



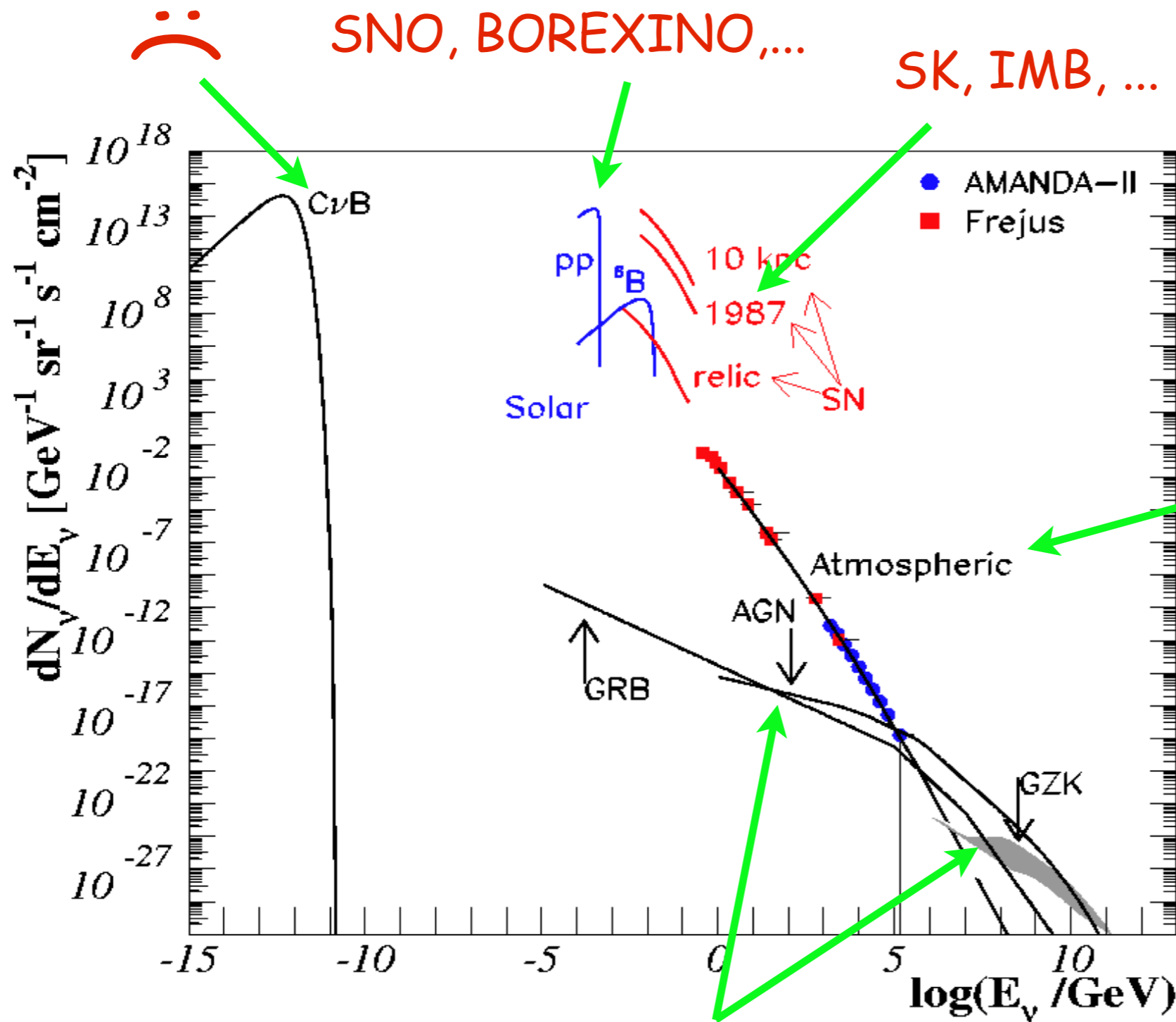
Interpreting the IceCube events by decaying dark matter hints and constraints

Arman Esmaili

Laboratori Nazionali del Gran Sasso (LNGS)
Theory Group

20/Nov/2014

Neutrino Sky



SNO, BOREXINO, ...

SK, IMB, ...

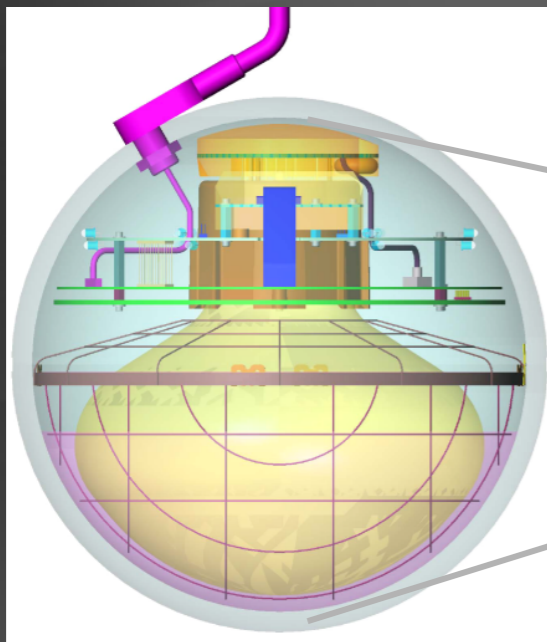
SK, AMANDA, IceCube...

Background for astrophysical neutrinos

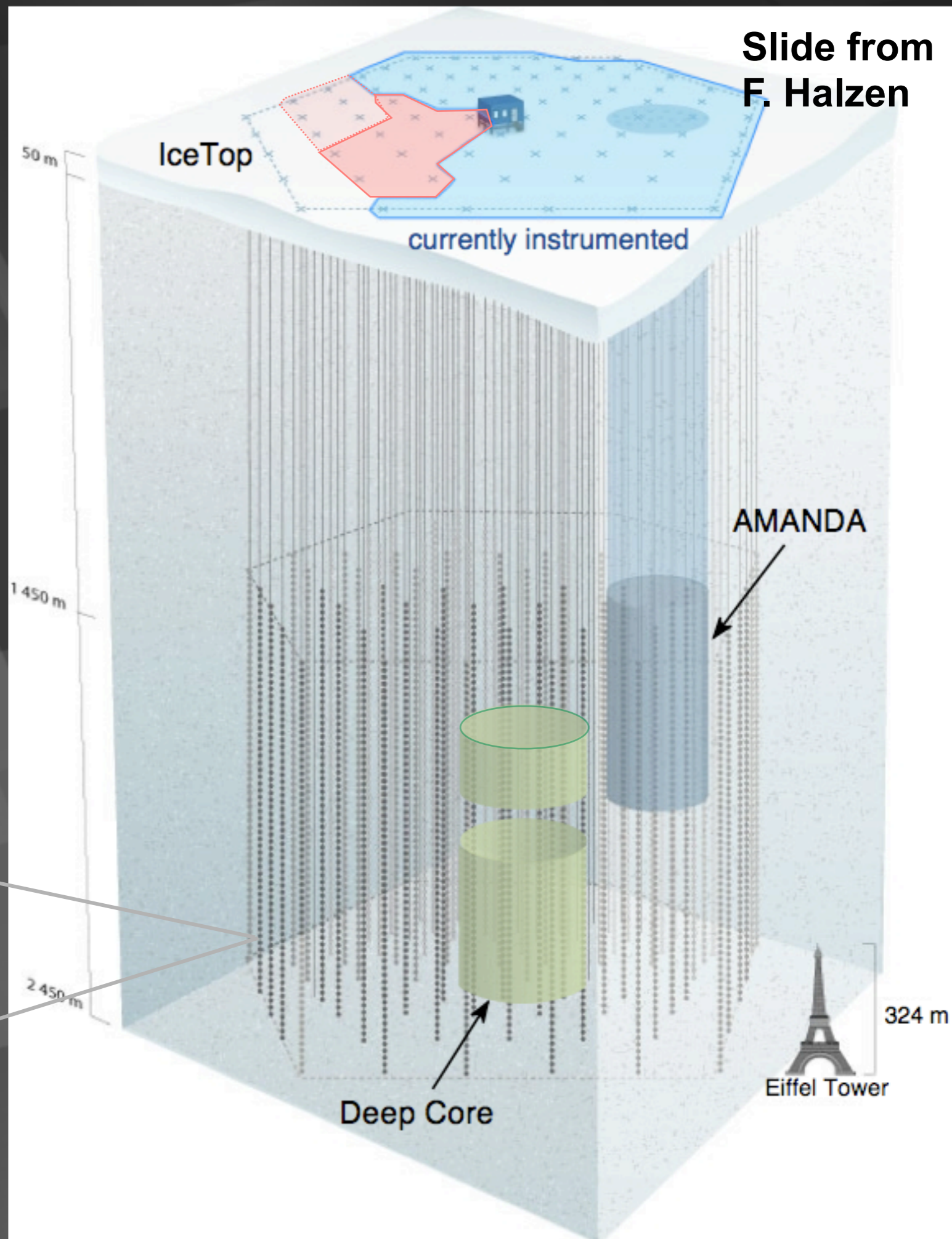
IceCube ?

IceCube / Deep Core

- 5320 optical modules on 86 strings (+ IceTop)
- detects ~ 220 neutrinos and 1.7×10^8 muons per day
- threshold 10 GeV
- angular resolution < 1 degree



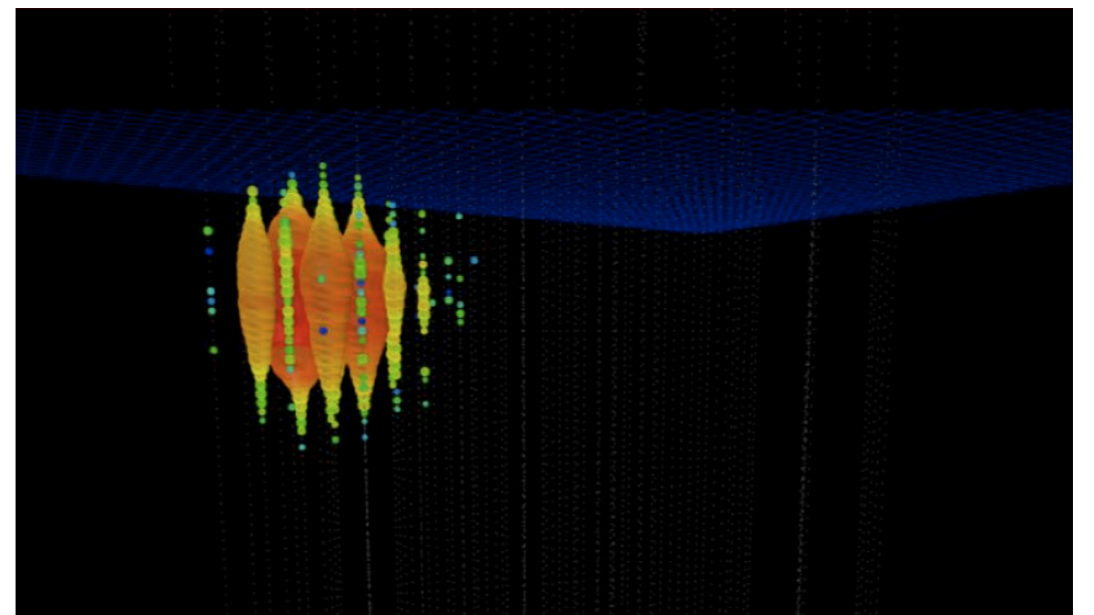
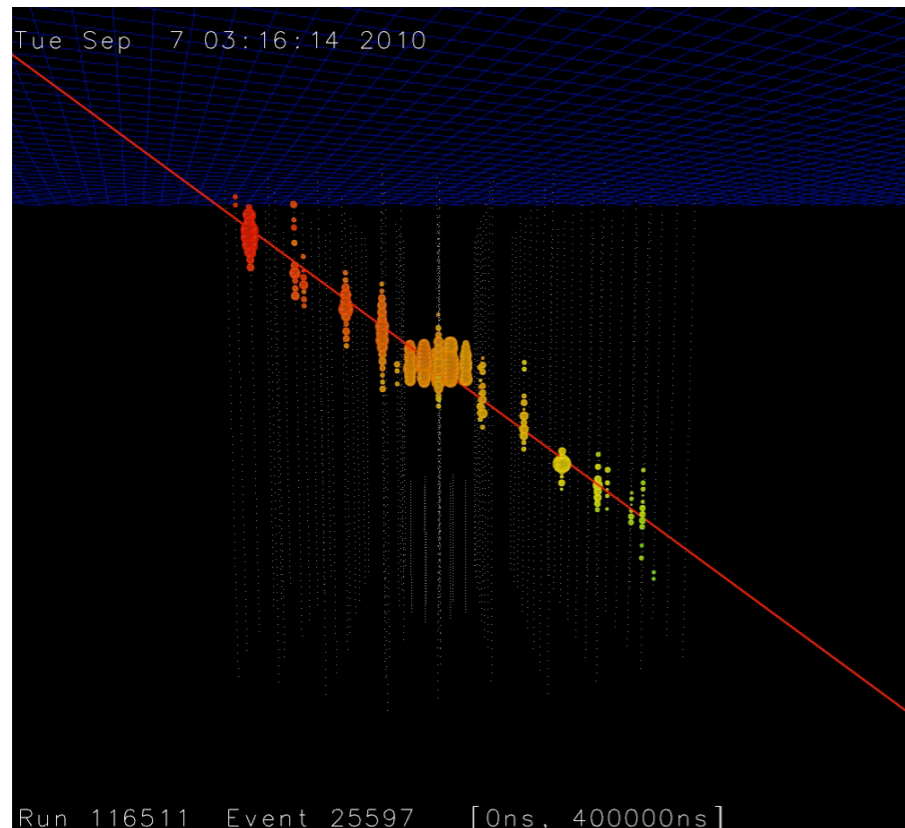
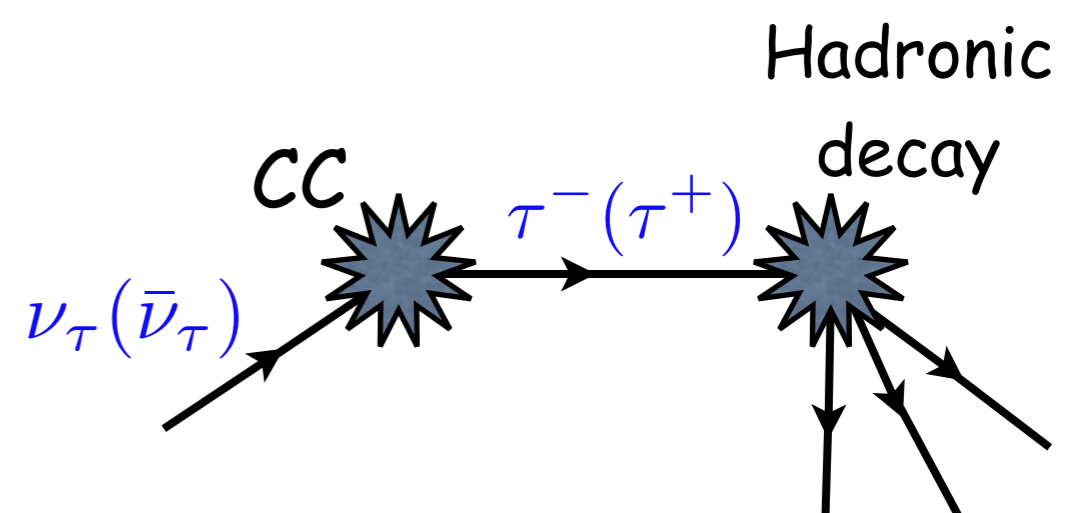
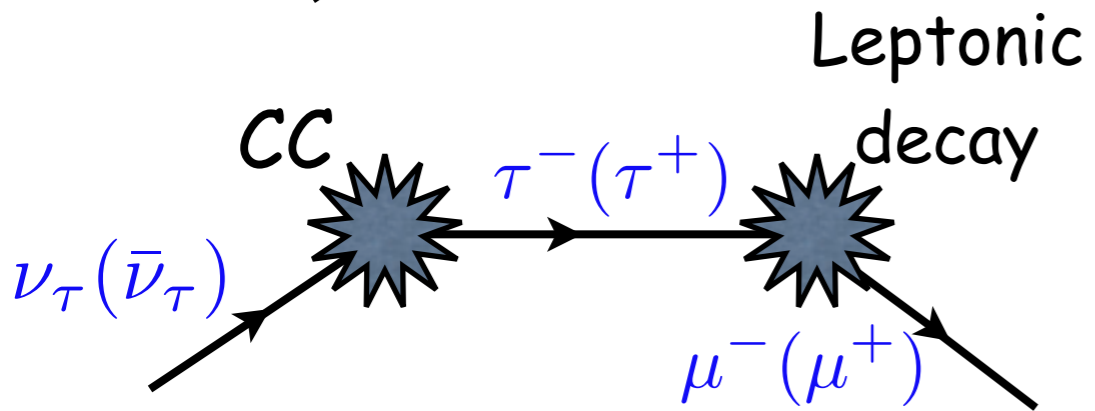
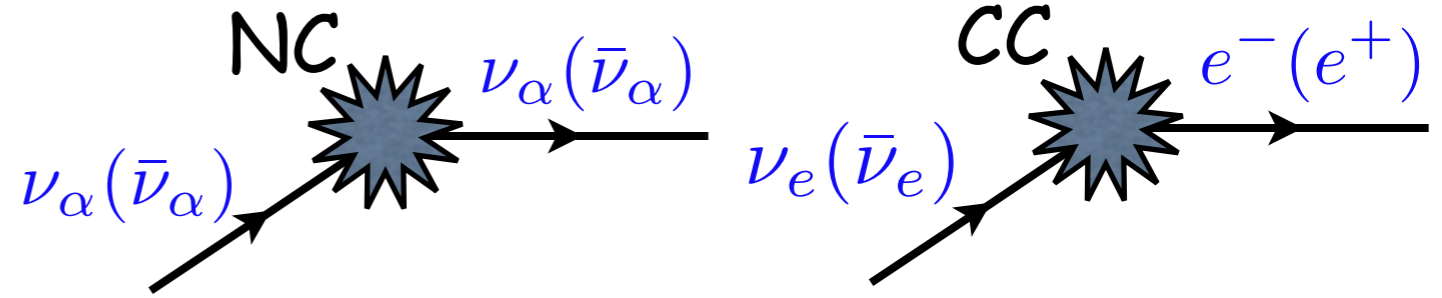
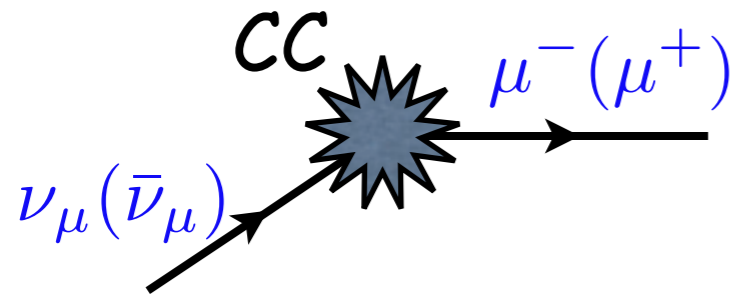
Digital Optical Module (DOM)



Flavoring at IceCube

muon-track events

cascade events



figures from IceCube website

Flavoring at IceCube

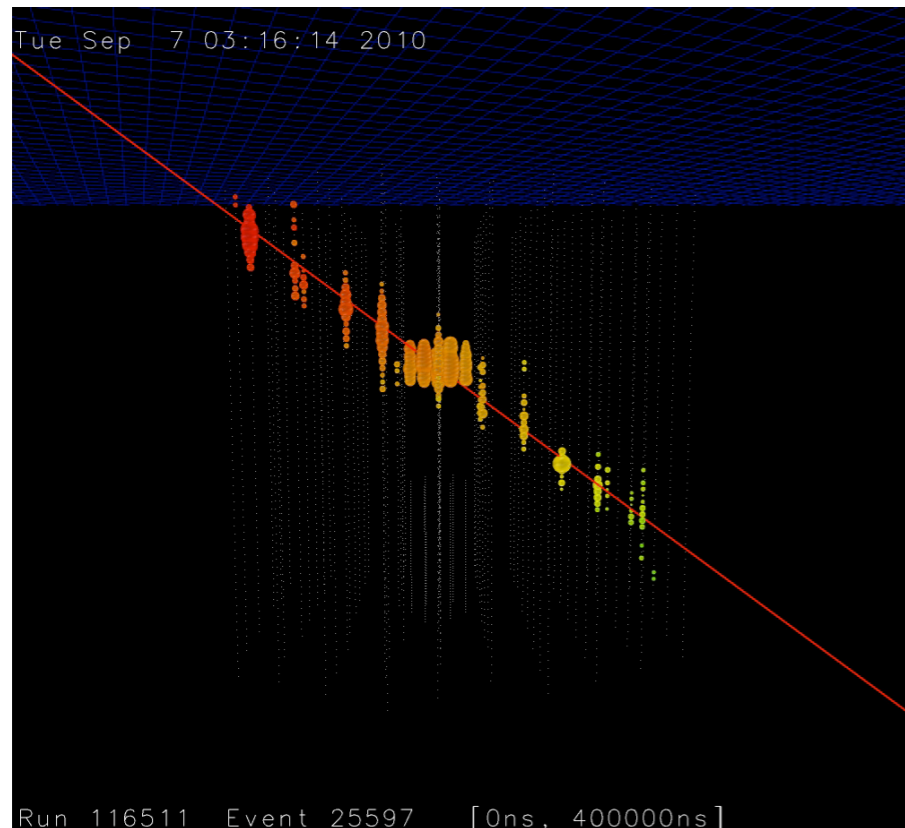
muon-track events

great angular resolution ($< 1^\circ$)

moderate energy resolution ($\sigma_E \sim E$)

ν_τ

μ ($\bar{\mu}$)



cascade events

poor angular resolution ($< 10^\circ - 20^\circ$)

great energy resolution ($\sigma_E \sim 0.15 \times E$)

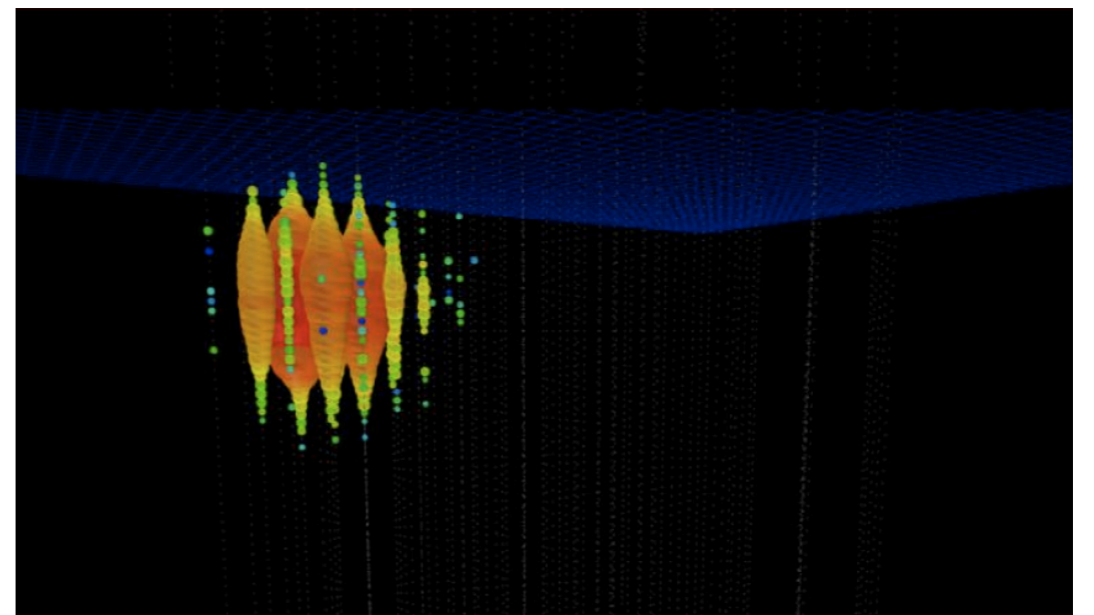
ν_α ($\bar{\nu}_\alpha$)

(e^+)

mic

y

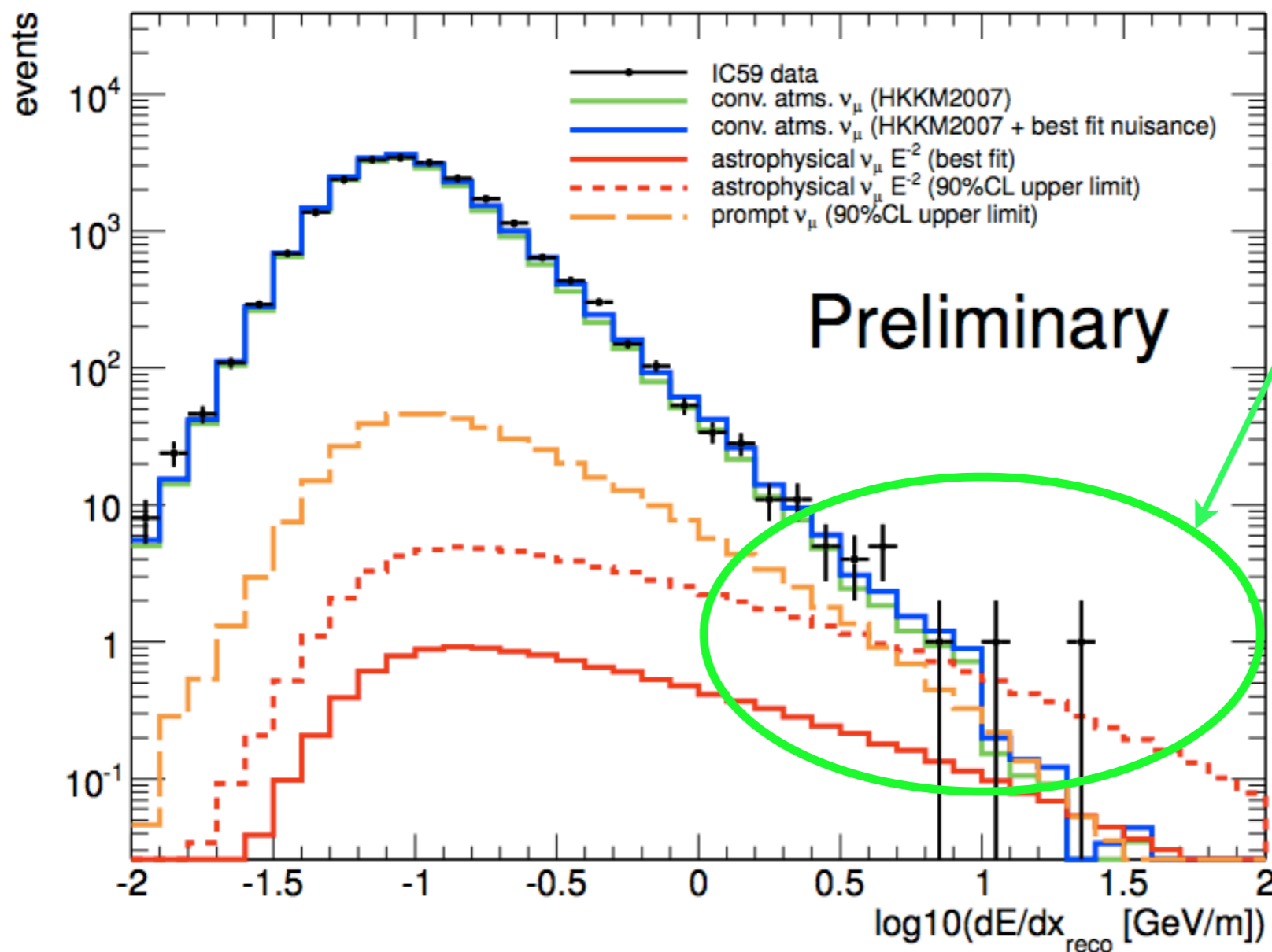
ν_τ ($\bar{\nu}_\tau$)



figures from
IceCube
website

Mission for IceCube began !

✓ muon-track events at IceCube-59, 348 days livetime.



excess in high energy tail
($\sim 300 \text{ TeV}$)

prompt atm neutrinos ?
or
astrophysical neutrinos?

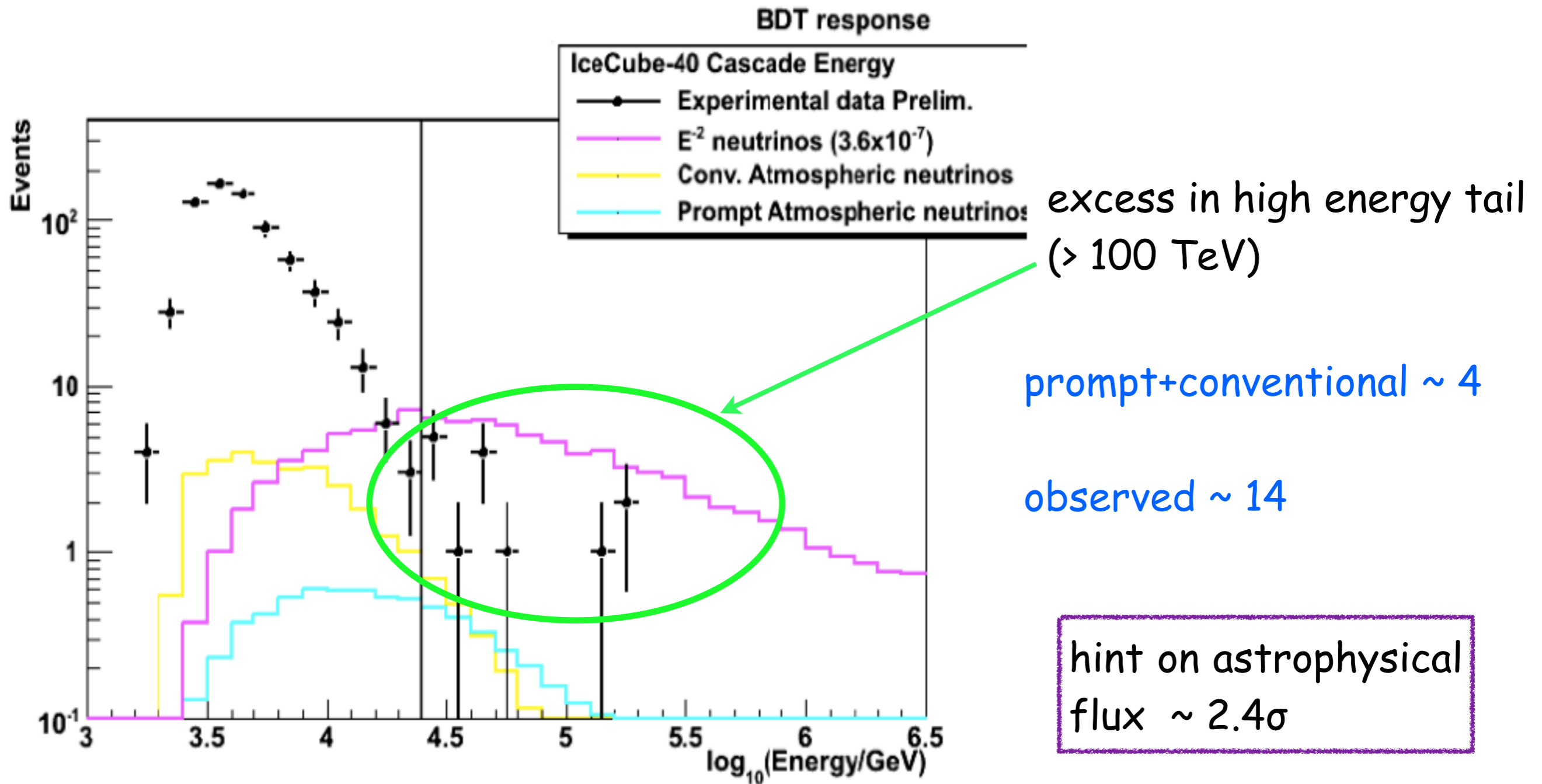
analysis shows preference
to astrophysical origin of
excess, significance $\sim 2.1\sigma$

A. Schukraft [IceCube Collaboration]

Nucl. Phys. Proc. Suppl. 266 (2013) [arXiv:1302.0127]

Mission for IceCube began !

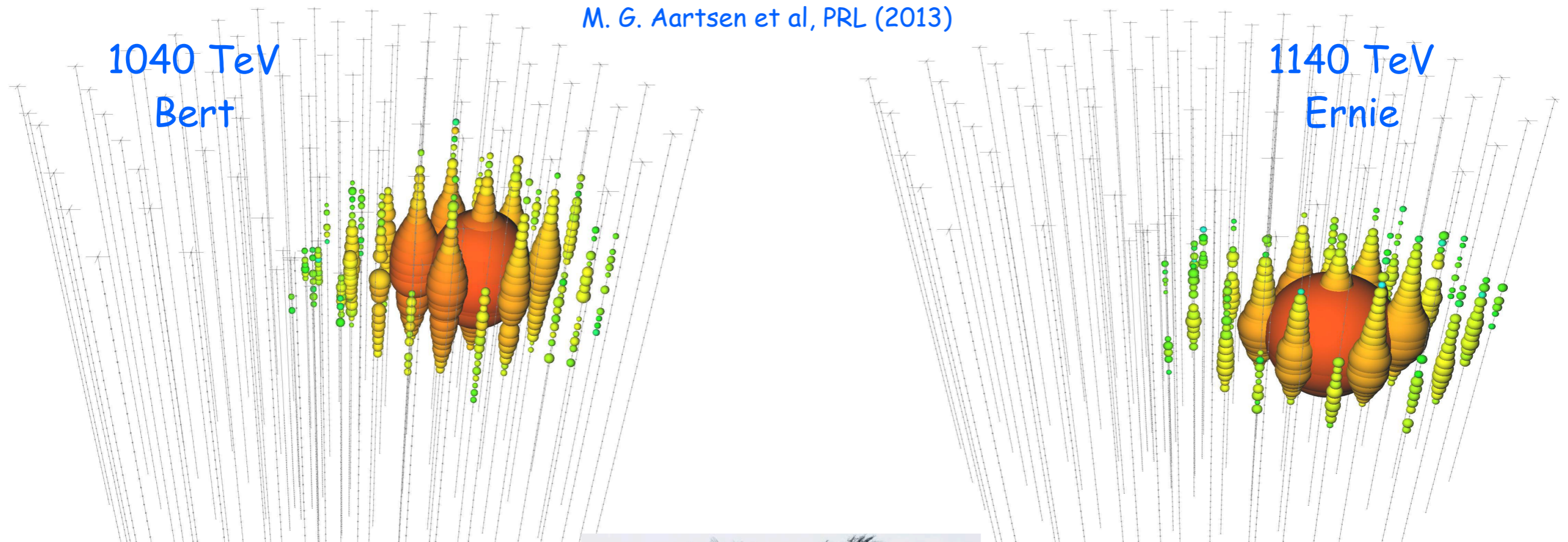
✓ cascade events at IceCube-40, 367 days livetime.



E. Middell, PhD thesis

Mission for IceCube began !

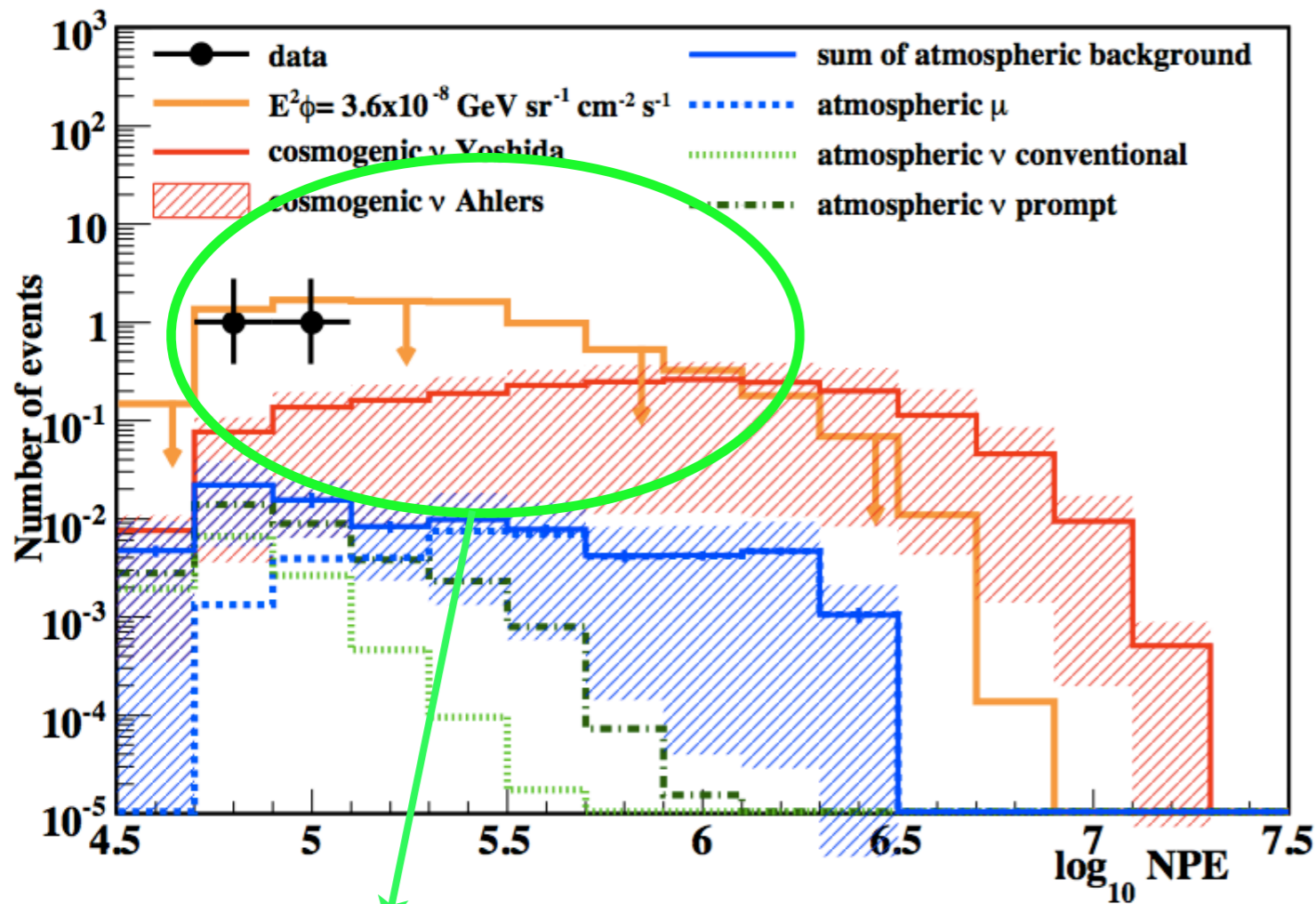
- ✓ The two PeV cascade events, 616 days livetime



Mission for IceCube began !

✓ The two PeV cascade events, 616 days livetime

M. G. Aartsen et al. [IceCube Collaboration],
Phys. Rev. Lett. 111 (2013), [arXiv:1304.5356]



expected bkg. (conventional+prompt)
~ 0.08(-0.057)(+0.041) sys.

excess of events ~ 2.8σ

GZK ? too low energy, more events
should be seen in higher energies

astrophysical ? an E^{-2} spectrum
would give ~ 9 more events in
higher energies

it is like a cut-off at ~ PeV

demands more statistics

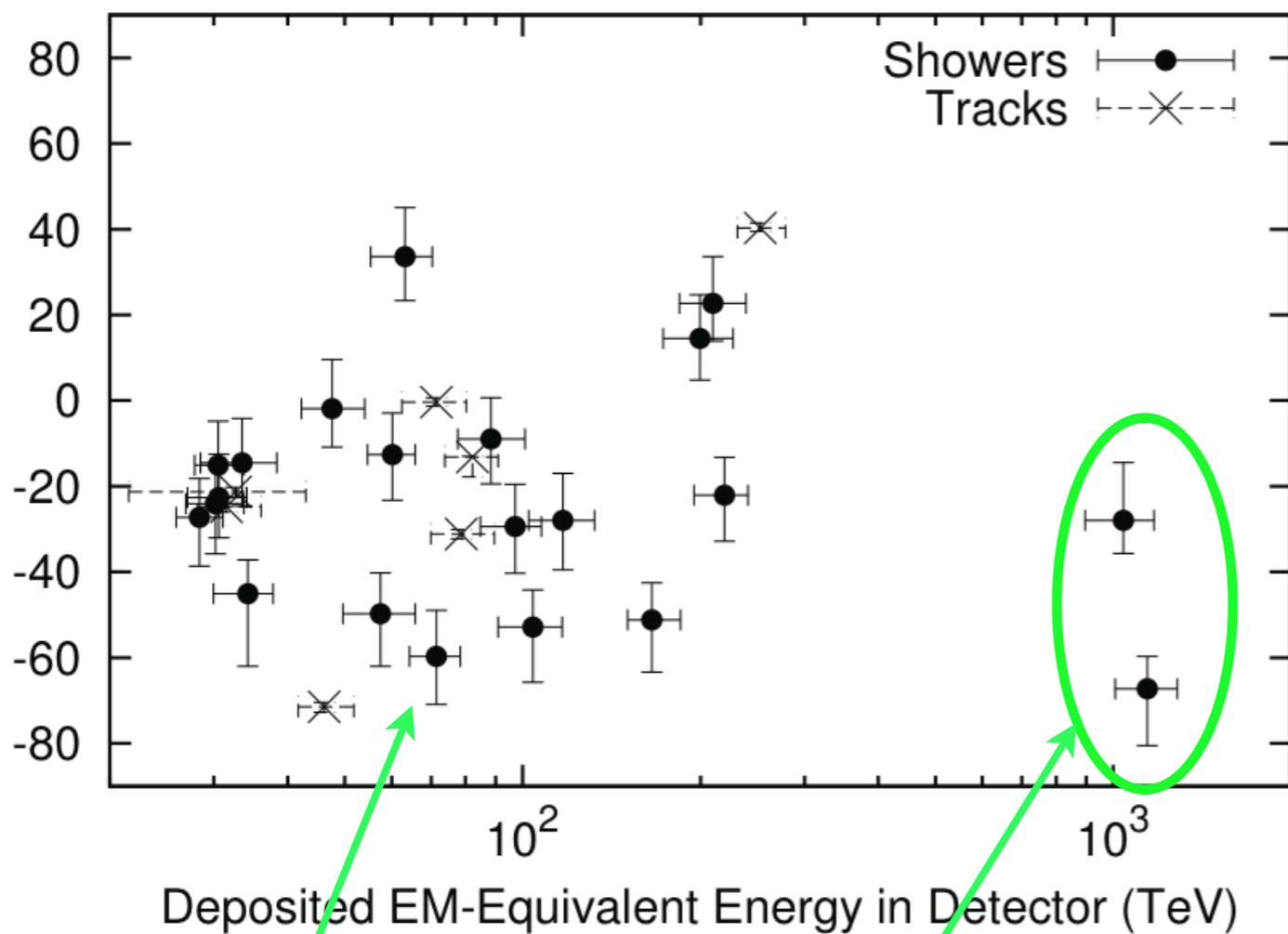
flavor composition ?
NC of ν_α or CC of ν_e

isotropy ?

Mission for IceCube began !

✓ Looking for lower energy contained events, 662 days livetime

M. G. Aartsen et al. [IceCube Collaboration],
Science 342 (2013), [arXiv:1311.5238]



The whole family!

previous PeV
cascade events
(Bert and Ernie)

✓ 26 more events

✓ all the new events are lower in energy

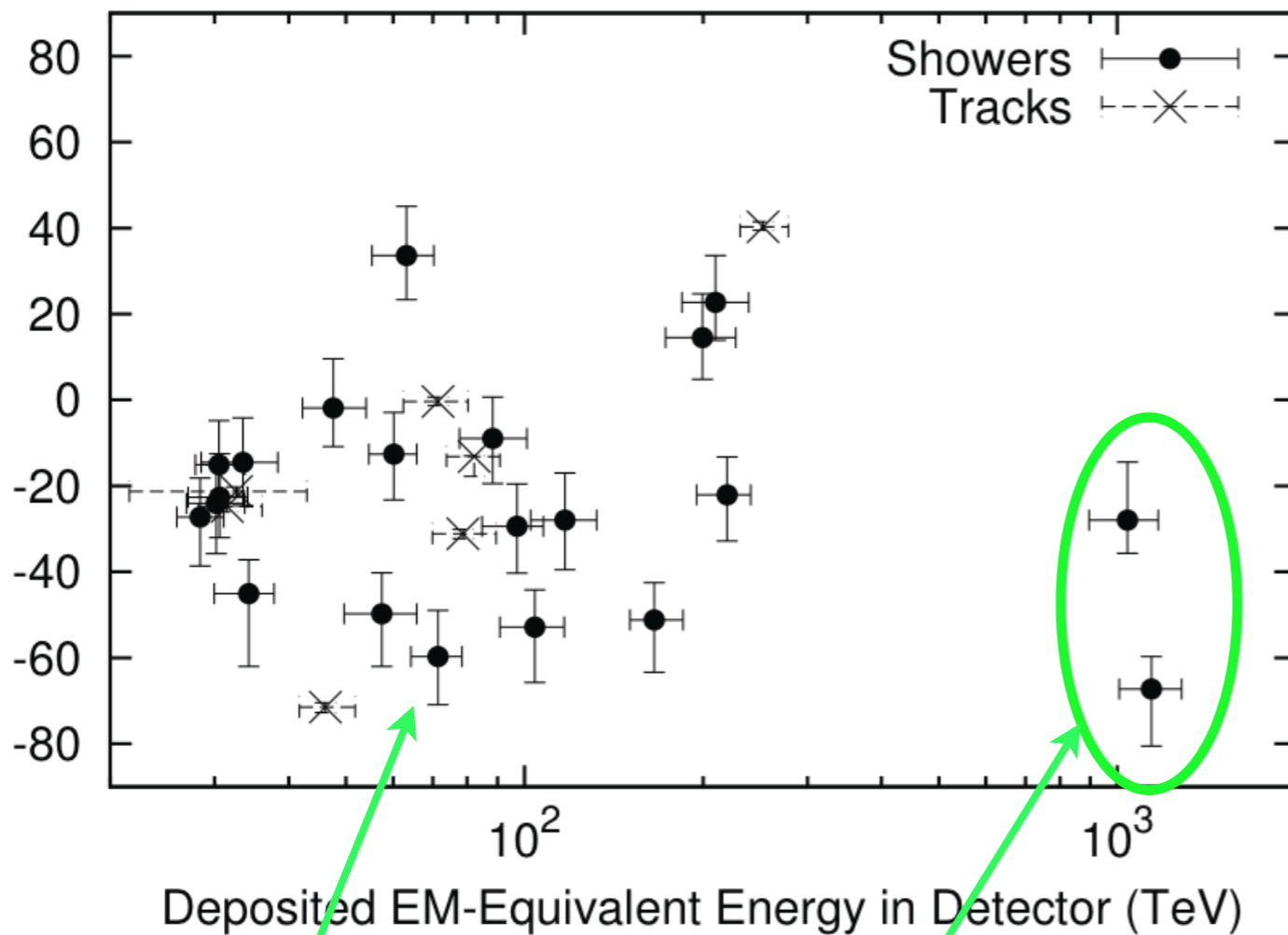
✓ expected bkg. (conventional+prompt)
 $\sim 10.6(+4.5)(-3.5)$ sys.

excess of events $\sim 4.3\sigma$

Mission for IceCube began !

✓ Looking for lower energy contained events, 662 days livetime

M. G. Aartsen et al. [IceCube Collaboration],
Science 342 (2013), [arXiv:1311.5238]



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The whole family!

previous PeV
cascade events
(Bert and Ernie)

which one?

atmospheric ?

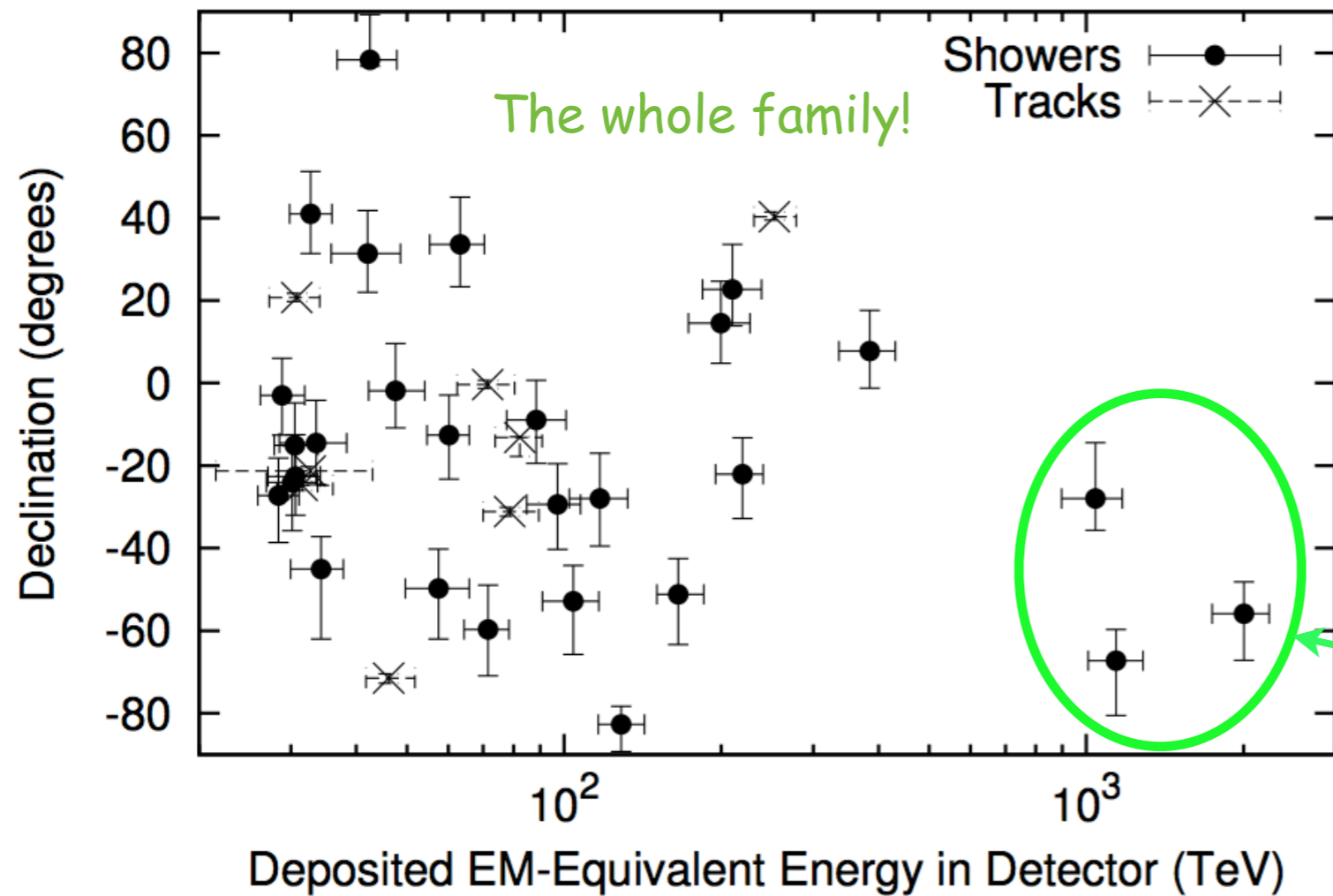
astrophysical ?

or something else ?

Mission for IceCube began !

✓ Looking for lower energy contained events, 988 days livetime

M. G. Aartsen et al. [IceCube Collaboration],
PRL 113 (2014), [arXiv:1405.5303]



✓ totally 37 events

✓ three events with energy \sim PeV

✓ expected bkg. (conventional+prompt)
 $\sim 15.6(+10.1)(-5.8)$ sys.

excess of events $\sim 5.7\sigma$

which one?

atmospheric ?

astrophysical ?

or something else ?



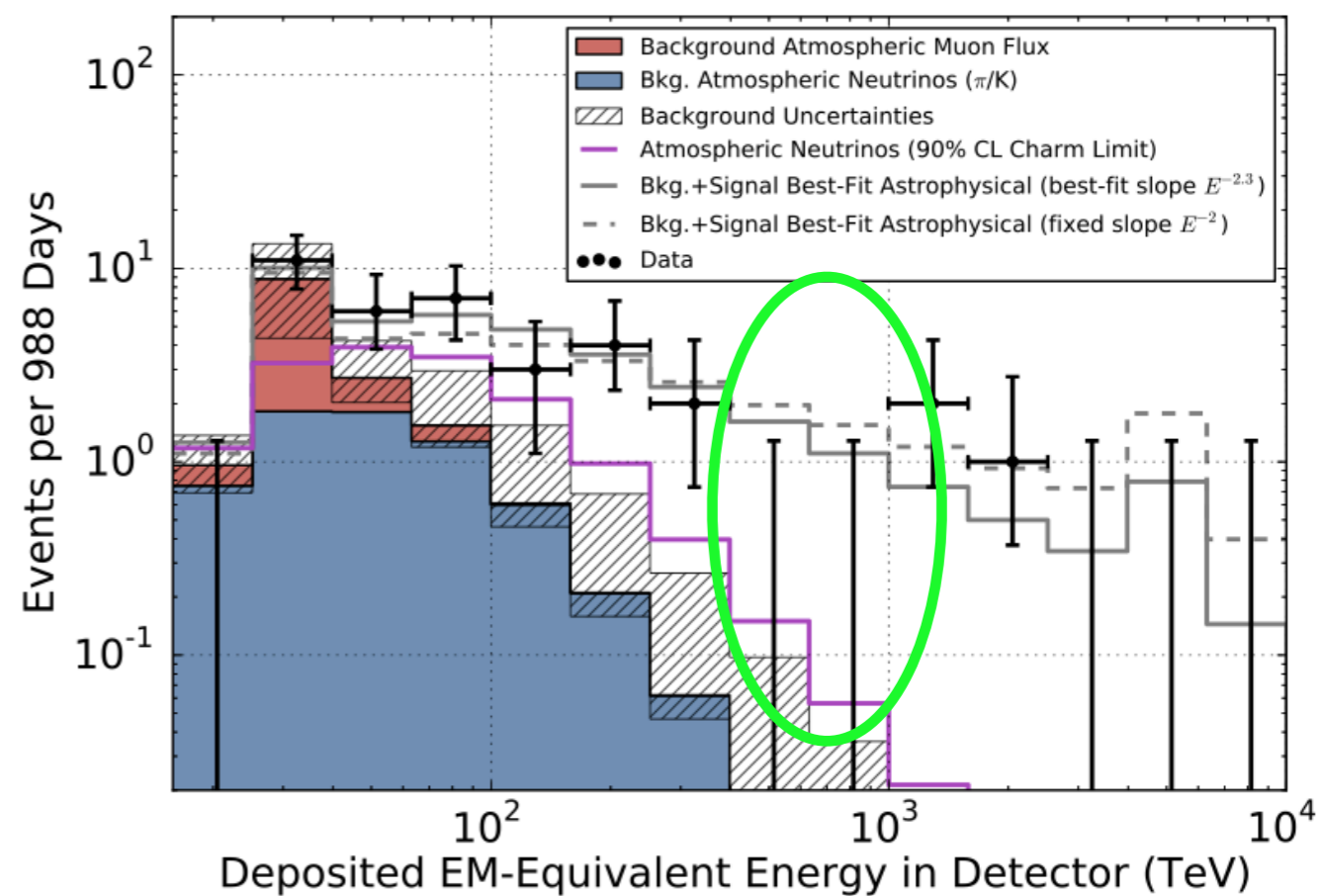
IceCube data

Problems with the astrophysical/atm interpretation of IceCube data

IceCube data

Problems with the astrophysical/atm interpretation of IceCube data

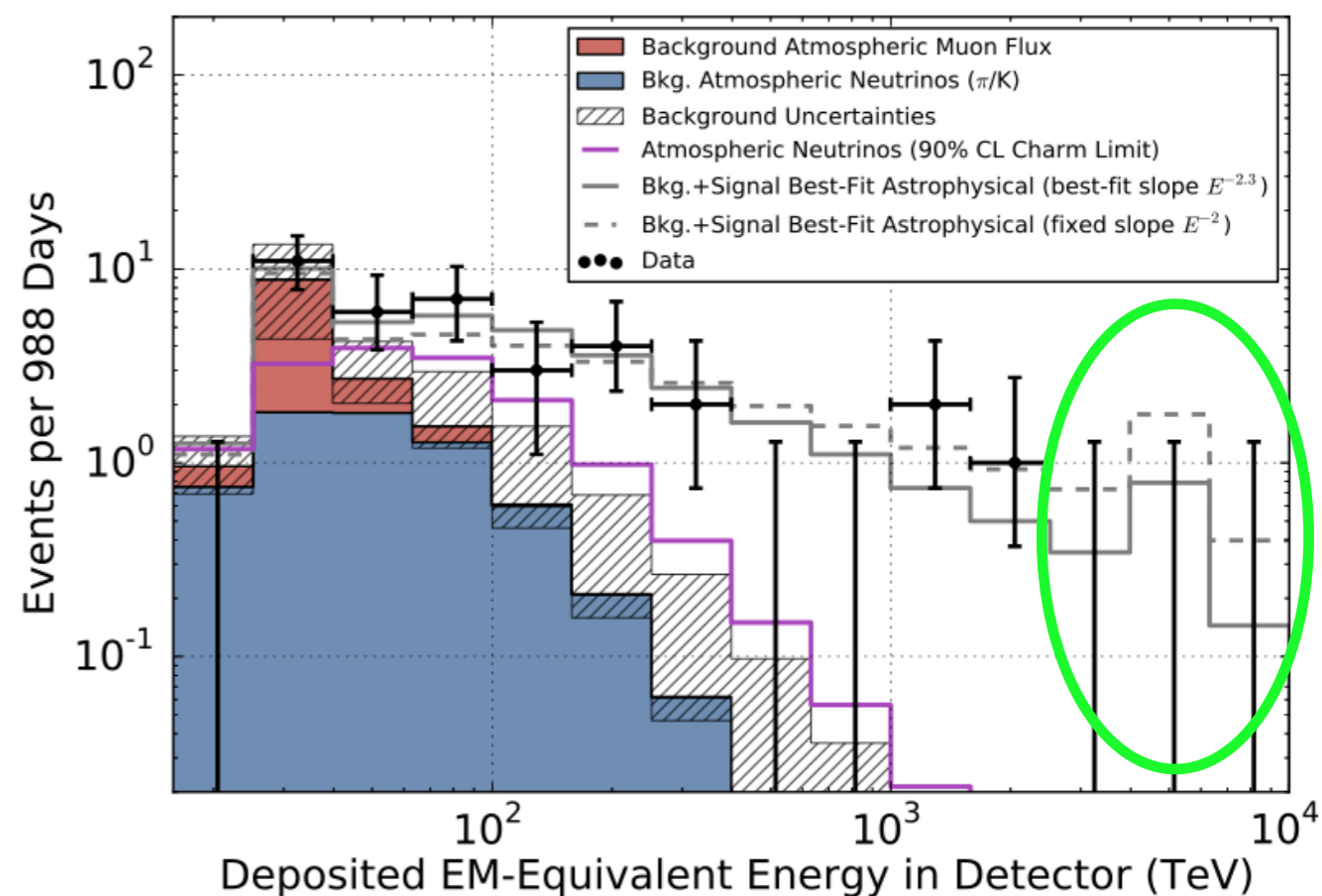
- ✓ deficit of events in the energy range $\sim (400 - 1000)$ TeV



IceCube data

Problems with the astrophysical/atm interpretation of IceCube data

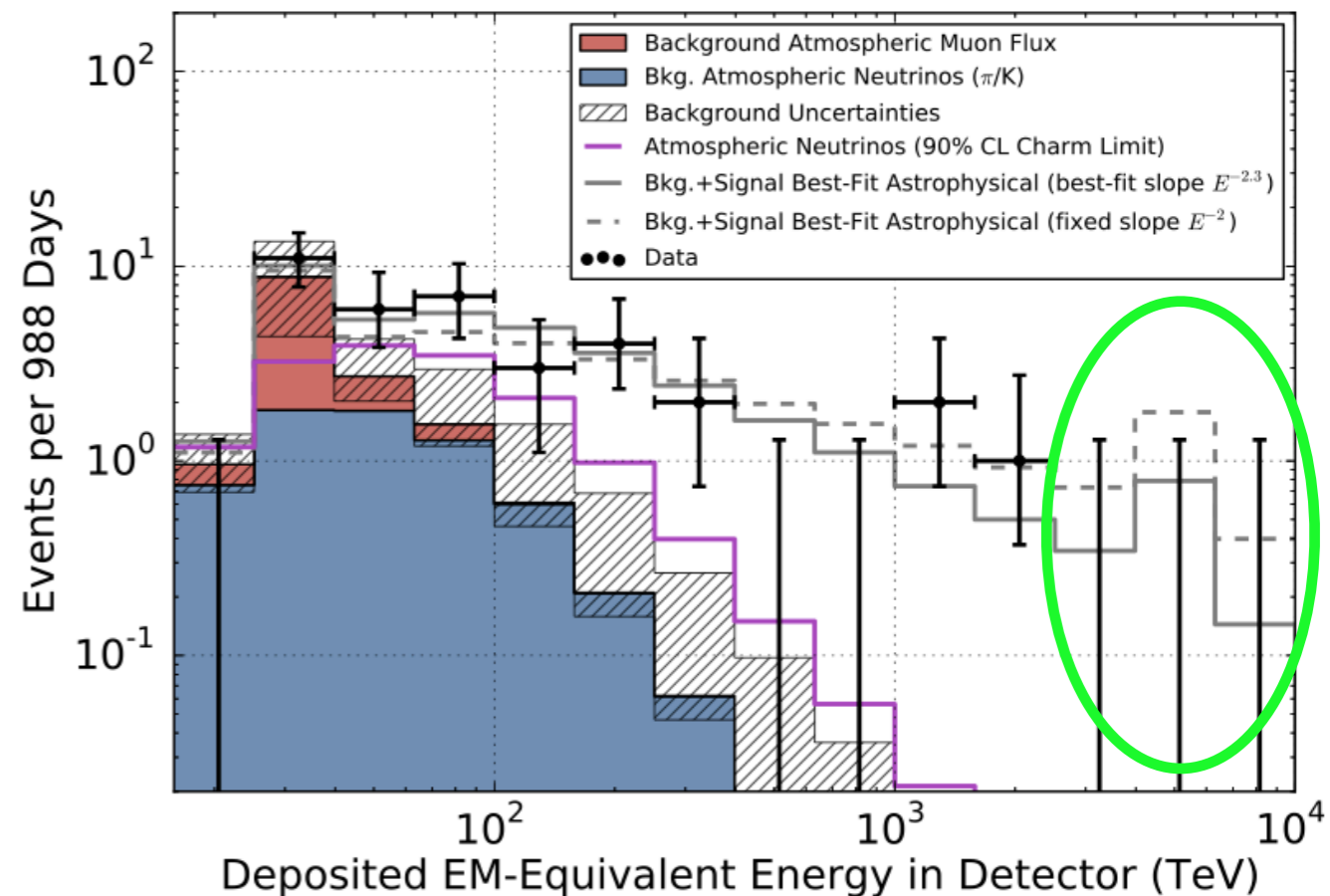
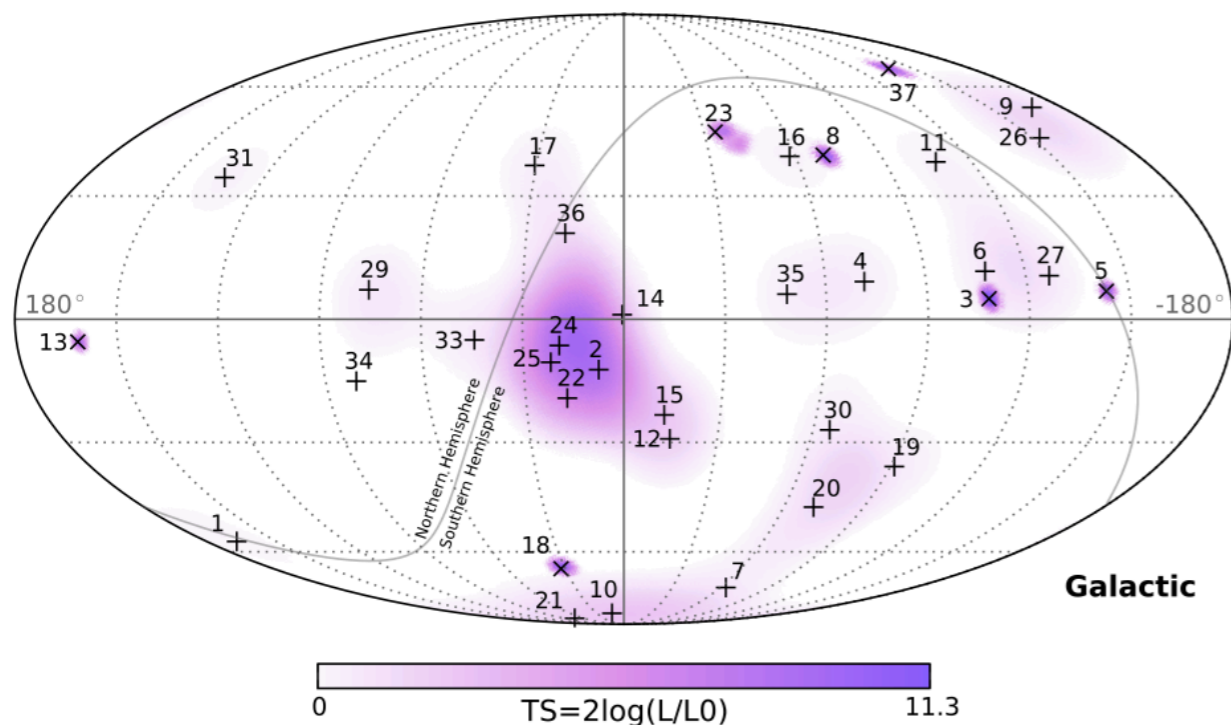
- ✓ deficit of events in the energy range $\sim (400 - 1000)$ TeV
- ✓ cut-off in events: no events observed with energy > 2 PeV



IceCube data

Problems with the astrophysical/atm interpretation of IceCube data

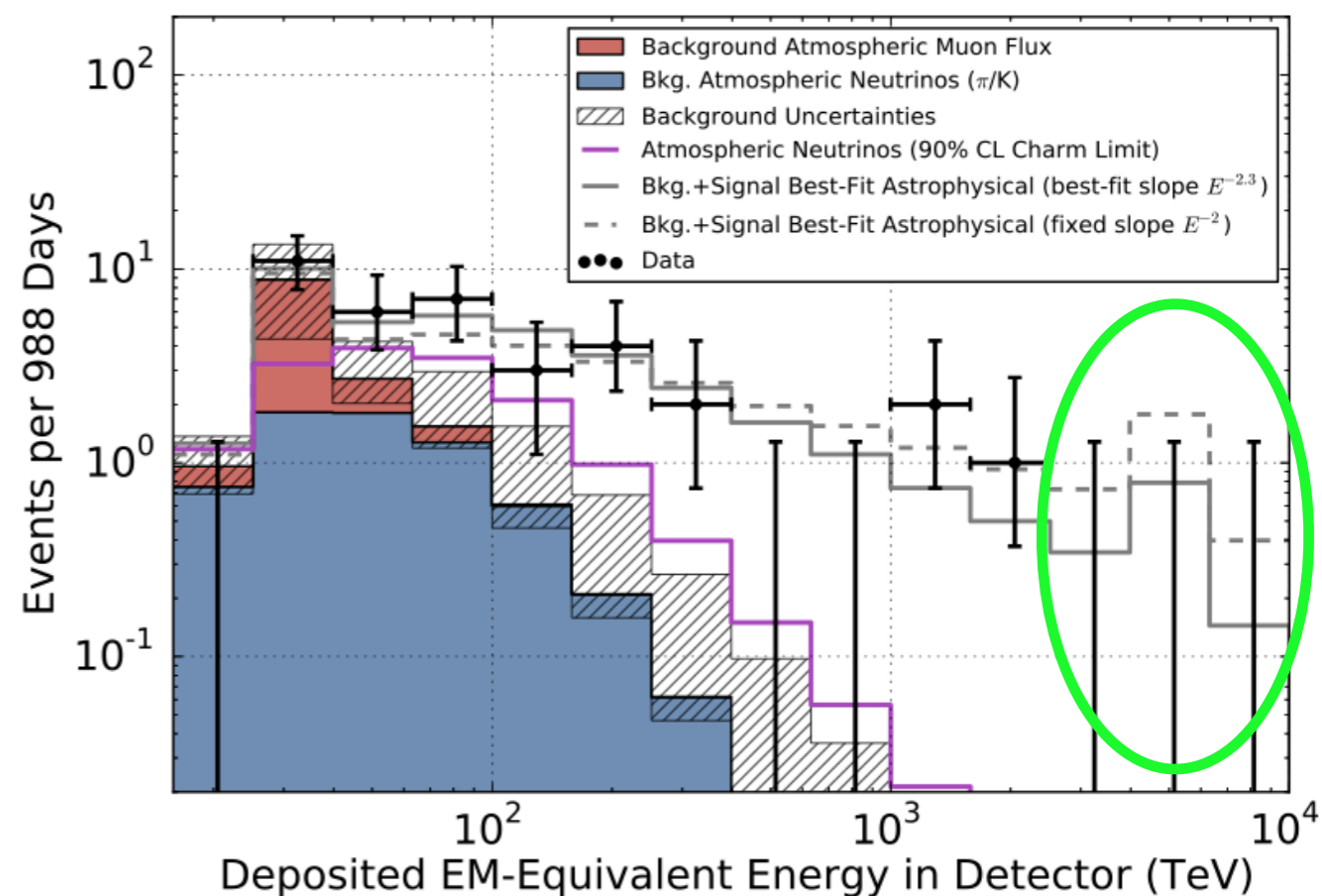
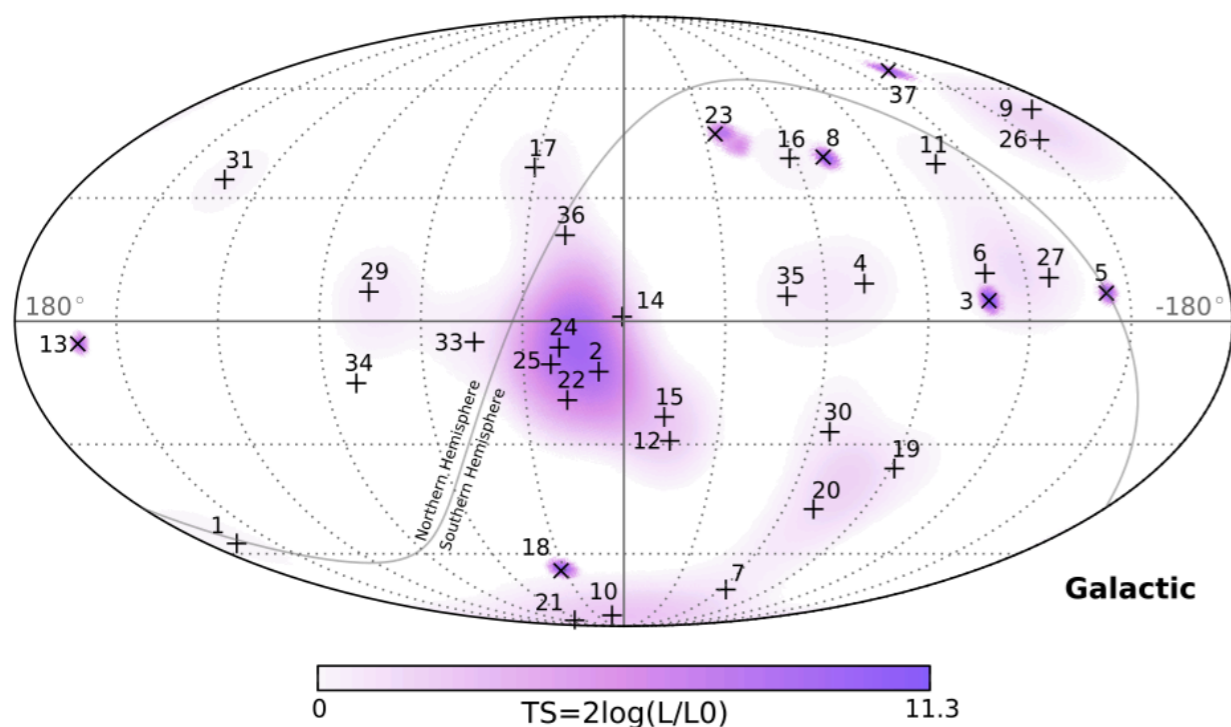
- ✓ deficit of events in the energy range $\sim (400 - 1000)$ TeV
- ✓ cut-off in events: no events observed with energy > 2 PeV
- ✓ angular distribution of events show mild anisotropies (enhanced toward GC)



IceCube data

Problems with the astrophysical/atm interpretation of IceCube data

- ✓ deficit of events in the energy range $\sim (400 - 1000)$ TeV
- ✓ cut-off in events: no events observed with energy > 2 PeV
- ✓ angular distribution of events show mild anisotropies (enhanced toward GC)
- ⚠ none of the above-mentioned issues are significant



Interpretations of IceCube data

✓ "Conventional" interpretations of IceCube data

Cosmic ray sources

GRBs

Galaxy clusters

Star-forming galaxies

AGNs

Fermi bubbles

Galactic Center activities

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M. D. Kistler, T. Stanev and H. Yuksel, arXiv:1301.1703 [astro-ph.HE]

K. Murase and K. Ioka, Phys. Rev. Lett. 111, no. 12, 121102 (2013) [arXiv:1306.2274 [astro-ph.HE]].

K. Murase, M. Ahlers and B. C. Lacki, Phys. Rev. D 88, no. 12, 121301 (2013) [arXiv:1306.3417 [astro-ph.HE]].

L. A. Anchordoqui, H. Goldberg, M. H. Lynch, A. V. Olinto, T. C. Paul and T. J. Weiler, arXiv:1306.5021 [astro-ph.HE].

R. Laha, J. F. Beacom, B. Dasgupta, S. Horiuchi and K. Murase, Phys. Rev. D 88, 043009 (2013) [arXiv:1306.2309 [astro-ph.HE]].

S. Razzaque, Phys. Rev. D 88, 081302 (2013) [arXiv:1309.2756 [astro-ph.HE]].

C. Y. Chen, P. S. Bhupal Dev and A. Soni, Phys. Rev. D 89, no. 3, 033012 (2014) [arXiv:1309.1764 [hep-ph]].

M. Ahlers and K. Murase, Phys. Rev. D 90, 023010 (2014) [arXiv:1309.4077 [astro-ph.HE]].

I. Tamborra, S. Ando and K. Murase, JCAP 1409, no. 09, 043 (2014) [arXiv:1404.1189 [astro-ph.HE]].

M. Kachelriess and S. Ostapchenko, Phys. Rev. D 90, 083002 (2014) [arXiv:1405.3797 [astro-ph.HE]].

M. Ahlers and F. Halzen, arXiv:1406.2160 [astro-ph.HE].

Y. Bai, A. J. Barger, V. Barger, R. Lu, A. D. Peterson and J. Salvado, Phys. Rev. D 90, 063012 (2014) [arXiv:1407.2243 [astro-ph.HE]].

A. Bhattacharya, R. Enberg, M. H. Reno and I. Sarcevic, arXiv:1407.2985 [astro-ph.HE].

C. Lunardini, S. Razzaque, K. T. Theodoseou and L. Yang, Phys. Rev. D 90, 023016 (2014) [arXiv:1311.7188 [astro-ph.HE]].

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For a review

L. A. Anchordoqui, V. Barger, I. Cholis, H. Goldberg, D. Hooper, A. Kusenko, J. G. Learned and D. Marfatia et al., Journal of High Energy Astrophysics 1-2, 1 (2014) [arXiv:1312.6587 [astro-ph.HE]].

Interpretations of IceCube data

✓ "New Physics" interpretations of IceCube data

Lepto-quarks

Y. Ema, R. Jinno and T. Moroi, Phys. Lett. B 733, 120 (2014) [arXiv:1312.3501 [hep-ph]].

Secret neutrino interactions

K. Ioka and K. Murase, PTEP 2014, (2014) [arXiv:1404.2279 [astro-ph.HE]].

K. C. Y. Ng and J. F. Beacom, Phys. Rev. D 90, 065035 (2014) [arXiv:1404.2288 [astro-ph.HE]].

resonant absorption on
cosmic neutrino background

M. Ibe and K. Kaneta, Phys. Rev. D 90, 053011 (2014) [arXiv:1407.2848 [hep-ph]].

V. Barger and W. Y. Keung, Phys. Lett. B 727, 190 (2013) [arXiv:1305.6907 [hep-ph]].

B. Feldstein, A. Kusenko, S. Matsumoto and T. T. Yanagida, Phys. Rev. D 88, no. 1, 015004 (2013) [arXiv:1303.7320 [hep-ph]].

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particles

Y. Bai, R. Lu and J. Salvado, arXiv:1311.5864 [hep-ph].

A. Bhattacharya, M. H. Reno and I. Sarcevic, JHEP 1406, 110 (2014) [arXiv:1403.1862 [hep-ph]].

J. Zavala, Phys. Rev. D 89, 123516 (2014) [arXiv:1404.2932 [astro-ph.HE]].

A. Bhattacharya, R. Gandhi and A. Gupta, arXiv:1407.3280 [hep-ph].

C. Rott, K. Kohri and S. C. Park, arXiv:1408.4575 [hep-ph].

T. Higaki, R. Kitano and R. Sato, JHEP 1407, 044 (2014) [arXiv:1405.0013 [hep-ph]].

Dark matter decay

A. Esmaili and P. D. Serpico, JCAP 1311, 054 (2013) [arXiv:1308.1105 [hep-ph]].

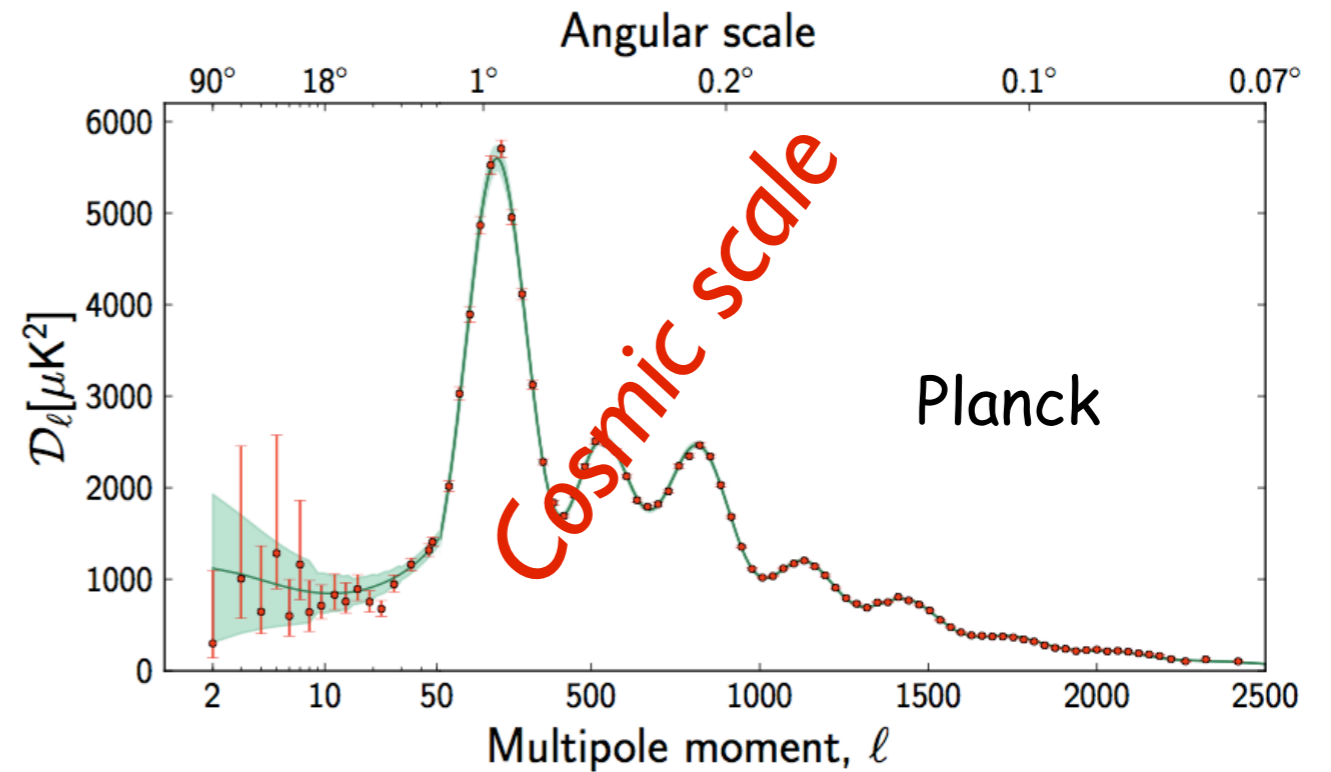
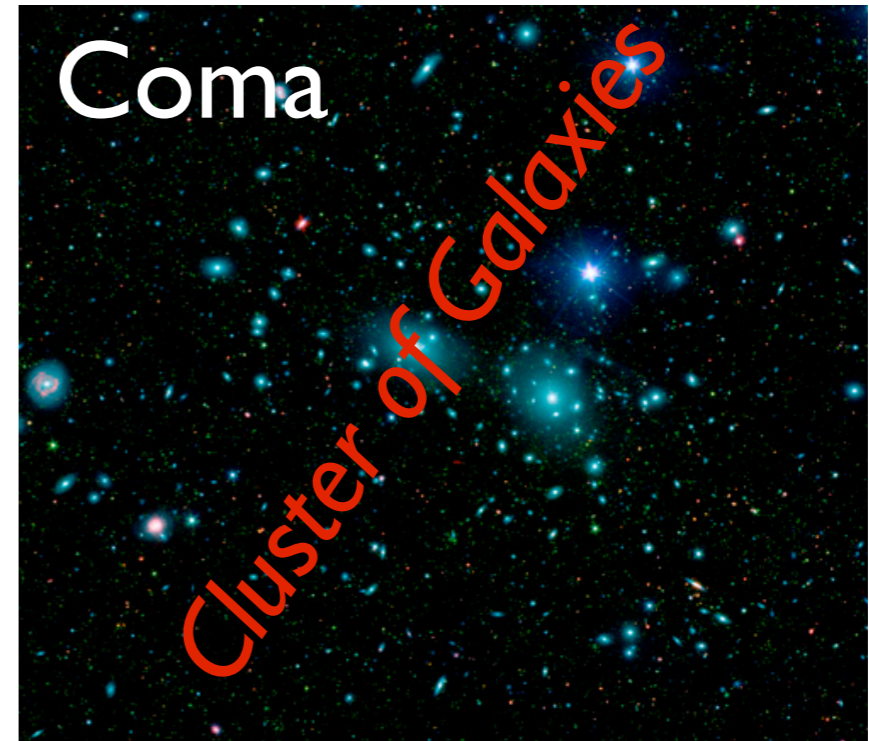
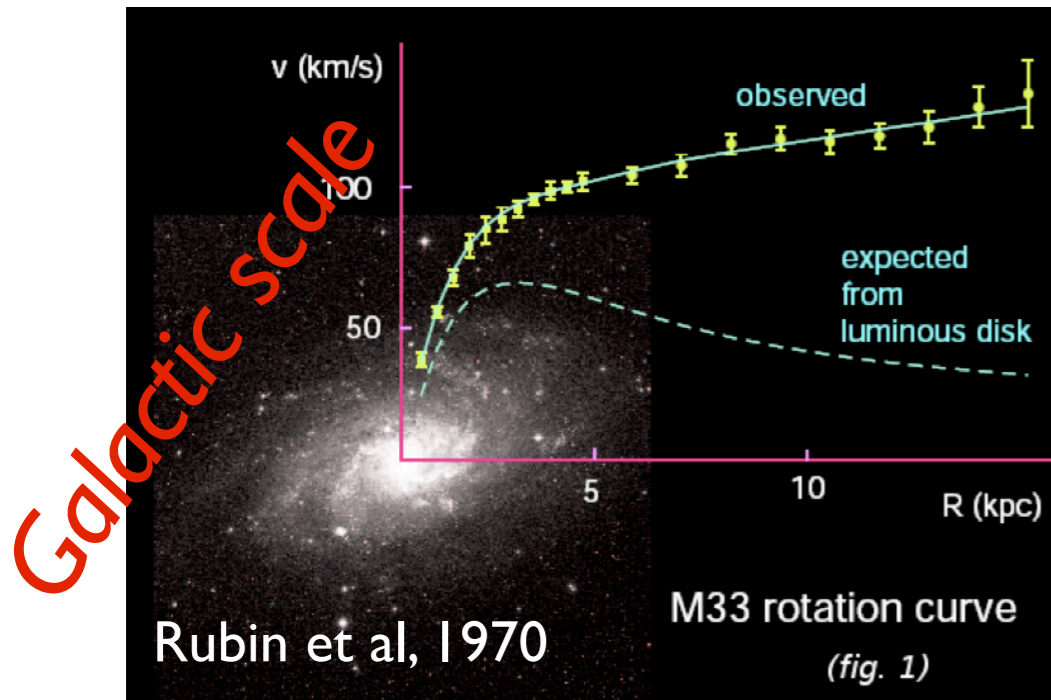
A. Esmaili, S. K. Kang and P. D. Serpico, arXiv:1410.5979 [hep-ph].

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A note on Dark Matter

DM exist!



A note on Dark Matter

DM exist!

What We Know?

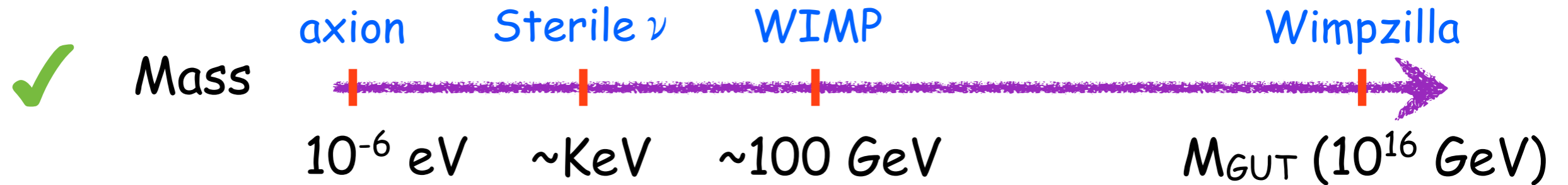
- ✓ Non Baryonic
- ✓ No electric and color charges
- ✓ Cold (or perhaps warm)
- ✓ Long lived (not necessarily stable)

All of these come from gravitational effects

A note on Dark Matter

DM exist!

What We Do Not Know?



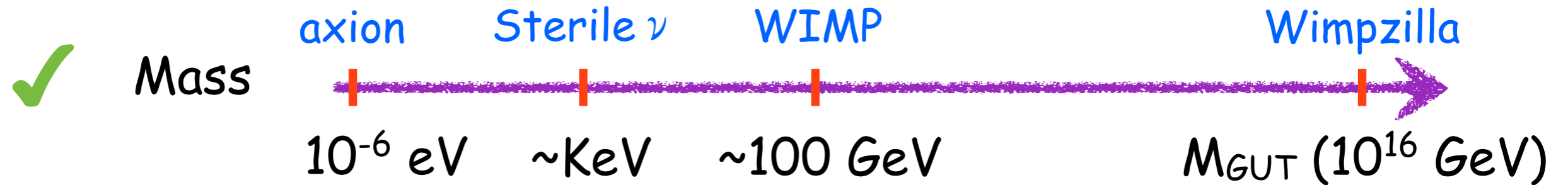
⚠ "WIMP" paradigm ?

Note that WIMP paradigm is a "particle physics" conjecture, needs to be validated at colliders

A note on Dark Matter

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What We Do Not Know?



⚠ "WIMP" paradigm ?

Note that WIMP paradigm is a "particle physics" conjecture, needs to be validated at colliders

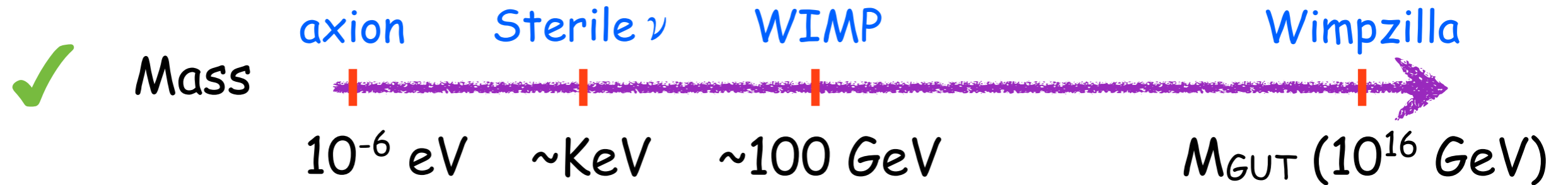
caution: street light effect



A note on Dark Matter

DM exist!

What We Do Not Know?



⚠ "WIMP" paradigm ?

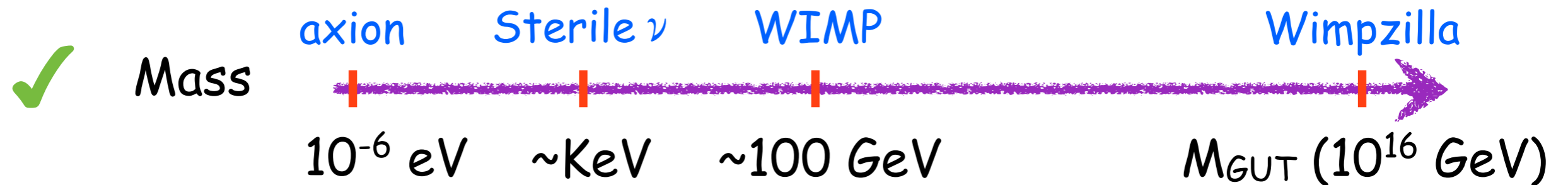
Note that WIMP paradigm is a "particle physics" conjecture, needs to be validated at colliders



A note on Dark Matter

DM exist!

What We Do Not Know?



"WIMP" paradigm ?

Note that WIMP paradigm is a "particle physics" conjecture, needs to be validated at colliders



Lifetime: stable (∞) or

$T_{DM} > 4.3 \times 10^{17}$ s (age of Universe)

$T_{DM} > 2.2 \times 10^{19}$ s (CMB) Y. Gong and X. Chen, PRD77 (2008), arXiv:0802.2296



Possible decay and/or annihilation channels



...

Limits on lifetime from neutrino experiments before recent IceCube data

$16 - 2.5 \times 10^3 \text{ TeV}$

AMANDA

$340 - 2 \times 10^5 \text{ TeV}$

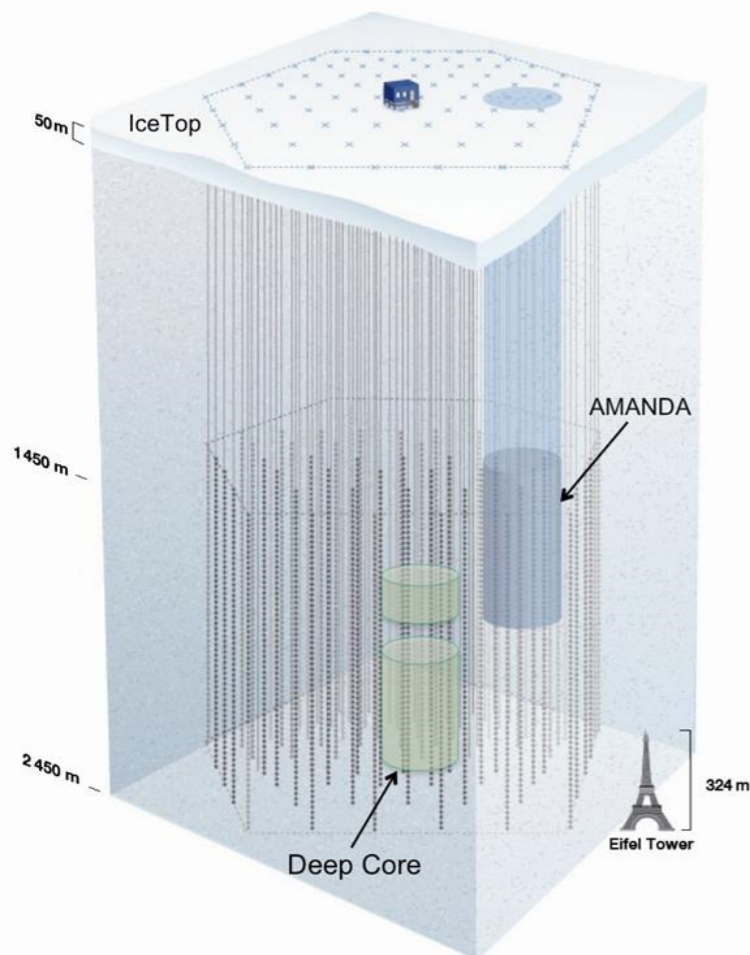
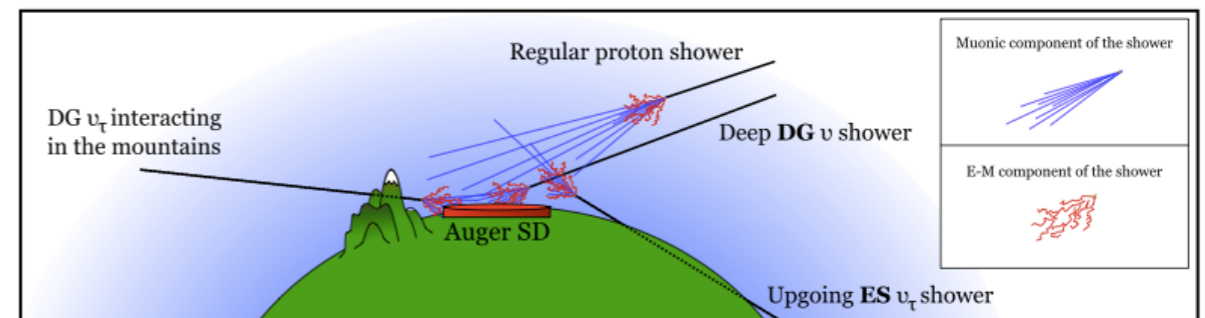
IceCube-22

$3 \times 10^3 - 6.3 \times 10^6 \text{ TeV}$

IceCube-40

$10^5 - 10^8 \text{ TeV}$

Auger



$10^6 - 3.2 \times 10^{11} \text{ TeV}$

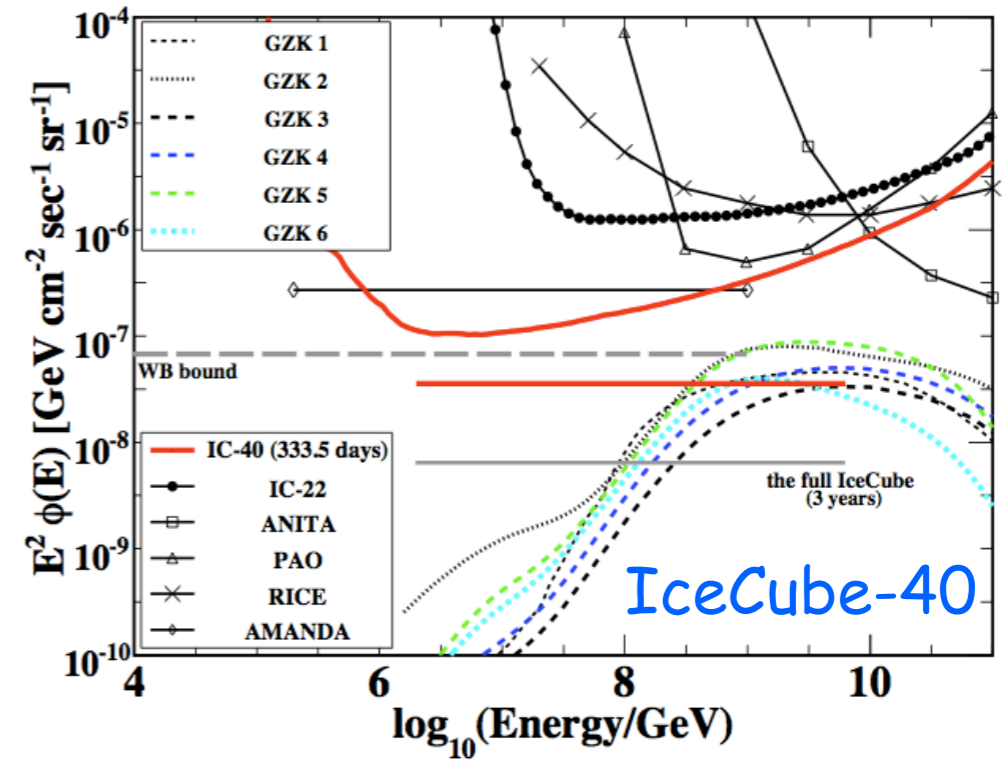
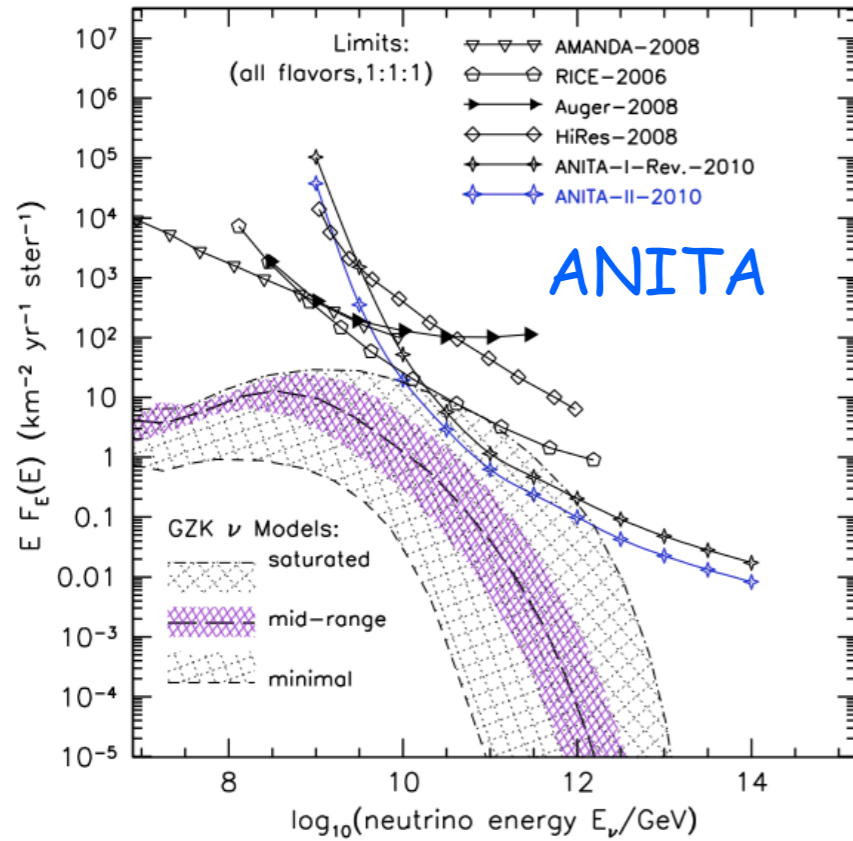
ANITA



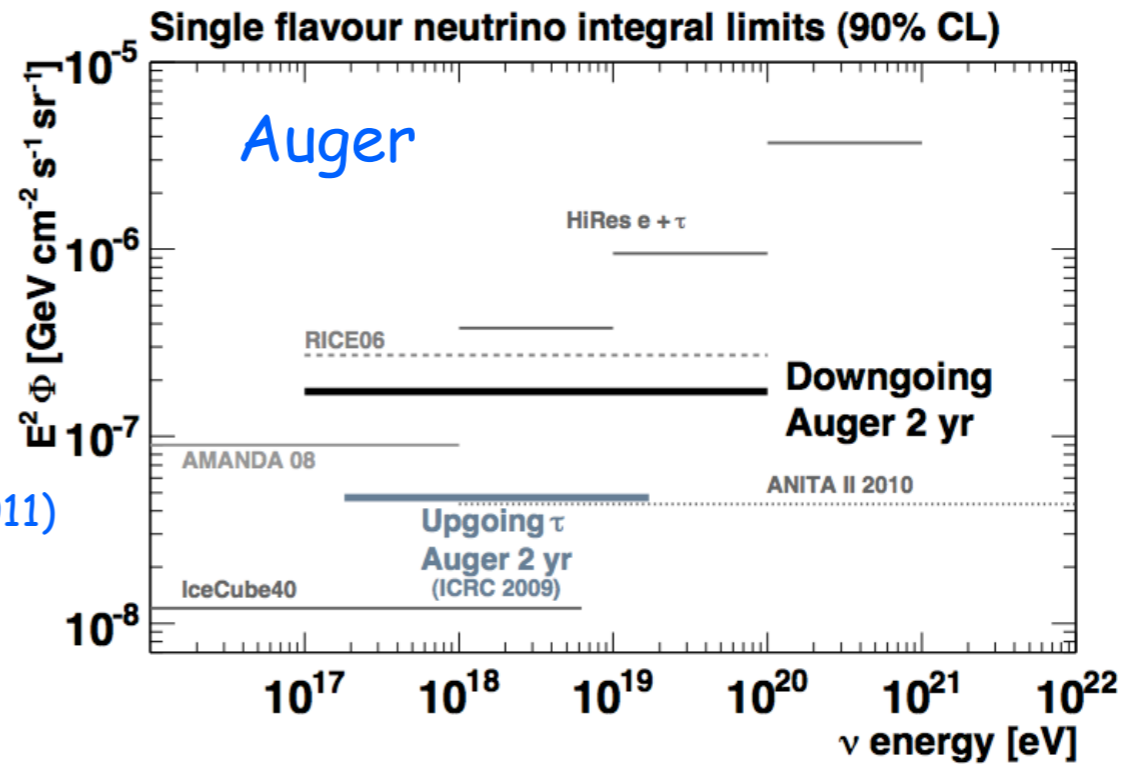
Experiments

R. Abbasi et al. [IceCube Collaboration], Phys.Rev.D83 (2011)
arXiv:1103.4250

P. Gorham et al. [ANITA Collaboration], Phys.Rev.D85 (2012)
arXiv:1011.5004, arXiv:1003.2961



P. Abreu et al. [Pierre Auger Collaboration], Phys.Rev.D84 (2011)
arXiv:1202.1493

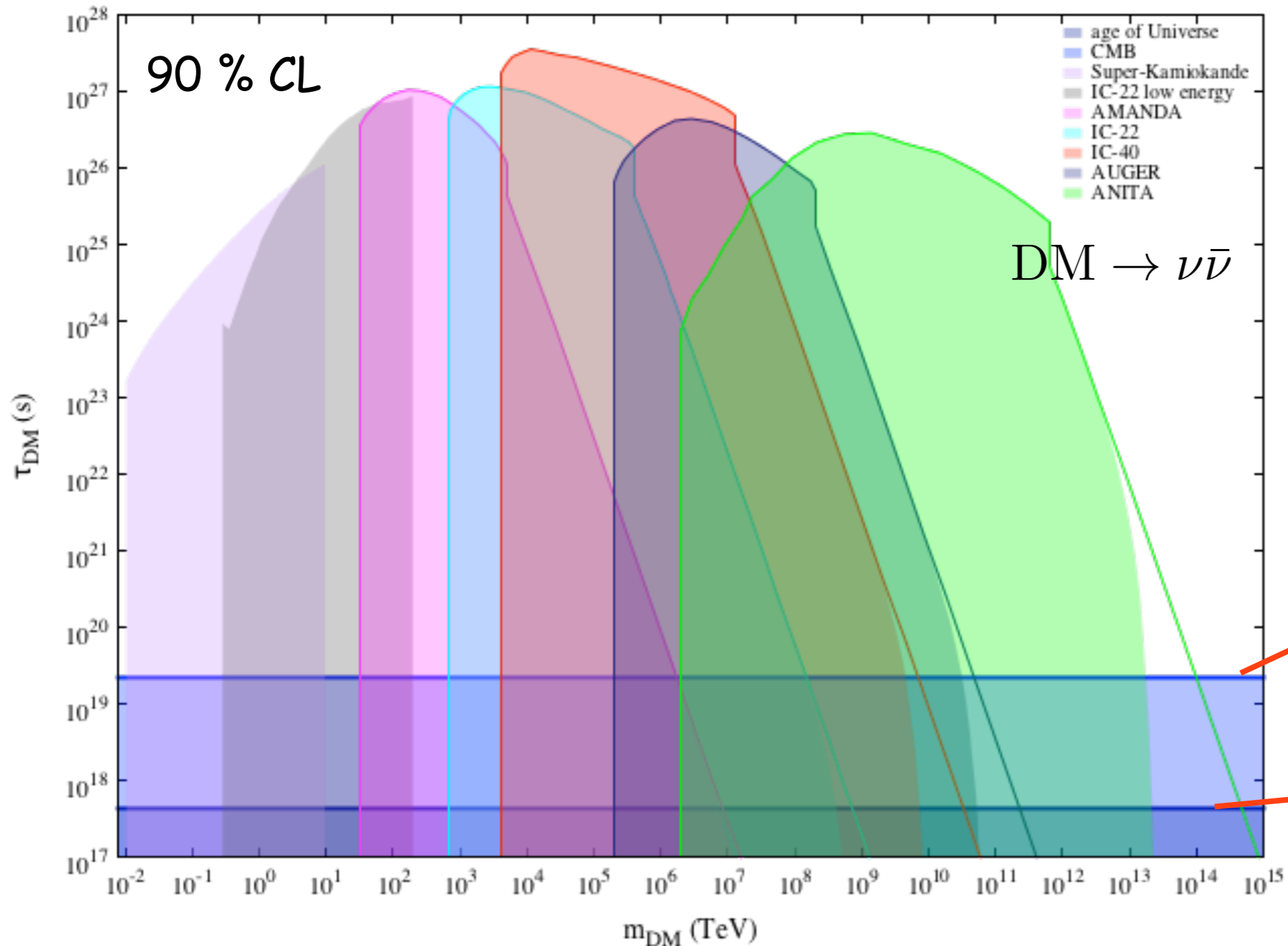


Limits on lifetime from neutrino experiments

before recent IceCube data

✓ Lifetime: stable (∞) or

A.E., Alejandro Ibarra and Orlando L. G. Peres
JCAP (2012) [arXiv: 1205.5281]



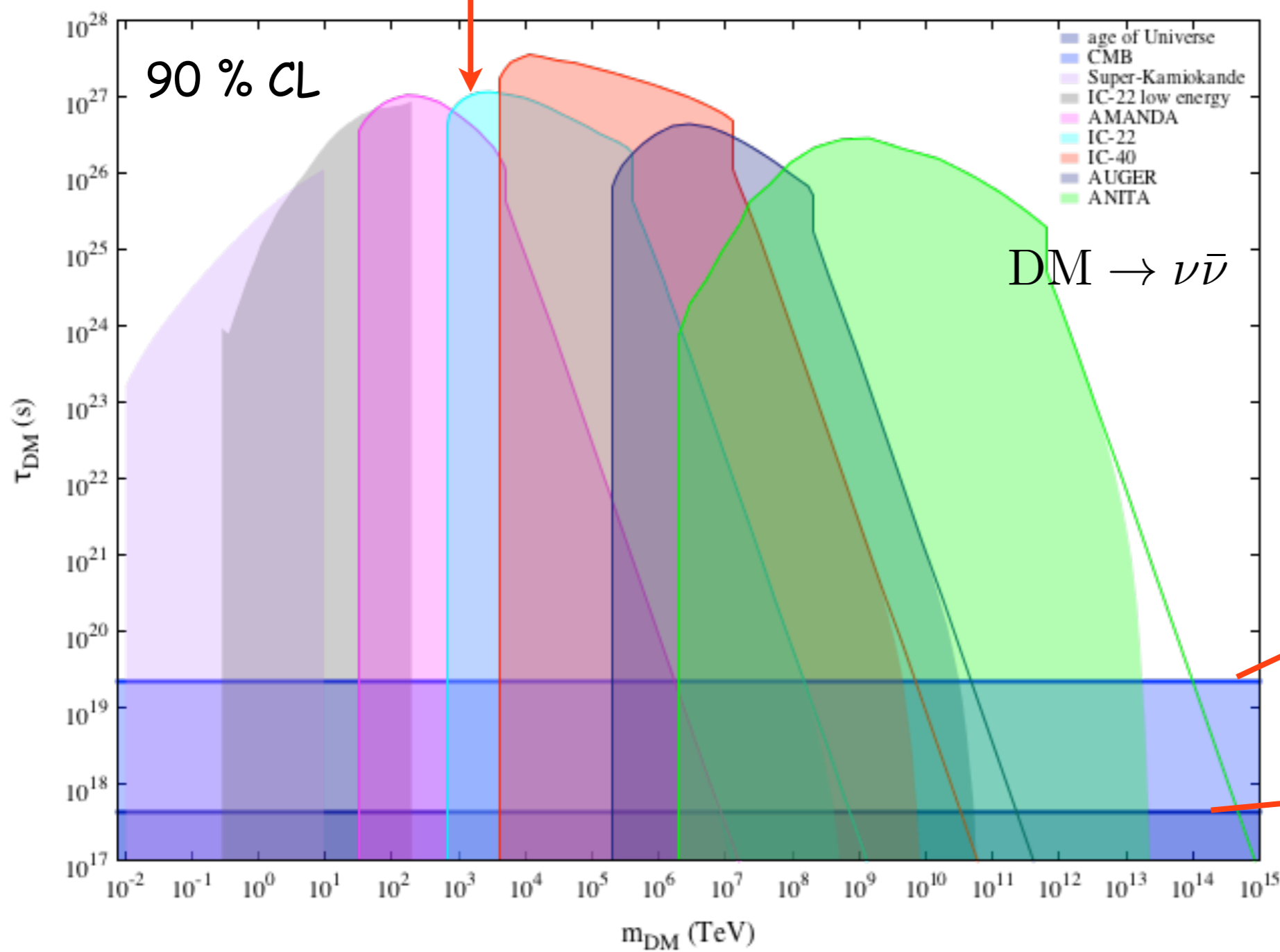
Using AMANDA, IceCube,
Auger and ANITA data

Limits on lifetime from neutrino experiments

before recent IceCube data

✓ Lifetime: stable (∞) or
this talk

A.E., Alejandro Ibarra and Orlando L. G. Peres
JCAP (2012) [arXiv: 1205.5281]



Using AMANDA, IceCube,
Auger and ANITA data

$T_{DM} > 2.2 \times 10^{19} \text{ s}$ (CMB)

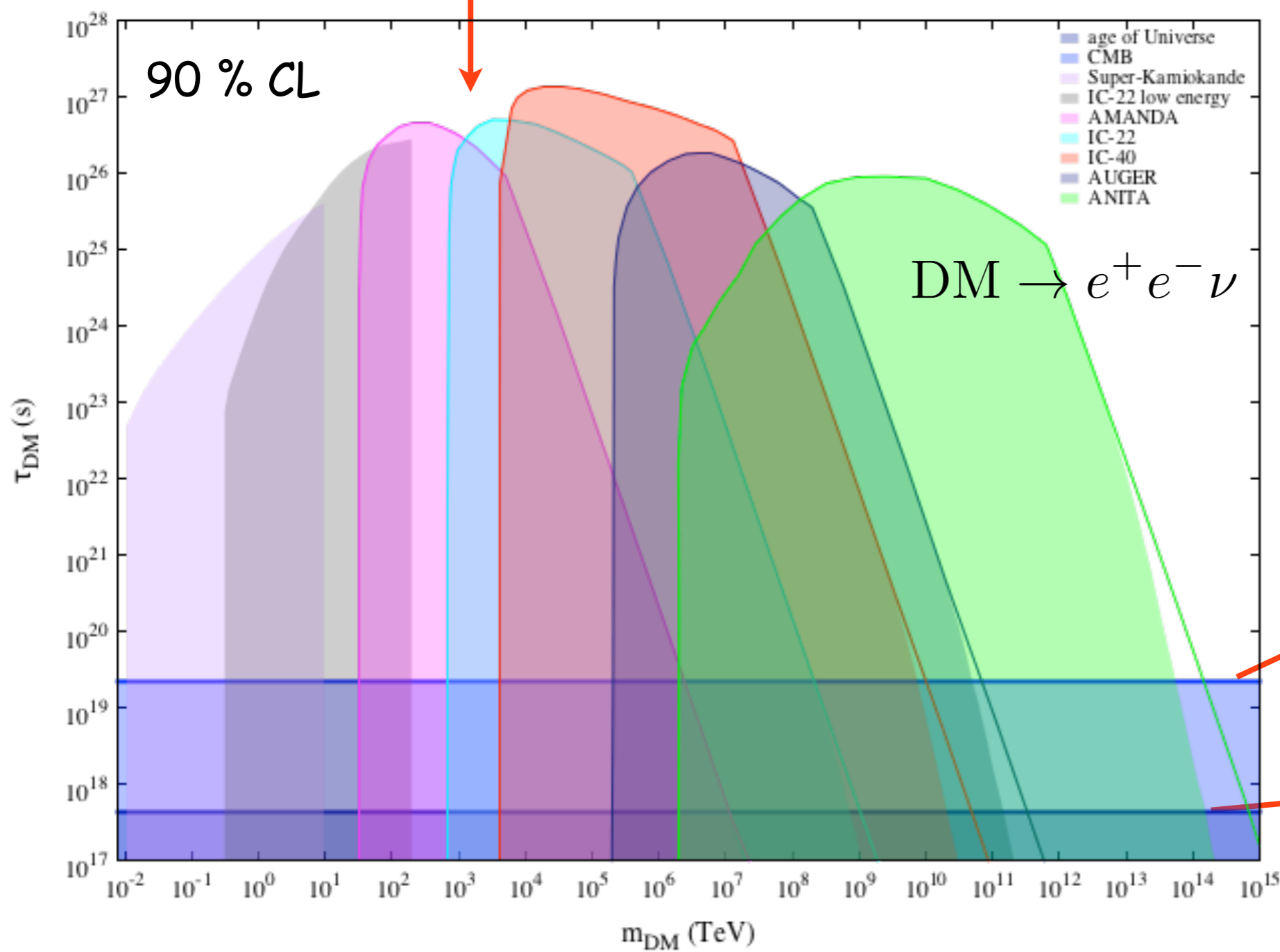
$T_{DM} > 4.3 \times 10^{17} \text{ s}$
(age of Universe)

Limits on lifetime from neutrino experiments

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✓ Lifetime: stable (∞) or
this talk

A.E., Alejandro Ibarra and Orlando L. G. Peres
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Using AMANDA, IceCube, Auger and ANITA data

$T_{DM} > 2.2 \times 10^{19} \text{ s}$ (CMB)

$T_{DM} > 4.3 \times 10^{17} \text{ s}$ (age of Universe)

Interpreting the IceCube events by decaying dark matter

Two main diagnostics:

- ✓ Energy distribution
- ✓ Angular distribution

Energy distribution of neutrinos from decaying DM

✓ Galactic contribution:

$$\frac{dJ_h}{dE_\nu}(l, b) = \frac{1}{4\pi m_{\text{DM}} \tau_{\text{DM}}} \frac{dN_\nu}{dE_\nu} \int_0^\infty ds \rho_h[r(s, l, b)]$$

NFW $\rho_{\text{halo}}(r) \simeq \frac{\rho_h}{r/r_c(1+r/r_c)^2}$

$$r(s, l, b) = \sqrt{s^2 + R_\odot^2 - 2sR_\odot \cos b \cos l}$$

Energy distribution of neutrinos from decaying DM

✓ Galactic contribution:

$$\frac{dJ_h}{dE_\nu}(l, b) = \frac{1}{4\pi m_{\text{DM}} \tau_{\text{DM}}} \frac{dN_\nu}{dE_\nu} \int_0^\infty ds \rho_h[r(s, l, b)]$$

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$r(s, l, b) = \sqrt{s^2 + R_\odot^2 - 2sR_\odot \cos b \cos l}$

✓ extragalactic contribution:

$$\frac{dJ_{\text{eg}}}{dE_\nu} = \frac{\Omega_{\text{DM}} \rho_c}{4\pi m_{\text{DM}} \tau_{\text{DM}}} \int_0^\infty dz \frac{1}{H(z)} \frac{dN_\nu}{dE_\nu} [(1+z)E_\nu]$$

Energy distribution of neutrinos from decaying DM

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NFW $\rho_{\text{halo}}(r) \simeq \frac{\rho_h}{r/r_c(1+r/r_c)^2}$

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M. Cirelli et. al., JCAP (2011)

energy spectrum of neutrinos
at production point
(including the EW corrections)

quarks

$$\frac{dN_\nu}{dE_\nu} = (1 - b_H) \left. \frac{dN_\nu}{dE_\nu} \right|_S + b_H \left. \frac{dN_\nu}{dE_\nu} \right|_H$$

neutrinos,
charged leptons

Energy distribution of neutrinos from decaying DM

✓ Galactic contribution:

$$\frac{dJ_h}{dE_\nu}(l, b) = \frac{1}{4\pi m_{\text{DM}} \tau_{\text{DM}}} \frac{dN_\nu}{dE_\nu} \int_0^\infty ds \rho_h[r(s, l, b)]$$

NFW $\rho_{\text{halo}}(r) \simeq \frac{\rho_h}{r/r_c(1+r/r_c)^2}$

$$r(s, l, b) = \sqrt{s^2 + R_\odot^2 - 2sR_\odot \cos b \cos l}$$

✓ extragalactic contribution:

$$\frac{dJ_{\text{eg}}}{dE_\nu} = \frac{\Omega_{\text{DM}} \rho_c}{4\pi m_{\text{DM}} \tau_{\text{DM}}} \int_0^\infty dz \frac{1}{H(z)} \frac{dN_\nu}{dE_\nu} [(1+z)E_\nu]$$

M. Cirelli et. al., JCAP (2011)

energy spectrum of neutrinos
at production point
(including the EW corrections)

quarks

neutrinos,
charged leptons

$$\frac{dN_\nu}{dE_\nu} = (1 - b_H) \left. \frac{dN_\nu}{dE_\nu} \right|_S + b_H \left. \frac{dN_\nu}{dE_\nu} \right|_H$$

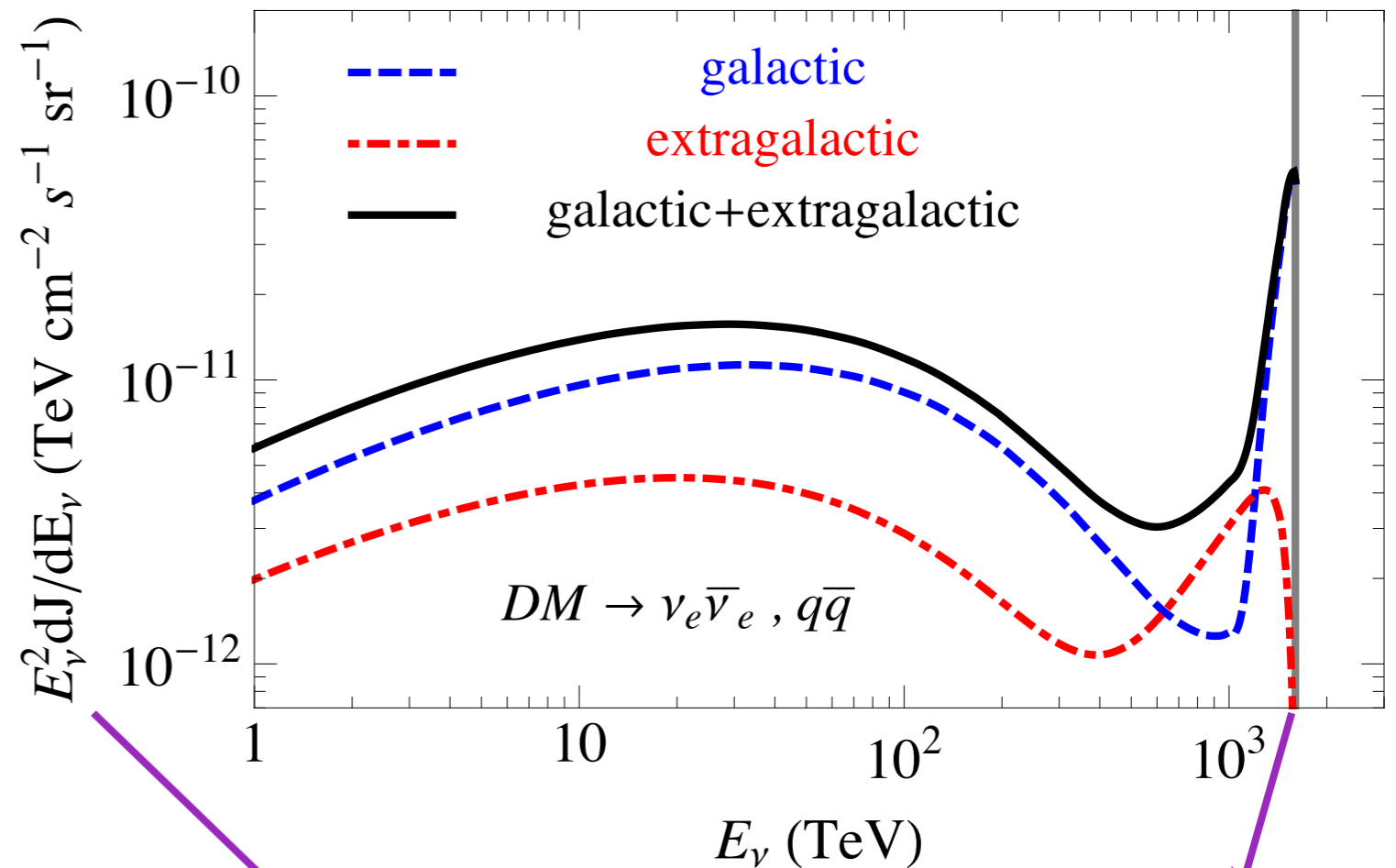
at the Earth $\begin{pmatrix} J_e \\ J_\mu \\ J_\tau \end{pmatrix} = \begin{pmatrix} P_{ee} & P_{e\mu} & P_{e\tau} \\ & P_{\mu\mu} & P_{\mu\tau} \\ & & P_{\tau\tau} \end{pmatrix} \begin{pmatrix} I_e \\ I_\mu \\ I_\tau \end{pmatrix}$ production point

decoherent oscillation

Flux of neutrinos from decaying DM

✓ an example:

A. E., Pasquale D. Serpico,
JCAP (2013) [arXiv:1308.1105]



$(U_e + U_\mu + U_\tau)/3$

$m_{DM}/2 = 1.6 \text{ PeV}$

$b_H = 0.12$ and $\tau_{DM} = 2 \times 10^{27} \text{ s}$

Flux of neutrinos from decaying DM

✓ an example:

intriguing features:

a cut-off at $m_{\text{DM}}/2$

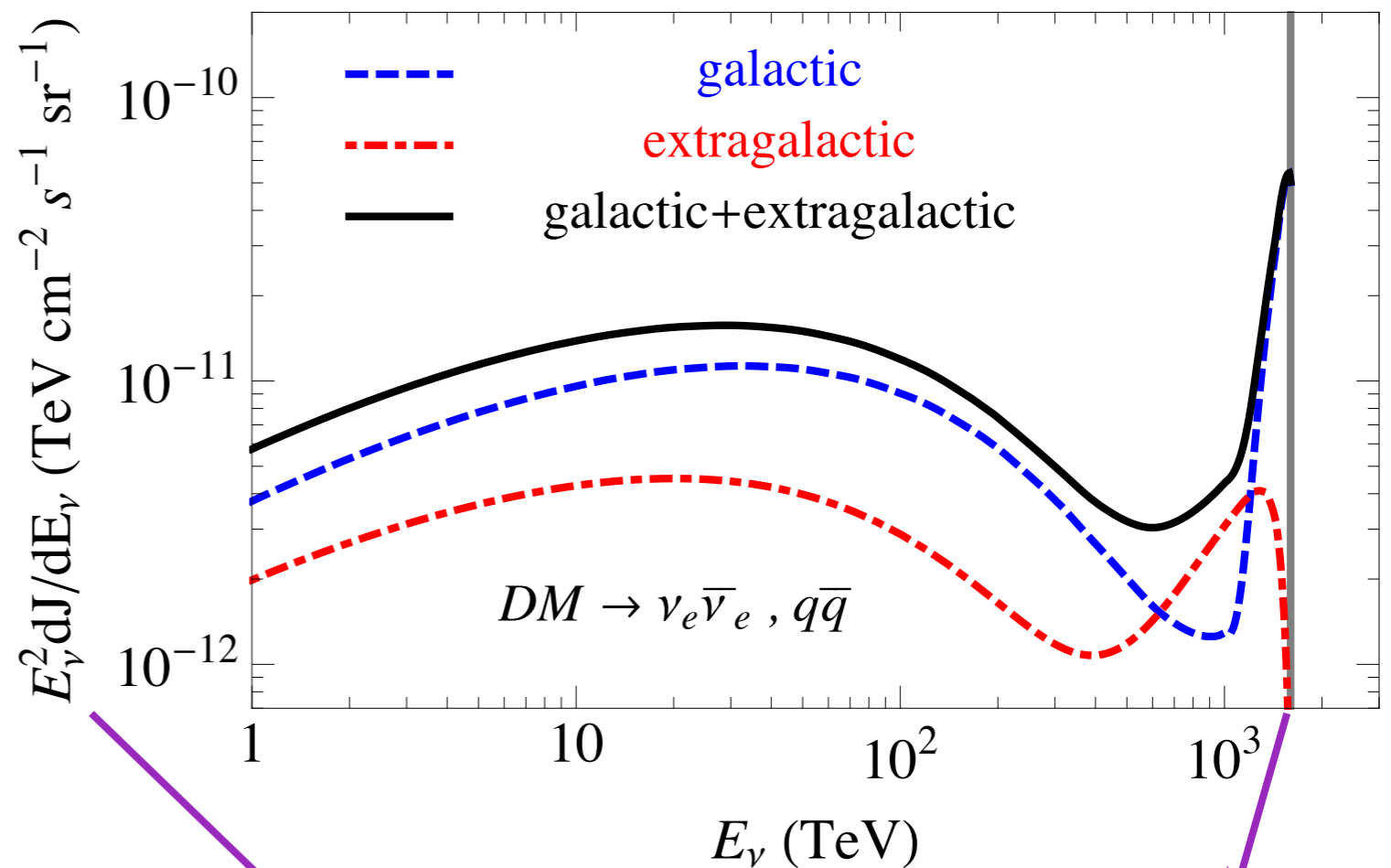
a peak in $\sim \text{PeV}$

a dip in $\sim (0.4-1) \text{ PeV}$

populated spectrum in $< 0.4 \text{ PeV}$

due to soft channel and EW cascades

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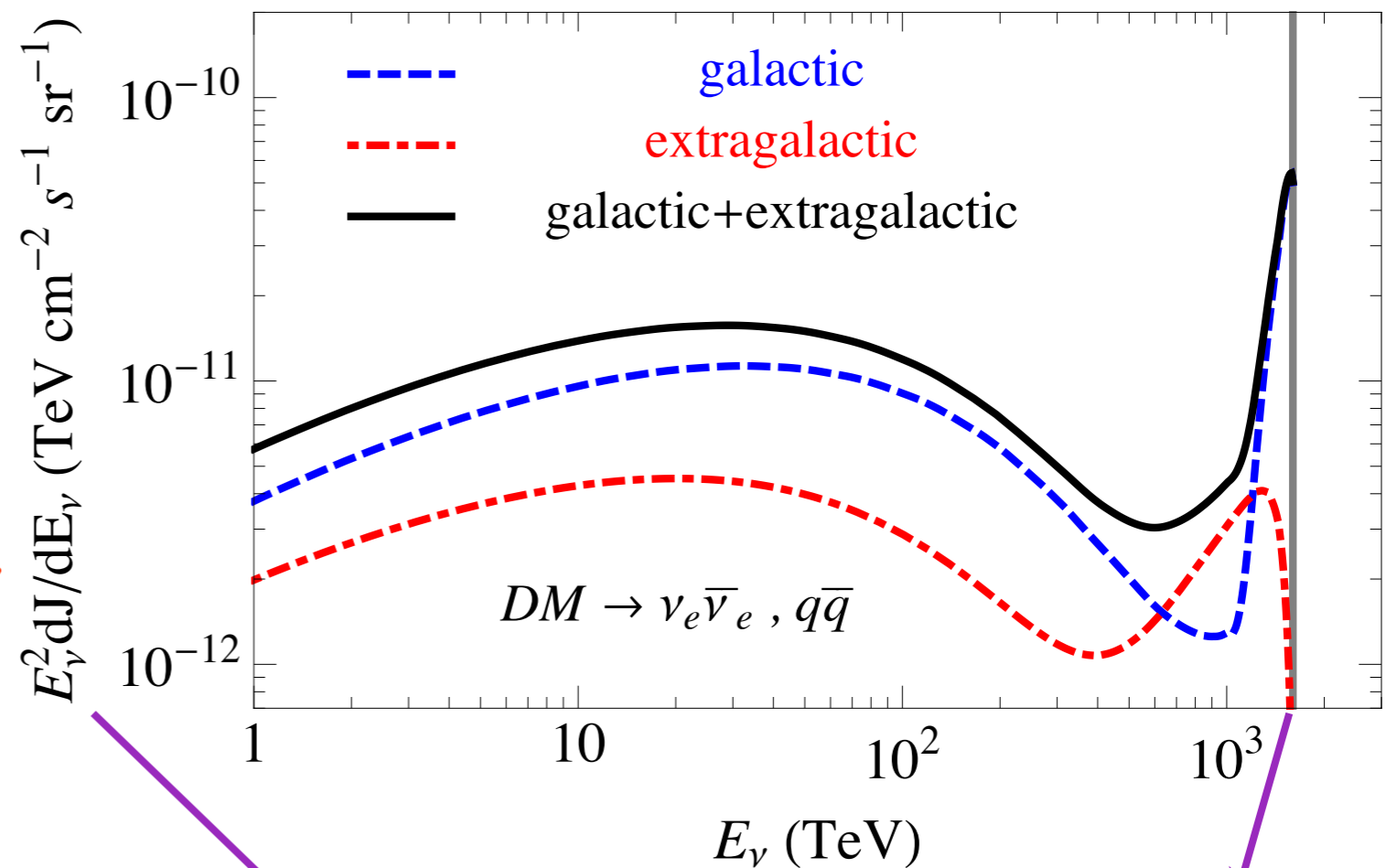
populated spectrum in $< 0.4 \text{ PeV}$

due to soft channel and EW cascades

b_H controls the peak height at $\sim \text{PeV}$

τ_{DM} controls the low energy population

A. E., Pasquale D. Serpico,
JCAP (2013) [arXiv:1308.1105]



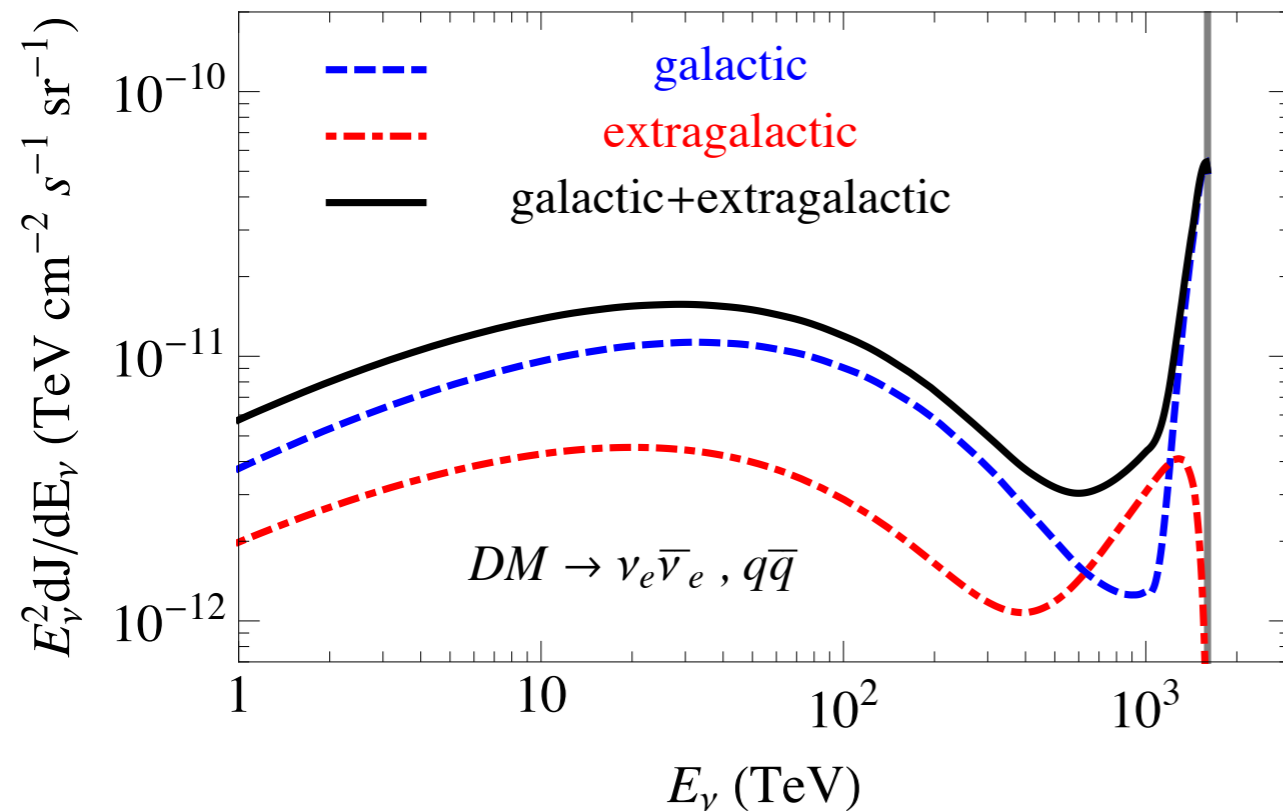
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Flux of neutrinos from decaying DM

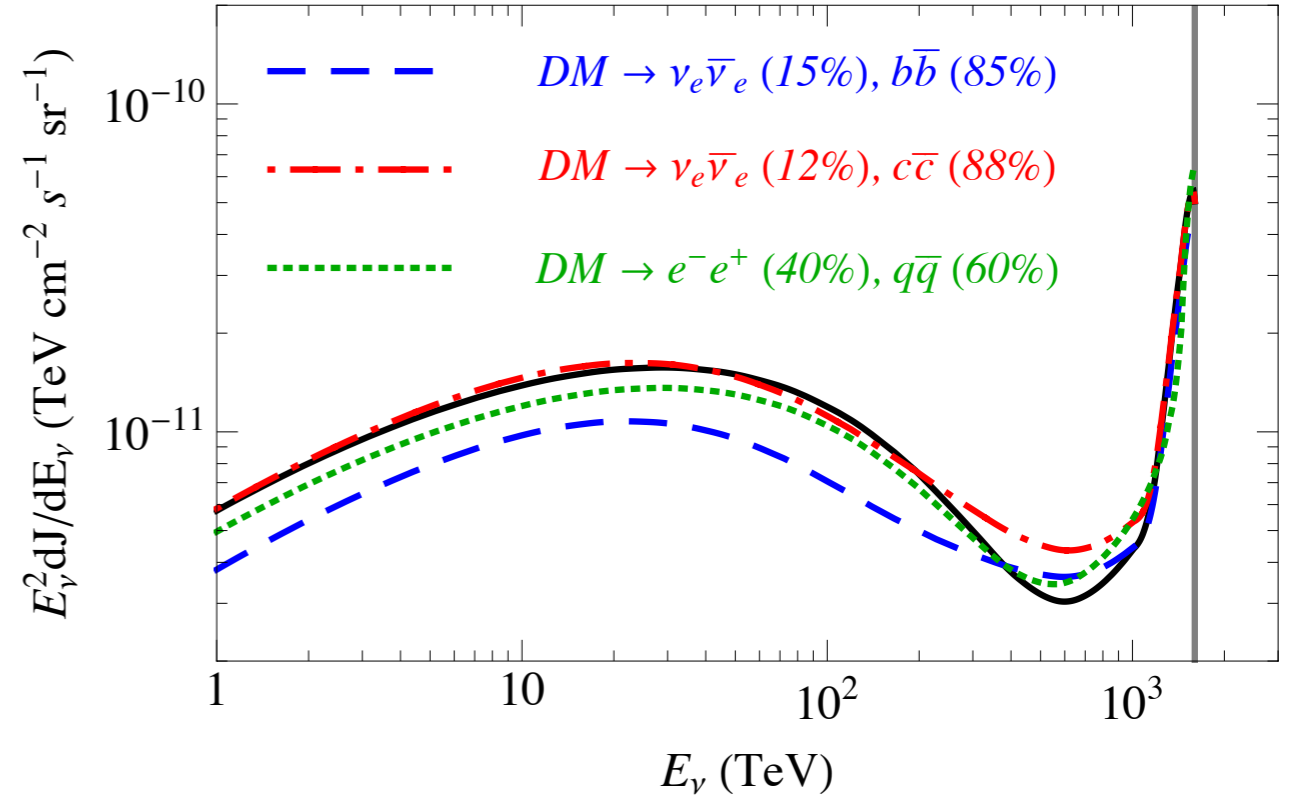
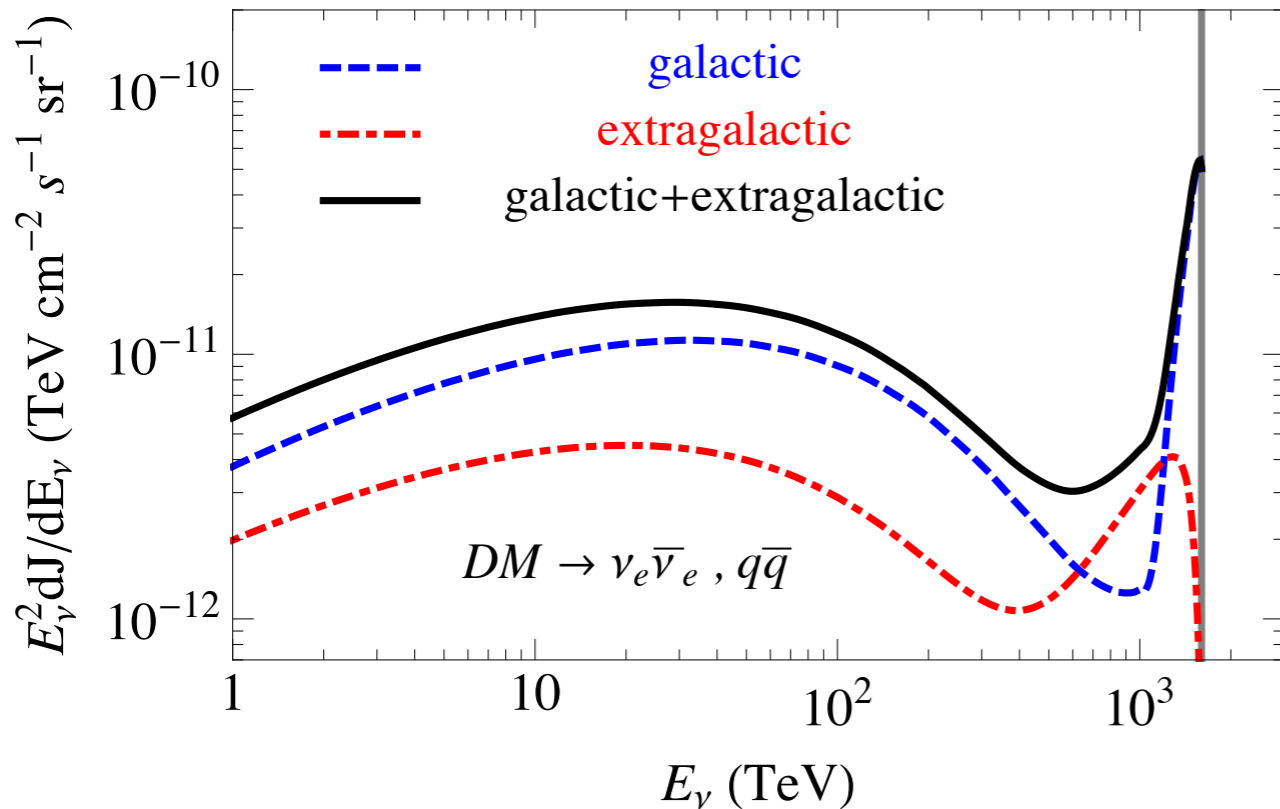
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Flux of neutrinos from decaying DM

✓ fine-tuned decay channels ?

$$\tau_{\text{DM}} = (1-3) \times 10^{27} \text{ s}$$

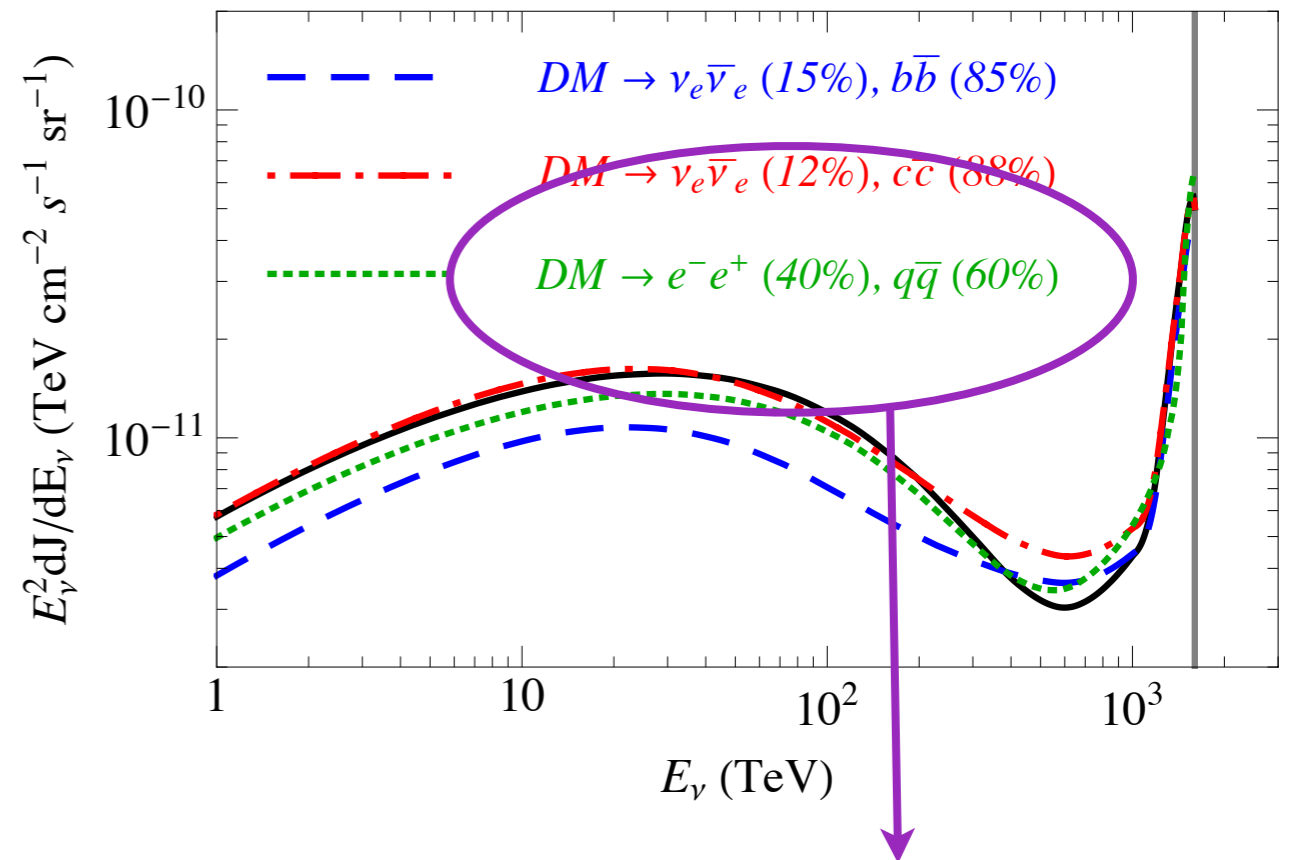
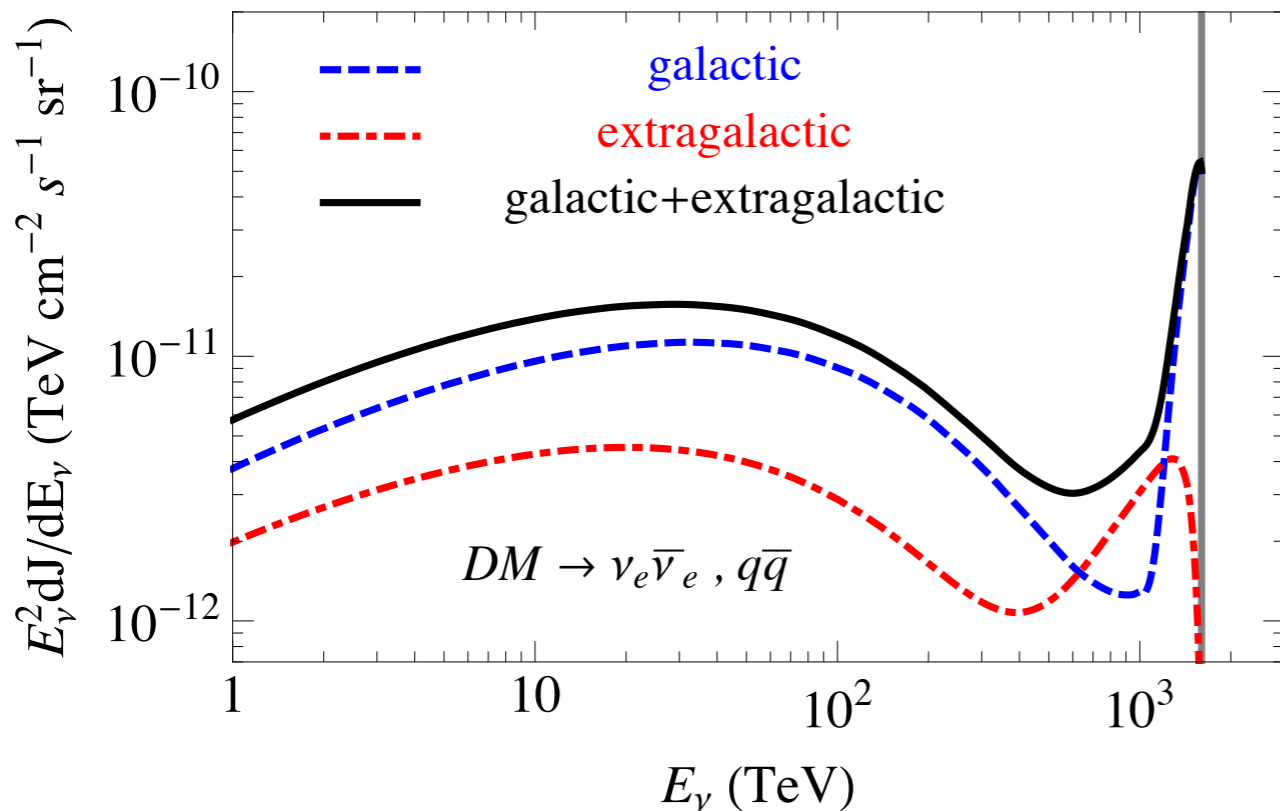


the intriguing features are generic

Flux of neutrinos from decaying DM

✓ fine-tuned decay channels ?

$$\tau_{DM} = (1-3) \times 10^{27} \text{ s}$$



the crucial role of EW cascades

the intriguing features are generic

Confronting with energy distribution of IceCube data

two years data set

✓ branching ratio b_H gives the two PeV events

✓ soft channel and lifetime τ_{DM} gives the upturn in low energy

the value of τ_{DM} is compatible with the bounds derived from neutrinos and gamma rays

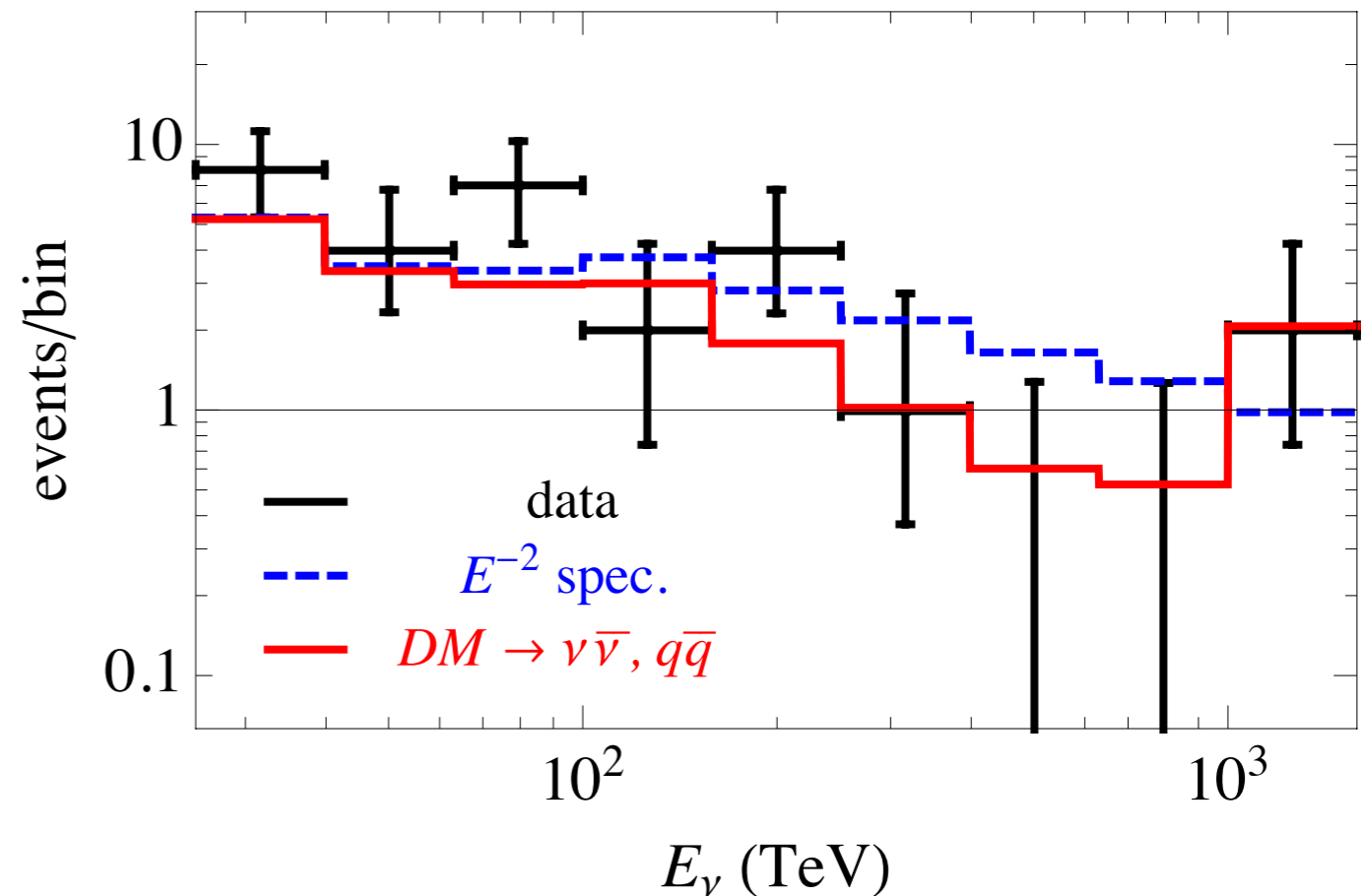
✓ natural explanation for the lack of events $> \text{PeV}$

the value of m_{DM} can be changed within the current uncertainty of the highest energy events

✓ the low energy bins contain large bkg. contribution

the important discriminators of DM vs astrophysical model are high energy bins, where clearly data shows preference to DM model

$$b_H = 0.12 \text{ and } \tau_{DM} = 2 \times 10^{27} \text{ s}$$



✓ different decay channels lead to qualitatively same result

Confronting with energy distribution of IceCube data

three years data set

SM sector  Dark sector

portal type:

$$\mathcal{L}_{\text{portal}} = \frac{\mathcal{O}_{\text{SM}} \mathcal{O}_{\text{DM}}}{\Lambda^{d-4}}$$

Confronting with energy distribution of IceCube data

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"neutrino" portal:

$$\mathcal{O}_{\text{SM}} \rightarrow HL$$

A. Falkowski, J. Juknevich and J. Shelton
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✓ $d = 4 :$ $\mathcal{O}_{\text{DM}} \rightarrow N$

heavy sterile neutrino, DM candidate

T. Higaki, R. Kitano and R. Sato, JHEP (2014)
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[arXiv:1405.0013 [hep-ph]].

UV completion:

$$SU(3)_c \times SU(2)_L \times U(1)_Y \times U(1)_{B-L}$$

$$m_\phi \sim 10^{13} \text{ GeV}$$

"Higgs" field ϕ_{B-L} plays the role of inflaton

$$T_R \sim 10^7 \text{ GeV}$$

Confronting with energy distribution of IceCube data

three years data set

Leptogenesis: $\phi \rightarrow N_2 N_2$ $M_2 \sim 10^{12}$ GeV $\rightarrow \frac{n_B}{s} \sim 10^{-10}$

DM abundance: $\Omega_{N_1} \simeq 0.2 \left(\frac{M_1}{4 \text{ PeV}} \right)^3 \left(\frac{T_R}{3 \times 10^7 \text{ GeV}} \right)^{-1}$

DM lifetime: $\tau_{N_1} \simeq 8 \times 10^{28} \text{ s} \left(\frac{M_1}{1 \text{ PeV}} \right)^{-1} \left(\frac{10^{-29}}{|y_N|^2} \right)$

DM decay channels: $\text{Br}(\ell^\pm W^\mp) = 2\text{Br}(\nu_\ell Z) = 2\text{Br}(\nu_\ell h) = |U_{\ell 1}|^2$ NH

$\text{Br}(\ell^\pm W^\mp) = 2\text{Br}(\nu_\ell Z) = 2\text{Br}(\nu_\ell h) = |U_{\ell 3}|^2$ IH

Confronting with energy distribution of IceCube data

three years data set

IH

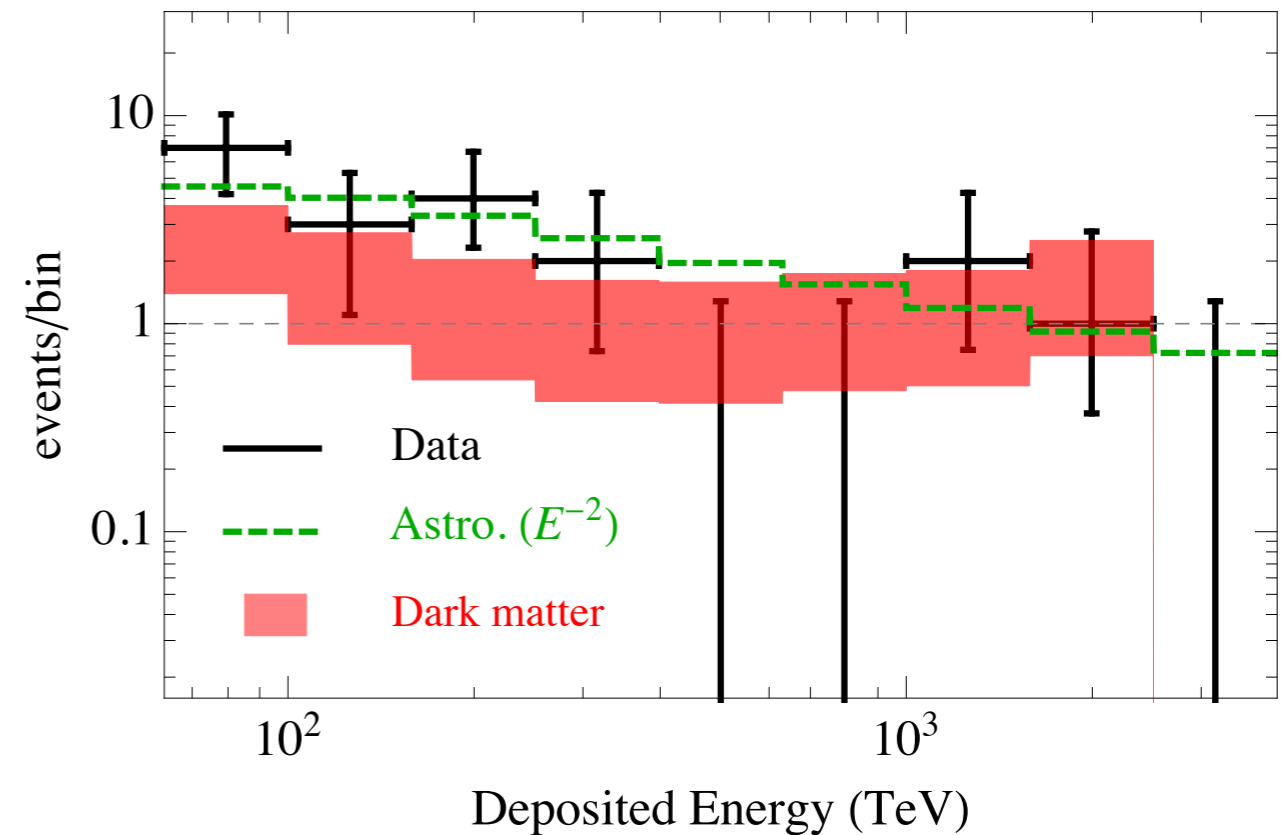
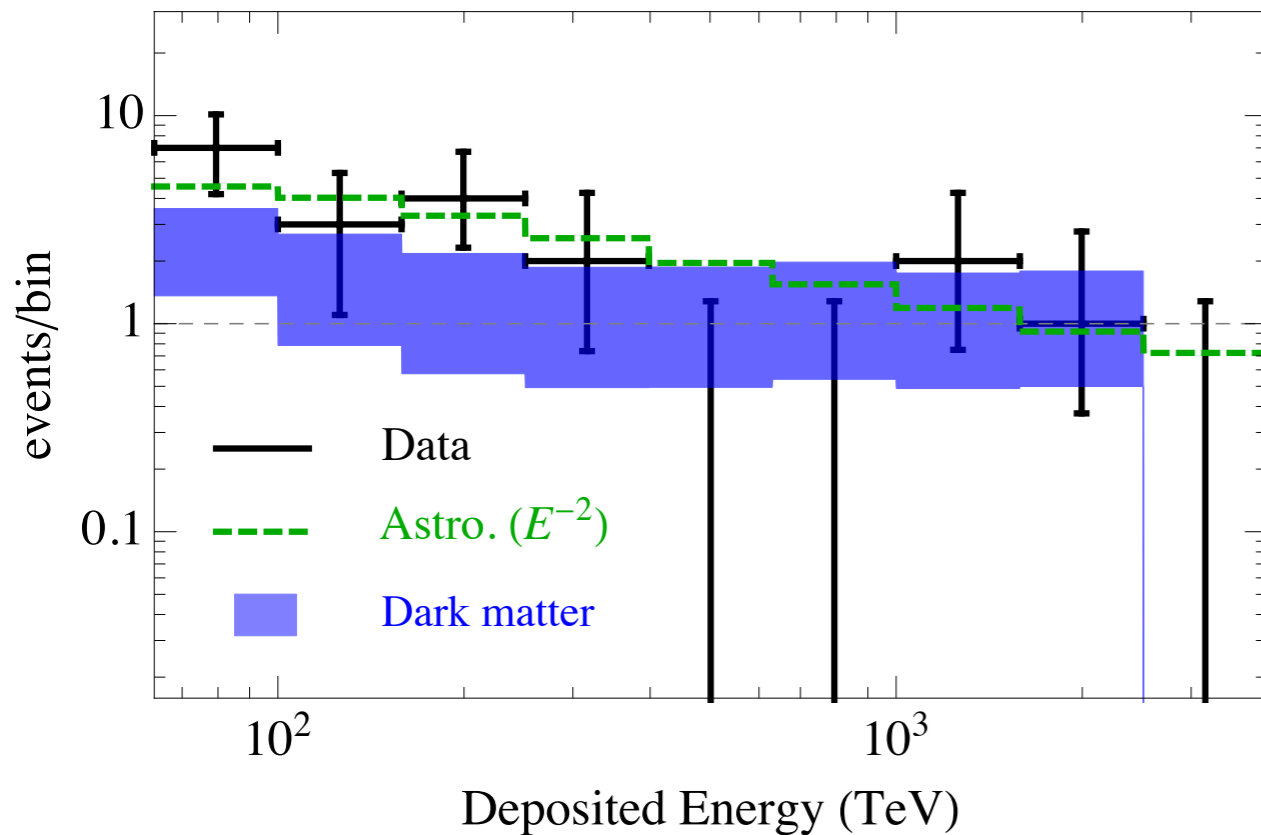
$$\tau_{DM} = 1.1 \times 10^{28} \text{ s}$$

shaded: $\pm 1\sigma$

NH

$$\tau_{DM} = 7.3 \times 10^{27} \text{ s}$$

shaded: $\pm 1\sigma$



$$m_{DM} = 4 \text{ PeV}$$

Confronting with energy distribution of IceCube data

three years data set

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arXiv:0908.1790 [hep-ph].

✓ $d=4:$ $\mathcal{O}_{\text{DM}} \rightarrow N$

production mechanism:

$$m_\phi \gg m_N$$

inflaton decay

$$m_\phi \ll m_N$$

freeze-in

$$g\phi NN, \quad g \simeq 10^{-6}$$

Confronting with energy distribution of IceCube data

three years data set

SM sector \longleftrightarrow Dark sector

portal type:

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A. Falkowski, J. Juknevich and J. Shelton
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✓ $d = 5$: $\mathcal{O}_{\text{DM}} \rightarrow \chi\phi$ singlet fermion and scalar
(Asymmetric DM)

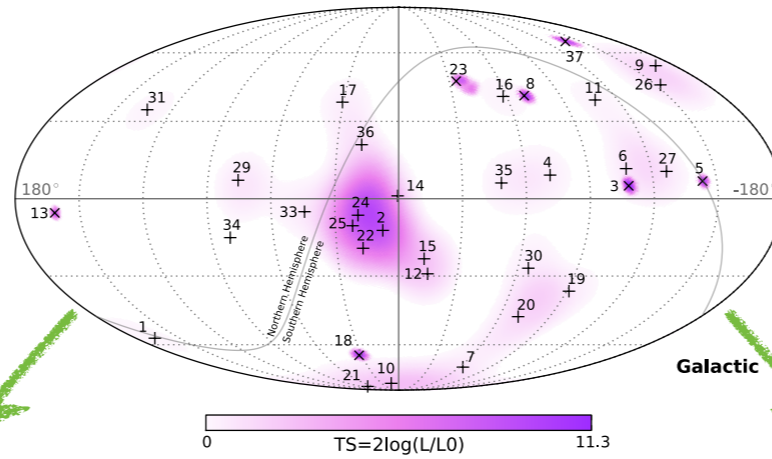
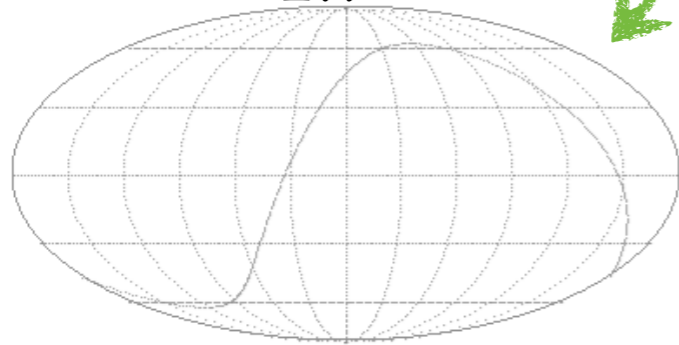
✓ $d = 6$: other portals

✓ For $d > 4$ there are more freedom in branching ratios. We have shown that for the most constrained model ($d=4$) a good fit to the data can be obtained. Obviously better fits can be achieved for $d > 4$.

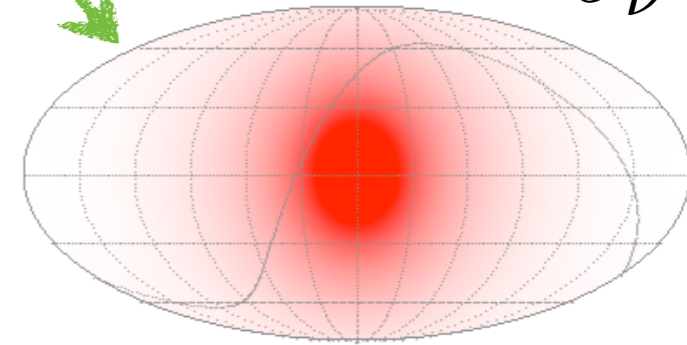
Angular distribution of neutrinos from decaying DM

✓ We would compare

$$p^{\text{iso}} = \frac{1}{4\pi}$$



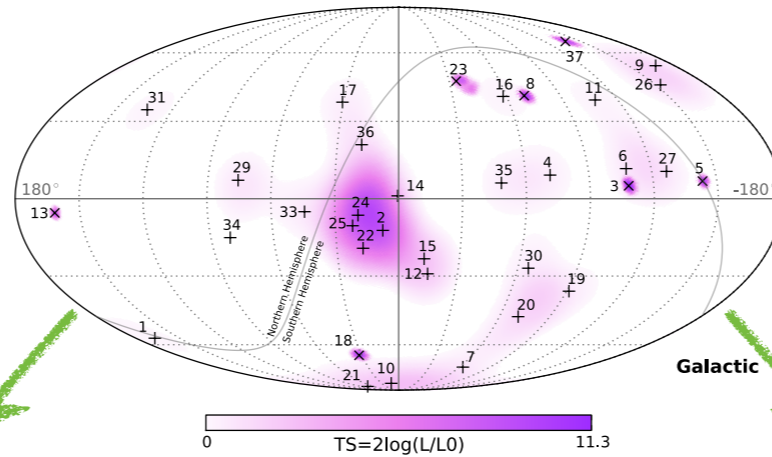
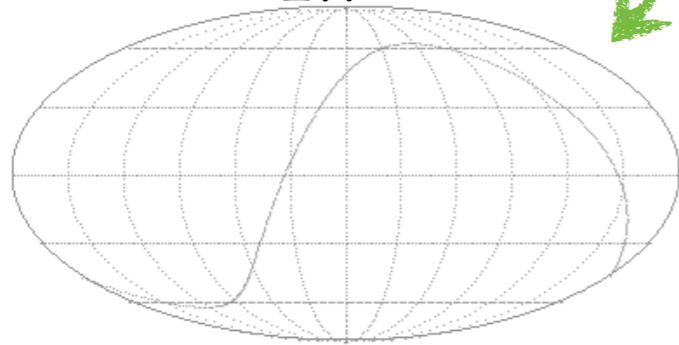
$$p^{\text{DM}} = \frac{1}{J_\nu} \frac{d^2 J_\nu}{db dl}$$



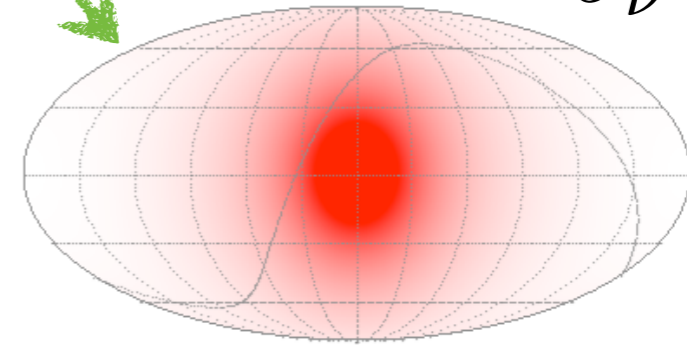
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PDF of data

$$p_i(b, l) = \frac{1}{2\pi\sigma_i^2} \exp\left[-\frac{|\vec{x} - \vec{x}_i|^2}{2\sigma_i^2}\right]$$

"flat sky"
approximation

PDF of
isotropic dis.

$$p^{\text{iso}} = \frac{1}{4\pi}$$

PDF of DM

$$p^{\text{DM}}(b, l) = \frac{1}{J_\nu} \frac{d^2 J_\nu}{db dl} = \frac{\int_0^\infty \rho[r(s, b, l)] ds + \Omega_{\text{DM}} \rho_c \beta}{4\pi(\eta + \Omega_{\text{DM}} \rho_c \beta)}$$

Angular distribution of neutrinos from decaying DM

✓ Likelihood analysis

Number of signal events

Test
Statistics

$$\text{TS}_{\text{like}} = 2 \sum_{i=1}^N (\ln f_i - \ln p_i^{\text{iso}}) = 2 \ln \left(\prod_{i=1}^N f_i \right) - 2N \ln \left(\frac{1}{4\pi} \right)$$

$$f_i = \int p_i(b, l) p^{\text{DM}}(b, l) \cos(b) \, db \, dl = \frac{1}{2\pi\sigma_i^2} \int e^{-\frac{|\vec{x}_i - \vec{x}|^2}{2\sigma_i^2}} p^{\text{DM}}(b, l) \cos(b) \, db \, dl$$

$N = 35$? too optimistic!

Angular distribution of neutrinos from decaying DM

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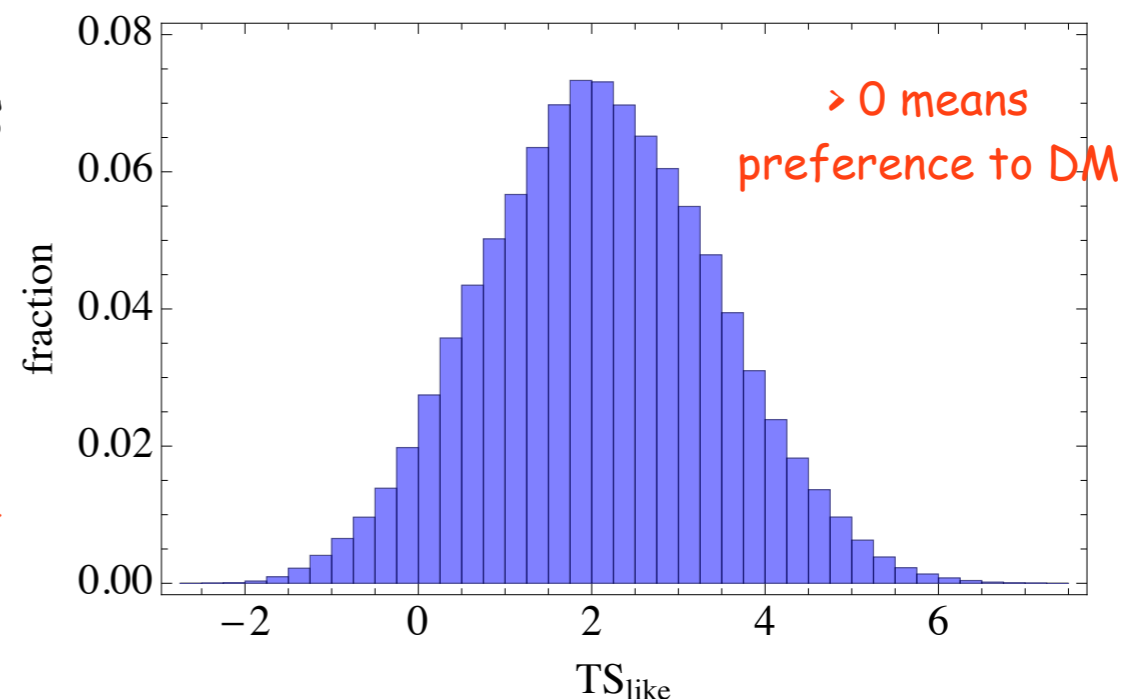
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$N = 35$? too optimistic!

✓ let's assume $N_b = 15$ and all the events with $E > 150$ TeV are signal events

→ $\binom{26}{15}$ ways of selecting the bkg events among the low energy events

Distribution of TS_{like} for all these realizations (mean value = 2.1)



Angular distribution of neutrinos from decaying DM

✓ Likelihood analysis

Number of signal events

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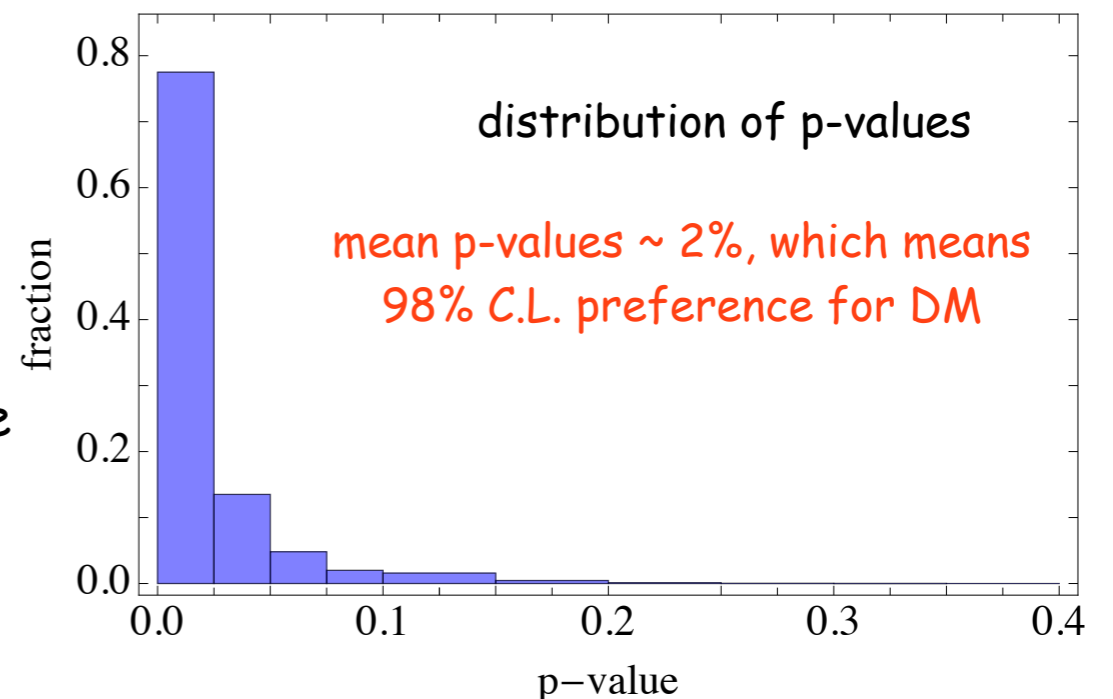
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Quantifying the preference

generating a sample (10^5) of isotropically distributed set of 20 events

for each realization of bkg choosing, p-value is the fraction of generated events which have smaller TS_{like} than the one computed by observed data

p-value



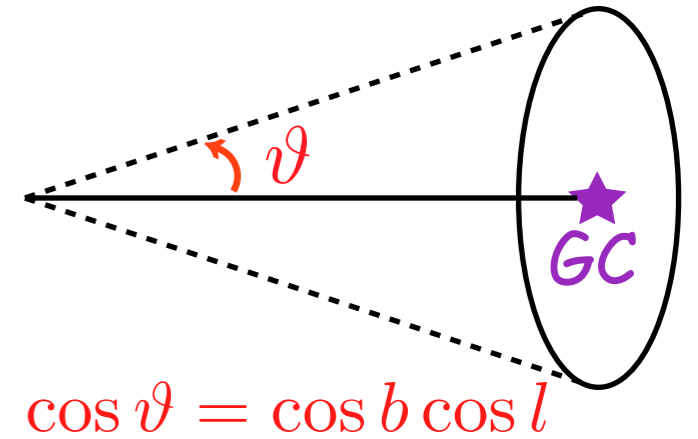
Angular distribution of neutrinos from decaying DM

✓ Kolmogorov-Smirnov test: a powerful non-parametric test

The 2-dim KS test have some ambiguities

$$p^{\text{iso}}(\vartheta) = \int_0^{2\pi} p^{\text{iso}}(\vartheta, \varphi) d\varphi = \int_0^{2\pi} \frac{1}{4\pi} d\varphi = \frac{1}{2}$$

$$p^{\text{DM}}(\vartheta) = \int_0^{2\pi} p^{\text{DM}}(\vartheta, \varphi) d\varphi = \frac{\int_0^\infty \rho[r(s, \vartheta)] ds + \Omega_{\text{DM}} \rho_c \beta}{2(\eta + \Omega_{\text{DM}} \rho_c \beta)}$$



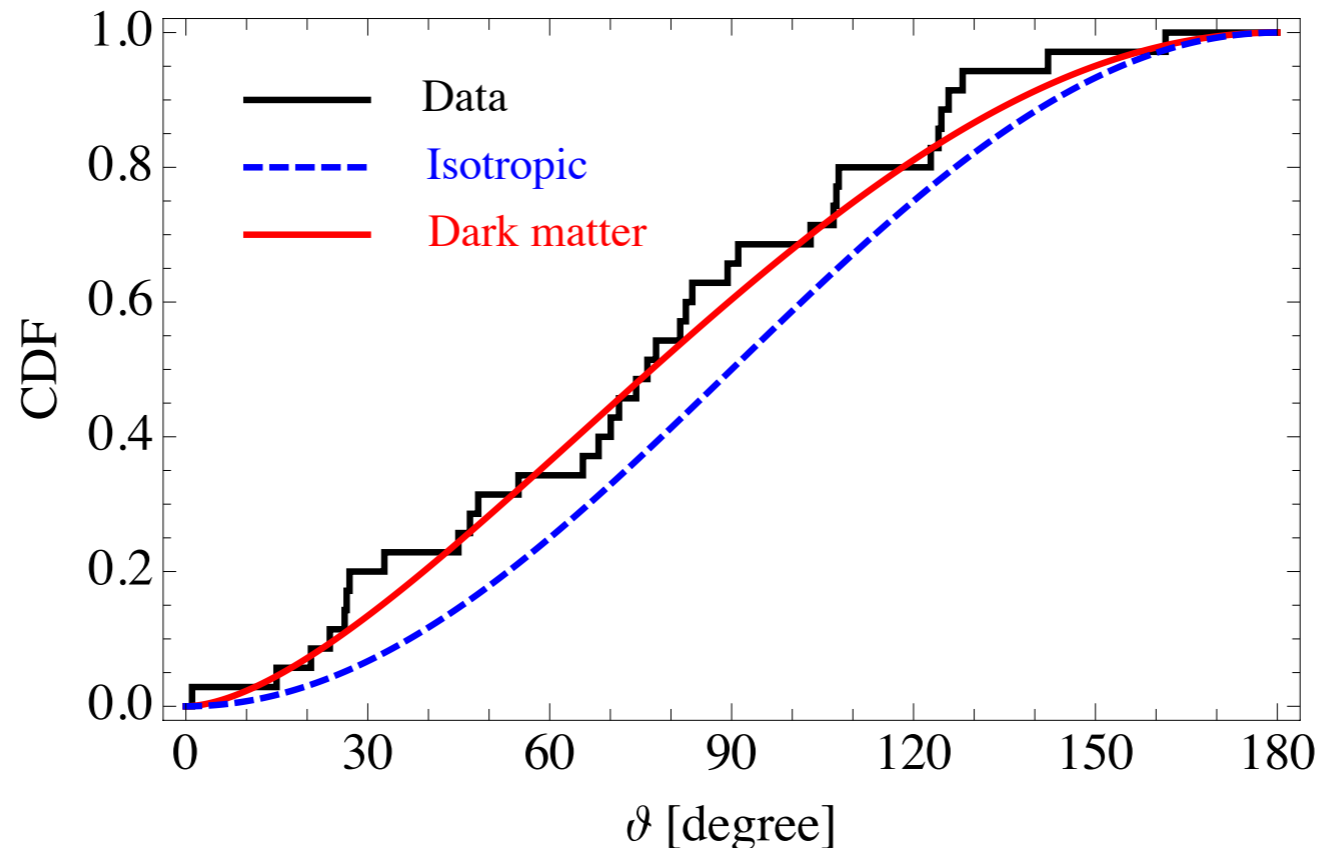
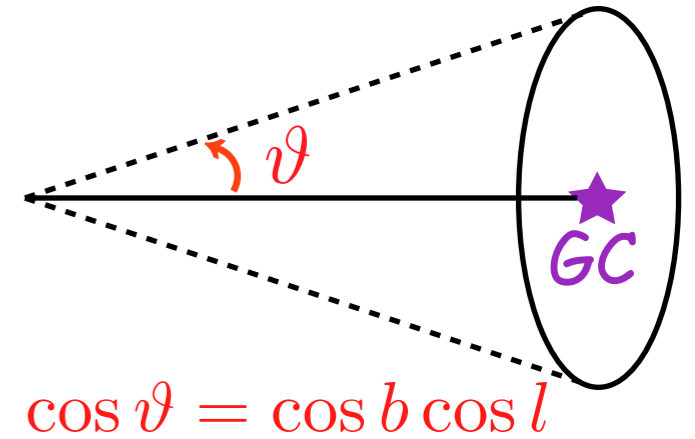
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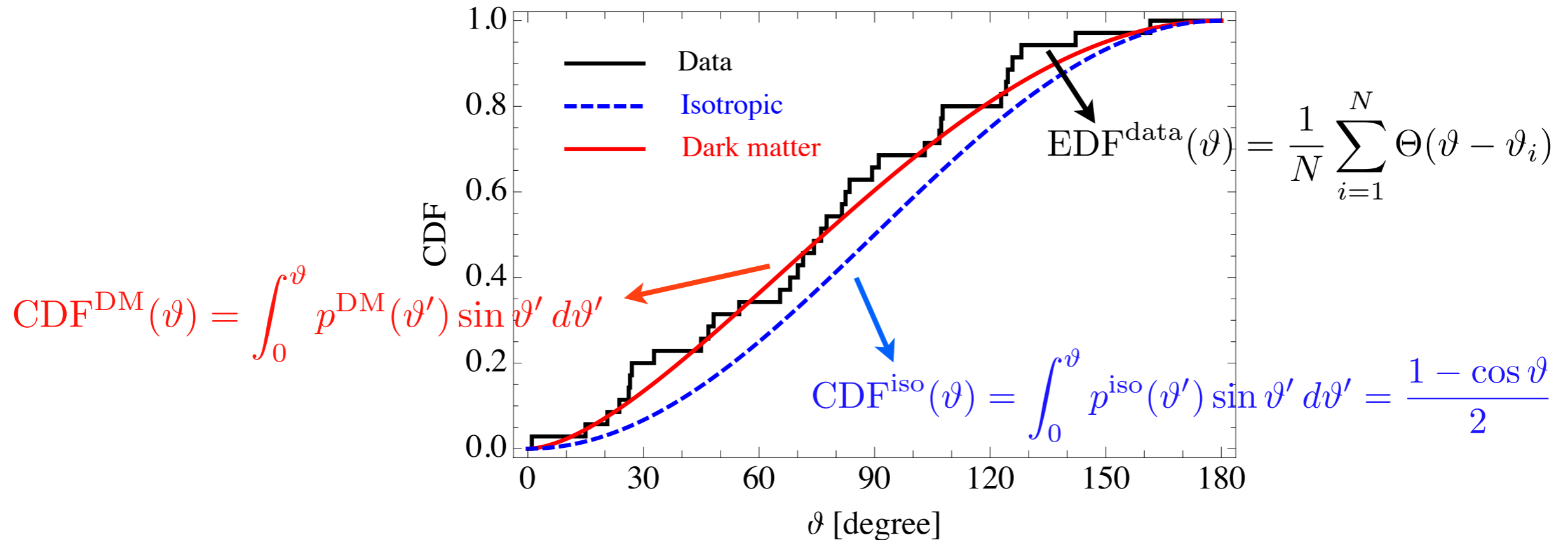
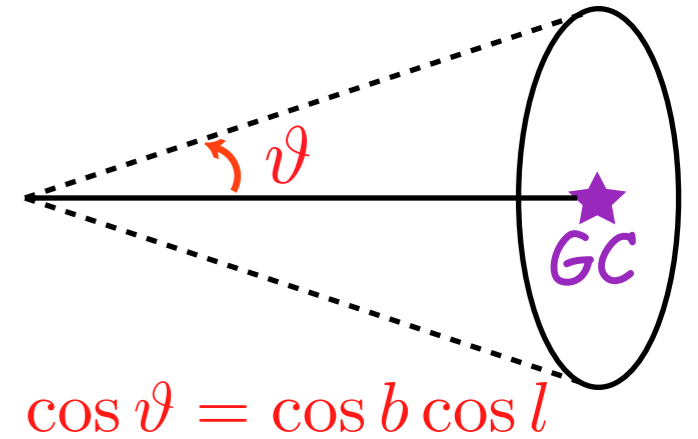
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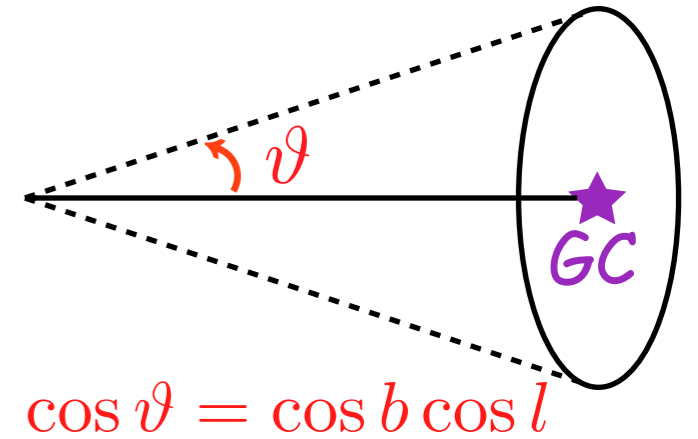
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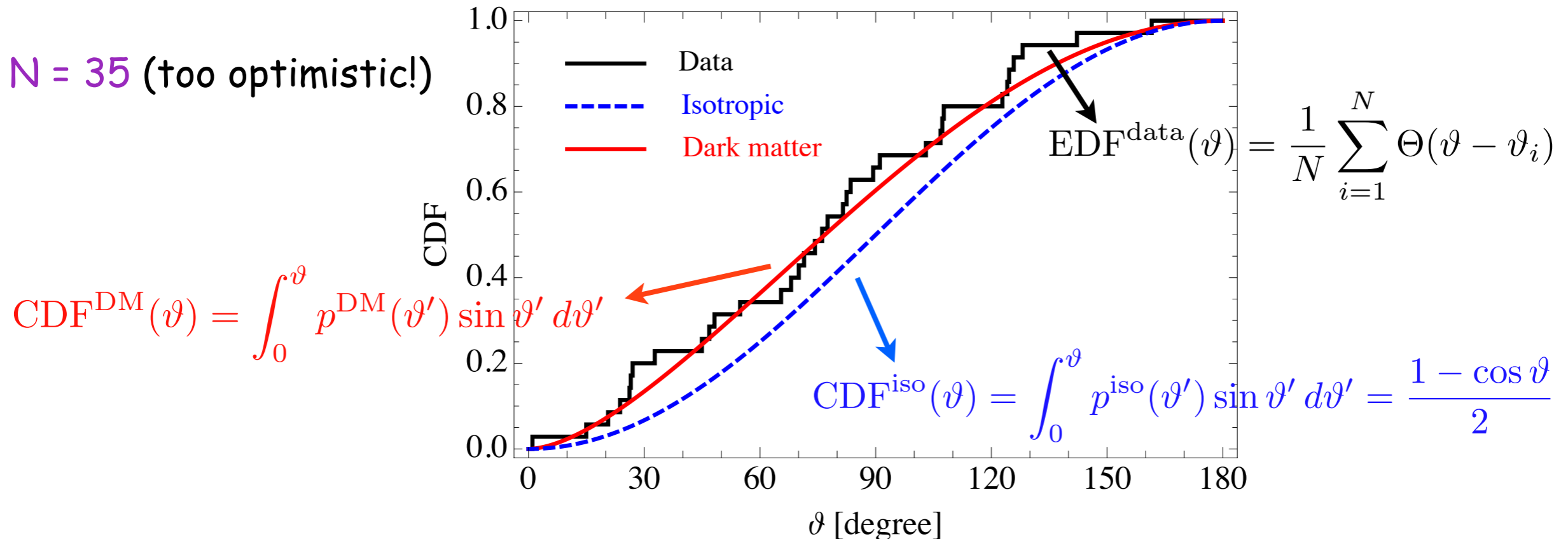
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$N = 35$ (too optimistic!)

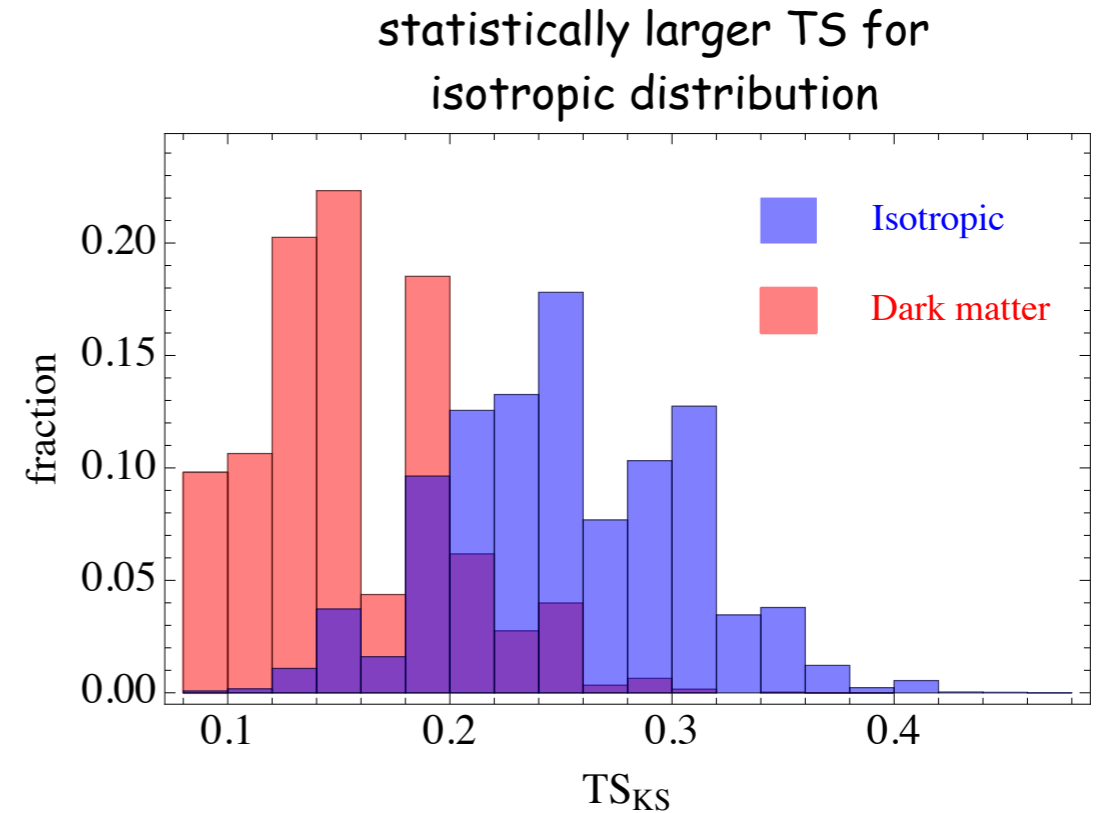


Angular distribution of neutrinos from decaying DM

✓ Kolmogorov-Smirnov test:

Test Statistics

$$TS_{KS} = \max_{1 \leq i \leq N} \left\{ CDF^{DM}(\vartheta_i) - \frac{i-1}{N}, \frac{i}{N} - CDF^{DM}(\vartheta_i) \right\}$$



Angular distribution of neutrinos from decaying DM

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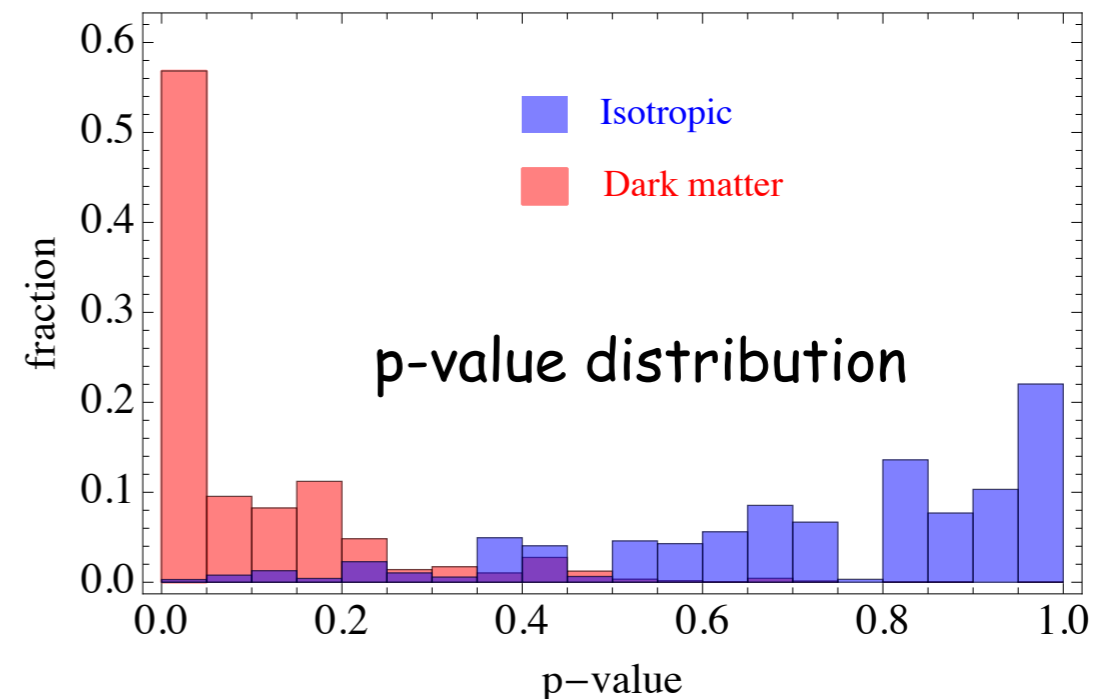
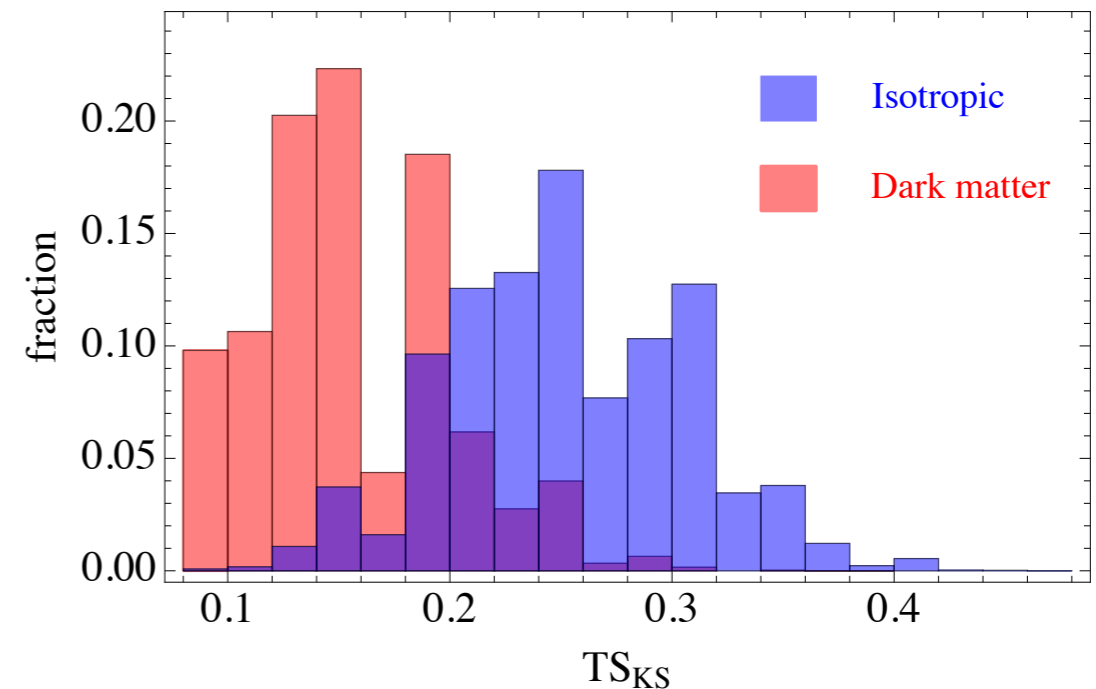
on the average, 10% of generated isotropic sample have smaller TS_{KS} than the values obtained for data vs DM dis.

for data vs isotropic dis. it is 73%



less than 2 σ preference for DM dis.

statistically larger TS for isotropic distribution



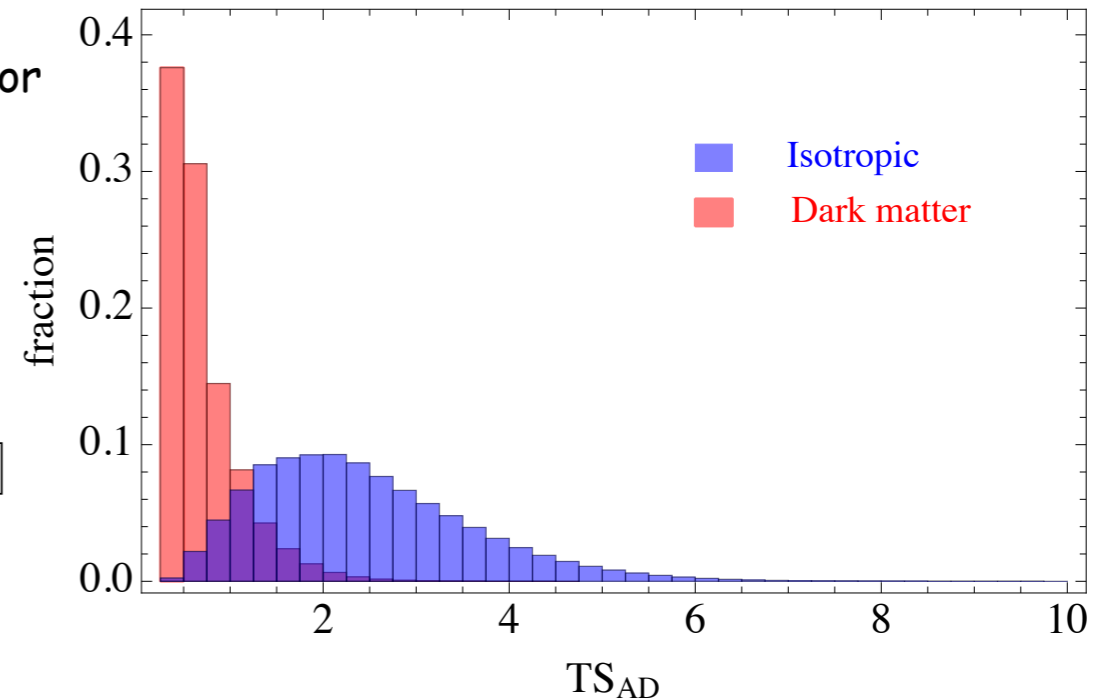
Angular distribution of neutrinos from decaying DM

- ✓ Anderson-Darling test: a powerful non-parametric test, especially sensitive to the end points

Test Statistics

$$TS_{AD} = -N - \frac{1}{N} \sum_{i=1}^N (2i - 1) [\ln(\text{CDF}^{\text{DM}}(\vartheta_i)) + \ln(1 - \text{CDF}^{\text{DM}}(\vartheta_{N+1-i}))]$$

statistically larger TS for
isotropic distribution



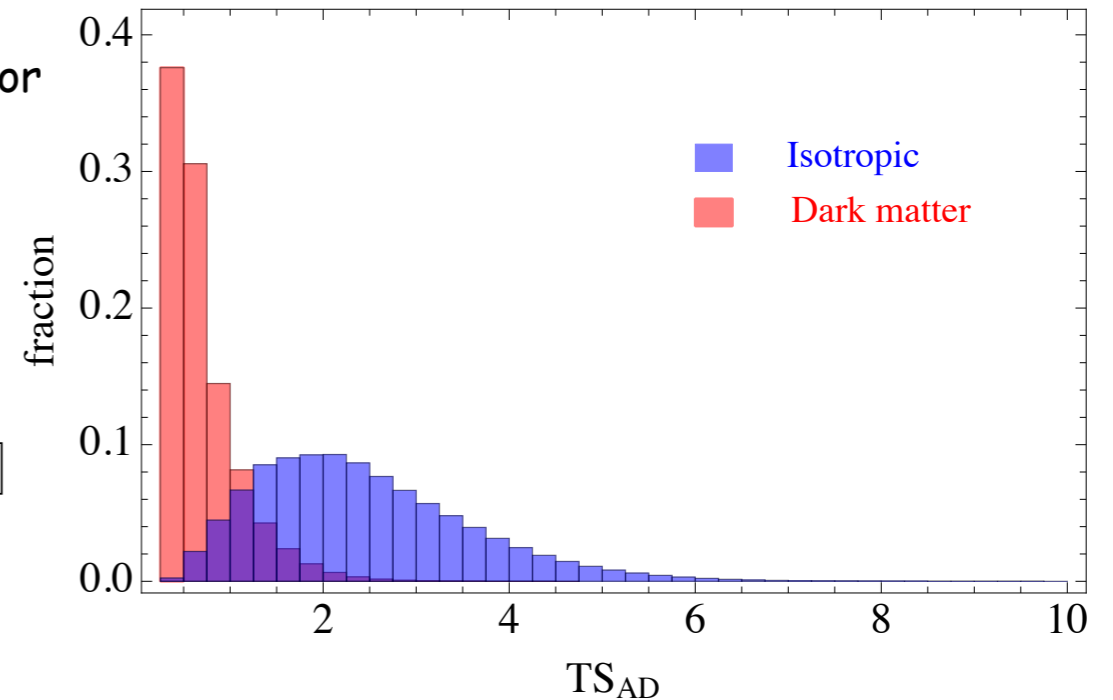
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statistically larger TS for isotropic distribution



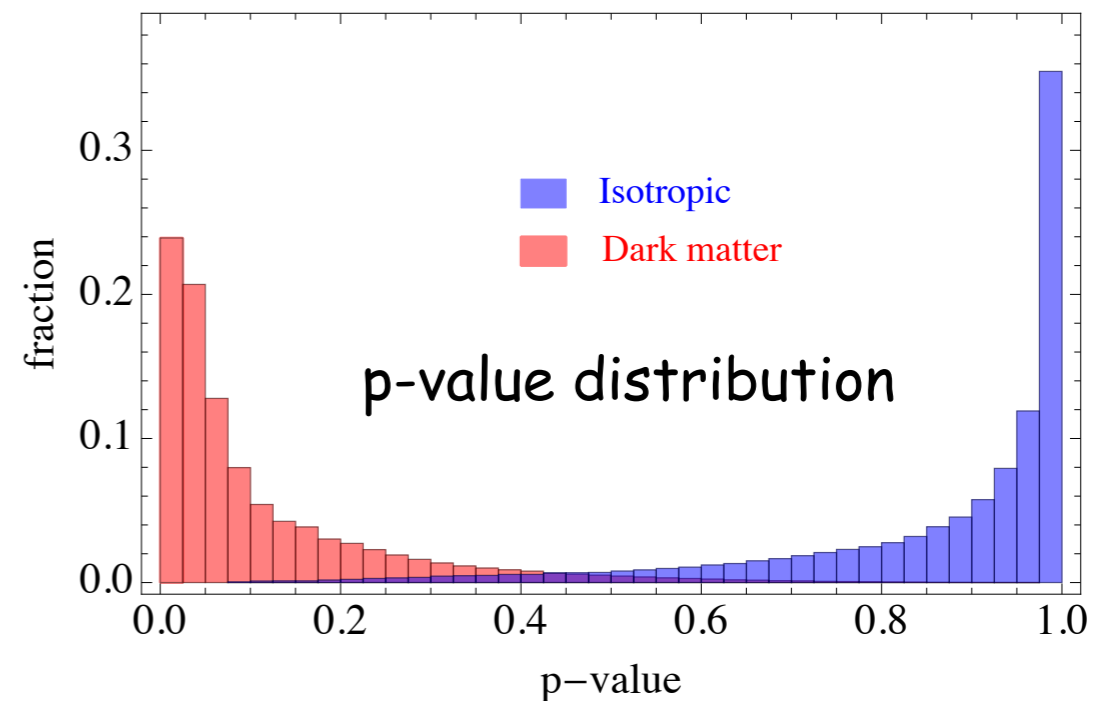
again, generating a sample (10^5) of isotropically distributed set of 20 events



on the average, 11% of generated isotropic sample have smaller TS_{KS} than the values obtained for data vs DM dis.
for data vs isotropic dis. it is 86%



less than 2σ preference for DM dis.



Angular distribution of neutrinos from decaying DM

✓ Some issues:

Angular resolution in KS
and AD tests?



even after shifting all the events
to higher ϑ values still the
preference for DM persist

Angular distribution of neutrinos from decaying DM

✓ Some issues:

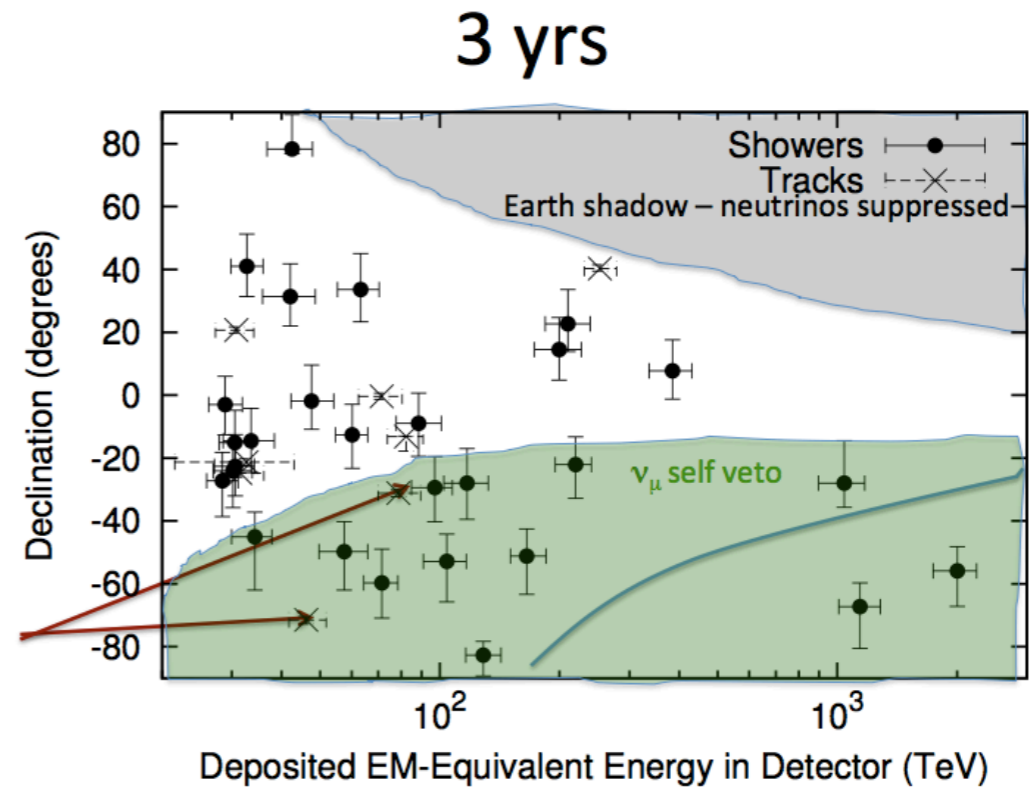
Angular resolution in KS
and AD tests?



even after shifting all the events
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Background rejection?

Figure from T. Gaisser



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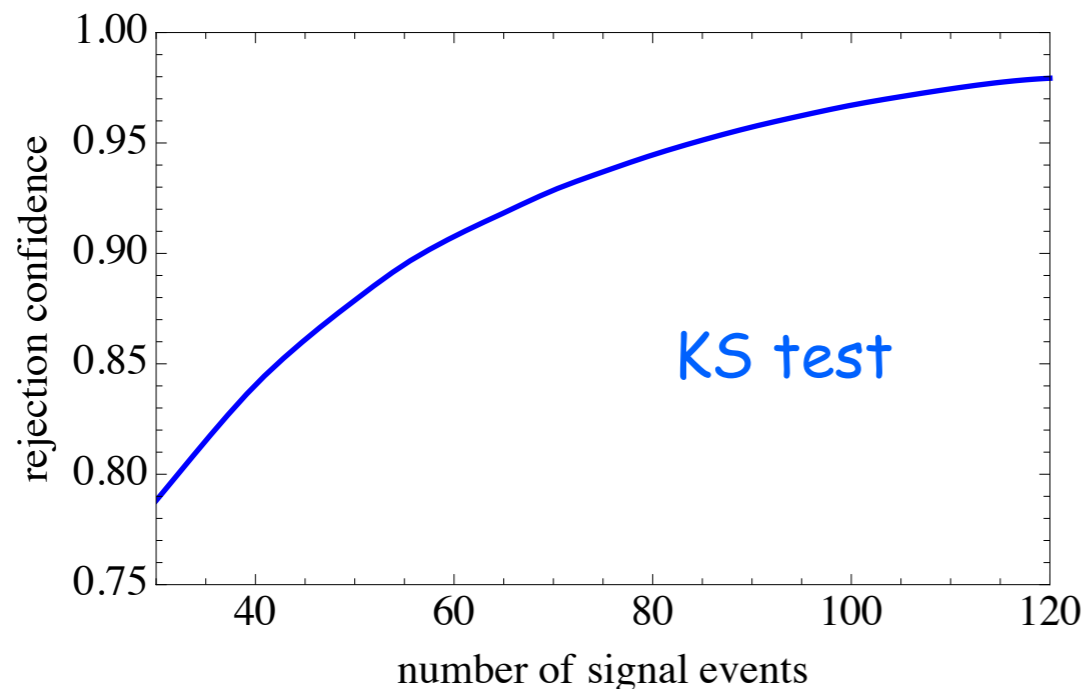
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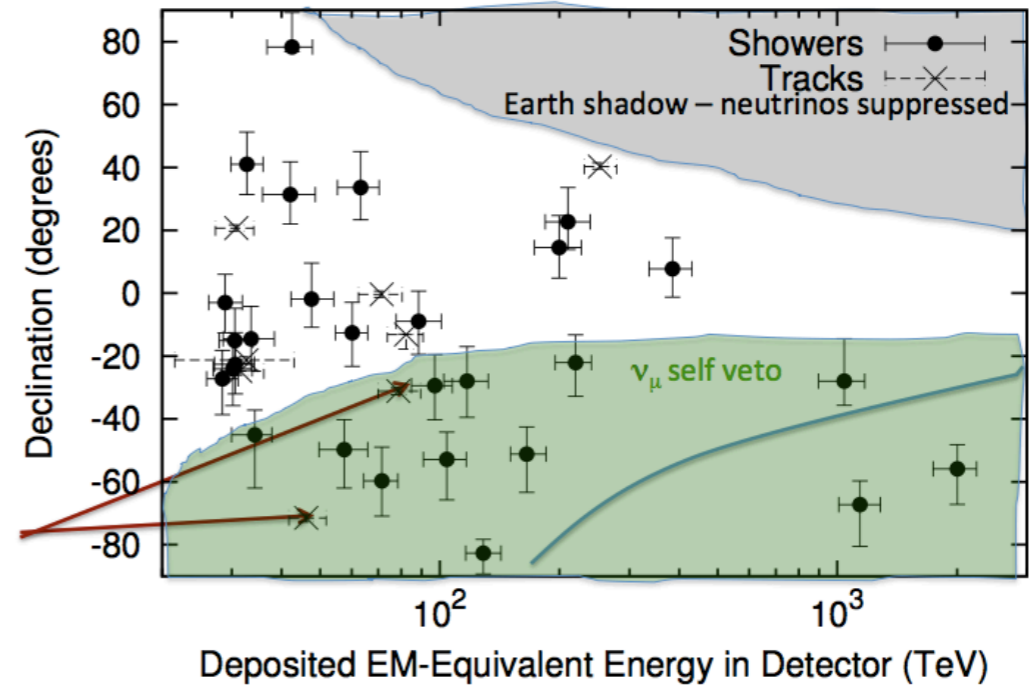
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Figure from T. Gaisser



3 yrs



How many events are
needed for a 3σ
discrimination?

Gamma ray bounds

Universe is opaque for
gamma-rays with $E > 1 \text{ TeV}$



cascades develop: gamma-ray
interaction with interstellar
radiation field and CMB



gamma-rays populate at
lower energies, 1 - 100 GeV

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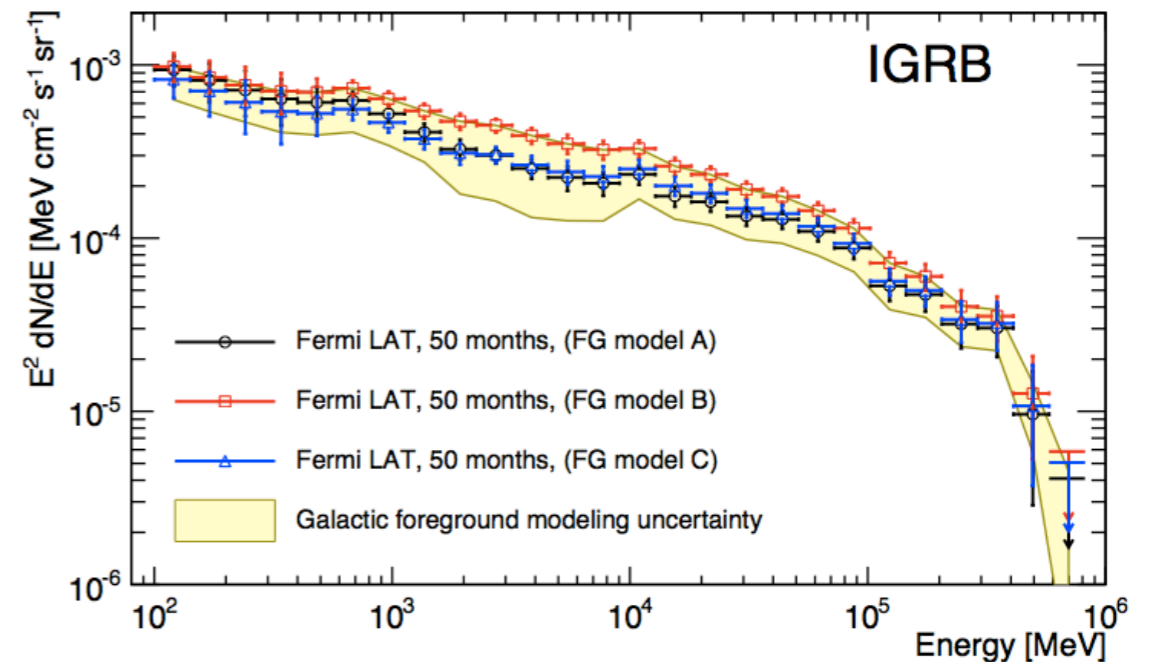
integrated energy density

$$\omega_\gamma = \frac{4\pi}{c} \int_{E_1}^{E_2} E_\gamma \frac{d\varphi_\gamma}{dE_\gamma} dE_\gamma \lesssim 4.4 \times 10^{-7} \text{ eV/cm}^3$$

$$E_1 \sim \mathcal{O}(1) \text{ GeV}$$

$$E_2 \sim \mathcal{O}(100) \text{ GeV}$$

M. Ackermann et al. [The Fermi LAT Collaboration], arXiv:1410.3696 [astro-ph.HE].



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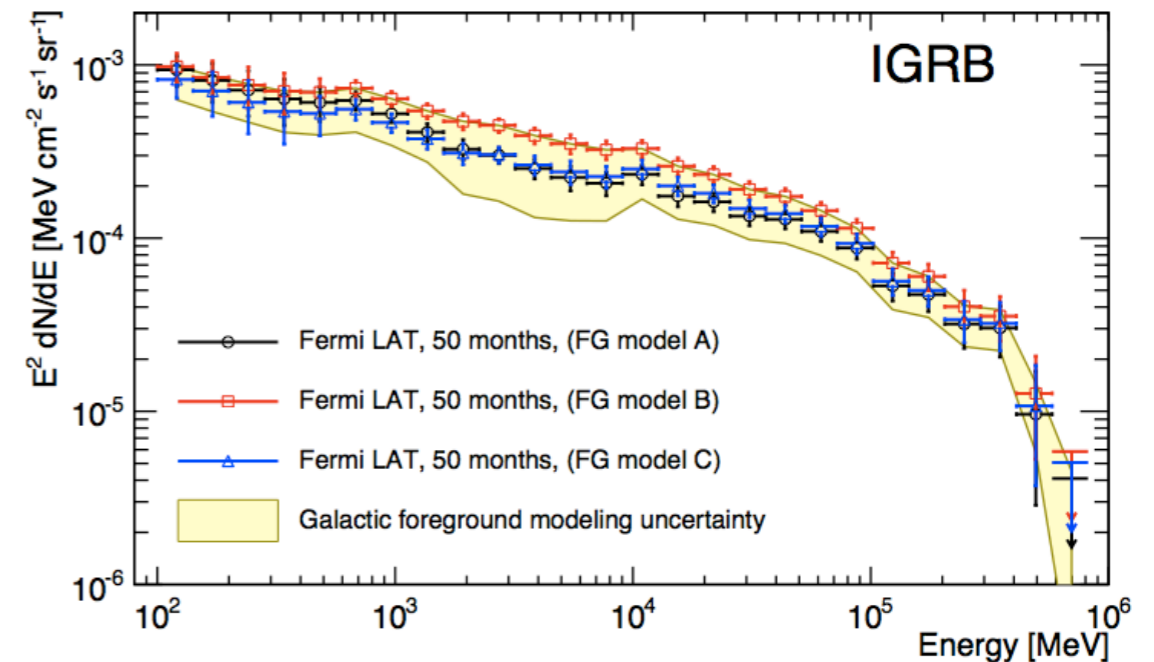
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total electromagnetic energy budget
(NH case)

$$\frac{4\pi}{c} \int \sum_{i=\text{gal,extragal}} \left[E_\gamma \frac{d\varphi_\gamma}{dE_\gamma} + E_e \frac{d\varphi_{e^\pm}}{dE_e} \right] dE \simeq 5.2 \times 10^{-8} \text{ eV/cm}^3$$



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Gamma ray bounds

✓ Galactic component

at \sim PeV, the absorption length of gamma-rays
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energy budget consideration :

the integrated flux of γ -ray above 330 TeV must be below $1.0 \times 10^{-13} \text{ cm}^{-2} \text{ s}^{-1} \text{ sr}^{-1}$

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For DM decay
(NH case)
unattenuated flux



the integrated flux of γ -ray above 330 TeV is $1.2 \times 10^{-14} \text{ cm}^{-2} \text{ s}^{-1} \text{ sr}^{-1}$

the integrated flux of γ -ray above 775 TeV is $8.8 \times 10^{-15} \text{ cm}^{-2} \text{ s}^{-1} \text{ sr}^{-1}$



What can we learn about DM if the IceCube events originate from conventional astrophysical flux?

constraining:

- ✓ DM lifetime
- ✓ annihilation cross section

Constraining DM properties

✓ DM lifetime

contribution of DM to the events in each bin should be smaller than N_{limit}

bin #	$\log_{10}(E_\nu/\text{TeV})$	$N_{\text{astro}}(E_\nu^{-2} \div E_\nu^{-2.3})$	N_{data}	$N_{\text{limit}}(E_\nu^{-2} \div E_\nu^{-2.3})$	N_{limit}
#1	1.4 – 1.6	9.46 ÷ 10	11	7.8 ÷ 7.46	16.6
#2	1.6 – 1.8	4.31 ÷ 5.3	6	6.53 ÷ 5.87	10.5
#3	1.8 – 2.0	4.55 ÷ 5.68	7	7.41 ÷ 6.58	11.8
#4	2.0 – 2.2	3.97 ÷ 4.82	3	3.98 ÷ 3.73	6.68
#5	2.2 – 2.4	3.32 ÷ 3.56	4	5.15 ÷ 5.01	8.00
#6	2.4 – 2.6	2.59 ÷ 2.42	2	3.65 ÷ 3.71	5.32
#7	2.6 – 2.8	1.96 ÷ 1.62	0	2.3 ÷ 2.3	2.3
#8	2.8 – 3.0	1.55 ÷ 1.1	0	2.3 ÷ 2.3	2.3
#9	3.0 – 3.2	1.2 ÷ 0.74	2	4.31 ÷ 4.64	5.32
#10	3.2 – 3.4	0.92 ÷ 0.5	1	3.3 ÷ 3.51	3.89
#11	3.4 – 3.6	0.73 ÷ 0.35	0	2.3 ÷ 2.3	2.3
#12	3.6 – 3.8	1.72 ÷ 0.76	0	2.3 ÷ 2.3	2.3

Poisson statistics:

at $q\%$ C.L.

$$\frac{q}{100} = \frac{\int_0^{N_{\text{limit}}^i} L(N_{\text{data}}^i, N) dN}{\int_0^\infty L(N_{\text{data}}^i, N) dN}$$

$$L(N_{\text{data}}^i, N) = \frac{(N + N_{\text{astro}}^i)^{N_{\text{data}}^i}}{N_{\text{data}}^i!} e^{-(N + N_{\text{astro}}^i)}$$

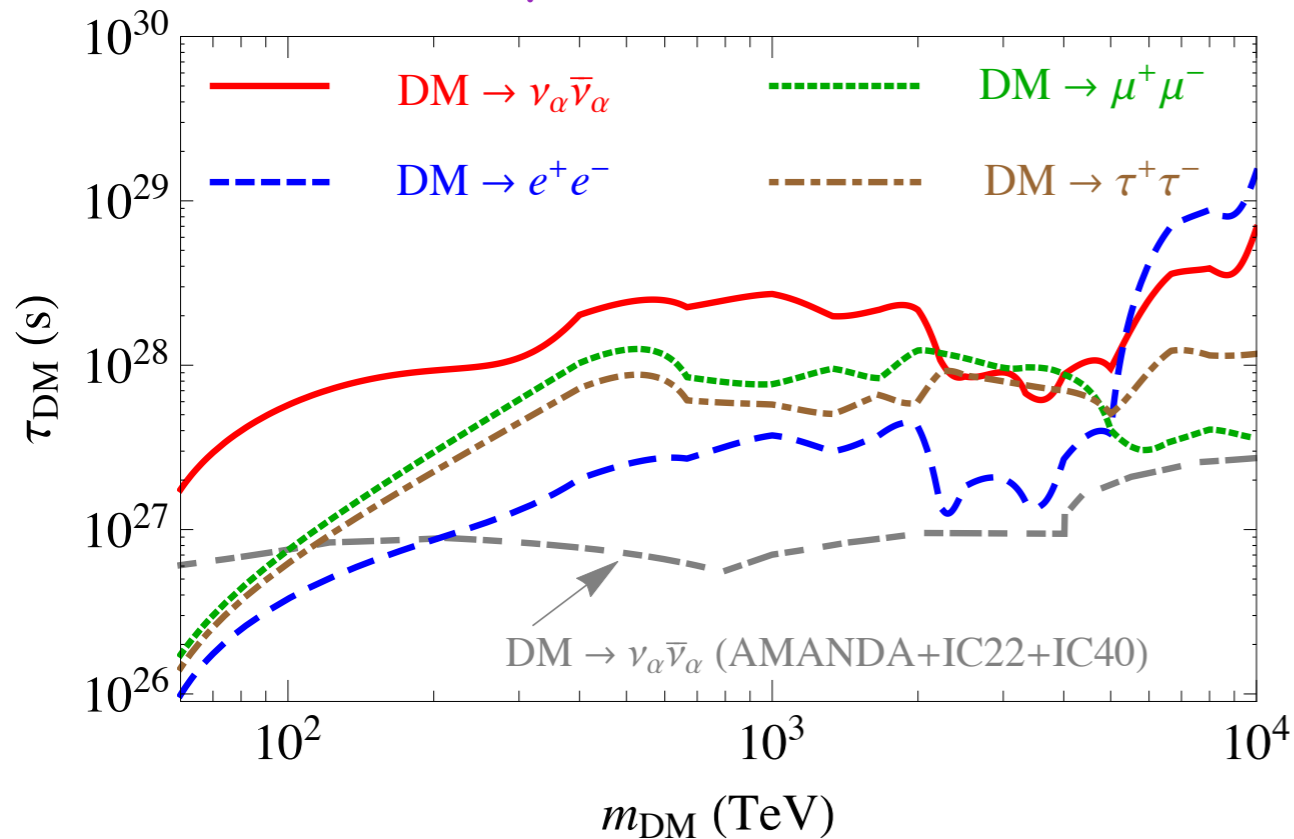
or

$$L(N_{\text{data}}^i, N) = \frac{(N)^{N_{\text{data}}^i}}{N_{\text{data}}^i!} e^{-N}$$

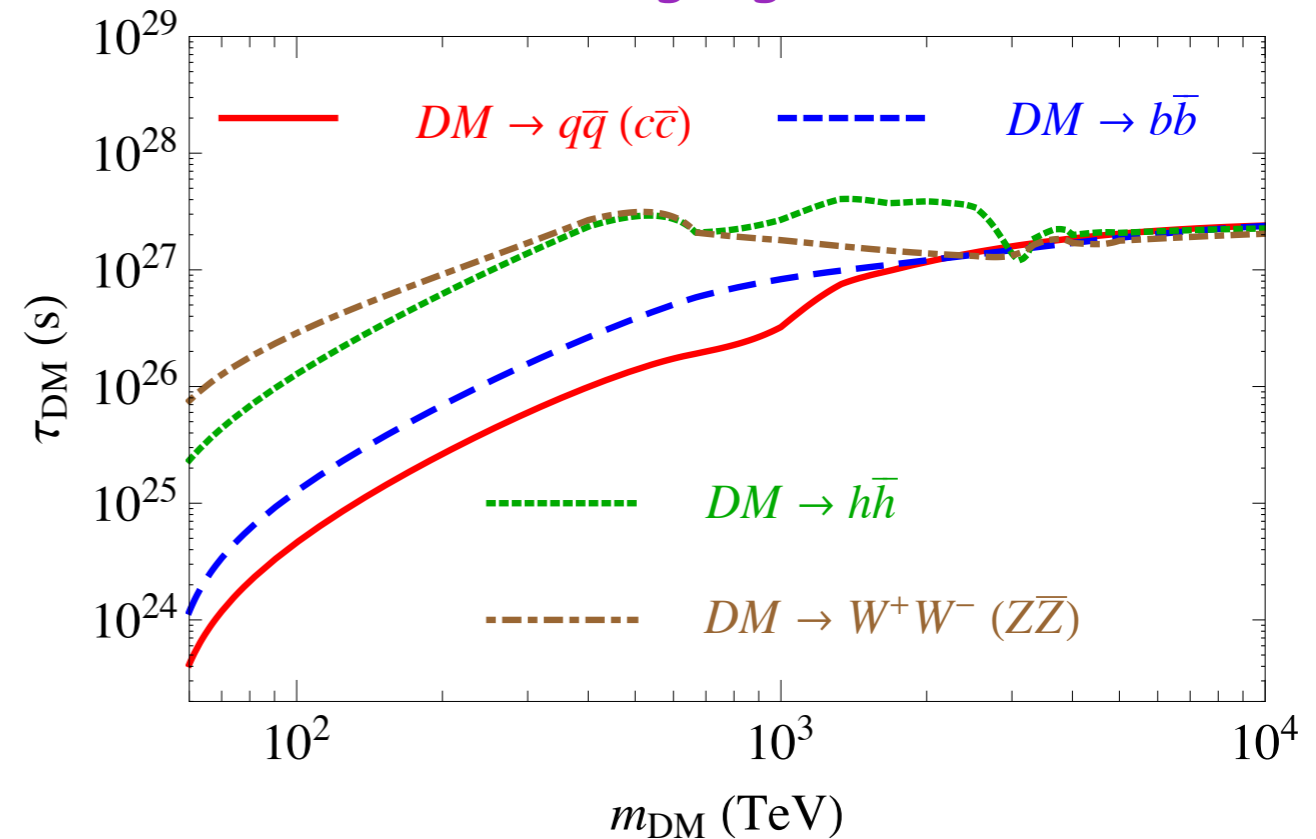
Constraining DM properties

✓ limits on DM lifetime (90% C.L.)

leptonic channels



hadronic/gauge channels

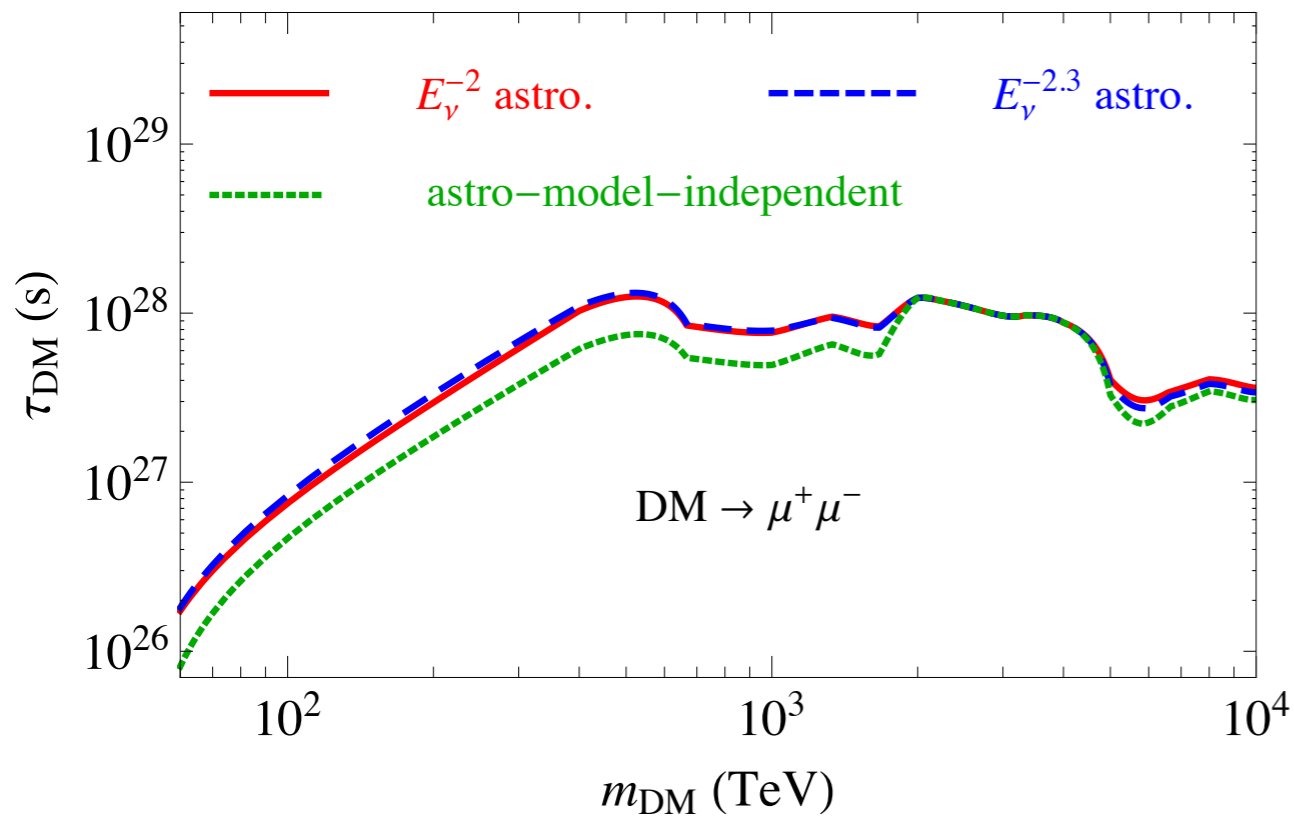
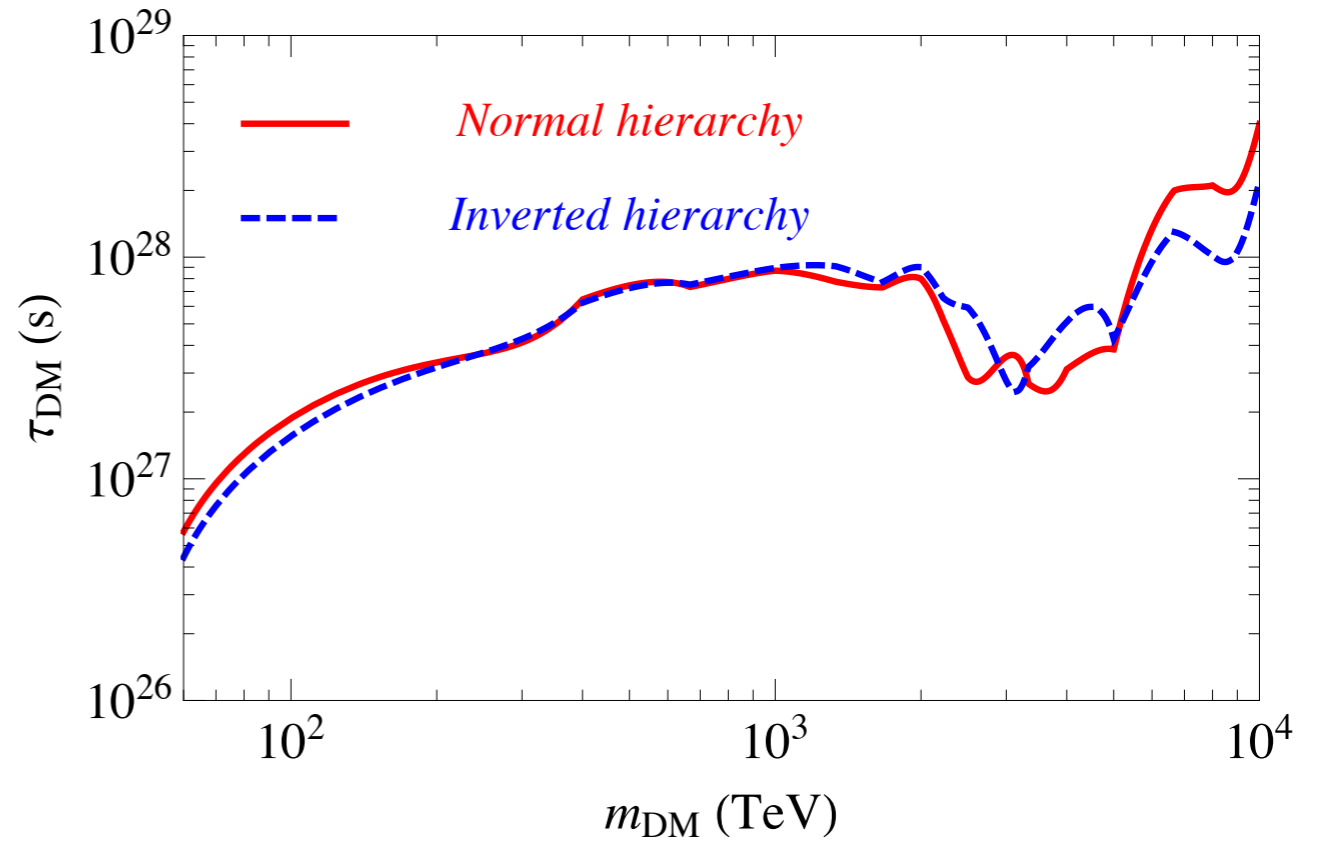


- ✓ at least one order of magnitude stronger lower limit on the DM lifetime, in the relevant DM mass range
- ✓ for a specific model, different channels should be scaled according to the corresponding branching ratios

Constraining DM properties

✓ limits on DM lifetime (90% C.L.)

NH and IH cases →



← dependence on the astro. model?

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The lower part (< 100 TeV) of the observed spectrum can be used to probe $\langle\sigma v\rangle$

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The isotropic components of neutrino flux from DM annihilation:

The residual isotropic flux from the Galactic halo (anti-GC direction)

$$\frac{dJ_{\text{iso}}^{\text{ann}}}{dE_\nu} = \frac{\langle\sigma v\rangle}{2} \frac{1}{4\pi m_{\text{DM}}^2} \frac{dN}{dE_\nu} (\text{l.o.s.})_{\text{anti-GC}} \quad \text{where} \quad (\text{l.o.s.})_{\text{anti-GC}} = \int_0^\infty \rho^2[r(s, b=0, l=\pi)] ds$$

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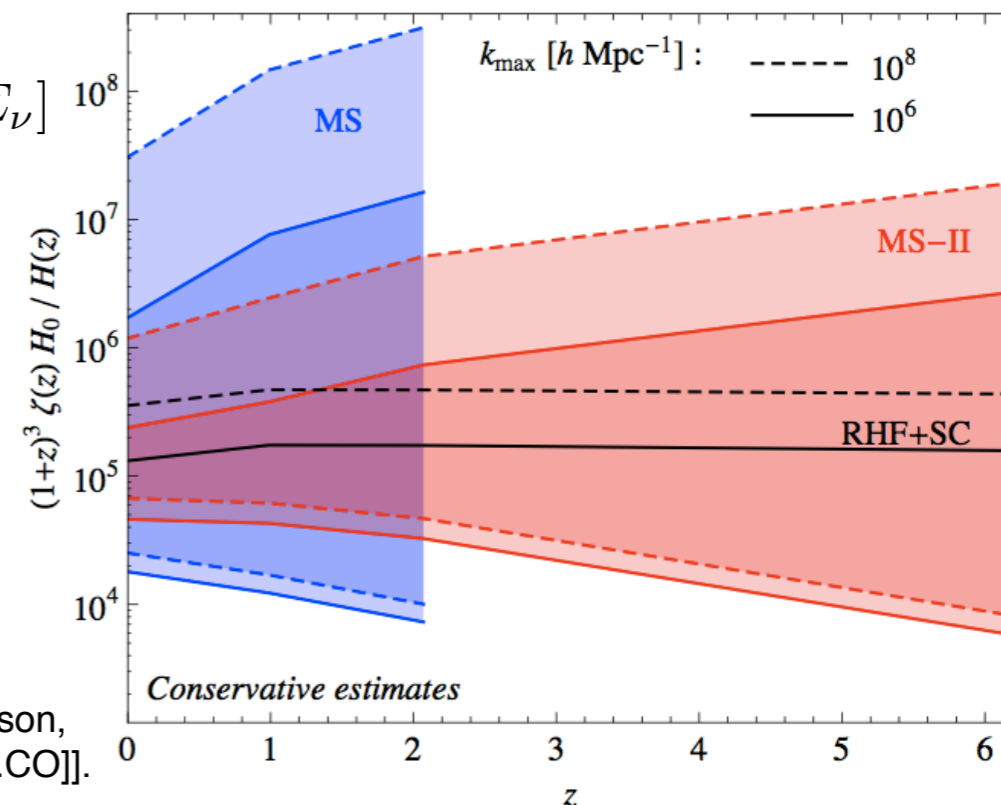
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The cosmic flux from all redshift

$$\frac{dJ_{\text{cos}}^{\text{ann}}}{dE_\nu} = \frac{\langle\sigma v\rangle}{2} \frac{\Omega_{\text{DM}}^2 \rho_c^2}{4\pi m_{\text{DM}}^2} \frac{c}{H_0} \int_0^\infty \frac{(1+z)^3 \zeta(z) dz}{\sqrt{\Omega_m (1+z)^3 + \Omega_\Lambda}} \frac{dN}{dE_\nu} [(1+z) E_\nu]$$

$\zeta(z)$ flux multiplier (DM clustering)



E. Sefusatti, G. Zaharijas, P. D. Serpico, D. Theurel and M. Gustafsson,
Mon. Not. Roy. Astron. Soc. 441, (2014) [arXiv:1401.2117 [astro-ph.CO]].

Constraining DM properties

✓ upper limits on annihilation cross section $\langle\sigma v\rangle$ (90% C.L.)

minimum \div maximum value used for $\zeta(z)$ unit of $\langle\sigma v\rangle$ is $10^{-22} \text{ cm}^3\text{s}^{-1}$

DM + DM \rightarrow m_{DM}	100 TeV	50 TeV	30 TeV
$\nu_\alpha \bar{\nu}_\alpha$	1.39 \div 0.22	1.21 \div 0.36	2.44 \div 0.88
$q\bar{q}$	489 \div 84.5	1427 \div 299	9934 \div 4603
$b\bar{b}$	185 \div 30.4	517 \div 106	3514 \div 1621
$c\bar{c}$	592 \div 100	1708 \div 348	11218 \div 5215
e^+e^-	14.7 \div 2.38	17.8 \div 5.06	41.3 \div 14.2
$\mu^+\mu^-$	4.47 \div 0.65	9.06 \div 1.6	23.7 \div 9.23
$\tau^+\tau^-$	5.84 \div 0.93	10.9 \div 2.3	28.5 \div 10.8
$h\bar{h}$	21.2 \div 3.36	53.4 \div 9.49	177 \div 76.5
$Z\bar{Z}$	11.9 \div 2.05	18.1 \div 4.09	40.7 \div 16.3
W^+W^-	14.4 \div 2.4	23.7 \div 4.96	54.5 \div 22.3

✓ for some final states (neutrinos, charged leptons) the limit is a bit stronger than the unitary bound

conclusions

✓ The excess of events observed by IceCube in the energy range ~ 30 TeV - 2 PeV is an evidence for astrophysical flux or other "New Physics" induced fluxes

✓ Several features of the observed events motivate us for a DM interpretation: cut-off at ~ 2 PeV, dip in (400 - 1000) TeV and anisotropy.

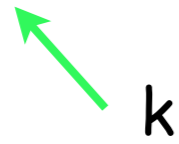
✓ We argued that a PeV-scale decaying DM, with generic decay channels, can naturally explain these features. The required lifetime is allowed by the current limits. Both the energy and angular distributions mildly prefer DM interpretation.

✓ With more statistics in the next few years, the DM interpretation of IceCube events can be tested: especially the angular distribution can resolve the issue. The persistence of the dip feature in $\sim (400 - 1000)$ TeV also can be supportive.

Thank you !

Decaying Dark Matter

$$\Phi = k E_\nu^{-2}$$



	k	$E^{\min}-E^{\max}$ (TeV)	N_{bg}	N_{sig}	N_{limit}
AMANDA	7.4×10^{-8}	16 - 2.5×10^3	6	7	5.4
IceCube-22	1.6×10^{-7}	340 - 2×10^5	0.6	3	6.1
IceCube-40	3.6×10^{-8}	2×10^3 - 6.3×10^6	0.1	0	2.3
Auger	1.7×10^{-7}	10^5 - 10^8	0	0	2.3
ANITA	1.3×10^{-7}	10^6 - 3.2×10^{11}	0.97	1	3.3

$$N_{\text{exp}} = T \Delta \Omega \sum_{\alpha} \left[\int_{E_{\nu}^{\min}}^{E_{\nu}^{\max}} \Phi_{\nu_{\alpha} + \bar{\nu}_{\alpha}} A_{\text{eff}}^{\alpha}(E_{\nu}) dE_{\nu} \right]$$

$$q/100 = \frac{\int_0^{N_{\text{limit}}} L(N_{\text{sig}}|N) dN}{\int_0^{\infty} L(N_{\text{sig}}|N) dN} \quad \text{where} \quad L(N_{\text{sig}}|N) = \frac{(N + N_{\text{bg}})^{N_{\text{sig}}}}{N_{\text{sig}}!} e^{-(N + N_{\text{bg}})}$$

Decaying Dark Matter

✓ extragalactic contribution:

$$H(z) = H_0 \sqrt{\Omega_\Lambda + \Omega_m (1+z)^3}$$

$$\frac{dJ_{\text{eg}}}{dE_\nu} = \frac{\Omega_{\text{DM}} \rho_c}{4\pi m_{\text{DM}} \tau_{\text{DM}}} \int_0^\infty dz \frac{1}{H(z)} \frac{dN_\nu}{dE_\nu} [(1+z)E_\nu] e^{-s_\nu(E_\nu, z)}$$

Opacity

$$s_\nu(E_\nu, z) = \begin{cases} 7.4 \times 10^{-17} (1+z)^{7/2} (E_\nu / \text{TeV}), & \text{for } 1 \ll z < z_{\text{eq}} \\ 1.7 \times 10^{-14} (1+z)^3 (E_\nu / \text{TeV}), & \text{for } z \gg z_{\text{eq}} \end{cases}$$

at the Earth

$$\begin{pmatrix} J_e \\ J_\mu \\ J_\tau \end{pmatrix} = \begin{pmatrix} P_{ee} & P_{e\mu} & P_{e\tau} \\ & P_{\mu\mu} & P_{\mu\tau} \\ & & P_{\tau\tau} \end{pmatrix} \begin{pmatrix} I_e \\ I_\mu \\ I_\tau \end{pmatrix} \quad \text{production point}$$

decoherent
oscillation

Event 8 (Track)

