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**Explaining the LHC flavour  
anomalies**

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# Outline:

- Introduction: The LHC flavour anomalies
- 2HDM with gauged  $L_{\hat{1}} - L_{\hat{0}}$  and vector-like quarks
  - $B \rightarrow K^* \mu^+ \mu^-$
  - $B \rightarrow K \mu^+ \mu^- / B \rightarrow K e^+ e^-$
  - $B_s \rightarrow \phi \mu^+ \mu^-$
  - $h \rightarrow \tau \mu$
  - $\tau \rightarrow \mu \mu \mu$
- 2HDM and 3HDM with gauged horizontal charged (no vector-like quarks)
- Conclusions

# The LHC flavour anomalies

# $B \rightarrow K^* \mu\mu, B_s \rightarrow \phi\mu^+ \mu$ and $R(K)$

See talk of Wolfgang Altmannshofer

- 2-3  $\sigma$  deviation from the SM mostly in  $C5'$
- 2.6  $\sigma$  deviation from the theoretically rather clean SM expectation for  $R(K)$ .
- 3.5  $\sigma$  in  $B_s \rightarrow \phi\mu\mu$
- Contribution to  $C_9^{\mu\mu}$  but not  $C_9^{ee}$  gives simultaneously a good fit  more than 4.3  $\sigma$  better than SM
- Explanation:
  - Leptoquarks see talk of Sophie Renner and Gudrun Hiller
  - Extra dimensions C. Niehoff, P. Stangl, D. Straub, arXiv:1503.03865.
  - Flavour non-universal  $Z'$  W. Altmannshofer, et al. 1403.1269



$$h \rightarrow \tau\mu$$

- 2.4  $\sigma$  difference from zero

$$\text{Br}[h \rightarrow \mu\tau] = \left(0.89^{+0.40}_{-0.37}\right)\% \quad \text{CMS-PAS-HIG-14-005}$$

- Can be explained in the effective field theory approach by

$$Q_{e\phi}^{fi} = \ell_f \phi e_i \phi^\dagger \phi$$

R. Harnik, J. Kopp, and J. Zupan, 1209.1397.  
G. Blankenburg, J. Ellis, and G. Isidori, 1202.5704.  
S. Davidson and P. Verdier, 1211.1248.

- No dominant contribution from vector-like fermions

A. Falkowski, D. M. Straub, and A. Vicente, 1312.5329

 **Extended  
Higgs sector**

A. Dery, et. al. 1408.1371.  
M. D. Campos, et. al., 1408.1652.  
A. Celis, et. al., 1409.4439.  
D. Aristizabal Sierra and A. Vicente, 1409.7690.  
C.-J. Lee and J. Tandean, 1410.6803.  
J. Heeck, et. al., 1412.3671.

# 2HDM with vector-quarks

Andreas Crivellin, Giancarlo D'Ambrosio and Julian Heeck

Explaining  $B \rightarrow K^* \mu \mu$ ,  $B \rightarrow K \mu \mu$  /  $B \rightarrow K e e$  and  $h \rightarrow \tau \mu$  in a two-Higgs-doublet model with gauged  $L_\mu$ - $L_\tau$   
PRL, arXiv:1501.00993.

# Gauged $L_i - L_{\hat{o}}$

- Vectorial U(1) gauge group:

$$Q(e) = 0, Q(\mu) = 1, Q(\tau) = -1$$

- Anomaly free

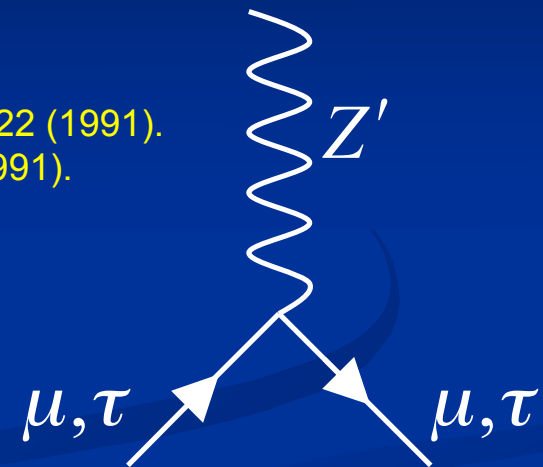
X. He, G. C. et al., Phys.Rev. **D43**, 22 (1991).  
R. Foot, Mod.Phys.Lett. **A6**, 527 (1991).

- Good zero order approximation to the PMNS matrix:

- maximal atmospheric and
- vanishing reactor neutrino mixing angle

$$M_\nu = \begin{pmatrix} X & 0 & 0 \\ 0 & 0 & Y \\ 0 & Y & 0 \end{pmatrix}$$

P. Binetruy, et al., hep-ph/9610481.  
N. F. Bell and R. R. Volkas, hep-ph/0008177.  
S. Choubey and W. Rodejohann, hep-ph/0411190.  
J. Heeck and W. Rodejohann, 1107.5238



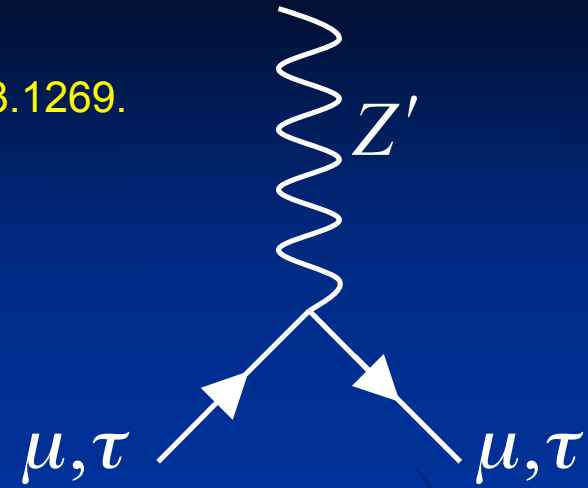
➔ Breaking necessary for a realistic neutrino sector

# The Model

W. Altmannshofer, S. Gori, M. Pospelov, and I. Yavin, 1403.1269.

- Gauged  $L_\mu - L_\tau$ :  $Z'$  boson with

$$-ig' \bar{l}_f \gamma^\mu Z'_\mu l_i \begin{pmatrix} 0 & 0 & 0 \\ 0 & 1 & 0 \\ 0 & 0 & -1 \end{pmatrix}_{fi}$$



- Vector-like quarks charged under  $L_\mu - L_\tau$

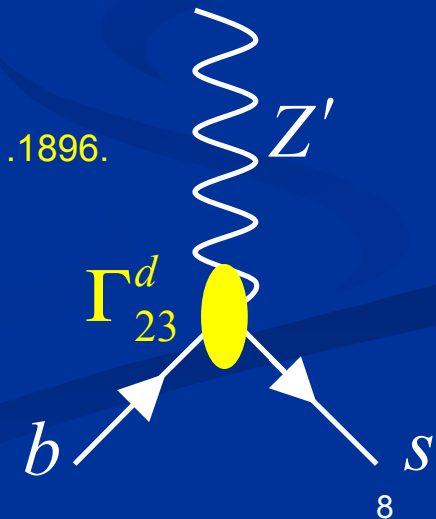
$$m_Q \bar{Q}_L \tilde{Q}_R + m_D \bar{D}_L \tilde{D}_R + m_U \bar{U}_L U_R + \text{h.c.}$$

- Effective  $Z'$  quark couplings

P. Langacker, 0801.1345., A. J. Buras, F. De Fazio, and J. Girrbach, 1211.1896.

$$ig' \gamma^\mu d_f \left( \Gamma_{fi}^L P_L + \Gamma_{fi}^R P_R \right) d_i Z'_\mu$$

$$\Gamma_{ij}^{dR} = -\frac{V_\Phi^2}{2m_D^2} (Y_i^D Y_j^{D*}), \quad \Gamma_{ij}^{dL} = \frac{V_\Phi^2}{2m_Q^2} (Y_i^Q Y_j^{Q*})$$





# 2<sup>nd</sup> Doublet breaks $L_i - L_{\hat{O}}$

J. Heeck, M. Holthausen, W. Rodejohann and Y. Shimizu, 1412.3671

- Two Higgs doublets

$$Q_{L_{\mu} - L_{\tau}}(\Psi_2) = 0 \quad Q_{L_{\mu} - L_{\tau}}(\Psi_1) = 2$$

- Yukawa couplings

$$L_Y \supset -\bar{\ell}_f Y_i^{\ell} \delta_{fi} \Psi_2 e_i - \xi_{\tau\mu} \bar{\ell}_3 \Psi_1 e_2 - \bar{Q}_f Y_{fi}^u \tilde{\Psi}_2 u_i - \bar{Q}_f Y_{fi}^d \Psi_2 d_i + \text{h.c.}$$

- Flavour changing SM-like Higgs coupling

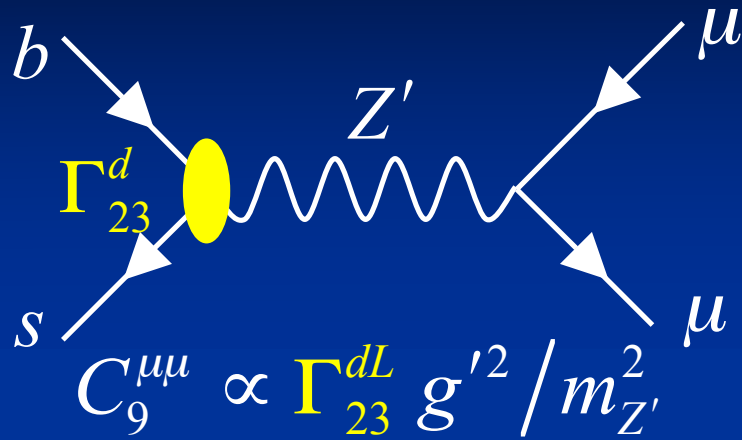
$$\Gamma_{\tau\mu}^h \bar{\tau} P_R \mu h^0 \approx \frac{m_{\tau}}{v} \frac{\cos(\alpha - \beta)}{\cos(\beta) \sin(\beta)} \theta_R \bar{\tau} P_R \mu h^0 \quad \begin{array}{l} \sin \theta_R \quad \frac{v}{\sqrt{2} m_{\tau}} \xi_{\tau\mu} \cos \beta \\ \sin \theta_L \quad 0 \end{array}$$

- Lepton flavour violating Z' couplings

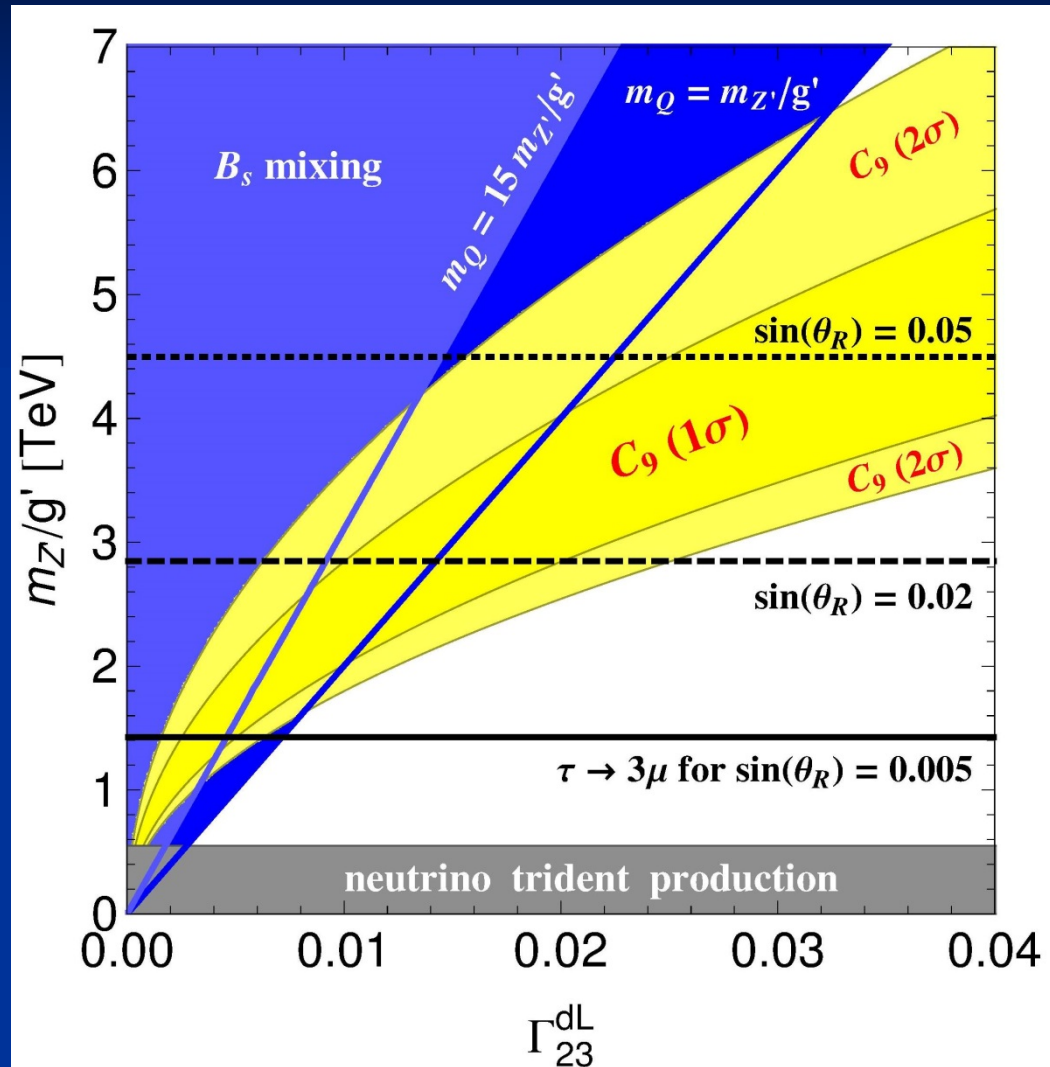
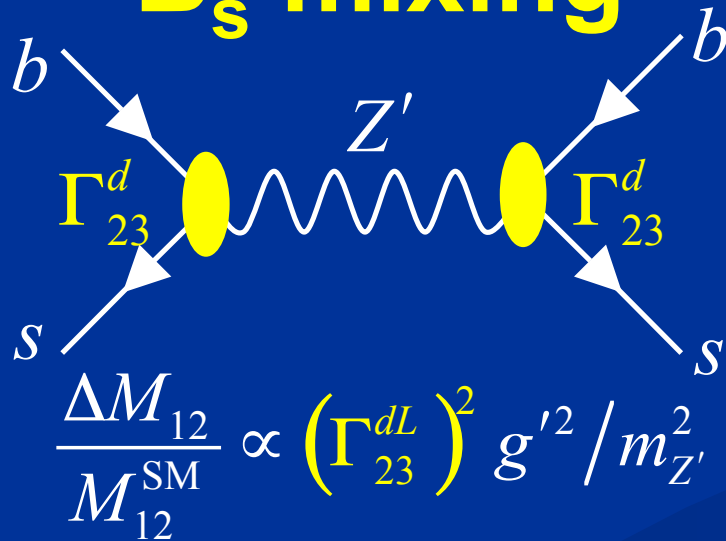
$$g' Z'(\bar{\mu}, \bar{\tau}) \begin{pmatrix} \cos 2\theta_R & \sin 2\theta_R \\ \sin 2\theta_R & -\cos 2\theta_R \end{pmatrix} \gamma^{\nu} P_R \begin{pmatrix} \mu \\ \tau \end{pmatrix}$$

$$m_D^2 \rightarrow \infty$$

$b \rightarrow s \mu \mu$



$B_s$  mixing

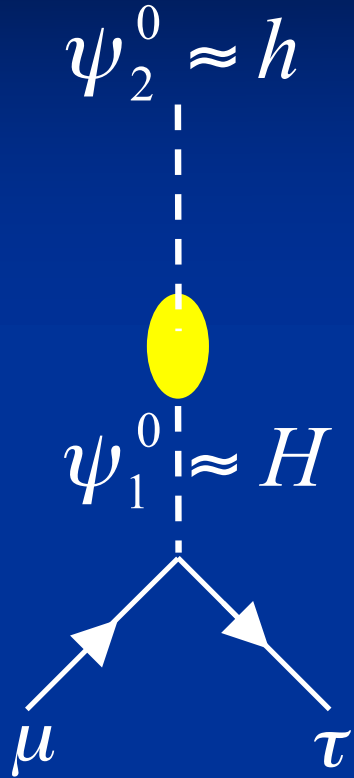


allowed regions

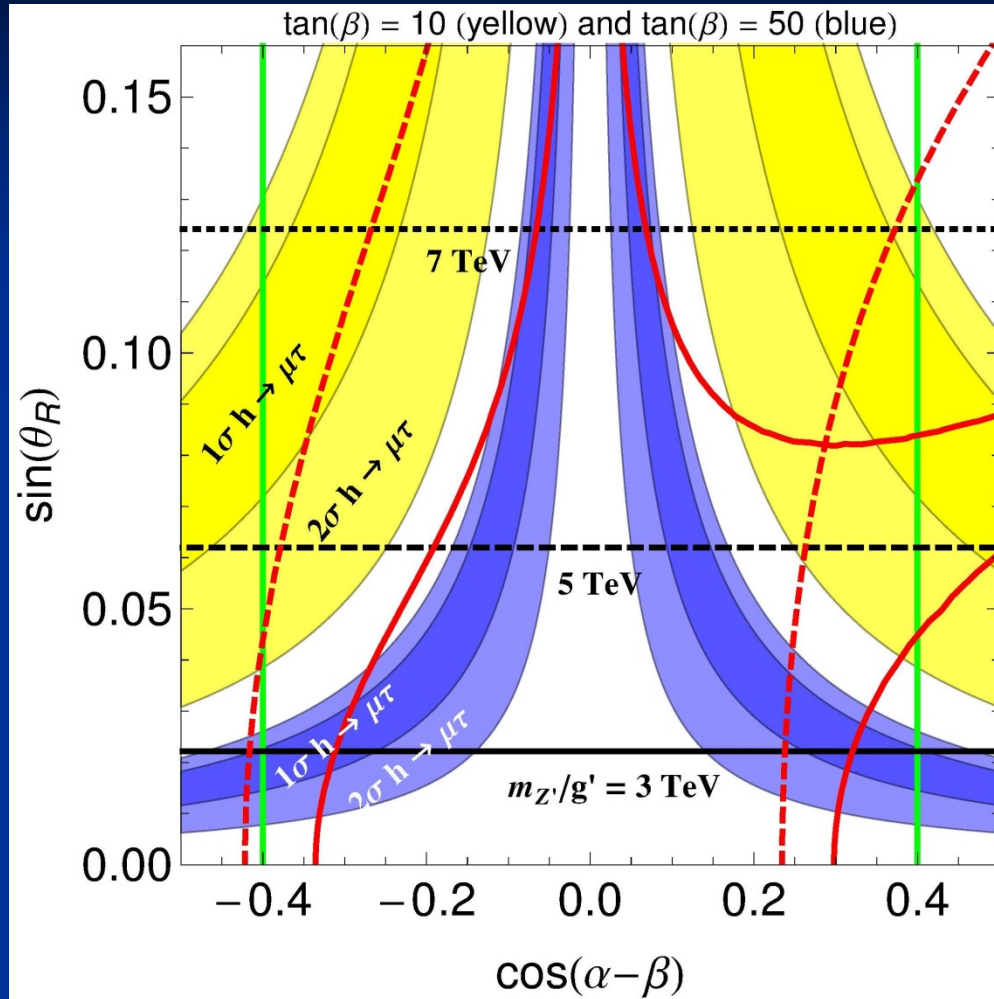
$h \rightarrow \mu\tau$

and

$\tau \rightarrow \mu\mu\mu$



$h \rightarrow \gamma\gamma$  etc.



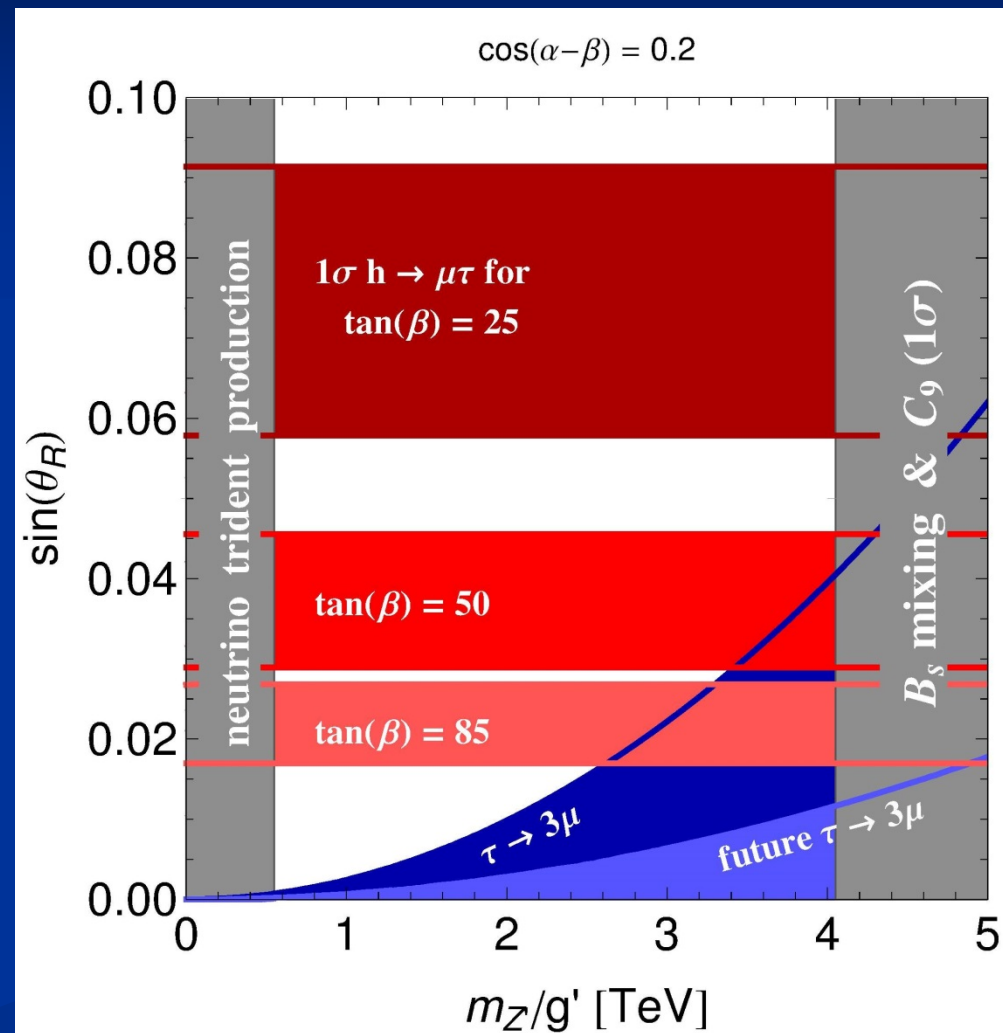
$h \rightarrow \tau\tau$

   $h \rightarrow \mu\tau$  ( $\tan\beta = 50$ )

   $h \rightarrow \mu\tau$  ( $\tan\beta = 10$ )

# $\tau \rightarrow \mu\mu\mu$ and $h \rightarrow \mu\tau$

- excluded
- allowed by  $h \rightarrow \tau\mu$
- allowed by  $\tau \rightarrow \mu\mu\mu$



# Horizontal charges

Andreas Crivellin, Giancarlo D'Ambrosio and Julian Heeck

Addressing the LHC flavour anomalies with horizontal gauge symmetries (arXiv:1503.03477)

# Charge assignment

- Avoid vector-like quarks by assigning charges to baryons as well

➔ same mechanism in the quark and lepton sector

- Use  $L_\mu - L_\tau$  in the lepton sector

➔ good symmetry for the PMNS matrix  
➔ effect in  $C_9^{\mu\mu}$  but not  $C_9^{ee}$

- First two quark generations must have the same charges because the large Cabibbo angle would lead to huge effect in Kaon mixing

- Anomaly free

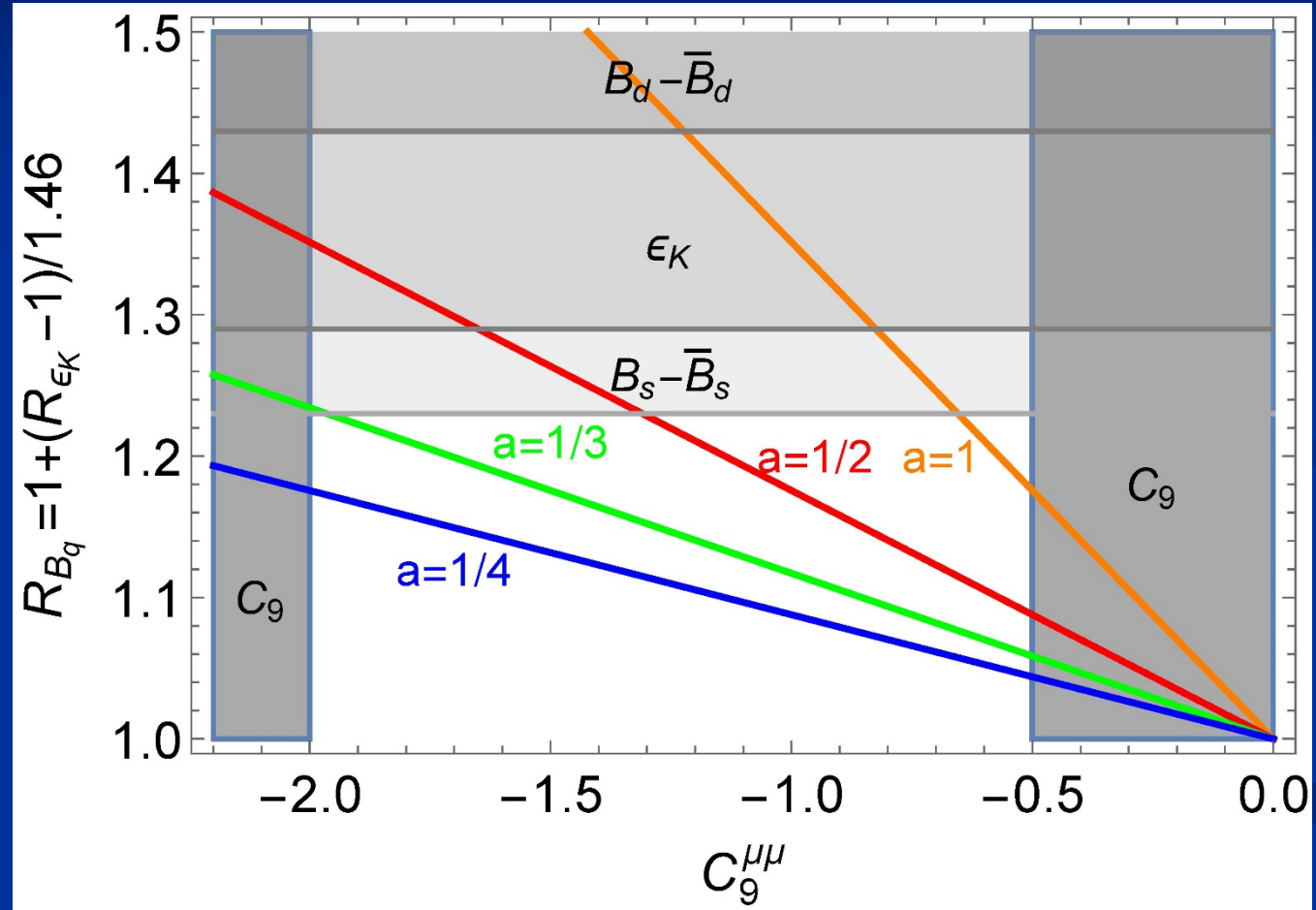
➔  $Q(B) = (-a, -a, 2a)$



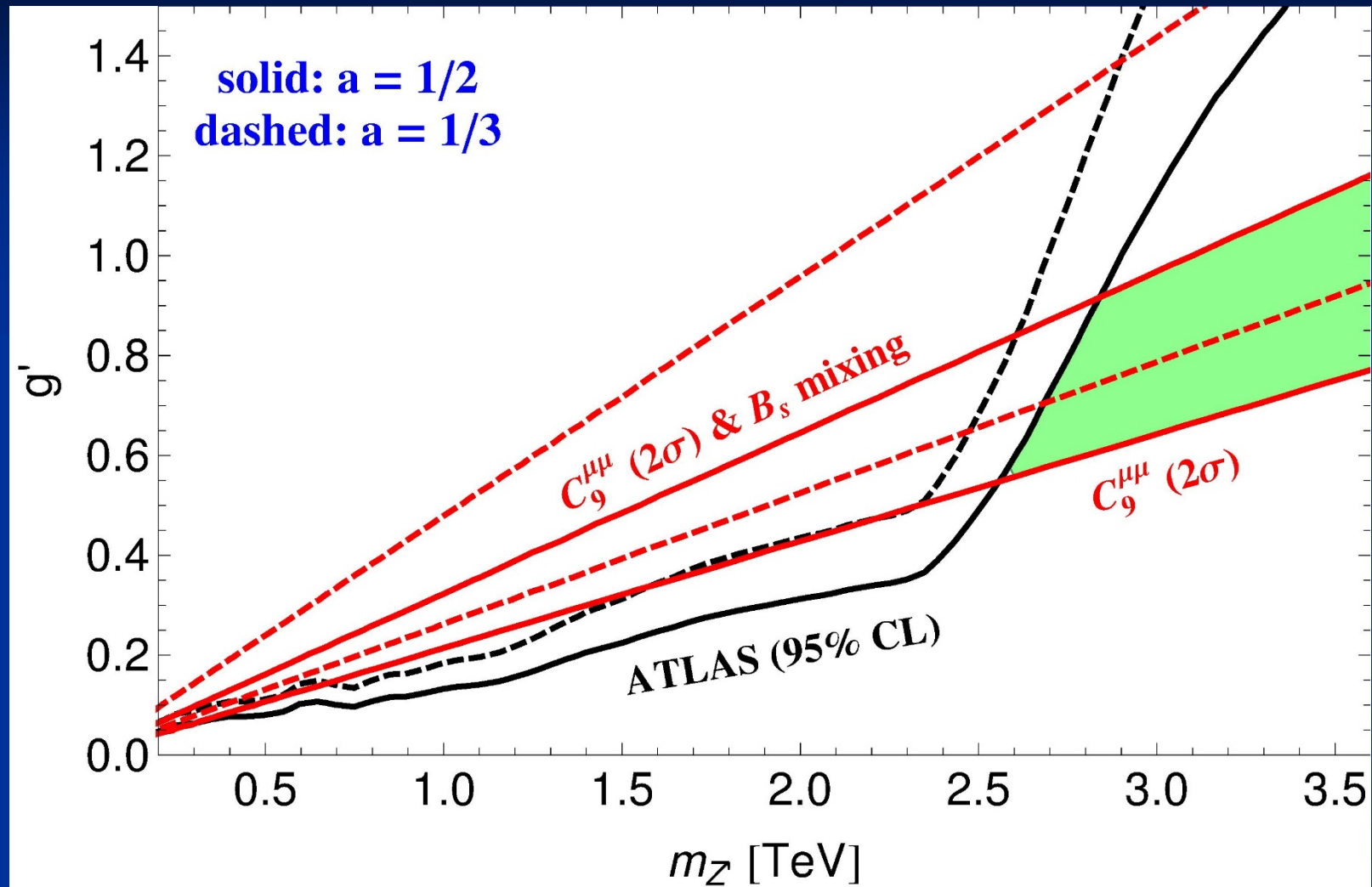
# $\Delta F=2$ : $Z'$ contribution

$$R_{B_q} = \frac{\Delta m_{B_q}}{\Delta m_{B_q}^{SM}}$$

$$R_{\epsilon_K} = \frac{\epsilon_K}{\epsilon_K^{SM}}$$



# LHC limits



ATLAS



$C_9^{\mu\mu} \& B_s - \bar{B}_s$



$a = 1/2$  allowed

# 3HDM

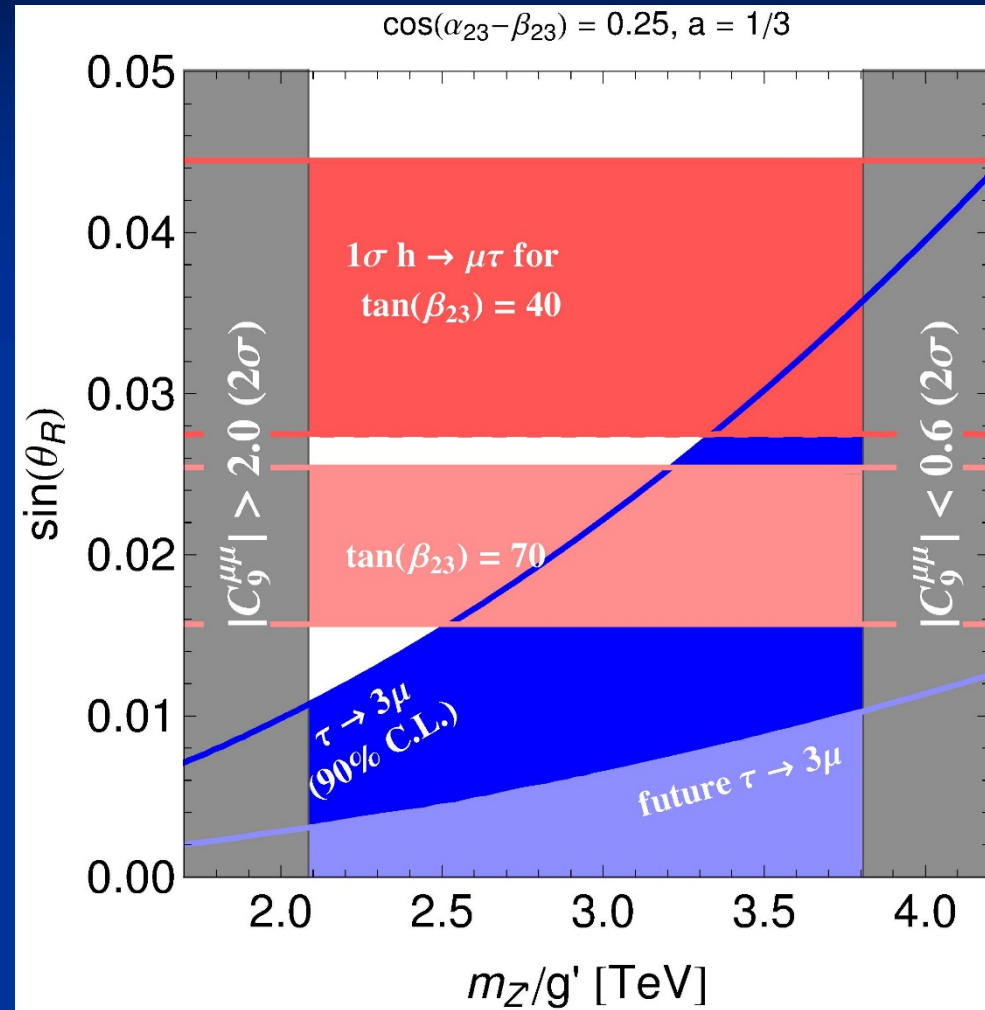
- Same effect in

$$\tau \rightarrow \mu\mu\mu$$

$$h \rightarrow \mu\tau$$

provided that the mixing among the doublets is small

- excluded
- allowed by  $h \rightarrow \tau\mu$
- allowed by  $\tau \rightarrow \mu\mu\mu$



# Conclusions

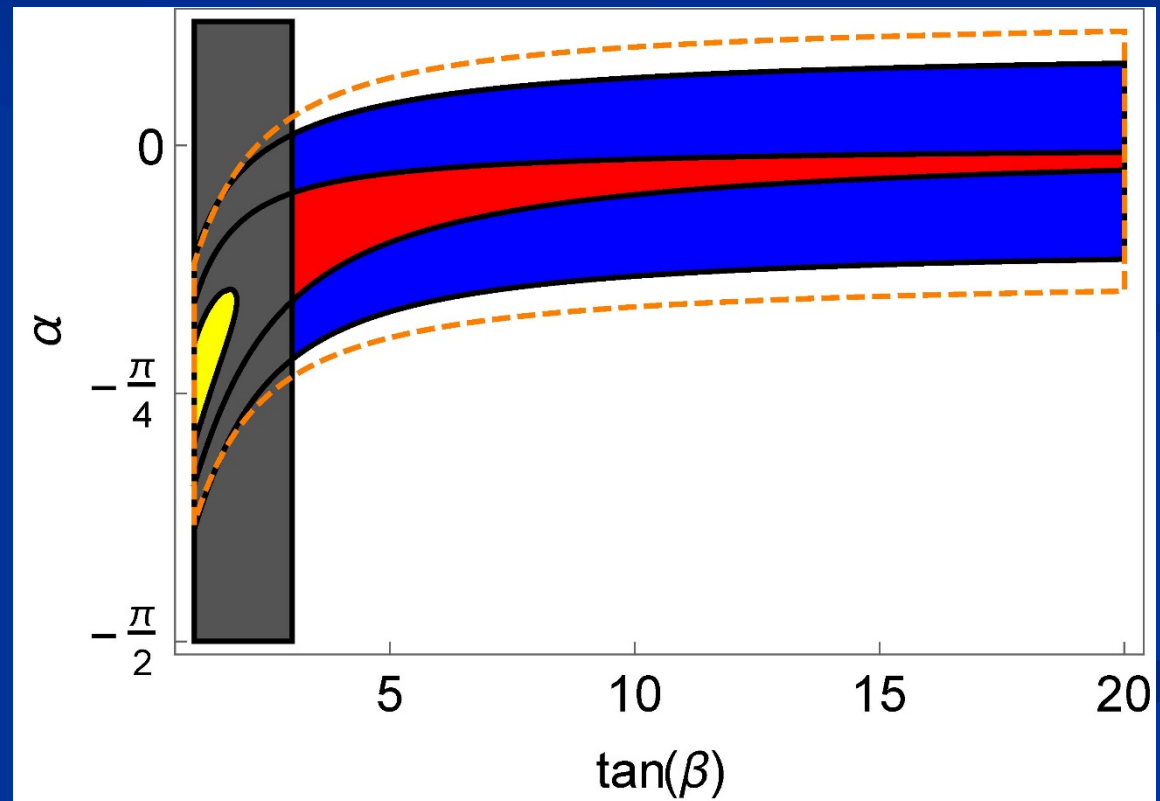
- The LHC found four anomalies in the flavour sector
  - $h \rightarrow \tau\mu$
  - $B \rightarrow K^* \mu^+ \mu^-$
  - $B \rightarrow K \mu^+ \mu^- / B \rightarrow K e^+ e^-$
  - $B_s \rightarrow \phi \mu^+ \mu^-$
- All four anomalies can be explained in a model with gauged  $L_i - L_\delta$ 
  - 2HDM with vector-like quarks
  - 3HDM with gauged flavour dependent B-L charges

Relation to LFUV in tauonic B decays?

# $\Delta F=2$ : Higgs contributions

$$m_H = 300 \text{ GeV}, \quad C_9^{\mu\mu} = -1.3$$

- $m_A = 350 \text{ GeV}$
- $m_A = 300 \text{ GeV}$
- $m_A = 250 \text{ GeV}$



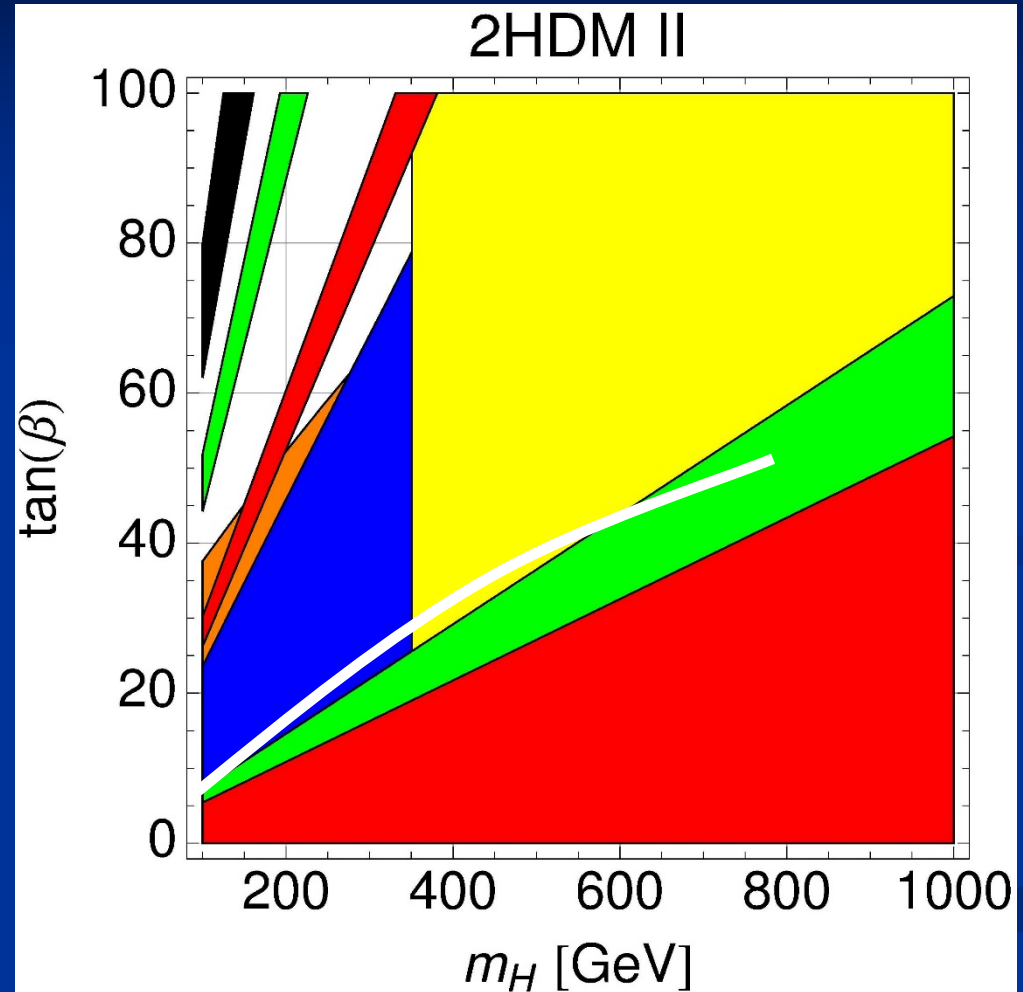
# Type-II 2HDM

■ Allowed

2 $\sigma$  regions from:

(superimposed)

- $b \rightarrow s\gamma$
- $B \rightarrow \tau\nu$
- $K \rightarrow \mu\nu / \pi \rightarrow \mu\nu$
- $B \rightarrow D\tau\nu$
- $B_s \rightarrow \mu^+\mu^-$
- $B \rightarrow D^*\tau\nu$
- LHC



➔ Tension from  $B \rightarrow D^*\tau\nu$



# $B \rightarrow D^{(*)} \tau \nu$ in the 2HDM III

- $B \rightarrow D^{(*)} \tau \nu$  and  $B \rightarrow D \tau \nu$  can be explained simultaneously using  $\epsilon_{32}^u$ .  $\rightarrow$  Check model via  $H^0, A^0 \rightarrow \bar{t}c$

Allowed regions from:

