Quantum simulation of lattice gauge theories in AMO systems: challenges and perspectives

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Joint work with **C. Laflamme, P. Zoller** (Innsbruck), **W. Evans, U.-J. Wiese** (Bern), **W. Bietenholz, U. Gerber, H. Megia-Diaz** (Mexico City)

<u>Related works</u> in collaboration with **E. Rico** (Bilbao), **S. Montangero's group** (Ulm), **U.-J. Wiese's group** (Bern)

How can quantum info and HEP interact?

Development of novel, quantum-information-based **numerical methods** (sign-

See talk by Pietro on Saturday

Rico, Pichler, MD et al. PRL 2014 Buyens et al., PRL 2014 Tagliacozzo et al., 2014 Silvi et al., 2014 Use quantum Inforelated concepts in the context of gauge theories

(topological) entanglement entropies, characterization of phase transitions,

Design of **quantum simulation platforms**



What is quantum simulation

Feynman's lecture 'Simulating Physics with Computers'



Today's dish: *Asymptotic freedom* and the continuum limit





Why shall we study **gauge theories** in synthetic system: **some motivations**

From lattice to field theories: the continuum limit

Observe <u>asymptotic freedom</u> in cold atom systems using <u>alkaline-</u> <u>earth atoms</u> exploiting the <u>continuum limit</u>



Gorshkov et al., Nat Phys. 2010; Cazalilla et al., NJP 2009.

Gauge theories are very interesting beasts...

Ubiquitous theoretical framework



Lee, Nagaosa and Wen, RMP 2006; *Introduction to Frustrated magnetism*, Springer 2011; Focus issue on quantum spin liquids, NJP2013. Montvay and Münster, *Quantum fields on a Lattice*

...however, not yet realized in the panorama of synthetic systems...

The reason being that ...

It is very hard to get a **gauge symmetry** in synthetic systems!!!

From a Hamiltonian perspective, a **gauge symmetry** is nothing but a set of **local constraints**

An example of a constrained system: Ising lattice gauge theory - aka the toric code

$$H = \sum_{\Box} \prod_{j \in \Box} \sigma_j^x$$
Hamiltonian

$$G_{+} = \prod_{j \in +} \sigma_{j}^{z} \quad \underbrace{}_{j \in +} \quad \mathbf{Constraints}$$

Flip all four spins around a plaquette. Since two spins the vertex are flipped, the parity is conserved

Wegner, 1971; Kitaev, 2003; Dissipative dynamics: Barreiro, Mueller et al, Nature 2011

Where we stand now - synthetic lattice gauge theories

Ultracold atoms

Tewari et al. (PRL2006); Kapit and Mueller, PRA 2010 Banerjee, MD et al. PRL 2012, PRL 2013 Zohar et al, PRL 2012, PRL 2013, PRA2013 Notarnicola et al., Meurice et al., arxiv.2015;

 $U(1), \mathbb{Z}_N, SU(N), U(N), SO(3), ...$

Digital approaches

Byrnes and Yamamoto, PRA 2006. Weimer et al., Nat. Phys. 2010. Tagliacozzo et al., Ann. Phys. 2012, Tagliacozzo et al., Nat. Comm. 2013

Circuit QED

Marcos et al., PRL 2013, Ann. Phys. 2014 Mezzacapo et al., arxiv2015

Rydberg atoms

A. Glätzle, MD et al., PRX 2014; arXiv2014

Trapped ions

Hauke, MD et al., PRX 2013; Nath, MD et al. arxiv.1504.01474

Quantum Zeno dynamics

Stannigel et al., PRL 2014

How to connect lattice models to field theories?

High Energy - field theories

$$\mathcal{L} = \frac{1}{2g^2} \mathrm{Tr} F_{\mu\nu} F_{\mu\nu} + \dots$$

Fundamental question: the continuum limit

$$H = \sum_{i,j} \psi_i^{\dagger} U_{ij} \psi_j + \dots \qquad \qquad H = \int dx \; \tilde{\psi}^{\dagger}(x) (D_x) \psi(x') + \dots$$

In Wilson's lattice field theory, usually achieved by taking the limit

 $a \rightarrow 0$

The dependence on *a* is present in the Hamiltonian/Lagrangian couplings in a controlled way

Open question: the continuum limit

Open question until now: *can we approach the continuum limit in synthetic quantum systems*?

(2) Unorthodox way: dimensional reduction and asymptotic freedom

The continuum limit: the unorthodox way

$$H = \sum_{i,j} \psi_i^{\dagger} U_{ij} \psi_j + \dots \qquad \qquad H = \int dx \; \tilde{\psi}^{\dagger}(x) (D_x) \psi(x') + \dots$$

Take a theory in one dimension more, and extend this additional dimension until a correlation lengths is much larger than its width

A simple example of dimensional reduction

Intuition: Neel states in 2D antiferromagnets

At T=0, Neel order is stable, and described by an O(2) order parameter

Full-rotation -> rotation around one axis

 $SO(3)/SO(2) \rightarrow S^2$

We get a field theory in 1D as (pseudo-)Goldstone bosons via *spontaneous symmetry breaking*

B.B. Beard, et al., Phys. Rev. Lett. 94, 010603 (2005); N. Read and S. Sachdev, Phys. Rev. Lett. 62, 1694 (1989).

The continuum limit: the unorthodox way

$$H = \sum_{i,j} \psi_i^{\dagger} U_{ij} \psi_j + \dots \qquad \qquad H = \int dx \; \tilde{\psi}^{\dagger}(x) (D_x) \psi(x') + \dots$$

Take a theory in one dimension more, and extend this additional dimension until a correlation lengths is much larger than its width

Why CP(N-1) models?

$$S[P] = \int_0^\beta dct \int_0^L dx \, \frac{1}{g^2} \text{Tr}[\partial_x P \partial_x P + \frac{1}{c^2} \partial_t P \partial_t P] - i\theta Q[P],$$

A.D'Adda, M. Luescher, and P. Di Vecchia, Nucl. Phys. B 146, 63 (1978).

Asymptotic freedom: what is that?

Interactions between particles gets weaker and weaker at short distances [= high energies = high momenta]

Formally, identified with the first coefficient of the beta function being negative

$$\beta(g) = \beta_0 * g^3$$
$$g^2(k) = \frac{g^2}{1 + \beta_0 g^2 \log[k^2/M^2]}$$

1974: Politzer, Wilczek and Gross: asymptotic freedom is fundamental in QCD!

Dimensional reduction in CP(N-1) models: microscopics

Dimensional reduction from a 2+1d SU(N) antiferromagnet

NB: the two representations coincide for the N=2 case / Heisenberg model.

Dimensional reduction in CP(N-1) models

Dimensional reduction from a 2+1d SU(N) antiferromagnet

$$= -J \sum_{a=1}^{N^2 - 1} \sum_{\langle xy \rangle} T_x^a T_y^{a*}$$

 $[T_x^a, T_y^b] = \mathrm{i}\delta_{xy}f_{abc}T_x^c$

H

For L'>>1, Goldstone bosons live in a coset space

$$SU(N)/U(N-1) = CP(N-1)$$

Which yields at low-energies a 1+1d theory:

$$S[P] = \int_0^\beta dct \int_0^L dx \, \frac{1}{g^2} \operatorname{Tr}[\partial_x P \partial_x P + \frac{1}{c^2} \partial_t P \partial_t P] - i\theta Q[P],$$
$$\theta = \pi n$$

B.B. Beard, et al., Phys. Rev. Lett. 94, 010603 (2005); N. Read and S. Sachdev, Phys. Rev. Lett. 62, 1694 (1989).

Intuition: Neel states in 2D antiferromagnets

At T=0, Neel order is stable, and described by an O(2) order parameter

Full-rotation -> rotation around one axis

$$SO(3)/SO(2) \rightarrow S^2 = CP(1)$$

B.B. Beard, et al., Phys. Rev. Lett. 94, 010603 (2005); N. Read and S. Sachdev, Phys. Rev. Lett. 62, 1694 (1989).

Implementing SU(N) AF magnets with alternate representations: N=3

Gorshkov et al., Nat Phys. 2010; Cazalilla et al., NJP 2009; for a review, see Cazalilla and Rey, Rep. Progr. Phys. 2014

How to observe asymptotic freedom (N=3)?

Emergence of asymptotic freedom

$$g^{2}(k) = \frac{g^{2}}{1 + \beta_{0}g^{2}\log[k^{2}/M^{2}]}$$
$$M \simeq 1/\xi \qquad g^{2}(k) \simeq 1/L'$$
$$\xi \propto \exp(4\pi L'\rho_{s}/N) \gg L'$$

QMC results using loop updates up to V=400 x 20

Can be measured by Bragg spectroscopy, microscope, noise correlations,...

B.B. Beard, et al., Phys. Rev. Lett. 94, 010603 (2005)

Theta angle and decay of false vacua

For odd number of chains, the system gets gapped and dimerization takes place

Conclusions

Synthetic systems are now approaching the level of control for realizing gauge theories

Synthetic quantum magnetism can be exploited to realize the continuum limit of CP(N) models

Asymptotic freedom of CP(N) can be observe using conventional probes

Thank you

Catherine Laflamme

Wynne Evans

Uwe-Jens Wiese

Urs Gerber

Hector Megia-Diaz

C. Laflamme, P. Zoller (Innsbruck), W. Evans, U.-J. Wiese (Bern), W. Bietenholz, U. Gerber, H. Megia-Diaz (Mexico City), arxiv.1507.06788

<u>Atoms U(1):</u> Banerjee, MD, Mueller, Rico, Stebler, Wiese, Zoller, PRL **109**, 175602; <u>Atoms SU(N)/U(N)</u>: Banerjee, Boegli, MD, Rico, Stebler, Wiese, Zoller, PRL **110**, 125303; <u>Quantum Zeno</u>: Stannigel, Hauke, Marcos, Hafezi, Diehl, MD, Zoller, PRL **112**, 120406; <u>Trapped Ions</u>: Hauke, Marcos, MD, Zoller, PRX **3**, 041018 (2013); <u>TN approaches</u>: Rico, Pichler, MD, Zoller, Montangero, PRL **112**, 201601, arXiv.2015 Rydbergs: Glätzle, MD, Nath, Rousochatzakis, Moessner, Zoller, PRX, **4**, 041037.