



Società Italiana di Fisica
Italian Physical Society

SOCIETÀ ITALIANA DI FISICA 101° CONGRESSO NAZIONALE ROMA, 21-25 SETTEMBRE, 2015



INGV

SEZIONE 4: Geofisica, fisica dell'ambiente

16/09/2015

CORRECTION'S METHOD OF THE ELECTRON DENSITY MODEL IN IONOSPHERE BY RAY TRACING TECHNIQUES

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ATTICON 8909

This abstract focuses on the lead role that is played by the variation of the electron density grid into the ray tracing algorithm, which is correlated to the change of the electron content along the ionospheric ray path, for obtaining a ray tracing as much reliable as possible.

In many cases of practical interest, the group delay time depends on the geometric length and the electron content of the ray path.

The issue is faced theoretically, and a simple analytical relation, between the variation of the electron content along the path and the difference in time between the group delays, calculated and measured, both in the ionosphere and in the vacuum, is obtained and discussed.

An example of how an oblique radio link can be improved by varying the electron density grid is also shown and discussed.

Keywords: *ionospheric ray tracing; electron density model; ray path correction; electron content.*

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1. Introduction

The group delay time t_c , calculated step by step along the ionospheric ray path, is particularly interesting for this talk.

The reason is that the calculated group delay time t_c is easily comparable with the effective measured group delay time t_m in some technological applications, such as the Over The Horizon Radar (OTHR), oblique synchronized sounding or ionospheric backscatter (Davies, 1990).

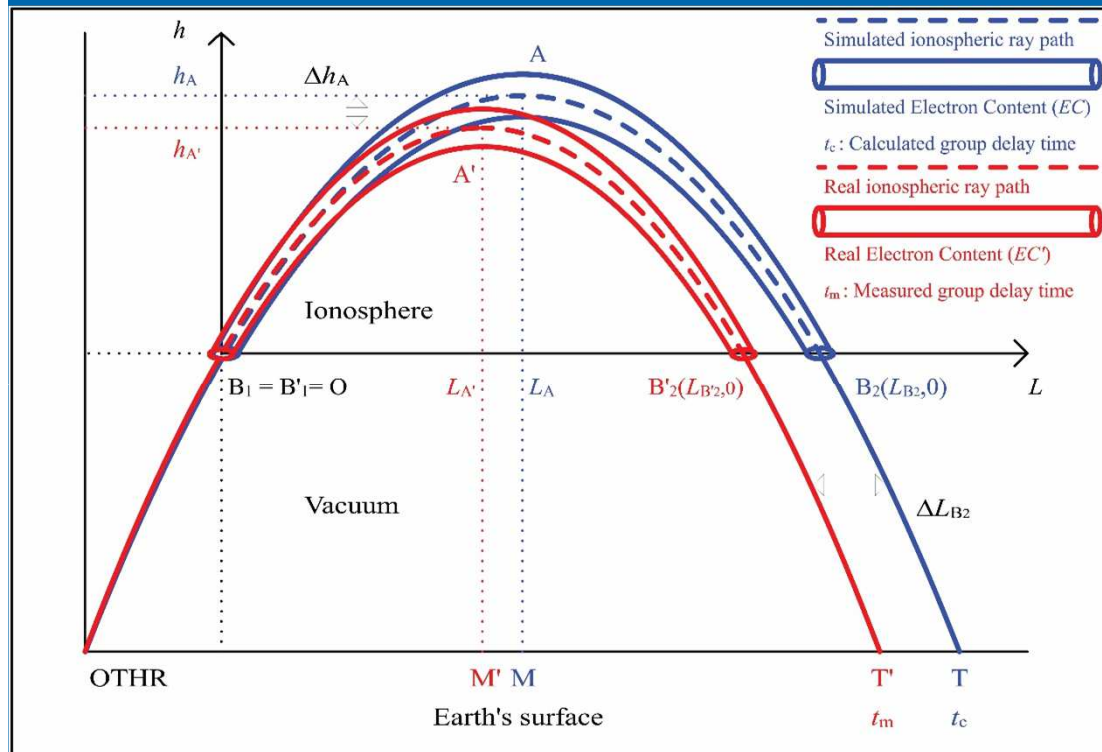


Figure 1 –

A little change of Electron Content (EC) can imply great refractive effects:

large variations for both the Apogee (A) location and the path length of ionospheric ray tracing.

The two ionospheric ray paths experience two different group delay times t_c and t_m , respectively the calculated and the measured group delay times.

2. Ray path and ionospheric models

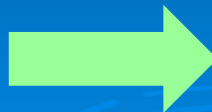
Ray tracing is a deterministic process of which the ionospheric ray path accuracy is arbitrarily chosen through the computational algorithm step (Haselgrove, 1955).

$$Z = 0$$



$$n^2 = 1 - \frac{X}{1 - \frac{1}{2} \frac{Y_T^2}{(1-X)} \pm \sqrt{\frac{Y_T^4}{4(1-X)^2} + Y_L^2}}$$

$$Y_L = Y_T = 0$$



$$n^2 = 1 - X = 1 - \frac{\omega_p^2}{\omega^2}$$

3. Group delay time calculation

For those radio waves propagating scarcely into the ionosphere, which is penetrated in correspondence to an oblique incidence, i.e. with low elevation angles, an approximate refractive index can hold throughout the HF band, i.e. 3-30 MHz.

$$\begin{aligned}
 t_c &= \int_{\text{OTHR}}^T \frac{dl}{v_g(l)} = \frac{1}{c} \int_{\text{OTHR}}^T n_g(l) dl = \\
 &= \frac{1}{c} \int_{\text{OTHR}}^T \frac{dl}{n(l)} = \frac{1}{c} \int_{\text{OTHR}}^T \frac{dl}{\sqrt{1 - \frac{\omega_p^2(l)}{\omega^2}}} \cong \\
 &\cong \frac{1}{c} \int_{\text{OTHR}}^T \left[1 + \frac{\omega_p^2(l)}{2\omega^2} \right] dl = \frac{l_0}{c} + \frac{e^2}{2cm\epsilon_0} \frac{1}{\omega^2} \int_{\text{OTHR}}^T N(l) dl
 \end{aligned}$$

$$t_c \cong t_{c0} + \frac{k}{f^2} EC$$

$$t_m \cong t_{m0} + \frac{k}{f^2} EC'$$



$$\begin{aligned}
 \Delta t = t_c - t_m &\cong \\
 &\cong \Delta t_0 + \frac{k}{f^2} \Delta EC = \\
 &= t_{c0} - t_{m0} + \frac{k}{f^2} (EC - EC')
 \end{aligned}$$

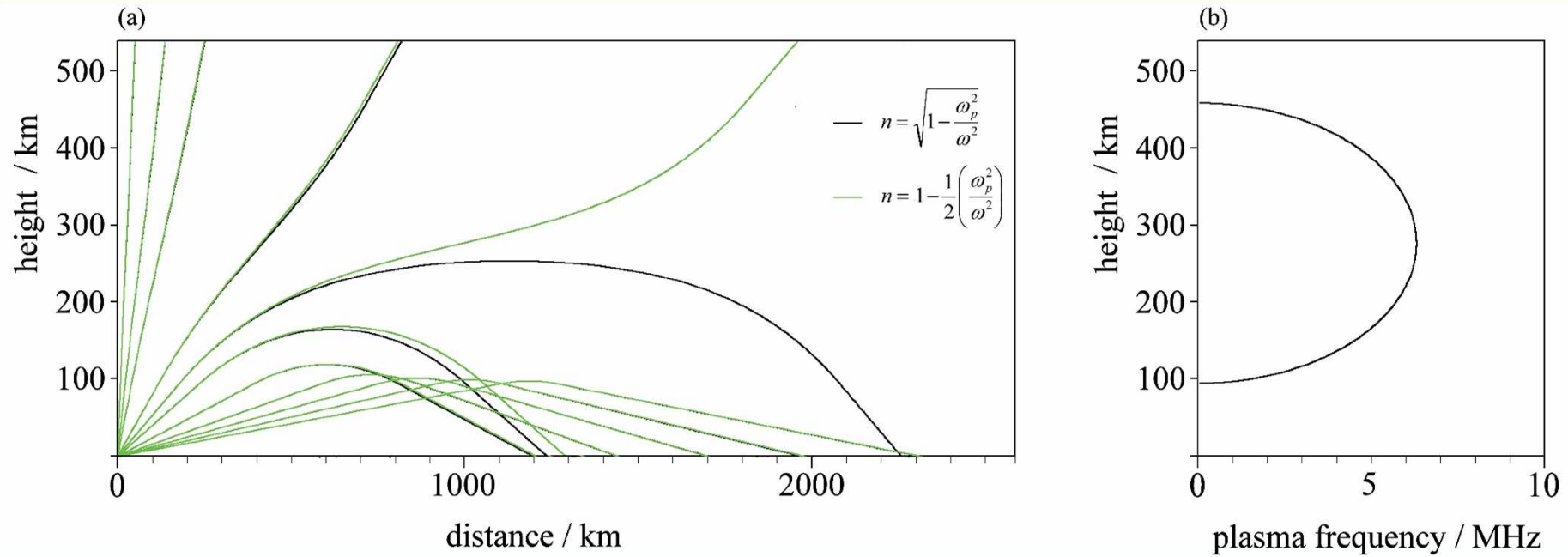


Figure 2 –

a) The ionospheric ray paths associated with a high frequency (HF) signal emitted from a transmitter (OTHR) at different elevation angles for a radio wave of frequency $f=13.4$ MHz.

The ray paths are computed with the exact and first-order Taylor's series expansion of the phase refractive index n .

b) The corresponding plasma frequency profile $f_p(h)$ of parabolic shape, with maximum $f_p^{(\max)}=6.784$ MHz at a height of 277 km, assuming $f_p(h)=0$ MHz under 95 km and above 459 km.

4. Correction of the electron density model

Such a condition is quite common when the frequency f is below the plasma frequency f_p and the radio wave is reflected beneath the electron density maximum (Davies, 1990):

$$\Delta EC \rightarrow 0 \Rightarrow \Delta t \approx \Delta t_0$$

Assuming that both the simulated and real ray paths can be modelled by parabolic curves:

$$\Delta L_{B_2} \ll L_{B_2}$$

$$\Delta h_A \ll h_A$$

$$\Delta L_{B_2} \cdot \Delta h_A \rightarrow 0$$

$$c\Delta t_0 \cong \frac{1}{8h_A^2} \left\{ 4L_{B_2} h_A \sqrt{1 + 16 \left(\frac{h_A}{L_{B_2}} \right)^2} \Delta h_A + \right. \\ \left. + L_{B_2} (2h_A \Delta L_{B_2} - L_{B_2} \Delta h_A) \ln \left[4 \frac{h_A}{L_{B_2}} + \sqrt{1 + 16 \left(\frac{h_A}{L_{B_2}} \right)^2} \right] \right\}$$

$$\Delta EC \approx \frac{f^2}{k} (\Delta t - \Delta t_0)$$

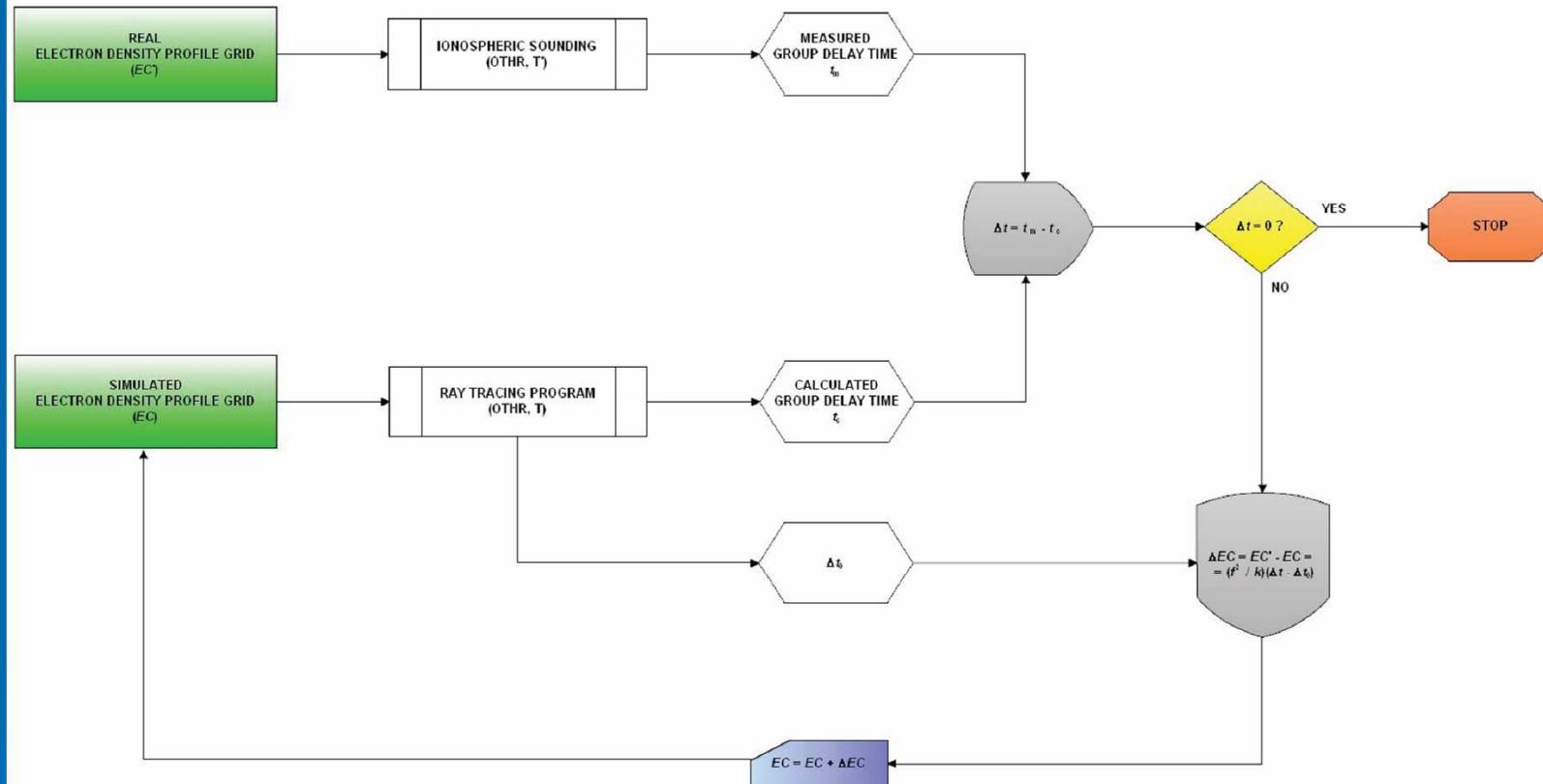


Figure 3 –

Logic flowchart of an algorithm for the correction's method of the electron density model in ionosphere by ray tracing techniques.

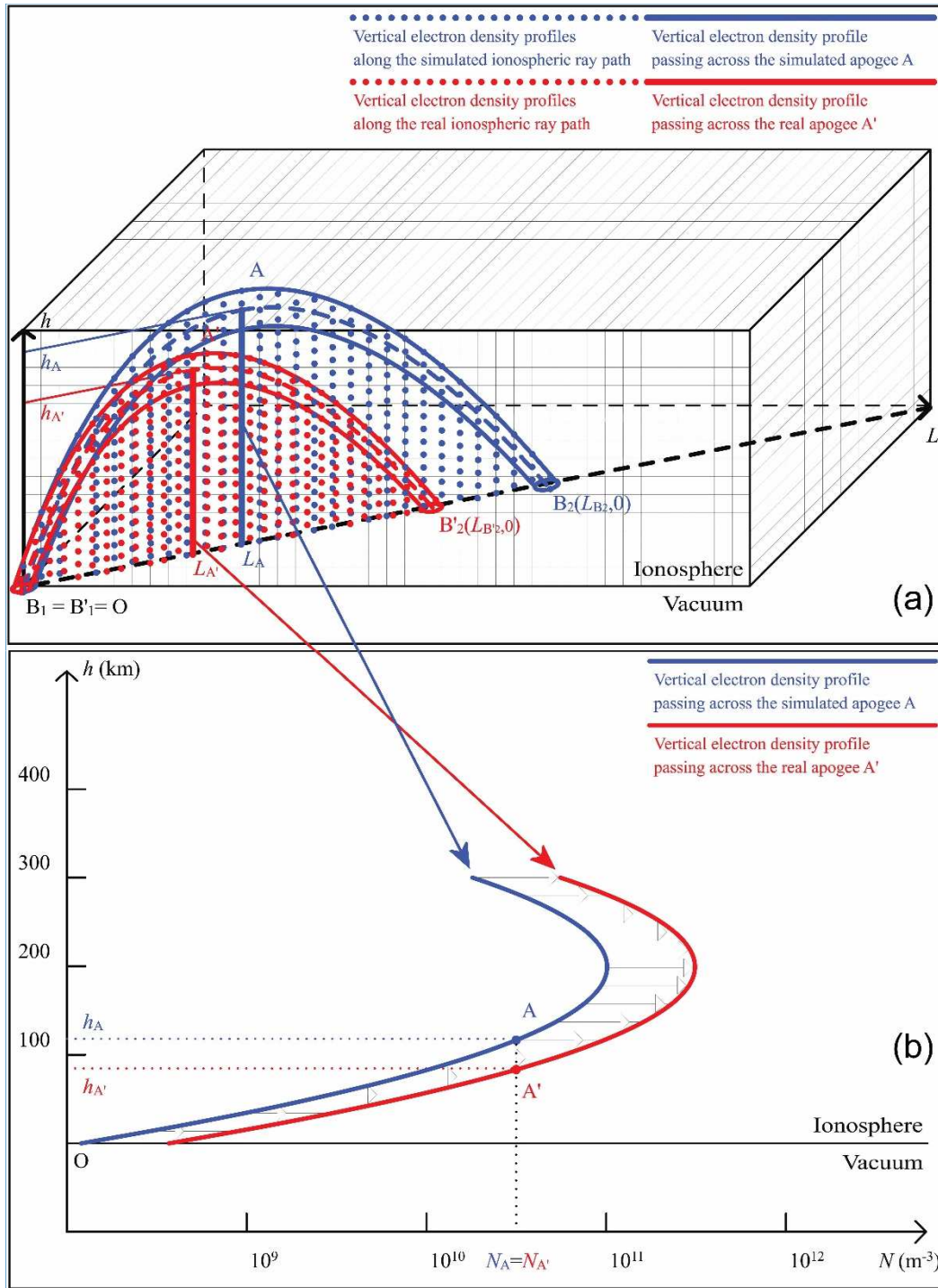


Figure 4 –

a) Vertical electron density profiles along the ionospheric ray path that have to be corrected compensating the time difference between the calculated and measured group delays in ionosphere.

b) The two vertical electron density profiles, plotted as a function of height, refer to the plumb lines passing across the simulated (A) or real (A') apogees. The two vertical profiles correspond to the simulated (N) or real (N') electron densities over the simulated (M) or real (M') midpoints between the transmitter (OTHR) and the simulated (T target) or real (T' target) receiver of Fig. 1.

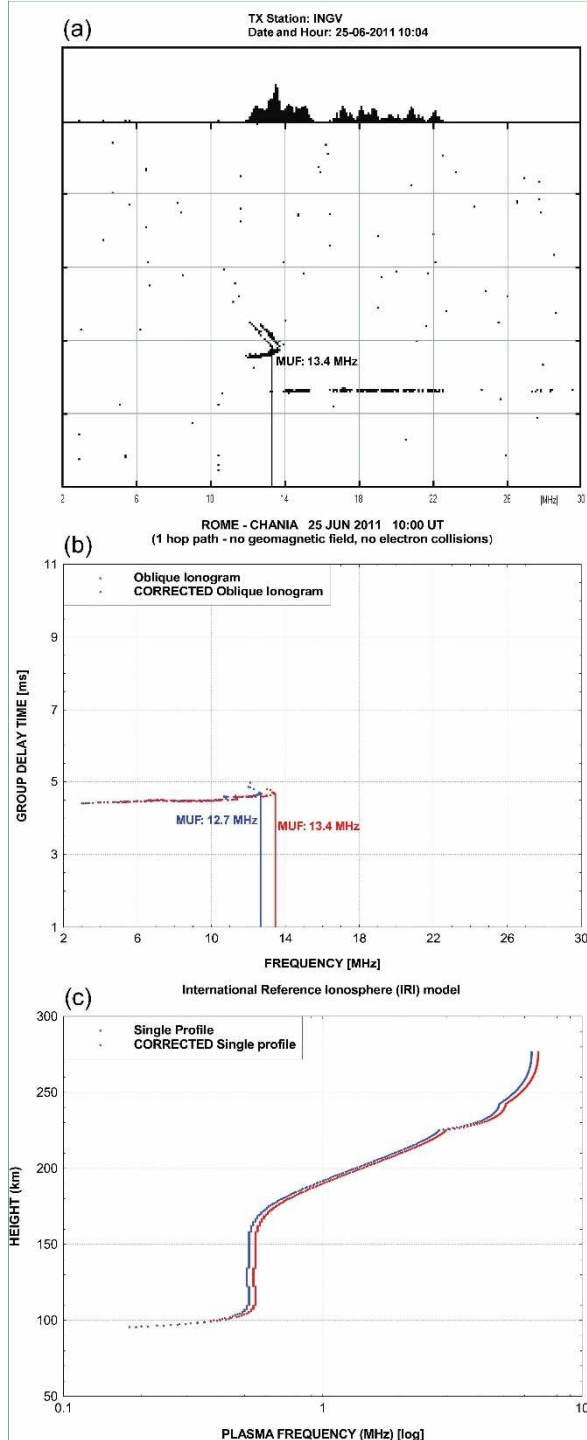


Figure 5 –

a) Oblique ionogram recorded on the 25 June 2011 at 10.00 UT along the Rome-Chania radio link, for which the maximum usable frequency MUF=13.4 MHz value is highlighted.

b) Oblique ionogram synthesized by the IONORT-IRI system showing two ordinary traces corresponding to a plasma frequency profile $f_p(h)$ (trace in blue colour), with a MUF=12.7 MHz (signed in blue colour), and to a corrected plasma frequency profile $f_p(h)+0.06 \cdot f_p(h)$ (trace in red), with a MUF=13.4 MHz (signed in red).

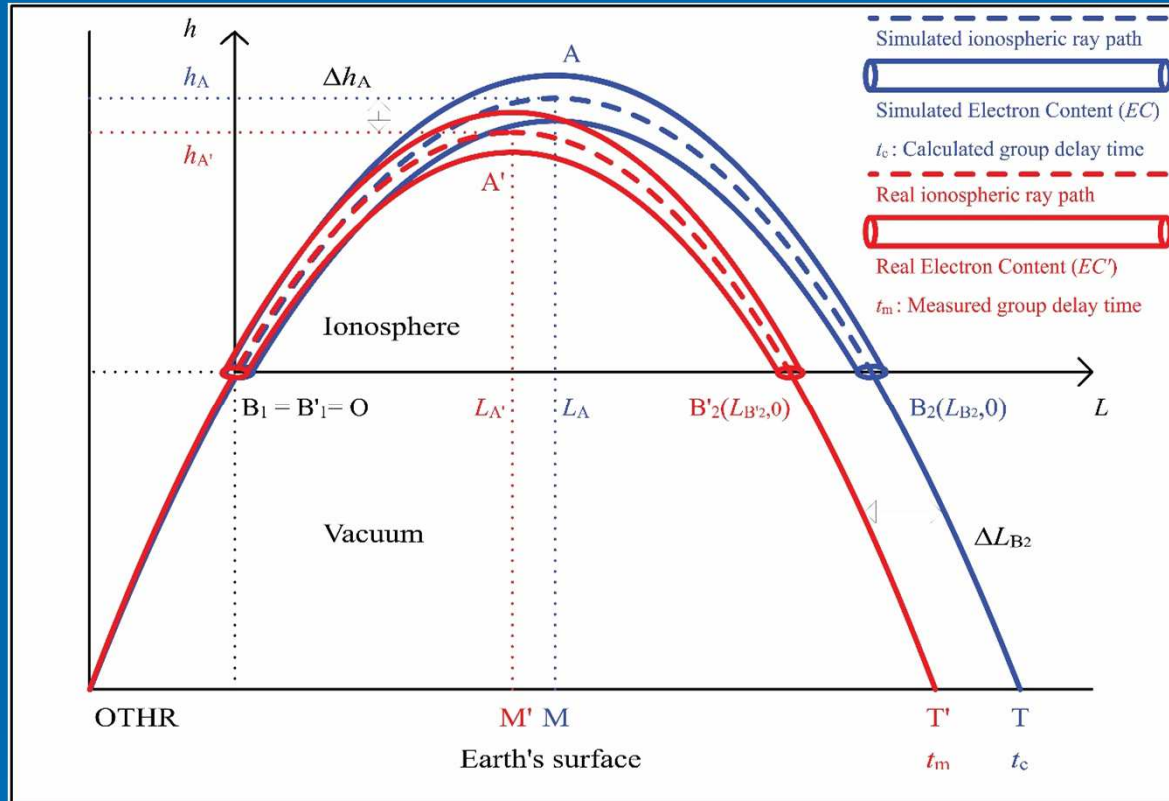
c) The plasma frequency profile $f_p(h)$ (in blue) and the corrected plasma frequency profile $f_p(h)+0.06 \cdot f_p(h)$ (in red) used by IONORT to synthesize the traces shown in b).

5. Conclusions

Being a RT program a deterministic computational procedure whose accuracy depends on the integration step, the only source of error is mainly referable to the electron density model representing the input parameter.

This talk demonstrated theoretically that, after the only possible measurement, which is the measured group delay time t_m , the correction can be performed: indeed, this quantity, compared with the calculated group delay time t_c , allows to calculate the time difference Δt between the group delays in ionosphere.

6. Appendix



$$\Delta L_{B_2} \ll L_{B_2}$$

$$\Delta h_A \ll h_A$$

$$\Delta L_{B_2} \cdot \Delta h_A \rightarrow 0$$

$$\Delta \Gamma = 0$$



$$\Delta L_{B_2} \approx \Delta h_A \left\{ \frac{1}{2} \frac{L_{B_2}}{h_A} - 2 \frac{\sqrt{1 + 16 \left(\frac{h_A}{L_{B_2}} \right)^2}}{\ln \left[4 \frac{h_A}{L_{B_2}} + \sqrt{1 + 16 \left(\frac{h_A}{L_{B_2}} \right)^2} \right]} \right\}$$

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[IMPACT FACTOR 1.238]