



UNIVERSITÀ
DEGLI STUDI
DI PALERMO

UNIVERSITÀ DELLA CALABRIA



Dipartimento di FISICA



LANDAUER'S PRINCIPLE IN MULTIPARTITE OPEN QUANTUM SYSTEM DYNAMICS

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- INFN - Gruppo collegato di Cosenza

in collaboration with: **F. Ciccarello, R. McCloskey, M. Paternostro & G. M. Palma**

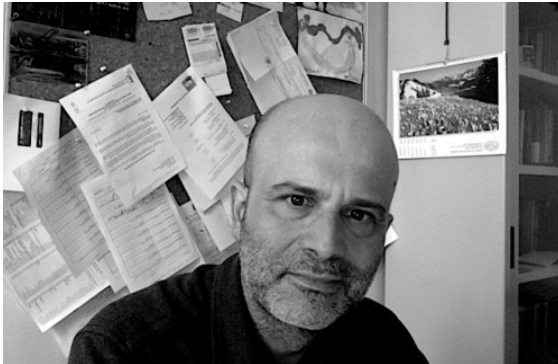
IQIS - 2015, Monopoli, 12 September 2015

in collaboration with

University of Palermo



F. Ciccarello



G. M. Palma



R. McCloskey



M. Paternostro

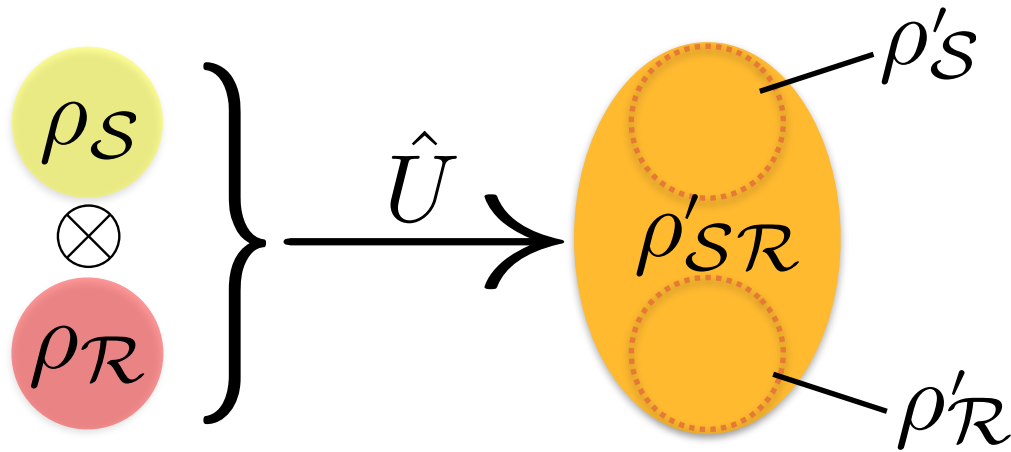
Queen's University Belfast

Outline

- Landauer's Principle
- Standard Quantum Collision Model
- Landauer's Principle for Fluxes
 - Recycling environment
 - Indirect erasure scheme
- Conclusions

Landauer's Principle

D. Reeb and M. Wolf, New J. Phys. 16, 103011 (2014)



β inverse temperature of \mathcal{R}

ΔQ heat transferred to \mathcal{R}

$\Delta \tilde{S}$ entropy decrease of \mathcal{S}

$$\rho_{SR} = \rho_S \otimes \rho_R$$

$$\rho_R = e^{-\beta H} / \text{Tr}[e^{-\beta H}]$$

$$\rho'_{SR} = \hat{U} \rho_{SR} \hat{U}^\dagger$$

$$\Delta \tilde{S} = S(\rho_S) - S(\rho'_S)$$

$$\Delta Q = \text{Tr}[H \rho'_R] - \text{Tr}[H \rho_R]$$

R. Landauer, IBM J. of Res. and Develop. 5 183 (1961)

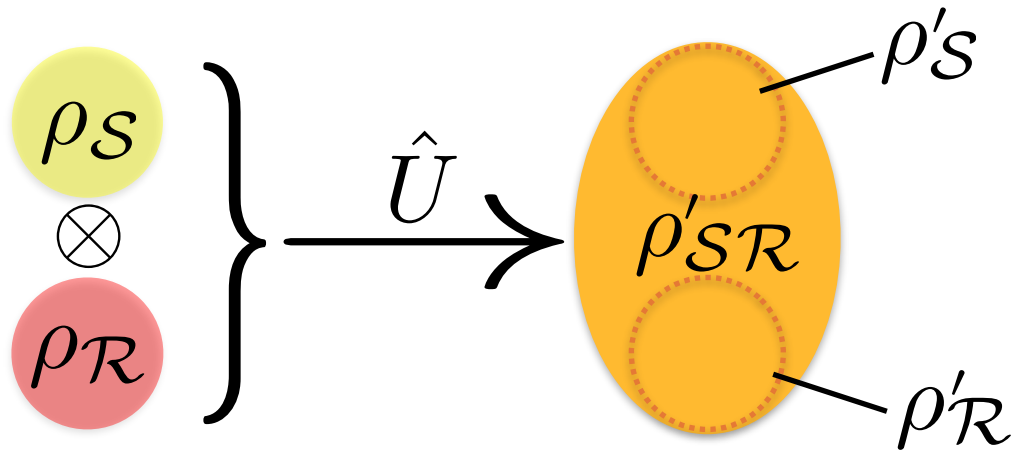
A. Berut, Nature 483, 187 (2012)

J.P.P. Silva et al., arXiv: 1412.6490

D. Reeb and M. Wolf, New J. Phys. 16, 103011 (2014)

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$$\rho'_{SR} = \hat{U} \rho_{SR} \hat{U}^\dagger$$

$$\beta \Delta Q \geq \Delta \tilde{S}$$

no Hamiltonian, no temperature for \mathcal{S} needed

$$\Delta \tilde{S} = S(\rho_S) - S(\rho'_S)$$

$$\Delta Q = \text{Tr}[H \rho'_R] - \text{Tr}[H \rho_R]$$

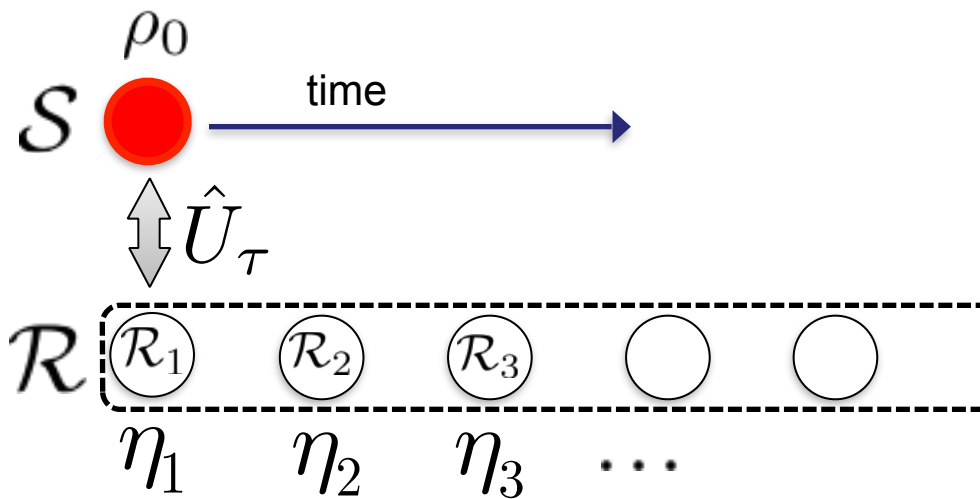
R. Landauer, IBM J. of Res. and Develop. 5 183 (1961)

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J.P.P. Silva et al., arXiv: 1412.6490

D. Reeb and M. Wolf, New J. Phys. 16, 103011 (2014)

Quantum Collision Model



τ =collision time

g =coupling strength

$$\hat{U}_\tau = e^{-igV\tau}$$

$$V = \sum_k \hat{S}_k \otimes \hat{R}_k$$

$$\begin{aligned} \rho_\tau &= \text{Tr}_{\mathcal{R}_1} [\hat{U}_\tau (\rho_0 \otimes \eta_1) \hat{U}_\tau^\dagger] = \mathcal{E}_1^\tau [\rho_0] \\ \rho_{2\tau} &= \text{Tr}_{\mathcal{R}_2} [\hat{U}_\tau (\rho_\tau \otimes \eta_2) \hat{U}_\tau^\dagger] = \mathcal{E}_2^\tau [\rho_\tau] \\ &\vdots \\ \rho_{n\tau} &= \text{Tr}_{\mathcal{R}_n} [\hat{U}_\tau (\rho_{(n-1)\tau} \otimes \eta_n) \hat{U}_\tau^\dagger] = \mathcal{E}_n^\tau [\rho_{(n-1)\tau}] \end{aligned}$$

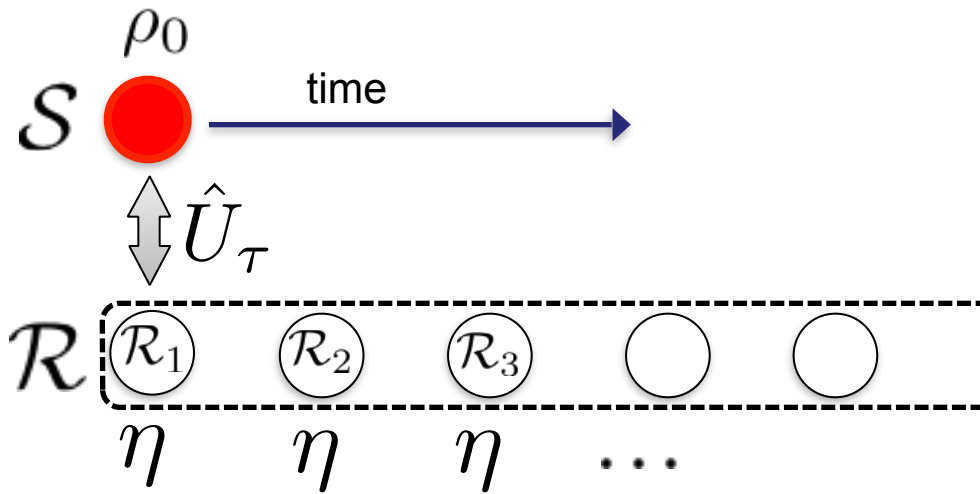
time

J. Rau, Phys. Rev. **129**, 1880 (1963)

M. Ziman et al., PRA **65**, 042105 (2002); V. Scarani et al., PRL **88**, 97905 (2002)

M. Ziman, P. Stelmachovic, and V. Buzek, Open Syst. and Inform. Dynam. **12**, 81 (2005)

Quantum Collision Model



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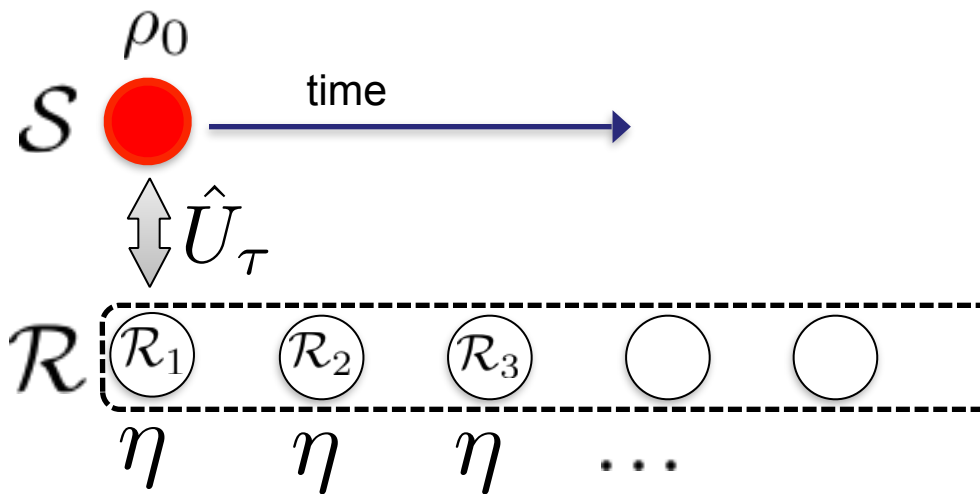
$$V = \sum_k \hat{S}_k \otimes \hat{R}_k$$

$$\eta_k = \eta \quad \forall k$$

$$\mathcal{E}_k^\tau := \mathcal{E}^{k\tau} \rightarrow \mathcal{E}_k^\tau \mathcal{E}_j^\tau = \mathcal{E}^{(k+j)\tau}$$

One parameter
semigroup

Quantum Collision Model



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 g = coupling strength

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$$\eta_k = \eta \quad \forall k$$

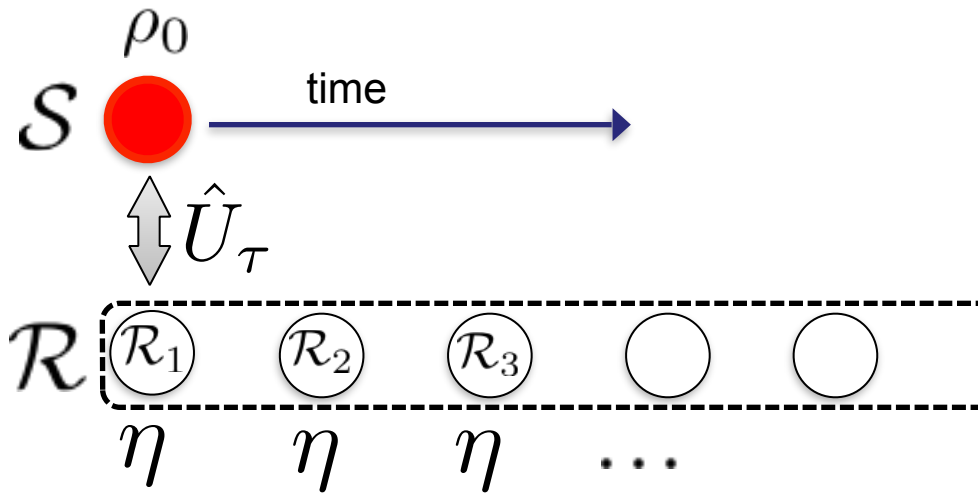
$$\mathcal{E}_k^\tau := \mathcal{E}^{k\tau} \rightarrow \mathcal{E}_k^\tau \mathcal{E}_j^\tau = \mathcal{E}^{(k+j)\tau}$$

One parameter semigroup

$$\rho_{n\tau} = \text{Tr}_{\mathcal{R}} [\hat{U}_\tau (\rho_{(n-1)\tau} \otimes \eta) \hat{U}_\tau^\dagger] = \mathcal{E}^\tau [\rho_{(n-1)\tau}]$$

$$\Delta \rho_{n\tau} = \rho_{n\tau} - \rho_{(n-1)\tau} = (\mathcal{E}^\tau - \mathbb{I}) [\rho_{(n-1)\tau}]$$

QCM in the continuous time limit



τ = collision time
 g = coupling strength

$$\hat{U}_\tau = e^{-igV\tau}$$

$$V = \sum_k \hat{S}_k \otimes \hat{R}_k$$

TAKING THE LIMIT

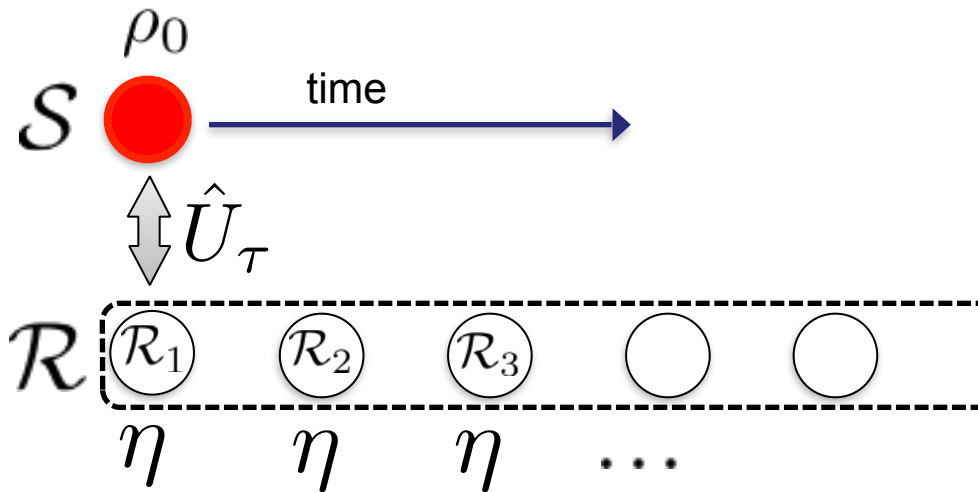
$$\begin{aligned}
 n &\rightarrow \infty \\
 \tau &\rightarrow 0
 \end{aligned}$$

IN SUCH WAY

THAT

$$\begin{aligned}
 n\tau &\rightarrow t \\
 g^2\tau &\rightarrow \gamma
 \end{aligned}$$

QCM in the continuous time limit



τ = collision time
 g = coupling strength

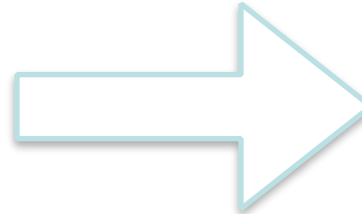
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IN SUCH WAY



THAT

$$\begin{aligned}
 n\tau &\rightarrow t \\
 g^2\tau &\rightarrow \gamma
 \end{aligned}$$

$$\frac{\Delta\rho_{n\tau}}{\tau} = \frac{(\mathcal{E}_\tau - \mathbb{I})}{\tau} [\rho_{(n-1)\tau}]$$

$$\frac{d\rho(t)}{dt} = \gamma \sum_k \langle R_k R_{k'} \rangle_\eta \left(S_k \rho(t) S_{k'} - \frac{1}{2} \{ S_{k'} S_k, \rho(t) \} \right) \quad \text{Lindblad ME}$$

Quantum Collision Models

Advantages:

1. Intuitive & easy to handle
2. Lindblad-type MEs arise naturally
3. Promising tools for tackling NM dynamics

T. Rybar, S. N. Filippov, M. Ziman, and V. Buzek, J. Phys. B 45, 154006 (2012)

N. K. Bernardes, A. R. R. Carvalho, C. H. Monken, M. F. Santos, Phys. Rev A 90, 032111 (2014)

V. Giovannetti, and G. M. Palma, Phys. Rev. Lett. 108, 040401 (2012)

F. Ciccarello, G. M. Palma, and V. Giovannetti, Phys. Rev. A 87, 040103(R) (2013)

F. Ciccarello and V. Giovannetti, Phys. Scrip. T153, 014010 (2013)

R. McCloskey and M. Paternostro, Phys. Rev. A 89, 052120 (2014)

A. Bodor, L. Diosi, Z. Kallus, and T. Konrad, Phys. Rev. A 87, 052113 (2013)

J. Jin, V. Giovannetti, R. Fazio, F. Sciarrino, P. Mataloni, A. Crespi, and R. Osellame, PRA 91, 012122 (2015)

N. K. Bernardes et al., arXiv: arXiv:1504.01602

QCM and Landauer's Principle

$\rho_{n\tau}$

\mathcal{S}

\hat{U}_τ

$\dots (\mathcal{R}_n) \dots$

η^{th}

$$\eta^{\text{th}} = e^{-\beta H_{\mathcal{R}}} / \text{Tr}[e^{-\beta H_{\mathcal{R}}}]$$

$\tau = \text{collision time}$

$g = \text{coupling strength}$

$$\hat{U}_\tau = e^{-igV\tau}$$

$$V = \sum_k \hat{S}_k \otimes \hat{R}_k$$

$$\rho_{n\tau} = \text{Tr}_{\mathcal{R}} [\hat{U}_\tau (\rho_{(n-1)\tau} \otimes \eta^{\text{th}}) \hat{U}_\tau^\dagger] = \mathcal{E}^\tau [\rho_{(n-1)\tau}]$$

$$\eta'_n = \text{Tr}_{\mathcal{S}} [\hat{U}_\tau (\rho_{(n-1)\tau} \otimes \eta^{\text{th}}) \hat{U}_\tau^\dagger] = \Lambda_{(n-1)\tau} [\eta^{\text{th}}]$$

FLUXES

$$\Delta \tilde{S}_{n\tau} = S(\rho_{(n-1)\tau}) - S(\rho_{n\tau})$$

$$\Delta E_{n\tau} = \text{Tr} \left[\hat{H}_{\mathcal{S}} (\mathcal{E}^\tau - \mathbb{I}) [\rho_{(n-1)\tau}] \right]$$

$$\Delta Q_{n\tau} = \text{Tr} \left[\hat{H}_{\mathcal{R}} (\Lambda_{(n-1)\tau} - \mathbb{I}) [\eta^{\text{th}}] \right]$$

...continuous time limit

$$\dot{\tilde{S}}(t)$$

$$\dot{E}(t)$$

$$\dot{Q}(t)$$

QCM and Landauer's Principle

$\rho_{n\tau}$

\mathcal{S}

\hat{U}_τ

$\dots \mathcal{R}_n \dots$

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FLUXES

$$\dot{E}(t) = \gamma \sum_{k,j} \langle \hat{R}_k \hat{R}_j \rangle_{\eta^{\text{th}}} \langle \hat{S}_k \hat{H}_{\mathcal{S}} \hat{S}_j - \frac{1}{2} \{ \hat{S}_k \hat{S}_j, \hat{H}_{\mathcal{S}} \} \rangle_{\rho(t)}$$

$$\dot{Q}(t) = \gamma \sum_{k,j} \langle \hat{S}_k \hat{S}_j \rangle_{\rho(t)} \langle \hat{R}_k \hat{H}_{\mathcal{R}} \hat{R}_j - \frac{1}{2} \{ \hat{R}_k \hat{R}_j, \hat{H}_{\mathcal{R}} \} \rangle_{\eta^{\text{th}}}$$

QCM and Landauer's Principle

$\rho_{n\tau}$

S

\hat{U}_τ

$\dots (\mathcal{R}_n) \dots$

η^{th}

$$\eta^{\text{th}} = e^{-\beta H_{\mathcal{R}}} / \text{Tr}[e^{-\beta H_{\mathcal{R}}}]$$

THERMALIZATION PROCESSES

$$[\hat{U}_\tau, \hat{H}_S + \hat{H}_{\mathcal{R}}] = 0$$

$$\rho^{\text{eq}} \equiv \rho(\infty) = e^{-\beta H_S} / \text{Tr}[e^{-\beta H_S}]$$

$$\Delta Q_{n\tau} = -\Delta E_{n\tau} \rightarrow \dot{Q}(t) = -\dot{E}(t)$$

$$\dot{S}(\rho(t) | \rho^{\text{eq}}) = \text{Tr}[\dot{\rho}(t)(\ln \rho(t) - \ln \rho^{\text{eq}})] = \dot{\tilde{S}} + \beta \dot{E}$$

≤ 0 for CPT maps

QCM and Landauer's Principle

$\rho_{n\tau}$

S

\hat{U}_τ

$\dots (\mathcal{R}_n) \dots$

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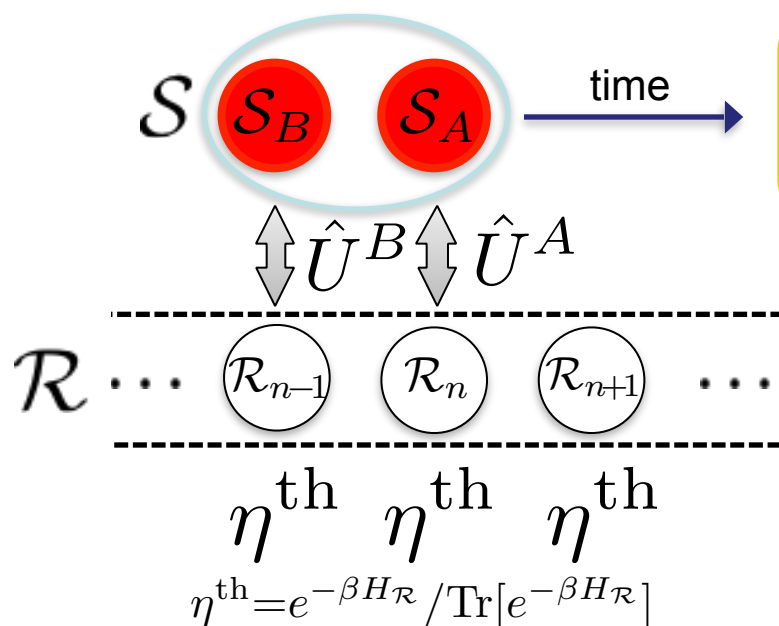
$$\dot{S}(\rho(t) | \rho^{\text{eq}}) = \text{Tr}[\dot{\rho}(t)(\ln \rho(t) - \ln \rho^{\text{eq}})] = \dot{\tilde{S}} + \beta \dot{E}$$

≤ 0 for CPT maps

Landauer's principle for fluxes

$$\beta \dot{Q}(t) \geq \dot{\tilde{S}}(\rho(t))$$

Recycling environment



$\tau = \text{collision time}$

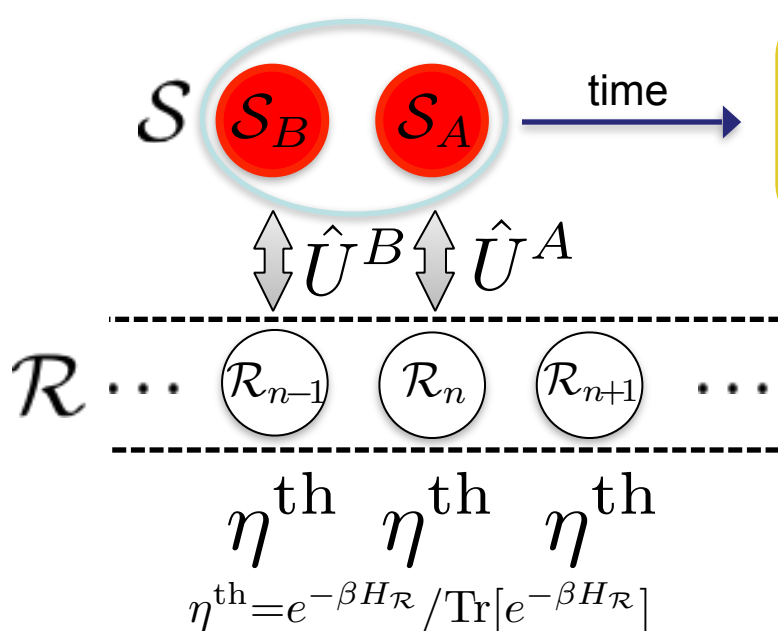
$$\hat{U}_{\tau} = \hat{U}^B \hat{U}^A$$

V. Giovannetti, and G. M. Palma, PRL. 108, 040401 (2012)
 SL et al., PRA **91**, 022121 (2015)

$$\dot{\rho}(t) = \sum_{X=A,B} \mathcal{L}_X[\rho(t)] + \mathcal{D}_{AB}^{\rightarrow}[\rho(t)]$$

$$\rho^{eq} = \frac{e^{-\beta H_A}}{Z_A} \otimes \frac{e^{-\beta H_B}}{Z_B}$$

Recycling environment



Landauer's principle for fluxes

$$\beta \dot{Q} \geq \dot{\tilde{S}}_{AB}$$

$$\beta \dot{Q} \geq 2\dot{\tilde{S}}_A + \dot{S}_{A|B} - \dot{S}_{B|A} + \dot{I}(A:B)$$

independent
erasure bound

intra-system correlations

$\tau = \text{collision time}$

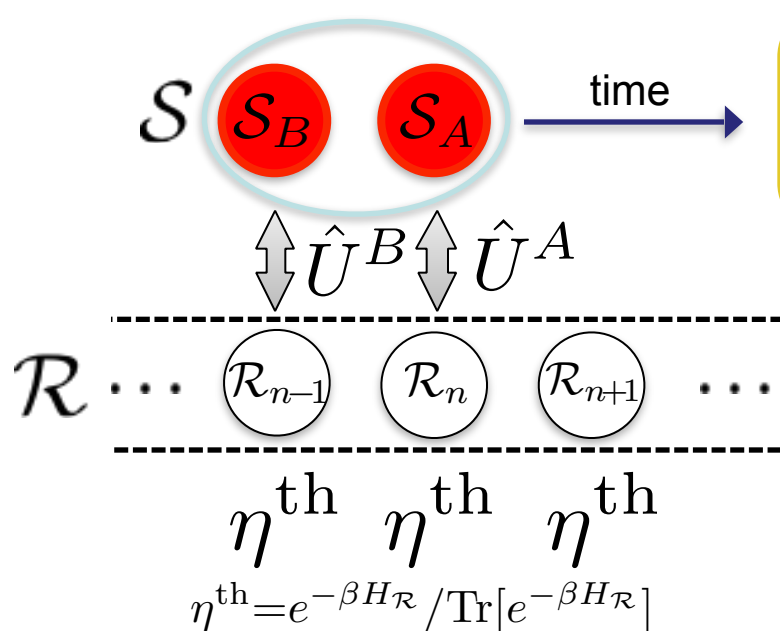
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Recycling environment



$\tau = \text{collision time}$

$$\hat{U}_\tau = \hat{U}^B \hat{U}^A$$

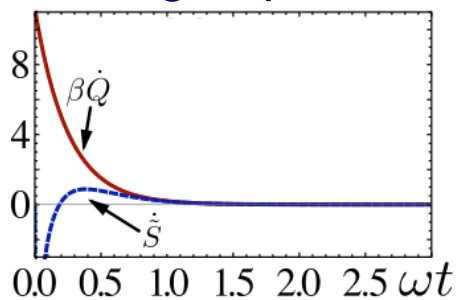
$$H_S = (\omega/2)\sigma_A^z + (\omega/2)\sigma_B^z$$

$$H_{\mathcal{R}_n} = (\omega/2)\sigma_n^z$$

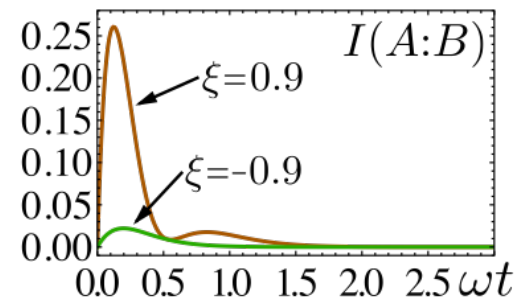
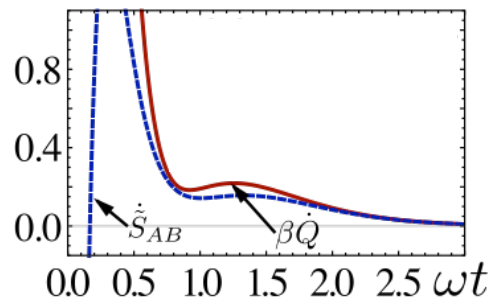
$$V_X = \sigma_X^x \sigma_n^x + \sigma_X^y \sigma_n^y$$

$$\rho_0 = |\uparrow\uparrow\rangle_{AB} \langle\uparrow\uparrow| \quad \eta^{\text{th}} = \begin{pmatrix} \frac{1-\xi}{2} & 0 \\ 0 & \frac{1+\xi}{2} \end{pmatrix}$$

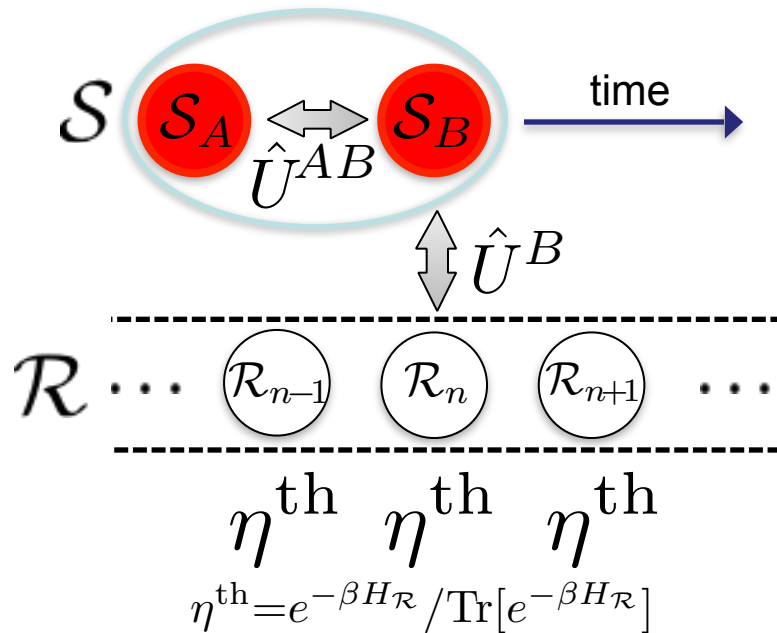
single qubit



2-qubit recycling scheme



Indirect Erasure



$\tau = \text{collision time}$

$$\hat{U}_\tau = \hat{U}^B \hat{U}^{AB}$$

SL, F. Ciccarello and G. M Palma, in preparation (2015)

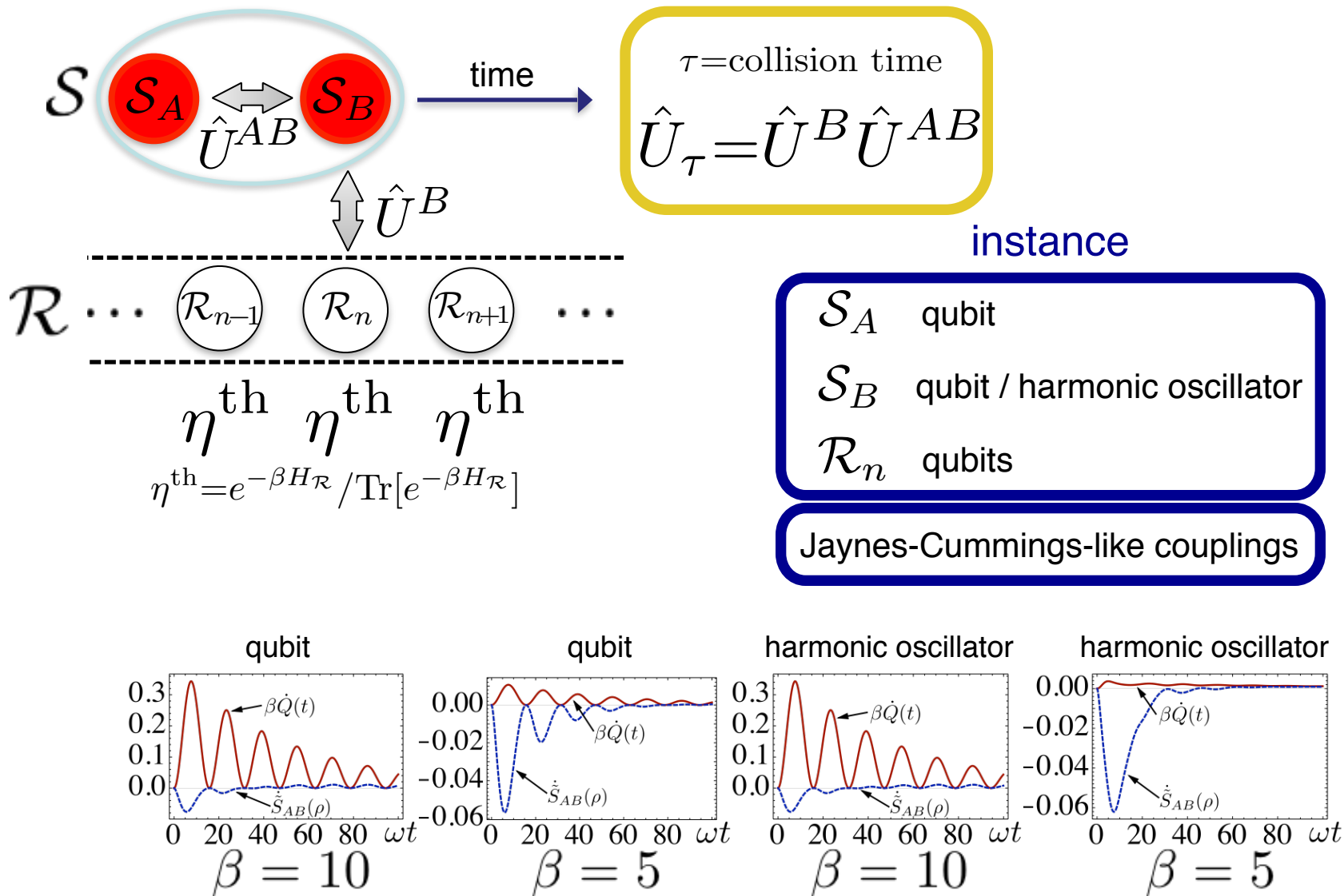
$$\dot{\rho}(t) = -i[H_{AB}, \rho(t)] + \mathcal{L}_B[\rho(t)]$$

$$\rho^{eq} = \frac{e^{-\beta H_A}}{Z_A} \otimes \frac{e^{-\beta H_B}}{Z_B}$$

Landauer's principle for fluxes

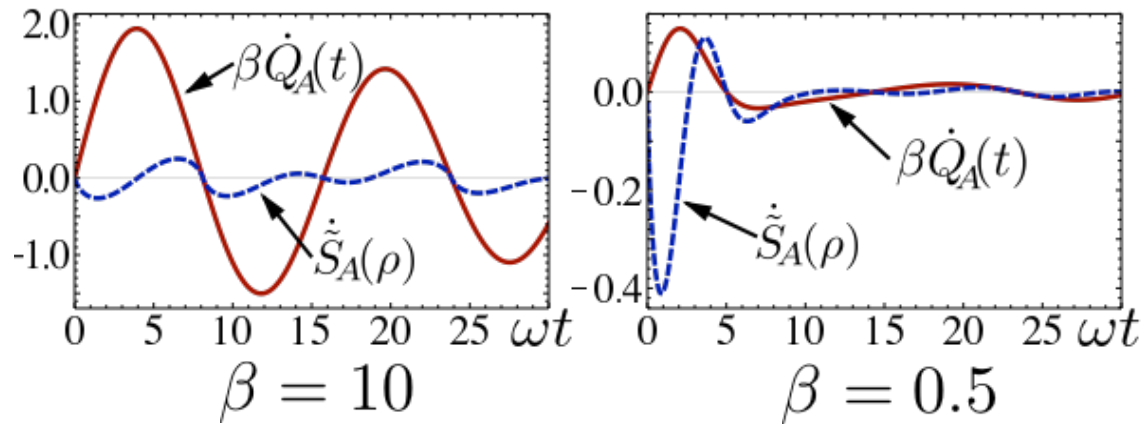
$$\beta \dot{Q} \geq \dot{\tilde{S}}_{AB}$$

Indirect Erasure



Indirect Erasure

...looking at S_A only:



Landauer's principle violation

CONCLUSIONS

Landauer's principle from a microscopic system-bath quantum model

Landauer's principle for fluxes of open quantum system

Link between Erasure process and intra-system correlations

preprint:

SL, R. McCloskey, F. Ciccarello, M. Paternostro & G. M. Palma, [arXiv: 1503.07837](https://arxiv.org/abs/1503.07837)

thank you!