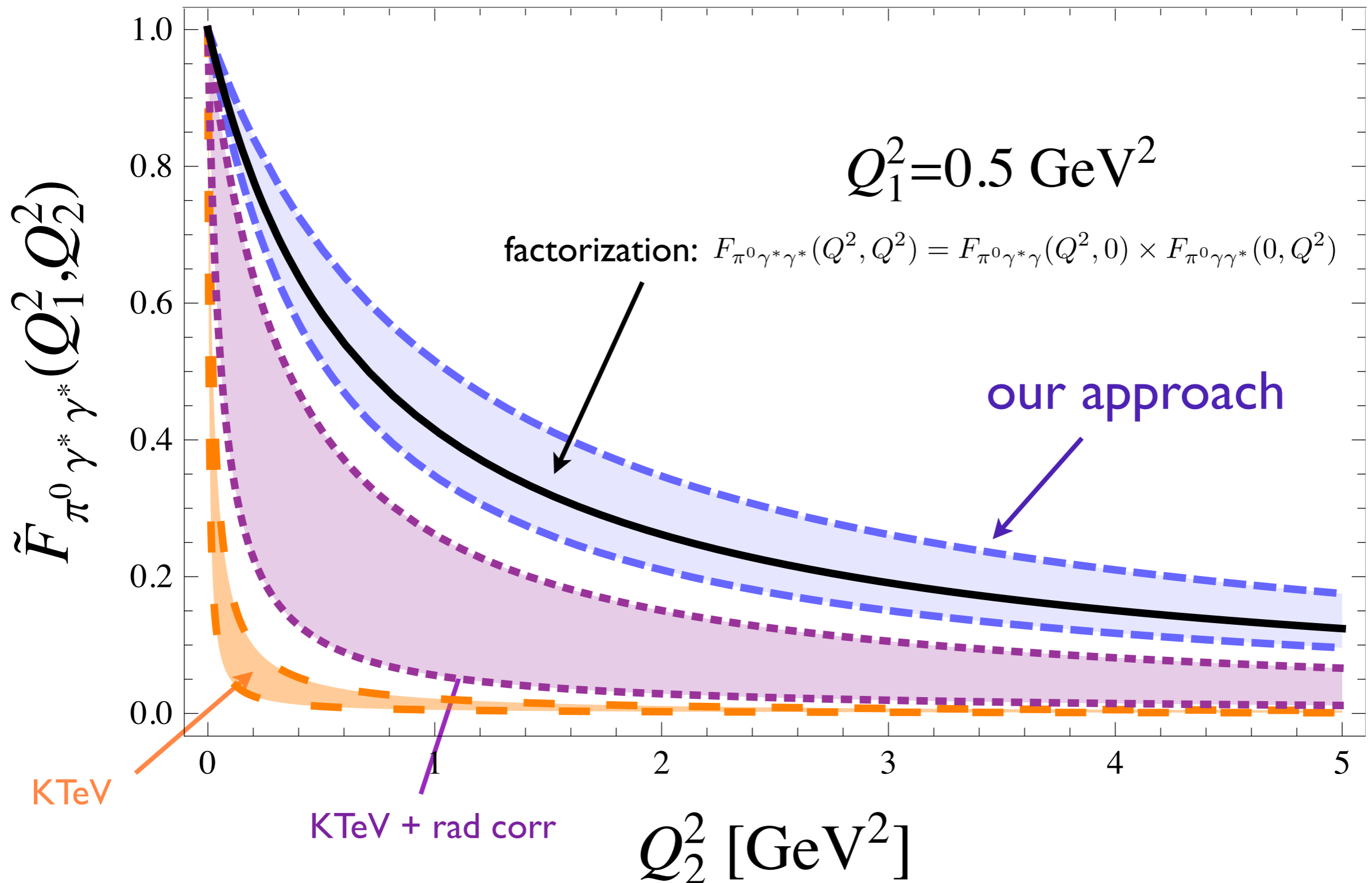
$$\Delta a_{\mu}^{HLBL} \sim 4 \times 10^{-10}$$

$$\Delta a_{\mu}^{HLBL; \pi^0 \rightarrow e^+e^-} \sim (2 - 3) \times 10^{-10}$$

+ similar effect for the η decay!

The role of doubly virtual TFF data

[P.M, P. Sanchez-Puertas, 15]



Dissection of $\eta \rightarrow l^+l^-$

[P.M, P. Sanchez-Puertas, in prep.]

PDG value dominated by the KTeV measurement

$$\frac{BR(P \rightarrow \bar{l}l)}{BR(P \rightarrow \gamma\gamma)} = 2 \left(\frac{\alpha m_l}{\pi m_P} \right)^2 \beta_l(m_P^2) |\mathcal{A}(m_P^2)|^2 = 5.8(8) \cdot 10^{-6} \quad (\mu^+\mu^-)$$
$$\leq 5.6 \cdot 10^{-6} \quad (e^+e^-)$$

$$\text{Unitary Bound for the } \mu\mu \text{ case} = 4.37 \cdot 10^{-6}$$

$$\text{SM calculations with } m_\eta^2/\Lambda^2 \sim 0 = 4.99 \cdot 10^{-6}$$

$$\text{Our result from SL+TL (full result)} = 4.51(2) \cdot 10^{-6}$$

Applications

1. Hadronic Light-by-Light contribution to muon ($g-2$)
2. PS decays into lepton pairs ($\pi^0 \rightarrow e^+e^-$)
3. η - η' mixing

η - η' mixing

η - η' mixing in the flavor basis

$$\begin{pmatrix} f_{\eta}^q & f_{\eta}^s \\ f_{\eta'}^q & f_{\eta'}^s \end{pmatrix} = \begin{pmatrix} f_q \cos[\phi] & -f_s \sin[\phi] \\ f_q \sin[\phi] & f_s \cos[\phi] \end{pmatrix}$$

From the TFFs we can determine f_q, f_s, ϕ

$$\Gamma_{\eta \rightarrow \gamma\gamma} = \frac{9\alpha^2}{32\pi^3} M_{\eta}^3 \left(\frac{C_q \cos[\phi]}{f_q} - \frac{C_s \sin[\phi]}{f_s} \right)^2$$

$$\Gamma_{\eta' \rightarrow \gamma\gamma} = \frac{9\alpha^2}{32\pi^3} M_{\eta'}^3 \left(\frac{C_q \sin[\phi]}{f_q} + \frac{C_s \cos[\phi]}{f_s} \right)^2$$

$$\lim_{Q^2 \rightarrow \infty} Q^2 F_{\eta\gamma\gamma^*}(Q^2) = f_{\eta}^q \frac{10}{3} + f_{\eta}^s \frac{2\sqrt{2}}{3},$$

$$\lim_{Q^2 \rightarrow \infty} Q^2 F_{\eta'\gamma\gamma^*}(Q^2) = f_{\eta'}^q \frac{10}{3} + f_{\eta'}^s \frac{2\sqrt{2}}{3}.$$

η - η' mixing

η - η' mixing in the flavor basis

$$\begin{pmatrix} f_{\eta}^q & f_{\eta}^s \\ f_{\eta'}^q & f_{\eta'}^s \end{pmatrix} = \begin{pmatrix} f_q \cos[\phi] & -f_s \sin[\phi] \\ f_q \sin[\phi] & f_s \cos[\phi] \end{pmatrix}$$

From the TFFs we can determine f_q, f_s, ϕ [R.Escribano, P.M., P. Sanchez-Puertas, '15]

$$\Gamma_{\eta \rightarrow \gamma\gamma} = \frac{9\alpha^2}{32\pi^3} M_{\eta}^3 \left(\frac{C_q \cos[\phi]}{f_q} - \frac{C_s \sin[\phi]}{f_s} \right)^2$$

$$\Gamma_{\eta' \rightarrow \gamma\gamma} = \frac{9\alpha^2}{32\pi^3} M_{\eta'}^3 \left(\frac{C_q \sin[\phi]}{f_q} + \frac{C_s \cos[\phi]}{f_s} \right)^2$$

$$\lim_{Q^2 \rightarrow \infty} Q^2 F_{\eta\gamma\gamma^*}(Q^2) = f_{\eta}^q \frac{10}{3} + f_{\eta}^s \frac{2\sqrt{2}}{3},$$

$$\lim_{Q^2 \rightarrow \infty} Q^2 F_{\eta'\gamma\gamma^*}(Q^2) = f_{\eta'}^q \frac{10}{3} + f_{\eta'}^s \frac{2\sqrt{2}}{3}.$$

[R.Escribano, P.M., P. Sanchez-Puertas, '14]

$$f_q = 1.07(1)f_{\pi}, \quad f_s = 1.39(14)f_{\pi}, \quad \phi = 39.3(1.3)^{\circ}$$

Update of Frere-Escribano '05 with PDG12 using 9 inputs

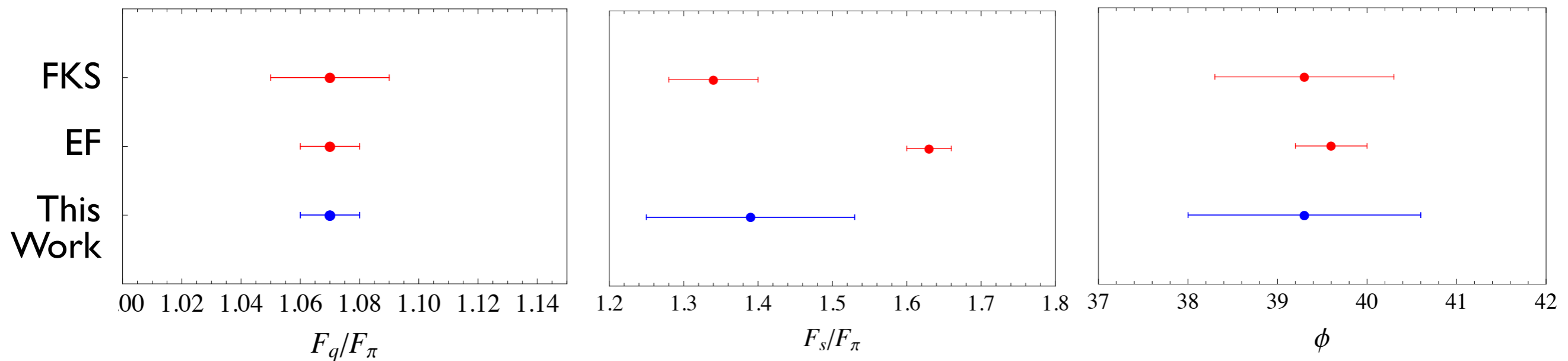
$$f_q = 1.07(1)f_{\pi}, \quad f_s = 1.63(2)f_{\pi}, \quad \phi = 40.4(0.3)^{\circ}$$

η - η' mixing

η - η' mixing in the flavor basis

From the TFFs we can determine F_q, F_s, ϕ

[R.Escribano, P.M., P. Sanchez-Puertas, '15]



FKS: Feldmann, Kroll, Stech, PLB 449, 339, (1999)

EF: Escribano, Frere, JHEP 0506, 029 (2005) updated in Escribano, P.M, Sanchez-Puertas, 2013.

Conclusions

- Transition Form Factors are a good laboratory to study meson properties (one and two virtualities)
- We propose Padé Approximants' method: is *easy, systematic* and can be *improved* upon by including new data.
- Considering space- and time-like data
 - provides very accurate LECs and asymptotic limits
 - provides insight in mixing scheme and meson structure
 - predicts $V\rho\gamma$, J/ψ , rare decays, continuum...
 - beautiful synergy experiment - theory

Outlook

- To THEORY:
 - rethink VMD, LMD, LMD+V, HQCD in terms of PAs: data-driven and systematic
- To DISPERSION Theory:
 - excellent interpolation tool: bring DRs to ∞ (link with pQCD)
 - progress in PA theory is undergoing: better access TL, use unitary...
- To LATTICE:
 - used recently for lattice fitting of HVP: suggestion use it correctly!
- To EXPERIMENT:
 - what energy and precision to measure: use PAs to identify!

Thanks!

Dissection of the HLBL contribution

a la Knecht-Nyffeler '01

Error budget:

$$F_{\pi^0\gamma^*\gamma^*}^{VMD}(q_1^2, q_2^2) = -\frac{N_c}{12\pi^2 f_\pi} \frac{M_V^2}{(q_1^2 - M_V^2)} \frac{M_V^2}{(q_2^2 - M_V^2)}$$

$$F_{\pi^0\gamma^*\gamma^*}^{LMD}(q_1^2, q_2^2) = \frac{f_\pi}{3} \frac{(q_1^2 + q_2^2) - cv}{(q_1^2 - M_V^2)(q_2^2 - M_V^2)}$$

$$F_{\pi^0\gamma^*\gamma^*}^{LMD+V}(q_1^2, q_2^2) = \frac{f_\pi}{3} \frac{q_1^2 q_2^2 (q_1^2 + q_2^2) + h_1 (q_1^2 + q_2^2)^2 + h_2 q_1^2 q_2^2 + h_5 (q_1^2 + q_2^2) + h_7}{(q_1^2 - M_{V_1}^2)(q_1^2 - M_{V_2}^2)(q_2^2 - M_{V_1}^2)(q_2^2 - M_{V_2}^2)}$$

$$\begin{aligned} \Delta F_\pi &\Rightarrow 2\Delta a_\mu^{\text{HLBL}, P} \\ \Delta \text{ slope} &\Rightarrow 0.75\Delta a_\mu^{\text{HLBL}, P} \\ \Delta \text{ curv.} &\Rightarrow 0.5\Delta a_\mu^{\text{HLBL}, P} \\ \Delta m_\rho = \Gamma/2 &\Rightarrow 1.3\Delta a_\mu^{\text{HLBL}, P} \end{aligned}$$

Current exp. precision:

$$\begin{aligned} \Delta F_\pi &\sim 1.1\% \\ \Delta \text{ slope} &\sim 13\% \\ \Delta \text{ curvature} &\sim 25\% \end{aligned}$$

Chiral limit

$$F_0 \rightarrow F_\pi \sim 5\%$$

1/Nc

$$\Delta m_\rho \sim 10\%$$

$$\Delta a_\mu^{\text{HLBL}, \pi} \sim 15\%$$

Dubna contribution: corrections $m_e/m_\pi, m_e/\Lambda$

Dorokhov and Ivanov, '08

$$\mathcal{O}\left(\frac{m_e}{\Lambda}\right)^2 \quad \mathcal{O}\left(\frac{m_e}{\Lambda} \log \frac{m_e}{\Lambda}\right)^2$$

Dorokhov, Ivanov and Kovalenko '09

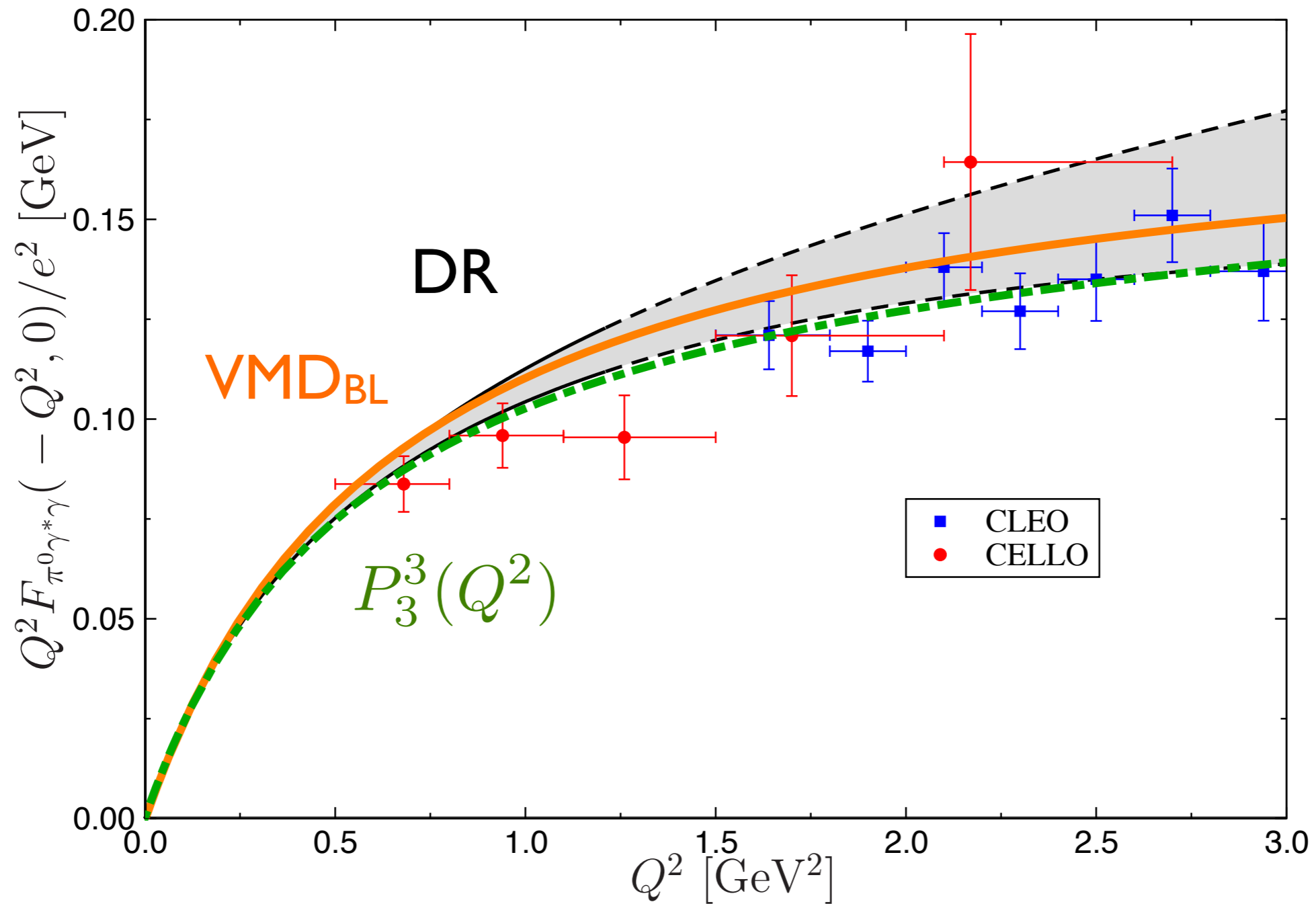
$$\mathcal{O}\left(\frac{m_\pi}{\Lambda}\right)^2 \quad \mathcal{O}\left(\frac{m_e}{m_\pi}\right)^2$$

Λ
the cut-off
or
VMD "mass"

Resummation of power corrections using Mellin-Barnes techniques.
Conclusion: corrections negligible!

$$BR_{\text{SM}}(\pi^0 \rightarrow e^+ e^-) = (6.23 \pm 0.09) \times 10^{-8} \sim 3\sigma$$

PA vs DR



PA vs DR

