

Lattice Gauge Tensor Networks

Simone Montangero Institute for Complex Quantum Systems & Institute for Integrated Quantum Science and Technologies Ulm University

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Wo

Non-perturbative formulation of high-energy models

High Energy

Lattice Gauge Theories

Condensed Matter

Quantum Information

Strongly correlated quantum many body systems Exotic phases of matter to support QI processing

Traditional numerical methods (Quantum Monte Carlo,...)

High Energy

Quantum Technologies for Lattice Gauge Theories

Condensed, Matter Quantum Information

Tensor networks (classical simulations) AMO physics (quantum simulations)



Tensor Networks

Brief Introduction



U. Schwollöck, Rev. Mod. Phys. (2005)

"Is all that useful at all?"

Tensor network algorithms

Describe ground states (finite chemical potential) and real time dynamics of (1D, short range) many-body quantum systems.

$$\boldsymbol{H} = \sum_{i,j} \boldsymbol{H}_{i,j} + H(t)$$

Access to local quantities, correlators (space and time), entanglement $\langle \hat{O}_i \rangle = \langle \hat{O}_i \hat{O}_j \rangle = \text{Tr} |\Psi \rangle \langle \Psi |$

Extension to open quantum systems dynamics

Q. trajectories, MPO, and A. Werner et.al. 1412.5746

Optimal control of many-body quantum system $H = H_D + \lambda(t)H_C \qquad \max_{\lambda(t)} F(\lambda, H, \dots)$

P. Doria, T. Calarco, SM PRL (2010), S. Lloyd, SM PRL (2014)

Applications



Quantum Phase Transition dynamics

Light-harvesting

dynamics





Entanglement/Squeezing manipulation



Optimal experimental protocols

"How is all that possibly working?"

Area laws

Quantum correlations

 $S = \text{Tr}\{\rho_A \log \rho_A\}$ $S \propto N^{D-1}(\log N)$

Representable entanglement in MPS states

 $S \propto \log m$ $m \propto O(Poly(N))$

Ground states of local Hamiltonians obey area laws

J.Eisert, M. Cramer, M.B. Plenio Rev. Mod. Phys. 2009

Time evolution

Highly excited dynamics:



P. Calabrese and J. Cardy, J. Stat. Mech. (2005)

+ localization effects:

 $S \propto \log(t)$

G. De Chiara, SM, P. Calabrese, R. Fazio, J.Stat.Mech. (2006)

Adiabatic or quasi-adiabatic dynamics:

 $S \sim \mathrm{const}$

What can be simulated can be controlled

S. Lloyd, SM PRL (2014)



Lattice Gauge Tensor Networks

Models

Local degrees of freedom



- Kogut-Susskind Hamiltonian formulation of LGT
- Condensed matter models
- Benchmarking of quantum simulations

- Charge field: spinless fermions/bosons
- Electric field: two-level system (qubit)

Effective move:

Matter jumps right/left, gauge field flips to the left/right.



$$H_{\text{int}}^{[\text{QED}]} = \sum_{x,\vec{a}} \psi_x^{\dagger} \,\sigma_{x,x+\vec{a}}^+ \,\psi_{x+\vec{a}} + (-1)^q \,\psi_x \,\sigma_{x,x+\vec{a}}^- \,\psi_{x+\vec{a}}^{\dagger}$$

Simplest example



$$\rho - \vec{\nabla} \cdot \vec{E} = 0 \quad \Longrightarrow \quad G_x |\varphi_{\rm phys}\rangle = \left(\psi_x^{\dagger} \ \psi_x - \frac{1}{2} \sum_{\vec{a}} \sigma_{x,x+\vec{a}}^z\right) |\varphi_{\rm phys}\rangle = 0$$

Lattice gauge generator (Gauss' law)



Dynamics commutes with the gauge symmetry

$$\left[H_{\text{int}}^{[\text{QED}]}, G_x\right] = 0 \quad \forall x$$



Quantum link formulation

Schwinger representation

 $U_{x,x+\vec{a}} \longrightarrow c_{x,\vec{a}} c^{\dagger}_{x+\vec{a},-\vec{a}}$

Schwinger represenation





Link operator

Electric field [U(1) generator]



 $U_{x,y} \equiv S_{x,y}^+ = c_y^\dagger c_x$

 $\{c_x, c_y^{\dagger}\} = \delta_{x,y}$ Schwinger fermions (rishons) $[c_x, c_y^{\dagger}] = \delta_{x,y}$ Schwinger bosons

Spin representation:

$$N_{x,y} = c_y^{\dagger} c_y + c_x^{\dagger} c_x \qquad \left[\vec{S}_{x,y}\right]^2 \equiv \frac{N_{x,y}}{2} \left[\frac{N_{x,y}}{2} + 1\right]$$

Spin-¹/₂: E=1/2 → E=-1/2 → Spin-1: E=0 -E=+1

Quantum link + Tensor Networks

- Standard approach with additional symmetry
- Gauge satisfied via local Hilbert space definition (also non abelian):

Reduced Hilbert space size

 Additional Abelian symmetry (rishons per site) that can be exploited for computational speedup.

Speed up increase with theory complexity

Gauge invariant TN

Gauge invariant local bases

$$G_x|i\rangle = 0$$

$$|j_x\rangle_r = \sum_{s_\psi, s_R, s_L} A^{[x]j}_{s_R, s_\psi, s_L} |s_R, s_\psi, s_L\rangle_x$$

Projection over rishons number

$$Q_{x,x+1} = \sum_{s_R,s_L,q} B_{s_L,q}^{[x]} C_{q,s_R}^{[x+1]} |s_L\rangle \langle s_L|_x \otimes |s_R\rangle \langle s_R|_{x+1} \qquad F_{j_x,j_x'}^{[x]q_{x-1},q_x} = \delta_{j_x,j_x'} \cdot Z^{[x]q_{x-1},j_x} \cdot V^{[x]q_x,j_x}$$



g=1.00 g=1.25 g=1.50 0.5 × 50 × 50 × 50 -0.5 -1 time time time time time time 4 6

Numerical Results

QEDsinkingeden incarto, ed.



 $|\mu| \gg t$

QEBhasendigenang41)-d



T. Pichler, E. Rico, M. Dalmonte, P. Zoller, SM PRL (2014)

Real-time dynamics

Spin 1 representation

$$H = -t \sum_{x} \left(\psi_{x}^{\dagger} S_{x,x+1}^{+} \psi_{x+1} + \psi_{x+1}^{\dagger} S_{x,x+1}^{-} {}^{\dagger} \psi_{x} \right)$$
$$+ m \sum_{x} (-1)^{x} \psi_{x}^{\dagger} \psi_{x} + \frac{g^{2}}{2} \sum_{x} (S_{x,x+1}^{3})^{2},$$



Exhibits string-breaking

T. Pichler, E. Rico, M. Dalmonte, P. Zoller, SM 1505.???

Gauge invariant Hilbert space



Local states

Many-body classical states



Real-time dynamcis



Electric field





m=0, t=1

State diagram







Mesons scattering

real-time simulation



against gauge-invariant errors in quantum simulation platforms

$$H_{I} = \xi \sum_{x} n_{x} \left(S_{x-1,x}^{z} + S_{x,x+1}^{z} \right)$$

Conclusions & Outlook

- Versatile tools that enables real-time lattice gauge investigations (1+1d)
- Non abelian LGT: SU(2) in progress
- Tensor network extensions in 2+1d
- Condensed matter models
- Real-time scattering dynamics

Thank you for your attention!

Tommaso Calarco

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Daniel Jaschke

Pietro Silvi

Simplified access

JIV

 Future and Emerging Technologies (FET) - Skills and career development (Marie Curie) Research infrastructures

Common rules, toolkit of funding schemes

Dissemination & kno

ICT

Funds:



SFB/TRR21 Co.Co.Mat.

IP-SIQS IP-DIADEMS IP-RYSQ





Numerics:



GRiD

Enrique Rico

Peter Zoller



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Simone Montangero - ICQ & IQST, Ulm University