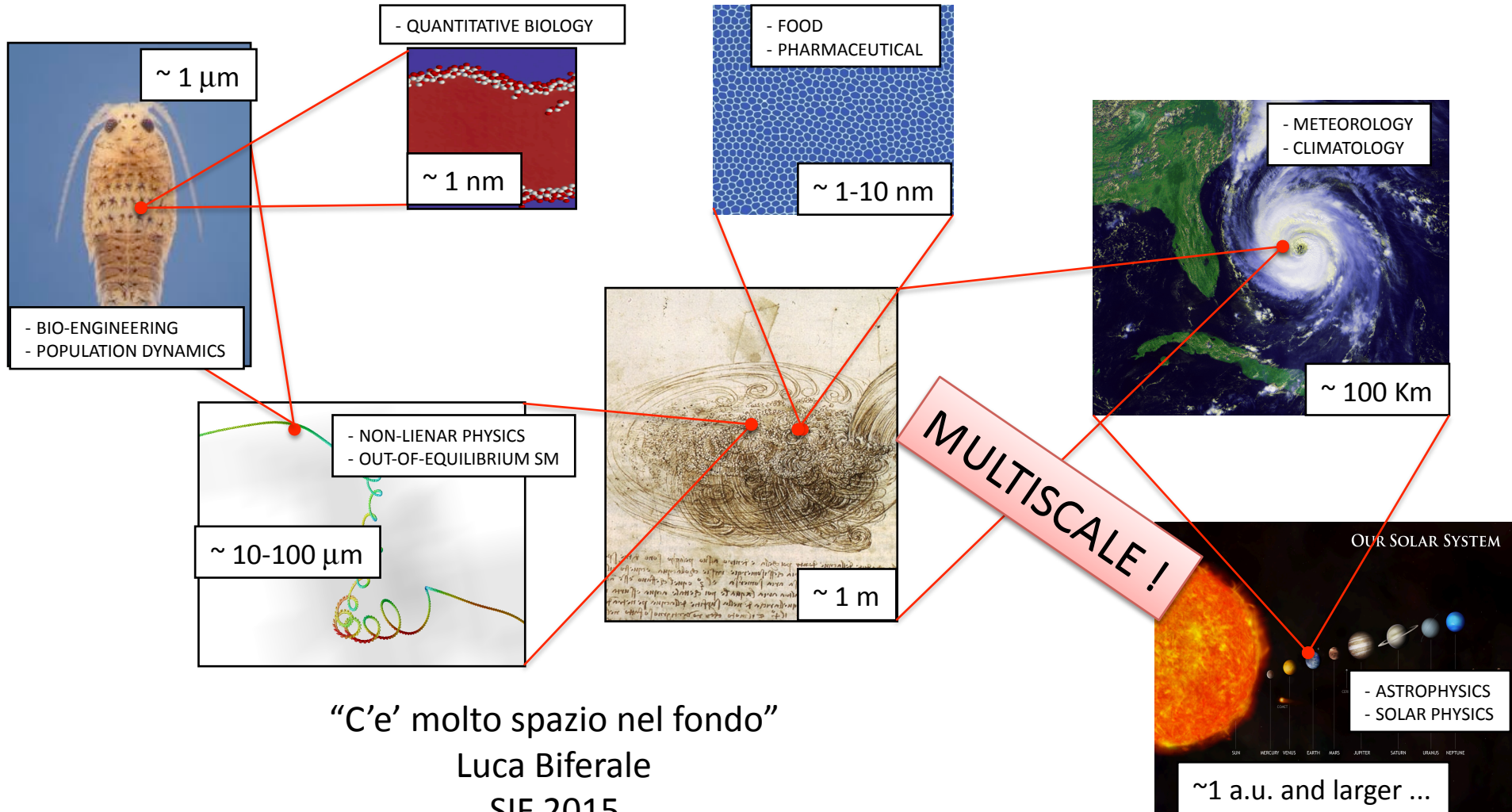


Πάντα ρει

inside, outside, above, below and around us



“C’e’ molto spazio nel fondo”

Luca Biferale

SIF 2015

Dept. Physics, INFN and CAST, University of Rome ‘Tor Vergata’

-A guided short tour around a mesoscopic/kinetic numerical tool:
Lattice Boltzman Equations

-Fluctuating Hydrodynamics
-Softy Glass Material

Credits:

M. Sbragaglia, R. Benzi, A. Gupta, A. Scagliarini (Dept. Physics, University of Rome “Tor Vergata”)

S. Succi & M. Bernaschi (IAC-CNR, Rome)

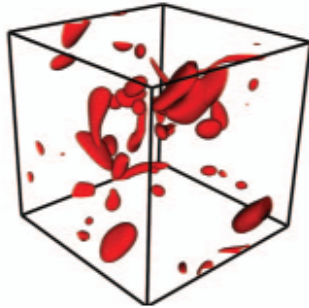
F. Toschi (University of Eindhoven & IAC-CNR, Rome)

M. Pierno, S. Varagnolo, G.P. Mistura (Dept. Physics, University of Padova)

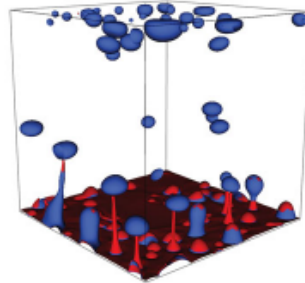


LBM for Complex Fluid Dynamics: Why & Where ?

$10^{-3} - 10^0 m$



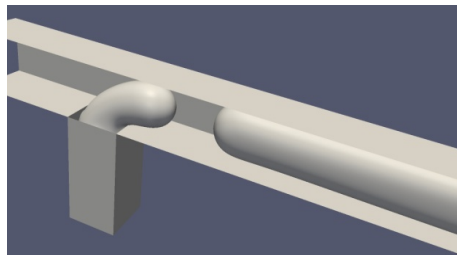
Droplets/Particles in Turbulence



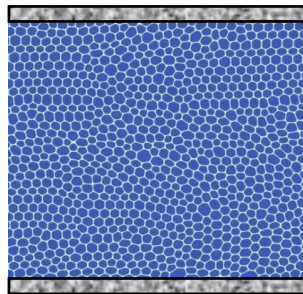
Boiling Fluids

- ✓ How the Heat Transport in a Rayleigh-Bènard Cell is affected by the presence of bubbles ?
- ✓ How Droplets interact with small scales Turbulence ?

$10^{-6} - 10^{-3} m$



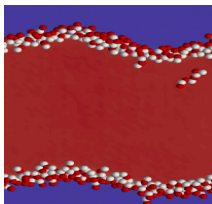
Droplets in Microfluidics



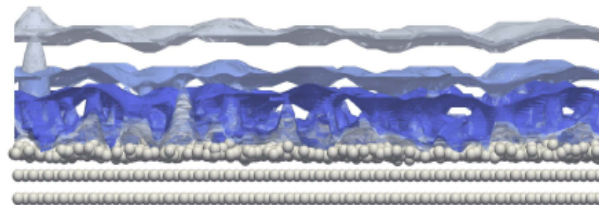
Soft-Glassy Materials

- ✓ Dynamics of Single Droplets: bounded vs unbounded geometries ?
- ✓ How The Rheology of a Collection of Droplets is affected by confinement?

$10^{-9} - 10^{-6} m$



Membranes



Nanofluids

- ✓ How to correctly control surfactant dynamics at the interfaces?
- ✓ How the slip at the liquid-solid interfaces is affected by roughness and wettability ?

Where to Sit Across the scales?



**MACROSCOPIC
(NAVIER-STOKES)**

$$\begin{aligned}\vec{\nabla} \cdot \vec{u} &= 0 \\ \rho D_t \vec{u} &= -\vec{\nabla} p + \eta \Delta \vec{u} \\ &\dots\end{aligned}$$

Continuum Description



Chapman-Enskog
(Multi-Scale Lattice Theory)



**MESOSCOPIC
(LATTICE BOLTZMANN)**

$$\partial_t f(\vec{x}, \vec{v}) + \vec{v} \cdot \vec{\partial}_x f(\vec{v}, \vec{x}) = -\frac{1}{\tau} (f(\vec{x}, \vec{v}) - f^{eq}(\vec{u}, \rho))$$

Kinetic '*Inspired*' Equations



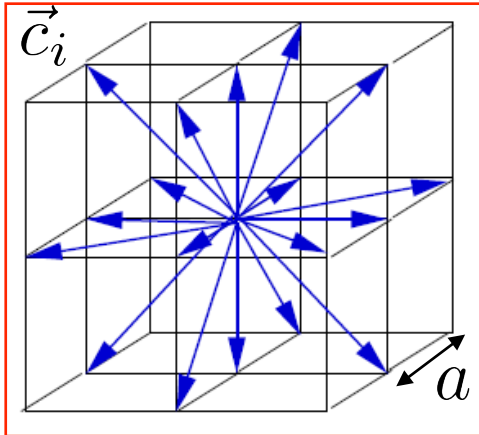
Supra-molecular
Framework for lattice
interactions



**MICROSCOPIC
(MOLECULAR DYNAMICS)**

$$\begin{aligned}\mathcal{H} &= \sum_k \frac{p_k^2}{2m} + \sum_{i,k} V_{i,k} \\ \dot{q}_k &= \frac{\partial \mathcal{H}}{\partial p_k} \quad \dot{p}_k = -\frac{\partial \mathcal{H}}{\partial q_k}\end{aligned}$$

Lattice Boltzmann: Basics & Conservation Laws



$$\partial_t f(\vec{x}, \vec{v}) + \vec{v} \cdot \vec{\partial}_x f(\vec{v}, \vec{x}) = -\frac{1}{\tau} (f(\vec{x}, \vec{v}) - f^{eq}(\vec{u}, \rho))$$

$$\vec{v} \rightarrow \vec{c}_i$$

$$f(\vec{v}, \vec{x}) \rightarrow n_i(\vec{x})$$

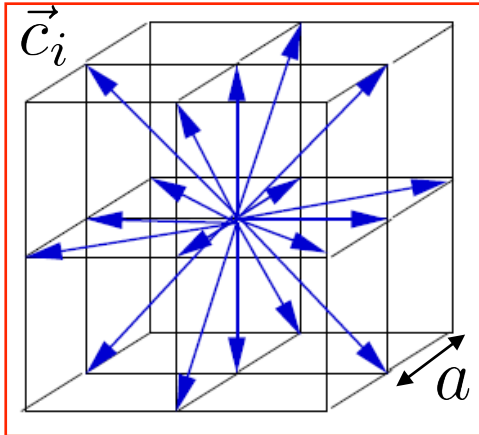


$$n_i(\vec{r} + \vec{c}_i, t + 1) = n_i^* = n_i(\vec{r}, t) + \Delta_i \{n_i(\vec{r}, t)\}$$

- \vec{c}_i **Small Set** of Velocities
- \vec{c}_i Connects 2 sites...**exact streaming** preserved
- $n_i(\vec{r}, t)$ Real Number: particle pdf on space-time location (\vec{r}, t)

How to access coarse grained variables?

Lattice Boltzmann: Basics & Conservation Laws



- Linearized Boltzmann Equation
- Fully Discretized (time and space)
- Sites \vec{r} , and Lattice Spacing a
- Time t , time step $\Delta t = 1$

$$n_i(\vec{r} + \vec{c}_i, t + 1) = n_i^* = n_i(\vec{r}, t) + \Delta_i \{n_i(\vec{r}, t)\}$$

Coarse Grained Fields:

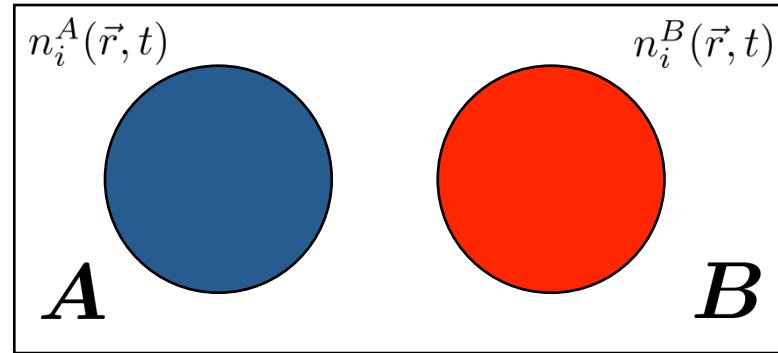
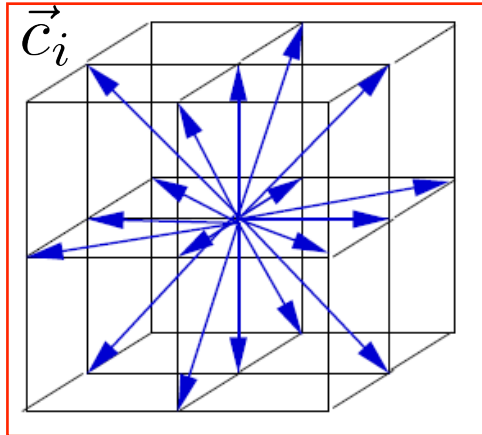
$$\rho(\vec{r}, t) = \sum_i n_i(\vec{r}, t) \quad (\text{Density})$$

$$\vec{j}(\vec{r}, t) = \sum_i \vec{c}_i n_i(\vec{r}, t) \quad (\text{Momentum})$$

$$\sum_i \Delta_i = \sum_i \vec{c}_i \Delta_i = 0$$

Mass Conservation!
Momentum Conservation!!
Locality !!!

Increasing Complexity: Multicomponent LBM



$$n_i^\sigma(\vec{r} + \vec{c}_i, t + 1) = n_i^{\sigma,*} = n_i^\sigma(\vec{r}, t) + \Delta_i^\sigma$$

$n_i^\sigma(\vec{r}, t)$ Real Number: particle pdf on space-time location (\vec{r}, t) component σ

Coarse Grained Fields:

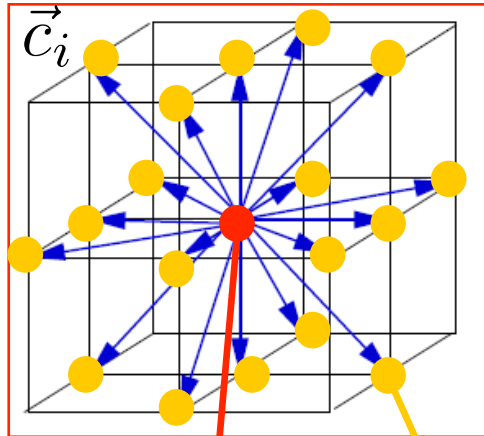
$$\rho_\sigma(\vec{r}, t) = \sum_i n_i^\sigma(\vec{r}, t) \quad (\text{Density})$$

Mass Conservation!
Global Momentum Conservation
Locality !!!

$$\vec{j}_\sigma(\vec{r}, t) = \sum_i \vec{c}_i n_i^\sigma(\vec{r}, t) \quad (\text{Momentum})$$

$$\sum_i \Delta_i^\sigma = 0 \quad \sum_i \vec{c}_i \Delta_i^\sigma \neq 0 \quad \sum_\sigma \sum_i \vec{c}_i \Delta_i^\sigma = 0$$

Non Ideal Lattice Fluids: Basic Ideas



$$\sigma = A, B$$

$$n_i^\sigma(\vec{r} + \vec{c}_i, t + 1) = n_i^\sigma(\vec{r}, t) + \Delta_i^\sigma \{ n_i^\sigma(\vec{r}, t) \}$$

$$\Delta_i^\sigma = \sum_j L_{ij}^\sigma n_i^{\sigma, neq} + \Delta_i^{F, \sigma}$$

Relax to Equilibrium + Forcing

Lattice Mean-Field Interactions (Shan-Chen)

$$\rho_A(\vec{r}) \sum_i w_i \rho_B(\vec{r} + \vec{c}_i) \vec{c}_i + (A \leftrightarrow B) \quad \xrightarrow{\times g_{AB}} \quad \vec{F} = -g_{AB} \rho_A(\vec{r}) \sum_i w_i \rho_B(\vec{r} + \vec{c}_i) \vec{c}_i + (A \leftrightarrow B)$$

Coupling Strength Parameter regulates the Force intensity

A Few Reference Works:

- X. Shan & H. Chen, *Physical Review E* **47**, 1815 (1993)
- X. Shan & H. Chen, *Physical Review E* **49**, 2941 (1994)
- X. Shan, *Physical Review E* **77**, 066702 (2008)
- M. Sbragaglia & X. Shan, *Physical Review E* **84**, 036703 (2011)
- R. Benzi, *J. Chem. Phys.* **131**, 104903 (2009)

How to rationalize the outcome of numerical simulations ?

Equilibrium Properties: Bulk Densities

$$\mathcal{L}(\rho_A, \rho_B, \nabla \rho_A, \nabla \rho_B) = f_b(\rho_A, \rho_B) - \frac{g_{AB}}{2} \nabla \rho_A \cdot \nabla \rho_B - \lambda_A \rho_A - \lambda_B \rho_B$$

Lagrange Multipliers

$$f_b(\rho_A, \rho_B) = c_s^2 \rho_A \log \rho_A + c_s^2 \rho_B \log \rho_B + g_{AB} \rho_A \rho_B$$

Bulk Free Energy

$$f_b(\phi, \rho) = \frac{c_s^2}{2} (\rho + \phi) \log \left(\frac{\rho + \phi}{2} \right) + \frac{c_s^2}{2} (\rho - \phi) \log \left(\frac{\rho - \phi}{2} \right) + \frac{1}{4} g_{AB} (\rho^2 - \phi^2)$$

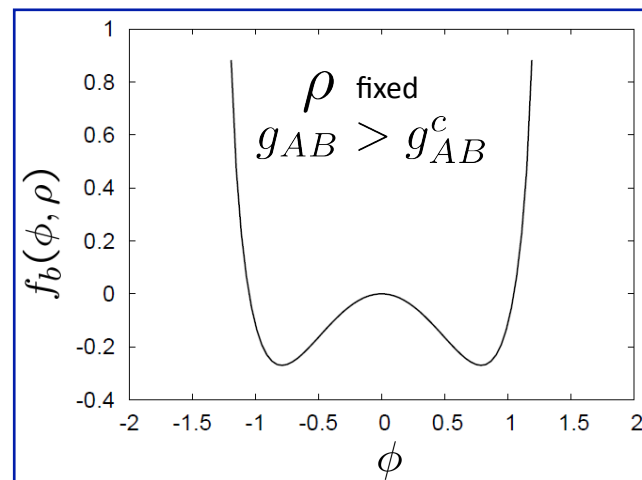
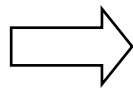
Double well structure

$$\rho = \rho_A + \rho_B$$

Total Density

$$\phi = \rho_A - \rho_B$$

Order Parameter

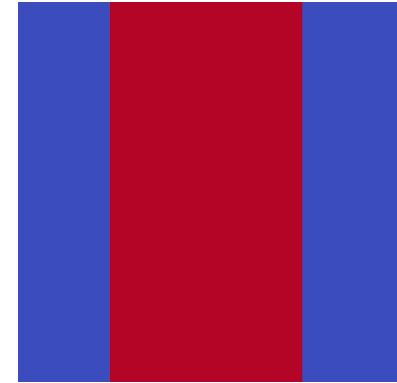


R. Benzi, *J. Chem. Phys.* **131**, 104903 (2009)

Fluctuating Multicomponent Lattice Boltzmann

Fluctuating Multicomponent Lattice Boltzmann

$$n_i^\sigma(\vec{r} + \vec{c}_i, t + 1) = n_i^\sigma(\vec{r}, t) + \Delta_i^\sigma + \Delta_{i,F}^\sigma + \hat{\Delta}_i^\sigma$$



$$\rho \left(\frac{\partial}{\partial t} \mathbf{u} + (\mathbf{u} \cdot \nabla) \mathbf{u} \right) = -\nabla p + \nabla \cdot (\boldsymbol{\sigma} + \hat{\boldsymbol{\sigma}}) + \sum_{\sigma} \mathbf{g}_{\sigma},$$

$$\frac{\partial}{\partial t} \rho_{\sigma} + \nabla \cdot (\rho_{\sigma} \mathbf{u}) = \nabla \cdot (\mathbf{d}_{\sigma} + \hat{\mathbf{d}}_{\sigma})$$

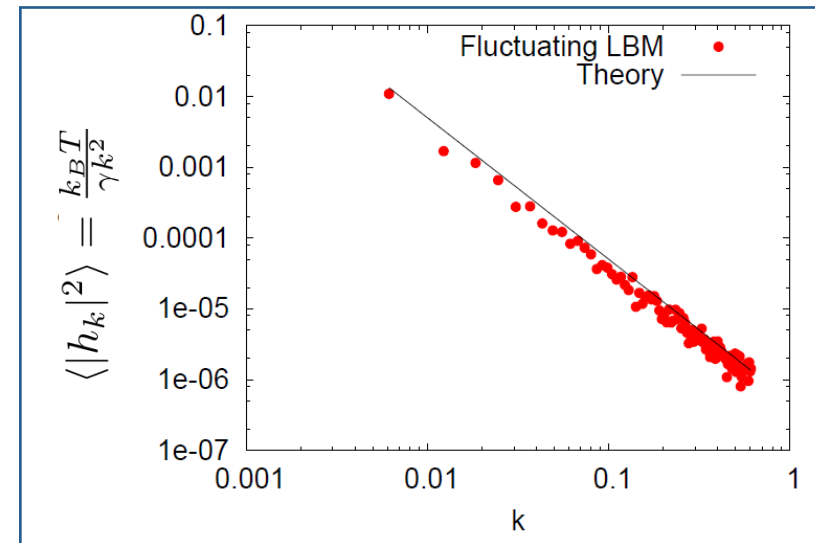
Fluctuating (Stress/Flux)

$$E = \gamma \int dx dy \left[\sqrt{1 + |\nabla h|^2} - 1 \right]$$

$$\approx \frac{\gamma}{2} \int dx dy |\nabla h|^2 = \frac{\gamma}{2} \sum_k k^2 |h_k|^2$$

Normal Modes Decomposition of Surface distortions

$$\langle |h_k|^2 \rangle = \frac{kT}{\gamma k^2}$$



Fluctuating Multicomponent Lattice Boltzmann

$$n_i^\sigma(\vec{r} + \vec{c}_i, t + 1) = n_i^\sigma(\vec{r}, t) + \Delta_i^\sigma + \Delta_{i,F}^\sigma + \hat{\Delta}_i^\sigma$$

Perturbation around equilibrium state at rest

$$\rho^\sigma(\vec{r}, t) = \rho_0^\sigma(\vec{r}) + \delta\rho^\sigma(\vec{r}, t)$$

$$\vec{v}(\vec{r}, t) = \vec{0} + \delta\vec{v}(\vec{r}, t)$$

How to Properly Formulate the Noise terms?

Mode decomposition of the fluctuations in the number density

$$\delta n_i^\sigma(\vec{r}, t) = n_i^\sigma(\vec{r}, t) - n_i^{\sigma,eq}(\rho_0^\sigma(\vec{r}), \vec{0}) \longrightarrow \delta m_a^\sigma(\vec{r}, t) = \sum_i e_{ai} \delta n_i^\sigma(\vec{r}, t)$$

$$\partial_t \delta \hat{m}_a^\sigma(\vec{k}, t) = \sum_{b,\zeta} \int d\vec{k}' \mathcal{L}_{ab}^{\sigma\zeta} \delta \hat{m}_b^\zeta(\vec{k}', t) + \hat{\zeta}_a^\sigma(\vec{k}, t)$$

Belardinelli et al., *Phys. Rev. E* **91**, 023313 (2015)

$$\partial_t a_i(t) = - \sum_j L_{ij} a_j(t) + \zeta_i(t) \quad \text{Linear Langevin Equation !!!}$$

$$G_{ij} = \langle a_i a_j^* \rangle$$

Equilibrium Correlations (Entropy Matrix)

Input for the Theory

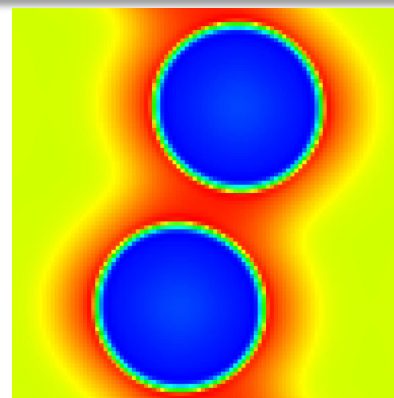
$$\langle \zeta_i(t) \zeta_j^*(t) \rangle = \sum_k (G_{ik} L_k^* + L_{ik} G_{kj})$$

Fixed by FDT (Fluctuation Dissipation Theorem)

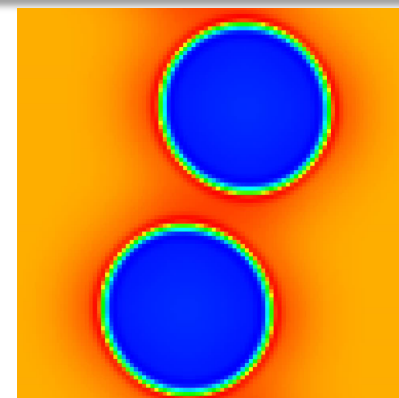
Modelling Soft-Glassy Materials

Frustration Mechanism at the interface

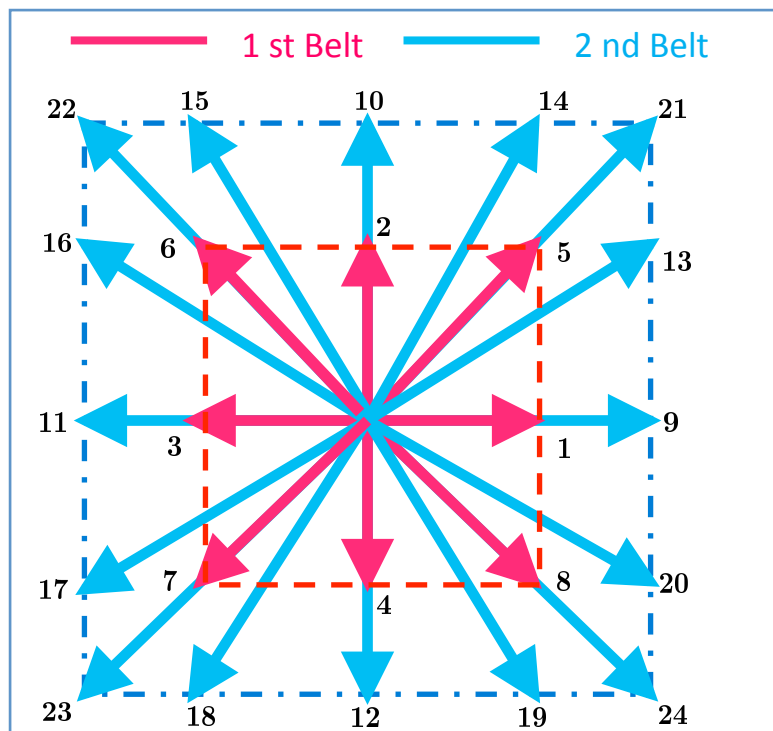
- ✓ “Simple” Lattice (competing) Interactions
- ✓ A variety of Rich Effects
- ✓ Control over these effects ?



Droplets Coalescence



Inhibition of Coalescence



Idea: Multirange Potentials NN & NNN

$$\vec{F}_i(\vec{r}) = -\mathcal{G}_1\psi(\vec{r}) \sum_{i \in NN} \psi(\vec{r} + \vec{c}_i)\vec{c}_i - \mathcal{G}_2\psi(\vec{r}) \sum_{i \in NNN} \psi(\vec{r} + \vec{c}_i)\vec{c}_i$$

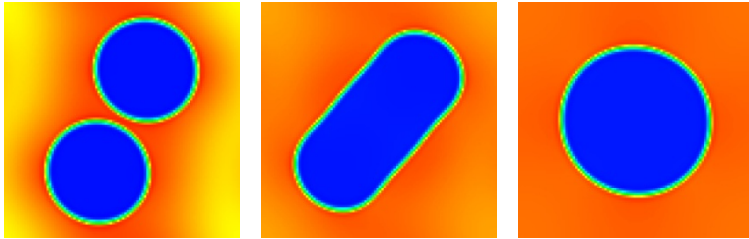
$$\psi(\vec{r}) = \psi[\rho(\vec{r})] \text{ Pseudo-Potential}$$

A Few Reference Works:

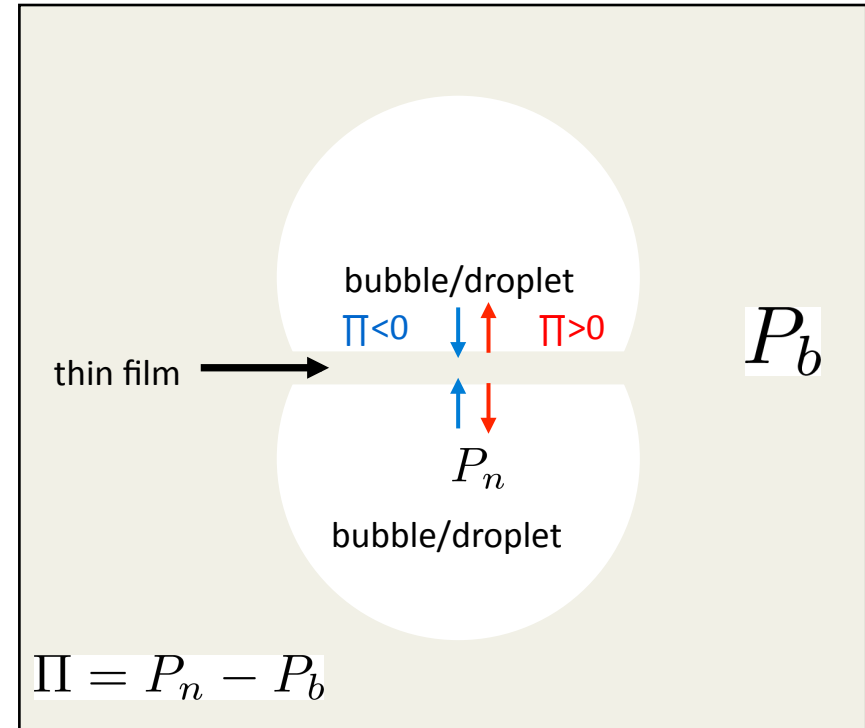
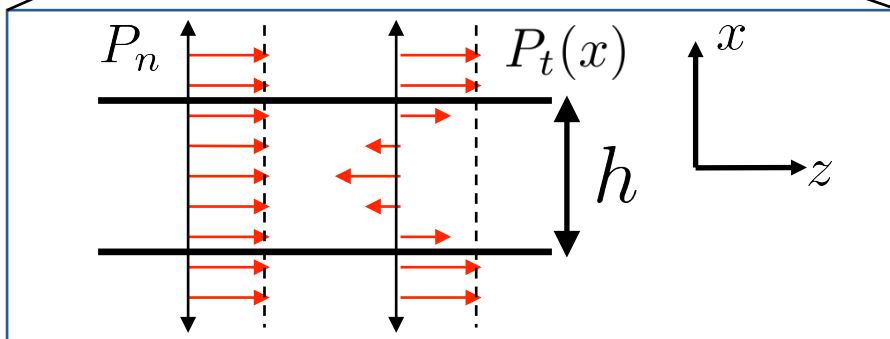
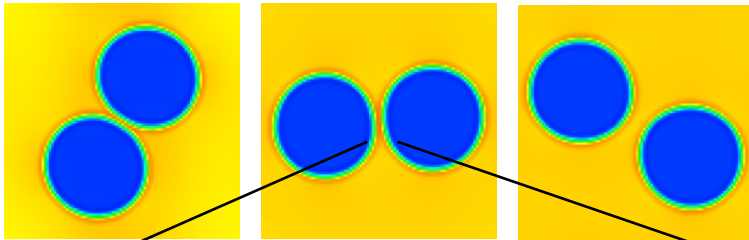
- R. Benzi et al, *J. Chem. Phys.* 131, 104903 (2009)
- R. Benzi et al, *Europhys. Lett.* 91, 14003 (2010)
- M. Sbragaglia et al, *Soft Matter* 8, 10773-10782 (2012)

Modelling of Disjoining Pressure

Droplets Coalescence



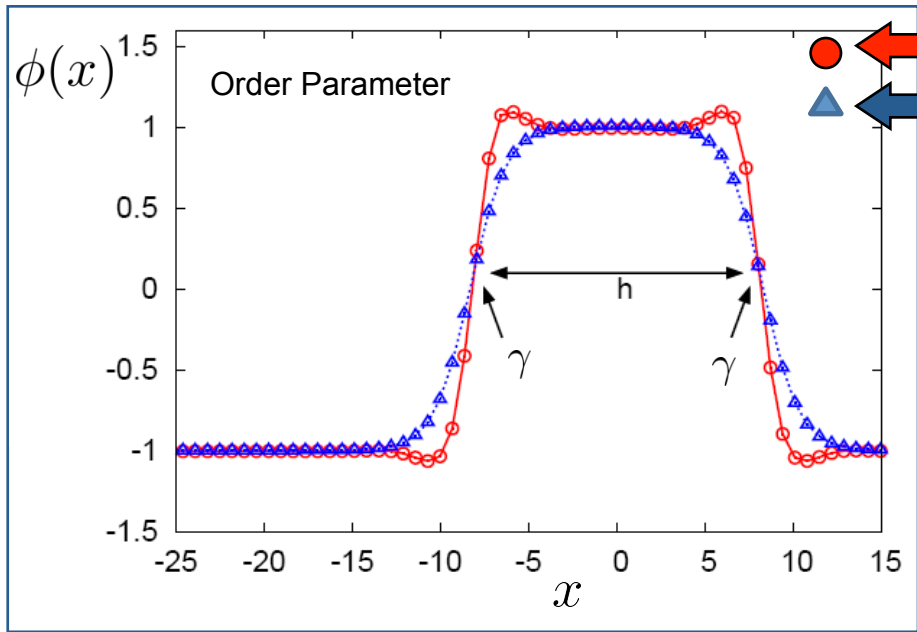
No Coalescence...why?



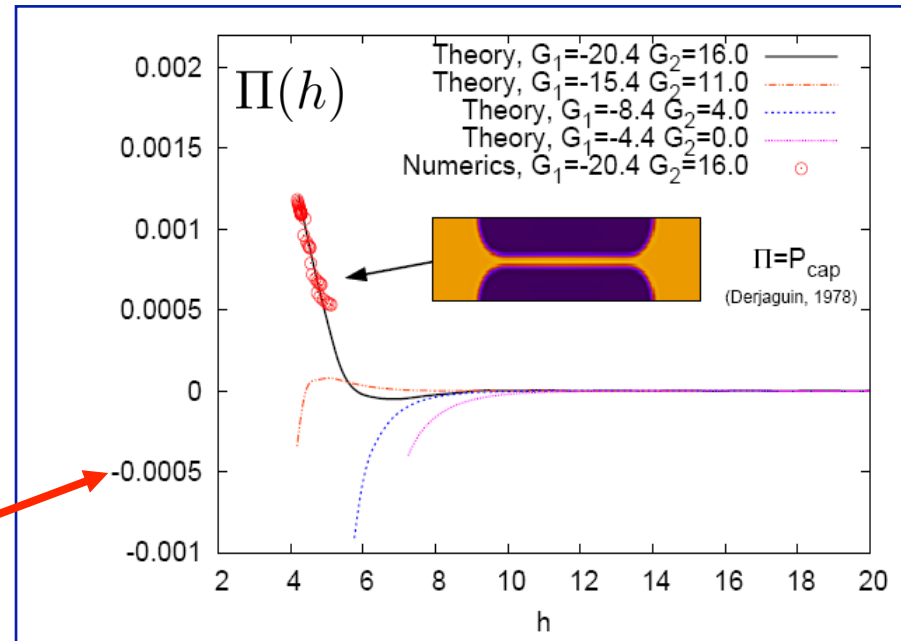
.....Because you have a Positive disjoining pressure

Derjaguin & Churaev, J. Colloid. Interface Sci. 66, 389 (1978)

Disjoining Pressure From Lattice Kinetic models



● ← (Competing NN & NNN) $\mathcal{G}_1 < 0; \mathcal{G}_2 > 0$
▲ ← (bare NN) $\mathcal{G}_1 < 0; \mathcal{G}_2 = 0$



$$\gamma_f = 2\gamma + \int_{\Pi(h=\infty)}^{\Pi(h)} h d\Pi$$

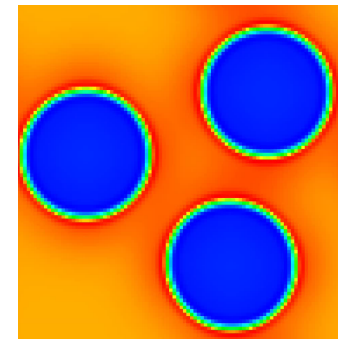
Disjoining Pressure definition based on 'film tension'



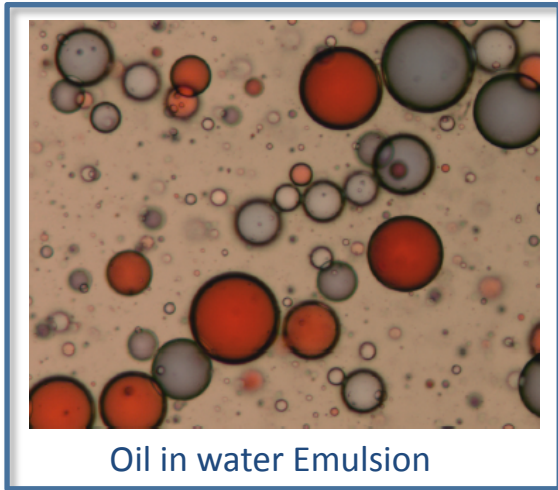
Exactly computed on the lattice

$$P_n = P_b(x) + C_1 \psi \frac{d^2 \psi}{dx^2} + \dots = \text{const}$$

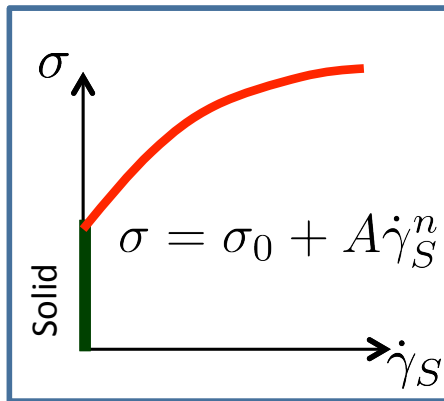
Analytical with the exact pressure tensor
(differential equation \rightarrow Profile)



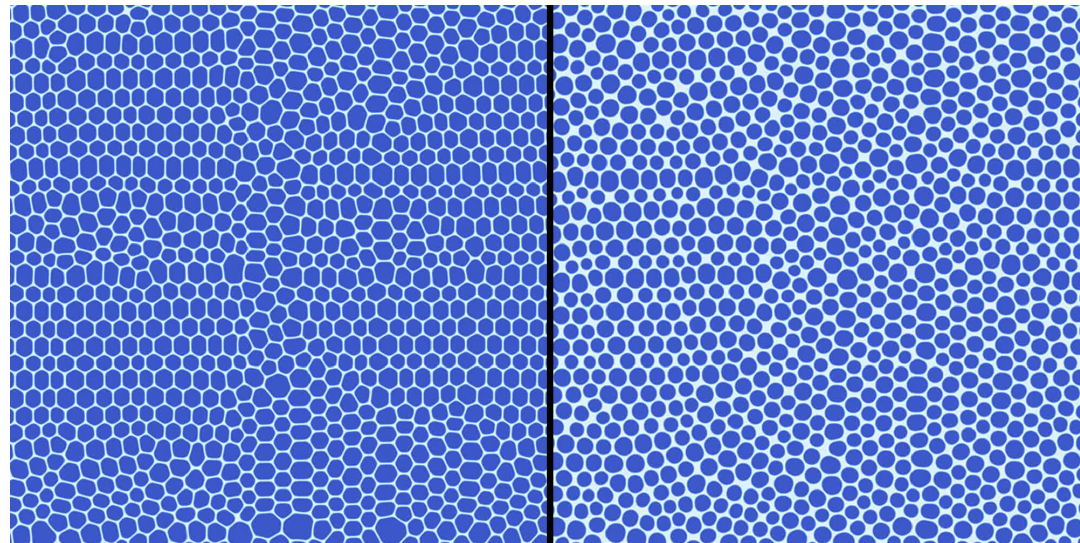
Rheology of Soft Glassy “Materials”



- ✓ Yield Stress (Solid below)
- ✓ Non-Newtonian (above yield Stress)
- ✓ Effect of Confinement (Cooperativity Effects)



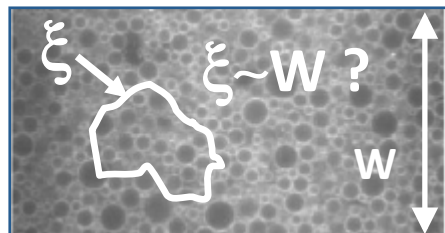
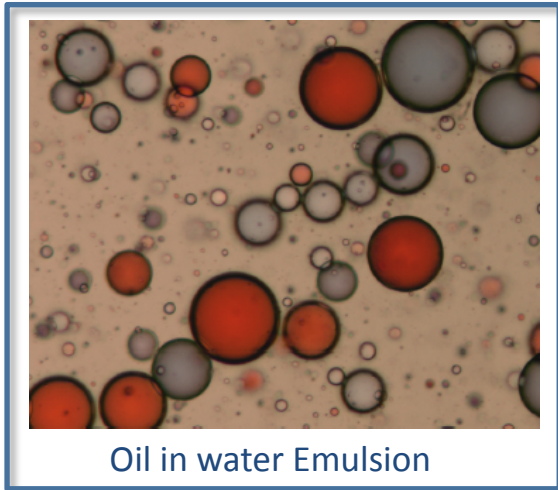
Herschel-Bulkley Model



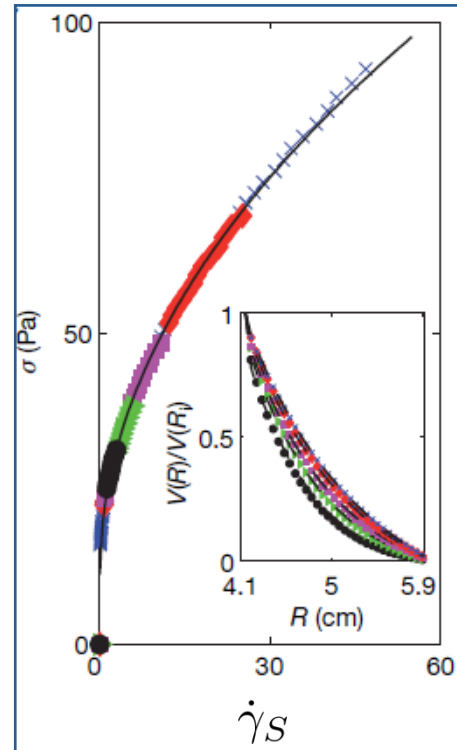
High Packing Fraction

Lower Packing Fraction

Rheology of Soft Glassy “Materials”

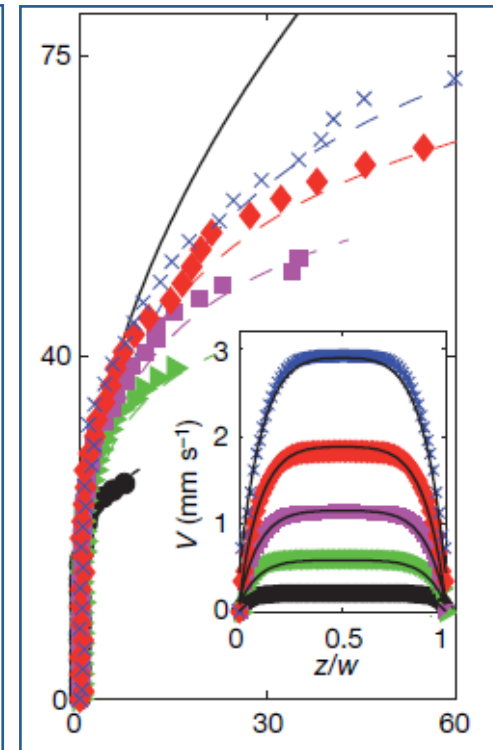


- ✓ Yield Stress (Solid below)
- ✓ Non-Newtonian (above yield Stress)
- ✓ Effect of Confinement (Cooperativity Effects)



$$W/\xi \gg 1$$

cooperativity does not matter

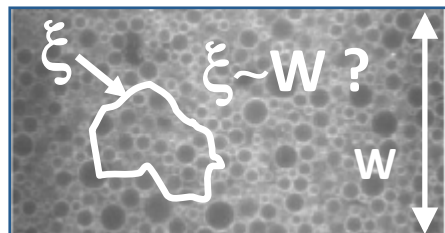
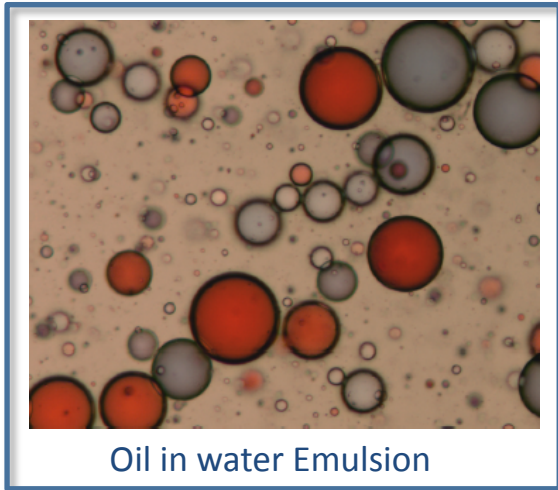


$$W/\xi \sim 1$$

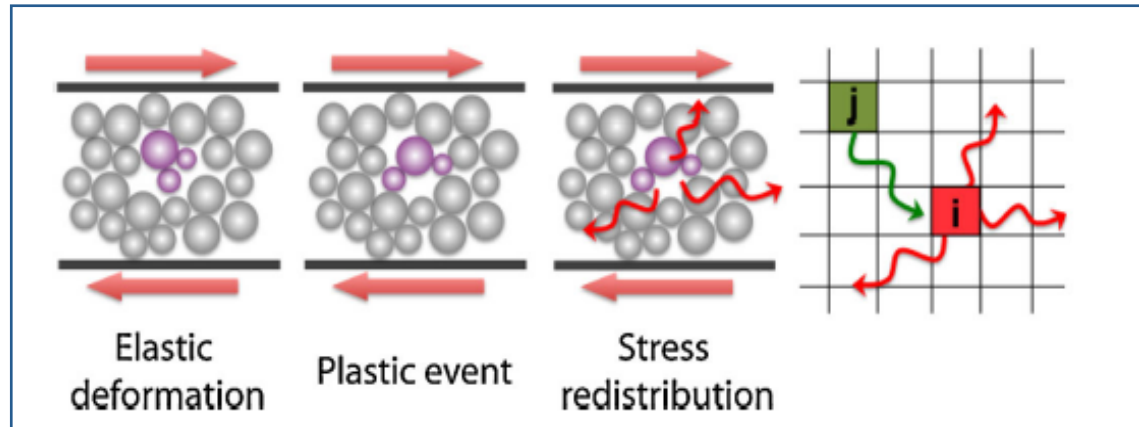
cooperativity matters

Goyon et al, Nature (2008)

Rheology of Soft Glassy "Materials"



- ✓ Yield Stress (Solid below)
- ✓ Non-Newtonian (above yield Stress)
- ✓ Effect of Confinement (Cooperativity Effects)



$$f = \frac{\dot{\gamma}_S}{\sigma}$$

Fluidity

$$\xi^2 \Delta f + (f_b - f) = 0$$

Diffusion equation for the Fluidity

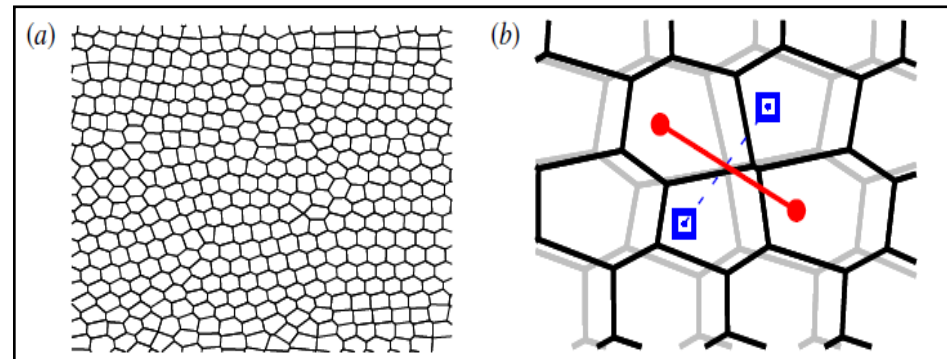
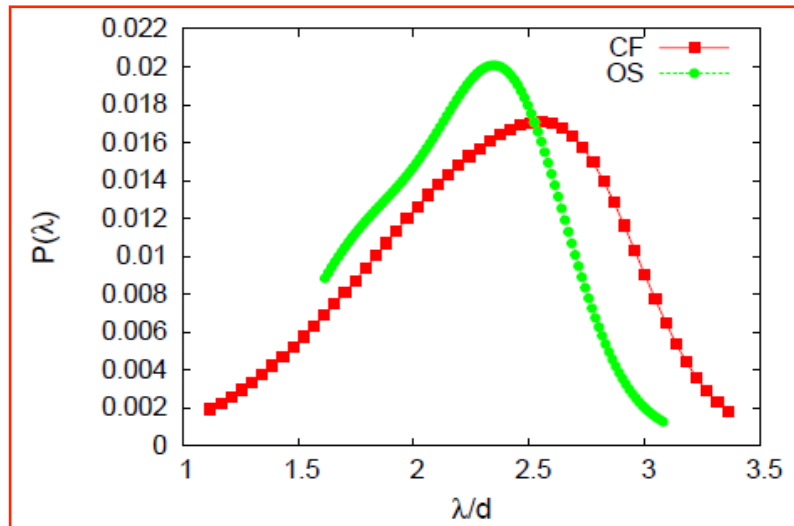
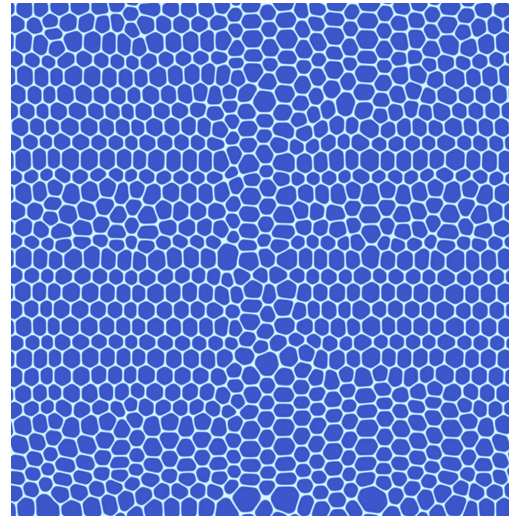


$$f_b = f_b(\sigma(\mathbf{x}))$$

"Bulk" Fluidity

Plastic Events: LBM with Different Load Conditions

Oscillating Strain (OS) Boundary Conditions



Dollet, Scagliarini & Sbragaglia, *Jour. Fluid. Mech* **776**, 556 (2015)

Probability Distribution of size of Plastic Rearrangement
(Different Flow Conditions: CF (Couette Flow); OS (Oscillatory Strain))

Conclusions

✓ **Lattice Boltzmann:**

✓ Direct Link with Hydrodynamics (Chapman Enskog)

✓ Easy Handling of Boundary Conditions

✓ Flexibility in modelling Non-Ideal interactions (Shan-Chen, Free-Energy, Lee-Fisher, Color Models, approaches)

✓ **Solid Background** (at least Single-Phase with/without Thermal Fluctuations ...see Shan et al. *Jour. Fluid. Mech.* (2006);Dunweg, Shiller, Ladd, *Phys. Rev. E* (2007))

✓ **Useful in Revealing Interesting Physics of Complex Fluid Dynamics problem...and Beyond:**

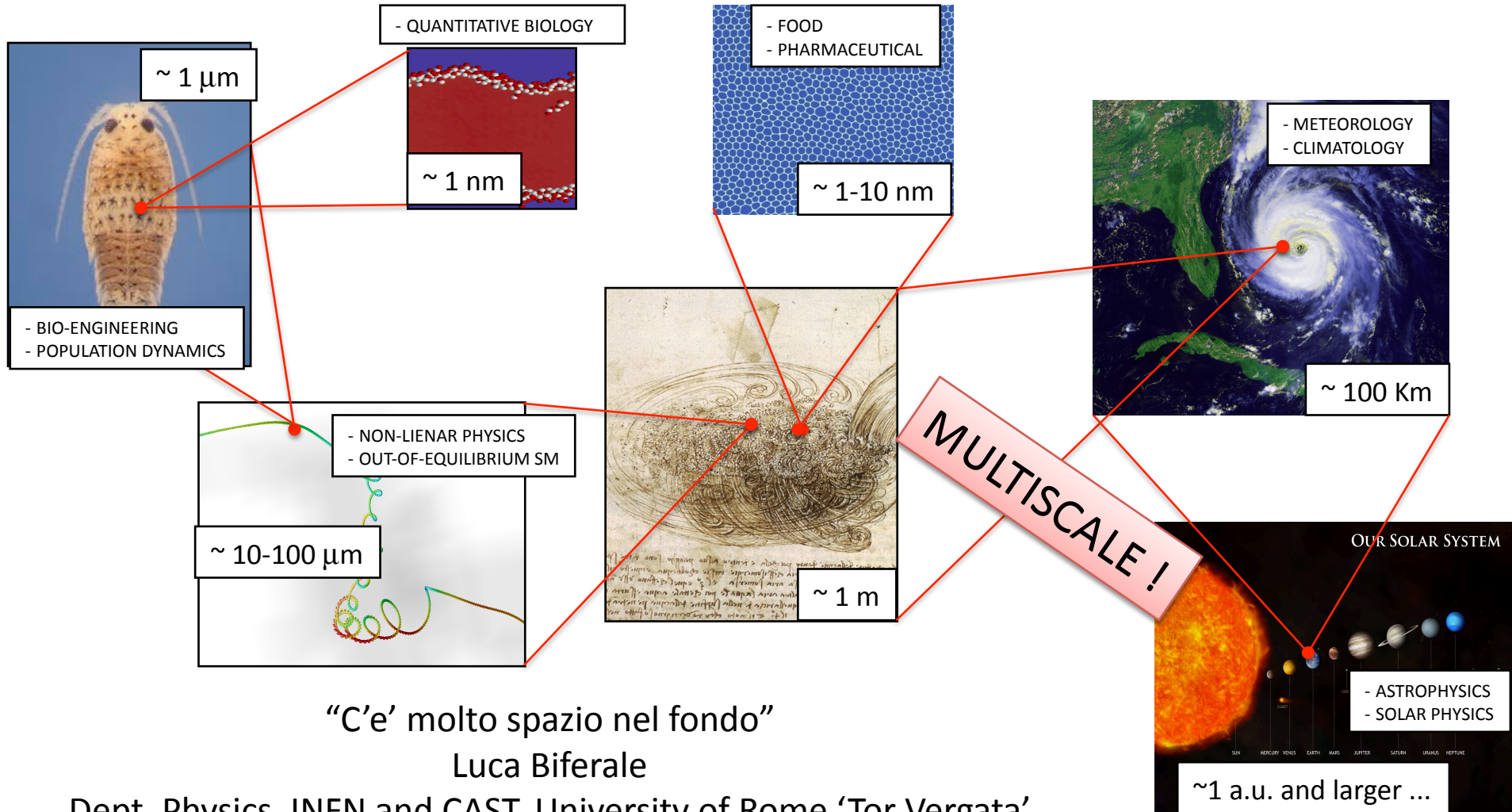
✓ **Microfluidics Droplets** (Surface tension, wettability Gradients, strong Confinement, Boundary Conditions, etc etc)

✓ **Soft-Glassy Rheology** (Onset of Motion, Fluidity, Spatial Stress Correlations, etc etc)

✓ **Coupling with Atomistic Scales** (Fluctuating Hydrodynamics)

Πάντα ρει

inside, outside, above, below and around us



“C’e’ molto spazio nel fondo”

Luca Biferale

Dept. Physics, INFN and CAST, University of Rome ‘Tor Vergata’

A simple (textbook) example on how to set Noise

Motion of particle with mass m in a fluid: **total Force** due to the **other particles** in the fluid

$$m \frac{dv(t)}{dt} = F_{tot}(t) \approx -\zeta v(t) \quad \text{Dominated by Friction: Is this the whole Story?}$$

$$v(t) = e^{-\zeta t/m} v(0) \rightarrow 0 \quad (t \rightarrow \infty)$$

$$\langle v^2 \rangle_{eq} \neq \frac{kT}{m} \quad \text{Mean Squared Velocity at equilibrium (t >> 0) is not reproduced !! **Modify above picture !!**}$$

Observed Randomness in the individual trajectory: **add Random Force**

$$m \frac{dv(t)}{dt} = F_{tot}(t) = -\zeta v(t) + \delta F(t) \quad \text{Langevin Equation for Brownian Particle}$$

$$\langle \delta F(t) \rangle = 0 \quad \langle \delta F(t) \delta F(t') \rangle = 2B \delta(t - t') \quad \text{More commonly presented view}$$

Both **Friction** and **Random Force** come from the bath: **Any fundamental relation?**

$$\langle v^2(t) \rangle = e^{-2\zeta t/m} v^2(0) + \frac{B}{\zeta m} (1 - e^{-2\zeta t/m})$$

$$\langle v^2(t) \rangle \rightarrow \langle v^2(t) \rangle_{eq} = \frac{kT}{m} \quad (t \rightarrow \infty)$$

$$B = \zeta kT$$

Fluctuation Dissipation Relation

Where to Sit Across the scales?



**MACROSCOPIC
(NAVIER-STOKES)**

$$\begin{aligned}\vec{\nabla} \cdot \vec{u} &= 0 \\ \rho D_t \vec{u} &= -\vec{\nabla} p + \eta \Delta \vec{u} \\ &\dots\end{aligned}$$

✓ Can we explore “Complex Hydrodynamics” with
Capillarity Phenomena via a **Coarse grained** description?

YES !!! A possibility is the **Lattice Boltzmann Method**



**MICROSCOPIC
(MOLECULAR DYNAMICS)**

$$\begin{aligned}\mathcal{H} &= \sum_k \frac{p_k^2}{2m} + \sum_{i,k} V_{i,k} \\ \dot{q}_k &= \frac{\partial \mathcal{H}}{\partial p_k} \quad \dot{p}_k = -\frac{\partial \mathcal{H}}{\partial q_k}\end{aligned}$$