

Electromagnetic spatial dispersion in periodic metamaterials



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Outline

❖ Brief introduction on Periodic Metamaterials

❖ Spatial dispersion:
artificial chirality/bi-anisotropy

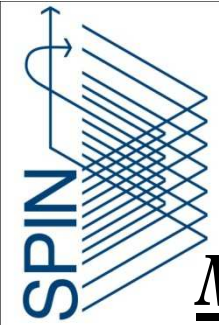
&

Artificial magnetism

❖ Homogenization theory: multi-scale approach

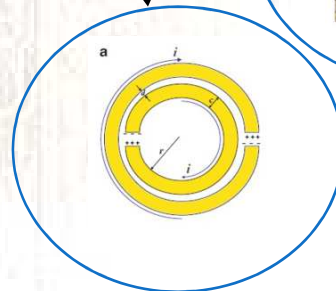
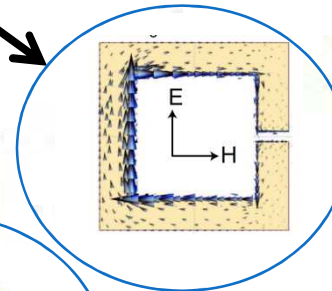
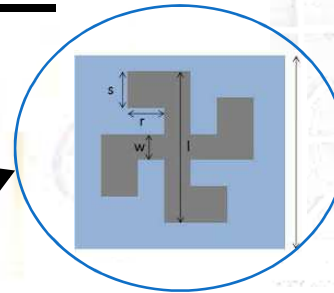
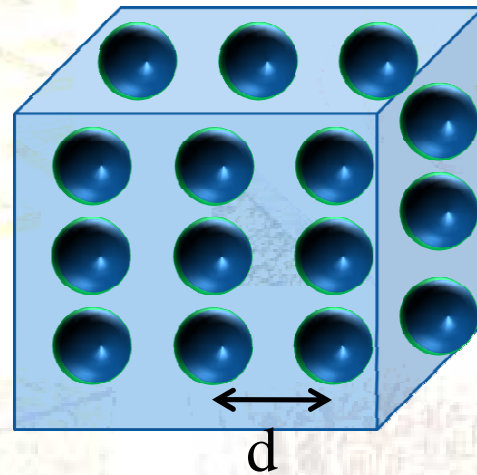
❖ 1D chiral metamaterials



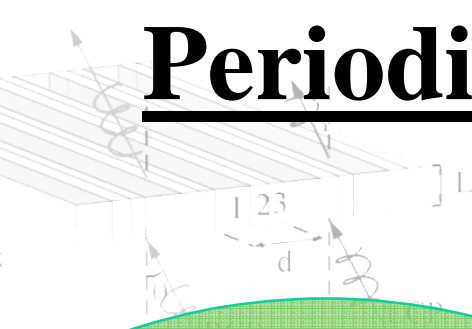
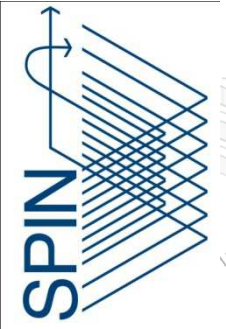


Metamaterials

Metamaterials are an arrangement of artificial structural elements, designed to achieve advantageous and/or unusual (electromagnetic) properties



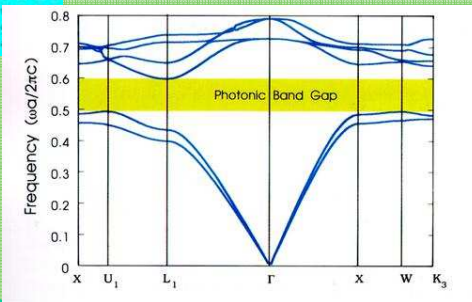
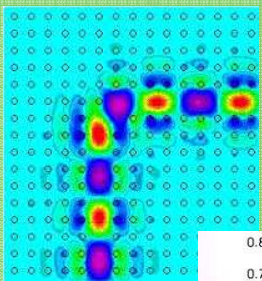
Periodic Metamaterials



Band Gap materials
(photonic crystals)

$$d \sim \lambda$$

Bragg scattering

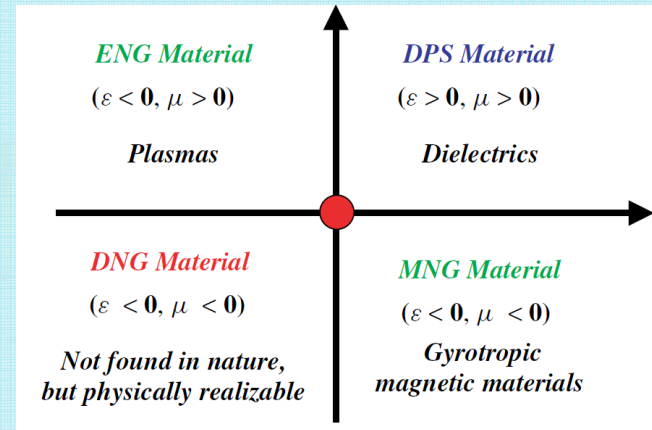
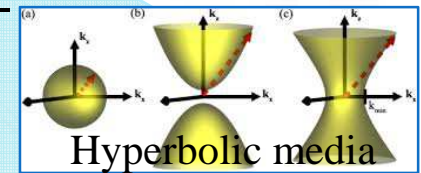


$$d < \lambda$$

Artificial magnetism
Artificial Bianisotropy

Effectively continuous metamaterials

$$d \ll \lambda$$

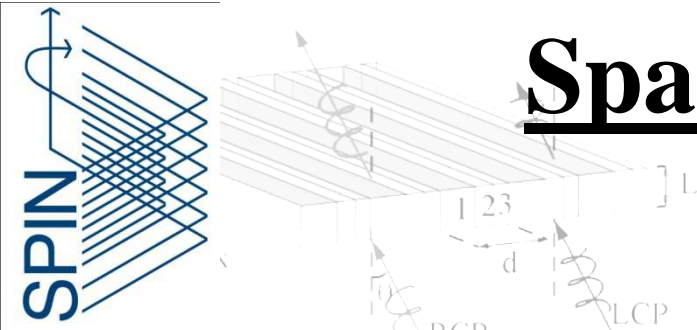


Structure organization of conventional media

Spatial dispersion- optical nonlocality



Spatial dispersion



$$\mathbf{D}(\mathbf{r}) = \int_V d\mathbf{r}' K(\mathbf{r}, \mathbf{r}') \mathbf{E}(\mathbf{r}')$$

Spatial dispersion- optical nonlocality



Second-order spatial dispersion

$$D_i = \epsilon_0 \left(\epsilon_{ij} E_j + \alpha_{ijk} \partial_k E_j + \beta_{iklj} \partial_k \partial_l E_j \right)$$

Maxwell's equations are invariant with respect to transformation:

$$\mathbf{D}' = \mathbf{D} + \nabla \times \mathbf{Q}, \quad \mathbf{H}' = \mathbf{H} - i\omega \mathbf{Q}$$

$\mu \neq 1$
(non magnetic inclusions!!)

First-order spatial dispersion
Bi-anisotropic response

$$\mathbf{D}' = \epsilon_0 \epsilon \mathbf{E} - \frac{i}{c} \mathbf{k}^T \mathbf{H}', \quad \mathbf{B} = \frac{i}{c} \mathbf{k} \mathbf{E} + \mu_0 \mu \mathbf{H}'$$

(non bi-anisotropic inclusions!!)

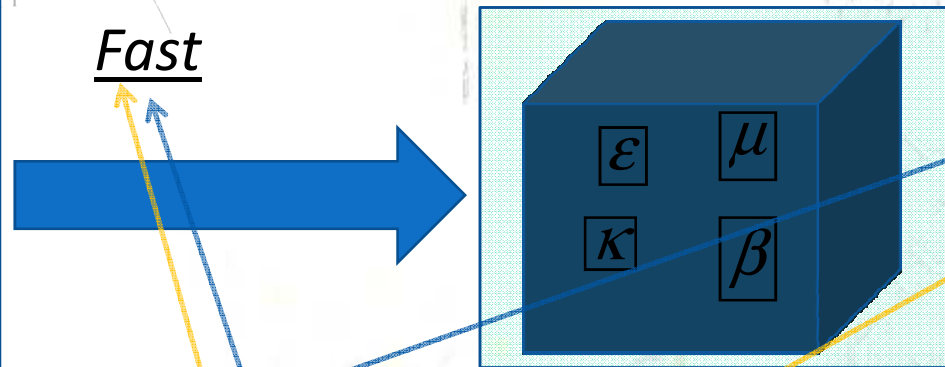
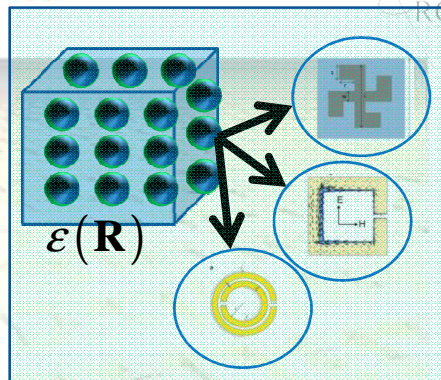
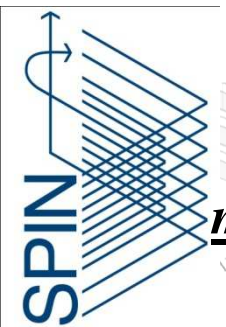
Chiral tensor



Homogenization theory



A suitable homogenization theory is the key ingredient to develop metamaterials allowing to mold the flow of electromagnetic waves in unprecedented ways.



Slow

“Microscopic” equations

$$\begin{aligned} \nabla \times \mathbf{E} &= i\omega\mu_0 \mathbf{H}, & \nabla \times \mathbf{H} &= -i\omega\mu_0 \mathbf{D} \\ \mathbf{D} &= \epsilon_0 \epsilon(\mathbf{R}) \mathbf{E}, & \mathbf{B} &= \mu_0 \mathbf{H} \end{aligned}$$

$$\epsilon(\mathbf{R}) \quad \mathbf{R} = \frac{\mathbf{r}}{\eta} \quad \eta = d/\lambda \ll 1$$

$$\mathbf{A} = \sum_{n=0}^{\infty} \left[\mathbf{A}_n(\mathbf{r}) + \tilde{\mathbf{A}}_n(\mathbf{r}, \mathbf{R}) \right] \eta^n$$

“Macroscopic” equations

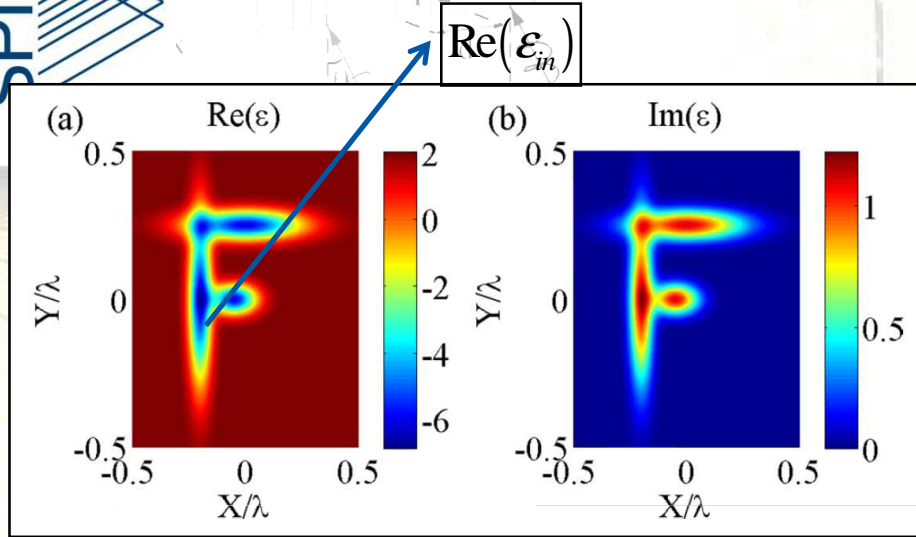
$$\begin{aligned} \nabla \times \bar{\mathbf{E}} &= i\omega\mu_0 \bar{\mathbf{H}}', & \nabla \times \bar{\mathbf{H}}' &= -i\omega\mu_0 \bar{\mathbf{D}}' \\ \bar{\mathbf{D}}' &= \epsilon_0 \epsilon^{(eff)} \bar{\mathbf{E}} - \frac{i}{c} \kappa^{(eff)T} \bar{\mathbf{H}}', & \bar{\mathbf{B}} &= \frac{i}{c} \kappa^{(eff)} \bar{\mathbf{E}} + \mu_0 \bar{\mathbf{H}}' \end{aligned}$$

$$\begin{aligned} \nabla_{\mathbf{R}} \cdot (\epsilon \nabla_{\mathbf{R}} f_j) &= -\partial_j \epsilon, & \tilde{\mathbf{E}}_0 &= \hat{e}_i \partial_i f_j \bar{\mathbf{E}}_{0j} \\ \epsilon_{ij}^{(eff)} &= \overline{Q_{ij} + Q_{ji}}, & Q_{ij} &= \epsilon (\delta_{ij} + \partial_i f_j) \\ \kappa_{ij}^{(eff)} &= \eta k_0 \left[e_{imj} \overline{\epsilon f_m} + \left(e_{imn} \delta_{jq} + \frac{1}{2} e_{mqn} \delta_{ij} \right) \overline{\epsilon f_m \partial_q f_n} \right] \end{aligned}$$

New!!!



2D chiral metamaterials



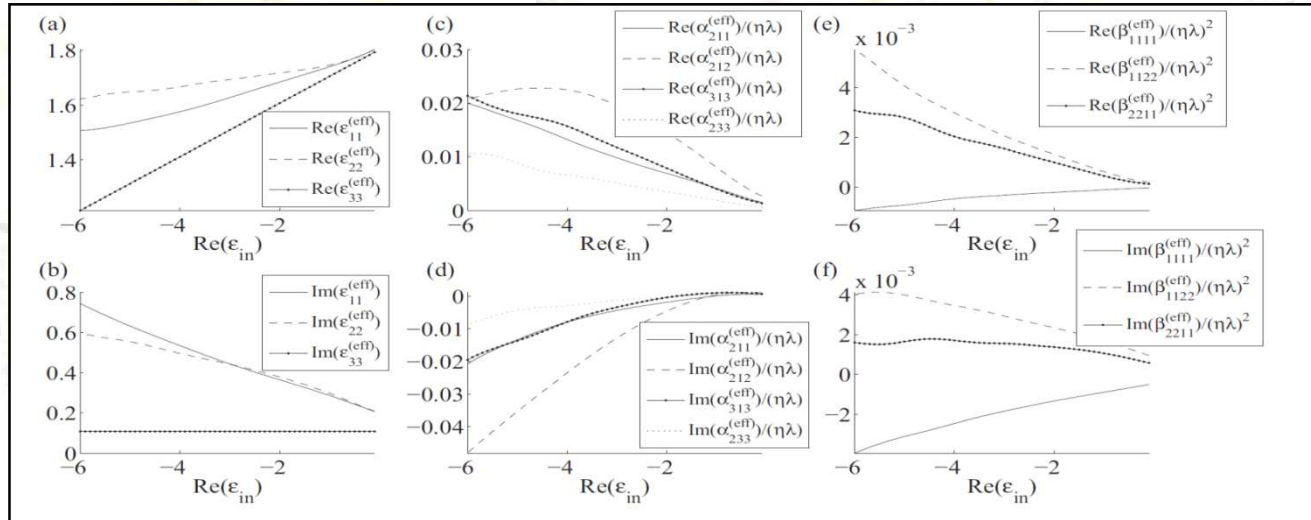
$$\nabla_{\mathbf{R}} \cdot (\boldsymbol{\varepsilon} \nabla_{\mathbf{R}} f_j) = -\partial_j \mathcal{E}$$

periodic boundary condition

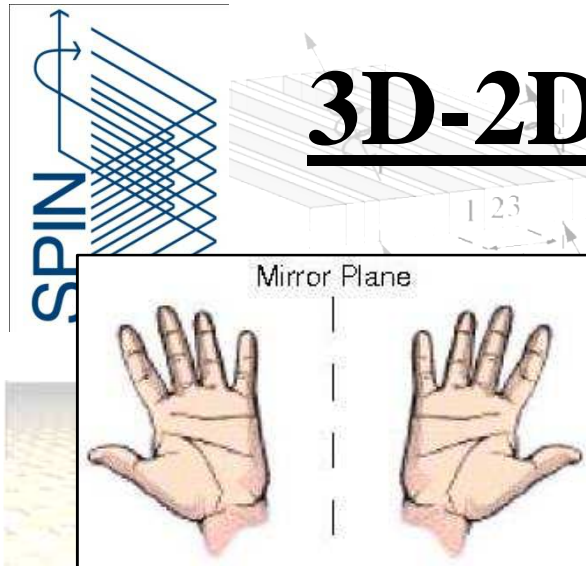
$$\bar{D}_i = \varepsilon_0 \left(\varepsilon_{ij}^{(eff)} E_j + \alpha_{ijk}^{(eff)} \partial_k E_j + \beta_{iklj}^{(eff)} \partial_k \partial_l E_j \right)$$

$$\bar{D}_i = \varepsilon_0 \varepsilon_{ij}^{(eff)} \bar{E}_j - \frac{i}{c} \kappa_{ji}^{(eff)} \bar{H}_j + \beta_{iklj}^{(eff)} \partial_k \partial_l \bar{E}_j,$$

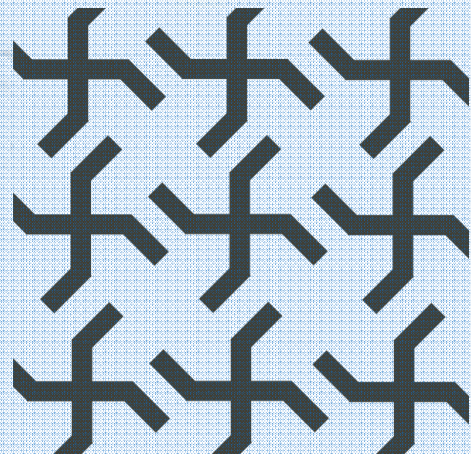
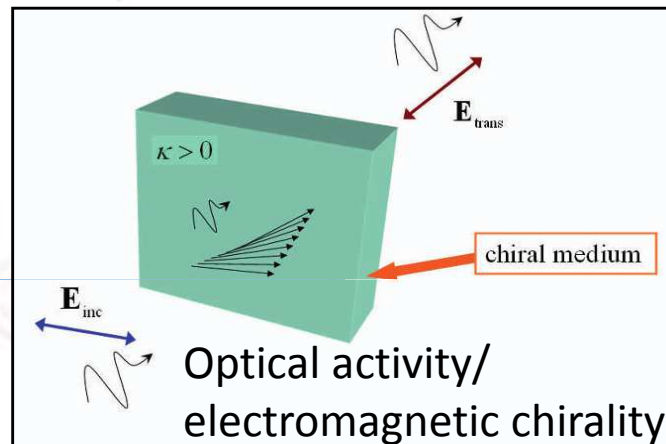
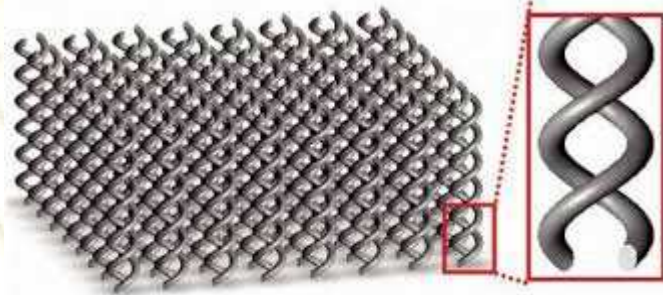
$$\bar{B}_i = \frac{i}{c} \kappa_{ij}^{(eff)} \bar{E}_j + \mu_0 \bar{H}_i$$



3D-2D chiral metamaterials



3D chirality: A left and right hand are mirror images of one another. Unlike a ball and its mirror image, a hand and its mirror image are not identical.

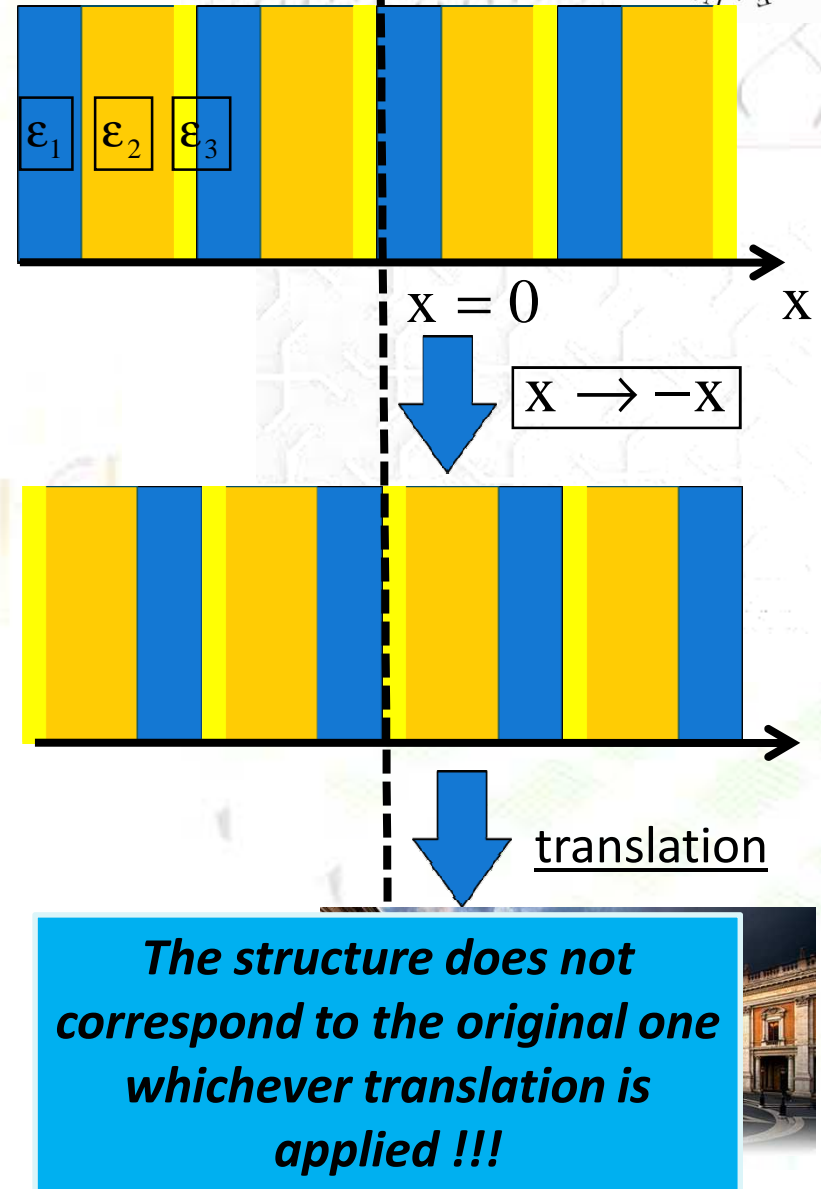
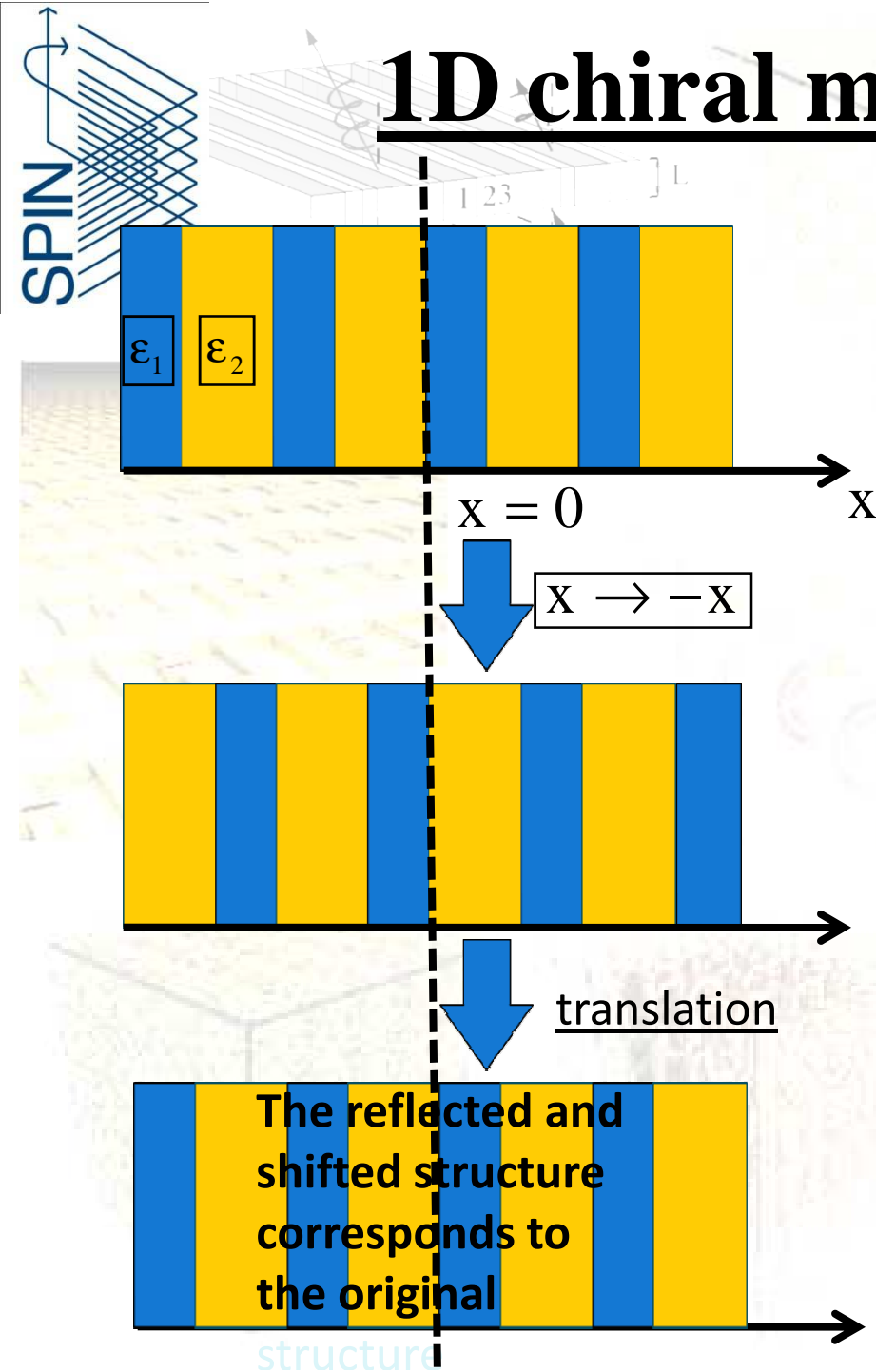


Planar object is said to be **2D chiral** if it cannot be brought into congruence with its mirror image unless it is lifted from the plane. **2D chirality produces bi-anisotropic effects**

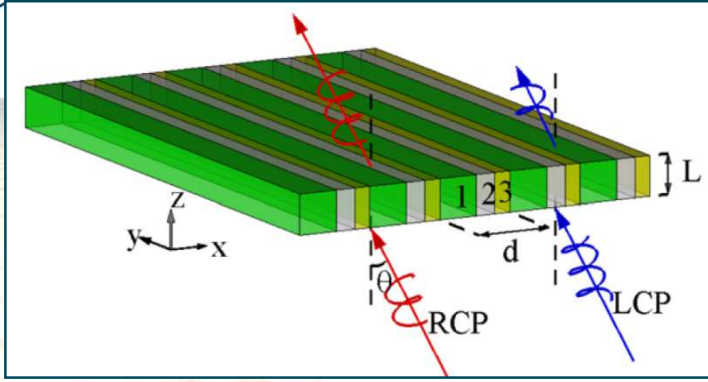
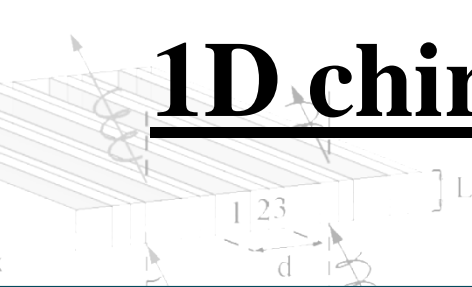
A. Papakostas *et al.*, "Optical Manifestations of Planar Chirality"
PRL 90, 107404 (2003)



1D chiral metamaterials



1D chiral Metamaterials

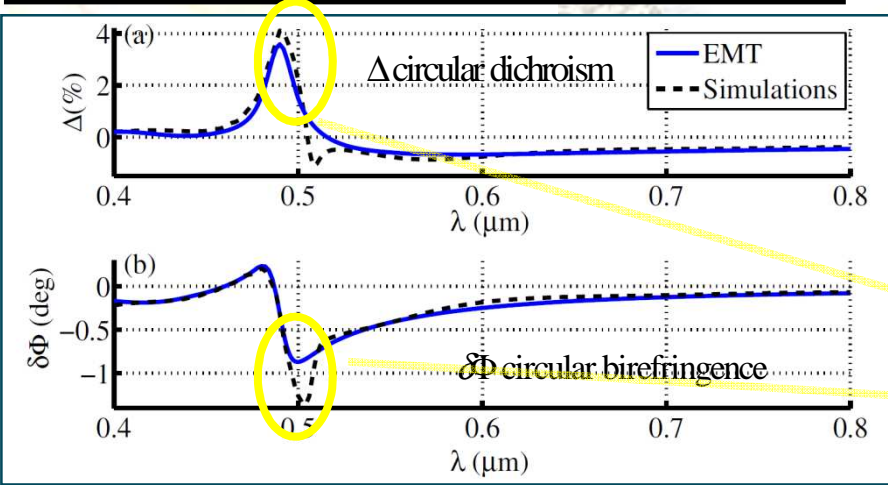
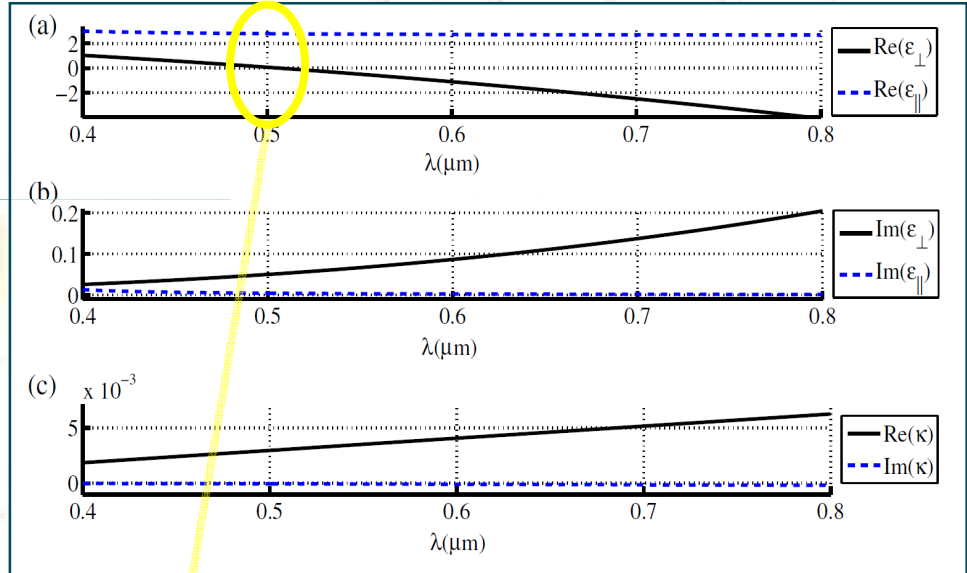


$$\epsilon^{(eff)} = \epsilon_0 \text{diag}(\epsilon_{\parallel}, \epsilon_{\perp} + \kappa^2, \epsilon_{\perp} + \kappa^2), \mu^{(eff)} = \mu_0 I$$

$$\kappa^{(eff)} = \begin{pmatrix} 0 & 0 & 0 \\ 0 & 0 & \kappa \\ 0 & -\kappa & 0 \end{pmatrix}$$

Omega-type response

Periodic 3-layer structure
 volume fractions: $f_1 = 0.18, f_2 = 0.33, f_3 = 1 - f_1 - f_2, d(\text{period}) = 25 \text{ nm}$
 Ag (1), SiO_2 (2), PMMA(3)



**Epsilon-near-zero
Enhancement !!!**



Conclusions

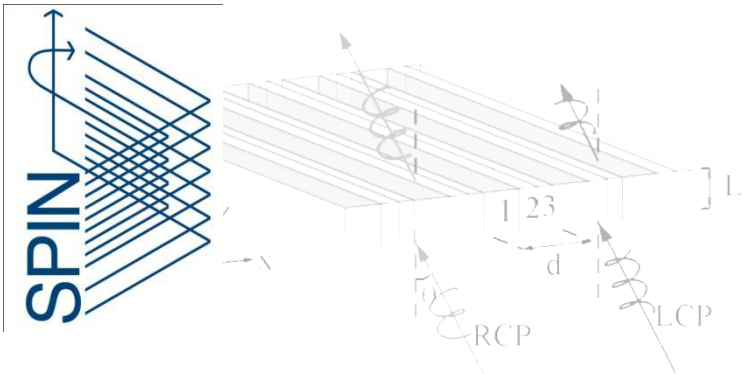
❖ Nonlocal homogenization theory in metamaterials

❖ 1D chirality & 1D chiral metamaterials

✓ THz reconfigurable all-optical metamaterials
(in collaboration with THz group, University of Bari)

✓ Metasurfaces, high dielectric contrast, effects of second-order dispersion





Thanks for your attention



Questions?

