Hadron Reaction and Spectroscopy Studies at JPAC

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Evidence of new hadrons

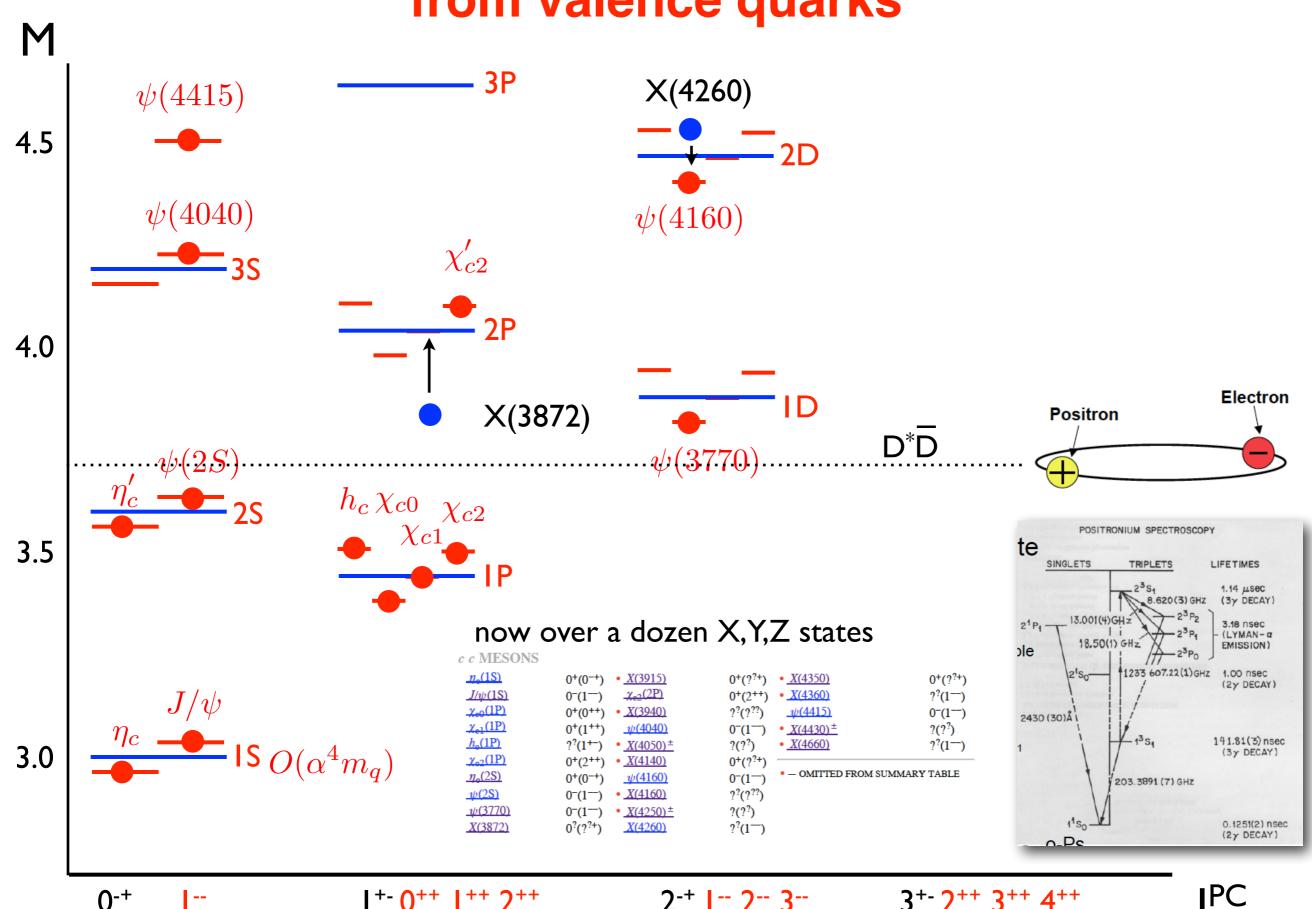
Role of reaction theory

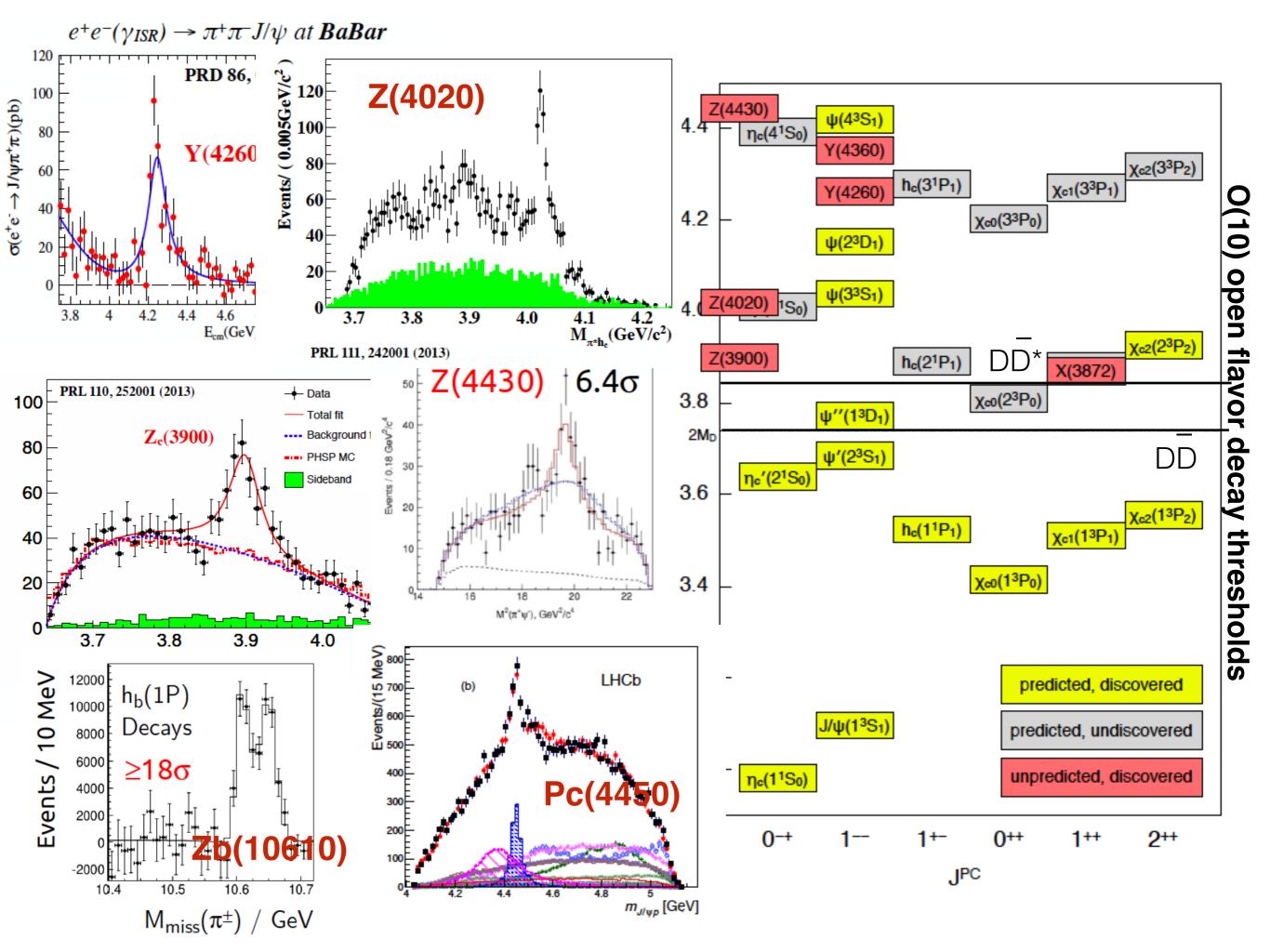
The XYZ's and exotics

Joint Physics Analysis
Center

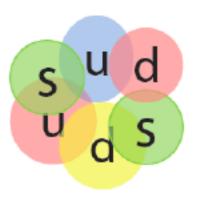
supported by US DOE, US NSF

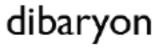
Long time ago hadrons were made from valence quarks

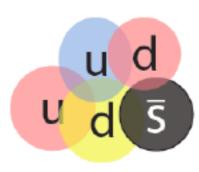




before we can address the following question...







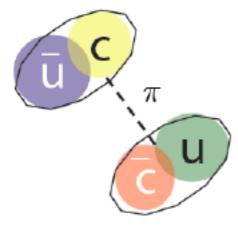
pentaquark



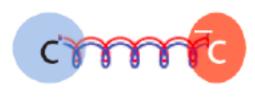
glueball



diquark + di-antiquark

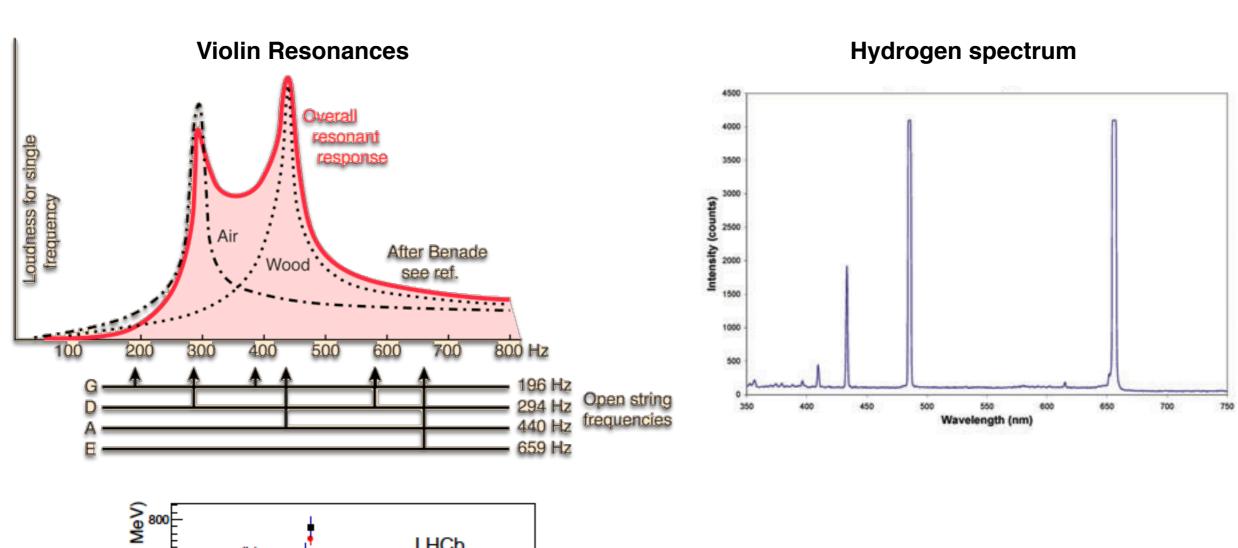


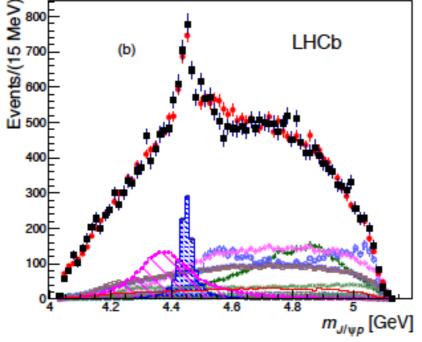
dimeson molecule



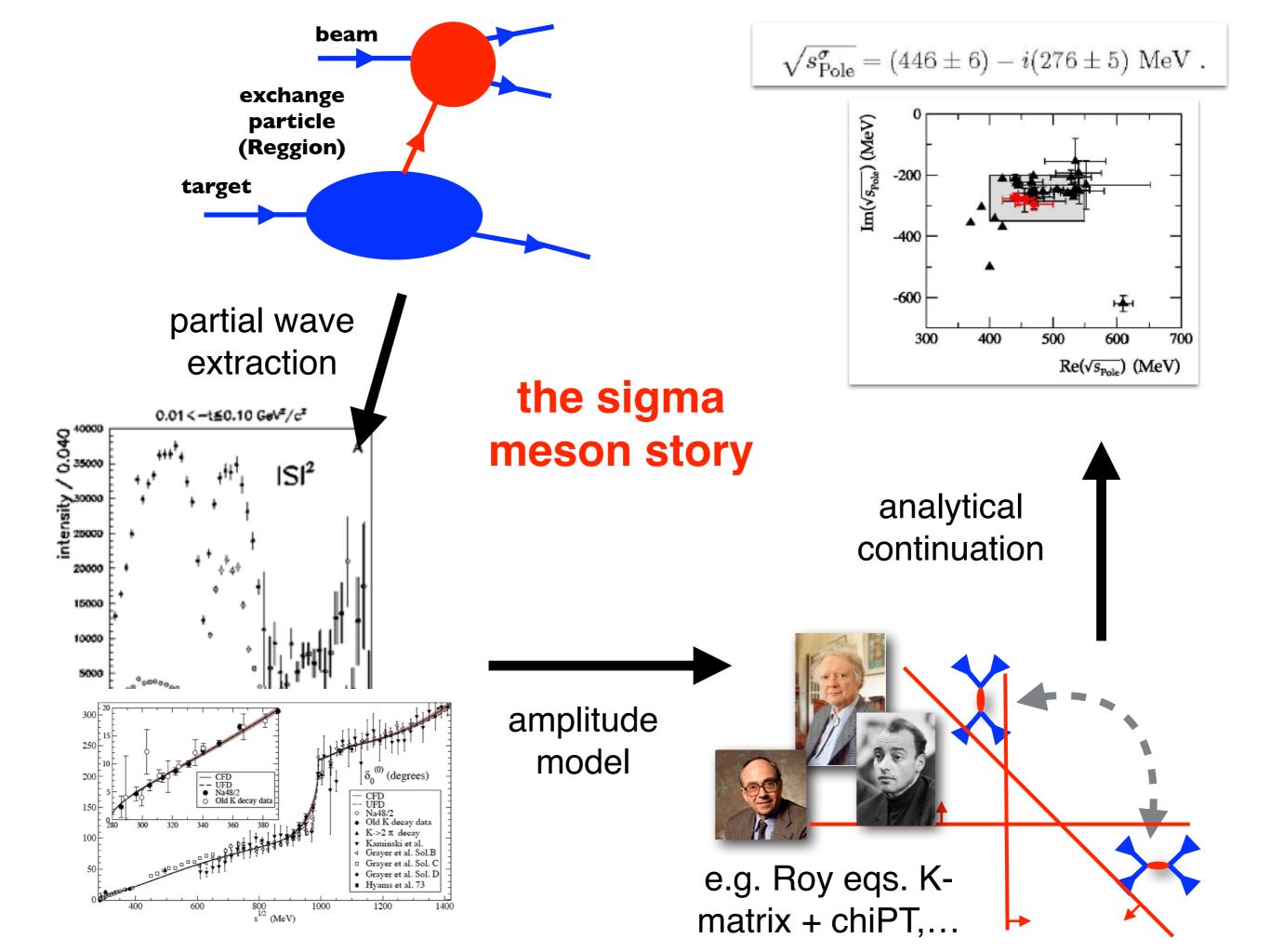
 $q \, \bar{q} \, g$ hybrid

...we need to know how to interpret "peaks"





... but state of the art analyst is based on analogy with a "very simple world"



Exploring the σ pole

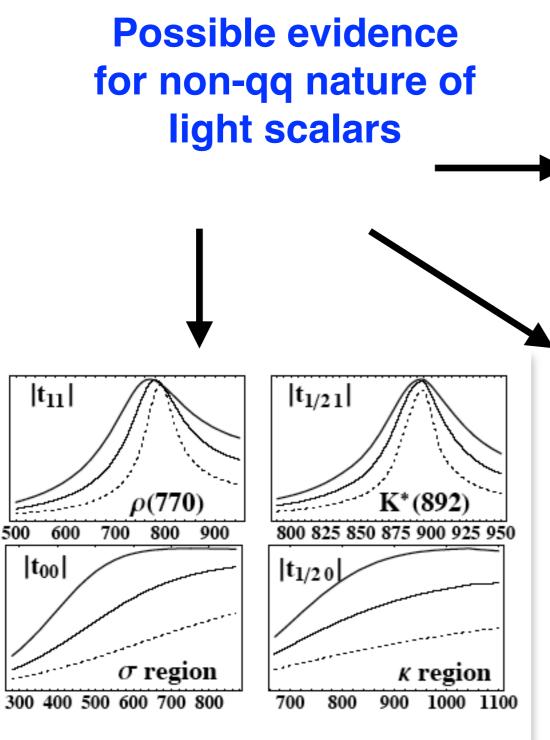
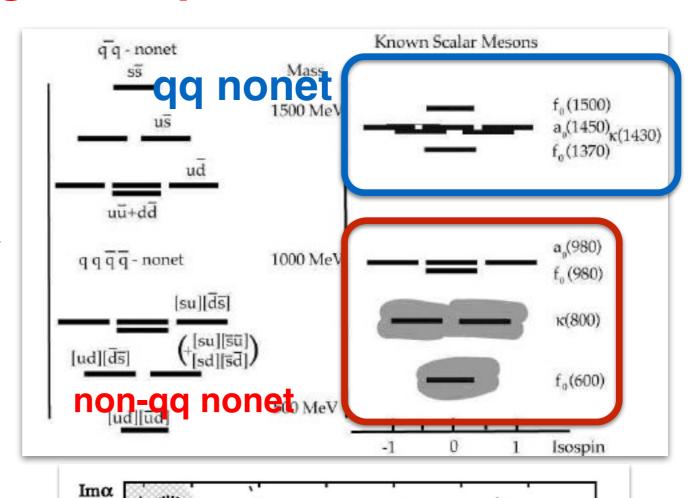


FIG. 1: Modulus of amplitudes in different meson-meson channels for $N_c = 3$ (thick line) $N_c = 5$ (thin continuous line) and $N_c = 10$ (thin dotted line), scaled at $\mu = 770 \,\mathrm{MeV}$.

Pelaez

Yukawa potential 0.3 0.2 0.1 ordinary trajectory 0.5 non-ordinary σ Reα

0.4



 $\alpha_{\sigma}(s)$, s<2 GeV²

 $\alpha_p(s)$, s<2 GeV²

 $\alpha_{\sigma}(s)$, s>2 GeV²

 $\alpha_p(s)$, s>2 GeV²

Striking similarity with

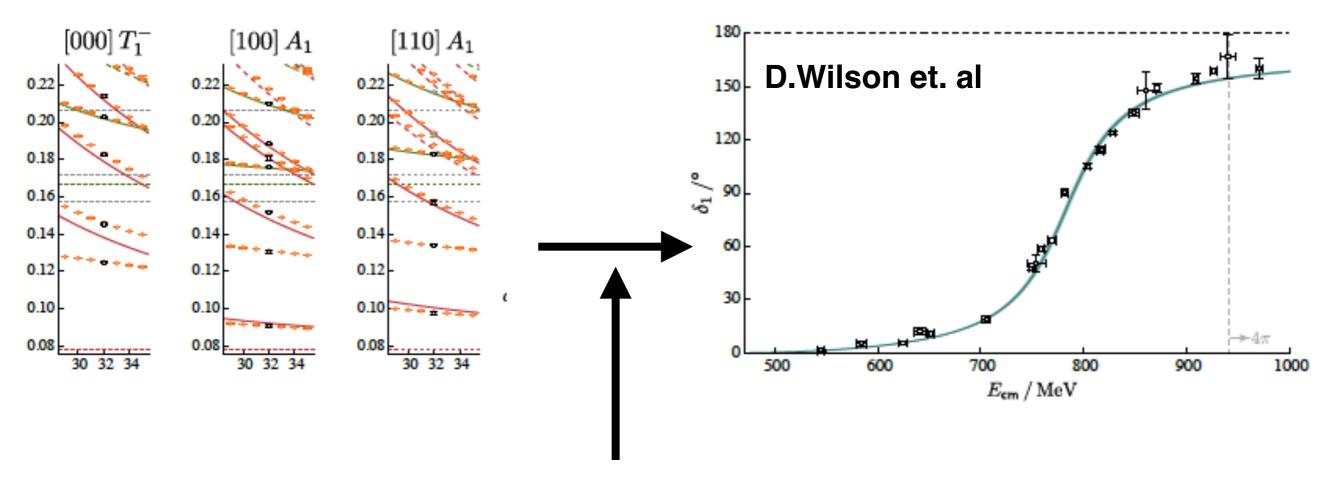
QCD on the Lattice: simulated scattering experiment

(known kinematical function)

$$Z(E_i) = T(E_i)$$

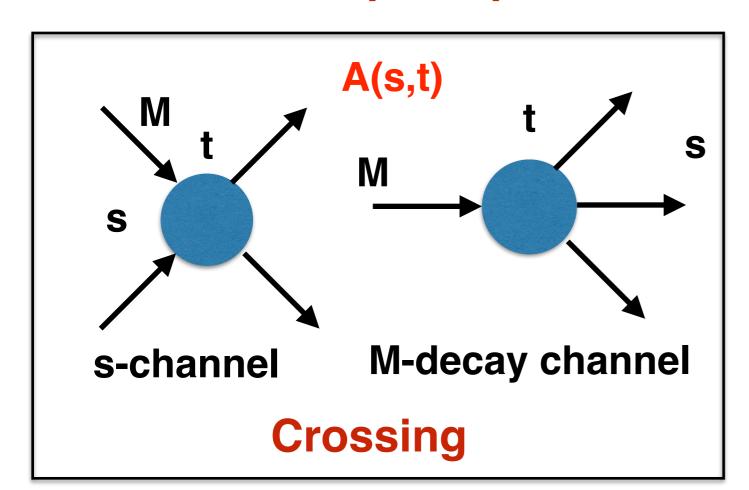
(infinite volume amplitude)

 E_i = discrete energy spectrum of states in the lattice

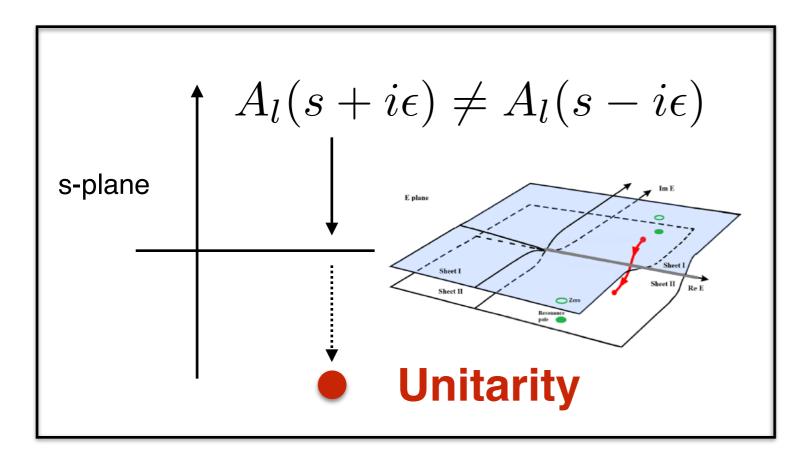


in general "solution" of the Lusher condition requires an analytical model for T

S-matrix principles: Crossing, Analyticity, Unitarity

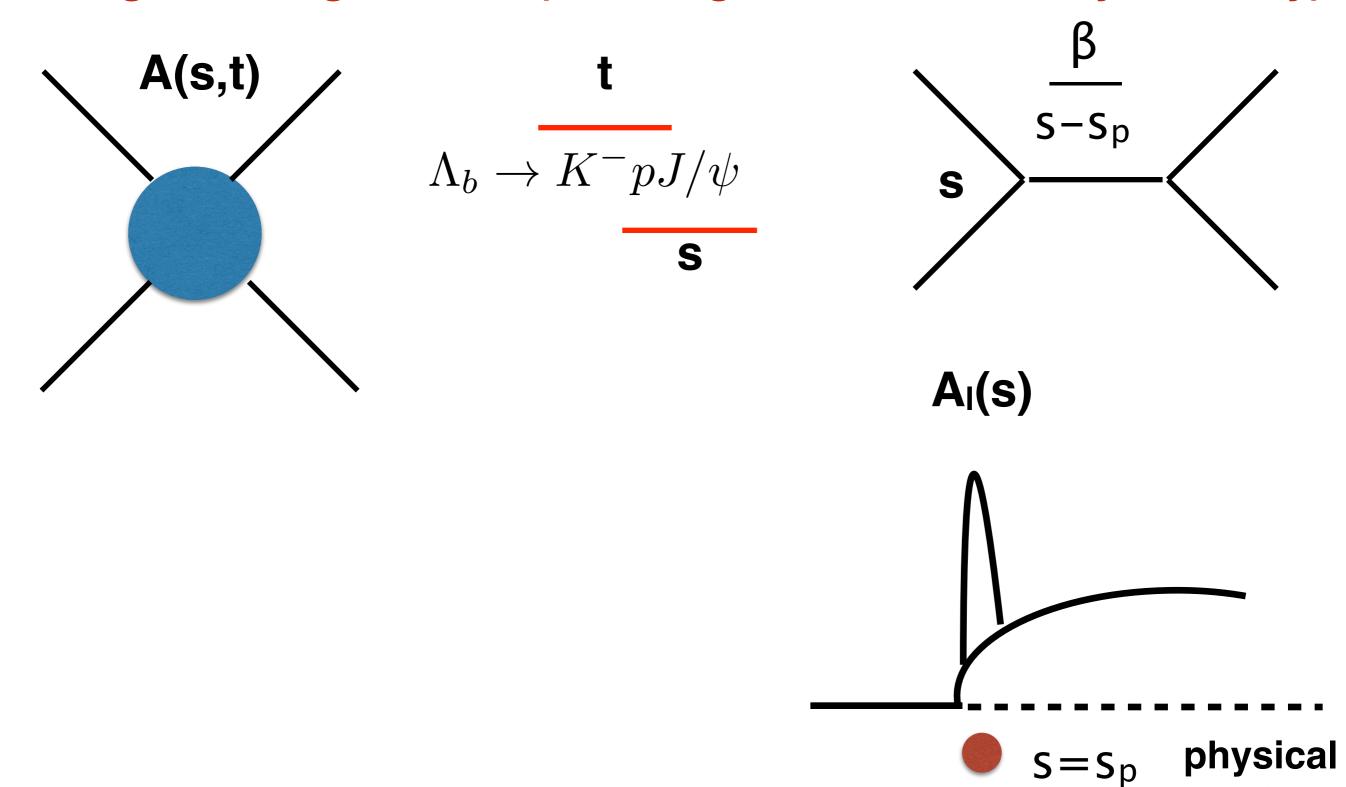


$$A(s,t) = \sum_l A_l(s) P_l(z_s)$$
Analyticity
 $A_l(s) = \lim_{\epsilon o 0} A_l(s+i\epsilon)$



bumps/peaks on the real axis (experiment) come from singularities in the complex domain.

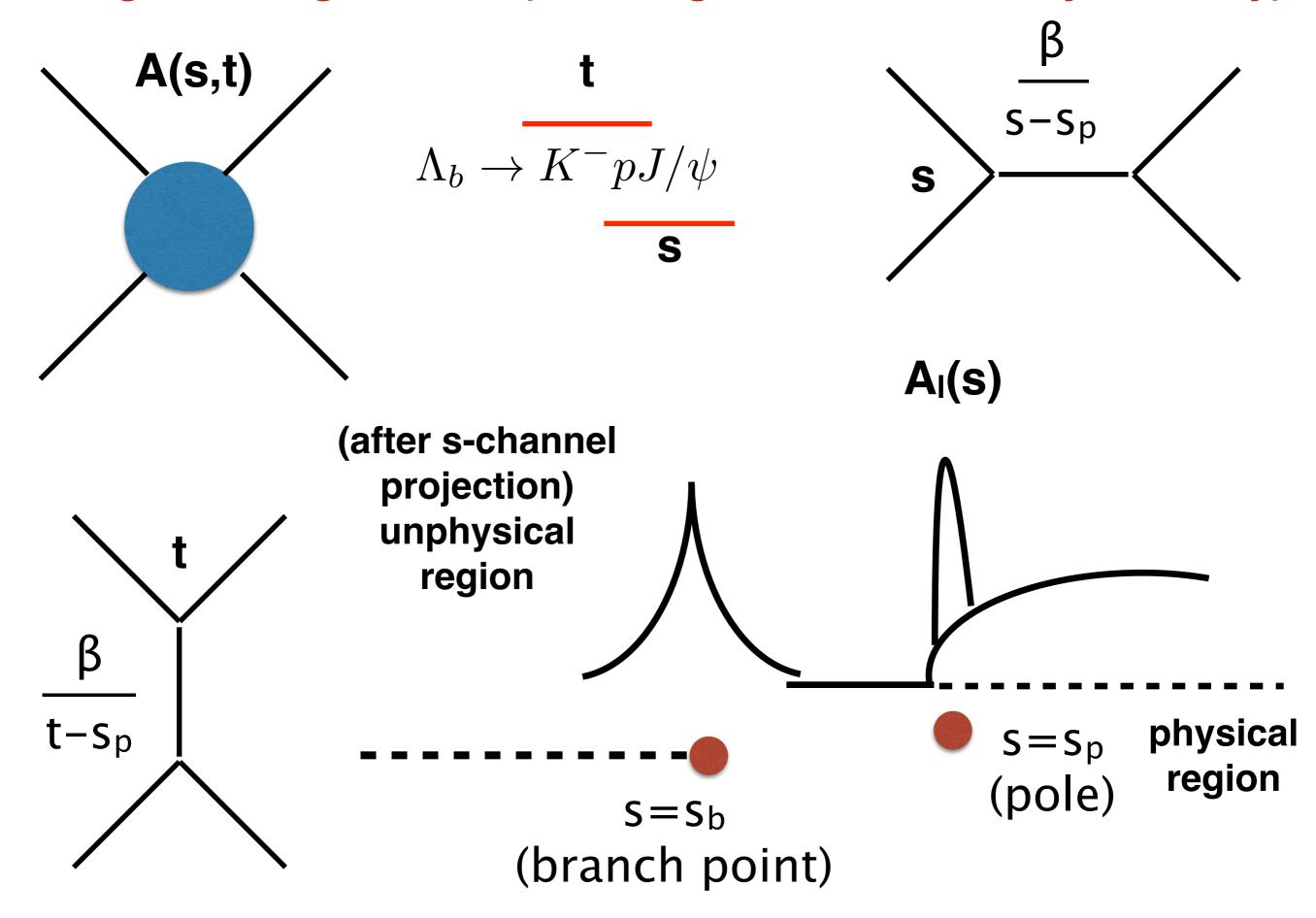
Origin of singularities (exchanges constrained by unitarity)



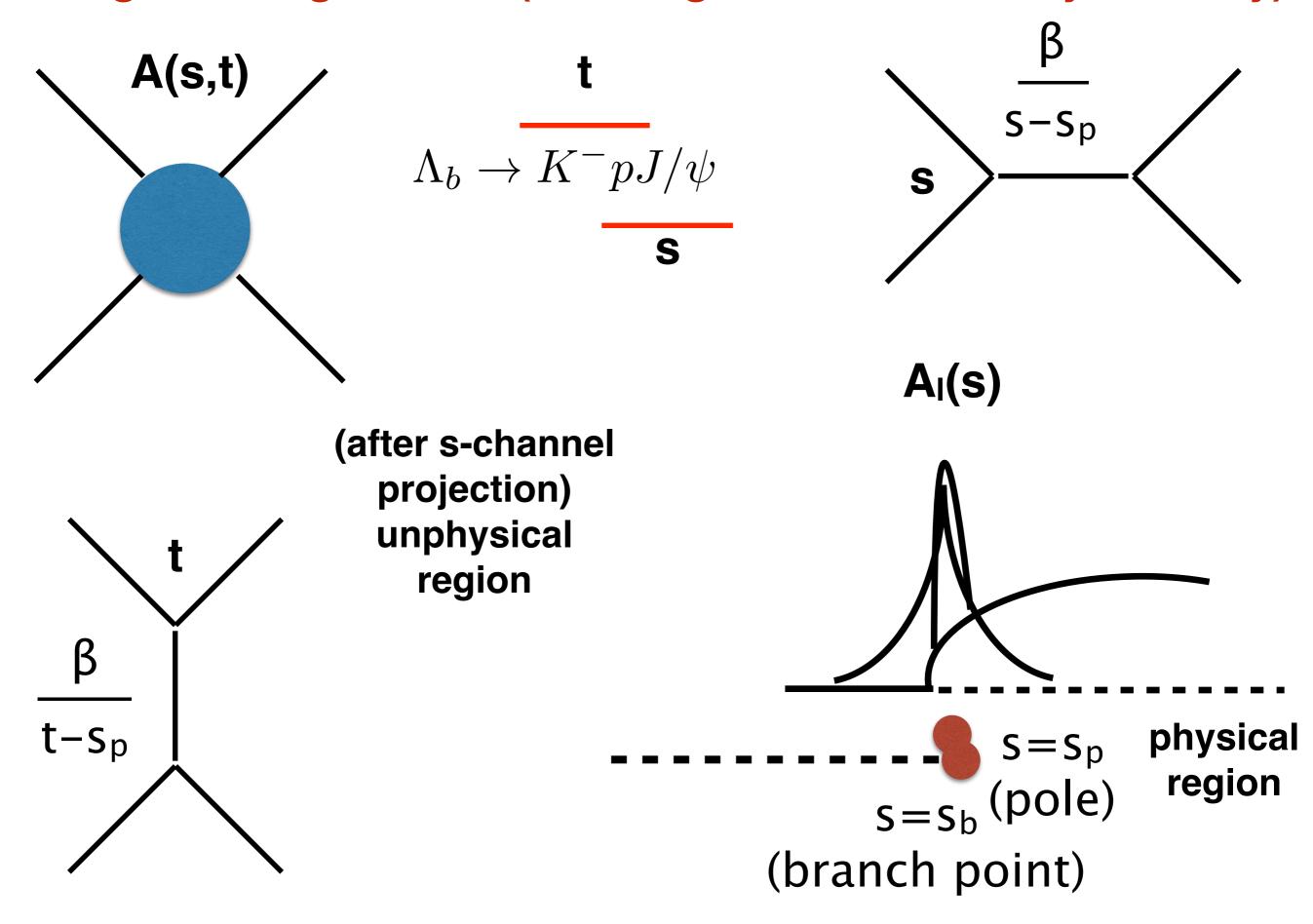
region

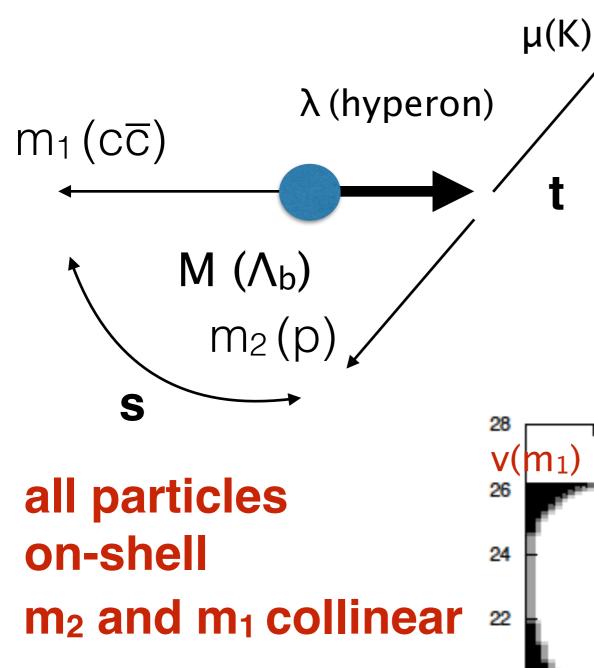
(pole)

Origin of singularities (exchanges constrained by unitarity)



Origin of singularities (exchanges constrained by unitarity)





 $v(m_2) > v(m_1)$

18

Coleman-Norton theorem

t-channel resonance can produce s-channel "band" if:

$$\lambda_{\text{min}}^2 = t \otimes \text{Streshold}$$

$$\text{and } u_{\text{max}}$$

$$\lambda \text{ off-shell}$$

$$v(m_1) = v(m_2)$$

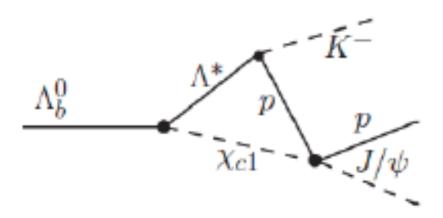
$$v(m_1) > v(m_1) = 0$$

$$v(m_1) > v(m_2)$$

$$v(m_1) = 0$$

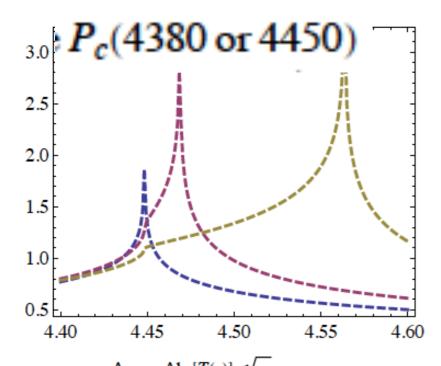
$$\lambda_{\text{max}}^2 = (M-m_1)^2 = t_{\text{max}}$$

The key to the YZ phenomena are the many nearby channels



$$M_{\Lambda_b^0} = 5.6195$$
, $\mu_{K^-} = 0.4937$, $m_1 = m_{\chi_{c1}} = 3.510$, $m_2 = m_p = 0.93827$
 $\lambda = m_{\Lambda^*} = 1.89$ (they take)

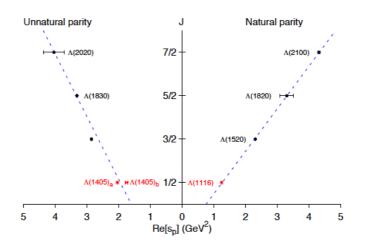
Coleman-Norton requires

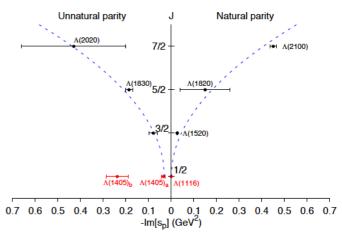


Axes: Abs[T(s)], \sqrt{s} Lines: blue ($\lambda = 1.89 \text{ GeV}$), red ($\lambda = 1.99 \text{ GeV}$), yellow ($\lambda = 2.09 \text{ GeV}$)

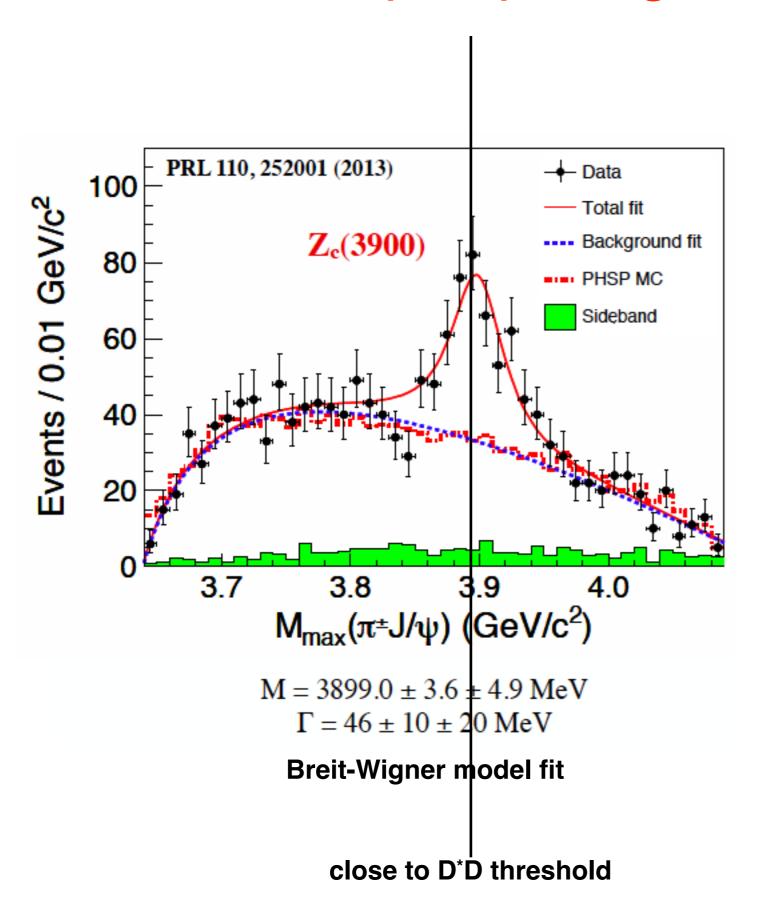
$$1.89 < \lambda < 2.11 \text{ GeV}$$

 $4.45 < \sqrt{s_{\text{neak}}} < 4.65 \text{ GeV}$



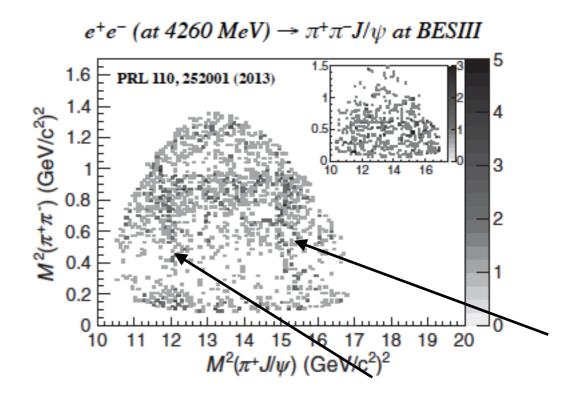


Z_c(3900) Charged charmonium?

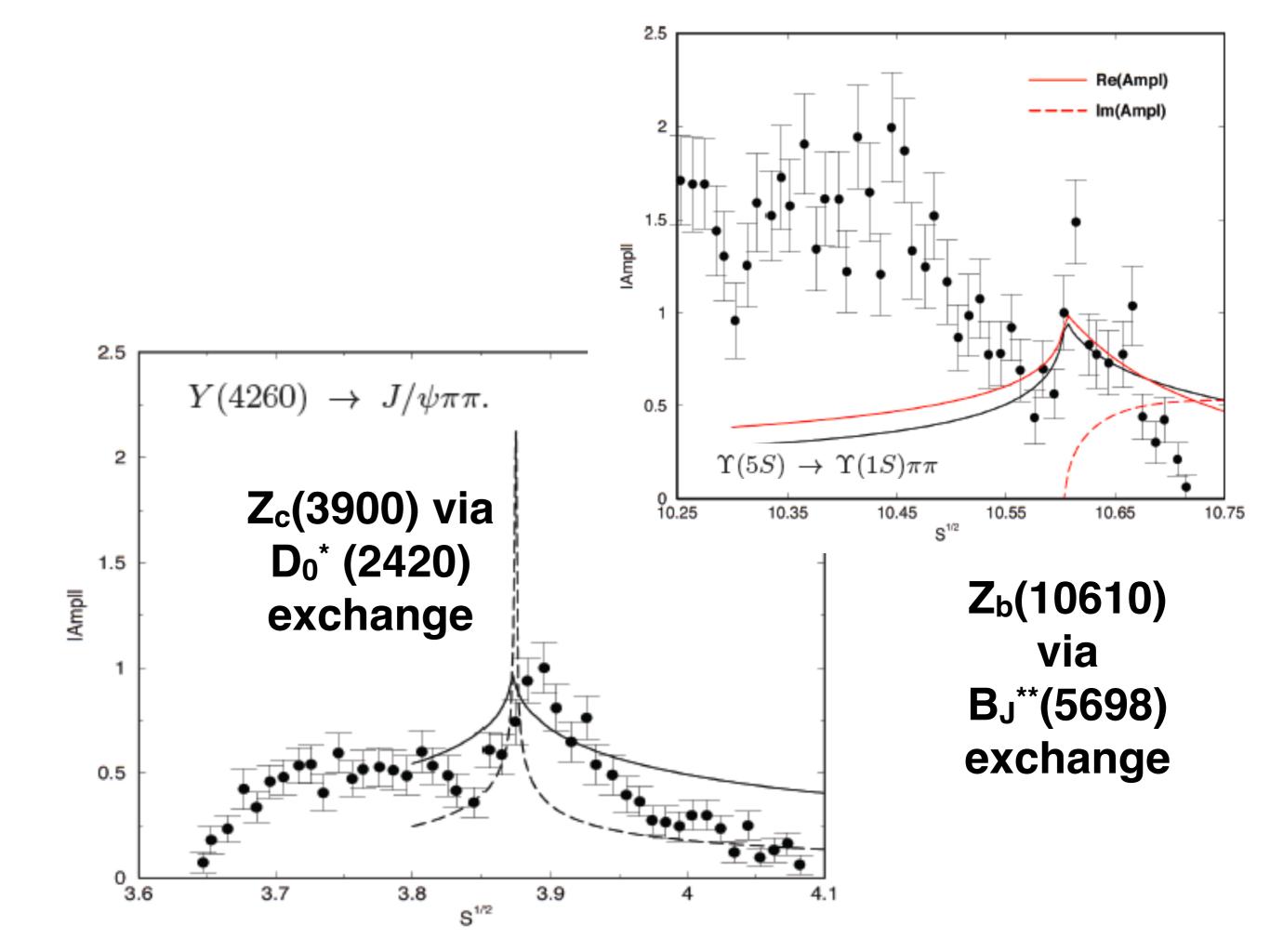


$$e^+e^- \rightarrow Y(4260)$$

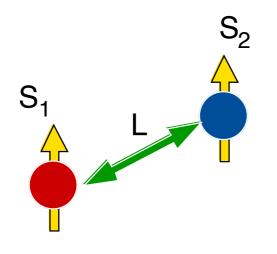
 $\rightarrow \pi^+ Z_c^- (3900)$
 $\rightarrow \pi^+ \pi^- J/\psi$



coupling to DD*

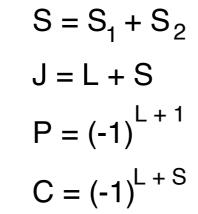


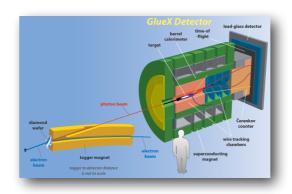
Gluon excitations



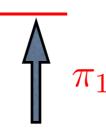




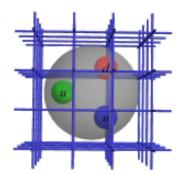




Mesons with $J^{PC} = 0^{--}, 0^{+-}, 1^{-+}, 2^{+-}$: Exotic Quantum Numbers

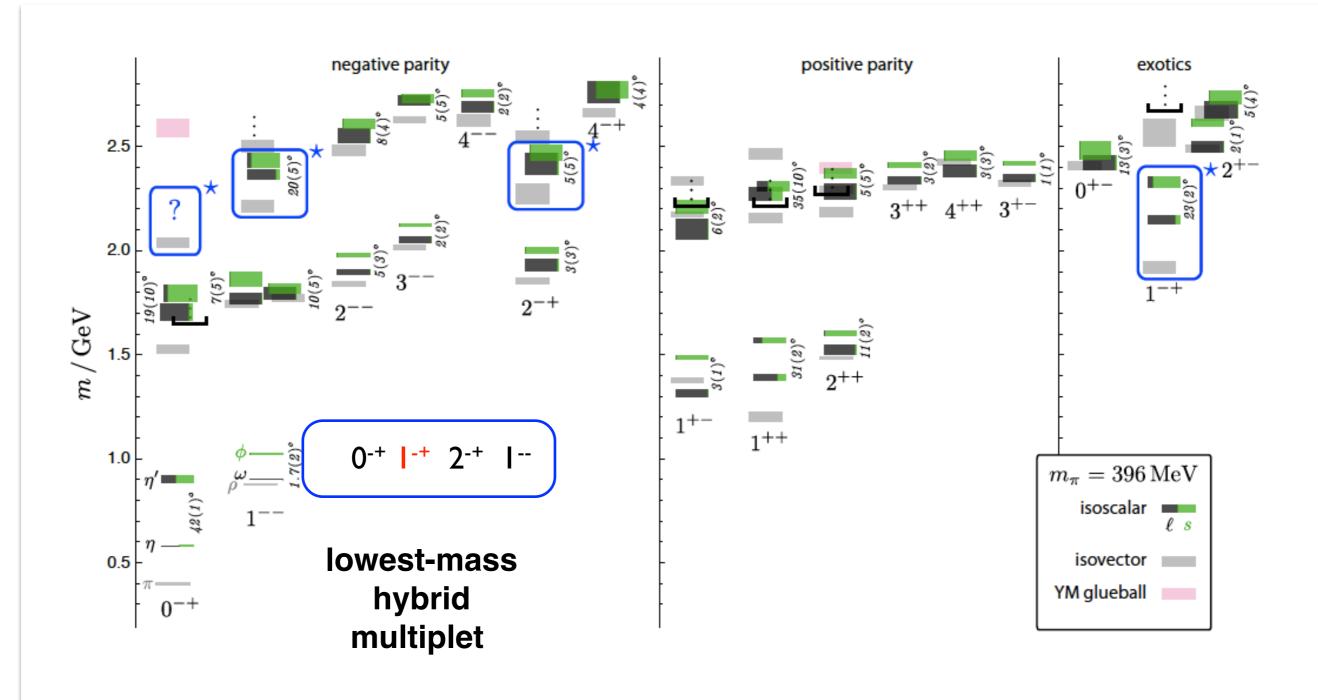


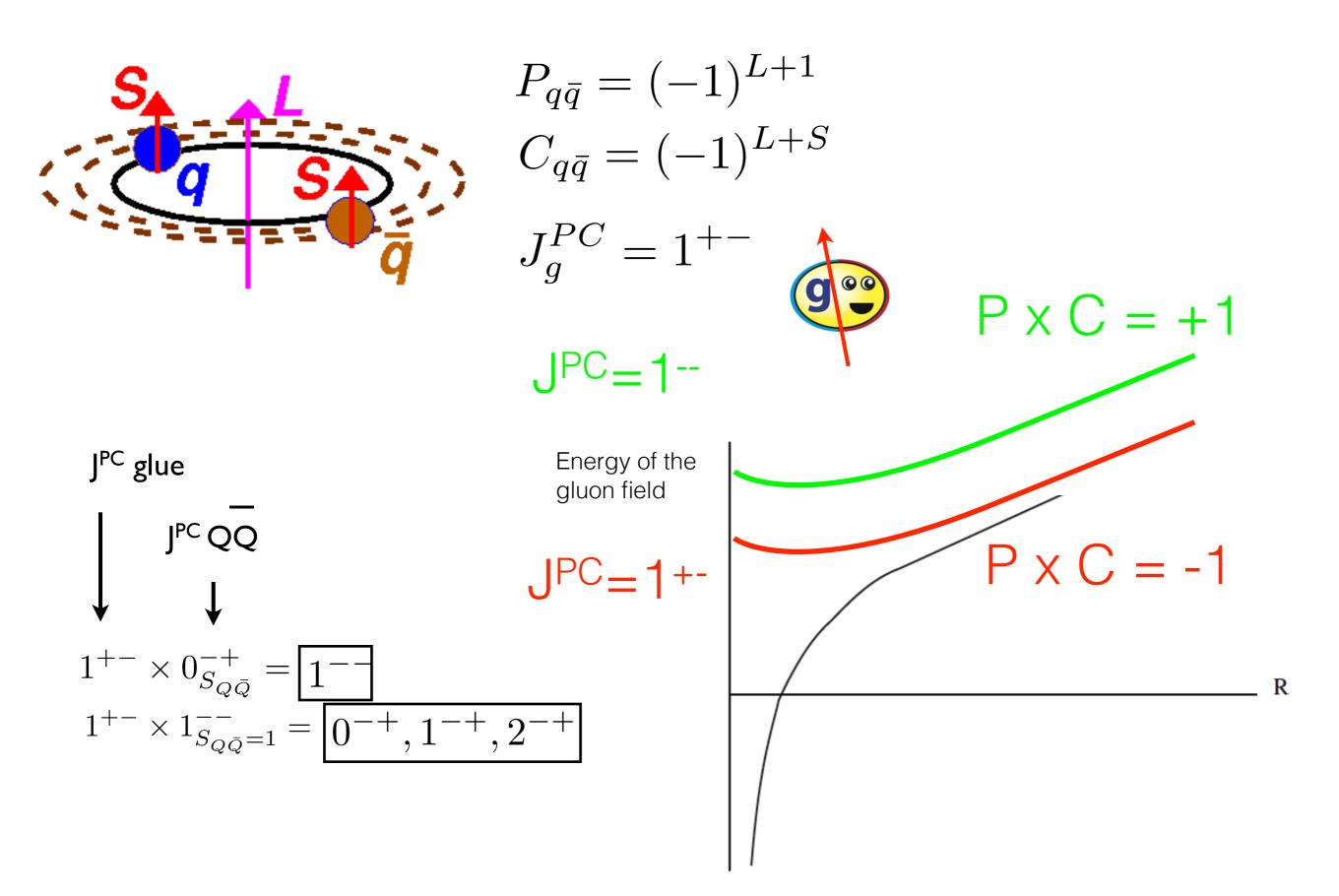
Expected to have very similar properties to ordinary QQ mesons



J. Dudek et al.

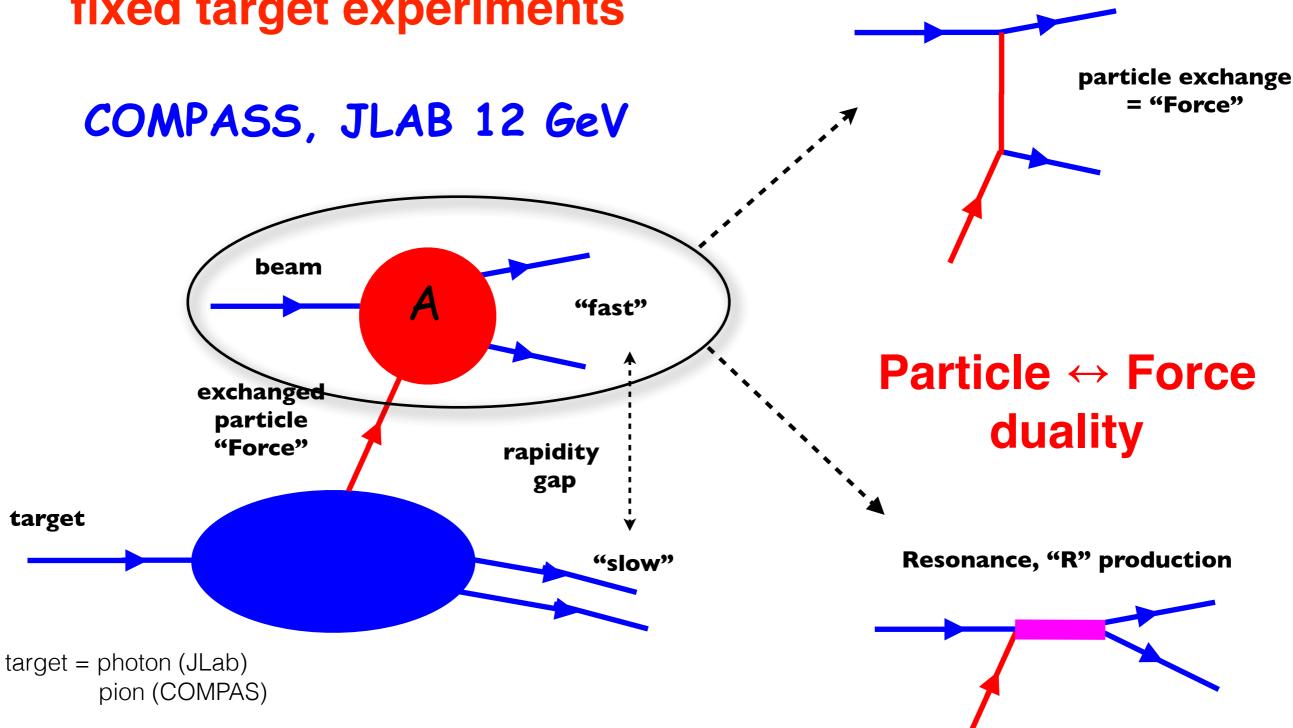
same pattern in ss, cc hybrid interpretation of the Y(4260)





Prediction of Coulomb gauge QCD

Hunting for Resonances in fixed target experiments

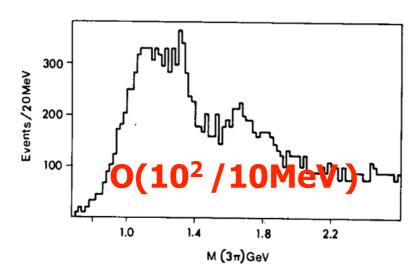


diffractive-dissociation

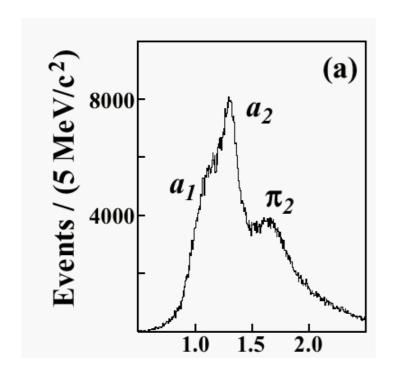
Evolution in statistics

$$\pi^- p \rightarrow \pi^- \pi^+ \pi^- p$$

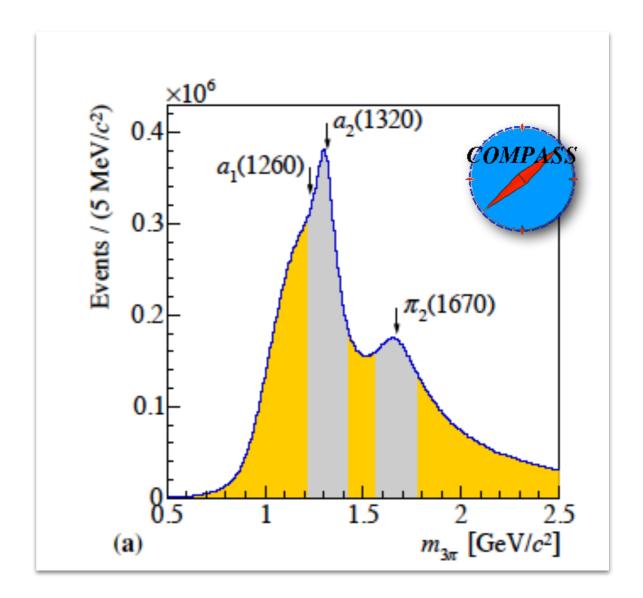
CERN ca. 1970



BNL (E852) ca 1995

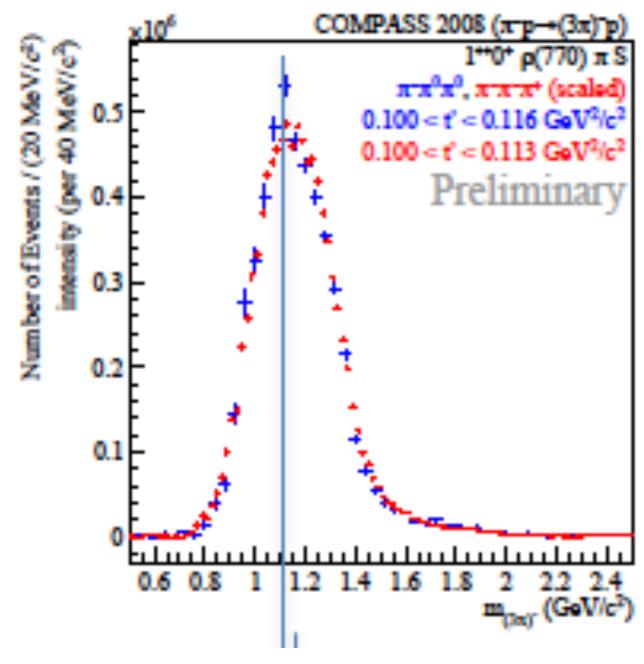


E852 (Full sample) $\sim 10^5/10 \text{MeV}$



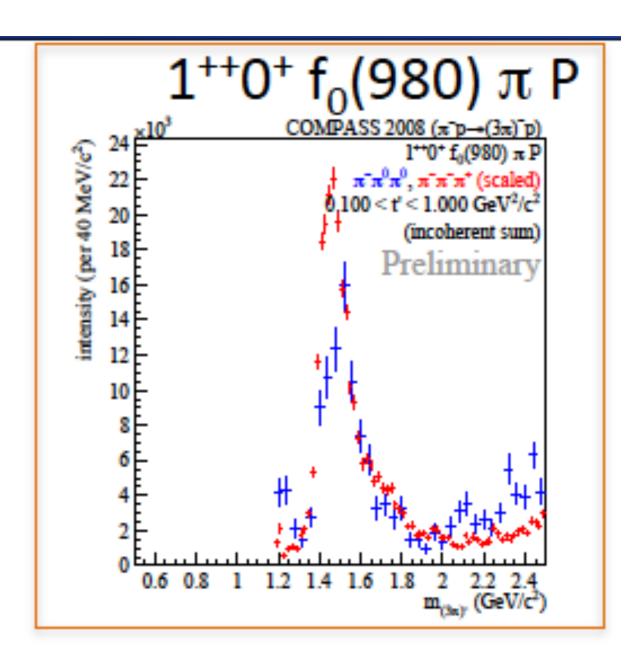
 $O(10^6/10MeV)$

COMPASS's a_1 meson(s) $J^{PC} = 1^{++}$ (L=1,S=1)



"Old" a1 in ρπ S-wave

New a1 in f0 π P-wave



Two Breit-Wigner resonances (COMPASS)

$$\Delta A_i(s) = \sum_{j=f_0\pi,
ho\pi} T^*_{ij}(s)
ho_j(s) A_j(s)$$

particle exchange = "Force"

 $A_i(s) = \sum_{j=f_0\pi,
ho\pi} T_{ij}(s) P_j(s)$

Resonance, "R" production

One resonance underlies

both peaks

1.5e+06

1e+06

5e+05

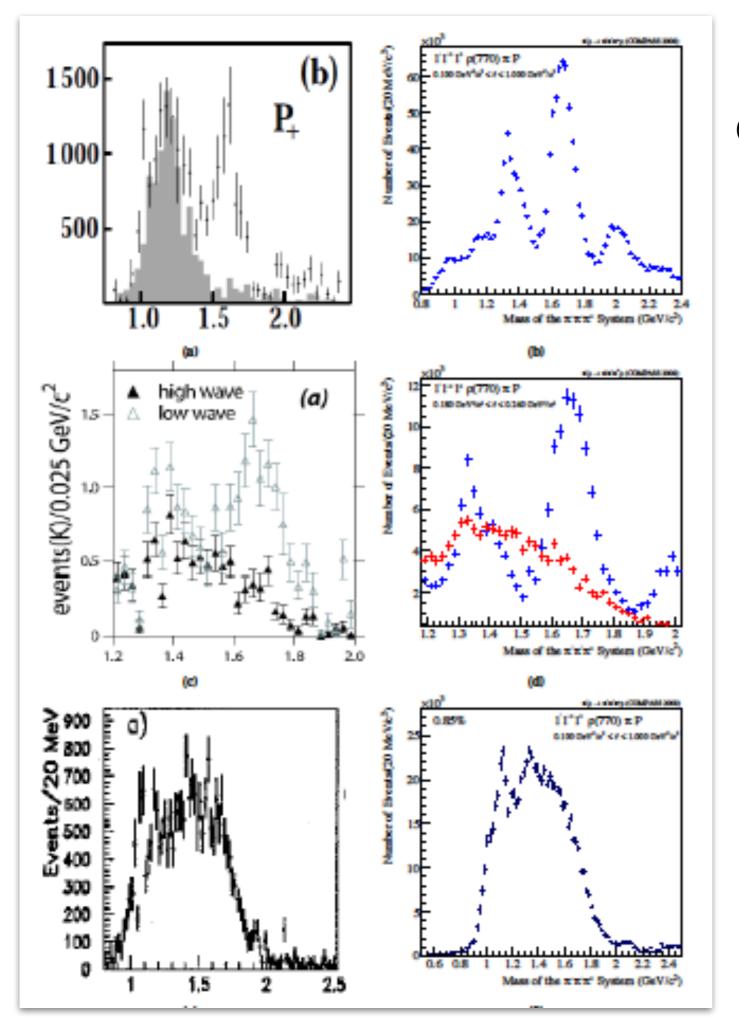
20000

15000

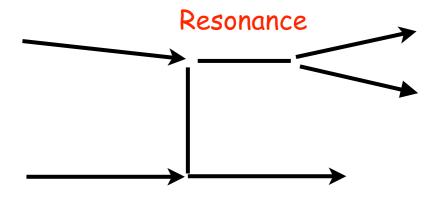
10000

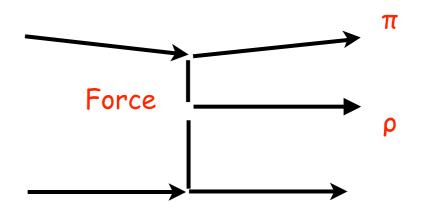
5000

1.5

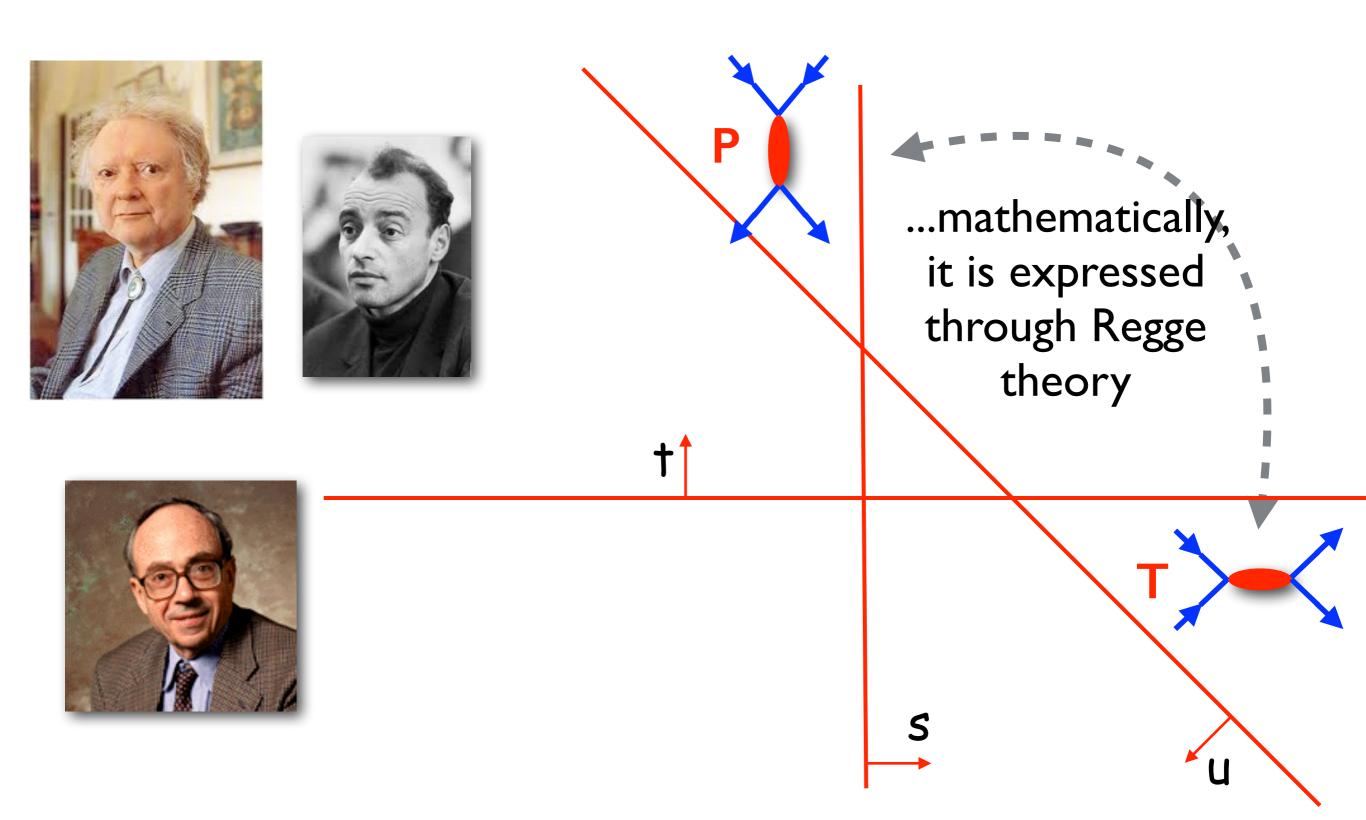


3pion spectrum in diffractive dissociation on hydrogen COMPASS vs E852/BNL and VES



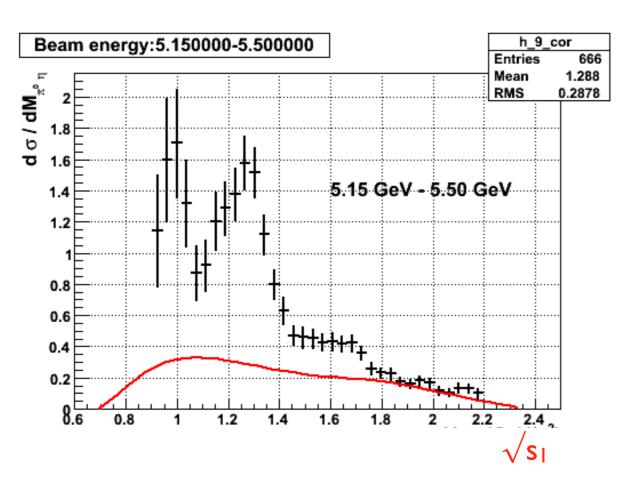


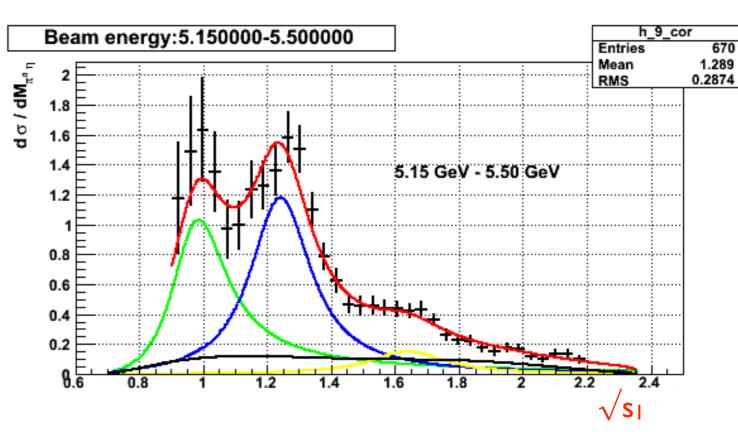
Forces (P) and Resonances (T) are not independent

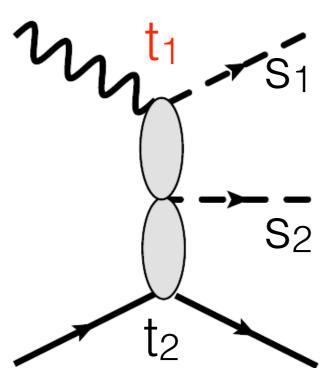


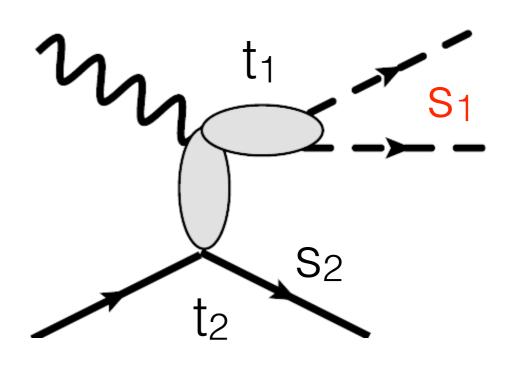
Unitarity enforces Force-Resonance duality = Regge theory

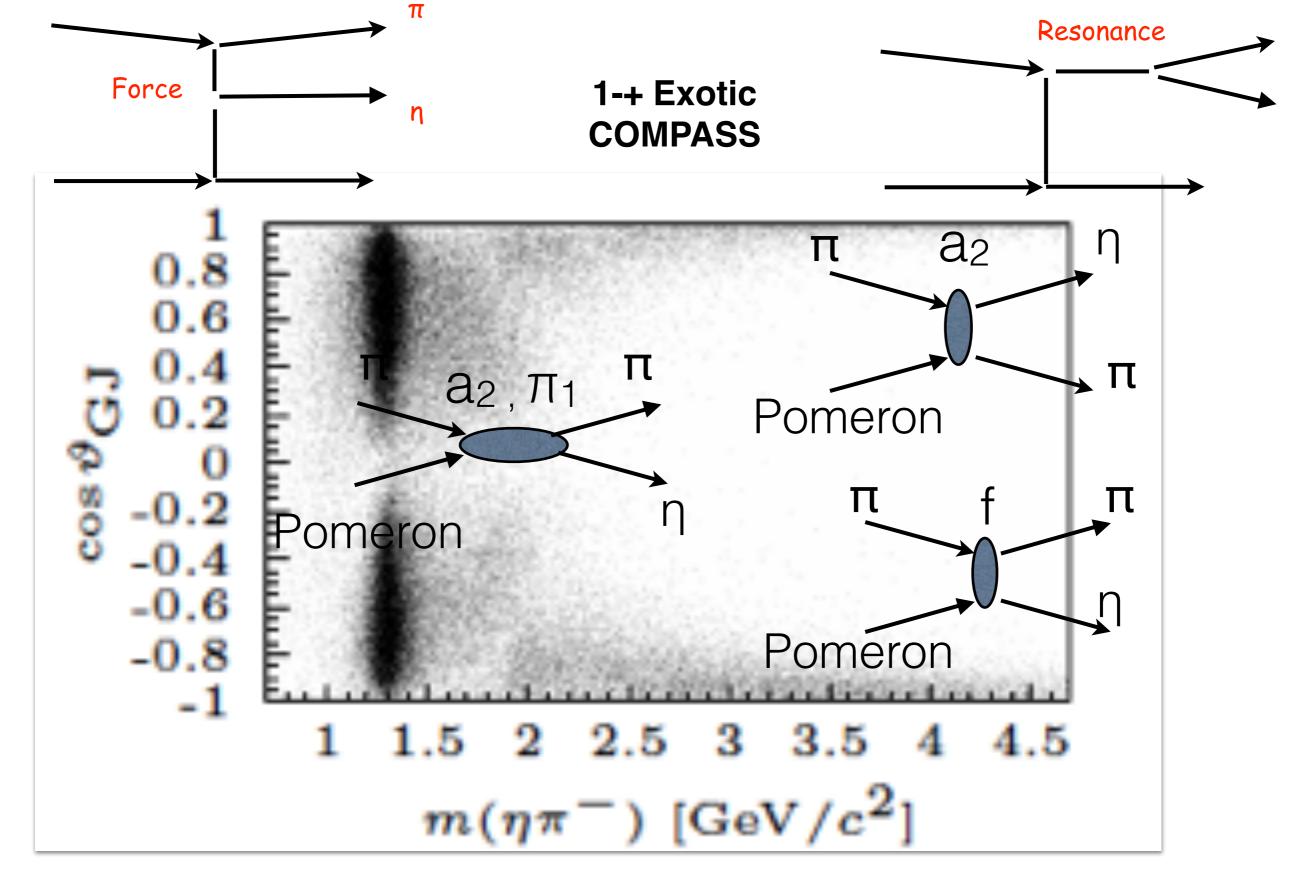
$Yp \rightarrow ηπ⁰p (g12)$











duality —> the production channel is dominated by spineven partial waves. What is the exotic (P-wave) dual to?

JPAC: helping to repair the gap



Develop theoretical, phenomenological/computational tools for hadron experiments

Experiment-theory collaboration

GLOBAL EFFORT

Create a vibrant community

JPAC: Specific Analyses

```
Light meson decays and light quark resonance ω/φ → 3π, πγ (Khuri-Treiman) ω → 3π (Veneziano, B4) η→3π, η→η'π π (Khuri-Treiman) J/Ψ→γπ0π0
```

Photoproduction: (production models, FESR and duality)

$$\gamma p \rightarrow \pi 0 p$$

 $\gamma p \rightarrow pK+K-$

Launched in the Fall of 2014 ~20 papers published

Exotica and XYZ's:

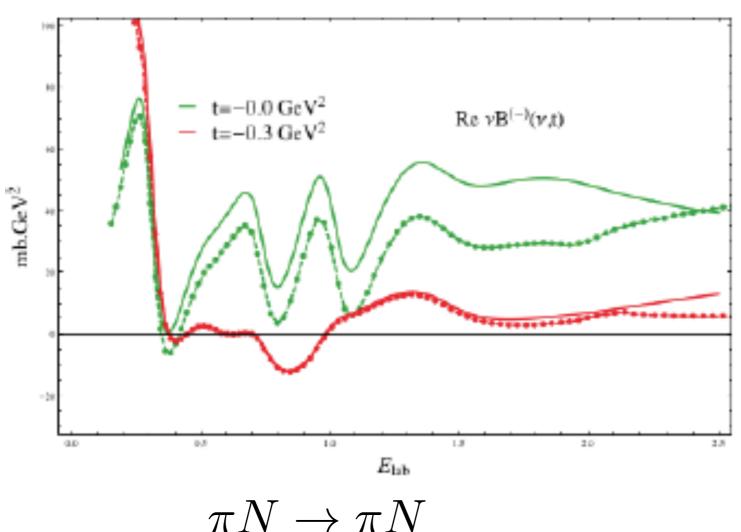
π-p → π-ηp & π-p→π-η'p (FESR)
B0 → Ψ' π- K+ u, Ψ(4260) → J/Ψ π+π-
J/Ψ →
$$3\pi$$
 (Veneziano, B4)

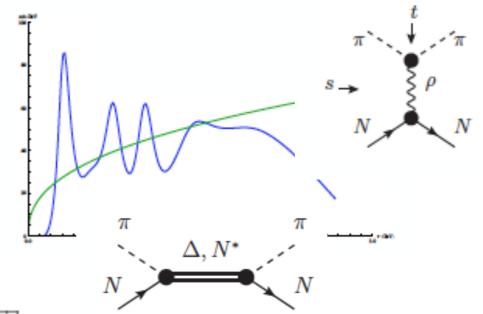
Resonance-Regge physics in meson-baryon scattering

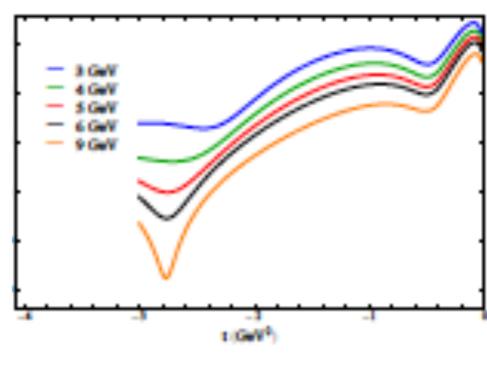
$$\int_{\nu_0}^{\Lambda} \text{Im } A^{(-)}(\nu', t) \nu'^{2k} d\nu' = \beta(t) \frac{\Lambda^{\alpha_{\rho}(t) + 2k + 1}}{\alpha_{\rho}(t) + 2k + 1}$$

Resonance

Regge





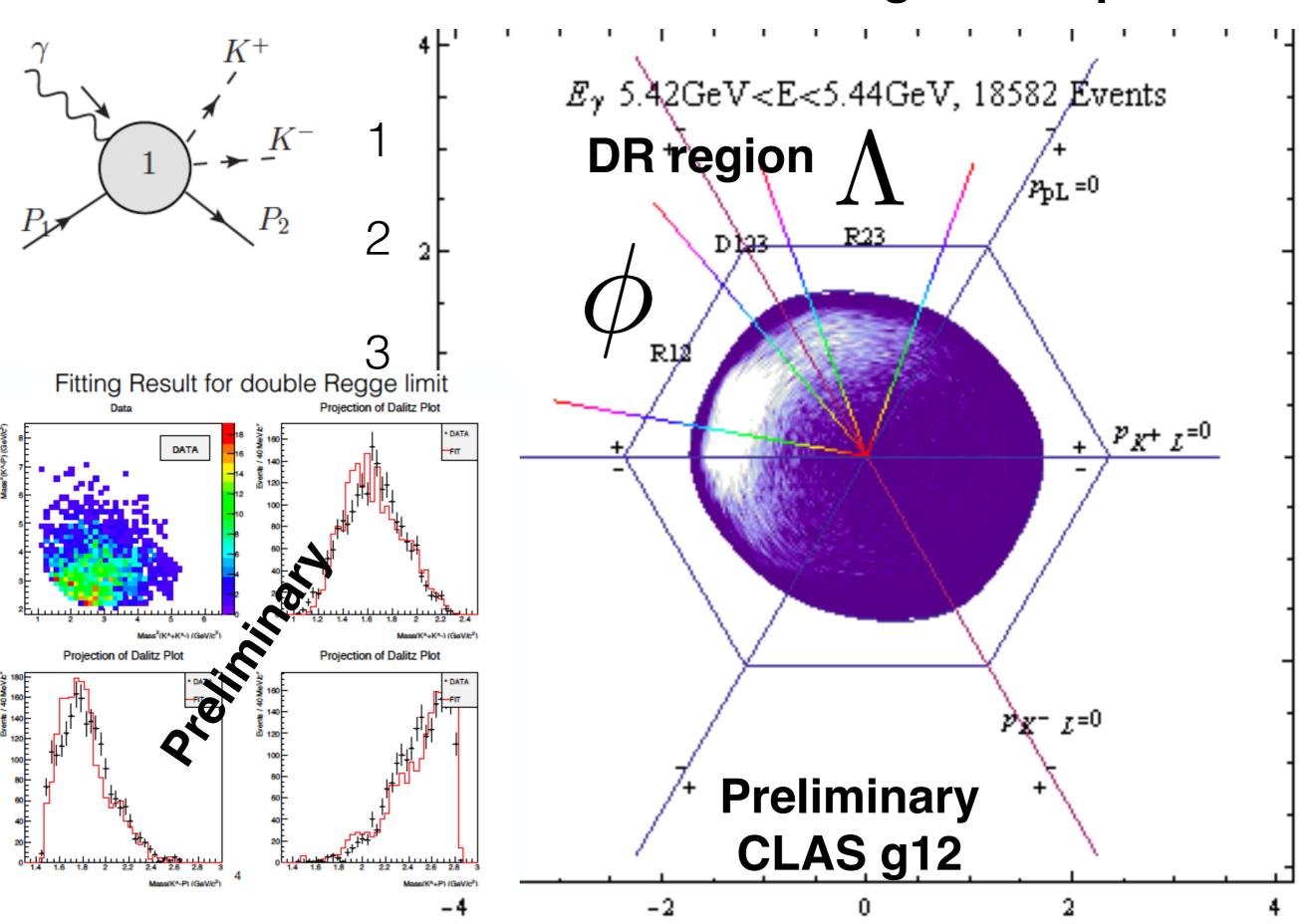


section at CLAS and GluoX energies.

$$\gamma p \to \pi^0 p$$

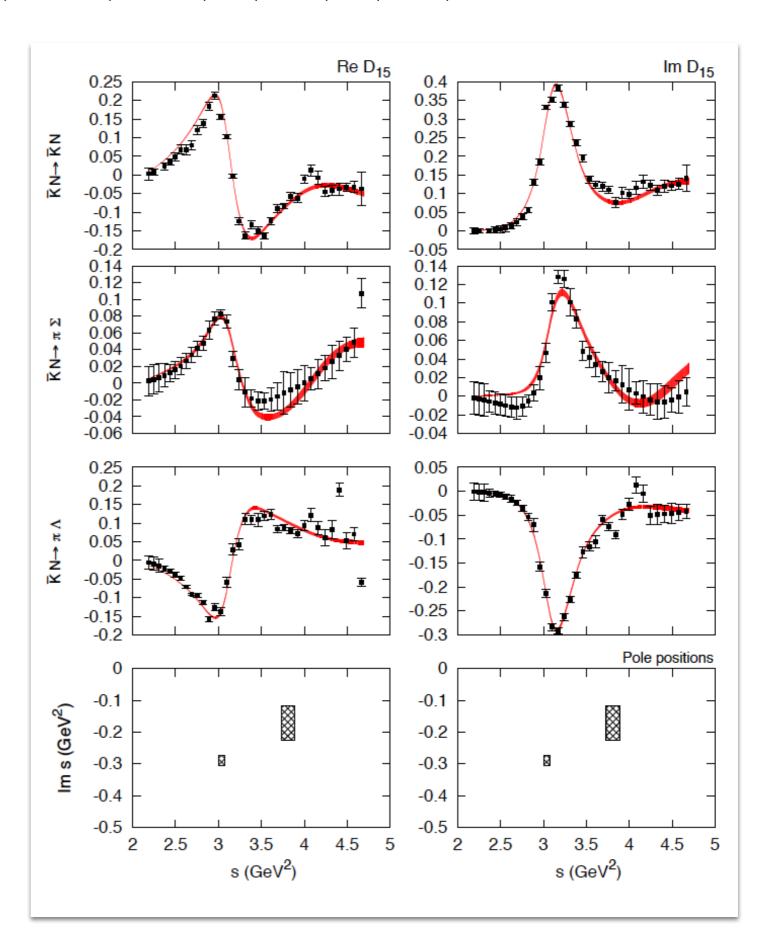
$\Phi(1020)$ 49% **Duality @ JLab** 46% $f_2(1270)$ 9% $f_1(1285)$ 49% $a_2(1320)$ B₅ amplitude description SK+K-6, $f_0(1370)$ seen $f_1(1420)$ dominant $\eta(1440)$ seen $\Lambda(1520)$ $a_0(1450)$ seen $f_0(1500)$ seen 89% $f_{2}(1525)$ $\Phi(1680)$ seen $\rho 3(1690)$ 1.6% ρ (1700) seen $f_0(1710)$ seen $\Phi_3(1850)$ seen $a_4(1450)$ seen 7% $f_4(2050)$ ф(1020) SK⁻p $\gamma p \rightarrow K+K-p (g11,g12)$ D.Schott, JPAC/GWU M. Shi JPAC/Pekin U.

van Hove longitudinal plot



$$(\bar{K}N), (\pi\Sigma), (\pi\Lambda) \to (\bar{K}N), (\pi\Sigma), (\pi\Lambda)$$

Hyperon production what is Lambda(1405)?

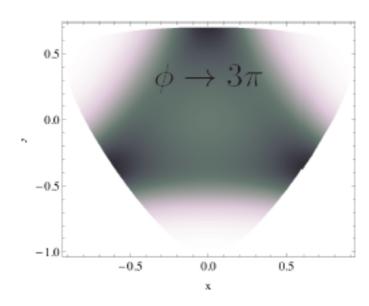


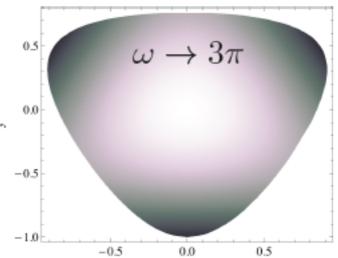
C.Fernandez Ramirez, JPAC

Light meson decays

$$\eta \rightarrow 3\pi \quad \omega \rightarrow 3\pi \quad \varphi \rightarrow 3\pi$$

- Constrained phase space,
- Effective (chiral) dynamics,
- Small number of partial waves,
- Amendable to dispersive methods





$$\eta \rightarrow 3\pi$$

$$A(s,t,u) = -\frac{1}{Q^2} \frac{M_K^2}{M_\pi^2} \frac{M_K^2 - M_\pi^2}{3\sqrt{3}F_\pi^2} M(s,t,u) \qquad \left(Q^2 \equiv \frac{m_s^2 - \hat{m}^2}{m_d^2 - m_u^2} \right)$$

$$\left(Q^2 \equiv \frac{m_s^2 - \hat{m}^2}{m_d^2 - m_u^2}\right)$$

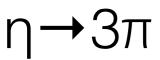


TABLE III: Dalitz plot parameters for $\eta \to 3\pi^0$. Set 1 and Set 2 correspond to (I, L) = (0, 0), (1, 1) and (I, L) = (0, 0), (2, 0), (1, 1) cases respectively (see Table I).

	α	β
GAMS-2000 [49]	-0.022 ± 0.023	-
Crystal Barrel, LEAR [50]	-0.052 ± 0.020	-
Crystal Ball, BNL [14]	-0.031 ± 0.004	-
SND [51]	-0.010 ± 0.023	-
CELSIUS-WASA [12]	-0.026 ± 0.014	-
WASA-at-COSY [52]	-0.027 ± 0.009	-
MAMI-B [53]	-0.032 ± 0.004	-
MAMI-C [46]	-0.032 ± 0.003	-
KLOE [16]	-0.0301 ± 0.0050	-
PDG average [1]	-0.0317 ± 0.0016	-
Theory		
Set 1	-0.023 ± 0.004	-0.000 ± 0.002
Set 2	-0.020 ± 0.004	-0.001 ± 0.003
NLO [21]	+0.013	-
NNLO [22]	$+0.013 \pm 0.032$	-
Kambor et al. [23]	-0.007 0.014	-
NREFT [28]	-0.025 ± 0.005	-0.004 ± 0.001
Kampf et al. [29]	-0.044 ± 0.004	-

P.Guo, I.Danilkin, HASPEC/JPAC

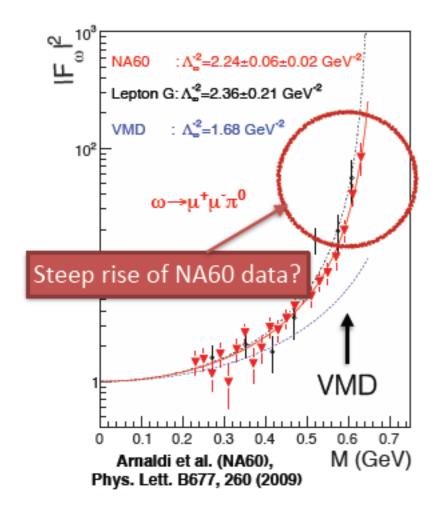
Ex:
$$\eta \rightarrow 3\pi$$
 $\Gamma_{\eta \rightarrow 3\pi} = 66 \text{ eV} + 94 \text{ eV} + ... = 296 \text{ eV}$

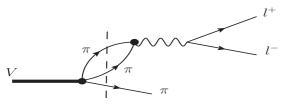
LO NLO Experiment

TABLE II: Dalitz plot parameters for $\eta \to \pi^+\pi^-\pi^0$. Set 1 and Set 2 correspond to (I,L) = (0,0), (1,1) and (I,L) = (0,0), (2,0), (1,1) cases respectively (see Table I).

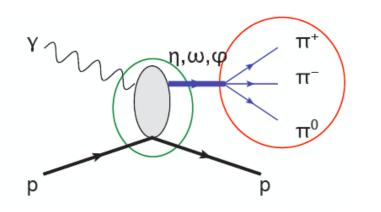
	a	b	d	f
WASA-at-COSY [11]	-1.144 ± 0.018	$0.219 \pm 0.019 \pm 0.037$	$0.086 \pm 0.018 \pm 0.018$	0.115 ± 0.037
KLOE [15]	$-1.090 \pm 0.005^{+0.008}_{-0.019}$	$0.124 \pm 0.006 \pm 0.010$	$0.057 \pm 0.006^{+0.007}_{-0.016}$	$0.14 \pm 0.01 \pm 0.02$
CBarrel [13]	-1.22 ± 0.07	0.22 ± 0.11	$0.06\pm0.04(\mathrm{fixed})$	-
Layter et al. [47]	-1.080 ± 0.014	0.03 ± 0.03	0.05 ± 0.03	_
Gormley et al. [48]	-1.17 ± 0.02	0.21 ± 0.03	0.06 ± 0.04	-
Theory				
Ce 1	1.110 ± 0.000	0.100 ± 0.010	0.047 ± 0.005	0.000 ± 0.004
Set 2	-1.117 ± 0.035	0.188 ± 0.014	0.079 ± 0.003	0.090 ± 0.003
NLO [21]	-1.371	0.452	0.053	0.027
NNLO [22]	-1.271 ± 0.075	0.394 ± 0.102	0.055 ± 0.057	0.025 ± 0.160
Kambor et al. [23]	-1.16	0.240.26	0.090.10	-
NREFT [28]	-1.213 ± 0.014	0.308 ± 0.023	0.050 ± 0.003	0.083 ± 0.019

ω→π⁰γ*

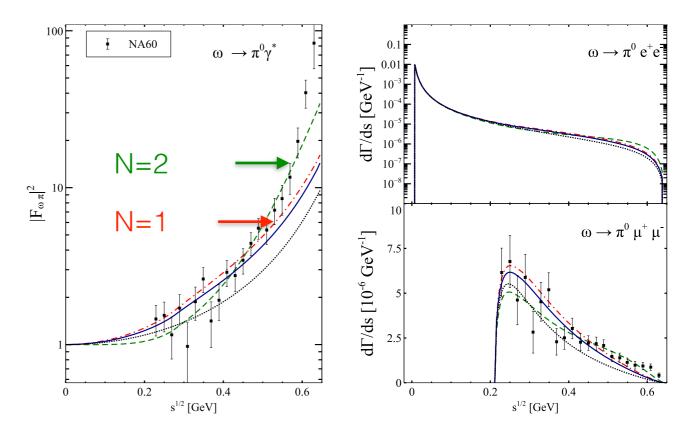




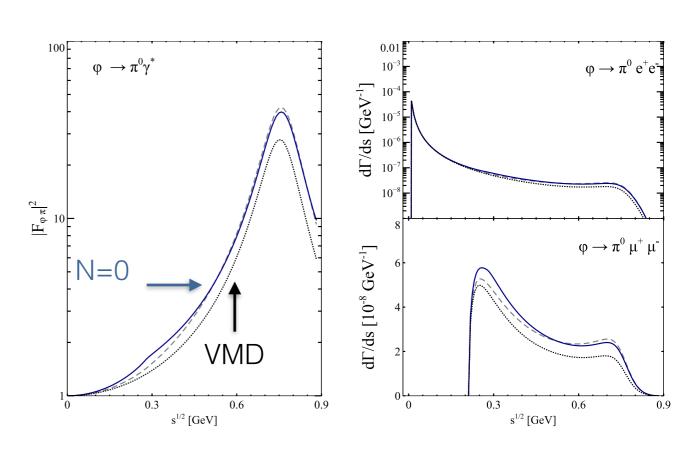
$$f_{V\pi}(s) = \int_{s_{\pi}}^{s_i} \frac{ds'}{\pi} \frac{\text{Disc } f_{V\pi}(s')}{s' - s} + \sum_{i=0}^{N} C_i \,\omega(s)^i$$



upcoming CLAS g12



I.Danilkin, JPAC



JPAC memebrs

Igor Danilkin

Andrew Jackura

Astrid Blin (Valencia)

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Meng Shi

Lingun Dai

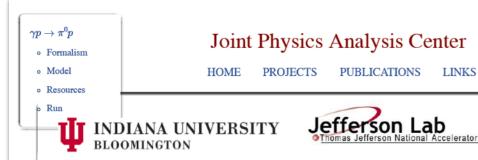
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https://wiki.jlab.org/jpac/





We present the model published in [Mat15a].

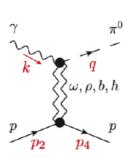
The differential cross section for $\gamma p \to \pi^0 p$ is computed with Regge amplitudes in the domain $E_\gamma \ge 4 \, \text{GeV}$ and $0.01 \le |t| \le 3 \, (\text{in GeV}^2)$.

The formulas can be extrapolated outside these intervals.

We use the CGLN invariant amplitudes A_i defined in [Chew57a].

See the section Formalism for the definition of the variables.

The fitting procedure is detailed in [Mat15a] . We report here only the main feature of the model.



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Formalism

The differential cross section is a function of 2 variables. The first is the beam energy in the laboratory frame E_{γ} (in GeV) or the total energy squared s (in GeV²). The second is the cosine of the scattering angle in the rest frame $\cos \theta$ or the momentum transferred squared t (in GeV²).

The differential cross section is expressed in term of the parity conserving helicity invariant amplitudes in the t -channel F_i

$$\frac{d\sigma}{dt} = \frac{389.4}{64\pi} \frac{k_t^2}{4M^2 E_7^2} \Big[2\sin^2\theta_t \left(t|F_1|^2 + 4p_t^2|F_2|^2 \right) + (1-\cos\theta_t)^2 |F_3 + 2\sqrt{t}p_t F_4|^2 + (1+\cos\theta_t)^2 |F_3 - 2\sqrt{t}p_t F_4|^2 \Big]$$

The differential cross section is expressed in μ b/GeV². We used $(\hbar c)^2 = 0.3894$ mb.GeV².

The t -channel is the rest frame of the process $\gamma_{\pi}^{0} \rightarrow p\bar{p}$.

In the t -channel, the momenta of the nucleon p_t and the pion k_t and the cosine of the scattering angle are

$$k_t = \frac{1}{2}\sqrt{t - 4M^2}, \qquad q_t = \frac{t - \mu^2}{2\sqrt{t}}, \qquad \cos\theta_t = \frac{s - u}{4p_t k_t}.$$
 (A.1)

The invariant amplitudes F_i are related through the CGLN A_i amplitudes [Chew57a] by

$$F_1 = -A_1 + 2MA_4, \quad \eta = +, \quad CP = +$$
 (A.2)

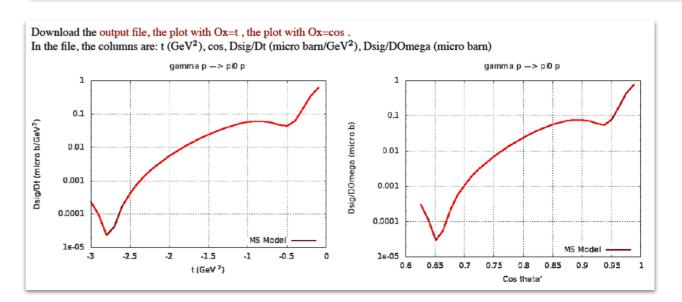
$$F_2 = A_1 + tA_2, \qquad \eta = -, \qquad CP = -$$
 (A.3)

$$F_3 = 2MA_1 - tA_4, \qquad \eta = +, \qquad CP = +$$
 (A.4)

$$F_4 = A_3 \qquad \qquad \eta = -, \qquad CP = + \tag{A.5}$$

The F_i amplitudes have good quantum numbers of the t-channel, the naturality $n = P(-1)^J$ and the product CP.

Run the code Choose the beam energy in the lab frame E_{γ} , the other variable (t or $\cos \theta$) and its minimal, maximal, and increment values. If you choose t ($\cos \theta$) only the min, max and step values of t ($\cos \theta$) are read. $E_{\gamma} \text{ in GeV } 9$ $t \cap \cos \theta$ $t \text{ in GeV}^2 \text{ (min max step) } -3$ $cos \theta$ t (min max step) 0.5 0.001



special thanks to Vincent Mathieu



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