Current status of carlomat_3.0, an automatic tool for low energetic electron-positron annihilation into hadrons

Karol Kołodziej

Institute of Physics University of Silesia Katowice

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The most promising theoretical frameworks for the descritpion of $e^+e^- \to {\rm hadrons}$ in the energy range below the J/ψ threshold are the Resonance Chiral Theory or Hidden Local Symmetry model which, among others, involve

- the photon-vector meson mixing
- a number of vertices of rather complicated Lorentz tensor structure that is not present in the Standard Model or scalar QED.

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- include the Feynman interaction vertices of new tensor structures,
- to compute the helicity amplitudes involving the mixing terms and new interaction vertices, like those predicted by the $R\chi T$ or HLS model and the effective Lagrangian of the EM interaction of nucleons,
- introduce new options to enable a better control over the effective models implemented.

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$$A^{\mu}(q) \qquad V^{\nu}(q) \\ \longleftarrow \qquad \equiv -e f_{AV}(q^2) g^{\mu V},$$

with
$$V=\rho^0, \omega, \phi, \rho_1=\rho(1450), \rho_2=\rho(1700).$$

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The actual names for that option in carlocom are: imrho, imome, imphi, imrh1, imrh2 for $\rho^0, \omega, \phi, \rho_1, \rho_2$ meson, respectively.

The complex factor c_j , j = 1, 2, ... is given by

$$c_j = w_j e^{i\phi_j} f_j(q^2),$$

where w_j is a positive weight, φ_j is an angle in degrees, which should be both specified for each possible particle mixing term in the main program for the MC computation carlocom, and $f_j(q^2)$ is a possible four momentum transfer dependence that is defined in subroutine weightfactor.

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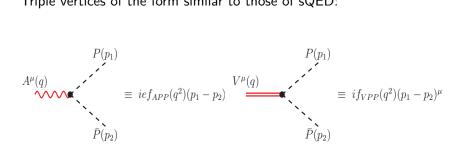
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q^2 -dependent interaction vertices

Triple vertices of the form similar to those of sQED:



where $V = \rho^0, \omega, \phi$ and $P = \pi^+, K^+, K^0$.

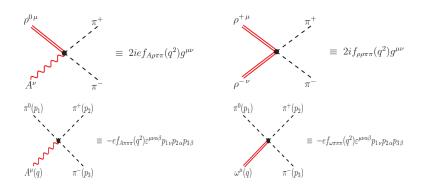
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Triple vertices of a more complicated tensor form:

where, in the top right corner, $V=\rho^0, \omega$, and in the bottom right corner $P=\pi^0$ and $V=\rho^0$ or $P=\pi^\mp$ and $V=\rho^\pm$.

q^2 -dependent interaction vertices

Quartic vertices of the HLS model:



The vertices in the first row have the same tensor form as the quartic vertex of the sQED or the quartic vertices of the Nambu-Goldstone boson – gauge boson interaction of the SM, which were implemented already in the first version of carlomat.

All subroutines that are used to compute the building blocks or the complete helicity amplitudes of the Feynman diagrams containing vertices of the HLS model have been supplied with the running coupling option.

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The options are to be specified in subroutine couplsm, where they are defined below the assignment for each particular coupling. icoupl_name=0/1,2,... if the fixed/running coupling is to be used in the computation, where choices 1,2,... corresponding to different running couplings f_{...}(q^2) should be added by the user as extra else if (ig == ...) then blocks in subroutine runcoupl.
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Many couplings of the R χ T or HLS model are not known well enough and must be adjusted in consecutive runs of the program in order to obtain satisfactory description of the experimental data. If there are no hints as to the form of the running couplings $f_{...}(q^2)$ then it is recommended to set the corresponding running coupling option to 0, which means that the fixed coupling is to be used in the computation.

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The main part of the MC computation program carlocom contains also a few flags: iarho, iaome, iaphi, iarho1 and iarho2 that allow to switch off and on the photon mixing with $\rho, \omega, \phi, \, \rho_1$ and ρ_2 vector mesons without a need of running the code generation program anew, provided that the corresponding mixing terms were included in a file vertices.dat when the MC code was generated.

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The subroutines for computation of the four vectors representing vector mesons have been in addition supplied with the running width option iwdth name, i.e. igmrh, igmom, igmph, igmr1, igmr2 for the running width of ρ^0 , ω , ϕ , ρ_1 , ρ_2 , respectively.

where choices 1,2,3 refer to different running width options in subroutine runwidth which again can easily be extended by the user. The options are controlled from carlocom, the main part of the MC computation program.

An important new option in the program that allows to test the EM gauge invariance for processes with one ore more external photons is igauge in carlocom.f:

where 1,2,... is the number of a photon, counting from left to right, whose polarization four vector is replaced with its four momentum.

To illustrate how this option can be used in practice, consider the following radiative processes:

$$\begin{array}{ccc} e^+e^- & \rightarrow & \pi^+\pi^-\mu^+\mu^-\gamma, \\ e^+e^- & \rightarrow & \pi^+\pi^-\pi^+\pi^-\gamma. \end{array}$$

Taking into account the Feynman rules of SM, without the Higgs couplings to electrons and muons, sQED, the $\gamma-\rho^0$ mixing and the vertices: $\gamma\pi^+\pi^-$, $\rho^0\pi^+\pi^-$, $\pi^0\gamma\gamma$, $\pi^0\gamma\rho^0$, $\gamma\rho^0\pi^+\pi^-$ and $\gamma\pi^0\pi^+\pi^ \Rightarrow$ 209 and 774 Feynman diagrams.

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The cross sections in pb at $\sqrt{s}=1$ GeV without and with contribution from the $\pi^{\mp}\gamma p^{\pm}$ vertices with the following cuts:

$$\theta_{\gamma I} > 5^{\circ}, \qquad \theta_{\gamma \pi} > 5^{\circ}, \qquad \textit{E}_{\gamma} > 10 \; \mathrm{MeV}. \label{eq:epsilon}$$

igauge	$\sigma(e^+e^- ightarrow\pi^+\pi^-\pi^+\pi^-\gamma)$		$\sigma(e^+e^- ightarrow \pi^+\pi^-\mu^+\mu^-\gamma)$	
0	11.86(5)	11.83(5)	0.0590(2)	0.0586(2)
1	0.124(2)e-30	0.441(1)e-10	0.636(9)e-33	0.973(1)e-9

The numbers in parentheses show the MC uncertainty of the last decimal.

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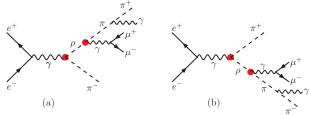
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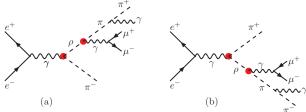


Amplitudes:

$$\begin{array}{lcl} M_{a} & = & \mathbf{g^{2}} \epsilon_{12\nu} \, \epsilon^{\nu\mu\alpha\beta} p_{12\alpha}(-q_{\beta}) \, \frac{-g_{\mu\rho} + \frac{q_{\mu}q_{\rho}}{M^{2}}}{q^{2} - M^{2}} \, \epsilon_{56\sigma} \, \epsilon^{\sigma\rho\gamma\delta}(-p_{56\gamma}) q_{\delta} \, s_{37}, \\ M_{b} & = & \mathbf{g^{2}} \epsilon_{12\nu} \, \epsilon^{\nu\mu\alpha\beta} p_{12\alpha}(-r_{\beta}) \, \frac{-g_{\mu\rho} + \frac{r_{\mu}r_{\rho}}{M^{2}}}{r^{2} - M^{2}} \, \epsilon_{56\sigma} \, \epsilon^{\sigma\rho\gamma\delta}(-p_{56\gamma}) r_{\delta} \, s_{47} \end{array}$$

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The scalars s_{37} and s_{47} representing the $\pi^+\pi^-\gamma$ and $\pi^-\pi^+\gamma$ vertex multiplied with the adjacent pion propagator, take the following form:

$$s_{37} = e^{\frac{(p_3 + p_7 - (-p_3))^{\mu} \varepsilon_{\mu}^*(p_7)}{(p_3 + p_7)^2 - m_{\pi}^2}} \bigg|_{\varepsilon(p_7) \to p_7} = e^{\frac{2p_3 \cdot p_7}{2p_3 \cdot p_7}} = e,$$

$$s_{47} = e^{\frac{(-p_4 - (p_4 + p_7))^{\mu} \varepsilon_{\mu}^*(p_7)}{(p_4 + p_7)^2 - m_{\pi}^2}} \bigg|_{\varepsilon(p_7) \to p_7} = -e^{\frac{2p_4 \cdot p_7}{2p_4 \cdot p_7}} = -e.$$

Hence

$$\begin{array}{lcl} \mathit{M}_{a} & = & \frac{eg^{2}}{q^{2}-\mathit{M}^{2}} \, \epsilon^{\mu\nu\alpha\beta} \, \epsilon_{\mu\sigma\gamma\delta} \, \epsilon_{12\nu} \, p_{12\alpha} \, p_{4\beta} \, \epsilon_{56}^{\sigma} \, p_{56}^{\gamma} (p_{3}+p_{7})^{\delta}, \\ \\ \mathit{M}_{b} & = & -\frac{eg^{2}}{r^{2}-\mathit{M}^{2}} \, \epsilon^{\mu\nu\alpha\beta} \, \epsilon_{\mu\sigma\gamma\delta} \, \epsilon_{12\nu} \, p_{12\alpha} \, p_{3\beta} \, \epsilon_{56}^{\sigma} \, p_{56}^{\gamma} (p_{4}+p_{7})^{\delta}. \end{array}$$

 \Rightarrow M_a and M_b neither vanish separately nor cancel each other.

Preparation for running and program usage

carlomat_3.0 is distributed as a single tar.gz archive carlomat_3.0.tgz which can be downloaded from: http://kk.us.edu.pl/carlomat.html.

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 Go to mc_computation, choose the centre of mass energy and required options in carlocom.f and execute make mc in the command line.

Whenever the Fortran compiler is changed, or a compiled program is transferred to another computer with a different processor then all the object and module files should be deleted by executing the commands:

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The basic output of the MC run is written to file tot_name, where name is created automatically if the assignment prcsnm='auto'

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The output files for processes

$$e^+e^- \rightarrow \pi^+\pi^-\mu^+\mu^-\gamma,$$

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with the preselected parameters and options should reproduce those delivered in directory test_output0.

The basic output of the MC run is written to file tot_name, where name is created automatically if the assignment prcsnm='auto'

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If the differential cross sections/distributions are required then set idis=1

in carlocom.f.

The number of distributions to be calculated must be specified in distribs.f and their parameters should be defined in calcdis.f. The output will be stored in data files db#_name and dl#_name which can be plotted with boxes and lines, respectively, with gnuplot.

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Restrictions

As in previous versions of the program the number of particles in the final state is limited to 10 which exceeds typical numbers of particles observed in the exclusive low energy e^+e^- -annihilation processes.

However, in the presence of photon–vector meson mixing, the Feynman diagrams proliferate, for example, with currently implemented Feynman rules, there are 90672 diagrams of $e^+e^- \rightarrow 3(\pi^+\pi^-)$.

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Many couplings of the effective models are not known well enough and must be adjusted in consecutive runs of the program in order to obtain satisfactory description of the experimental data. Thank you for your attention