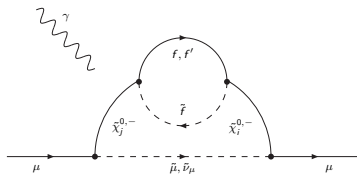


Magnetic moment $(g - 2)_\mu$ and Supersymmetry

Dominik Stöckinger

TU Dresden



Capri meeting “Flavour changing and conserving processes”,
September 2015

SM prediction too low by $\approx (30 \pm 8) \times 10^{-10}$

Can be due to new physics!

Plan:

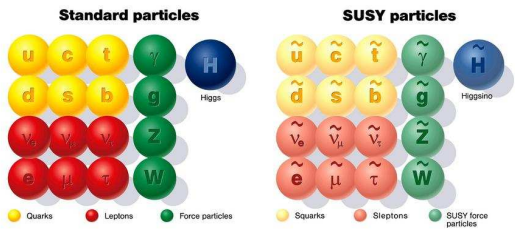
- General overview
- Specific topics in SUSY phenomenology/precision calculations

Outline

1 Overview: New Physics in General

SUSY well motivated

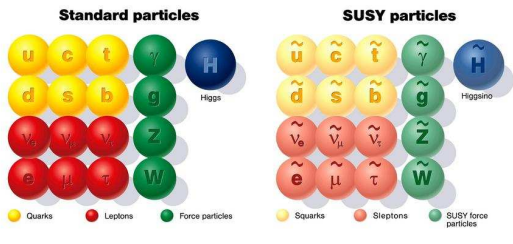
There must be new physics: Baryogenesis, ν -masses, dark matter...



SUSY well motivated

SUSY: maximal extension of Poincaré symmetry (Haag, Lopuszanski, Sohnius)

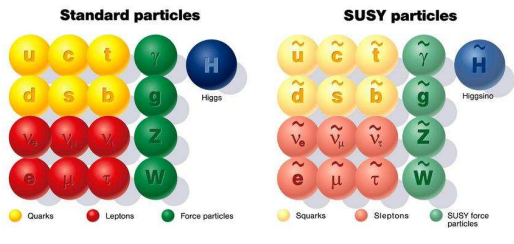
- improves quantum field theories
- connection to gravity
- explains existence of scalars



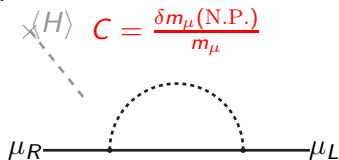
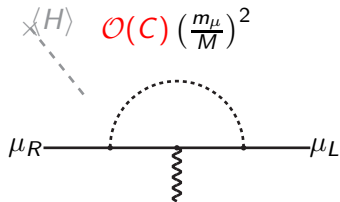
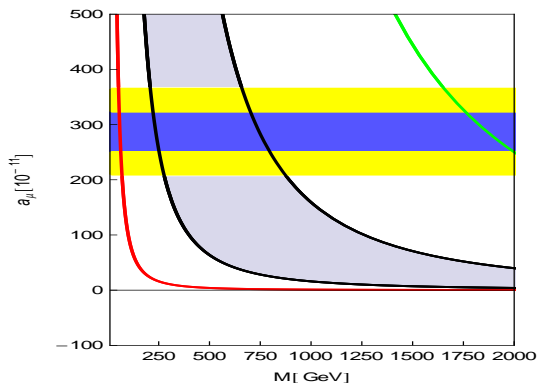
SUSY well motivated

SUSY at TeV-scale motivated ... and not excluded by LHC!

- can improve EW naturalness
- can generate/explain Mexican-hat potential and $M_h = 125$ GeV
- GUT
- dark matter



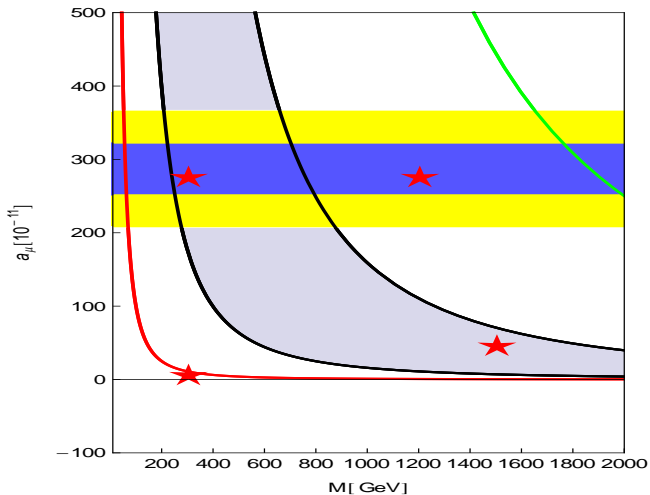
Very different contributions to a_μ : $C \sim \mathcal{O}(\alpha/4\pi) \dots \mathcal{O}(1)$



Outline

- 2 SUSY: Phenomenology
 - Illustrate wide range of possibilities
 - Novel, special scenarios

Plan: “Standard” and “Non-standard” scenarios



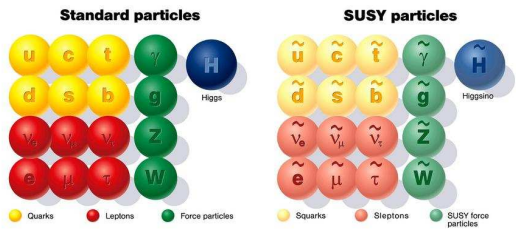
(SUSY is a framework not a concrete model! There is not “The SUSY prediction for a_μ ”, but prediction depends on scenario)

Standard SUSY and the MSSM

- MSSM: superpartners for all SM fields, two Higgs doublets
- free parameters: \tilde{p} masses and mixings, μ and $\tan\beta$

$$a_{\mu}^{\text{SUSY}} \approx 12 \times 10^{-10} \tan\beta \text{sign}(\mu) \left(\frac{100\text{GeV}}{M_{\text{SUSY}}} \right)^2$$

SUSY could be the origin of the observed $(30 \pm 8) \times 10^{-10}$ deviation!



Chirality flip and $\tan\beta$ -enhancement

Two Higgs and Higgsino doublets:

$$\tan\beta = \frac{\langle H_u \rangle}{\langle H_d \rangle}, \quad \mu = H_u - H_d \text{ transition}$$

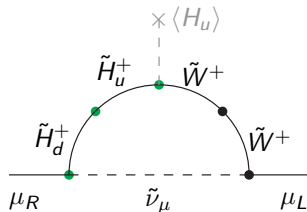
Large Yukawa as in 2HDM

$$m_\mu^{\text{tree}} = y_\mu \langle H_d \rangle$$

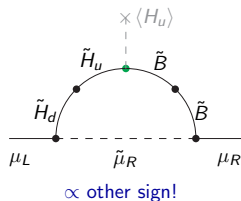
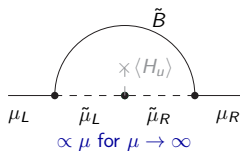
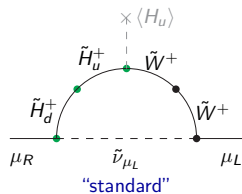
Diagram enhanced by Yukawa and large “other” vev

$$\propto y_\mu \langle H_u \rangle \mu = m_\mu \tan\beta \mu \quad \rightarrow a_\mu^{\text{SUSY}} \propto \tan\beta \text{sign}(\mu) \frac{m_\mu^2}{M_{\text{SUSY}}^2}$$

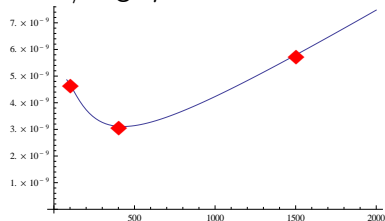
enhancement $\propto \tan\beta = 1 \dots 50$ (and $\propto \text{sign}(\mu)$)



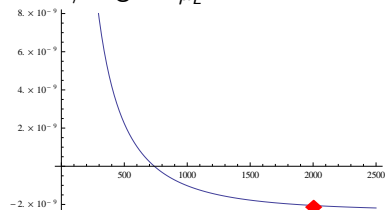
Further diagrams can sometimes dominate



small/large μ

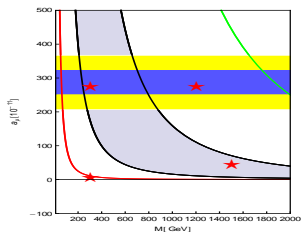


small/large $M_{\tilde{\mu}_L}$



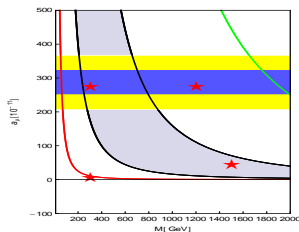
[Fargnoli, Gendiger, Passehr, DS, Stöckinger-Kim '13]

Hence: two examples within standard SUSY



Hence: two examples within standard SUSY

Obviously, the LHC rules out specific SUSY models as explanations of $g-2$.



e.g. “Constrained MSSM”.

Universality \Rightarrow masses strongly correlated

- Observed Higgs mass
- no gluinos/squarks at LHC
 \Rightarrow all sparticles heavy
- hence, a_μ contributions negligible

Hence: two examples within standard SUSY

Many other scenarios exist

many model building studies

[Endo, Hamaguchi, Ibe, Yanagida, D.P. Roy, ...]

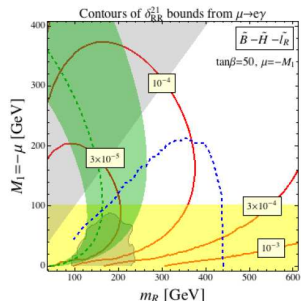
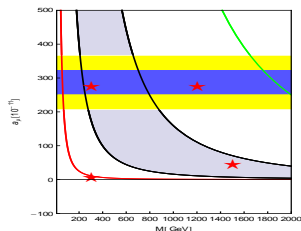
- large a_μ for $m_{\text{weak}} \sim 100 \dots 400$ GeV

[Fargnoli, Gendiger, Passet, DS, Stöckinger-Kim '13]

- evade LHC by $m_{\text{coloured}} \gg 1$ TeV
- or by compact spectra → talk by Patel

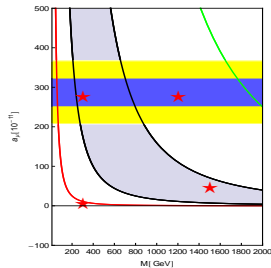
- can carry out dedicated LHC studies

[Endo; Yanagida; Roy; Calibbi; Roszkowski. ...] → talk by Paradisi



[Calibbi et al '15]

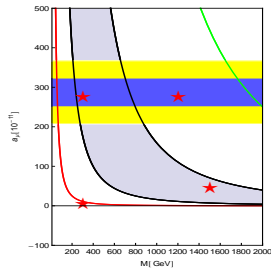
3rd Example: large a_μ for large masses



Is it possible to explain the a_μ deviation with **TeV-scale SUSY**?

What is the largest possible **SUSY** contribution to a_μ ?

3rd Example: large a_μ for large masses



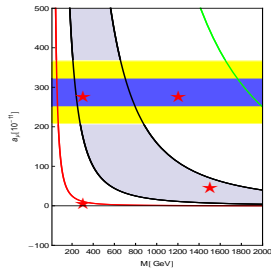
Idea: radiative muon mass in SUSY

$$m_\mu^{\text{tree}} = y_\mu v_d$$

set $v_d \rightarrow 0$, $\tan \beta \rightarrow \infty$

[Dobrescu, Fox '10][Altmannshofer, Straub '10] (see also [Davies, March-russell, Mccullough '11])

3rd Example: large a_μ for large masses



Idea: radiative muon mass in SUSY

$$m_\mu^{\text{full}} = y_\mu v_d +$$

$$+ \dots$$

set $v_d \rightarrow 0$, $\tan \beta \rightarrow \infty$

[Dobrescu, Fox '10][Altmannshofer, Straub '10] (see also [Davies, March-russell, McCullough '11])

Large a_μ in MSSM for $\frac{v_u}{v_d} = \tan \beta \rightarrow \infty$

[Bach, Park, DS, Stöckinger-Kim, '15]

“standard case” (equal masses, 1-loop)

$$m_\mu = v_d y_\mu \left(1 + \underbrace{\frac{v_u}{v_d} \Delta_\mu^{\text{red}}}_{\text{normally neglected}} \right)$$

$$a_\mu^{\text{SUSY}} \approx \underbrace{12 \times 10^{-10}}_{y_\mu \dots} \tan \beta \text{ sign}(\mu) \left(\frac{100 \text{ GeV}}{M_{\text{SUSY}}} \right)^2$$

Large a_μ in MSSM for $\frac{v_u}{v_d} = \tan \beta \rightarrow \infty$

[Bach, Park, DS, Stöckinger-Kim, '15]

actually, including higher order effects

[Marchetti, Mertens, Nierste, DS '08]

$$m_\mu = v_d y_\mu \left(1 + \underbrace{\frac{v_u}{v_d} \Delta_\mu^{\text{red}}}_{\text{normally neglected}} \right)$$

$$a_\mu^{\text{SUSY}} \approx \frac{12 \times 10^{-10} \tan \beta \text{ sign}(\mu)}{1 - 0.0018 \tan \beta \text{ sign}(\mu)} \left(\frac{100 \text{ GeV}}{M_{\text{SUSY}}} \right)^2$$

Large a_μ in MSSM for $\frac{v_u}{v_d} = \tan \beta \rightarrow \infty$

limit $\tan \beta \rightarrow \infty$

$$m_\mu = v_d y_\mu \left(1 + \underbrace{\frac{v_u}{v_d} \Delta_\mu^{\text{red}}}_{\text{normally neglected}} \right)$$

$$a_\mu^{\text{SUSY}} \approx -70 \times 10^{-10} \left(\frac{1000 \text{ GeV}}{M_{\text{SUSY}}} \right)^2$$

$\tan \beta$ and $\text{sign}(\mu)$ drop out, large contributions for $M_{\text{SUSY}} \sim \text{TeV}$!

Large a_μ in MSSM for $\frac{v_u}{v_d} = \tan \beta \rightarrow \infty$

limit $\tan \beta \rightarrow \infty$

$$m_\mu = v_d y_\mu \left(1 + \underbrace{\frac{v_u}{v_d} \Delta_\mu^{\text{red}}}_{\text{normally neglected}} \right)$$

$$a_\mu^{\text{SUSY}} \approx -70 \times 10^{-10} \left(\frac{1000 \text{ GeV}}{M_{\text{SUSY}}} \right)^2$$

“standard” case: sign wrong! But note: all masses were set equal!

Large a_μ in MSSM for $\tan \beta \rightarrow \infty$, different masses

[Bach,Park,DS,Stöckinger-Kim, '15]

Generally: $a_\mu^{\text{SUSY}} \rightarrow \frac{a_\mu^{\text{red}}}{\Delta_\mu^{\text{red}}}$

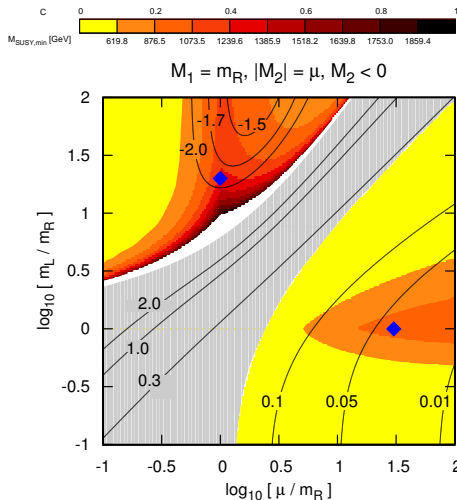
coloured: a_μ positive

Sample TeV-scale masses:

μ	M_1	M_2	m_L	m_R	$a_\mu/10^{-9}$
15	1	-1	1	1	3.01
1.3	1.3	-1.3	26	1.3	2.90

Experimental constraints ok: B-physics,

Higgs-physics, vacuum stability

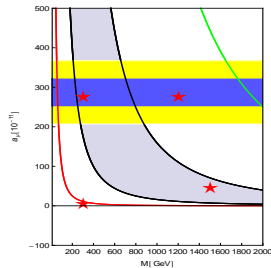


4th Example: small a_μ , small masses

Questions:

Can all SUSY scenarios give **large contributions**?

Try **opposite reaction to LHC**: not less, but
more SUSY, more symmetry, more light sparticles!



4th Example: small a_μ , small masses

Questions:

Can all SUSY scenarios give **large contributions**?

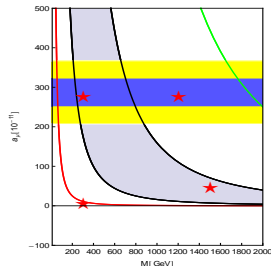
Try **opposite reaction to LHC**: not less, but **more SUSY, more symmetry, more light particles!**

R-symmetry, MRSSM [Kribs, Poppitz, Weiner '08]

- particles charged under conserved R-charge
- gauge bosons have two gauginos ($R = \pm 1$) and one scalar superpartner
- R-symmetry suppresses LHC cross sections, lighter sparticles viable

[Kribs, Martin '10]

- compatible with Higgs and LEP constraints [Diessner, Kalinowski, Kotlarski, DS '14]



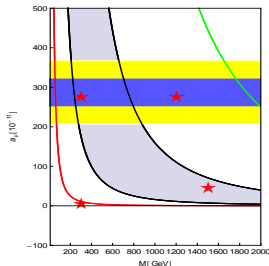
4th Example: small a_μ , small masses

Questions:

Can all SUSY scenarios give **large contributions**?

Try **opposite reaction to LHC**: not less, but
more SUSY, more symmetry, more light sparticles!

- kills $g - 2$ (no $\tan \beta$ enhancement)
- all $\tan \beta$ -enhanced contributions: $\propto \mu M_{1,2}$
- both μ and Majorana gaugino masses forbidden by R-symmetry
predictive!

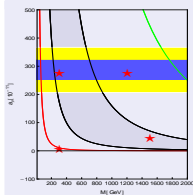


Complementarity/correlations

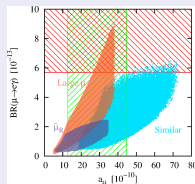
Strong correlations between observables are interesting

But they typically depend strongly on the scenario

all observables needed to uncover model details



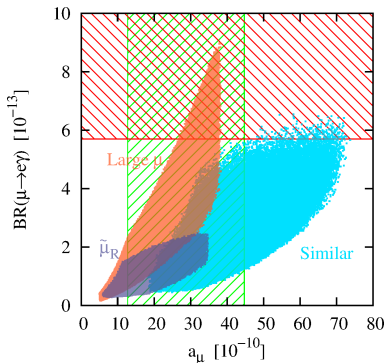
LHC \oplus a_μ : masses



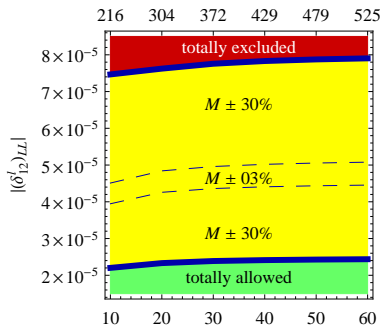
LFV \oplus a_μ : LFV parameters

$\mu \rightarrow e\gamma$: correlation with $g - 2$ depends on scenario

[Kersten, Park, DS, Velasco-Sevilla '14]
 M [GeV]



parameter regions where single diagrams dominate



bounds on δ_{LL} assuming M is fixed to accommodate a_μ

- study correlation for fixed $\delta_{LL} = m_{L12}^2 / \sqrt{m_{L11}^2 m_{L22}^2} = 2 \times 10^{-5}$
 - ▶ correlation often studied [Chacko, Kribs'01; Isidori, Mescia, Paradisi, Temes '07]
 - ▶ but depends on mass pattern [Kersten, Park, DS, Velasco-Sevilla '14]
- still, can constrain δ_{LL} assuming a_μ and mass pattern

Outline

3 SUSY: Precision

Some Numbers

$$M_W^{\text{exp}} = 80.385(15) \text{ GeV} \quad 0.02\%$$

$$\text{MSSM} \quad 1\text{Loop} \quad \mathcal{O}(0 \dots 10\sigma)$$

$$M_h^{\text{exp}} = 125.09(0.24) \text{ GeV} \quad 0.2\%$$

$$\text{MSSM} \quad 1\text{Loop} \quad > 30\sigma$$

$$a_\mu^{\text{exp}} = 11\,659\,208.9(6.3) \times 10^{-10} \quad 0.5\text{ppm}$$

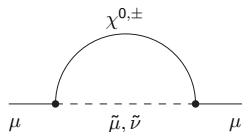
$$\text{MSSM} \quad \text{loop effects} \quad \mathcal{O}(0 \dots \pm 10\sigma)$$

Current SUSY theory uncertainty: 3×10^{-10}

Status of SUSY prediction

1-Loop

$$\propto \tan \beta$$



[Fayet '80],...

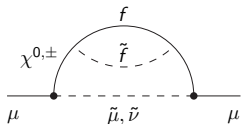
[Kosower et al '83],[Yuan et al '84],...

[Lopez et al '94],[Moroi '96]

complete

2-Loop (SUSY 1L)

$$\text{e.g. } \propto \log \frac{M_{\text{SUSY}}}{m_\mu}$$



[Degrassi, Giudice '98]

[Marchetti, Mertens, Nierste, DS '08]

[Schäfer, Stöckinger-Kim,

v. Weathershausen, DS '10]

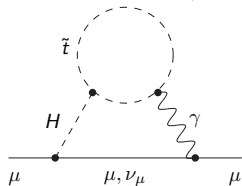
[Fagnoli, Gnendiger, Passehr, DS,

Stöckinger-Kim '13]

QED, $\tan^2 \beta$
 f, \tilde{f} -loops

2-Loop (SM 1L)

$$\text{e.g. } \propto \tan \beta \mu m_t$$



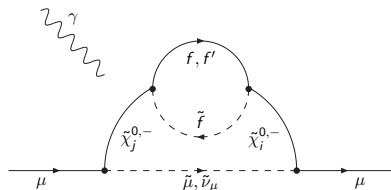
[Chen, Geng '01][Arhrib, Baek '02]

[Heinemeyer, DS, Weiglein '03]

[Heinemeyer, DS, Weiglein '04]

complete

The new contributions with $f\tilde{f}$ loops



Motivation:

- Split spectra / Big step towards full 2-loop calculation
- remaining class with dependence on squarks
- maximum complexity: 5 heavy + 2 light scales
- computed exactly, including renormalization
- contains large logs, $\Delta\rho$

Contributions involving $\Delta\rho$

One-loop ambiguity

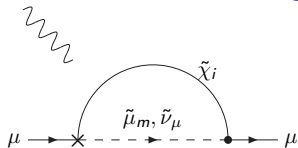
$$\left. \begin{aligned} a_{\mu}^{1L} &= \alpha(0) \dots = 29.4 \\ a_{\mu}^{1L} &= \alpha(M_Z) \dots = 31.6 \\ a_{\mu}^{1L} &= \alpha(G_F) \dots = 30.5 \end{aligned} \right\}$$

Fixed by full $2L\tilde{f}$ calculation

differ by $\Delta\alpha, \Delta\rho$: $2L\tilde{f}$ -terms

(for SPS1a, unit: 10^{-10})

Contributions involving $\Delta\rho$



$$\begin{aligned}
 &= a_\mu^{1L} \times \left(\dots + \frac{\delta(e^2/s_W^2)}{e^2/s_W^2} \right) \\
 &= a_\mu^{1L} \times \left(\Delta\alpha - \frac{c_W^2}{s_W^2} \Delta\rho + \dots \right)_{f, \tilde{f}\text{-loops}}
 \end{aligned}$$

One-loop ambiguity

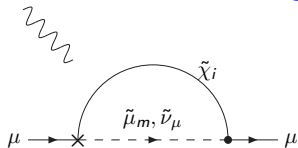
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Contributions involving $\Delta\rho$



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 &= a_\mu^{1L} \times \left(\Delta\alpha - \frac{c_W^2}{s_W^2} \Delta\rho + \dots \right)_{f, \tilde{f}\text{-loops}}
 \end{aligned}$$

One-loop ambiguity

$$\left. \begin{aligned}
 a_\mu^{1L} &= \alpha(0) \dots = 29.4 \\
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 a_\mu^{1L} &= \alpha(G_F) \dots = 30.5
 \end{aligned} \right\}$$

Fixed by full $2L\tilde{f}$ calculation

$$a_\mu^{1L+2L\tilde{f}} = 32.2$$

differ by $\Delta\alpha, \Delta\rho$: $2L\tilde{f}$ -terms

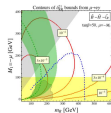
(for SPS1a, unit: 10^{-10})

Outline

4 Conclusions

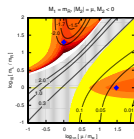
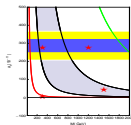
Summary: a_μ and SUSY

- New physics/SUSY not excluded and still well motivated



- a_μ^{SUSY} very model-dependent

- ▶ can be $\mathcal{O}(\pm 1 \dots 50) \times 10^{-10}$
- ▶ special SUSY scenarios $\tan\beta \rightarrow \infty$, R-symmetry



- New measurement will have strong impact

- ▶ constraints, model discriminator
- ▶ precise (2-loop) predictions important

