

## Boson Sampling with integrated photonics

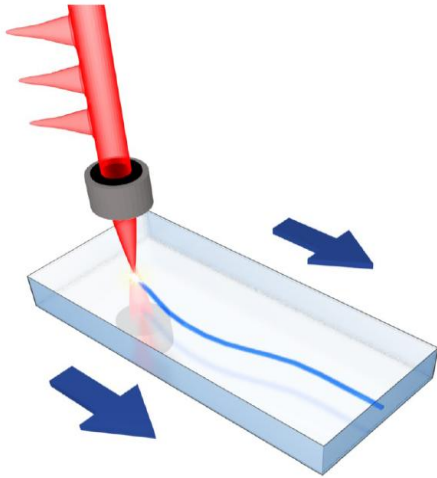
*Paolo Mataloni*

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Sapienza Università di Roma

<http://quantumoptics.phys.uniroma1.it>

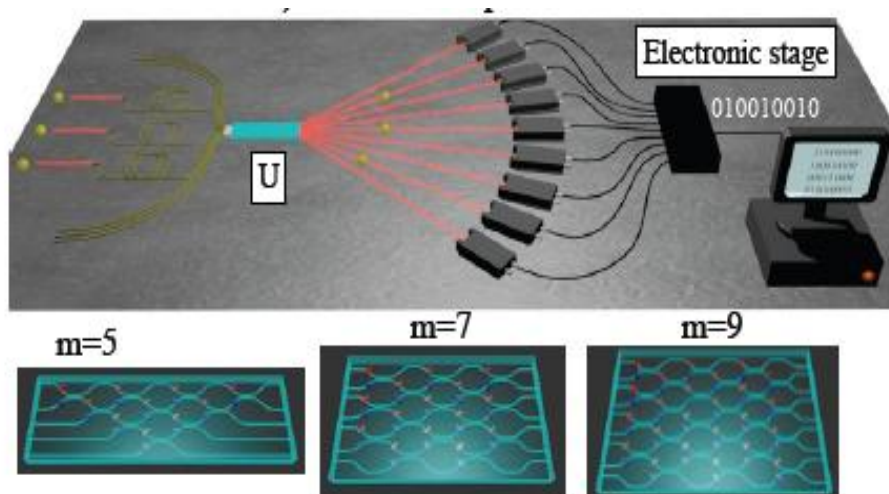
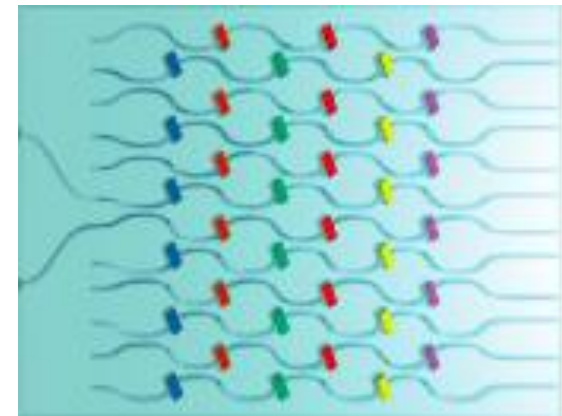
*IQIS 2015, Monopoli, September 10, 2015*

# Summary



## 1) Femtosecond laser writing: main features

## 2) Complex integrated photonics networks. Discrete quantum walks



## 3) BosonSampling: a promising route towards *Quantum Supremacy*

# Femtosecond laser writing

**3-D capability**

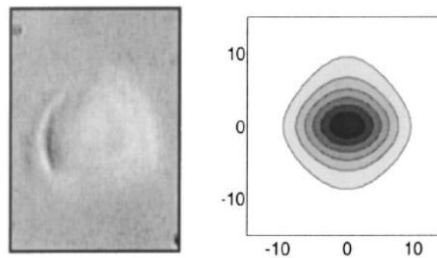
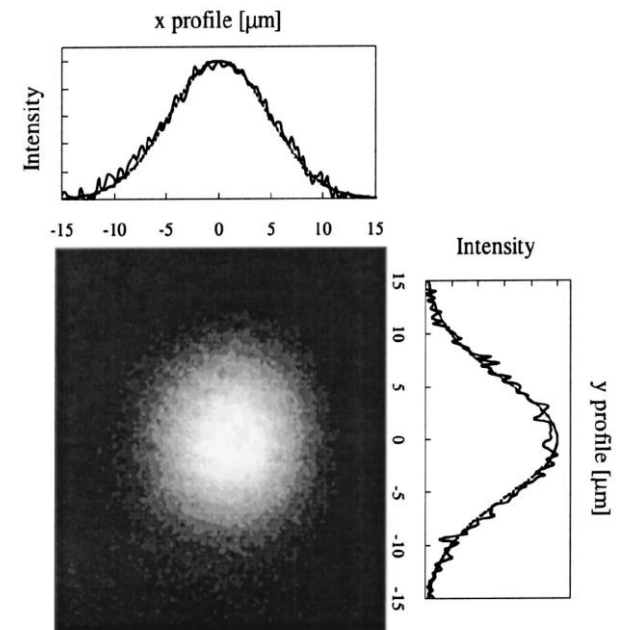
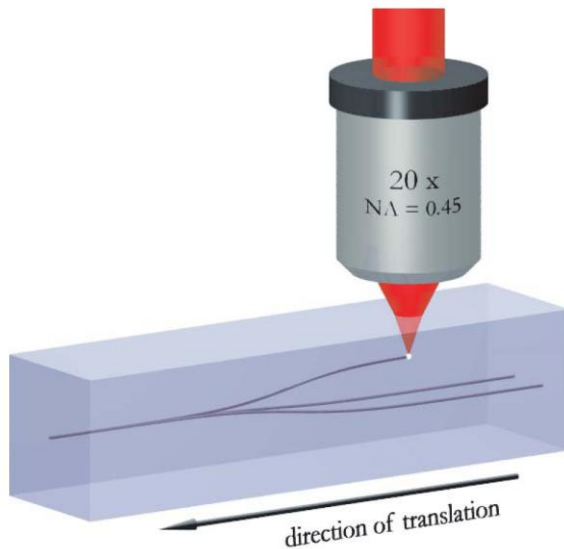
Rapid device prototyping:  
writing speed = 4 cm/s

Propagation of circular  
gaussian modes

**Main features**

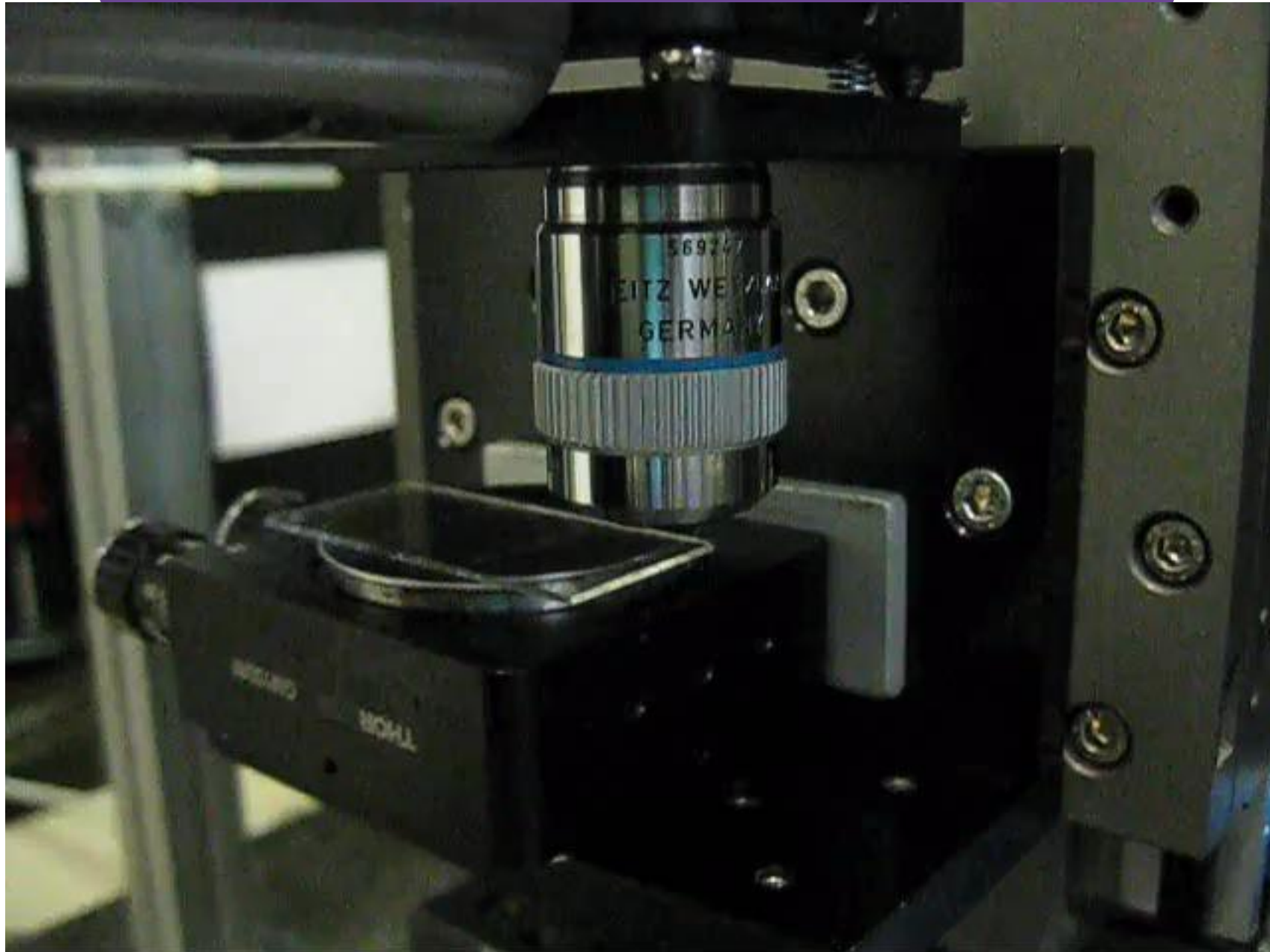
Circular waveguide  
transverse profile

**Low birefringence**

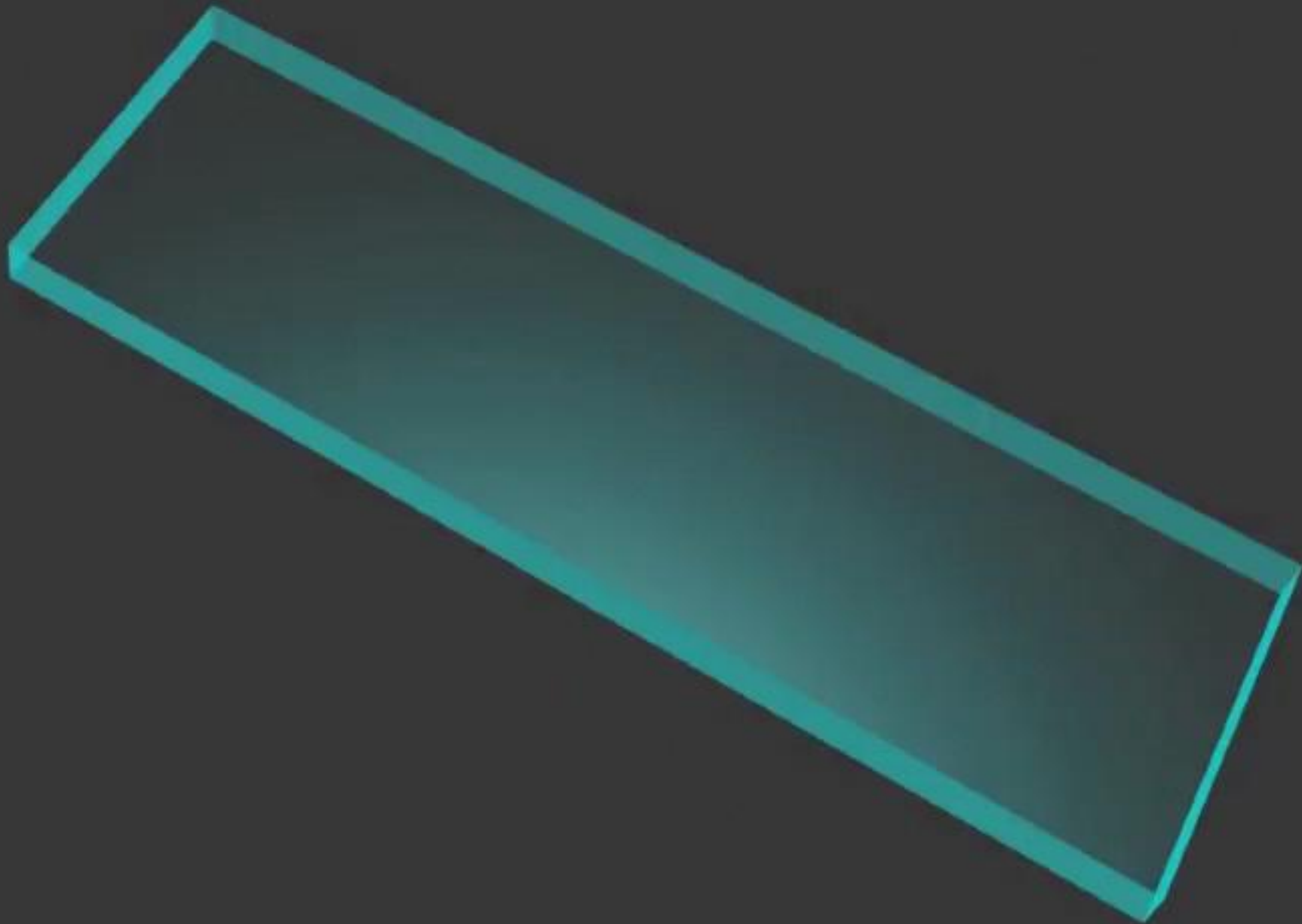


***ABLE TO SUPPORT ANY POLARIZATION STATE***

# Femtosecond laser writing



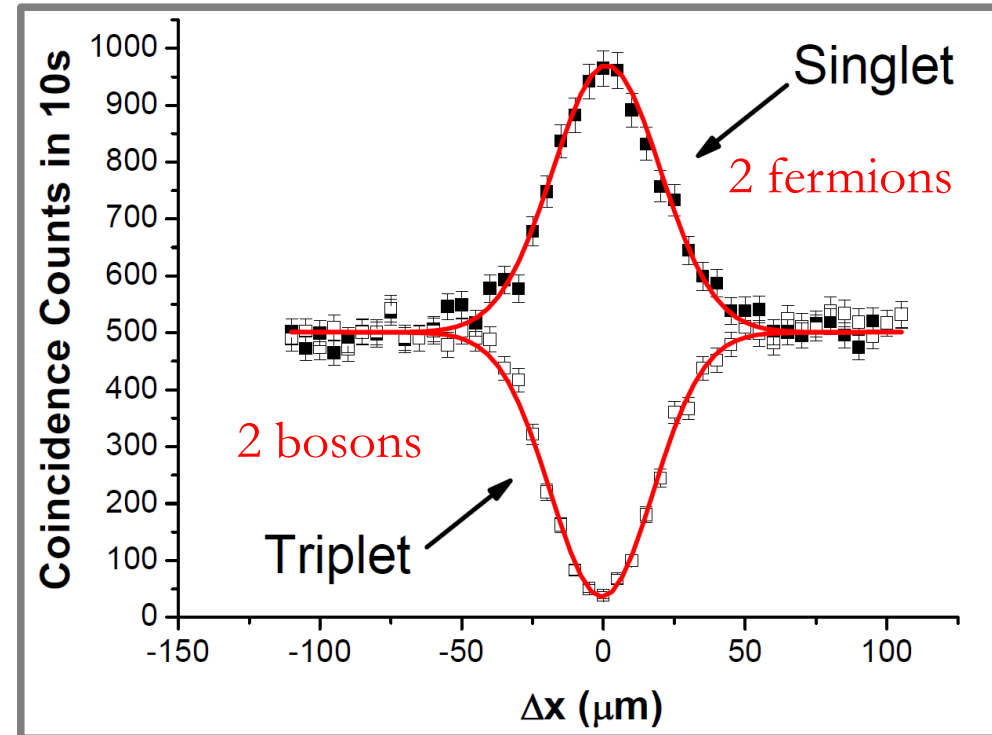
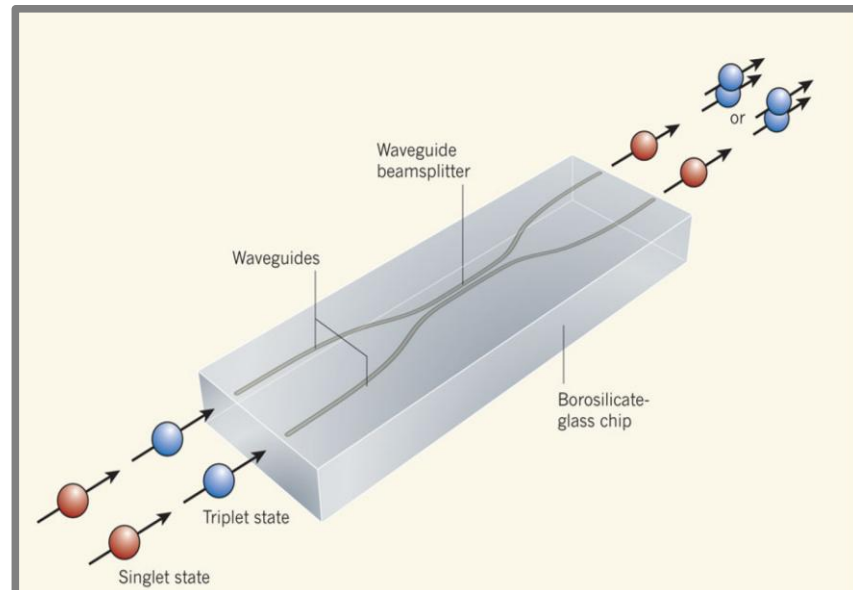
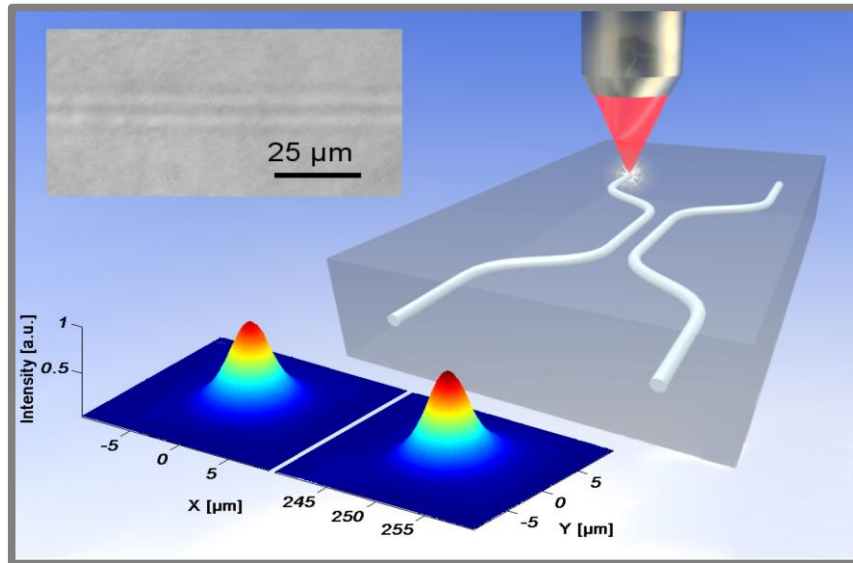
# Femtosecond laser writing





# Integrated beam splitter

$$|\Psi^\phi\rangle = \frac{1}{\sqrt{2}}(|H\rangle_A|V\rangle_B + e^{i\phi}|V\rangle_A|H\rangle_B)$$



L. Sansoni *et al.* *Phys. Rev. Lett.* **105**, 200503 (2010)

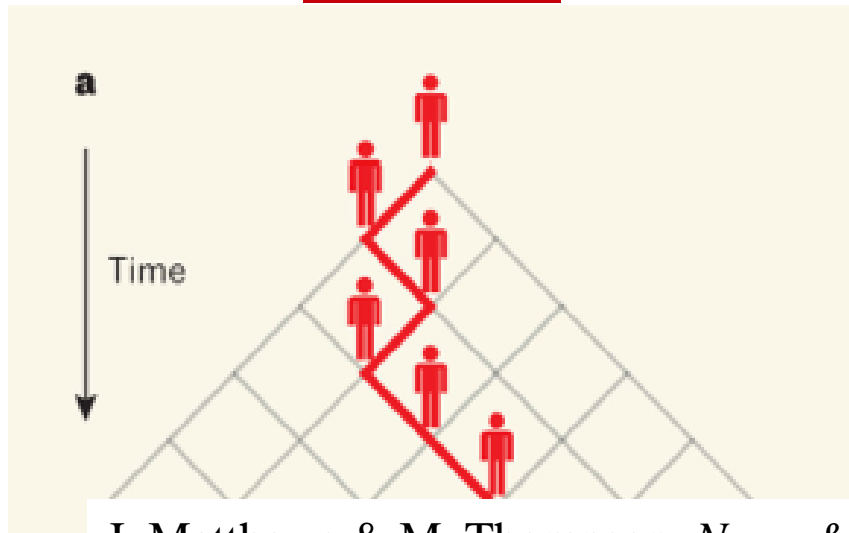


**Goal: to realize complex integrated photonics structures of beam splitters.**

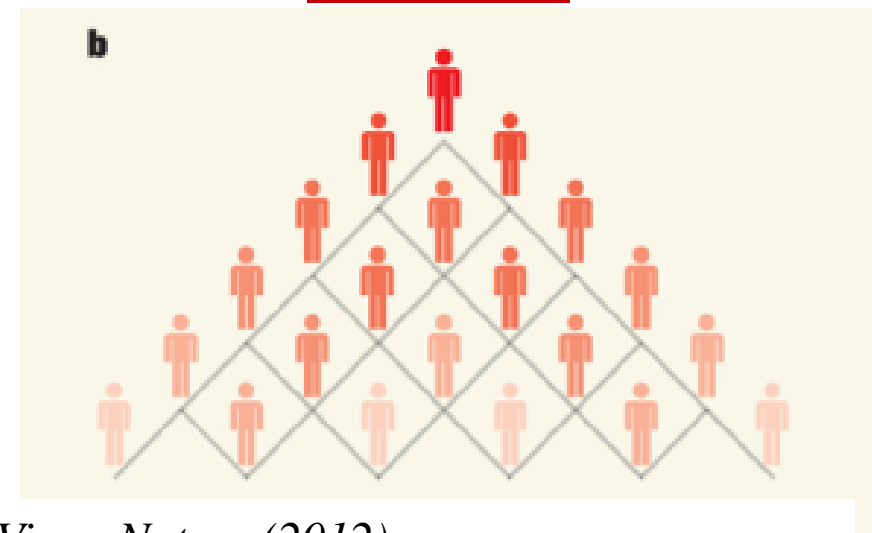
# Quantum Walk

Quantum walk: extension of the classical random walk:  
a walker on a lattice “jumping” between different sites with given probability

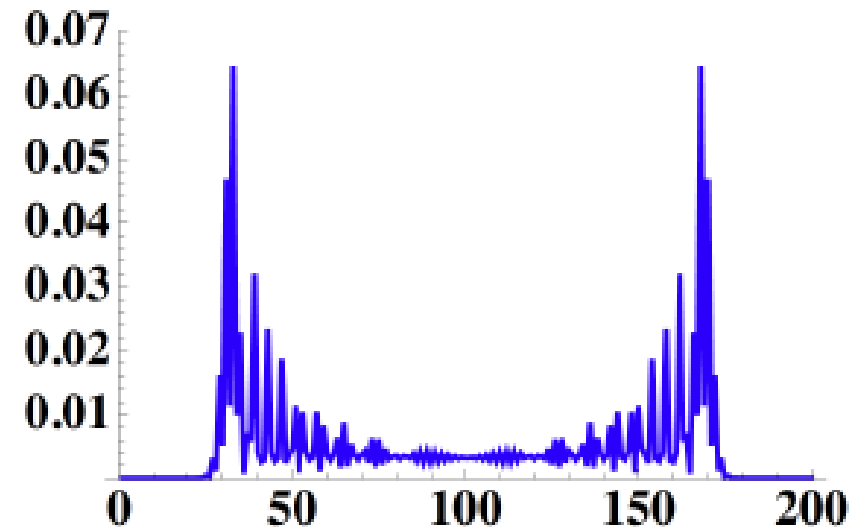
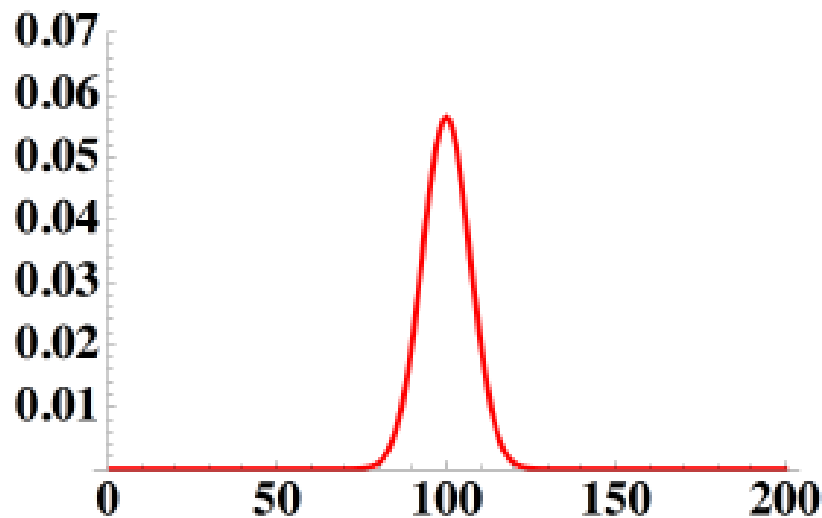
## Classical



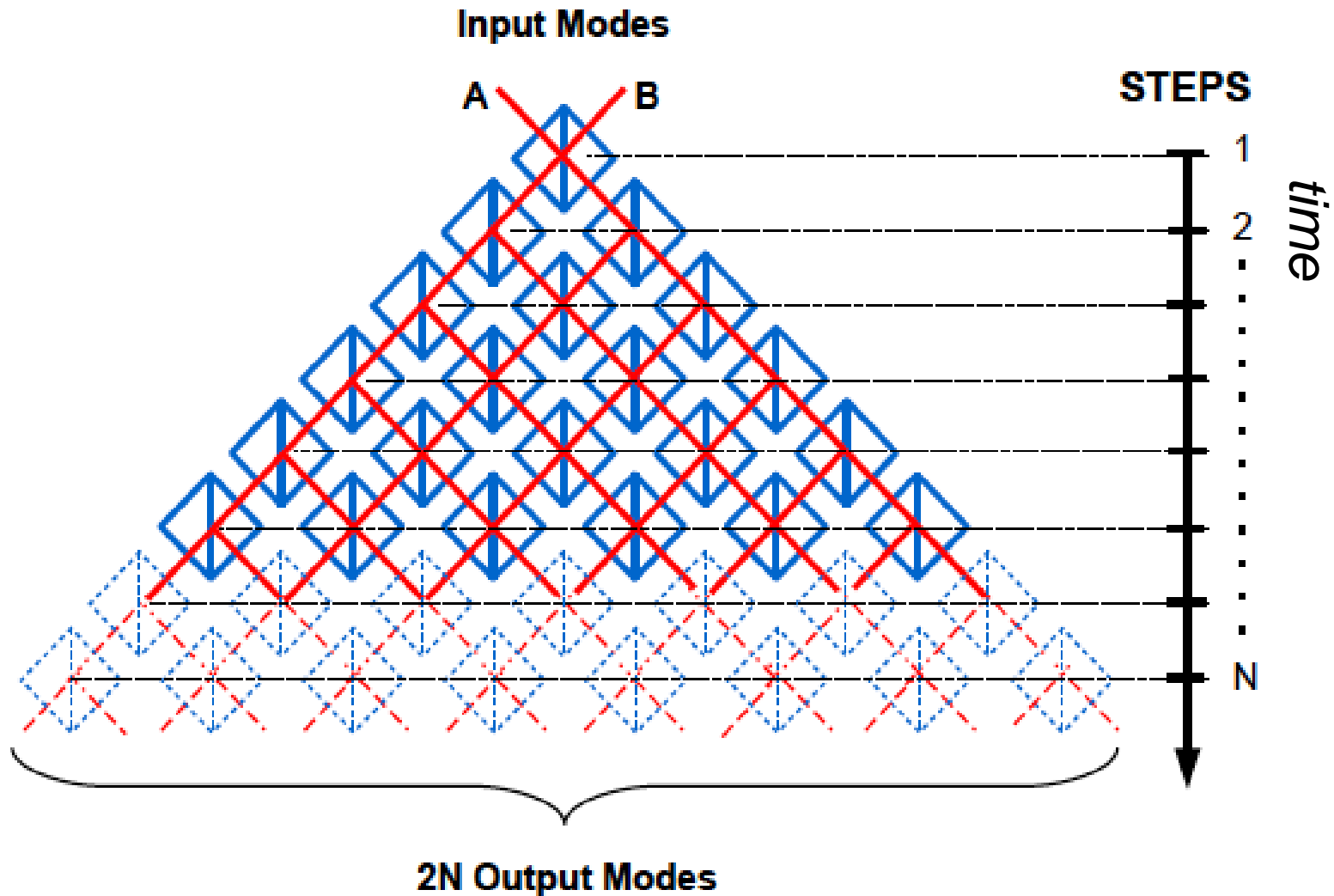
## Quantum



J. Matthews & M. Thompson, *News & Views Nature* (2012)



# Implementing QW with photons



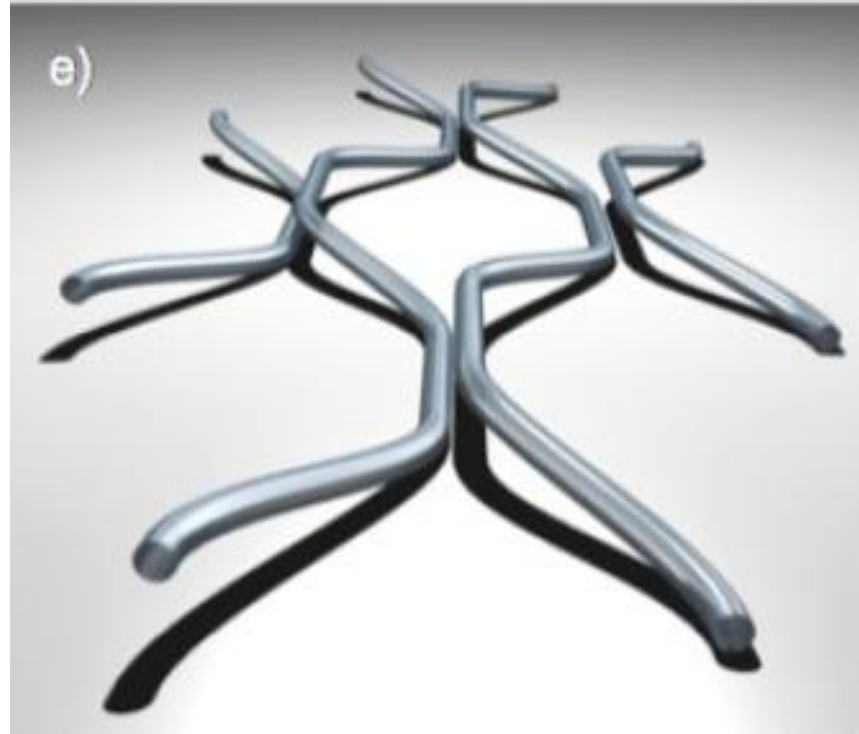
Discrete QW: realized by concatenating many beam splitters and phase shifters



**Feasible with integrated photonic waveguides**



# Polarization independent QW



- Polarization independent lattices made possible by 3D writing capability
- Path lengths controlled up to few nanometers



# Main QW realizations

Ordered QW



*Study the statistics of  
bosonic/fermionic wavefunction*  
PRL (2012)

Disordered QW depending

- on site

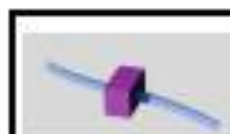
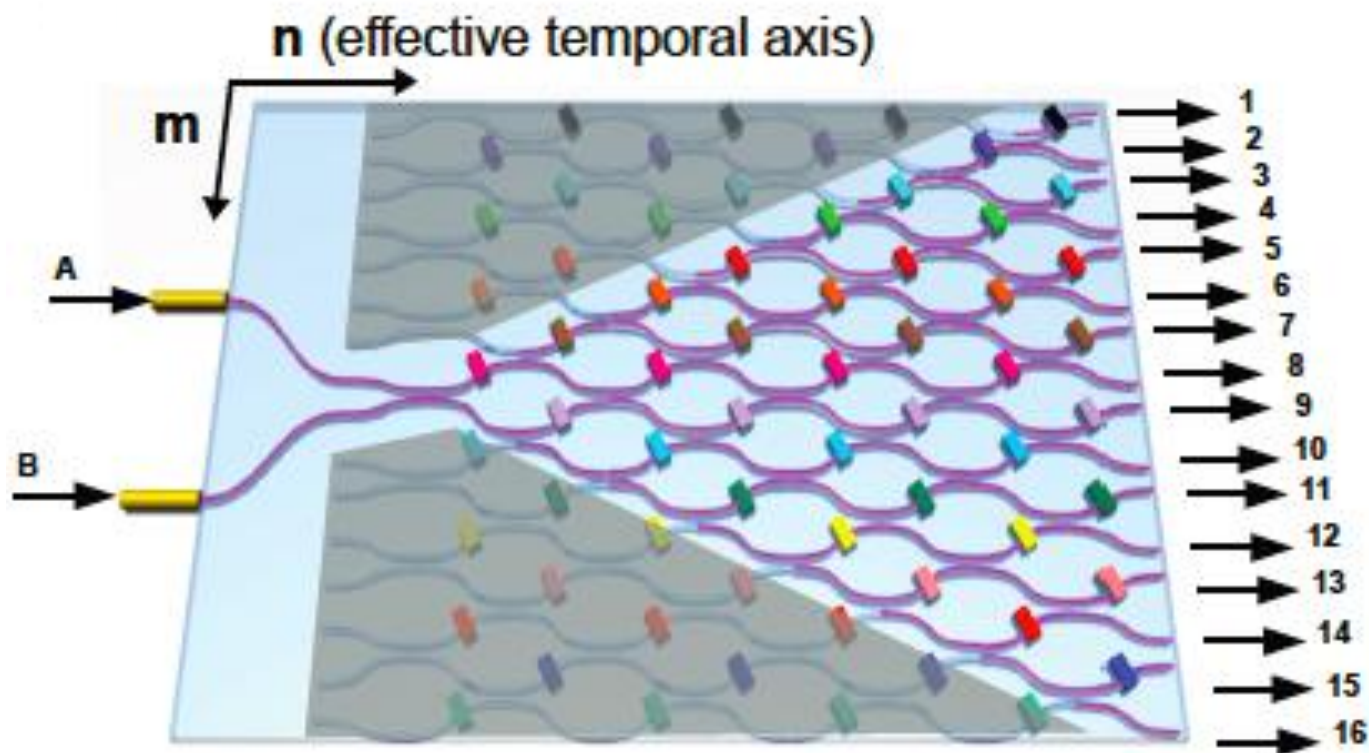
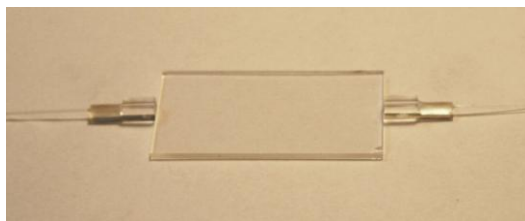
- but **NOT** on time



***Anderson localization** of the  
quantum particle wavefunction*  
Nat Photon (2013)

Up to 16 output modes  
and 64 polarization  
independent BSs and  
phase-shifters

**On chip BS array**



Phase  
shift



Directional  
coupler

# The BosonSampling problem

HOW TO ACHIEVE  
QUANTUM  
SUPREMACY ??



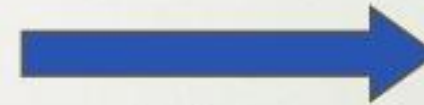
John Preskill  
@preskill

Segui

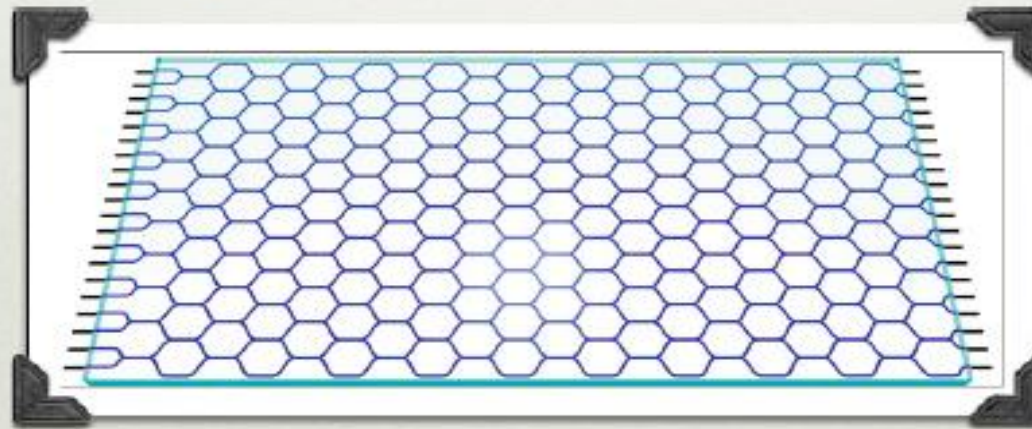
Proposed "quantum supremacy" for controlled quantum systems surpassing classical ones. Please suggest alternatives.

## BOSON SAMPLING

propagation on the chip with  $m$  modes



Input:  
 $n$  bosons



Output:  
 $n$ -photon state

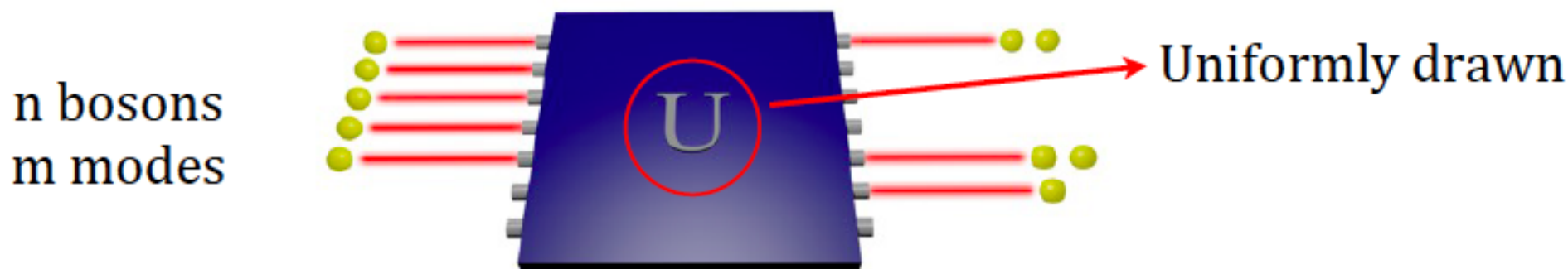
*Can a classical computer simulate  
the distribution of the output mode numbers ?*

**Answer: NO!!**



# BosonSampling

Sampling the output distribution (*even approximately*) of non-interacting bosons evolving through a linear network is hard to do with classical resources



Why? Transition amplitudes are related to the permanent of square matrices

$$\langle T | U_F | S \rangle = \frac{\text{Per}(U_{S,T})}{\sqrt{s_1! \dots s_m! t_1! \dots t_m!}}$$

$$\text{Per}(A) = \sum_{\sigma \in S_n} \prod_{i=1}^n a_{i, \sigma_i}$$
classically hard

input

	0	1	1	0	1
output 0	0.212	$-0.018 + 0.165i$	$-0.238 - 0.18i$	$-0.429 + 0.32i$	$-0.715 + 0.2i$
1	$-0.193 - 0.388i$	$-0.045 - 0.379i$	$0.19 + 0.311i$	$0.328 - 0.269i$	$-0.594 + 0.03i$
1	$-0.723 + 0.363i$	$0.087 - 0.09i$	$-0.076 - 0.155i$	$0.206 + 0.443i$	$-0.153 - 0.193i$
1	$-0.092 + 0.045i$	$-0.148 - 0.645i$	$-0.588 + 0.184i$	$-0.369 - 0.086i$	$0.167 + 0.025i$
0	$0.318 - 0.009i$	$-0.144 - 0.594i$	$0.452 - 0.405i$	$0.037 + 0.387i$	$0.071 + 0.025i$

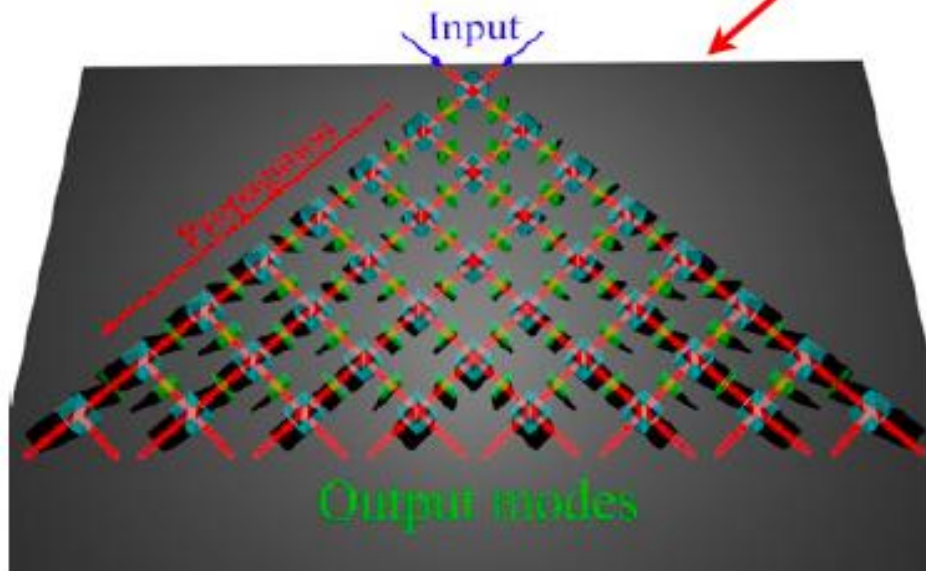
# BosonSampling

Photons naturally solve the BosonSampling problem

Experimental platform: photons in linear optical interferometers

- Required resources:
- Single-photon inputs
  - Multimode interferometers
  - Detection

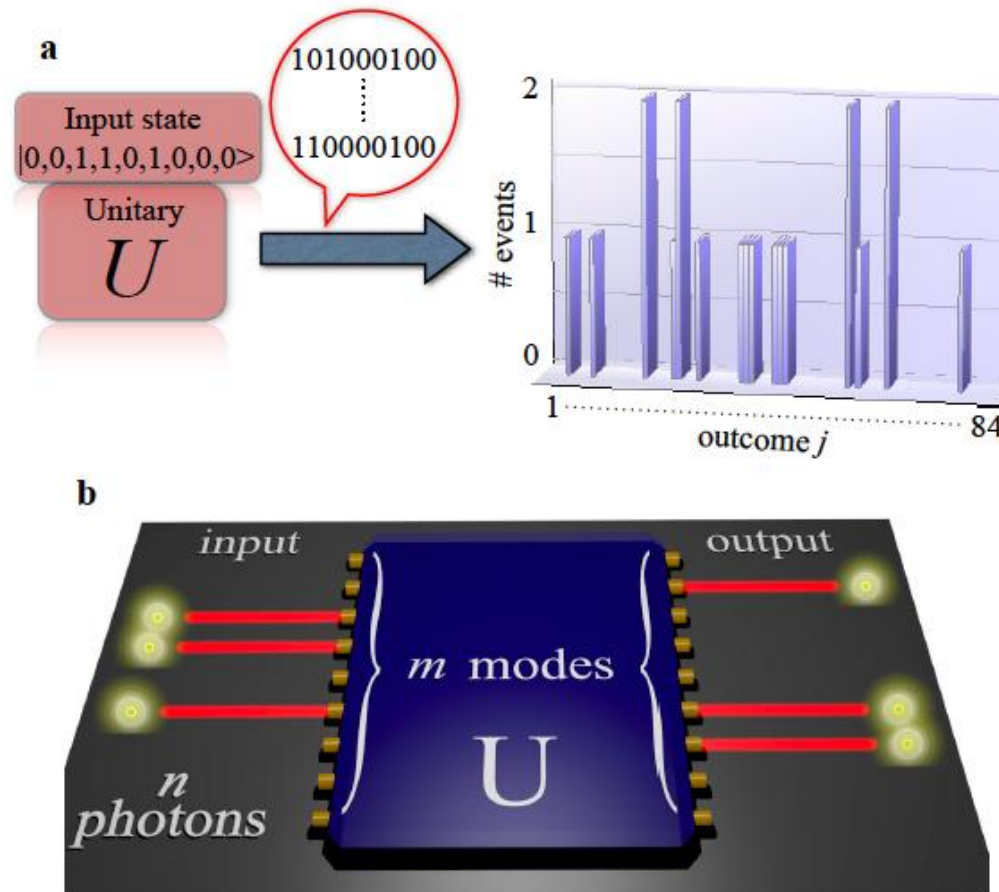
n photons  
m modes



Hard to implement with bulk optics

Require a technological step recently available due to integrated photonics

# BosonSampling



« Small-scale quantum computers made from an array of interconnected waveguides on a glass chip can now perform a task that is considered hard to undertake on a large scale by classical means. »

T. Ralph, News & Views, *Nature Photonics* 7, 514 (2013)



# BosonSampling



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OXFORD



nature  
photonics

LETTERS

PUBLISHED ONLINE: XX XX 2013 | DOI: 10.1038/NPHOTON.2013.112

## Integrated multimode interferometers with arbitrary designs for photonic boson sampling

Andrea Crespi<sup>1,2</sup>, Roberto Osellame<sup>1,2\*</sup>, Roberta Ramponi<sup>1,2</sup>, Daniel J. Brod<sup>3</sup>, Ernesto F. Galvão<sup>3\*</sup>, Nicolò Spagnolo<sup>4</sup>, Chiara Vitelli<sup>4,5</sup>, Enrico Maiorino<sup>4</sup>, Paolo Mataloni<sup>4</sup> and Fabio Sciarrino<sup>4\*</sup>

1 The evolution of bosons undergoing arbitrary linear unitary transformations quickly becomes hard to predict using classical computers as we increase the number of particles and modes. 2 Photons propagating in a multiport interferometer naturally solve this so-called boson sampling problem<sup>1</sup>, thereby motivating the development of technologies that enable precise control of multiphoton interference in large interferometers<sup>2-4</sup>. Here, we use novel three-dimensional manufacturing techniques to achieve simultaneous control of all the parameters describing

proportional to the permanent of a matrix associated with the interferometer (see Methods for details), and the permanent is a function that is notoriously hard to compute<sup>10</sup>. In ref. 1 it was estimated that a system of approximately 20 photons in  $m \approx 400$  modes would already take noticeably long to simulate classically. At present, the most promising technology for achieving this regime involves inputting Fock states into multimode integrated photonic chips<sup>2-4,11-13</sup>.

In this Letter we report on the experimental implementation of

## Boson Sampling on a Photonic Chip

Justin B. Spring<sup>1\*</sup>, Benjamin J. Metcalf<sup>1</sup>, Peter C. Humphreys<sup>1</sup>, W. Steven Kolthammer<sup>1</sup>, Xian-Min Jin<sup>1,2</sup>, Marco Barbieri<sup>1</sup>, Animesh Datta<sup>1</sup>, Nicholas Thomas-Peter<sup>1</sup>, Nathan K. Langford<sup>1,3</sup>, Dmytro Kundys<sup>4</sup>, James C. Gates<sup>4</sup>, Brian J. Smith<sup>1</sup>, Peter G. R. Smith<sup>4</sup>, Ian A. Walmsley<sup>1\*</sup>

Although universal quantum computers ideally solve problems such as factoring integers exponentially more efficiently than classical machines, the formidable challenges in building such devices motivate the demonstration of simpler, problem-specific algorithms that still promise a quantum speedup. We constructed a quantum boson-sampling machine (QBSM) to sample the output distribution resulting from the nonclassical interference of photons in an integrated photonic circuit, a problem thought to be exponentially hard to solve classically. Unlike universal quantum computation, boson sampling merely requires indistinguishable photons, linear state evolution, and detectors. We benchmarked our QBSM with three and four photons and analyzed sources of sampling inaccuracy. Scaling up to larger devices could offer the first definitive quantum-enhanced computation.

modes (18). Such circuits can be rapidly reconfigured to sample from a user-defined operation (19, 20). Importantly, boson sampling requires neither nonlinearities nor on-demand entanglement, which are substantial challenges in photonic universal quantum computation (21). This clears the way for experimental boson sampling with existing photonic technology, building on the extensively studied two-photon Hong-Ou-Mandel interference effect (22).

A QBSM (Fig. 1) samples the output distribution of a multiparticle bosonic quantum state  $|\Psi_{\text{out}}\rangle$ , prepared from a specified initial state  $|\mathbf{T}\rangle$  and linear transformation  $\mathbf{A}$ . Unavoidable losses in the system imply  $\mathbf{A}$  will not be unitary, although lossy QBSMs can still surpass classical computation (12, 23). A trial begins with the input state  $|\mathbf{T}\rangle = |T_1 \dots T_M\rangle \propto \prod_{i=1}^M (a_i^\dagger)^{T_i} |0\rangle$ , which describes  $N = \sum_{i=1}^M T_i$  particles distributed in  $M$  input modes in the occupation-number representation. The output state  $|\Psi_{\text{out}}\rangle$  is generated

Universal quantum computers require physical systems that are well isolated from unitary transformation  $U$  is thought to be exponentially hard to sample from classically (12). The

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nature  
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## Experimental boson sampling

Max Tillmann<sup>1,2\*</sup>, Borivoje Dakić<sup>1</sup>, René Heilmann<sup>3</sup>, Stefan Nolte<sup>3</sup>, Alexander Szameit<sup>3</sup> and Philip Walther<sup>1,2\*</sup>

Universal quantum computers<sup>1</sup> promise a dramatic increase in speed over classical computers, but their full-size realization remains challenging<sup>2</sup>. However, intermediate quantum computational models<sup>3-5</sup> have been proposed that are not universal but can solve problems that are believed to be classically hard. Aaronson and Arkhipov<sup>6</sup> have shown that interference of single photons in random optical networks can solve the hard problem of sampling the bosonic output distribution. Remarkably, this computation does not require measurement-based interactions<sup>7,8</sup> or adaptive feed-forward techniques<sup>9</sup>

photons. Randomly chosen instances of this problem are strongly believed to be hard to solve by classical means. Instances of boson sampling can be realized with quantum systems composed of non-interacting photons that are processed through randomly chosen networks of physical modes. The bosonic nature of the photons leads to non-classical interference, producing an output



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To implement a circuit, the subgraphs representing circuit elements are connected by paths. Figure 4 depicts a graph corresponding to a simple two-qubit computation. Timing is important: Wave packets must meet on the vertical paths for interactions to occur. We achieve this by choosing the numbers of vertices on each of the segments in the graph appropriately, taking into account the different propagation speeds of the two wave packets [see section S4 of (32)]. In section S3.1 of (32), we present a refinement of our scheme using planar graphs with maximum degree four.

By analyzing the full  $(n+1)$ -particle interacting many-body system, we prove that our algorithm performs the desired quantum computation up to an error term that can be made arbitrarily small (32). Our analysis goes beyond the scattering theory discussion presented above; we take into account the fact that both the wave packets and the graphs are finite. Specifically, we prove that by choosing the size of the wave packets, the number of vertices in the graph, and the total

## Photonic Boson Sampling in a Tunable Circuit

Matthew A. Broome<sup>1,2\*</sup>, Alessandro Fedrizzi<sup>1,2</sup>, Saleh Rahimi-Keshari<sup>2</sup>, Justin Dove<sup>3</sup>, Scott Aaronson<sup>3</sup>, Timothy C. Ralph<sup>2</sup>, Andrew G. White<sup>1,2</sup>

Quantum computers are unnecessary for exponentially efficient computation or simulation if the Extended Church-Turing thesis is correct. The thesis would be strongly contradicted by physical devices that efficiently perform tasks believed to be intractable for classical computers. Such a task is boson sampling: sampling the output distributions of  $n$  bosons scattered by some passive, linear unitary process. We tested the central premise of boson sampling, experimentally verifying that three-photon scattering amplitudes are given by the permanents of submatrices generated from a unitary describing a six-mode integrated optical circuit. We find the protocol to be robust, working even with the unavoidable effects of photon loss, non-ideal sources, and imperfect detection. Scaling this to large numbers of photons should be a much simpler task than building a universal quantum computer.

A major motivation for scalable quantum computers are realistic physical devices, then the

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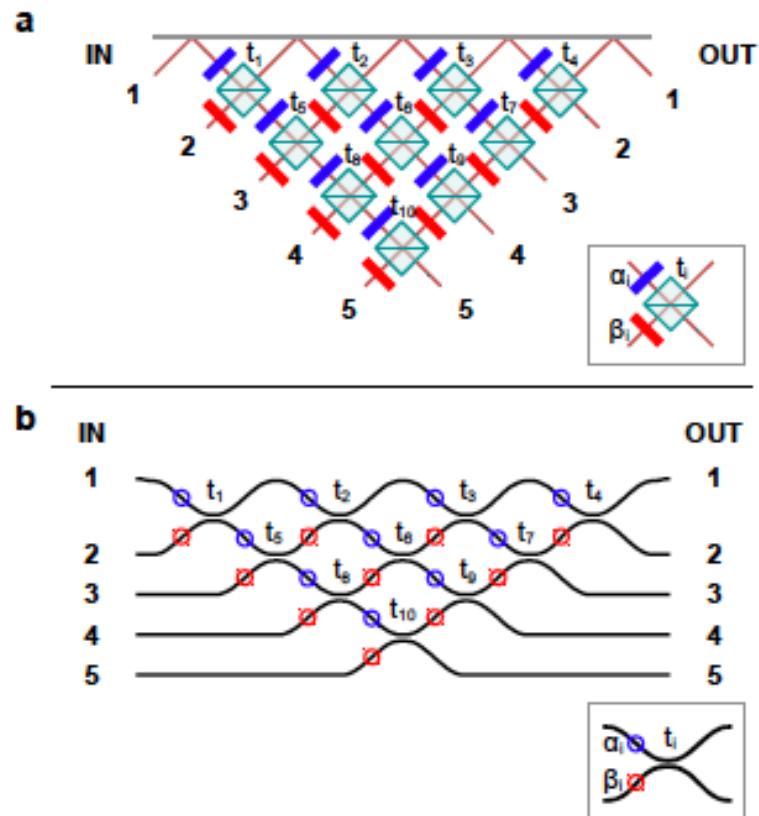
# BosonSampling: the chip

Requirement for Boson Sampling -  
design arbitrary interferometers

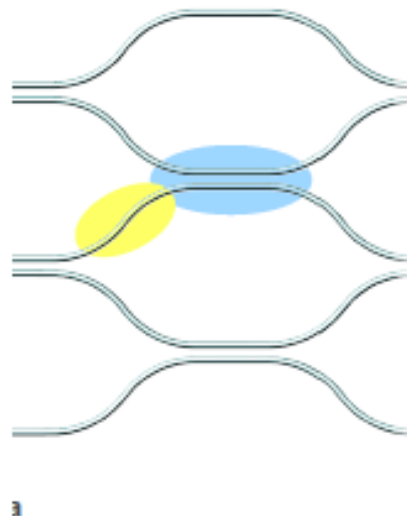


Requires independent control of  
phases and beam-splitter operation

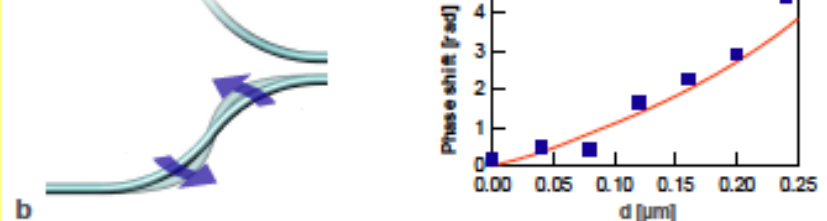
## Architecture for arbitrary unitary



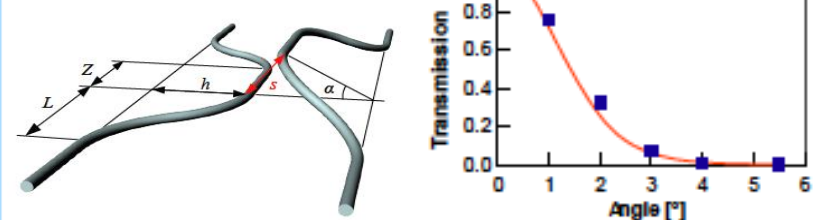
## Fabrication process



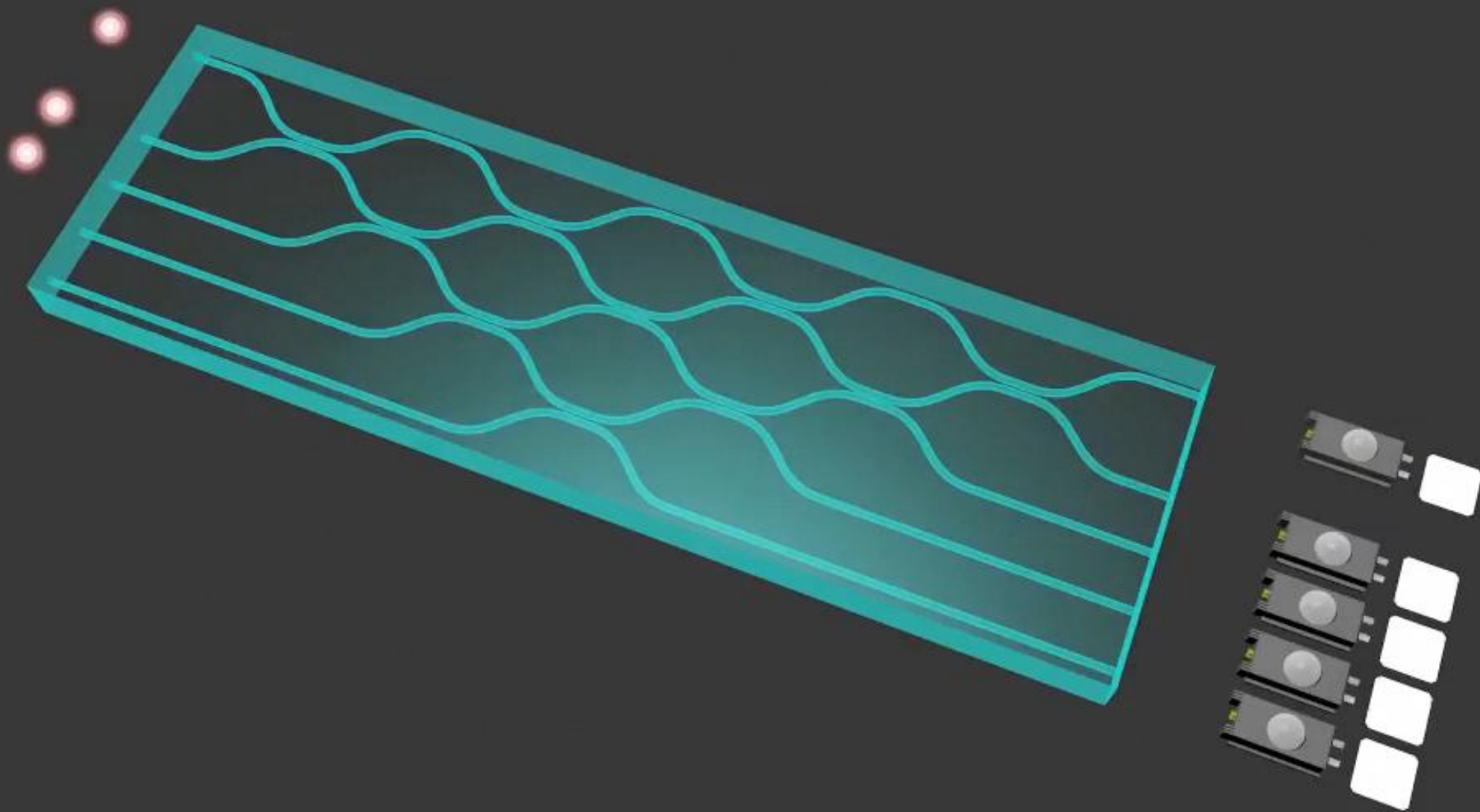
### PHASE-SHIFT CONTROL



### TRANSMISSION CONTROL

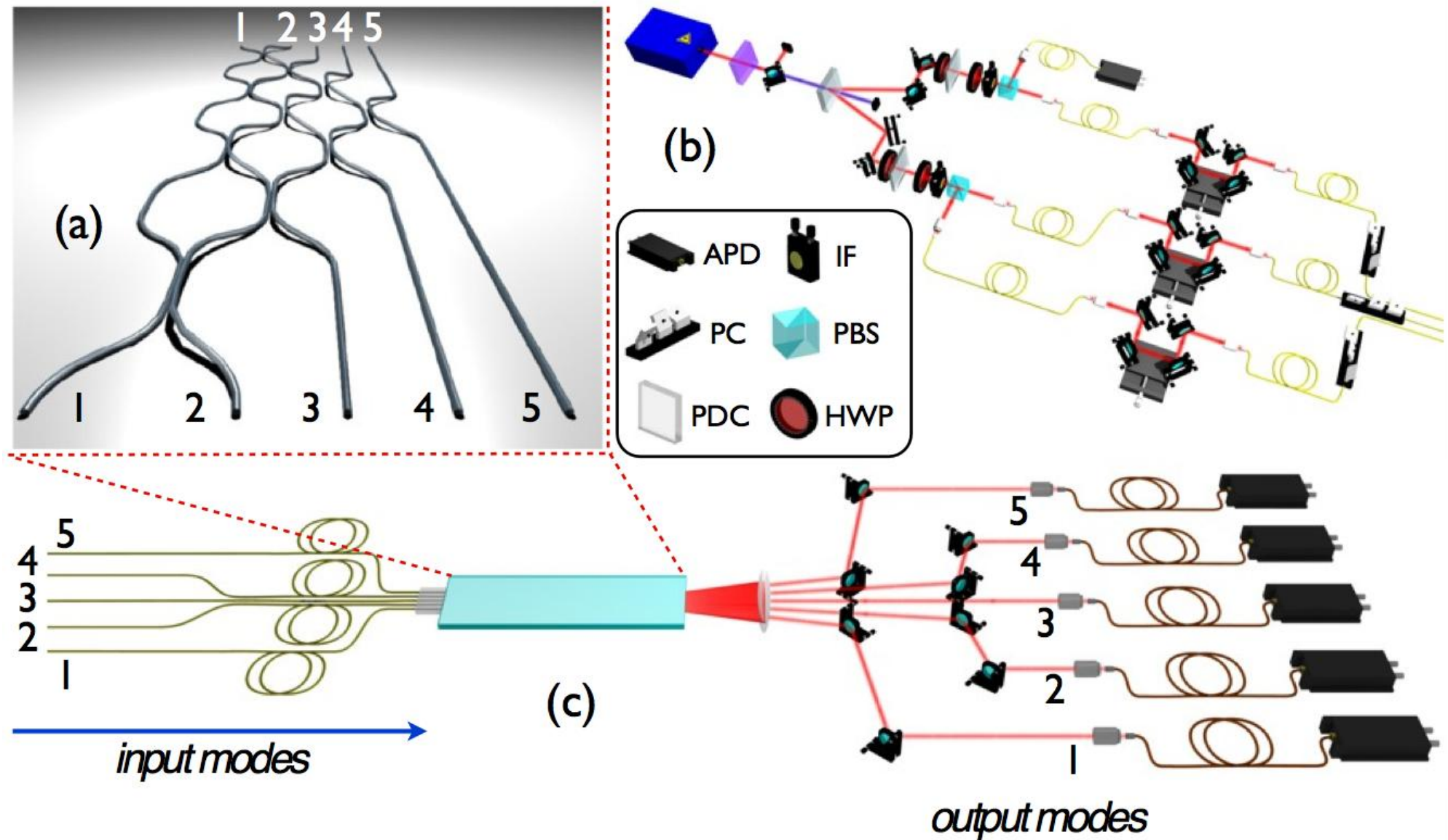


# BosonSampling



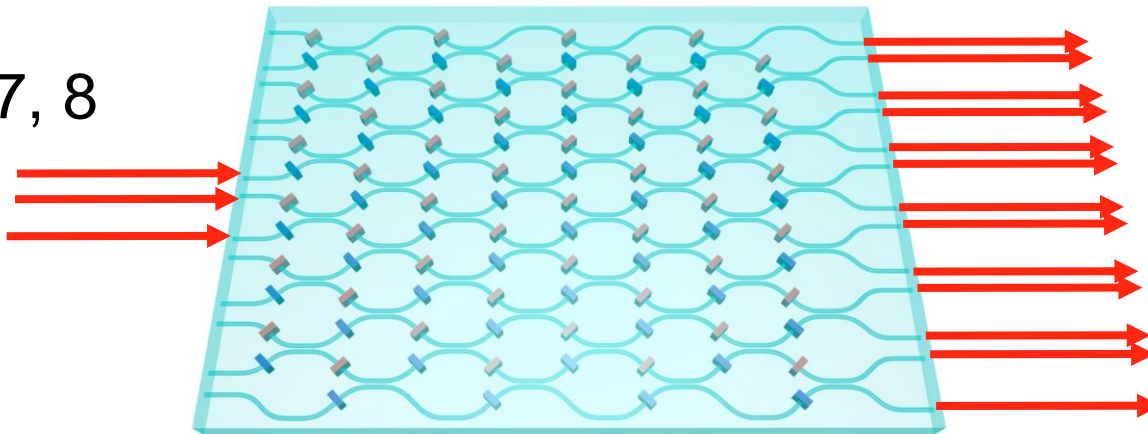


# BosonSampling: the setup

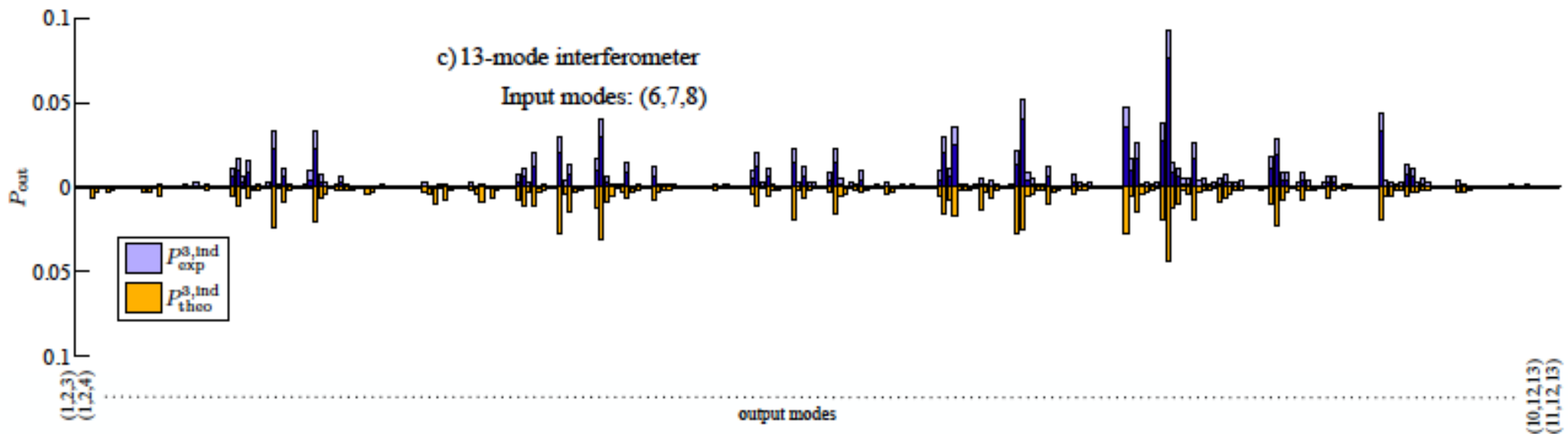


# Up to 13 mode systems investigated

Input: 6, 7, 8



Output: 286 different possible no-bunching configurations



## Can Boson Sampling be validated?

It has been argued that due to the high complexity, BosonSampling output in the hard-computational regime cannot be distinguished from the random output of a uniform distribution

C. Gogolin et al. *arXiv:1306.3995*

## The Theorists' Answer

For each single registered event, which identifies the output state, calculate the quantity

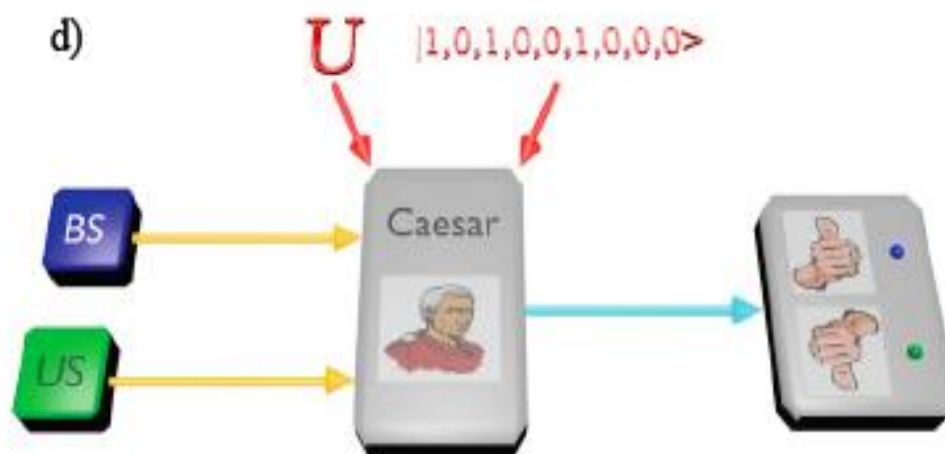
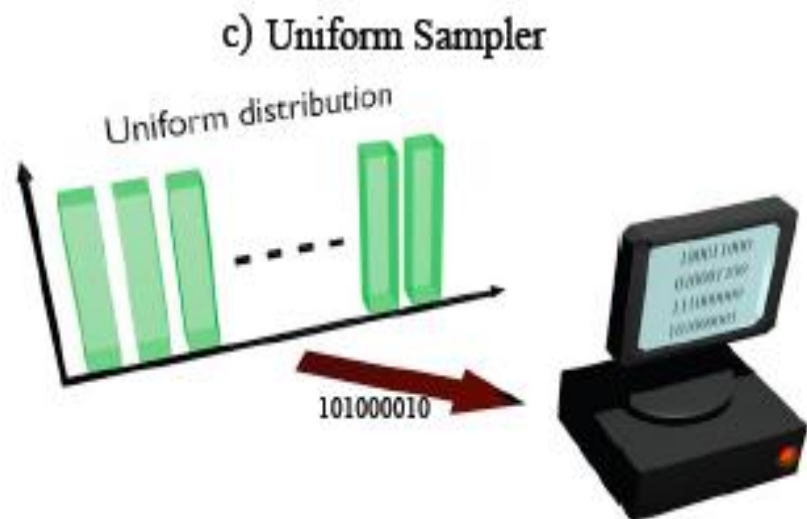
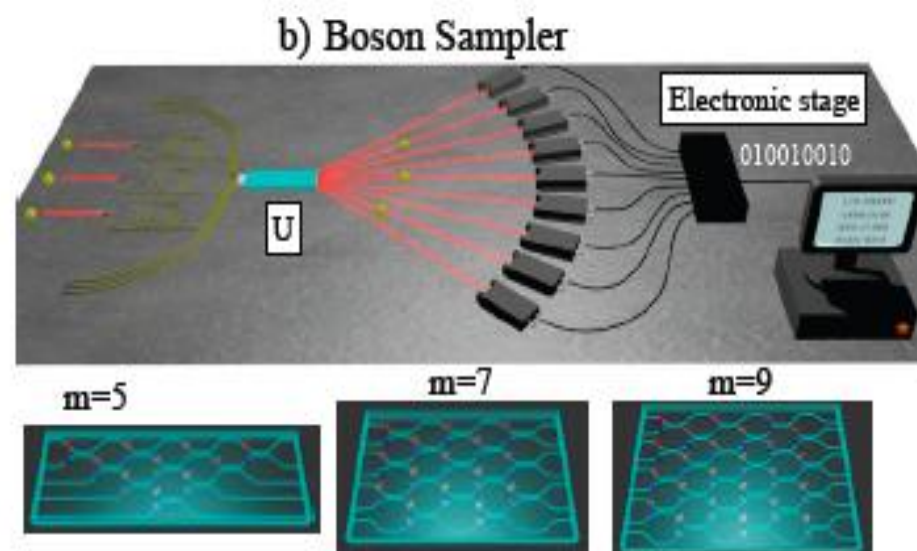
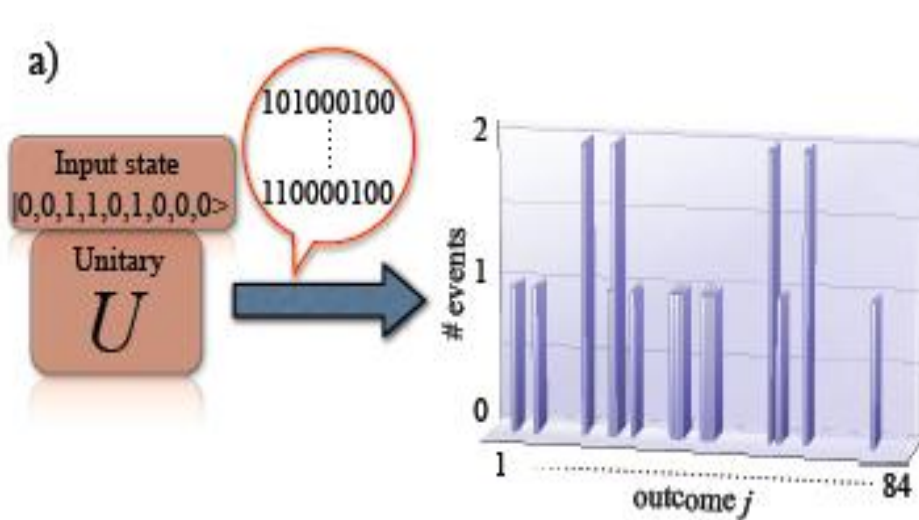
$$P = \prod_{i=1}^n \sum_{j=1}^n |A_{i,j}|^2$$

whith  $A_{i,j}$  = submatrix of  $U$  depening on the input and output states, and compare this value to its counterpart for a uniform distribution  $P_u$ .

If  $P > P_u$ , you can guess that the single event has been produced by a BosonSampler

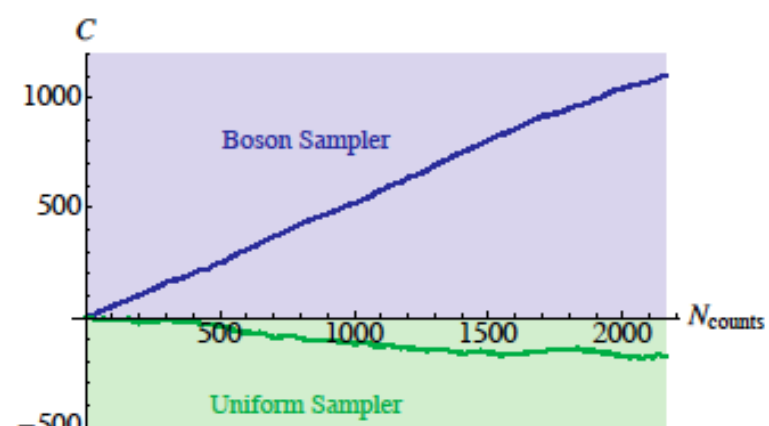
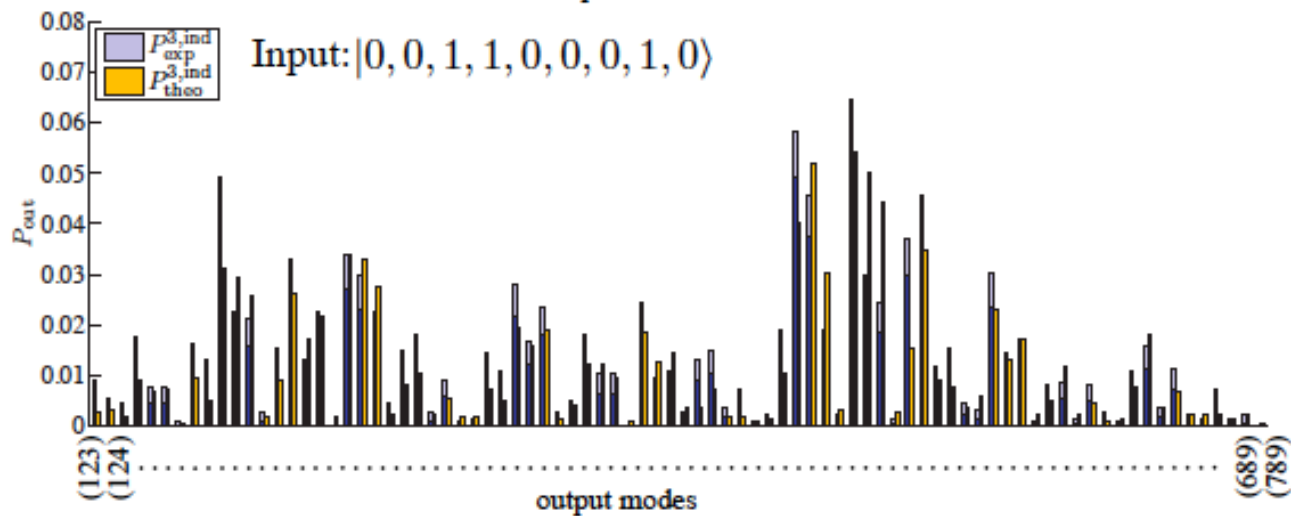
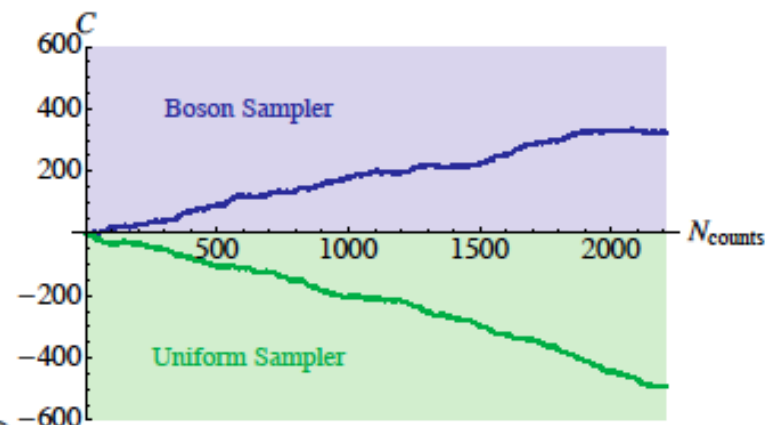
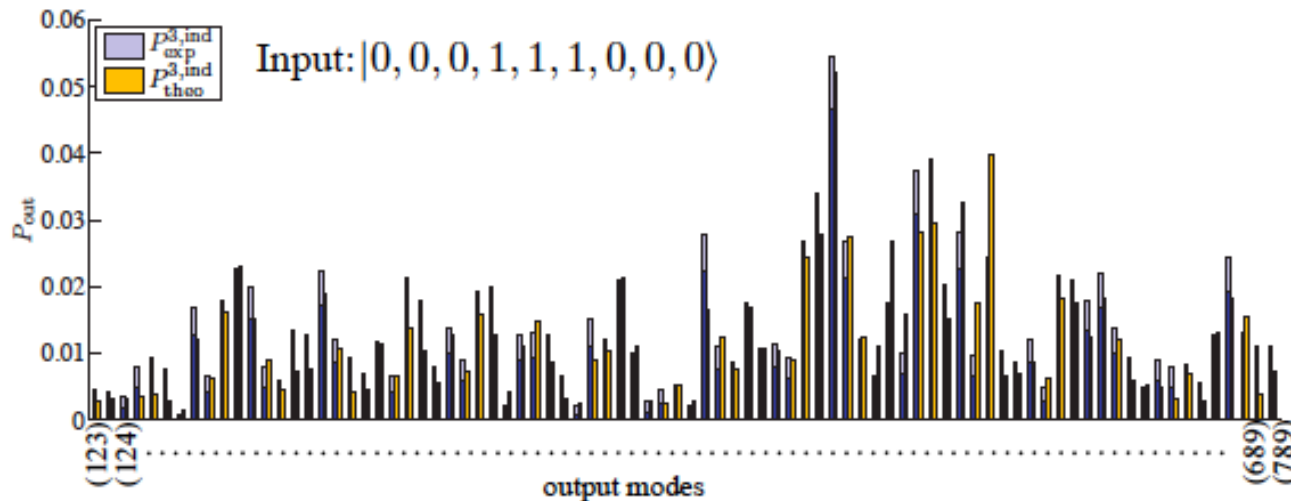


# Validation of BosonSampling



# Validation of BosonSampling with 9 modes

c) 9-mode interferometer



# Towards Quantum Supremacy

The extended Church-Turing (ECT) Thesis:

*Everything feasibly computable in the physical world is feasibly computable by a (probabilistic) Turing machine*

**GOAL:** Achieve Boson Sampling with  $n = 10\text{-}20$  photons  
and  $m = 100\text{-}200$  modes

## Open questions

- *Measure BS complexity*
- *Other equivalent experimental schemes*
- *Certify the functioning of a BS experiment*
- *How noise/imperfections affect a complex BS*

## Challenges

- *Efficient single photon sources*
- *Reconfigurable photonic circuits*
- *Efficient single photon detectors*

# Scattershot BosonSampling

$p$  = probability of generating a photon pair in a single source  
(typical values  $p=0.01-0.015$ )

$p^n$  probability of generating the  $n$ -photon  
input

Scattershot Boson Sampling,  $n$ -photon  
term

$p^n(1-p)^{m-n}$  probability of generating one of  
the  $n$ -photon input configurations

$\binom{m}{n}$  number of possible output  
configurations

Total generation  
rate:  $\sim p^n(1-p)^{m-n} \binom{m}{n}$

Sample both from the *input*  
and the *output modes*



**Potentially huge increase  
of the brightness of the  
quantum hardware**

**See Niko Viggianiello's Poster on Scattershot BosonSampling!!!**

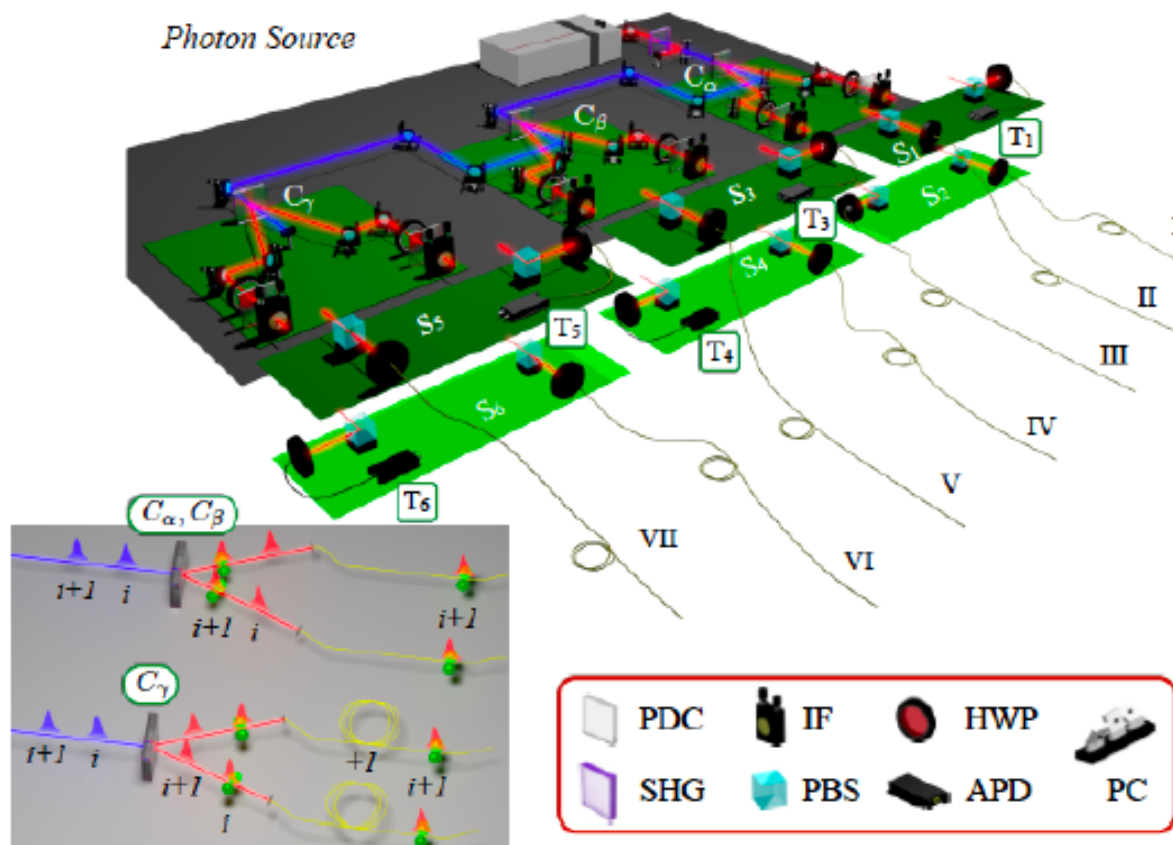
# Scattershot BosonSampling: generation

## Experimental setup - 1

*Photon source*

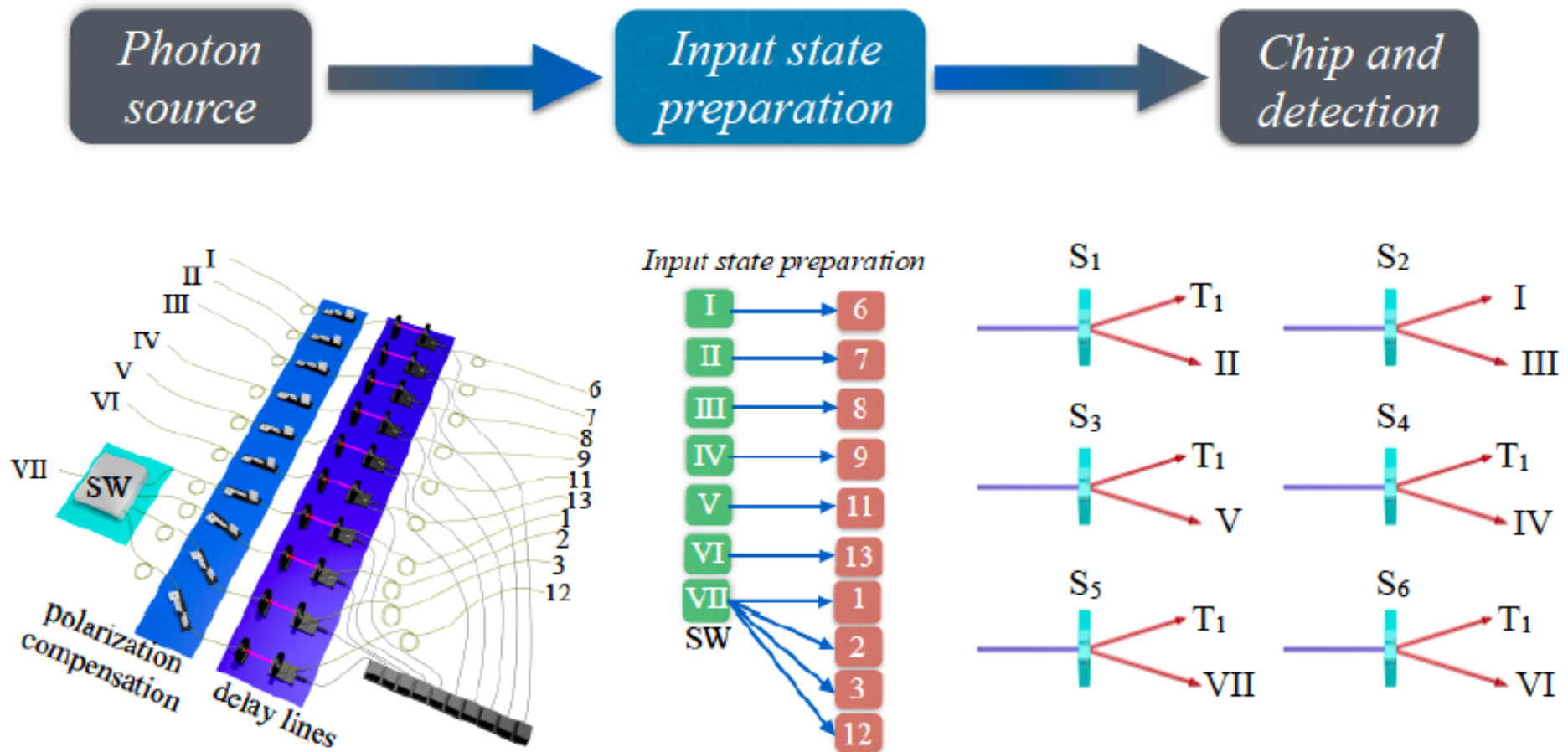
*Input state preparation*

*Chip and detection*



# Scattershot BosonSampling: preparation

## Experimental setup - 2



Three photon events:

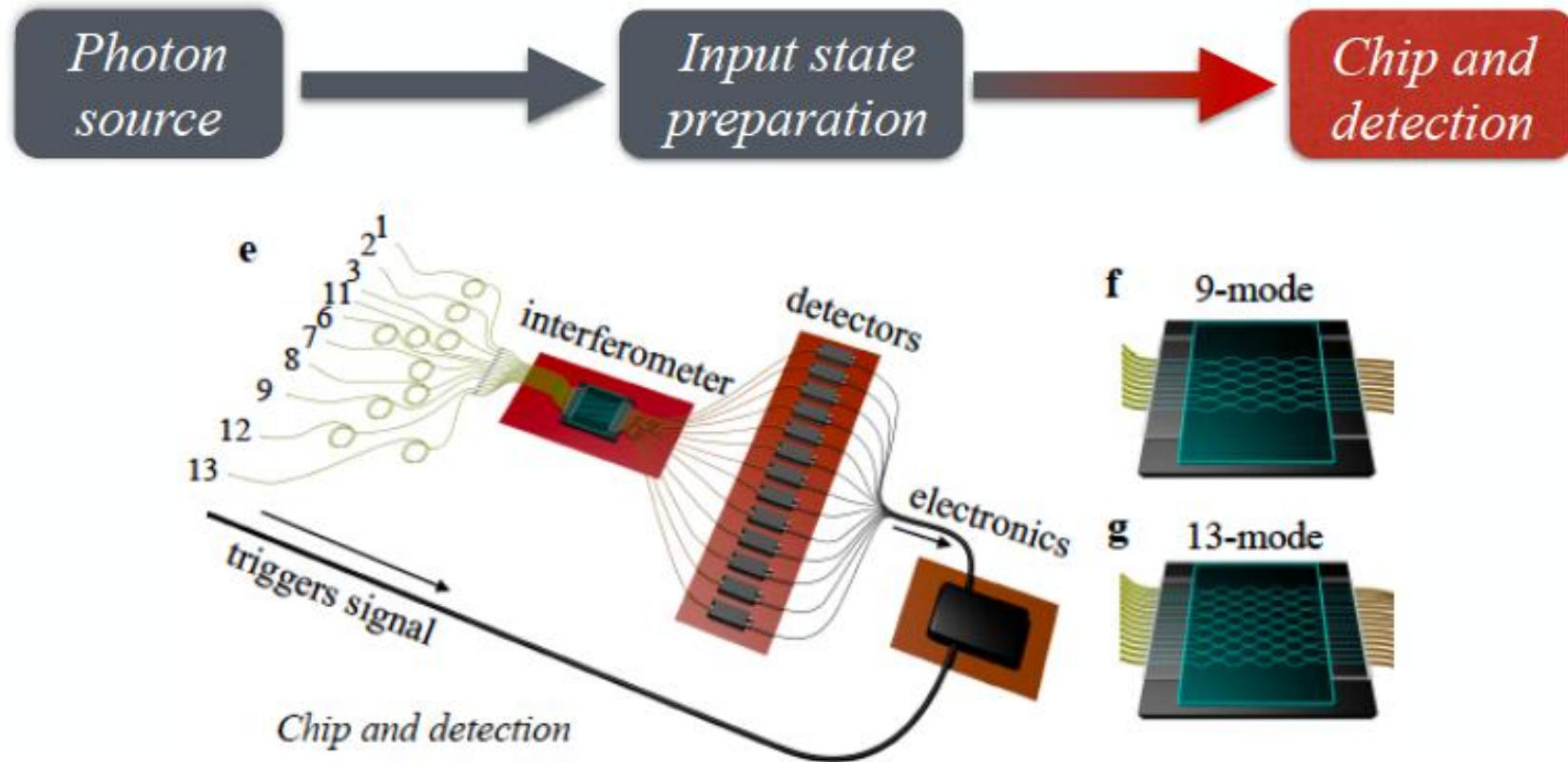
- 1) Photon I (input 6) [fixed]
- 2) Photon III (input 8) [fixed]
- 3) Random input heralded by  $T_i$

Input randomness further enhanced by sequential switching of photon VII



# Scattershot BS: chip and detection

## Experimental setup - 3



Evolution through  $m=9$  and  $m=13$  interferometers with random (but known) structure

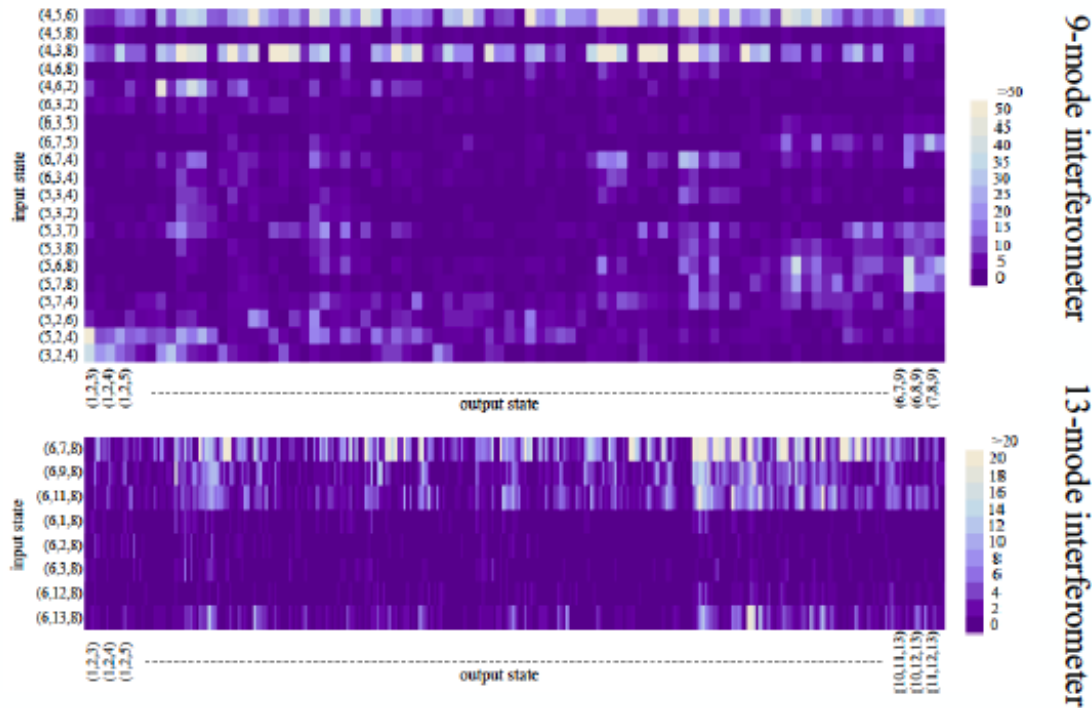
Coincidence detection for:

- Three-photon events with one heralding trigger
- Two-photon events with two heralding triggers

# Scattershot BS: random input

## Scattershot - sampling with random input

### Three-photon events



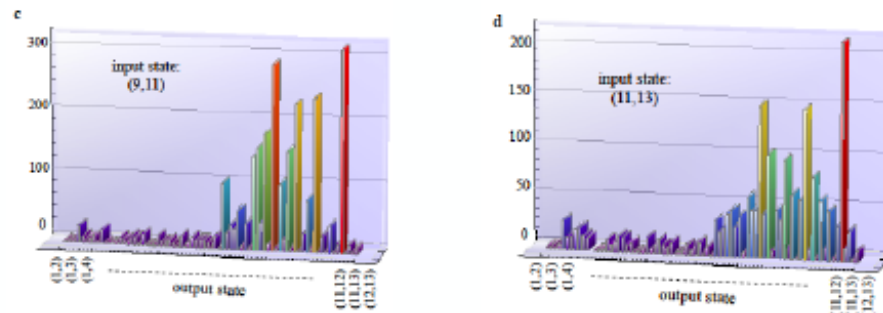
### Number of input/output configurations

m=9 interferometer  
2288 combinations

m=13 interferometer  
1680 combinations

Few events per input/  
output configurations

### Two-photon events



Data with three-photon and  
two-photon input collected  
simultaneously

# Implementing the Fast Fourier Transform

Boson Sampling experiments pose serious problem of certification of the result's correctness in the computationally-hard regime.

Use 3-D photonic chips to test true  $n$ -photon interference in a multimode device [by M.C. Tichy *et al.* (Phys. Rev. Lett., 2014)].

Proposal based on the suppression of specific output configurations in an interferometer implementing an  $n^D$ -dimensional Quantum Fourier Transform (QFT) matrix.

Generalization of the 2-photon/2-modes Hong-Ou-Mandel (HOM) effect, used to test a wide range of photonic platforms.

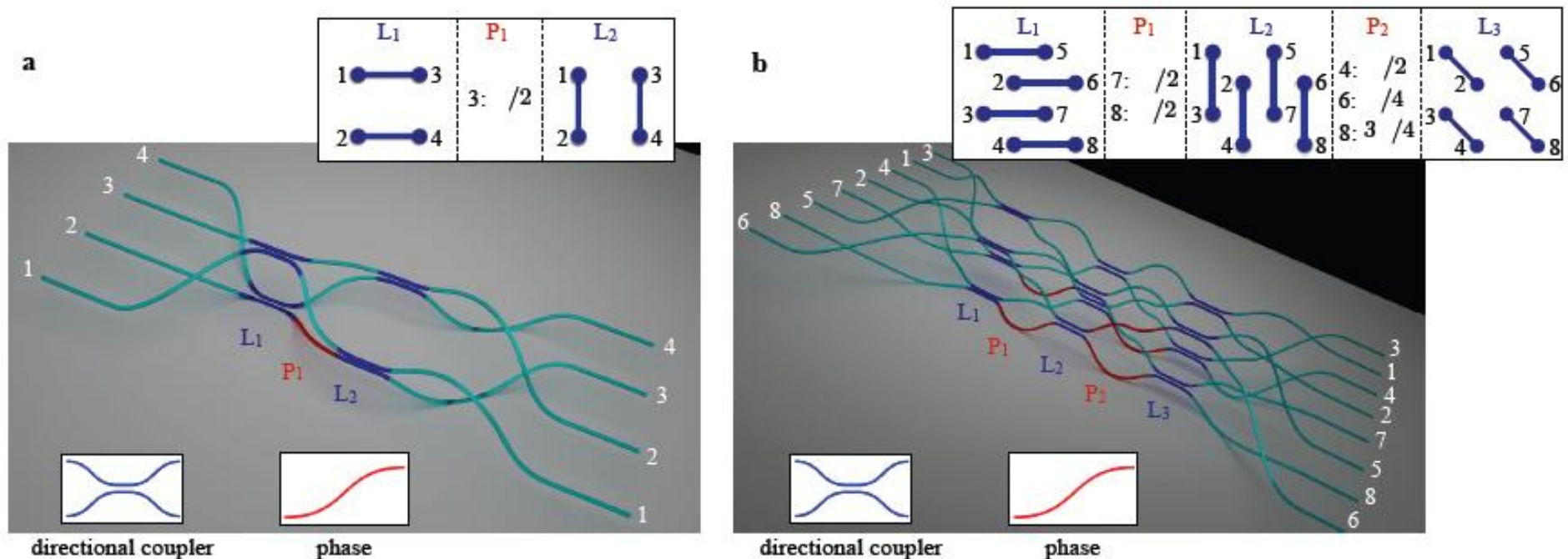
# Quantum Fourier Transform in a 3D chip

Quantum interference in multimode interferometers may determine suppression of a large fraction of the output configurations.

Study the evolution of particular input states through the network implementing the QFT described by the unitary matrix:

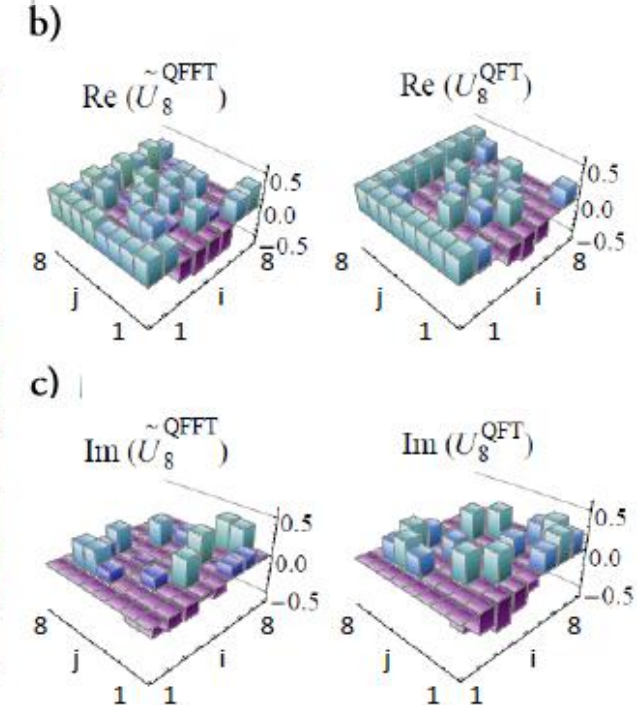
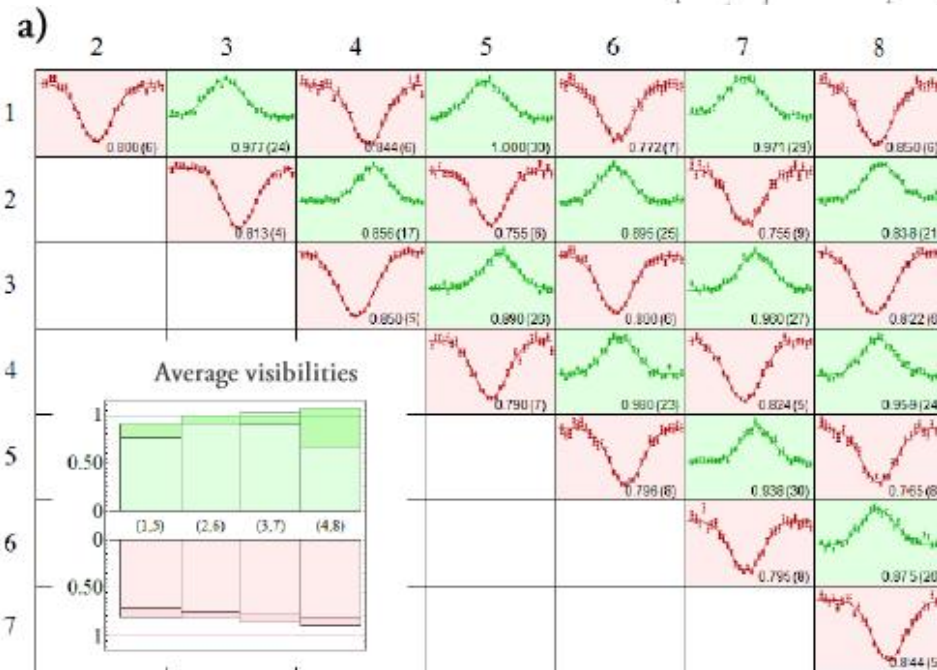
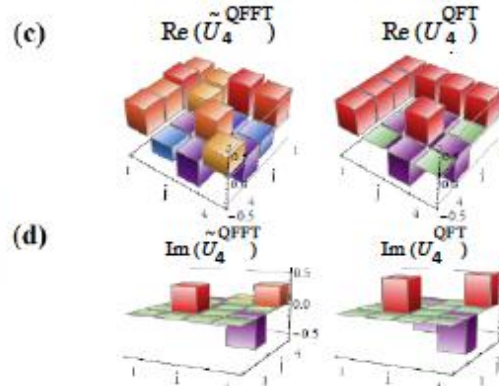
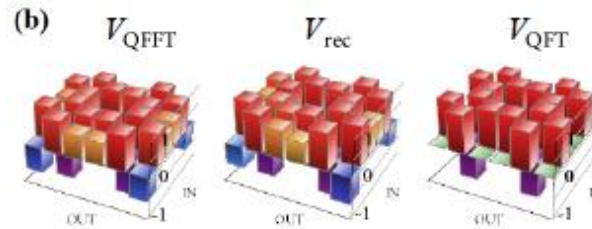
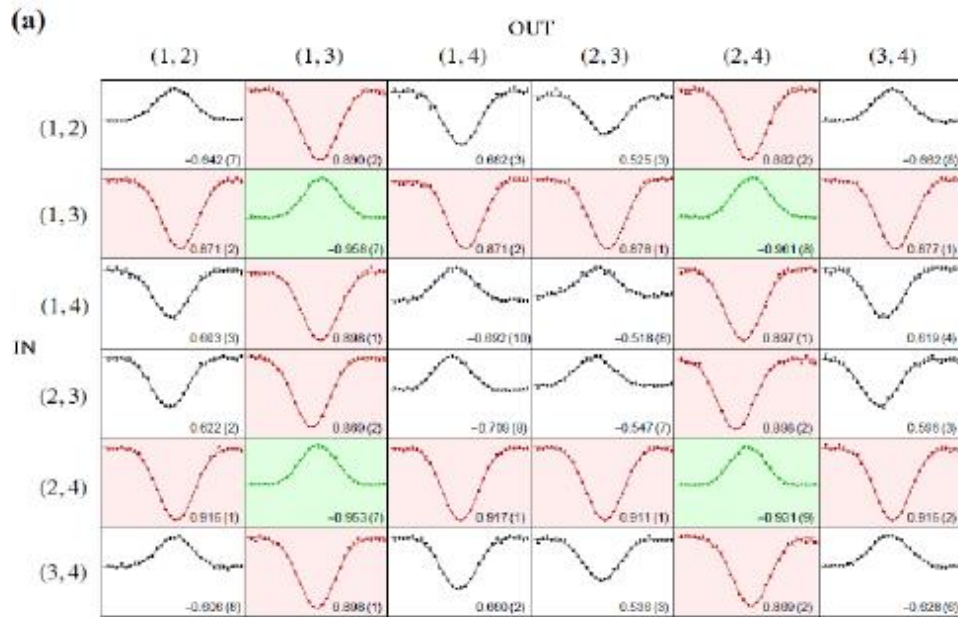
$$U_{l,q}^{\text{QFT}} = \frac{1}{\sqrt{m}} e^{i \frac{2\pi l q}{m}}$$

Test performed with 2 photon and 4- and 8- mode interferometers





# Quantum suppression law in a 3D chip



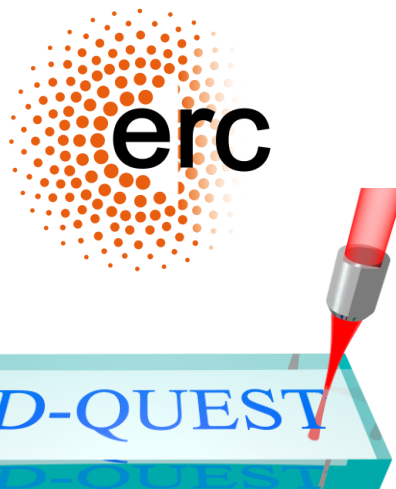
# Conclusions

- Flexibility of waveguide integrated circuits  
(in particular 3D capabilities of fsec laser writing).
- Potential for quantum simulations and quantum walks  
(2 non-interacting bosons/fermions)
- BosonSampling (proof-of-principle test, validation, scattershot BS, Quantum Suppression Law)





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