

Hadronic Light-by-Light corrections to $(g-2)_\mu$



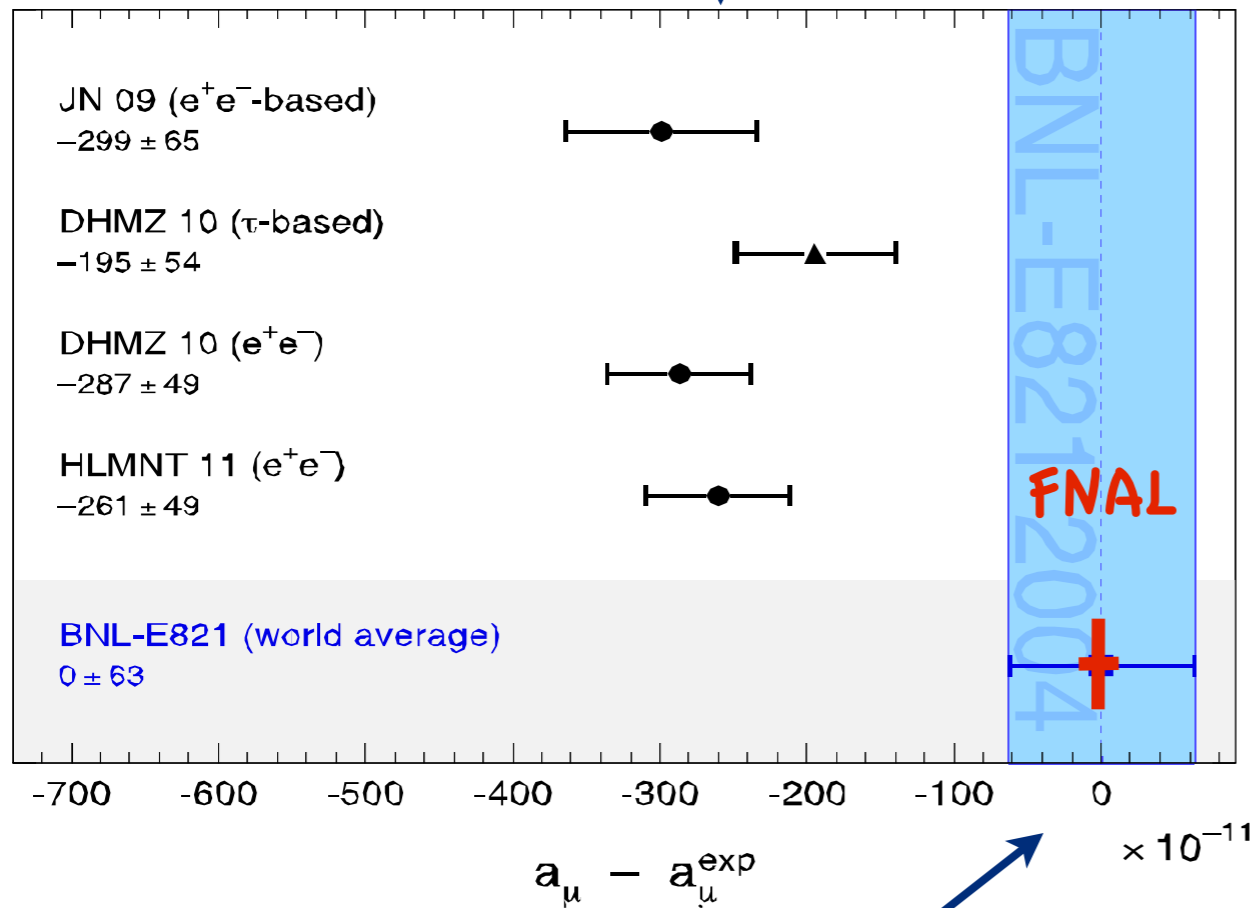
Marc Vanderhaeghen

FCCP 2015, September 10- 12, 2015

Anacapri, Italy

$(g-2)_\mu$: theory vs experiment

SM predictions for a_μ



BNL-E821 measurement of a_μ

$$a_\mu^{\text{exp}} = (11\,659\,208.9 \pm 6.3) \times 10^{-10}$$

$$a_\mu^{\text{exp}} - a_\mu^{\text{SM}} = (26.1 \pm 5.0_{\text{th}} \pm 6.3_{\text{exp}}) \times 10^{-10}$$

Hagiwara et al. (2011)

3 - 4 σ deviation from SM value !

Errors or new physics ?

New FNAL experiment (2016)

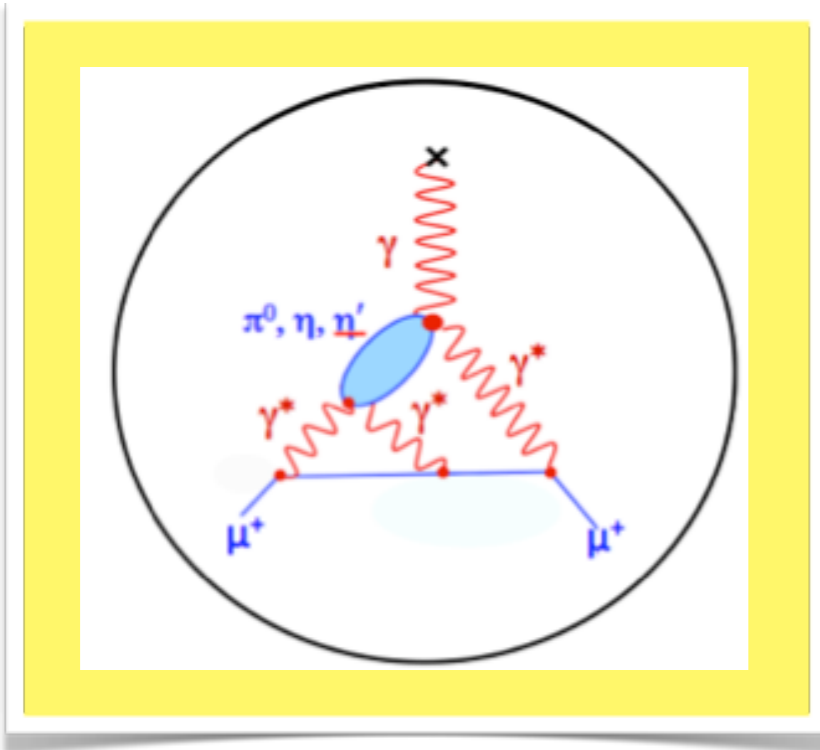
$$\delta a_\mu^{\text{FNAL}} = 1.6 \times 10^{-10}$$

factor 4 improvement in exp. error

-> Improve theory !

hadronic LbL corrections to $(g-2)_\mu$

New FNAL and J-Parc $(g-2)_\mu$ expt. : $\delta a_\mu^{exp} = 1.6 \times 10^{-10}$



$$a_\mu^{\text{had, LbL}} = (11.6 \pm 4.0) \times 10^{-10}$$

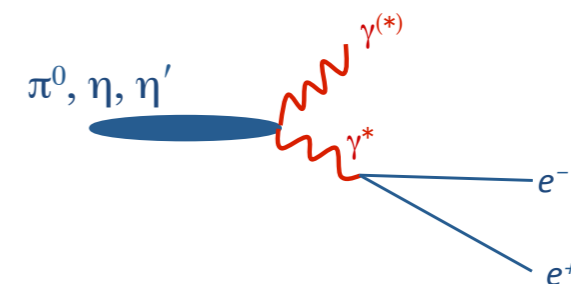
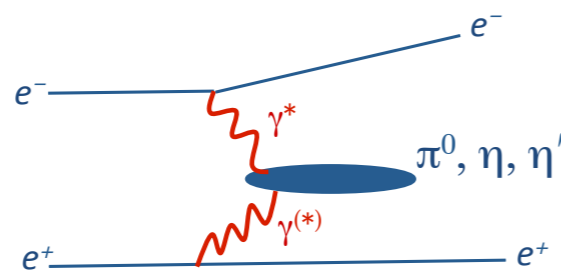
Jegerlehner, Nyffeler (2009)

$$a_\mu^{\text{had, LbL}} = (10.5 \pm 2.6) \times 10^{-10}$$

Prades, de Rafael, Vainshtein (2009)

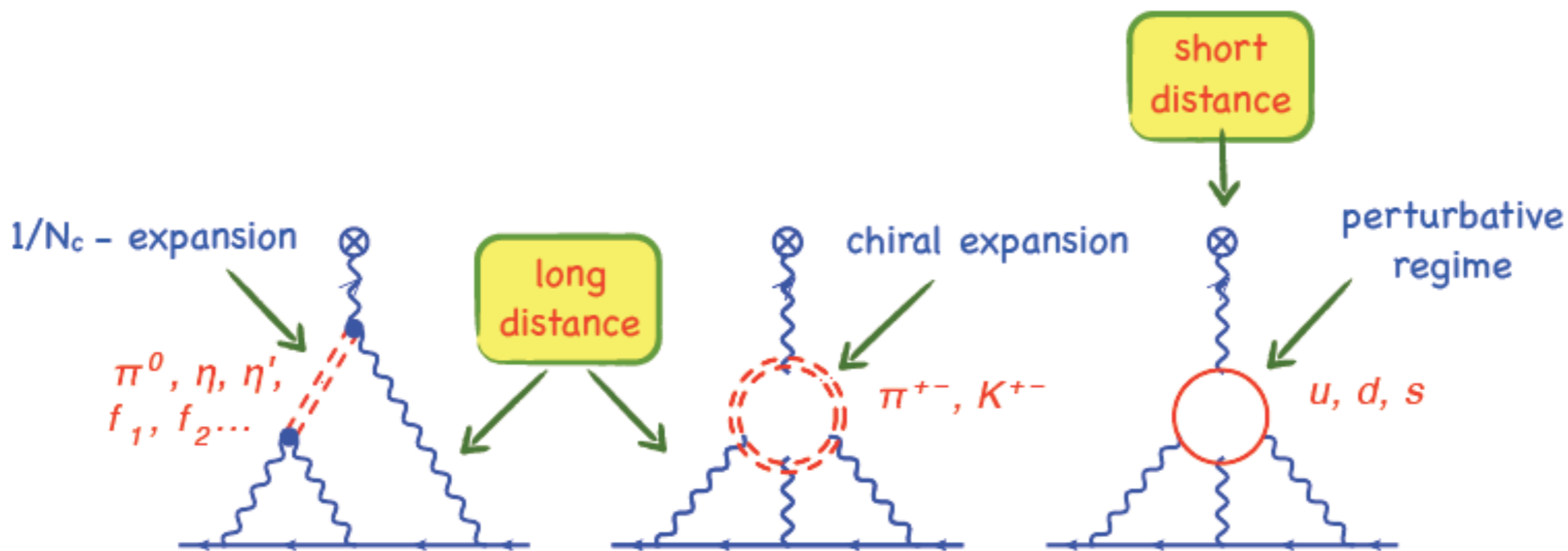
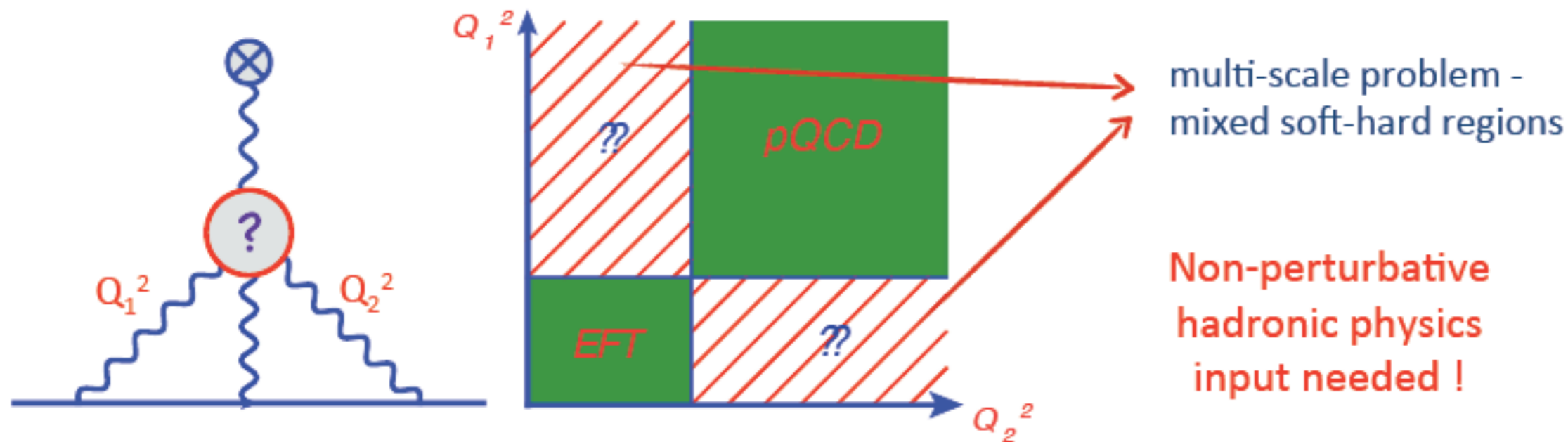
➔ **experimental input:** meson transition FFs, $\gamma^* \gamma^* \rightarrow$ multi-meson states, meson Dalitz decays

**BES-III,
MAMI,
...**

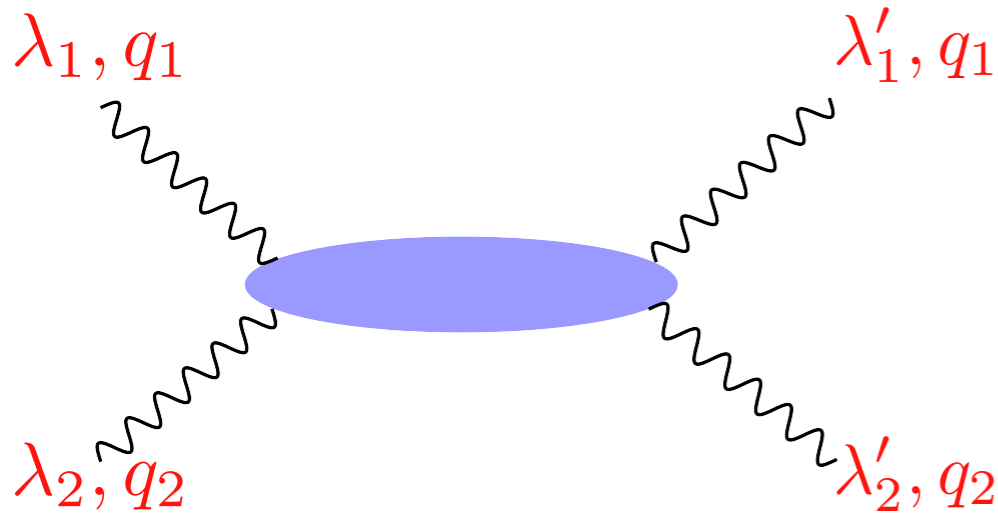


➔ **theory developments:** models, sum rules, dispersion relations, lattice, ...

hadronic LbL corrections to $(g-2)_\mu$: relevant contributions



Theory: sum rules for LbL scattering (I)



$$\gamma^*(\lambda_1, q_1) + \gamma^*(\lambda_2, q_2) \rightarrow \gamma^*(\lambda'_1, q_1) + \gamma^*(\lambda'_2, q_2)$$

kinematical invariants:

$$s = (q_1 + q_2)^2, \quad u = (q_1 - q_2)^2$$

$$\nu \equiv \frac{s - u}{4}, \quad Q_1^2 \equiv -q_1^2, \quad Q_2^2 \equiv -q_2^2$$

helicity amplitudes:

$$M_{\lambda'_1 \lambda'_2, \lambda_1 \lambda_2}(\nu, Q_1^2, Q_2^2)$$

$$\lambda = 0, \pm 1$$

discrete symmetries:



8 independent amplitudes:

$$P : M_{\lambda'_1 \lambda'_2, \lambda_1 \lambda_2} = M_{-\lambda'_1 - \lambda'_2, -\lambda_1 - \lambda_2}$$

$$T : M_{\lambda'_1 \lambda'_2, \lambda_1 \lambda_2} = M_{\lambda_1 \lambda_2, \lambda'_1 \lambda'_2}$$

$$M_{++, ++}, M_{+-, +-}, M_{++, --},$$

$$M_{00, 00}, M_{+0, +0}, M_{0+, 0+}, M_{++, 00}, M_{0+, -0}$$

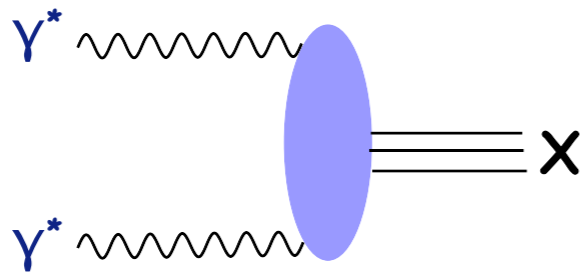
T

T and L

sum rules for LbL scattering (II)

➔ **Unitarity:** link to $\gamma^* \gamma^* \rightarrow X$ cross sections

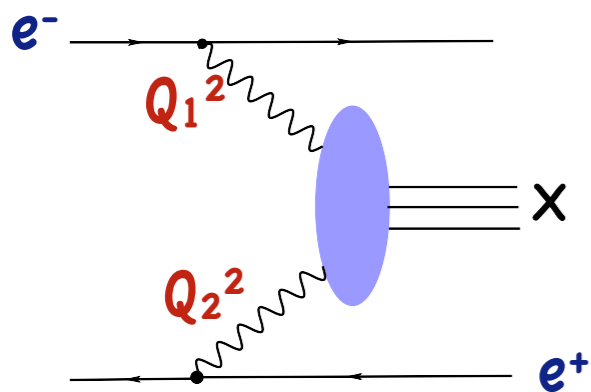
$$W_{\lambda'_1 \lambda'_2, \lambda_1 \lambda_2} \equiv \text{Im } M_{\lambda'_1 \lambda'_2, \lambda_1 \lambda_2}$$



$$\begin{aligned} W_{++,++} + W_{+-,+-} &\equiv 2\sqrt{X} (\sigma_0 + \sigma_2) = 2\sqrt{X} (\sigma_{\parallel} + \sigma_{\perp}) \equiv 4\sqrt{X} \sigma_{TT}, \\ W_{++,++} - W_{+-,+-} &\equiv 2\sqrt{X} (\sigma_0 - \sigma_2) \equiv 4\sqrt{X} \tau_{TT}^a, \\ W_{++,--} &\equiv 2\sqrt{X} (\sigma_{\parallel} - \sigma_{\perp}) \equiv 2\sqrt{X} \tau_{TT}, \\ W_{00,00} &\equiv 2\sqrt{X} \sigma_{LL}, \\ W_{+0,+0} &\equiv 2\sqrt{X} \sigma_{TL}, \\ W_{0+,0+} &\equiv 2\sqrt{X} \sigma_{LT}, \\ W_{++,00} + W_{0+,-0} &\equiv 4\sqrt{X} \tau_{TL}, \\ W_{++,00} - W_{0+,-0} &\equiv 4\sqrt{X} \tau_{TL}^a. \end{aligned}$$

$$X \equiv \nu^2 - Q_1^2 Q_2^2$$

➔ **Experiment:** $e^- e^+ \rightarrow e^- e^+ X$ cross sections

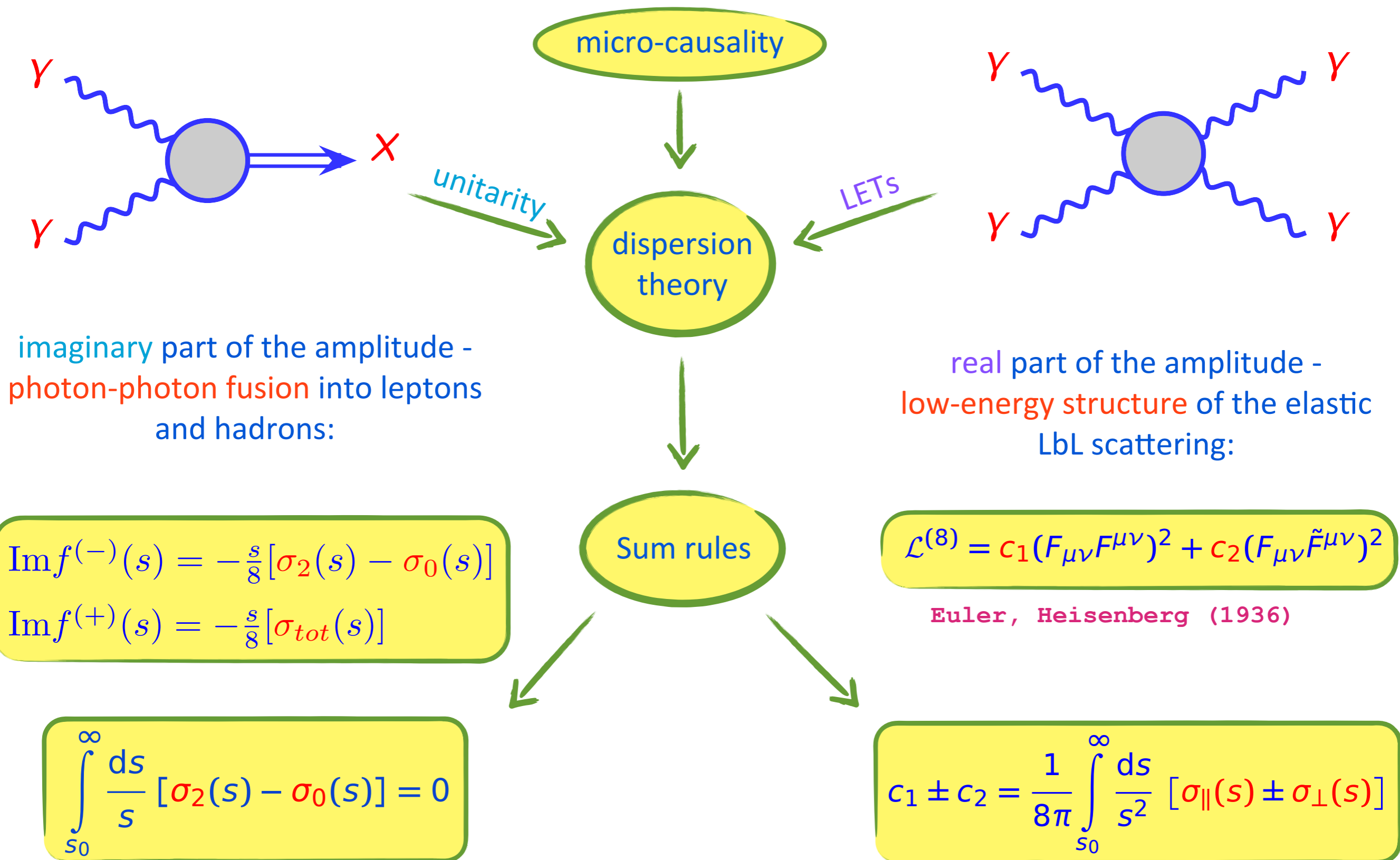


$$\begin{aligned} d\sigma = & \frac{\alpha^2}{16\pi^4 Q_1^2 Q_2^2} \frac{2\sqrt{X}}{s(1-4m^2/s)} \cdot \frac{d^3\vec{p}'_1}{E'_1} \cdot \frac{d^3\vec{p}'_2}{E'_2} \\ & \times \left\{ 4\rho_1^{++}\rho_2^{++}\sigma_{TT} + \rho_1^{00}\rho_2^{00}\sigma_{LL} + 2\rho_1^{++}\rho_2^{00}\sigma_{TL} + 2\rho_1^{00}\rho_2^{++}\sigma_{LT} \right. \\ & + 2(\rho_1^{++}-1)(\rho_2^{++}-1)(\cos 2\tilde{\phi})\tau_{TT} + 8 \left[\frac{(\rho_1^{00}+1)(\rho_2^{00}+1)}{(\rho_1^{++}-1)(\rho_2^{++}-1)} \right]^{1/2} (\cos \tilde{\phi})\tau_{TL} \\ & \left. + h_1 h_2 4 [(\rho_1^{00}+1)(\rho_2^{00}+1)]^{1/2} \tau_{TT}^a + h_1 h_2 8 [(\rho_1^{++}-1)(\rho_2^{++}-1)]^{1/2} (\cos \tilde{\phi})\tau_{TL}^a \right\} \end{aligned}$$

lepton beam polarization

ρ 's, ϕ : kinematical quantities

sum rules for LbL scattering (III)



sum rules for LbL scattering (IV)

3 superconvergent relations:

helicity difference
sum rule

Pascalutsa, Pauk, Vdh (2012, 2014)

$$0 = \int_{s_0}^{\infty} ds \frac{1}{(s + Q_1^2)} [\sigma_0 - \sigma_2]_{Q_2^2=0}$$

for $Q^2 = 0$: GDH sum rule
Gerasimov, Moulin
(1975),
Brodsky, Schmidt
(1995)

sum rules involving
longitudinal photons

$$0 = \int_{s_0}^{\infty} ds \frac{1}{(s + Q_1^2)^2} \left[\sigma_{\parallel} + \sigma_{LT} + \frac{(s + Q_1^2)}{Q_1 Q_2} \tau_{TL}^a \right]_{Q_2^2=0}$$

$$0 = \int_{s_0}^{\infty} ds \left[\frac{\tau_{TL}}{Q_1 Q_2} \right]_{Q_2^2=0}$$

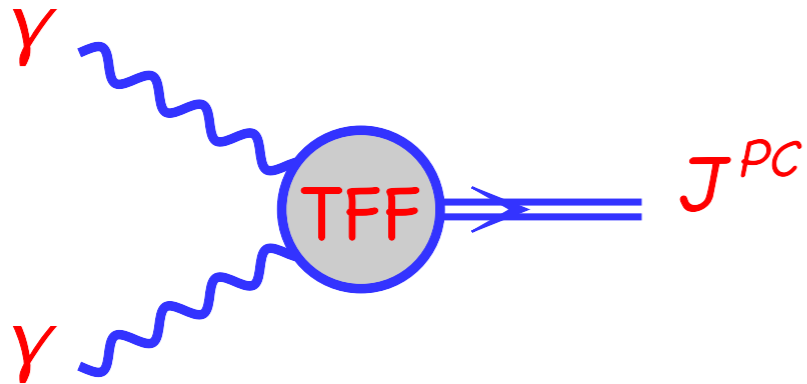
SRs involving LbL
low-energy constants:

$$c_1 \pm c_2 = \frac{1}{8\pi} \int_{s_0}^{\infty} \frac{ds}{s^2} [\sigma_{\parallel}(s) \pm \sigma_{\perp}(s)]$$

+ 6 new LECs at next order

➔ sum rules have been tested in perturbative QFT both at tree-level and 1-loop level

single meson production in $\gamma\gamma$ collisions (I)



- two-photon state: produced meson has $C=+1$
- both photons are real: $J=1$ final state is forbidden (Landau-Yang theorem);
- the main contribution comes from
- $J=0$: 0^{-+} (pseudoscalar) and 0^{++} (scalar)
- and $J=2$: 2^{++} (tensor)

- the SRs hold separately for channels of given intrinsic quantum numbers: isoscalar and isovector mesons, $c\bar{c}$ states

- input for the absorptive part of the SRs: $\gamma\gamma$ -hadrons response functions, can be expressed in terms of $\gamma\gamma \rightarrow M$ transition form factors

$$\sigma_{\Lambda}^{\gamma\gamma \rightarrow M}(s) \approx (2J+1)16\pi^2 \frac{\Gamma_{\gamma\gamma}}{m_M} \delta(s - m_M^2)$$

meson contribution to the cross-section in the narrow-resonance approximation

$$\Gamma_{\gamma\gamma}(P) = \frac{\pi\alpha^2}{4} m^3 |F_{M\gamma^*\gamma^*}(0,0)|^2$$

two-photons decay rate for the meson

single meson production in $\gamma\gamma$ collisions (II)

the $I=0$ channel

	$\int \frac{ds}{s} (\sigma_2 - \sigma_0)$ [nb]	c_1 [10^{-4}GeV^{-4}]	c_2 [10^{-4}GeV^{-4}]
η	-191 ± 10	0	0.65 ± 0.03
η'	-300 ± 10	0	0.33 ± 0.01
$f_0(980)$	-19 ± 5	0.020 ± 0.005	0
$f'_0(1370)$	-91 ± 36	0.049 ± 0.019	0
$f_2(1270)$	449 ± 52	0.141 ± 0.016	0.141 ± 0.016
$f'_2(1525)$	7 ± 1	0.002 ± 0.000	0.002 ± 0.000
$f_2(1565)$	56 ± 11	0.012 ± 0.002	0.012 ± 0.002
Sum	-89 ± 66	0.22 ± 0.03	1.14 ± 0.04

dominant contribution to c_2 comes from η, η' and $f_2(1270)$

dominant contribution to c_1 comes from $f_2(1270)$

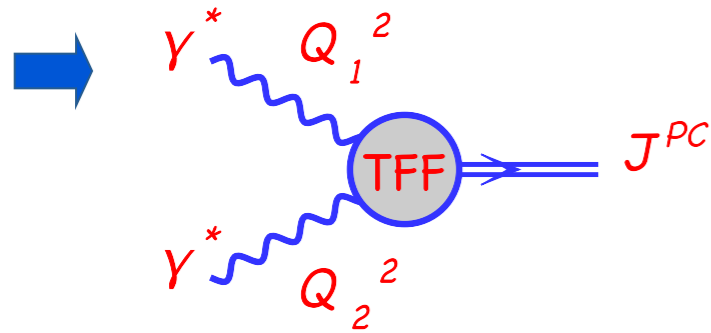
the $I=1$ channel

	$\int \frac{ds}{s} (\sigma_2 - \sigma_0)$ [nb]	c_1 [10^{-4}GeV^{-4}]	c_2 [10^{-4}GeV^{-4}]
π^0	-195 ± 13	0	10.94 ± 0.70
$a_0(980)$	-20 ± 8	0.021 ± 0.007	0
$a_2(1320)$	134 ± 8	0.039 ± 0.002	0.039 ± 0.002
$a_2(1700)$	18 ± 3	0.003 ± 0.001	0.003 ± 0.001
Sum	-63 ± 17	0.06 ± 0.01	10.98 ± 0.70

dominant contribution to c_2 comes from π^0

Pascalutsa, Pauk, Vdh (2012)

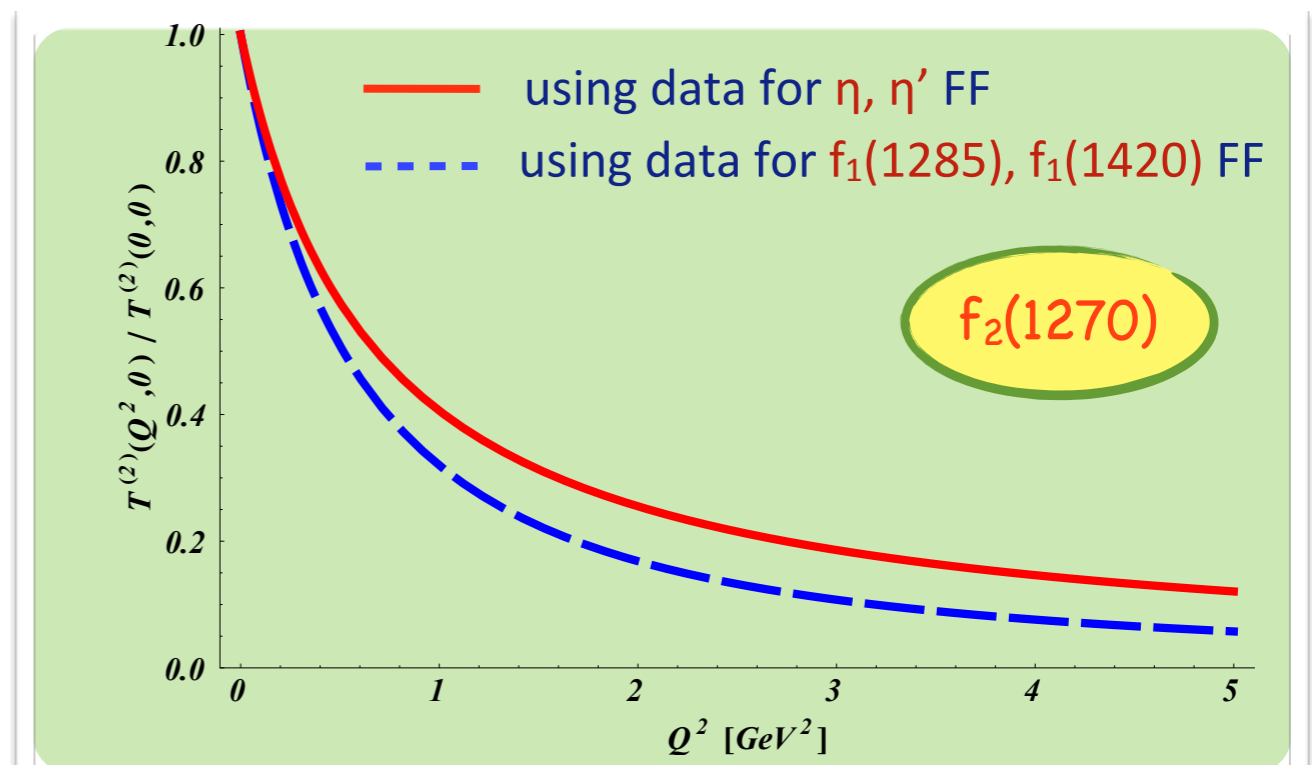
single meson production in $\gamma\gamma$ collisions (III)



- one photon is virtual Q_1^2 , second is quasi-real $Q_2^2 \approx 0$:
- axial-vector mesons 1^{++} are allowed
- $f_1(1285)$, $f_1(1420)$ transition FFs constrained from LEP (L3) data

	m_M [MeV]	$\Gamma_{\gamma\gamma}$ [keV]	$\int \frac{ds}{s^2} \sigma_{\parallel}(s)$ [nb / GeV ²]	$\int ds \left[\frac{1}{s} \frac{\tau_{TL}^a}{Q_1 Q_2} \right]_{Q_i^2=0}$ [nb / GeV ²]	$\int ds \left[\frac{1}{s^2} \sigma_{\parallel} + \frac{1}{s} \frac{\tau_{TL}^a}{Q_1 Q_2} \right]_{Q_i^2=0}$ [nb / GeV ²]
$f_1(1285)$	1281.8 ± 0.6	3.5 ± 0.8	0	-93 ± 21	-93 ± 21
$f_1(1420)$	1426.4 ± 0.9	3.2 ± 0.9	0	-50 ± 14	-50 ± 14
$f_0(980)$	980 ± 10	0.29 ± 0.07	20 ± 5	0	20 ± 5
$f_0'(1370)$	1200 – 1500	3.8 ± 1.5	48 ± 19	0	48 ± 19
$f_2(1270)$	1275.1 ± 1.2	3.03 ± 0.35	138 ± 16	≥ 0	138 ± 16
$f_2'(1525)$	1525 ± 5	0.081 ± 0.009	1.5 ± 0.2	≥ 0	1.5 ± 0.2
$f_2(1565)$	1562 ± 13	0.70 ± 0.14	12 ± 2	≥ 0	12 ± 2
Sum					76 ± 36

sum rules allow to constrain so far unmeasured contributions, e.g. $\gamma^* \gamma^* \rightarrow$ tensor mesons

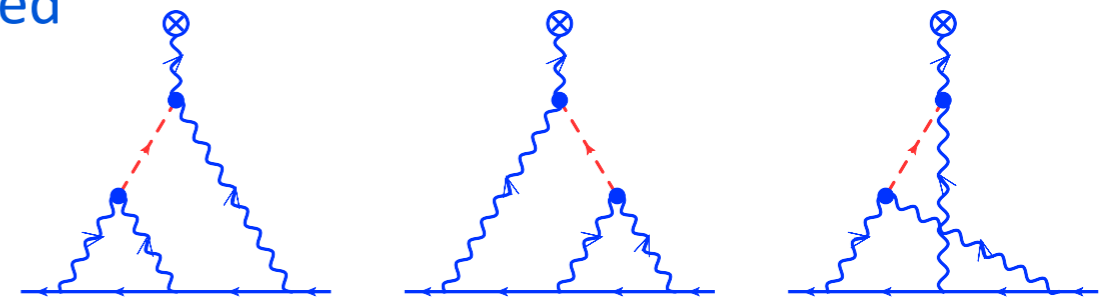


single meson contributions to a_μ (I)



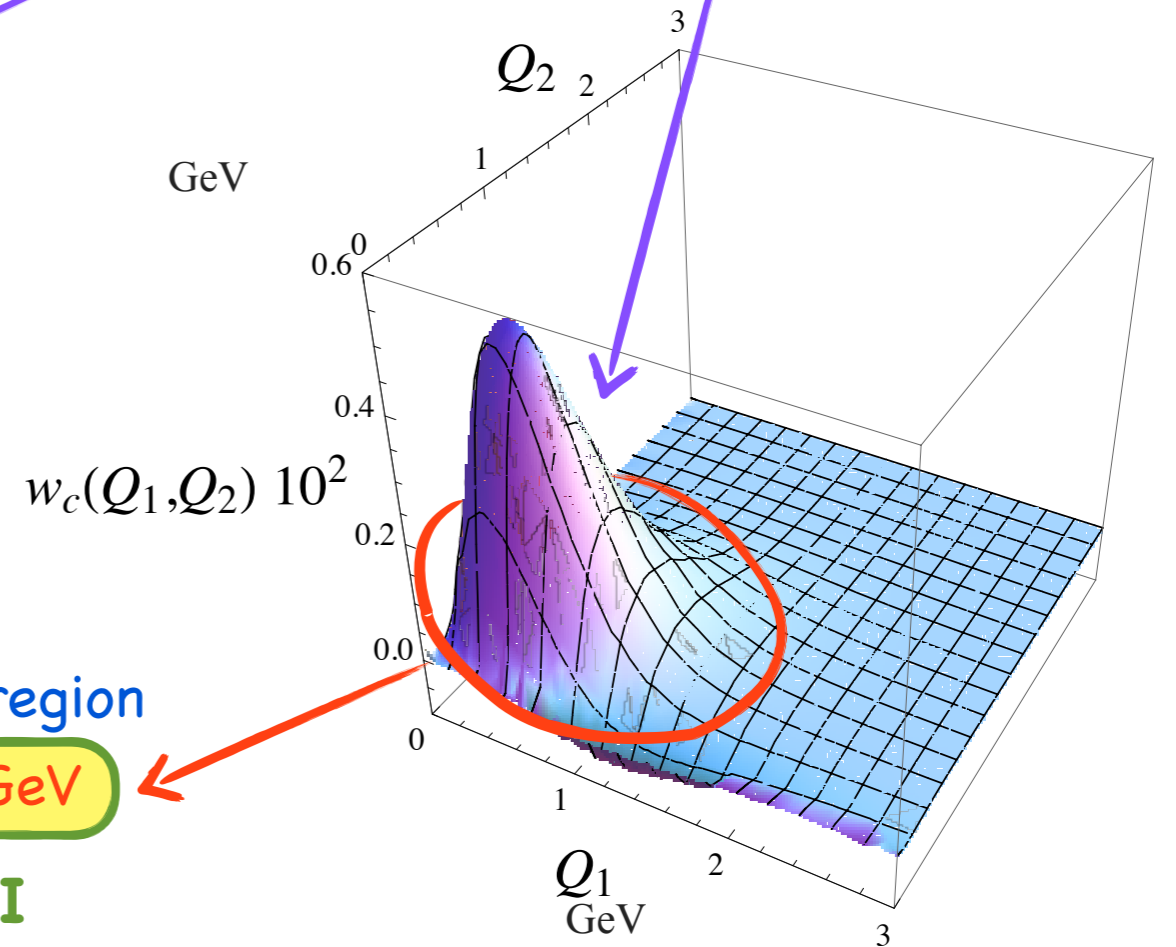
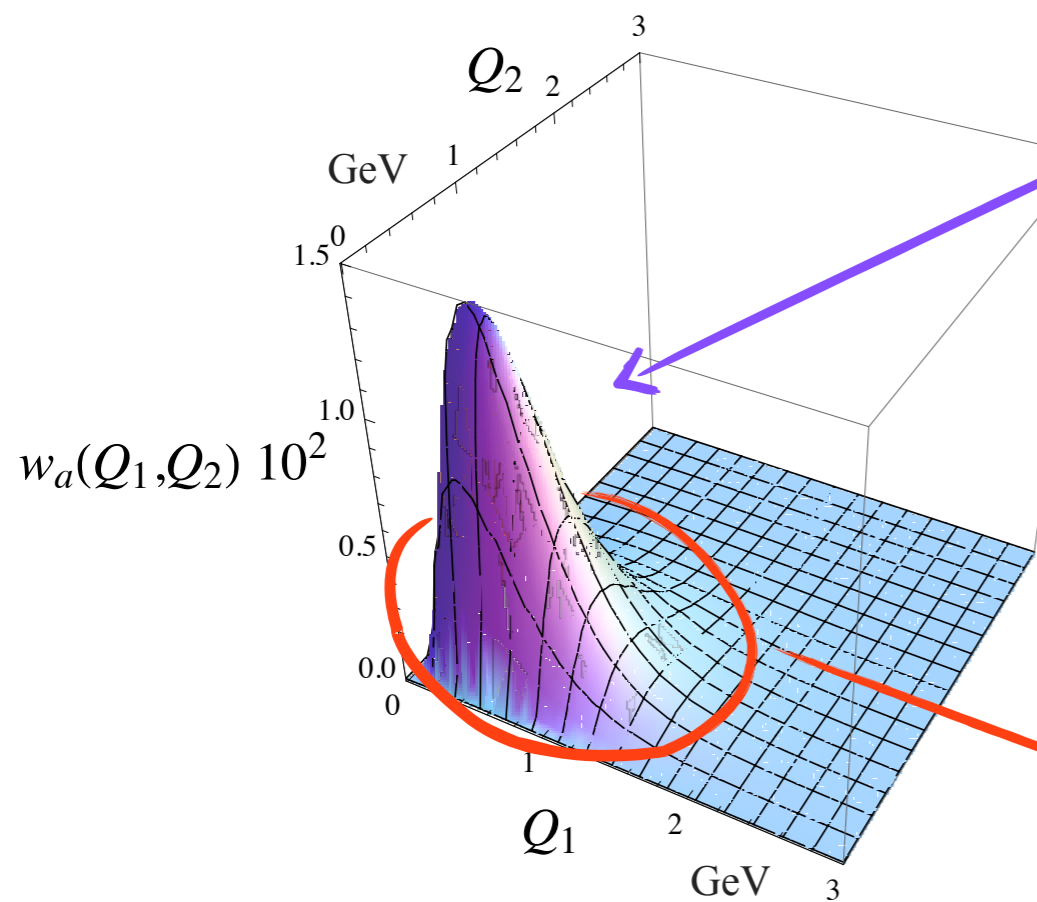
axial-vector meson contribution to a_μ re-evaluated

- Landau-Yang theorem constraint implemented
- $f_1(1285)$, $f_1(1420)$ transition FFs from L3 data



Pauk, Vdh (2013)

$$a_\mu^{LbL} = \frac{\alpha}{(2\pi)^2} \frac{\Lambda_{A1}^6 \Lambda_{A2}^6 \tilde{\Gamma}_{\gamma\gamma}(A)}{m m_A^5} \int dQ_1 \int dQ_2 [2w_a(Q_1, Q_2) + w_c(Q_1, Q_2)]$$



dominating region

$Q_1 \sim Q_2 \sim 1 \text{ GeV}$

BES III

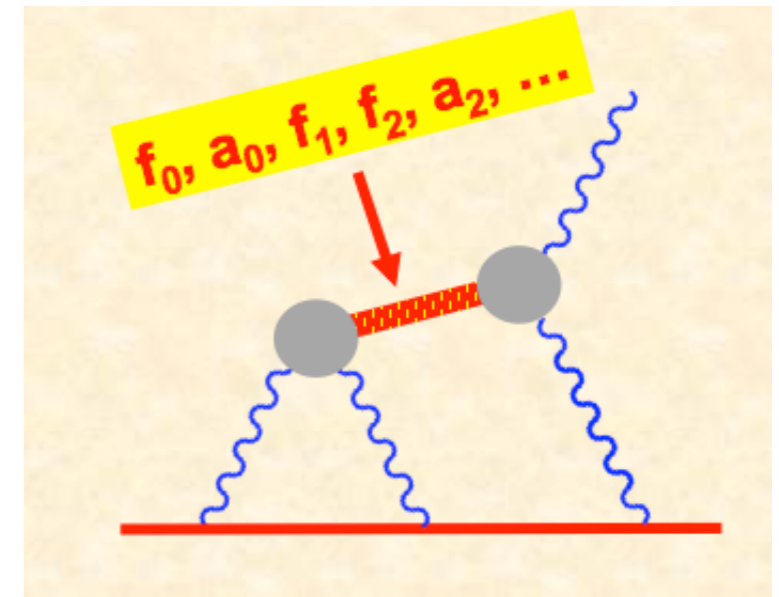
single meson contributions to a_μ (II)

➔ **axial-vector meson** re-evaluation was reported in 2 works

- implementation of Landau-Yang theorem

constraint leads to difference with previous results

➔ **tensor mesons** evaluated for first time



	pseudo-scalars	axial-vectors	scalars	tensors
BPP	85 ± 13	2.5 ± 1.0	-7 ± 2	-
HKS	82.7 ± 6.4	1.7 ± 1.7	-	-
MV	114 ± 10	22 ± 5	-	-
KN	83 ± 12	-	-	-
J	93.9 ± 12.4	~ 7	-6.0 ± 1.2	-
this work	-	6.4 ± 2.0	$-(0.9 \sim 3.1) \pm 0.8$	1.1 ± 0.1

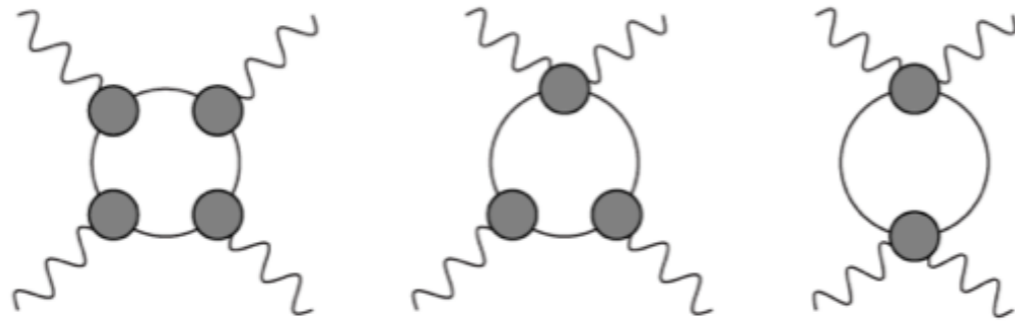
Jegerlehner

Pauk, Vdh

total $(6.6 \sim 4.4) \pm 2.9 \times 10^{-11}$

multi-meson production in $\gamma\gamma$ collisions (I)

➔ new estimate for **pion loop** contribution (with full VMD FF) **Bijnens (2014)**

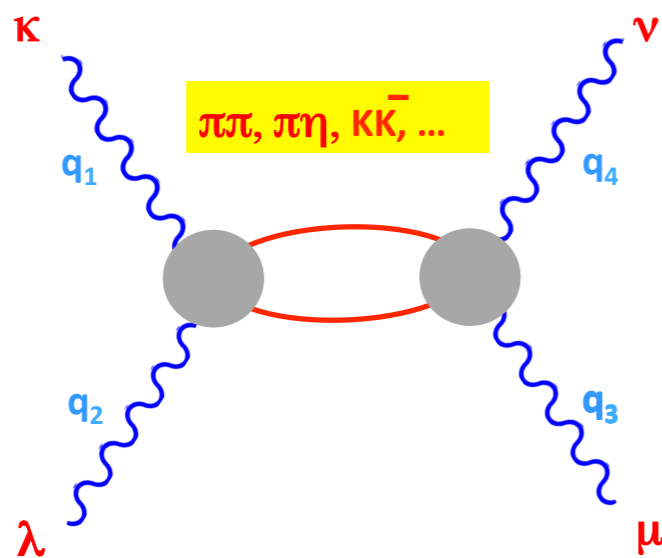


$$a_{\mu}^{\text{LbL } \pi\text{-loop}} = (-2.0 \pm 0.5) \times 10^{-10}$$

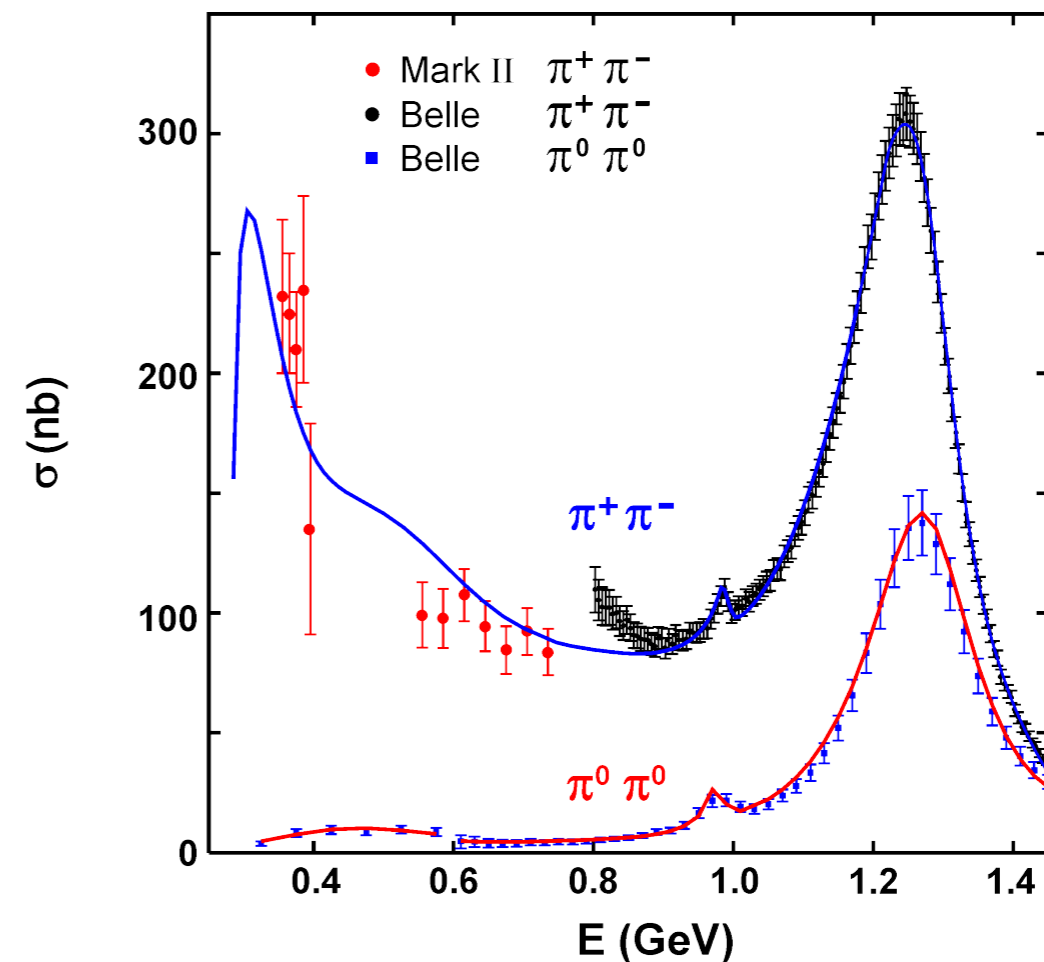
integrating momenta in loop up to 1 GeV

➔ contribution of **multi-meson channels**

Dai, Pennington (2014)



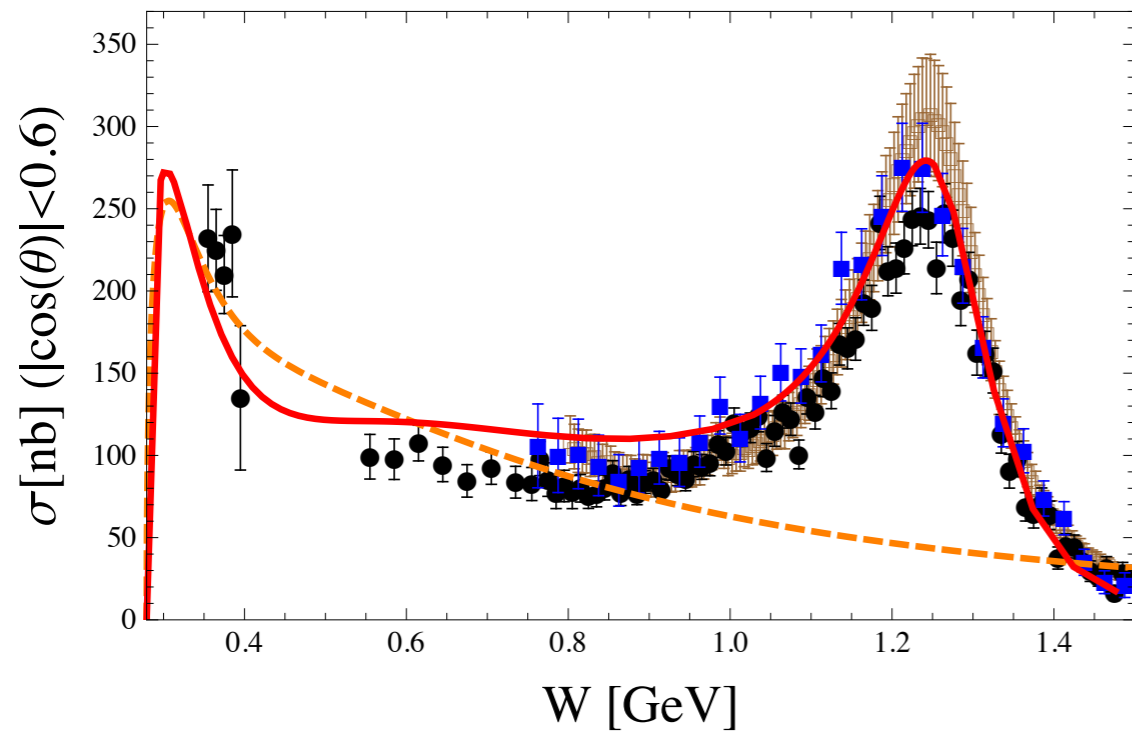
sum rules may be used as a consistency check of models



multi-meson production in $\gamma\gamma$ collisions (II)

➔ new dispersion formalism for $\gamma^* \gamma \rightarrow \pi\pi$

$\gamma\gamma \rightarrow \pi^+\pi^-$



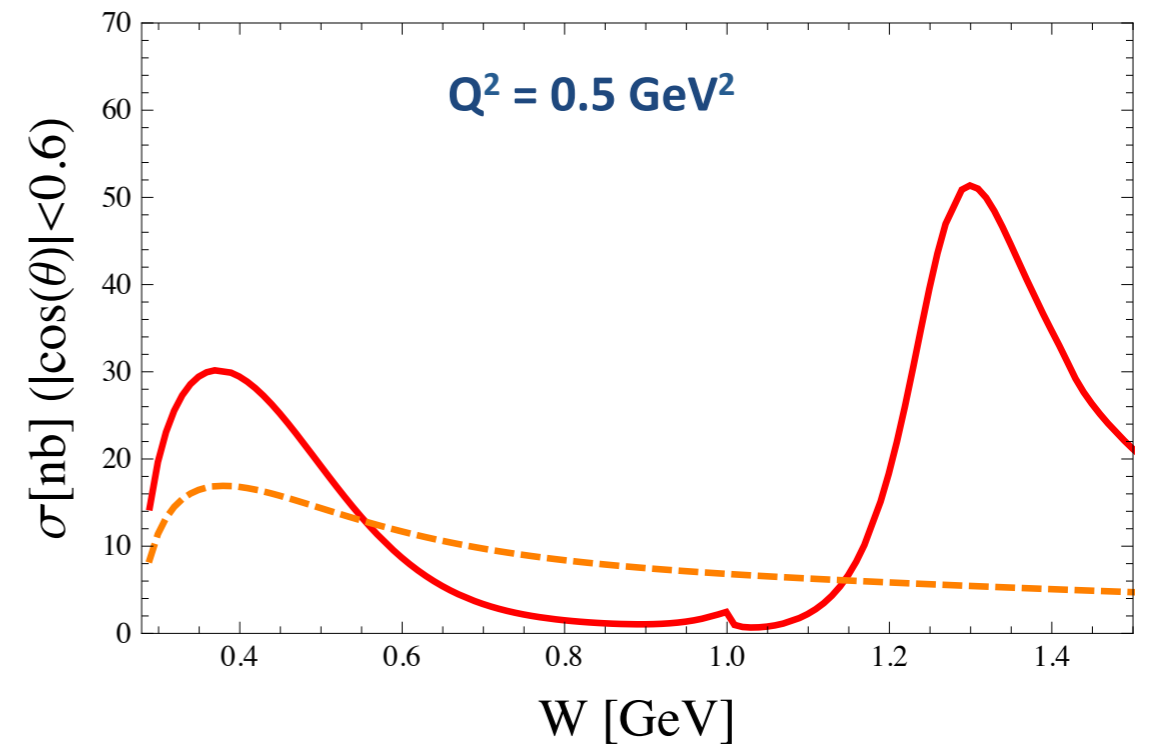
Masjuan, Sanchez Puertas
et al. (2015)



Born terms

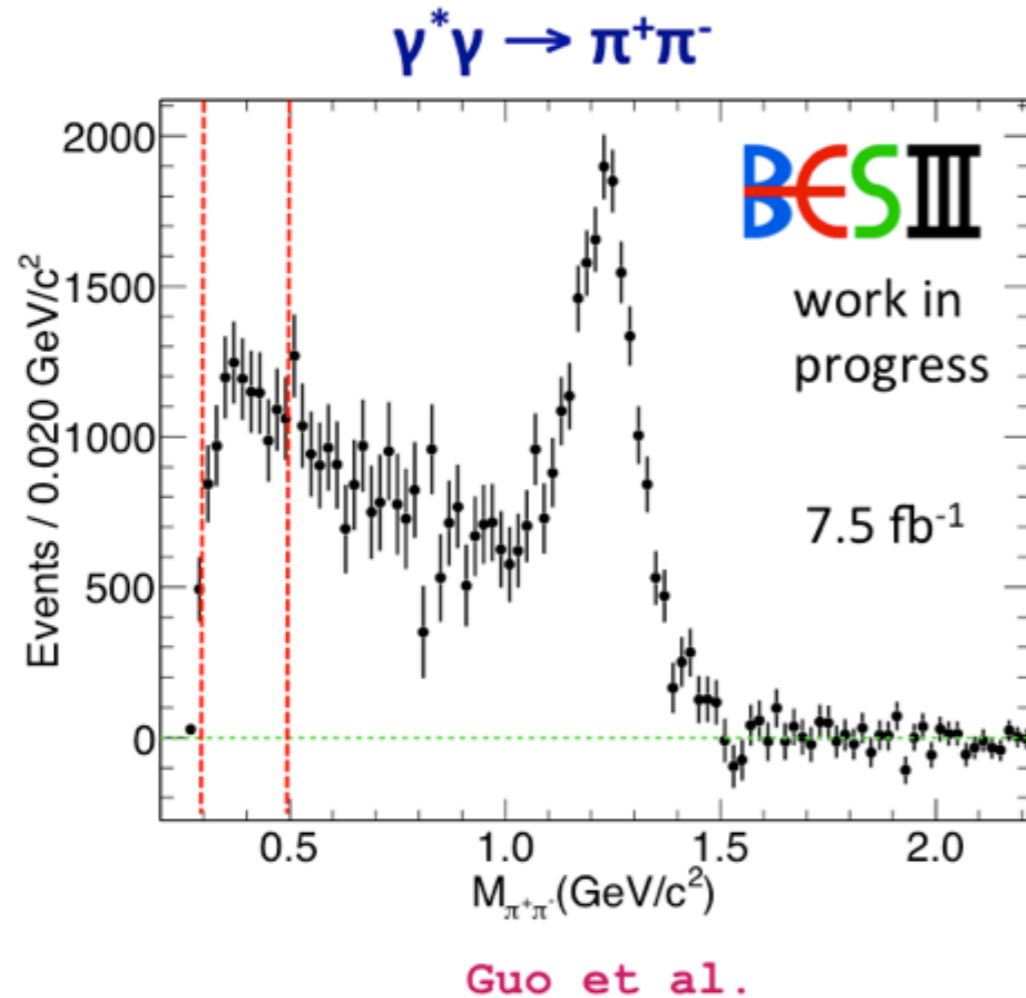
Unitarized s- and d-waves

$\gamma^* \gamma \rightarrow \pi^+\pi^-$

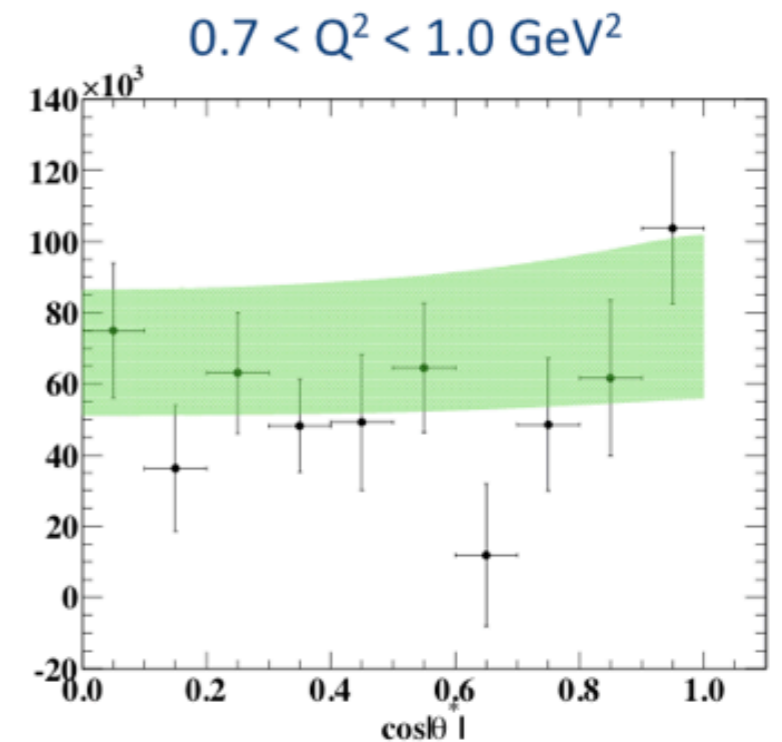
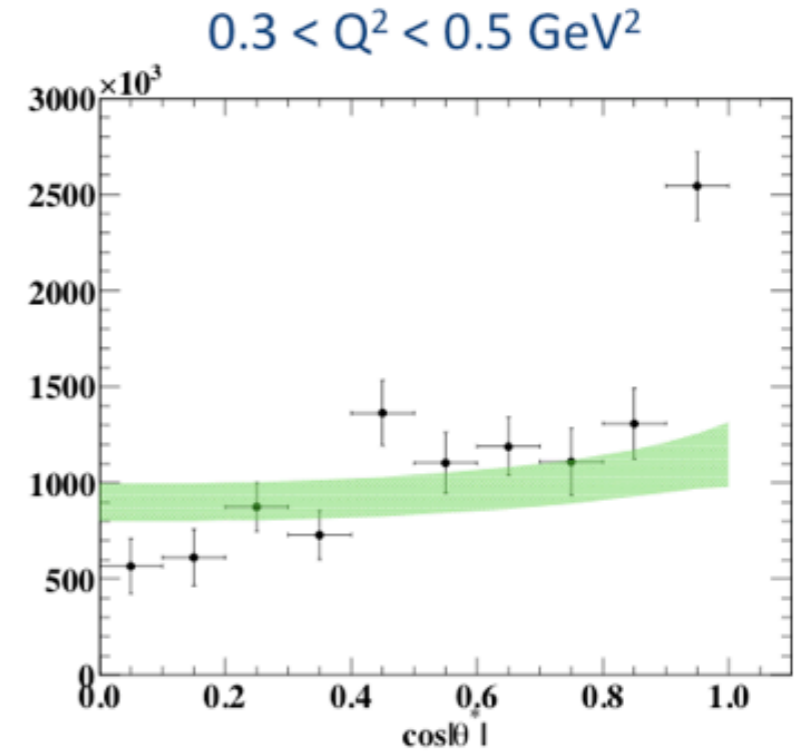


$Q^2 = 0.5 \text{ GeV}^2$

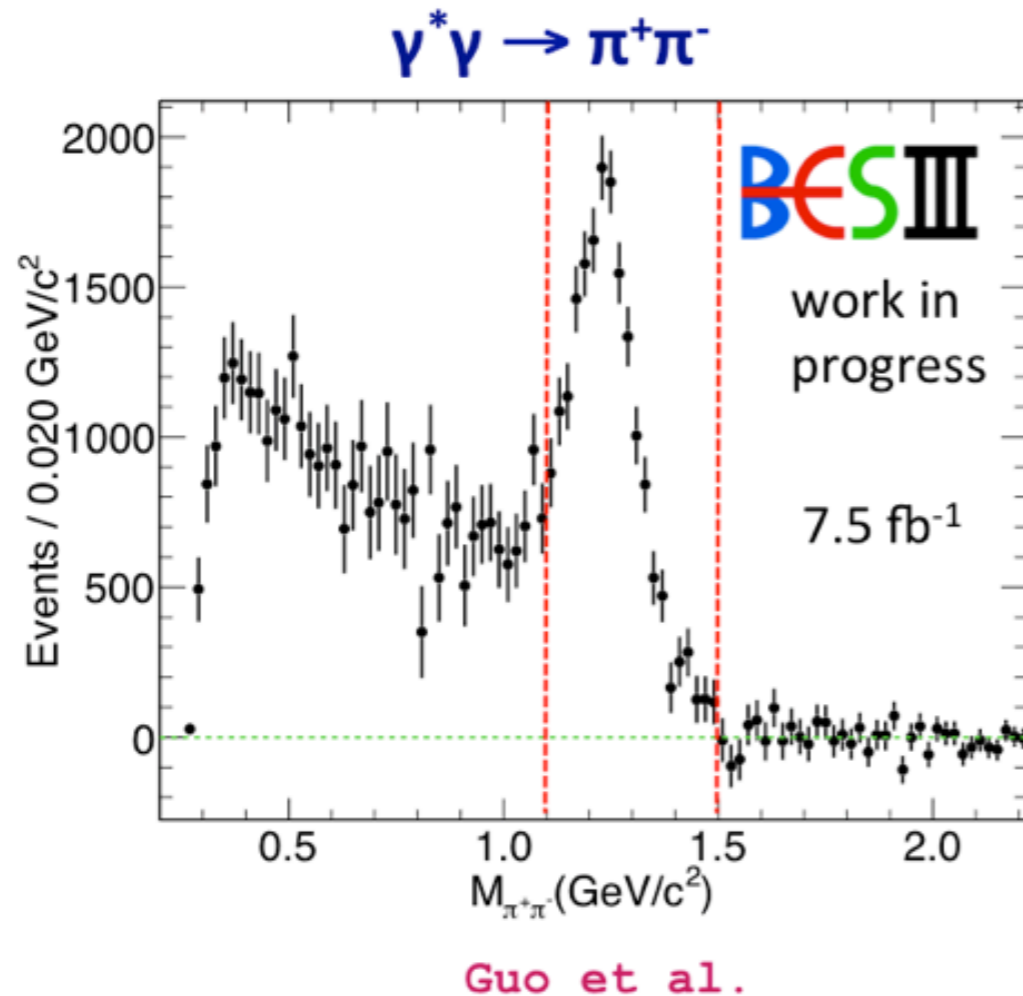
multi-meson production in $\gamma\gamma$ collisions (III)



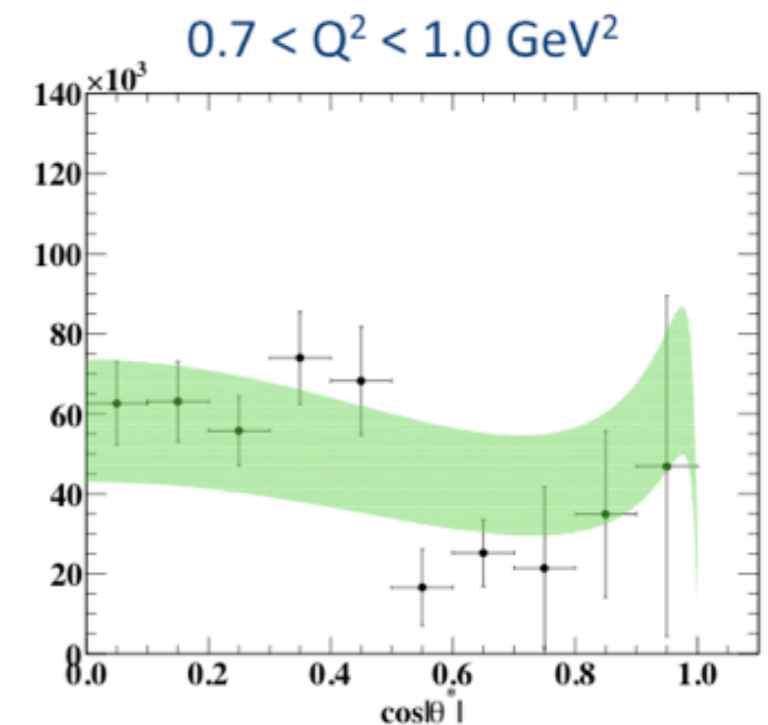
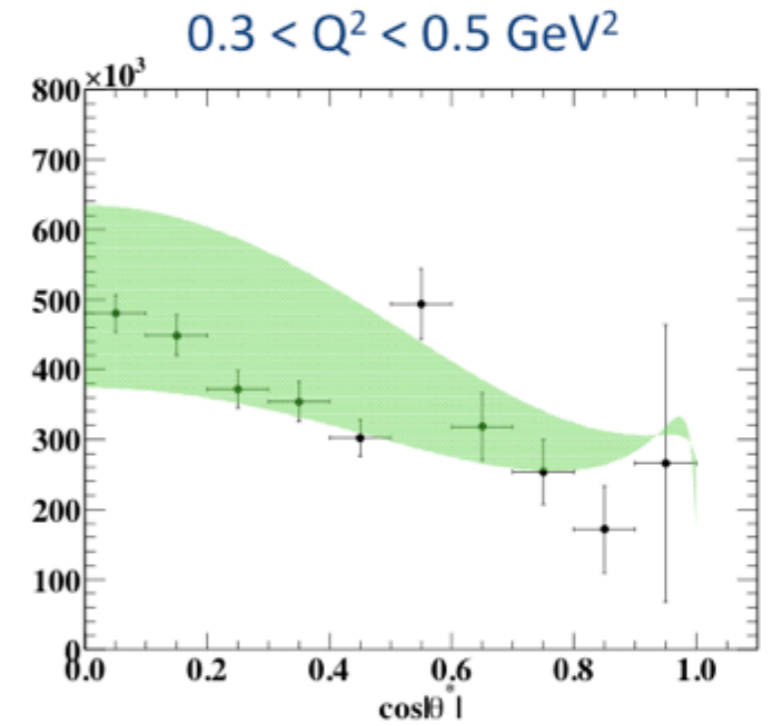
- First single tag measurements of $\pi^+\pi^-$ channel
- Successful μ, π separation: neural network



multi-meson production in $\gamma\gamma$ collisions (IV)



- First single tag measurements of $\pi^+\pi^-$ channel
- Successful μ, π separation: neural network



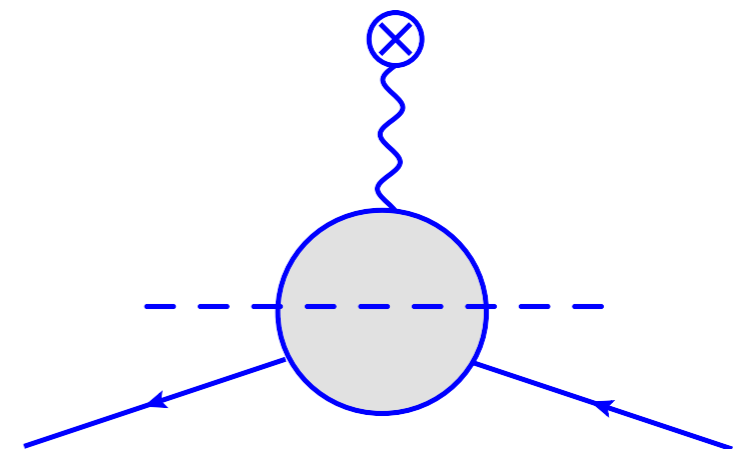
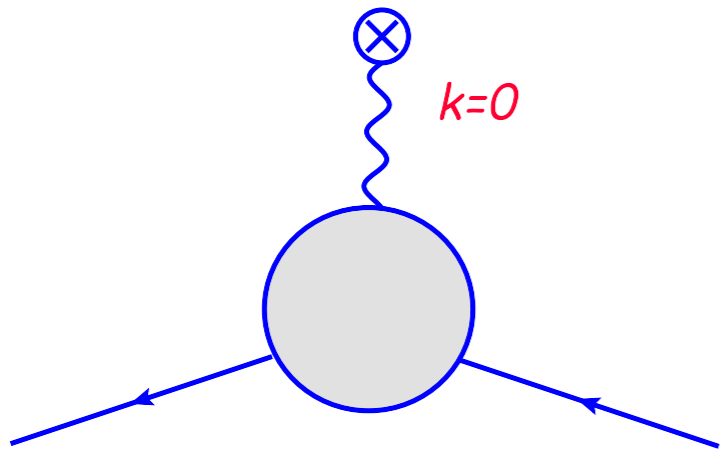
dispersion relation approaches for a_μ (I)

➔ dispersion formalism directly for a_μ Pauk, Vdh (2014)

$$a_\mu = \lim_{k \rightarrow 0} F_2(k^2, (p+k)^2, p^2)$$

$$F_2(0) = \frac{1}{2\pi i} \int \frac{dk^2}{k^2} \text{Abs } F_2(k^2)$$

$$a_\mu = F_2(0)$$



$$F_2(k^2) = e^6 \sum_{\lambda_1, \lambda_2, \lambda_3, \lambda} (-1)^{\lambda + \lambda_1 + \lambda_2 + \lambda_3} \int \frac{d^4 q_1}{(2\pi)^4} \int \frac{d^4 q_2}{(2\pi)^4} L_{\lambda_1 \lambda_2 \lambda_3 \lambda}(p, p', q_1, k - q_1 - q_2, q_2)$$

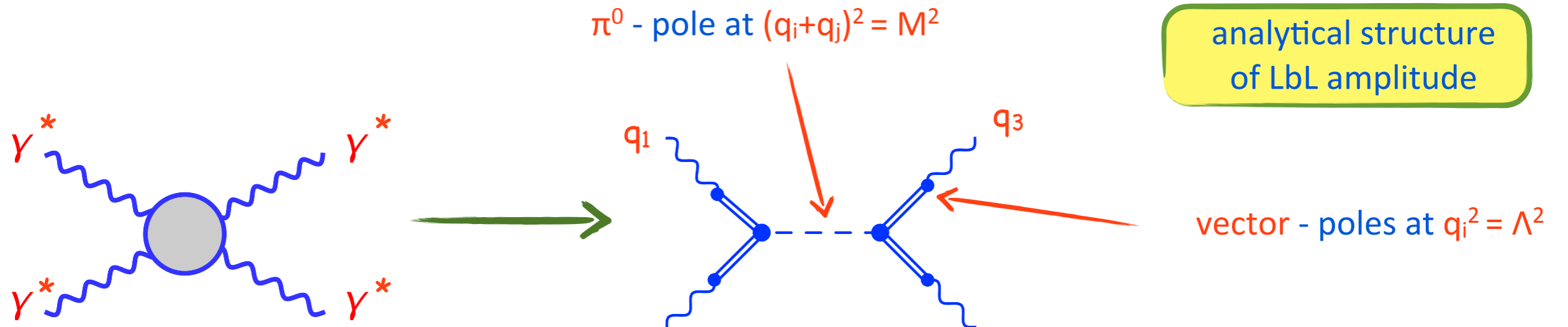
weighting functions (entire)

analytic structure \longrightarrow $\times \frac{\Pi_{\lambda_1 \lambda_2 \lambda_3 \lambda}(q_1, k - q_1 - q_2, q_2, k)}{q_1^2 q_2^2 (k - q_1 - q_2)^2 [(p + q_1)^2 - m^2] [(p + k - q_2)^2 - m^2]}$

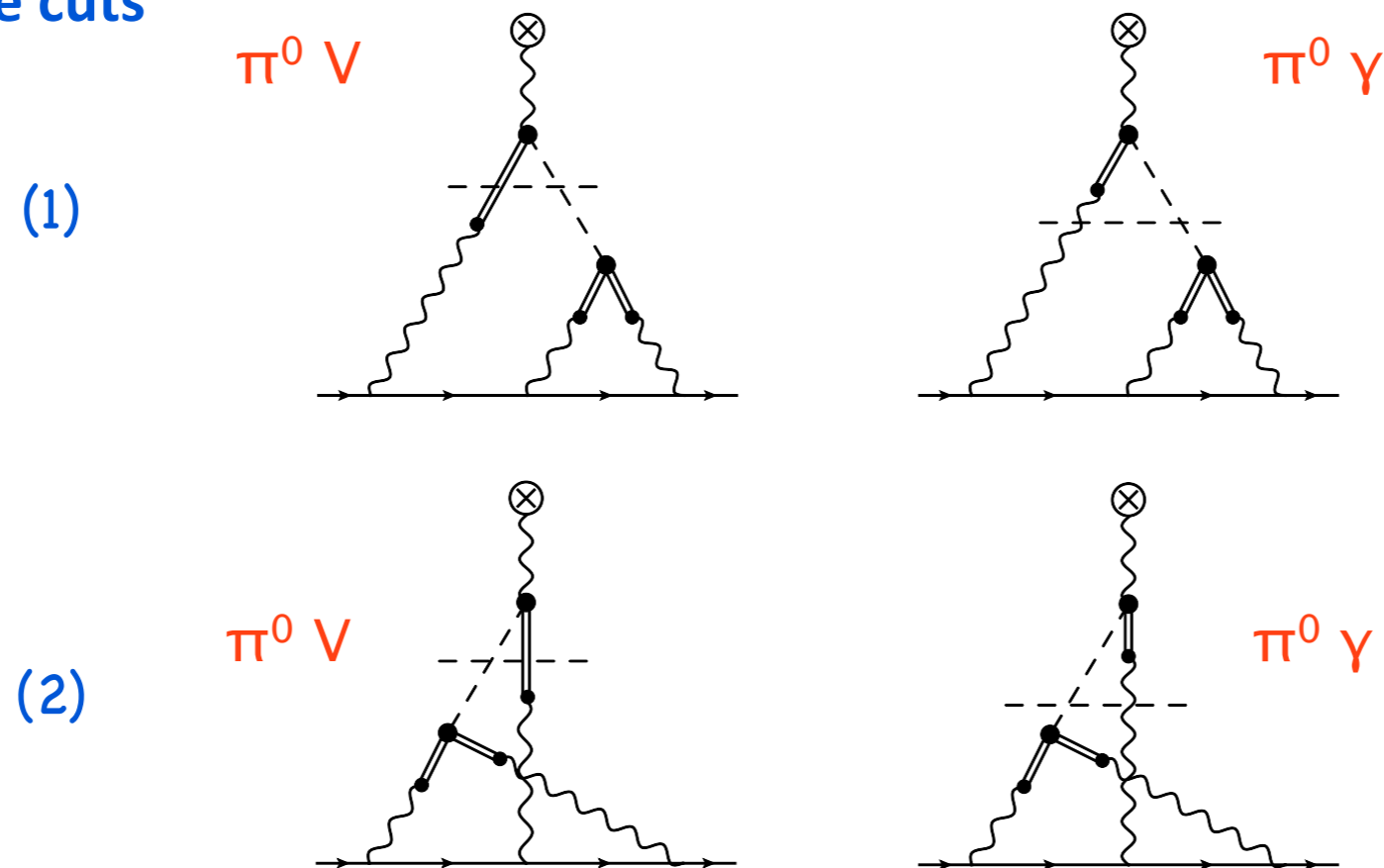
$$\Pi_{\lambda_1 \lambda_2 \lambda_3 \lambda_4}(q_1, q_2, q_3) = \epsilon^\mu(q_1, \lambda_1) \epsilon^\nu(q_2, \lambda_2) \epsilon^\lambda(q_3, \lambda_3) \epsilon^\rho(q_4, \lambda_4) \Pi_{\mu\nu\lambda\rho}(q_1, q_2, q_3)$$

dispersion relation approaches for a_μ (II)

➔ proof of principle: pole contributions

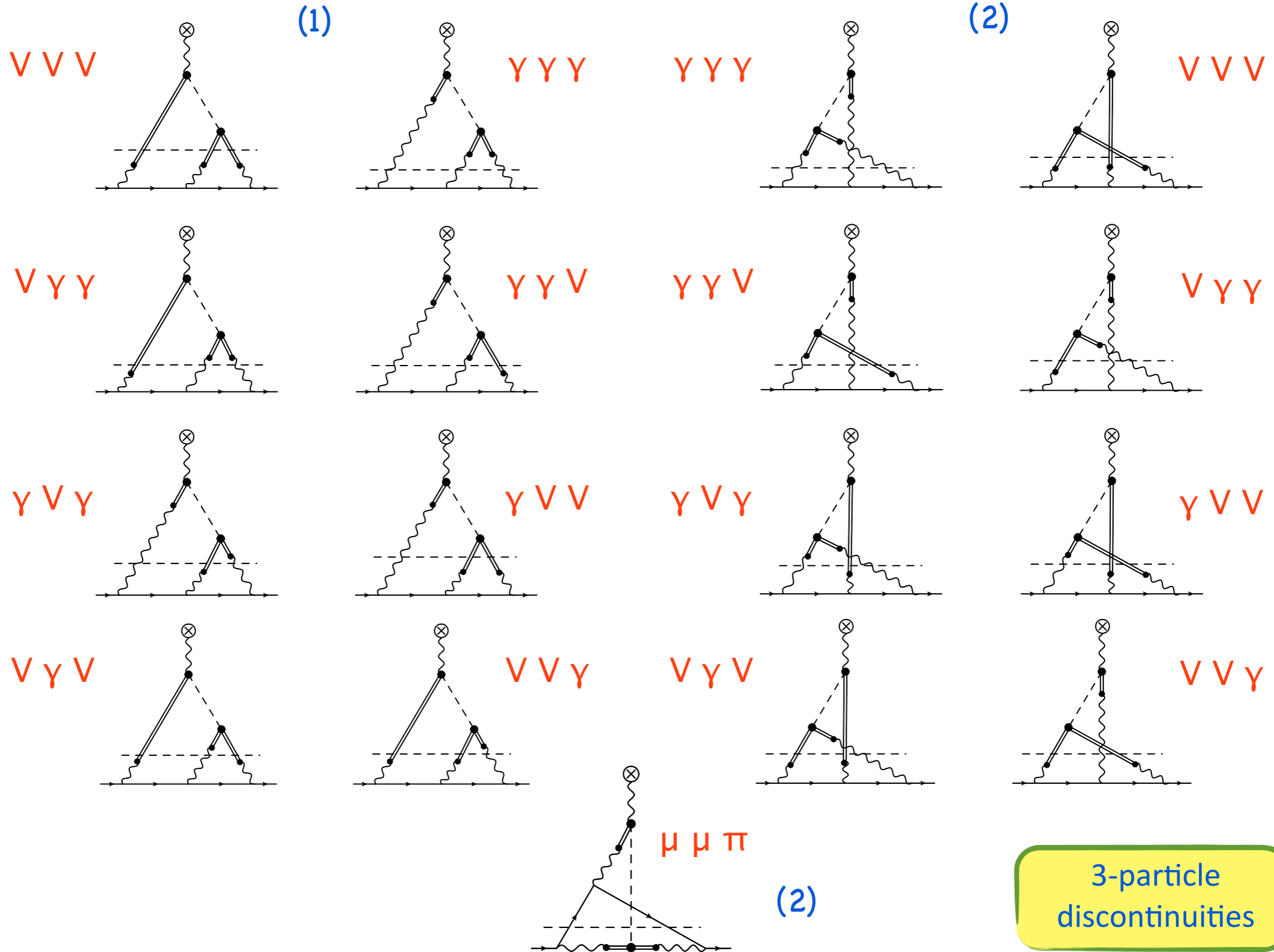


➔ 2-particle cuts



2-particle discontinuities

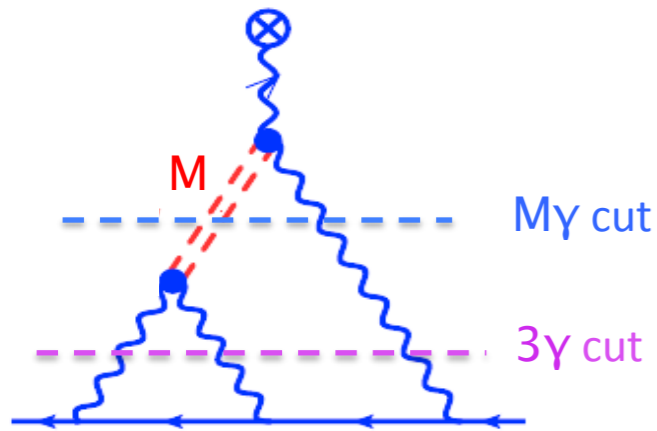
dispersion relation approaches for a_μ (III)



3-particle discontinuities

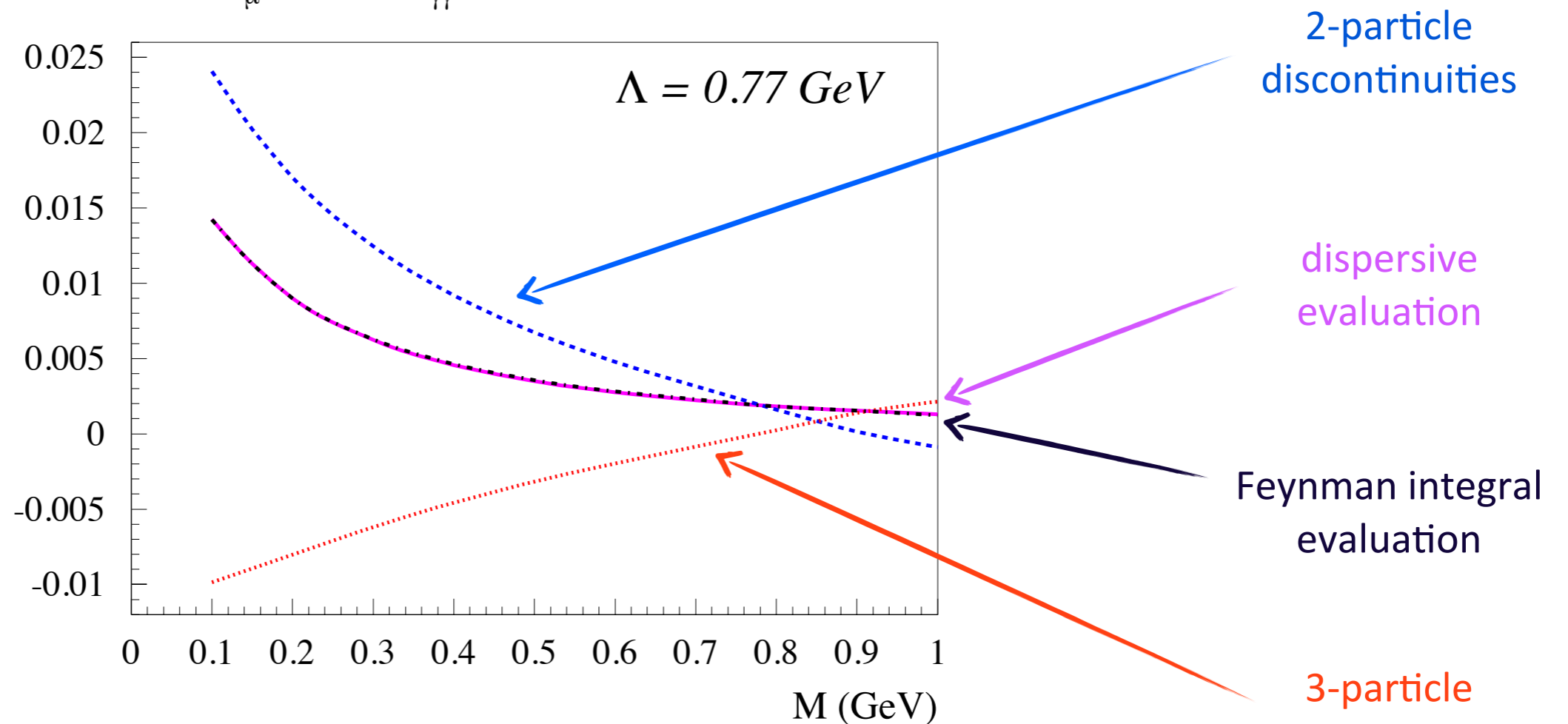
dispersion relation approaches for a_μ (IV)

reconstruction of a_μ from dispersion integral: **proof of principle**



$$F_2^{(i)}(0) = \frac{1}{2\pi i} \int_{M^2}^{\infty} \frac{dt}{t} \text{Disc}_t^{(2)} F_2^{(i)}(t) + \frac{1}{2\pi i} \int_0^{\infty} \frac{dt}{t} \text{Disc}_t^{(3)} F_2^{(i)}(t)$$

$a_\mu * M^3 / (\alpha \Gamma_{\gamma\gamma})$ (in GeV^2): diagram a



Pauk, Vdh (2014)

Summary and outlook

- ➔ **HVP:** new experimental program at BES-III
 - first results for $e^+ e^- \rightarrow \pi^+ \pi^- \gamma$
 - aim: hadronic cross sections to 1% accuracy
 - strong lattice effort to rival experimental accuracy
- ➔ **HLbL: new theoretical tools** for $\gamma^* \gamma^* \rightarrow X$
 - **sum rules, dispersive frameworks** for transition FFs: allow to include experimental constraints
 - new evaluation of **heavier meson contributions:** $a_\mu = (6.6 \sim 4.4) \pm 2.9 \times 10^{-11}$
 - pioneering lattice efforts
- ➔ **new dispersion relation** frameworks for **HLbL** to a_μ :
 - > require close collaboration with experiment (spacelike, timelike, meson decays)
- ➔ Outcome of Mainz workshop:
draft of roadmap for a data driven approach also in HLbL
- ➔ **goal: realistic error estimate on a_μ / reduce to 2×10^{-10} (20 % of HLbL)**