

The neutrinoless double beta-decay: a brief overview

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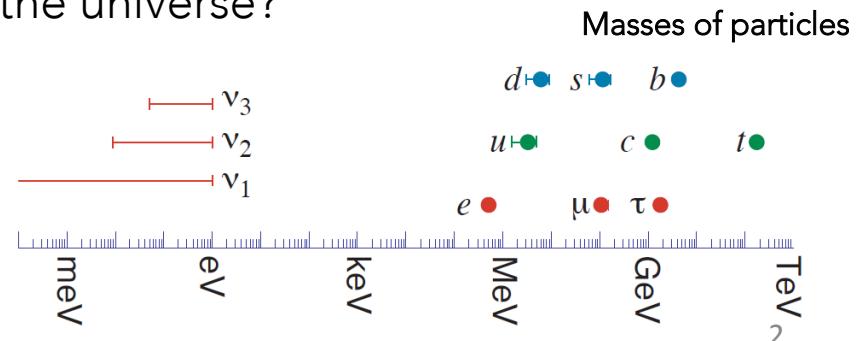
7th May 2015



"The Most Tiny Quantity of Reality ..."

F. Reines, ca. 1956

- ... is the **most abundant** matter particle in the universe
- ...plays important roles in :
 - Cosmology : BBN, CMB anisotropies & structure formation
 - Astrophysics : 99% of energy released in supernovae carried by neutrinos
- ...has a small but **non-zero** mass,
very different in scale from the other fermions : $\frac{m_\nu}{m_e} \sim 10^{-7}$
- Could the **very** small neutrino mass be related to physics at extremely **high** energy scales ? → "See-saw" mechanism ???
- Is there CP violation in the lepton sector and could this be related to the observed matter-antimatter asymmetry in the universe?

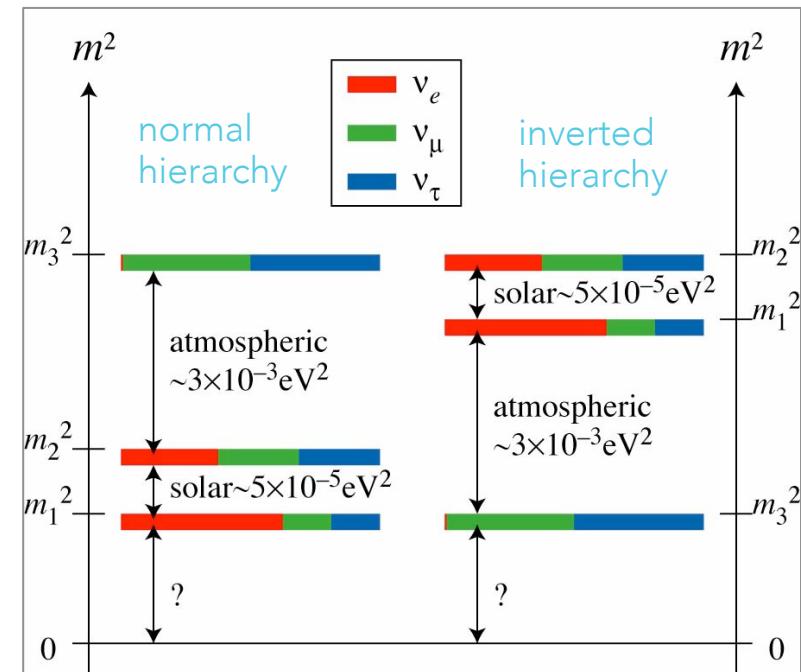


What we can learn from $0\nu\beta\beta$?

- neutrinos have mass and they mix
- precision measurements of mixing angles and Δm^2
- nature of neutrinos :
Dirac ($\nu \neq \bar{\nu}$) or Majorana ($\nu = \bar{\nu}$)
- absolute neutrino mass scale :
Only limits so far :
 - $m_{\nu_e} < 2.2$ eV (Tritium end-point)
 - $\sum m_{\nu_i} < 0.3$ eV (Cosmology)
- neutrino mass-hierarchy :
 - normal : $m_1 < m_2 < m_3$
 - inverted : $m_3 < m_1 < m_2$
- CP-violation in neutrino sector :
 - Dirac phase : $\delta \neq 0, \pi$
 - Majorana phases : $\alpha_{21}, \alpha_{31} \neq 0, \pi$

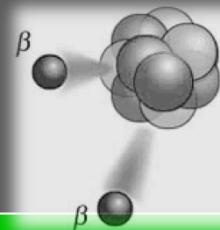
PMNS mixing matrix :

$$U = \begin{pmatrix} U_{e1} & U_{e2} & U_{e3} \\ U_{\mu 1} & U_{\mu 2} & U_{\mu 3} \\ U_{\tau 1} & U_{\tau 2} & U_{\tau 3} \end{pmatrix} \times \boxed{\begin{pmatrix} 1 & 0 & 0 \\ 0 & e^{i\alpha_{21}/2} & 0 \\ 0 & 0 & e^{i\alpha_{31}/2} \end{pmatrix}}$$



Graphics taken from DOE Nuclear Science Advisory Committee report on $0\nu\beta\beta$ (24 April 2014)

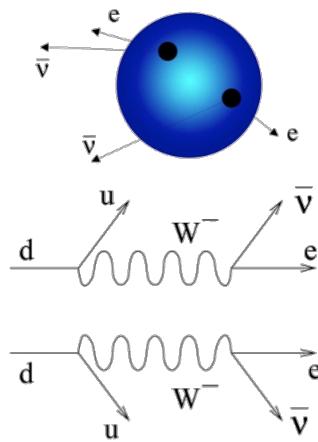
Double Beta Decay



2-neutrino double beta decay

$$(A, Z) \rightarrow (A, Z + 2) + 2e^- + 2\bar{\nu}$$

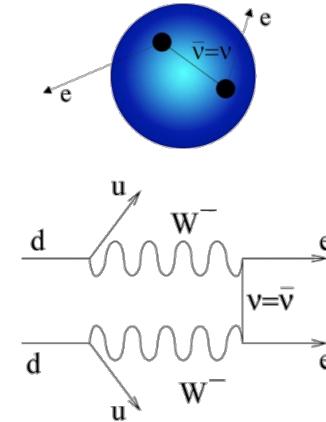
[Maria Goeppert-Meyer in 1935]



0-neutrino double beta decay

$$(A, Z) \rightarrow (A, Z + 2) + 2e^-$$

[Furry in 1939]



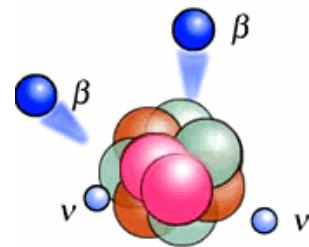
- Lepton number violating process $\Delta L=2$
- **forbidden in SM** → new physics !!!
- rate ($0\nu\beta\beta$) << rate ($2\nu\beta\beta$)
- **no positive detection** so far, with the exception of Klapdor- Kleingrothaus *et al.*, however result been disproved.

- 2nd order process allowed within SM
- **first observed in 1987** (Elliott *et al.*)
- today measured in several isotopes

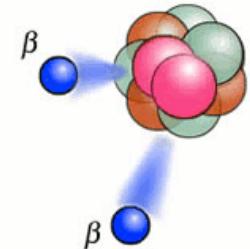
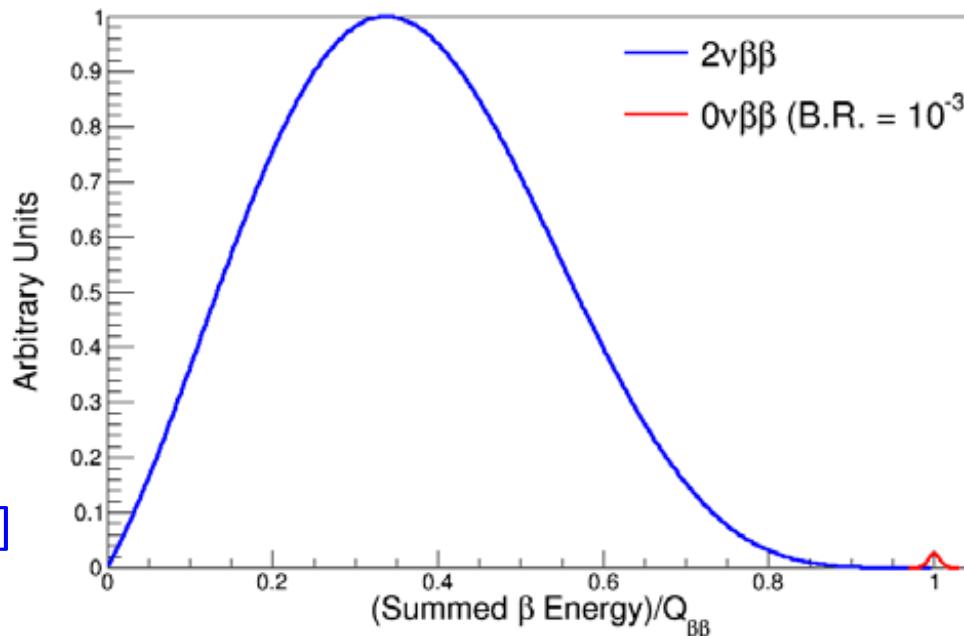
Experimental signature

Experiments measure the sum of the kinetic energies of the two emitted β 's

Signature: monochromatic line at the Q-value of the decay



$$(E_1 + E_2) \in [0, Q_{\beta\beta}]$$



$$(E_1 + E_2) / Q_{\beta\beta} \approx 1$$

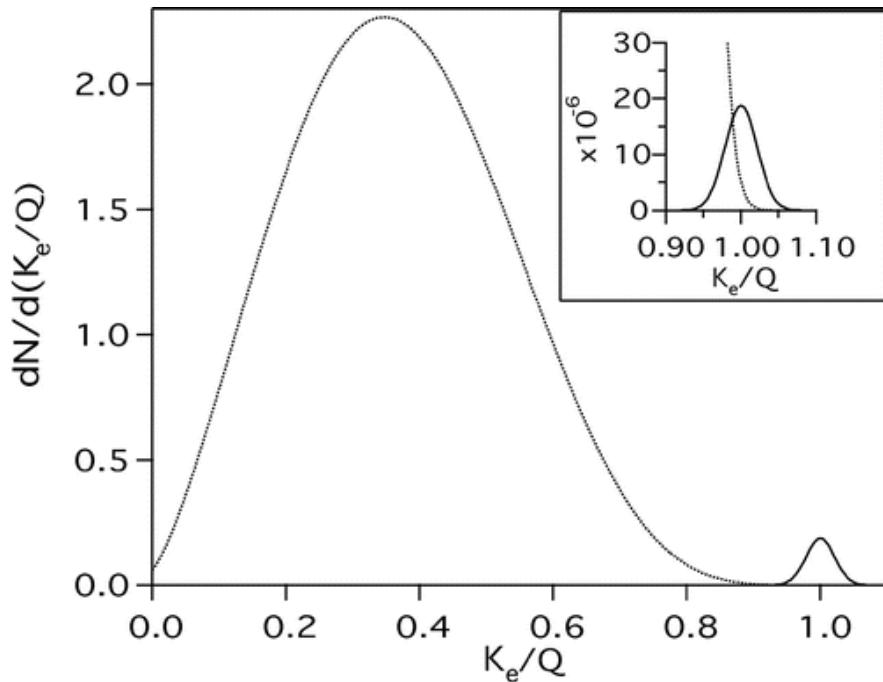
Energy resolution!!

Irreducible background

Experiments measure the sum of the kinetic energies of the two emitted β s

Signature: monochromatic line at the Q-value of the decay

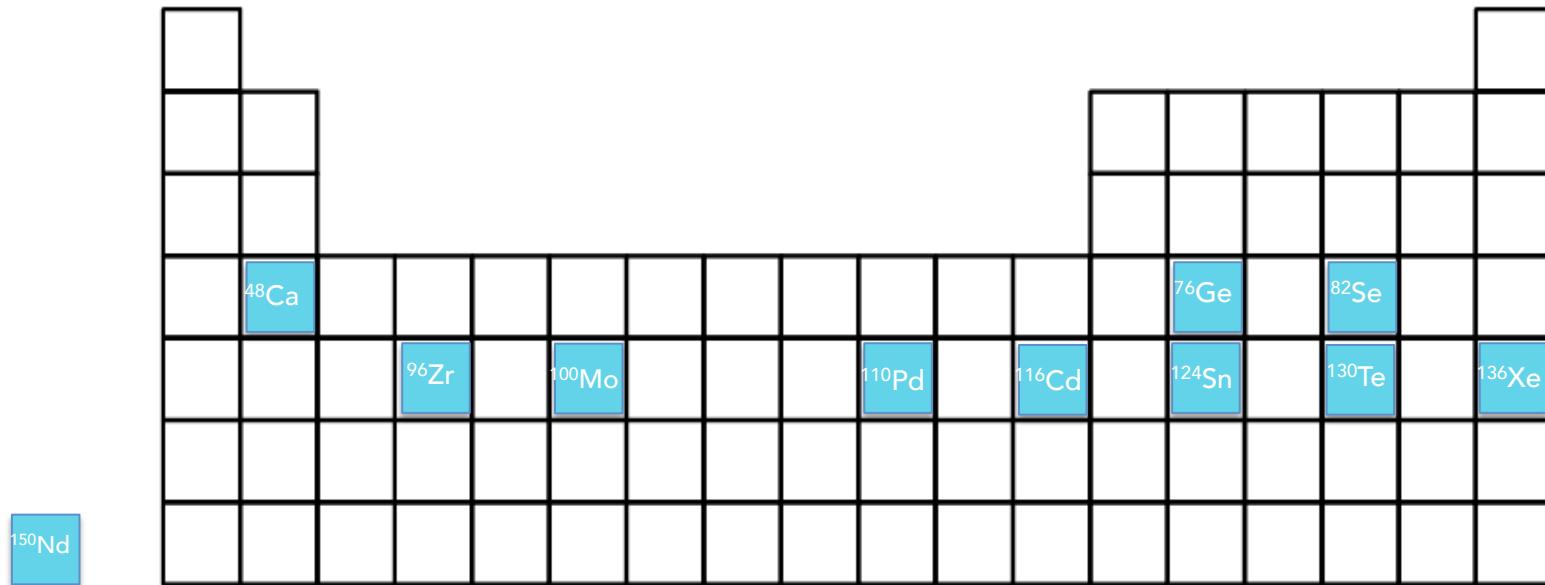
$2\nu\beta\beta$ irreducible background:



good energy
resolution ΔE
necessary!

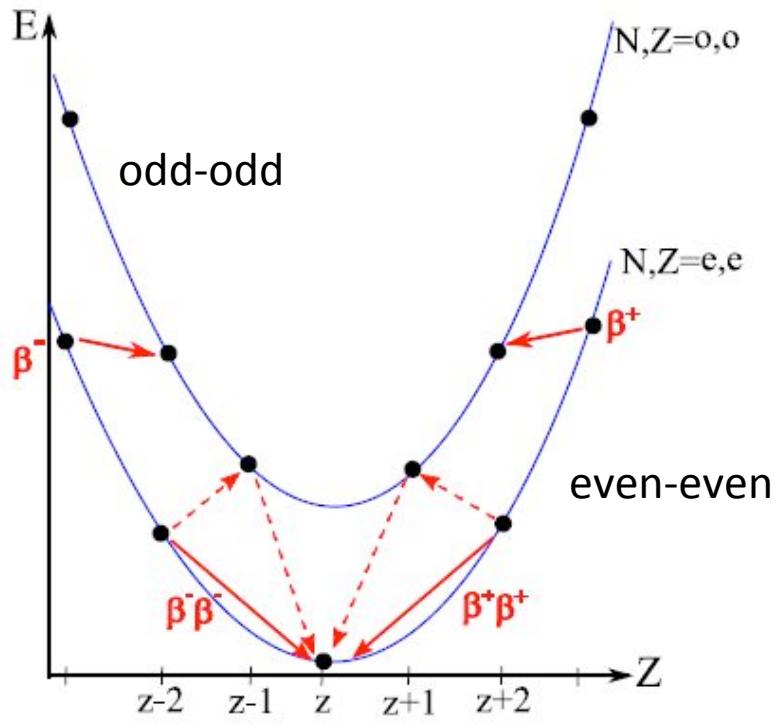
Who can do $0\nu\beta\beta$?

35 isotopes available, interesting for $0\nu\beta\beta$ searches only 11:



Isotopes of choice

only even-even nuclei can undergo a double beta decay



| Candidate Isotope | Q-value [MeV] | Natural abundance [%] |
|-------------------|---------------|-----------------------|
| ^{48}Ca | 4.272 | 0.187 |
| ^{76}Ge | 2.039 | 7.8 |
| ^{82}Se | 2.996 | 9.2 |
| ^{96}Zr | 3.350 | 2.8 |
| ^{100}Mo | 3.034 | 9.6 |
| ^{110}Pd | 2.004 | 11.8 |
| ^{116}Cd | 2.814 | 7.6 |
| ^{124}Sn | 2.530 | 5.6 |
| ^{130}Te | 2.530 | 34.5 |
| ^{136}Xe | 2.459 | 8.9 |
| ^{150}Nd | 3.371 | 5.6 |

higher Q-value
→ less background !!

Enrichment possible
but always expensive !!

Effective Neutrino Mass (light ν exchange)

REMEMBER: If neutrino has a non-zero mass, mass eigenstates (ν_1, ν_2, ν_3) and weak interactions eigenstates (ν_e, ν_μ, ν_τ) do not necessarily have to coincide!

$$\nu_l = \sum_i U_{li} \nu_i$$

- β -decay: incoherent sum
real neutrino
- cosmology: simple sum
pure kinematic effect
- $0\nu\beta\beta$ -decay: coherent sum
virtual neutrino
Majorana phases

$$\langle m_\beta \rangle \equiv \left(\sum_{i=1}^3 m_i^2 |U_{li}|^2 \right)^{1/2}$$

$$\Sigma \equiv \sum_{i=1}^3 m_i$$

$$\langle m_{\beta\beta} \rangle \equiv \left(\sum_{i=1}^3 m_i |U_{li}|^2 e^{i\alpha_i} \right)^{1/2}$$

α_i = Majorana phases

0vDBD experiments are sensitive to an **effective** neutrino mass!

Half-life of the decay

The decay half-life $T_{1/2}^{0\nu}$ is the **measurable quantity**. It depends on the effective Majorana mass $\langle m_{\beta\beta} \rangle^2$ of the neutrino exchanged between the two electron vertexes

Depending on g_A
half-lives can vary by a factor of 4-16

$$\frac{1}{T_{1/2}^{0\nu}} = g_A^4 G_{0\nu} |M_{0\nu}|^2 \left| \frac{\langle m_{\beta\beta} \rangle}{m_e} \right|^2$$

half-life

axial vector coupling constant

phase space integral $\sim Q^5$
atomic physics

nuclear matrix element (NME)
nuclear physics

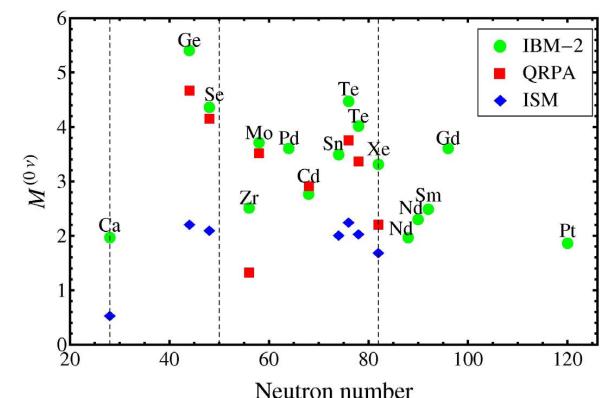
effective Majorana mass
particle physics

$$m_{\beta\beta} = f(\theta_{12}, \theta_{23}, \theta_{13}, \Delta m_{12}, \pm \Delta m_{13}, \Sigma_{cosm}, e^{i\alpha}, e^{i\beta})$$

K. Schäffner - IAPS @ LNGS

Several methods have been used to evaluate $M_{0\nu}$:
QRPA (Quasiparticle Random Phase Approximation)
ISM (Shell Model)

IBM-2 (Interacting Boson Model)
EDF (Density Functional Theory)



IBM-2 from J. Barea and F. Iachello, Phys. Rev. C 79, 044301 (2009);
J. Barea, J. Kotila and F. Iachello, Phys. Rev. C 87, 014315 (2013). MS-SRC.
QRPA from F. Šimkovic et al., Phys. Rev. C 77, 045503 (2008).
ISM from E. Caurier et al., Phys. Rev. Lett. 100, 052503 (2008).

Experimental sensitivity

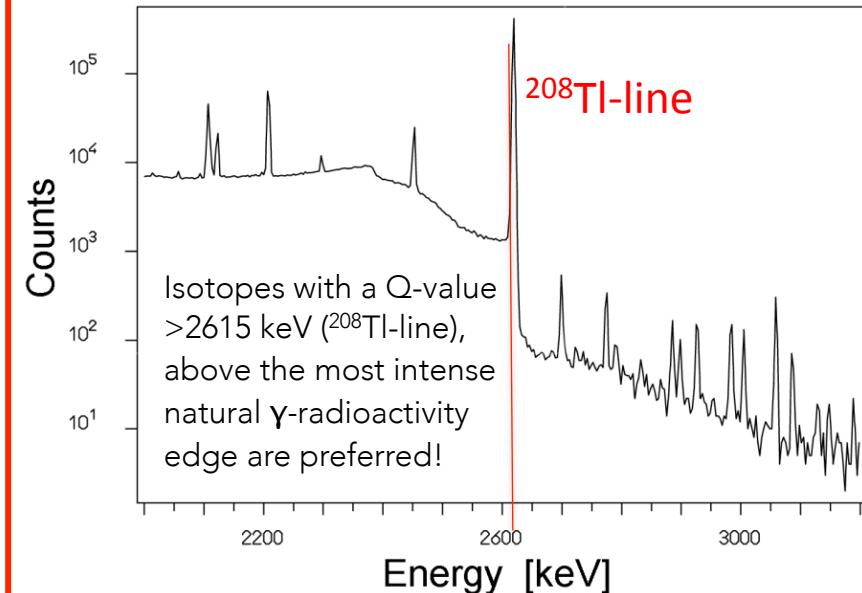
$$S_{0\nu} = \ln(2) N_A \frac{\eta \cdot \epsilon}{W} \sqrt{\frac{M \cdot T}{\Delta E \cdot B}}$$

STRATEGY

- low background
 - underground experiment
 - material selection, purification
- high Q-value
 - reject most of natural radioactivity
- high number of DBD nuclei
 - enrichment: ^{136}Xe easy, ^{48}Ca difficult
 - ^{130}Te natural high abundance
- excellent energy resolution to limit effect from $2\nu\beta\beta$

- M = detector mass [kg]
- T = live time [y]
- B = background in ROI [counts $\text{keV}^{-1} \text{kg}^{-1} \text{y}^{-1}$]
- W = molecular weight
- η = isotopic abundance
- N_A = Avogadro number
- ϵ = detector efficiency
- ΔE = energy resolution at Q-value

Environmental "underground" HPGE bck:
 ^{238}U and ^{232}Th trace contaminations



How to build a $0\nu\beta\beta$ detector?

detector
≠
source

Tracker

^{82}Se , ^{150}Nd

PRO

full final state kinematics
“Smoking-gun” signature for $0\nu\beta\beta$
Isotope flexible

TPC
= Time Projection Chambers

^{136}Xe

detector
= source

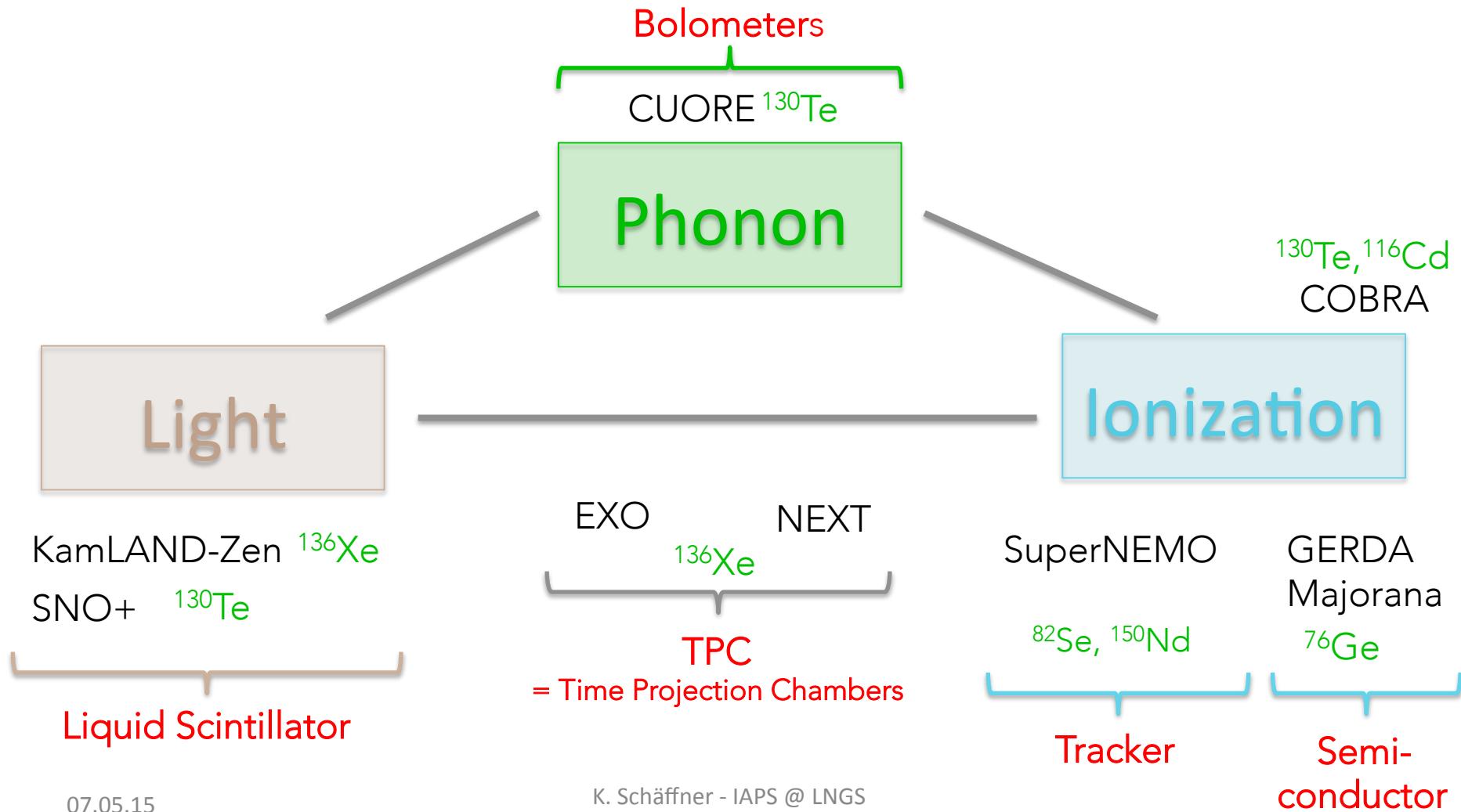
Bolometers
Semiconductors
Liquid Scintillators
Scintillator

^{76}Ge ^{130}Te ^{116}Cd

PRO

excellent energy resolution
compact
high detection efficiency

Present detection technologies

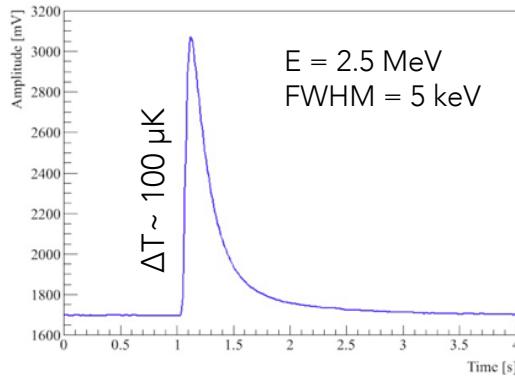
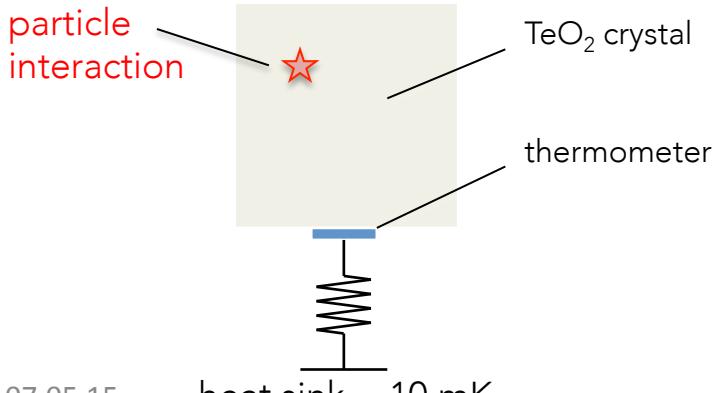
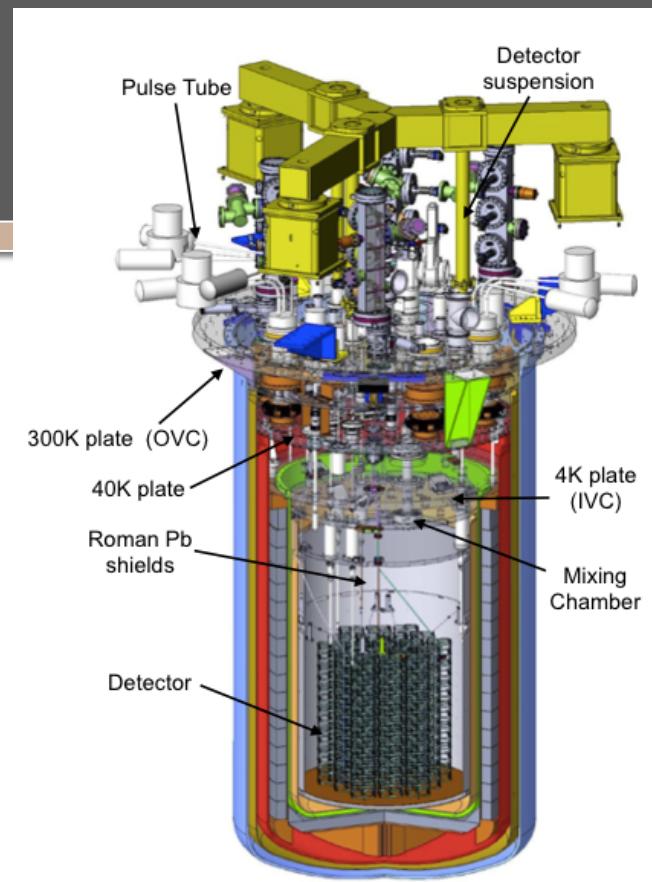


CUORE

- located at LNGS in Italy, start within 2015
- cryogenic bolometer experiment @ 10mK
- target: 988 crystals of TeO_2
- 206 kg fiducial mass of ^{130}Te
- Q-value: 2527 keV
- $\Delta E \sim 5\text{keV in ROI}$

GOAL

- background < 10 counts / (keV ton y)
- 5 years sensitivity: $T_{1/2}^{0\nu} > 9.5 \times 10^{25} \text{ y}$
 $\langle m_{\beta\beta} \rangle < 50 - 130 \text{ meV}$



Advantages of bolometers:

- excellent energy resolution
- large mass array
- low background in ROI
- wide choice of target material

GERDA phase II

- located at LNGS in Italy
- Ge diodes enriched to 86% with ^{76}Ge
- Q-value: 2039 keV
- $\Delta E \sim 3 \text{ keV in ROI}$
- GERDA phase I completed (2011-2013):
 $T_{1/2}^{0\nu} > 2.1 \times 10^{25} \text{ y}$ (90% C.L.)
excluded claim from Heidelberg-Moscow experiment
- GERDA phase II:
 - increase target mass 18kg + **20 kg**
 - PSD to distinguish multi-site events from single-site events
 - Compton veto using scintillation light detection from LAr



GOAL

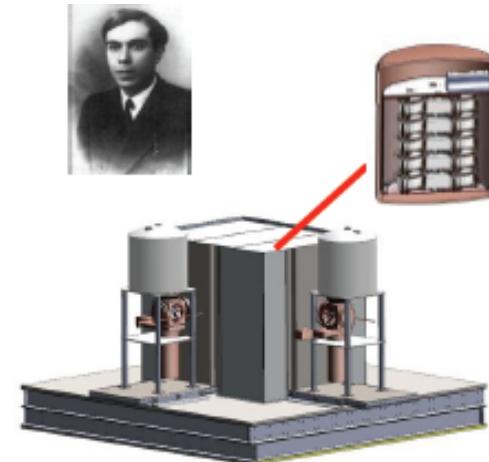
- background $< 0.001 \text{ counts / (keV kg y)}$
- sensitivity after 4 years: $T_{1/2}^{0\nu} > 1.5 \times 10^{26} \text{ y}$
 $\langle m_{\beta\beta} \rangle < 80 - 190 \text{ meV}$

Advantages of Ge-detectors:

- known technology
(enrichment + diode production)
- best energy resolution
- low bck in ROI
- flat background at Q-value

Majorana Demonstrator

- located at Sanford Underground Research Facility (SURF), South Dakota
- modular instrument composed of two cryostats
- cryostat built from ultra-pure electroformed copper
- p-type point contact Ge diodes:
 - 10 kg natural Ge
 - 30 kg enriched to 87% with ^{76}Ge
- Q-value: 2039 keV
- $\Delta E \sim 3 \text{ keV}$ in ROI (4keV wide window)

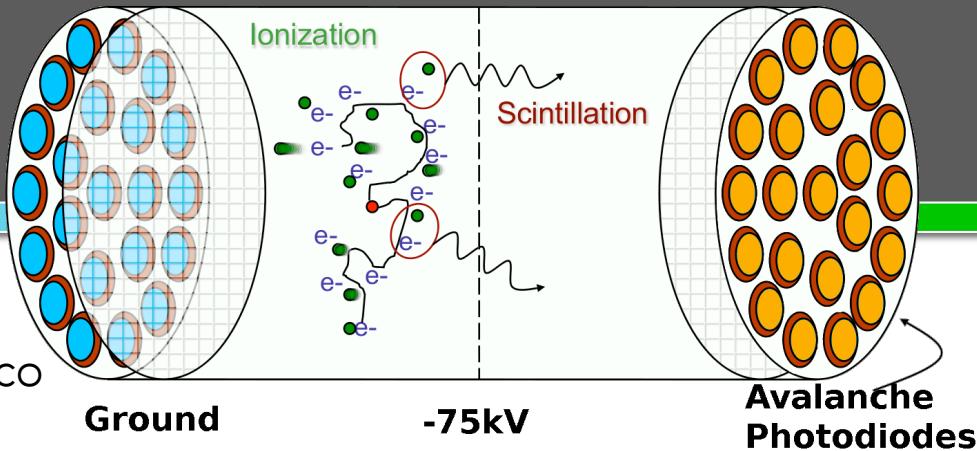


GOAL

- background ~ 3 counts / (ROI t y)
- collaborations aims to establish scalability to 1-tonne detector

EXO-200

- located at Waste Isolation Pilot Plant (WIPP) near Carlsbad New Mexico
- liquid filled TPC:
200 kg of enriched Xe (80.7% of ^{136}Xe)
of which 110 kg (79 kg fiducial isotope mass) are active in the detector
- discrimination of multi-scattering events from single-scatters
- Q-value: 2457.8 keV
- $\Delta E \sim 88.5$ keV in ROI
- EXO-200 (since 2011 in operation):
 $T_{1/2}^{0\nu} > 1.1 \times 10^{25} \text{ y}$ (90% C.L.)



GOAL

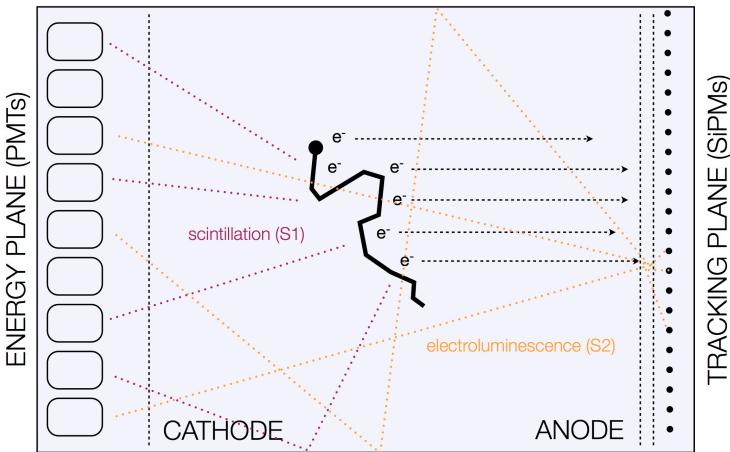
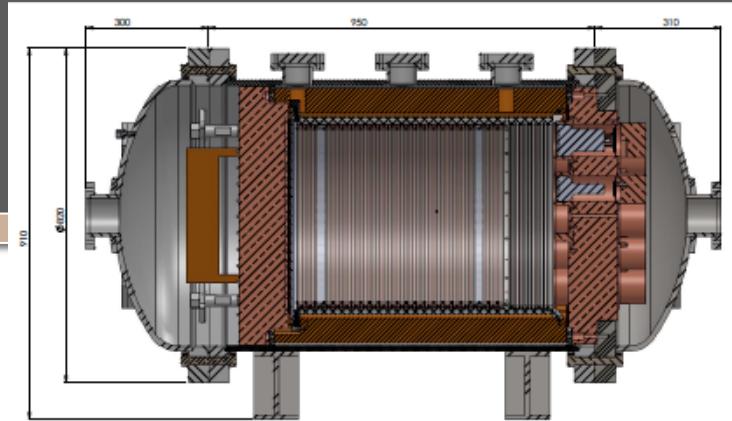
- reduce background from ^{214}Bi (Q=2447.8 keV) and ^{208}Tl (Q=2614.5 keV)
- reduce fiducial volume uncertainty
- improve energy resolution to about 56 keV
- aimed sensitivity after 3 years: $T_{1/2}^{0\nu} \sim 10^{26} \text{ y}$

Advantages of TPCs:

- easy to scale
- good self shielding if large

NEXT

- will be installed at Canfranc Underground Laboratory in Spain
- NEXT phase I presently in construction
- gaseous TPC (15 bar):
10 kg of ^{136}Xe for phase I
100 kg for NEXT-100 (scheduled for 2016)
- Q-value: 2457.8 keV
- tracking: silicon photomultipliers
- calorimetry: PMTs
- $\Delta E \sim 18$ keV in NEXT R&D



GOAL

- demonstrate that $\Delta E \sim 18$ keV are reachable
- increase fully tagged reconstruction efficiency (presently $\sim 30\%$)

Advantages of high pressure TPCs:

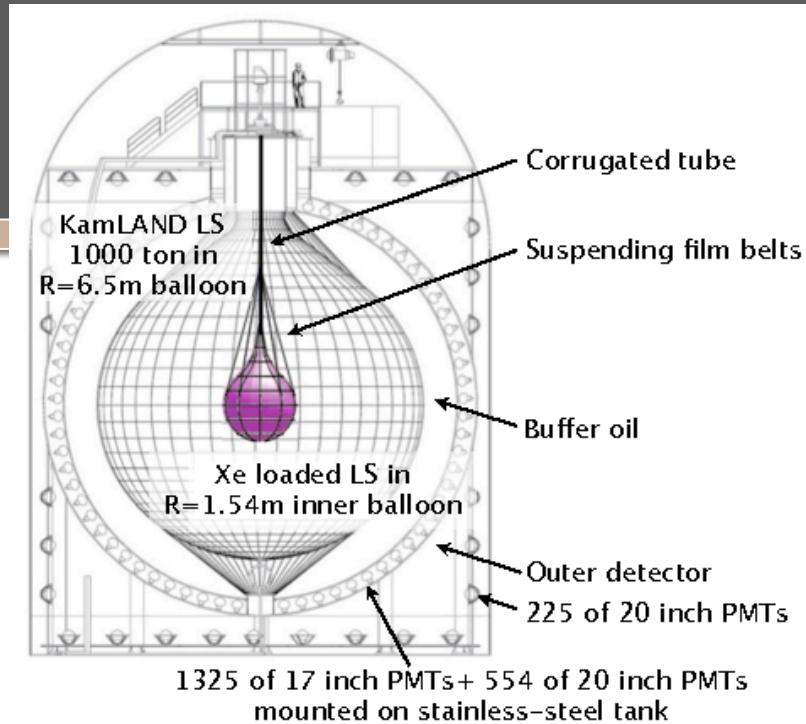
- easy to scale
- measure and image the final state kinematics of $0\nu\beta\beta$ event

KamLAND-Zen

- housed in the Kamioka mine in Japan
- mini-balloon filled with liquid scintillator loaded with enriched ^{136}Xe (90%)
- Q-value: 2457.8 keV
- $\Delta E \sim 250$ keV in ROI
- Phase I
 - 320 kg of ^{136}Xe (90%), background from $^{110\text{m}}\text{Ag}$
 - $T_{1/2}^{0\nu} > 1.9 \times 10^{25} \text{ y}$ (90% C.L.)
- Phase II (from Dec 2013)
 - 383 kg of ^{136}Xe (90%); LS purification; $^{110\text{m}}\text{Ag}$ reduced by a factor of 10
 - $T_{1/2}^{0\nu} > 1.3 \times 10^{25} \text{ y}$ (90% C.L.) → combined results: $T_{1/2}^{0\nu} > 2.6 \times 10^{25} \text{ y}$ (90% C.L.)
 $\langle m_{\beta\beta} \rangle < (140-280)\text{meV}$

GOAL

- Purification of liquid scintillator
- Increase mini-balloon volume to 600 kg
- sensitivity after 2 years: $T_{1/2}^{0\nu} > 2.0 \times 10^{26} \text{ y}$
 $\langle m_{\beta\beta} \rangle \sim 50 \text{ meV}$

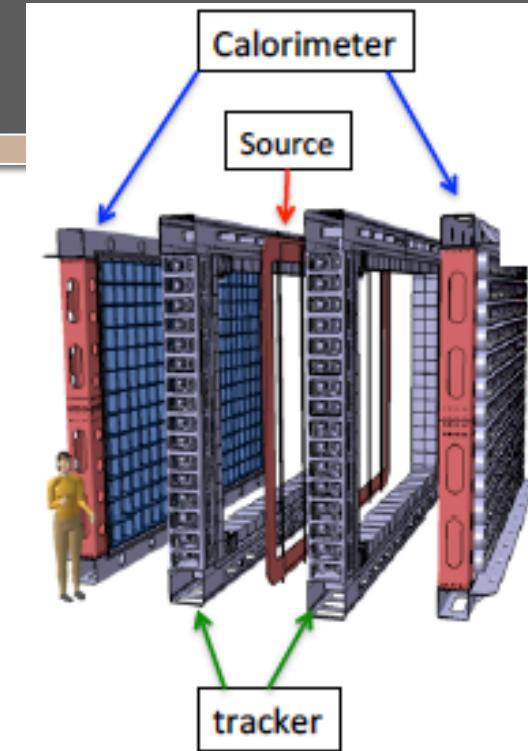


Advantages of Liquid Scintillators

- easy to scale
- possibility of loading LS with other isotopes

SuperNEMO

- located at the Modane Underground Laboratory at Fréjus
- the first "Demonstrator" module (7kg of ^{82}Se and/or ^{150}Nd) in commissioning and will run from 2015-2017
- 100-kg full SuperNEMO planned to start from 2017
- Tracking-and-calorimetry approach:
 - Central source frame: isotope on thin foils
 - Tracking: 2000 drift chambers
 - Segmented calorimetric elements (plastic scintillator) for measurement of the β^- energy and time-of-flight
- $\Delta E \sim 250 \text{ keV} @ 1\text{MeV}$



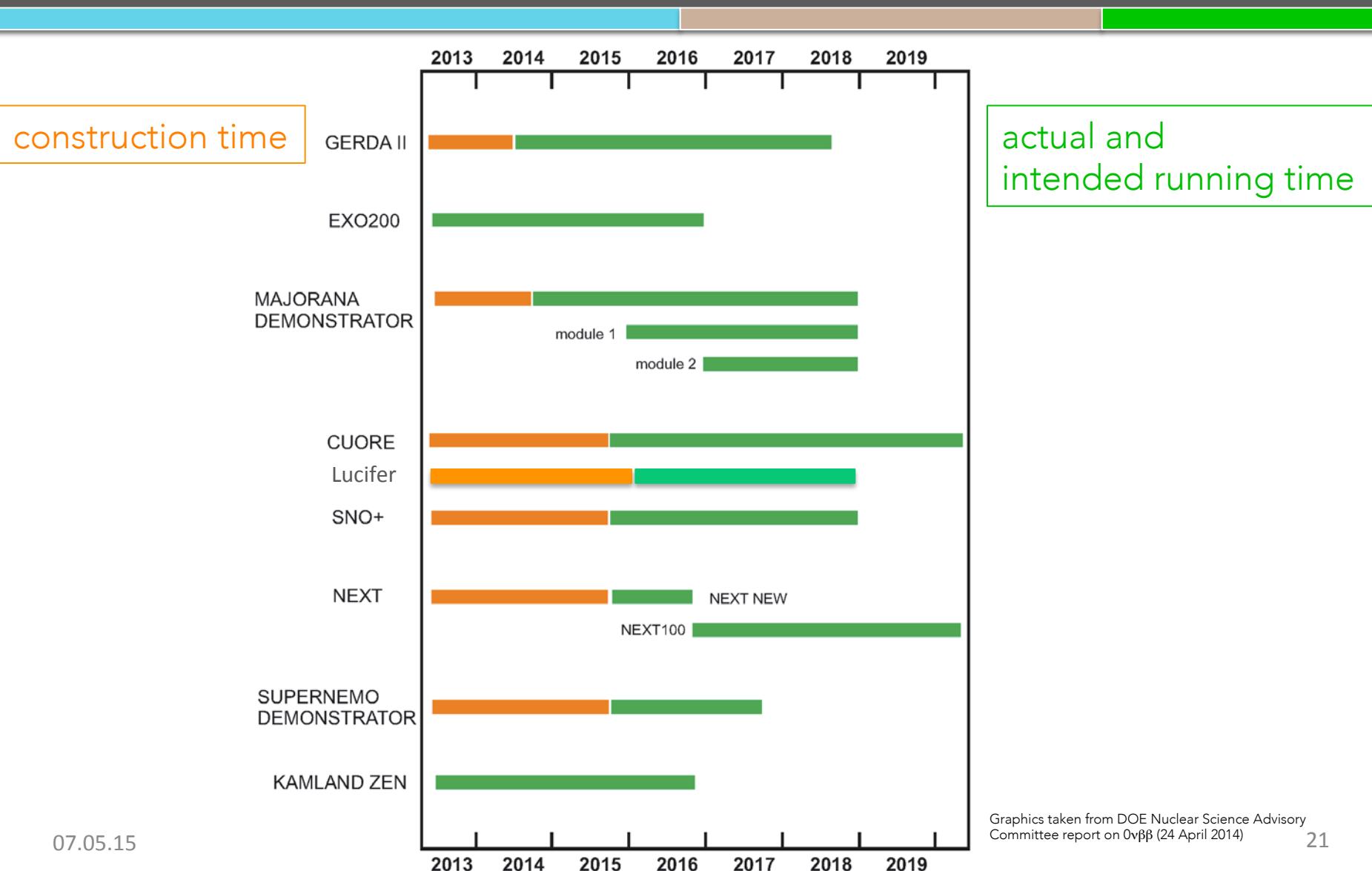
GOAL

- No background expected in ROI for 2 years of data
- sensitivity after 2 years: $T_{1/2}^{0\nu} > 6.6 \times 10^{24} \text{ y}$ for ^{82}Se
 $\langle m_{\beta\beta} \rangle \sim 160 - 440 \text{ meV}$

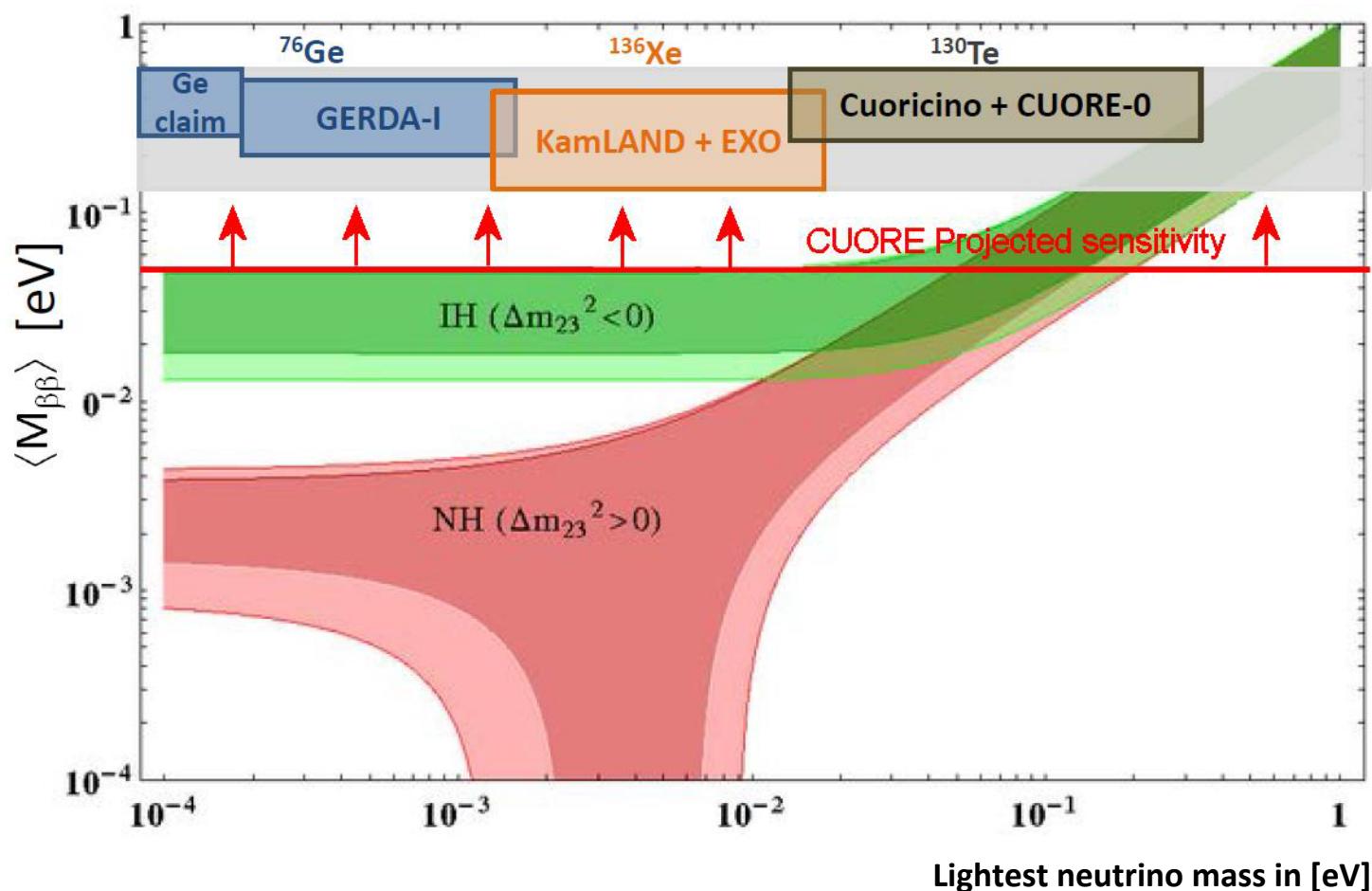
Advantages:

- allows to measure full kinematics of β^- products (angular distribution)
- scaling possible
- choose not mainstream isotopes: ^{48}Ca , ^{150}Nd , ^{82}Se (large PSF and low bck)

Timeline of present projects

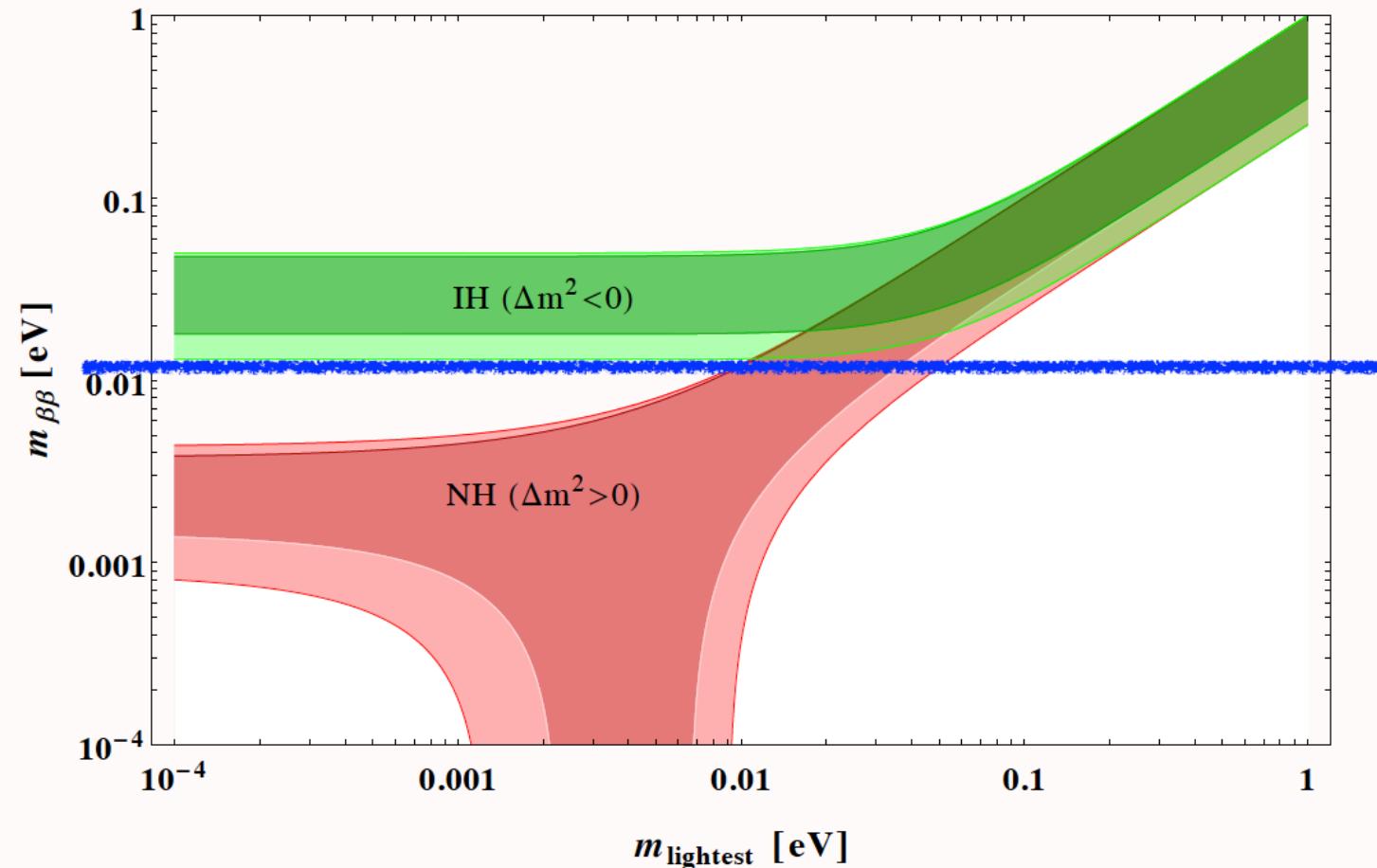


State of the art



What next ??

Dell'Oro, Marcocci e Vissani, Phys. Rev. D 90, 033005 (2014)



What is the ultimate strategy?

$$S_{0\nu} = \ln(2) N_A \frac{\eta \cdot \epsilon}{W} \sqrt{\frac{M \cdot T}{\Delta E \cdot B}}$$

Time T : not much to do other than wait

Increase in energy resolution ΔE :

- translates in an increase of the sensitivity of the experiments
- allows to better identify and account for background contributions from natural decay chains (^{214}Bi for ^{136}Xe experiments)
- allows to minimize the contribution from the irreducible contribution of the $2\nu\beta\beta$ decay

BUT: Ge-diodes and bolometers already show very good energy resolutions and liquid scintillator experiments are limited by light production and consequent light detection (R&D needed)

The ultimate experiment ...

.... is all about background!!

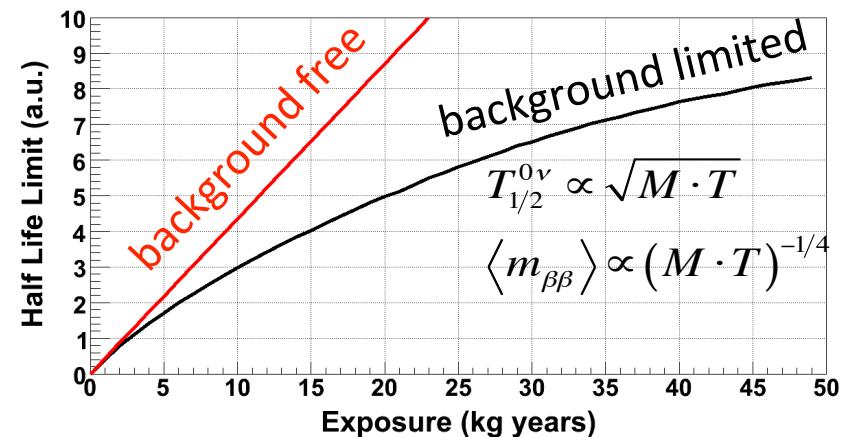
(1) ZERO BACKGROUND FRAMEWORK

$$S_{0\nu} = \ln(2)N_A \frac{\eta \cdot \epsilon}{W} [M \cdot T]$$

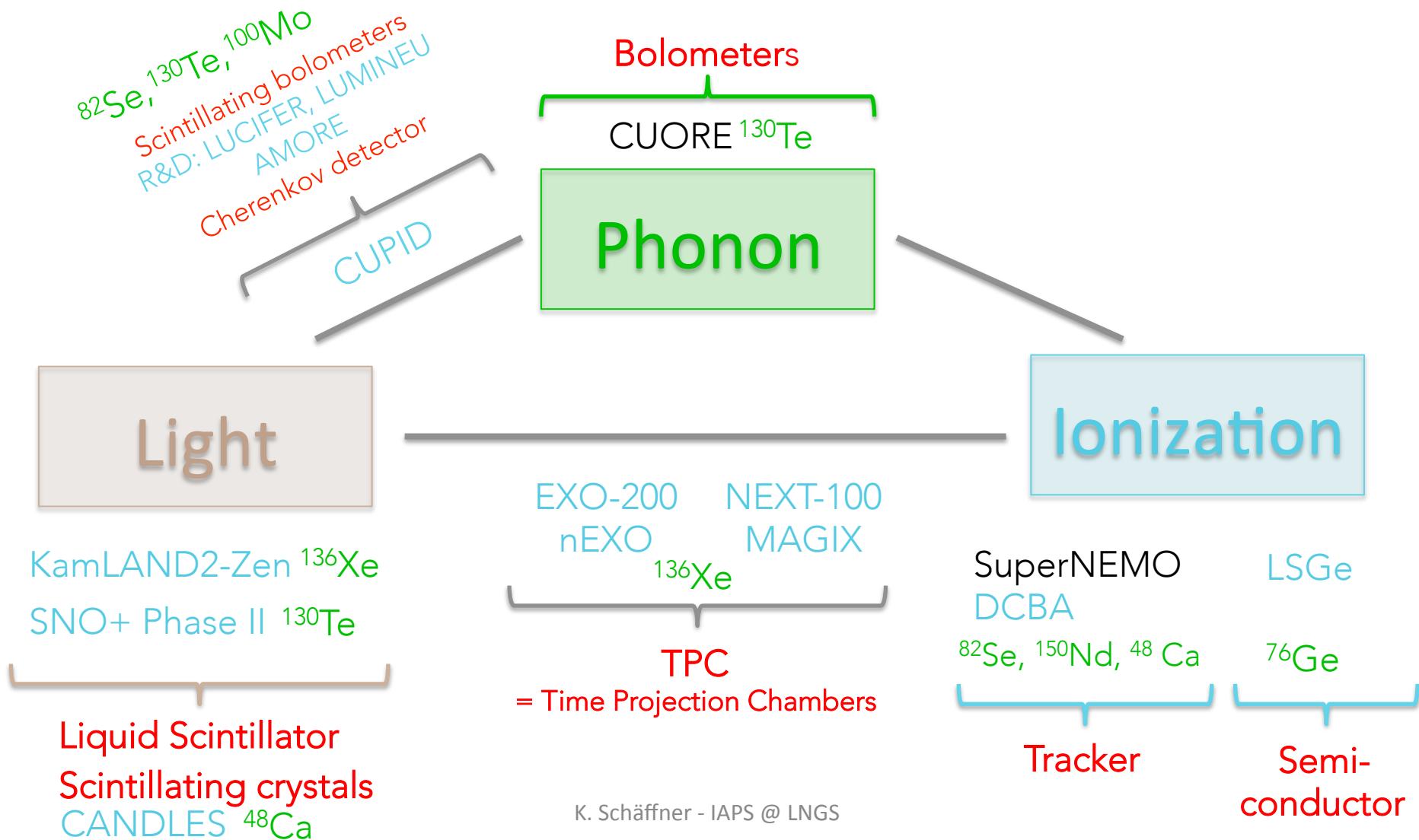
and

(2) increase the number of active DBD nuclei

BUT: translates to a huge necessary budget for enrichment !

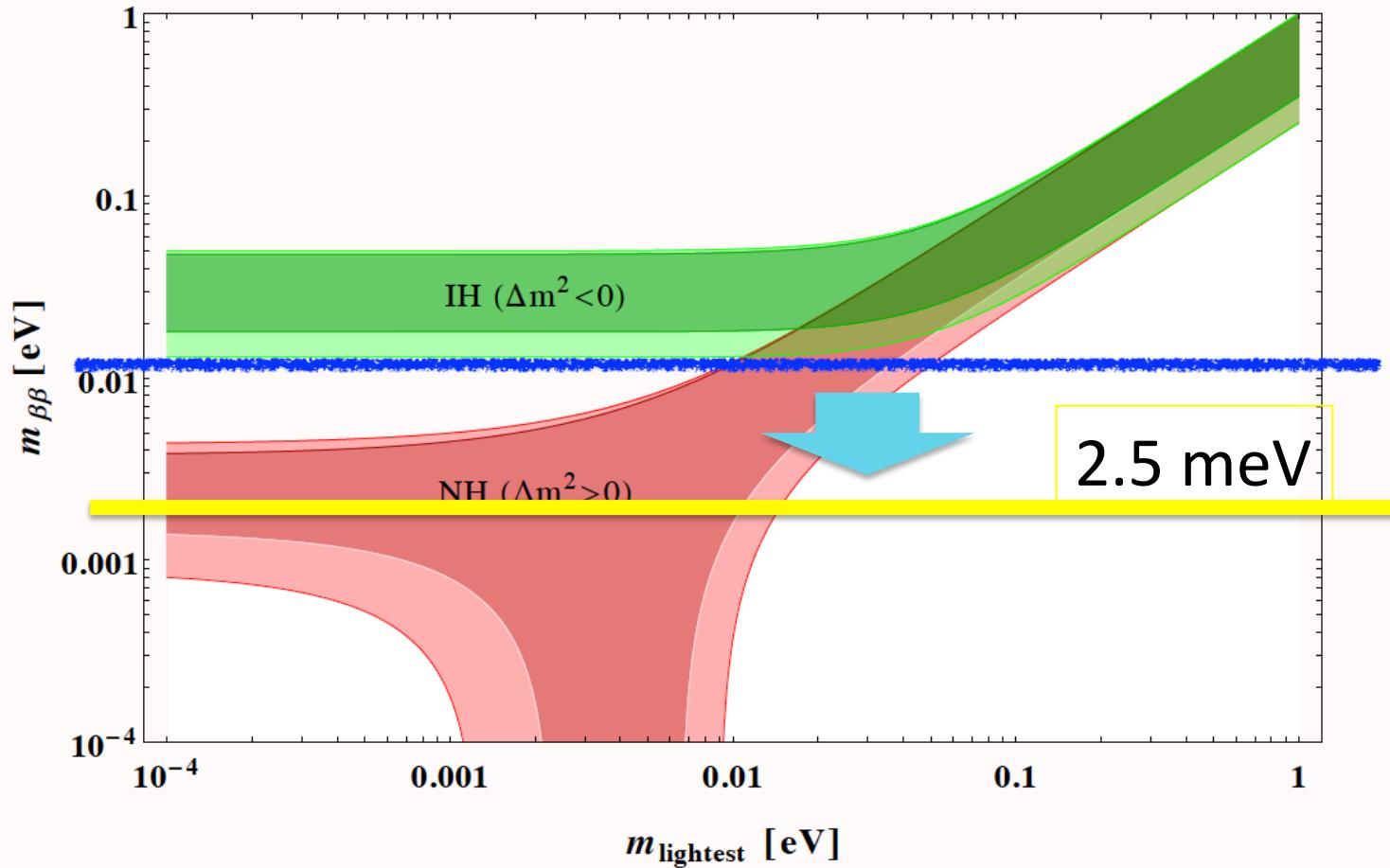


Future generation experiments



far future... assume

Dell'Oro, Marcocci e Vissani, Phys. Rev. D 90, 033005 (2014)



Isotope costs -normal hierarchy region

STEVEN D. BILLER

PHYSICAL REVIEW D **87**, 071301(R) (2013)

TABLE I. Properties of candidate $0\nu\beta\beta$ isotopes.

| Isotope | Q (MeV) | Percent natural abund. | Element cost [5] (\$/kg) | $G^{0\nu}$ ($10^{-14}/\text{yr}$) [6] | $M^{0\nu}$ (avg) [7] | $T_{1/2}^{0\nu}$ for 2.5 meV (10 ²⁹ yrs) | Tons of isotope for 1 ev/yr | Equivalent natural tons | Annual world production [5] (tons) | Natural element cost (\$M) | Enriched at \$20/g (\$M) | $0\nu/2\nu$ rate [2,8] (10 ⁻⁸) |
|-------------------|---------|------------------------|--------------------------|---|----------------------|---|-----------------------------|-------------------------|------------------------------------|----------------------------|--------------------------|--|
| ⁴⁸ Ca | 4.27 | 0.19 | 0.16 | 6.06 | 1.6 | 2.70 | 31.1 | 16380 | 2.4×10^8 | 2.6 | 622 | 0.016 |
| ⁷⁶ Ge | 2.04 | 7.8 | 1650 | 0.57 | 4.8 | 3.18 | 58.2 | 746 | 118 | 1221 | 1164 | 0.55 |
| ⁸² Se | 3.00 | 9.2 | 174 | 2.48 | 4.0 | 1.05 | 20.8 | 225 | 2000 | 39 | 416 | 0.092 |
| ⁹⁶ Zr | 3.35 | 2.8 | 36 | 5.02 | 3.0 | 0.93 | 21.4 | 763 | 1.4×10^6 | 27 | 427 | 0.025 |
| ¹⁰⁰ Mo | 3.04 | 9.6 | 35 | 3.89 | 4.6 | 0.51 | 12.2 | 127 | 2.5×10^5 | 4.4 | 244 | 0.014 |
| ¹¹⁰ Pd | 2.00 | 11.8 | 23000 | 1.18 | 6.0 | 0.98 | 26.0 | 221 | 207 | 5078 | 521 | 0.16 |
| ¹¹⁶ Cd | 2.81 | 7.6 | 2.8 | 4.08 | 3.6 | 0.79 | 22.1 | 290 | 2.2×10^4 | 0.81 | 441 | 0.035 |
| ¹²⁴ Sn | 2.29 | 5.6 | 30 | 2.21 | 3.7 | 1.38 | 41.2 | 736 | 2.5×10^5 | 22 | 825 | 0.072 |
| ¹³⁰ Te | 2.53 | 34.5 | 360 | 3.47 | 4.0 | 0.75 | 23.6 | 68 | ~ 150 | 24 | 471 | 0.92 |
| ¹³⁶ Xe | 2.46 | 8.9 | 1000 | 3.56 | 2.9 | 1.40 | 45.7 | 513 | 50 | 513 | 914 | 1.51 |
| ¹⁵⁰ Nd | 3.37 | 5.6 | 42 | 15.4 | 2.7 | 0.37 | 13.4 | 240 | $\sim 10^4$ | 11 | 269 | 0.024 |

approximate quantities and costs of material that would be required to achieve **a rate of one $0\nu\beta\beta$ decay per year** for an effective Majorana neutrino mass of 2.5 meV.
 (center of the normal hierarchy phase space as the eigenvalue of the smallest neutrino mass eigenstate approaches zero).

My Conclusions

- the search for the neutrinoless DBD addresses **urgent questions** of present day physics
- half-lives are unknown and a discovery might be around the corner
- many different detection technologies were developed within the last decade
- The results from the second generation of present experiments will allow us to:
 - start to scratch the inverted hierarchy region
 - check the discovery potential of the techniques for future generation experiments able to probe the complete inverted hierarchy region
- future generation experiments
 - target mass **of DBD nuclei** ~ 1 tonne
 - increase energy resolution
 - high costs for enrichment and isotope O(100 M Euro) -> **few exp. in the world**
 - in direction of **zero background** experiments (**the hardest task**)
- a **single experiment however will not be enough** for a discovery!!!
- investment in economic enrichment techniques for isotopes of interest ???
- **hope** for a small g_A , otherwise half-lives can change by a factor of 4-16 and a detection becomes very challenging
- a lot hands and brains needed ☺

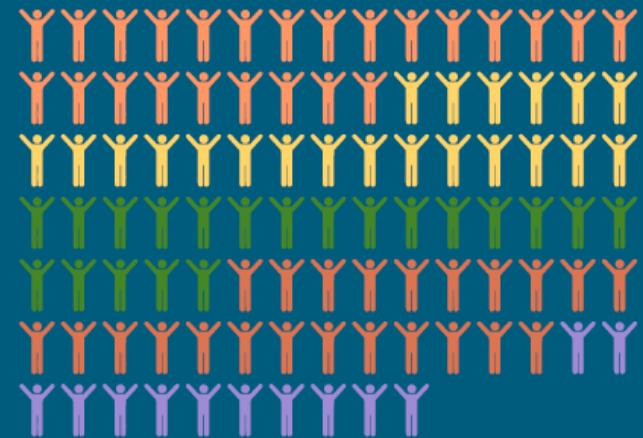
Double Beta Community

A Worldwide Effort



Underground labs
around the world

about 840
physicists worldwide



Germanium Bolometers TPC Scintillator
Tracking

Graphics adopted from L. Winslow: talk @ Neutrino 2014

Thank you for your attention



Additional material



Axial vector coupling constant g_A

The question of whether or not g_A in $0\nu\beta\beta$ is renormalized as much as in $2\nu\beta\beta$ is of much debate. In $2\nu\beta\beta$ only the 1^+ (GT) multipole contributes. In $0\nu\beta\beta$ all multipoles $1^+, 0^+, 2^-, 1^- \dots$ contribute.

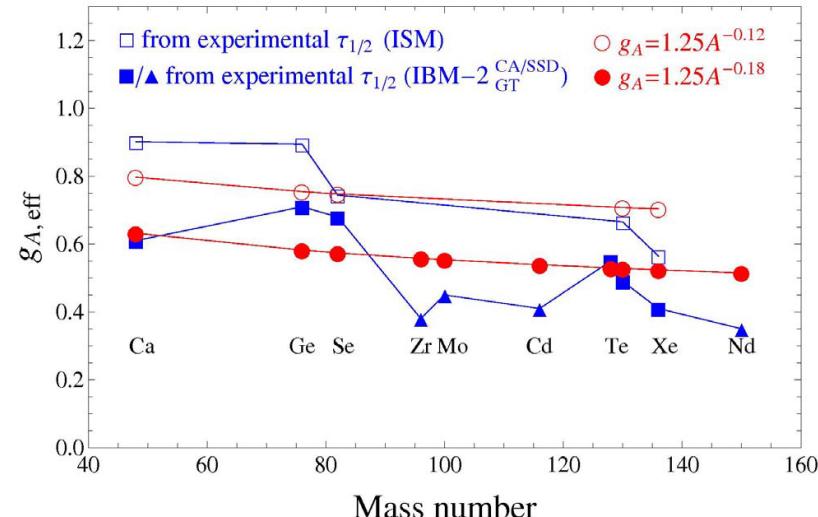
Some of these could be unquenched. However, even in $0\nu\beta\beta$, 1^+ intermediate states dominate. Hence, our current understanding is that g_A is renormalized in $0\nu\beta\beta$ as much as in $2\nu\beta\beta$.

This problem is currently being addressed from various sides. Experimentally by measuring the matrix elements to and from the intermediate odd-odd nucleus in $2\nu\beta\beta$ decay [1]. Theoretically, by using effective field theory (EFT) to estimate the effect of non-nucleonic degrees of freedom (two-body currents) [2].

[1] P. Puppe et al., Phys. Rev. C 86, 044603 (2012).

[2] J. Menendez, D. Gazit, and A. Schwenk, Phys. Rev. Lett. 107, 062501 (2011).

J. Barea, J. Kotila and F. Iachello, Phys. Rev. C 87, 014315 (2013).



Effective Neutrino Mass

The average light neutrino mass can be written as

$$\langle m_\nu \rangle = \left| c_{13}^2 c_{12}^2 m_1 + c_{13}^2 s_{12}^2 m_2 e^{i\varphi_2} + s_{13}^2 m_3 e^{i\varphi_3} \right|$$

$$c_{ij} = \cos \theta_{ij}, s_{ij} = \sin \theta_{ij}, \varphi_{2,3} = [0, 2\pi]$$

$$(m_1^2, m_2^2, m_3^2) = \frac{m_1^2 + m_2^2}{2} + \left(-\frac{\delta m^2}{2}, +\frac{\delta m^2}{2}, \pm \Delta m^2 \right)$$

A fit to oscillation experiments gives [§]

$$\sin^2 \theta_{12} = 0.312, \sin^2 \theta_{13} = 0.016, \sin^2 \theta_{23} = 0.466$$

$$\delta m^2 = 7.67 \times 10^{-5} \text{ eV}^2, \Delta m^2 = 2.39 \times 10^{-3} \text{ eV}^2$$

[§] G.L. Fogli *et al.*, Phys. Rev. D75, 053001(2007); D78, 033010 (2008).

[A recent result from Daya Bay, Phys. Rev. Lett. 108, 171803 (2012) gives $\sin^2 \theta_{13} = 0.024 \pm 0.005$, which slightly modifies the fit.]

$$\langle m_\nu \rangle = \left| \sum U_{ei}^2 m_i \right| = \left| U_{e1}^2 m_1 + U_{e2}^2 m_2 e^{i\alpha_{21}} + U_{e3}^2 m_3 e^{i\alpha_{31}} \right|$$

PMNS mixing matrix :

$$U = \begin{pmatrix} U_{e1} & U_{e2} & U_{e3} \\ U_{\mu 1} & U_{\mu 2} & U_{\mu 3} \\ U_{\tau 1} & U_{\tau 2} & U_{\tau 3} \end{pmatrix} \times \begin{pmatrix} 1 & 0 & 0 \\ 0 & e^{i\alpha_{21}/2} & 0 \\ 0 & 0 & e^{i\alpha_{31}/2} \end{pmatrix}$$

Majorana Phases

Why Ge?

Current experiments

| | | mass [kg]* (total/FV) | FWHM [keV] | background& [cnt/mol yr FWHM] | $T_{1/2}$ limit [10^{25} yr] after 4 yr | $\langle m_{ee} \rangle$ limit [meV] | date |
|-------------|----|--------------------------|---------------|----------------------------------|--|---|-------|
| Gerda II | Ge | 35/27 | 3 | 0.0004 | 15 | 80-190 | -2019 |
| MajoranaD | Ge | 30/24 | 3 | 0.0004 | 15 | 80-190 | -2019 |
| EXO-200 | Xe | 170/80 | 88 | 0.03 | 6 | 80-220 | -2019 |
| Kamland-Zen | Xe | 383/88 (600/?) | 250 | 0.03 | 20 | 44-120 | -2018 |
| NEXT | Xe | 100/80 | 17 | 0.0036 | 6 | 100-200 | -2020 |
| Cuore | Te | 600/206 | 5 | 0.02 | 9 | 50-200 | -2019 |
| SNO+ | Te | 2340/160 | 270 | 0.02 | 9 | 50-200 | -2020 |

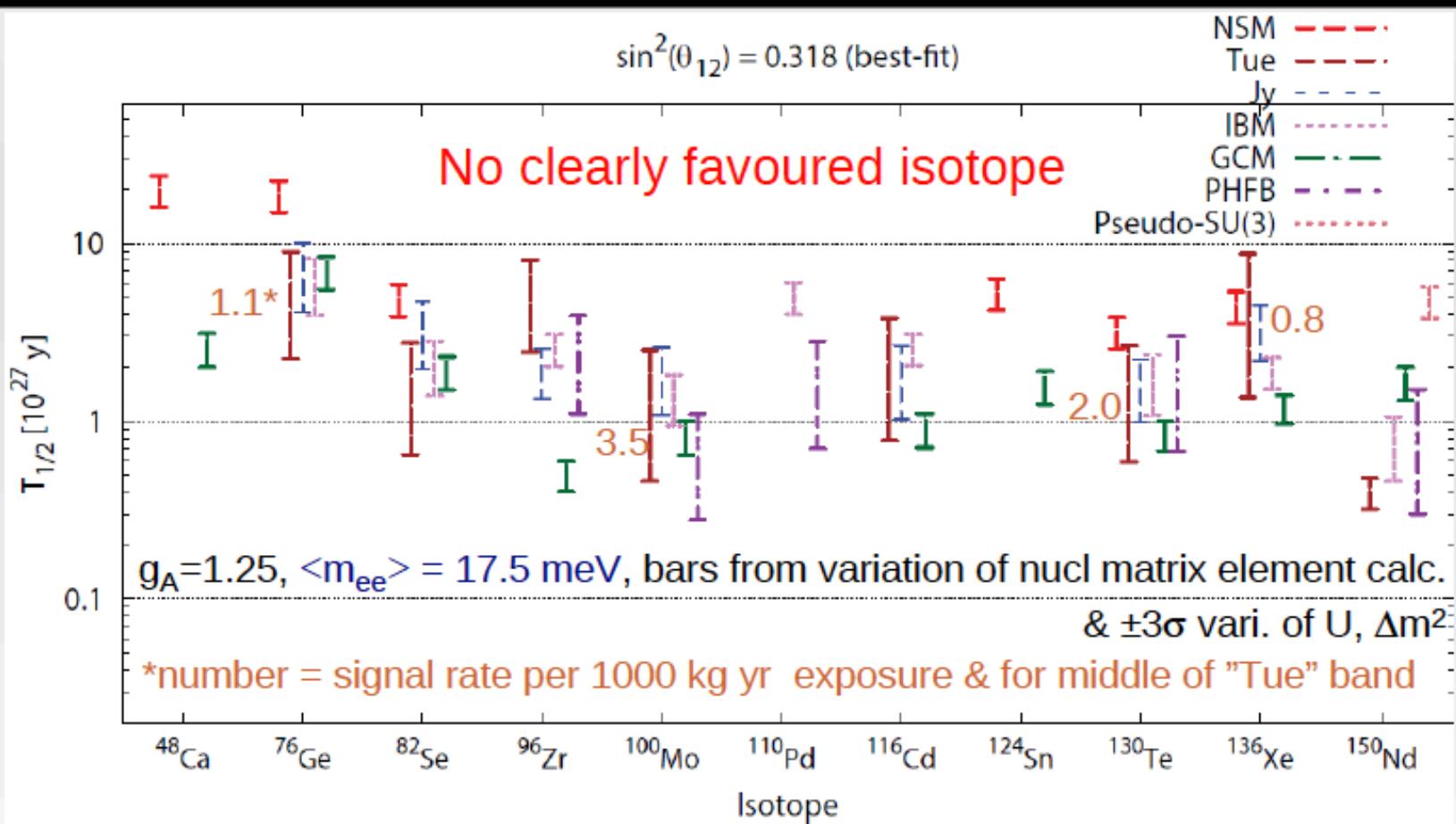
* total= element mass, FV= $0\nu\beta\beta$ isotope mass in fiducial volume (incl enrichment fraction)

& mol of $0\nu\beta\beta$ isotope in active volume and divided by $0\nu\beta\beta$ efficiency

Note: values are design numbers except for EXO-200 and Kamland-Zen

Ge experiments have lowest background → similar sensitivity despite small mass

Expected $T_{1/2}$ for different matrix elements



taken from DOE Nuclear Science Advisory Committee report on $0\nu\beta\beta$ (24 April 2014)
adopted from A. Dueck, W. Rodejohann and K. Zuber, Phys. Rev. D83 (2011) 113010

warning from theory colleague: unclear whether light Majorana neutrino exchange is dominating
& other ν properties unknown → motivation to cover inverted mass hierarchy is artificial

Spectrum of KamLAND-Zen

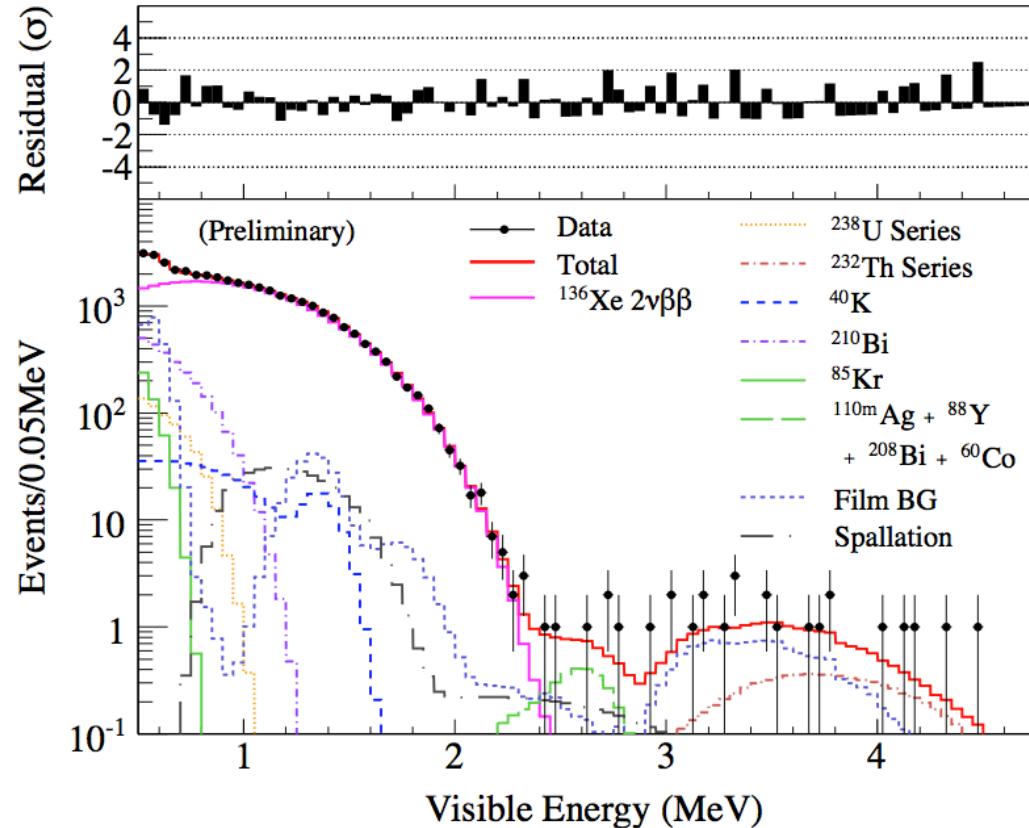


FIGURE 3. Preliminary energy spectrum of selected $\beta\beta$ candidates within 1.0 m fiducial radius is shown together with the best-fit backgrounds, with the $2\nu\beta\beta$ decay fit. The residuals from the best-fit are shown in the upper panel.

$\langle m_{\beta\beta} \rangle = 12 \text{ meV}$ and zero bck

- Objective: $\langle m_{\beta\beta} \rangle = 12 \text{ meV}$
- sensitivity of 1 sigma
- Nuclear matrix elements from IBM-2
- zero background condition (i.e. $M \cdot T \cdot B \cdot \Delta E \lesssim 1$)
- measuring time $T = 5 \text{ y}$
- enrichment level 90%
- actual costs for enrichment
- detector efficiency (fiducial volume) as present
- energy resolution as present
- $g_A = 1.269$

| | T1/2 0v [anni] | Massa isotopo [ton] | Efficienza (volume fiduciale) | Massa totale rivelatore [ton] | a.i. [%] | Costo isotopo/kg [k€] | Costo isotopo [M€] | $B^* \Delta E$ [c/kg/y] | ΔE FWHM [keV] | Fondo [c/keV/kg/y] | Q value [keV] | T1/2 2v [anni] | Rapporto eventi 0v/2v |
|----------------------------|----------------|---------------------|-------------------------------|-------------------------------|----------|-----------------------|--------------------|-------------------------|-----------------------|--------------------|---------------|----------------|-----------------------|
| 76Ge GERDA | 1,0E+28 | 0,365 | 0,82 | 0,49 | 7,8 | 70,00 | 34,64 | 5,5E-04 | 3 | 1,8E-04 | 2039 | 1,8E+21 | 6,3E+08 |
| 82Se Lucifer | 3,7E+27 | 0,145 | 0,80 | 0,36 | 9,2 | 70,00 | 14,11 | 1,4E-03 | 10 | 1,4E-04 | 2996 | 9,2E+19 | 4,4E+05 |
| 100Mo (ZnMO ₄) | 3,2E+27 | 0,151 | 0,80 | 0,48 | 7,6 | 100,00 | 20,97 | 1,3E-03 | 9 | 1,5E-04 | 3034 | 7,1E+18 | 8,0E+04 |
| 116Cd (CdWO ₄) | 5,4E+27 | 0,301 | 0,80 | 1,31 | 9,6 | 150,00 | 62,64 | 6,7E-04 | 6 | 1,1E-04 | 2814 | 2,8E+19 | 1,4E+06 |
| 130Te CUORE | 3,0E+27 | 0,189 | 0,87 | 0,30 | 34,2 | 13,00 | 3,14 | 1,1E-03 | 5 | 2,1E-04 | 2527 | 6,8E+20 | 1,1E+08 |
| 130Te SNO+ | 3,0E+27 | 0,189 | 0,20 | 350,61 | 34,2 | 13,00 | 13,67 | 1,1E-03 | 270 | 3,9E-06 | 2527 | 6,8E+20 | 4,3E-03 |
| 136Xe EXO | 4,3E+27 | 0,281 | 0,50 | 0,63 | 8,9 | 8,00 | 5,00 | 7,1E-04 | 58 | 1,2E-05 | 2458 | 2,1E+21 | 8,4E+01 |
| 136Xe Kam-Zen | 4,3E+27 | 0,281 | 0,30 | 42,72 | 8,9 | 8,00 | 8,34 | 7,1E-04 | 250 | 2,8E-06 | 2458 | 2,1E+21 | 1,3E-02 |
| 136Xe NEXT | 4,3E+27 | 0,281 | 0,30 | 1,04 | 8,9 | 8,00 | 8,34 | 7,1E-04 | 15 | 4,7E-05 | 2458 | 2,1E+21 | 2,8E+05 |