Precision top quark physics at a future linear e⁺e⁻ collider; mass & top EW couplings

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With special thanks to:

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Top is turning 20 this year Of course, parents love all their kids equally Top's maturing, Higgs is turning 3 this year





She's so cute; it's only natural that the new-born gets a bit more attention



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Top at lepton colliders

Only quark that escaped direct scrutiny at LEP/SLC

Top quark mass is an important parameter

- EW fit, Higgs loops, fate of the universe
- First chance to study top couplings to

neutral EW bosons

- − $q\bar{q} \rightarrow \gamma/Z \rightarrow t\bar{t}$ produced at Tevatron and LHC, but swamped by QCD production
- associated production $t\bar{t}\gamma$ and $t\bar{t}Z$ are (slowly) building up strength at the LHC e^+





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 Z/γ

Top quark pairs



See: Garcia, Perello, Ros, Vos, Study of single top production at high energy electron-positron colliders, arXiv:1411.2355

Must measure rate and properties of WbWb production. For a precise comparison of data and prediction more theory work is needed!



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top quark production at lepton colliders

For precision there is nothing like e⁺e⁻

QCD corrections calculated to N^2LO Scale variations at N^3LO estimated at ~ 0.3%. Electroweak corrections are sizable, though. Calibrate center-of-mass-energy to 1 in 10^4 and luminosity to 0.1%



Top at threshold

At threshold we have to include QCD bound-state corrections in calculations

Match threshold & continuum calculations and supply them in a generator (WHIZARD) F. Bach (DESY), A. Hoang (Vienna), M. Stahlhofen (DESY)

Parametric uncertainty due to top quark mass AND width are important at threshold (Martinez & Miquel extracted 4 parameters from fit)





F. Bach, preliminary

LL resummation (orange) and NLL resummation (blue) for FB asymmetry versus center-of-mass energy, m(1S) = 172 GeV, WHIZARD 2.2.3_beta_2



Influence of the top quark mass on x-sec and $A_{_{FR}}$

- very pronounced below $\sqrt{s} = 360 \text{ GeV}$
- 2.9%/GeV at \sqrt{s} = 380 GeV
- 1.3%/GeV at \sqrt{s} = 420 GeV
- 0.6%/GeV at $\sqrt{s} = 500 \text{ GeV}$

With the assumption of a 100 MeV pole mass measurement at threshold, the remaining uncertainty is one per mil or less above 420 GeV



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LC top physics – canonical programme

350 GeV:

Threshold: top quark mass to < 100 MeV (+width & Yukawa)

Kuhn, Acta Phys.Polon. B12 (1981) 347 Martinez, Miquel, EPJ C27, 49 (2003) Seidl, Simon, Tesar, Poss, EPJC73 (2013) 2530 A. Juste et al. ArXiv:1310.0799

500 GeV:

New physics: precise characterization of $t\bar{t}Z$ and $t\bar{t}\gamma$ vertices

M.S. Amjad et al., arXiv:1307.8102 F. Richard, arXiv:1403.2893

500-1500 GeV:

ttH direct access to top Yukawa coupling





Studies at relevant thresholds (tt, ttH) and at 500 GeV What's the potential at other sqrt(s)? What else can we do?





q 0.8

cross-section [9.0

0.2

0



Top quark mass



Top quark masss

Top quark mass measurements



Top quark mass at threshold



J. H. Kuhn, Acta Phys. Polon. B12 (1981) 347.



Top quark mass at an LC



Uncertainties on extracted top quark mass

Stat. error ~16/34 MeV (with/without polarization) No dependence on location of scan energy Non tt background (5%) \rightarrow 18 MeV Precision on \sqrt{s} (10⁻⁴) \rightarrow 30 MeV Uncertainty on lumi-spectrum \rightarrow ~10 MeV Uncertainty on theory x-sec \rightarrow few MeV Conversion 1S mass to $\overline{MS} \rightarrow$ 100 MeV

1S top mass and α_s combined 2D fit		
m_t stat. error m_t theory syst. (1%/3%)	34 MeV 5 MeV / 8 MeV	
α_s stat. error	0.0009	
α_s theory syst. (1%/3%)	0.0008 / 0.0022	

A very precise measurement: $\Delta m_t < 100 \text{ MeV}$

+ $\Delta \alpha s$ < 0.001 (+ $\Delta \Gamma_t$ < 30 MeV) (+ $\Delta y_t/y_t$ ~ 35% *)

* could claim 4.2% (with arXiv:1310.0563) if I insert a more precise value of α_{c}

Top-Z and top-photon couplings and new physics



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Top quark couplings in a nutshell

measure

 $\sigma(+) \quad A_{FB}(+)$ $\sigma(-) \quad A_{FB}(-)$

$$(+=e_{R}^{-}) \\ (-=e_{L}^{-}) \end{cases} \Rightarrow \begin{cases} F_{1V}^{\gamma} \\ F_{1V}^{Z} \\ F_{1V}^{Z} \end{cases}$$

$$F_{1V}^{\gamma} * F_{2V}^{\gamma}$$

$$F_{1V}^{Z} F_{1A}^{Z} F_{2V}^{Z}$$



Measure 2 observables for 2 beam polarizations:

- x-section
- FB asymmetry

Extract form factors in groups (assuming SM for remaining groups)

Assumptions:

LHC: 14 TeV, 300/fb LC: $\sqrt{s} = 500 \text{ GeV}$, L = 500/fb $P(e^{-}) = +/-80\%$, $P(e^{+}) = -/+30\%$ $\delta\sigma \sim 0.5\%$ (stat. + lumi) $\delta A_{_{EB}} \sim 1.8\%$ (stat., covers systematics?)

Polarization needed to disentangle photon and Z-boson form factors!

Especially for ttZ LC precision is better than existing (model-dependent) limits from top decay, LEP T-parameter, B-factories (full comparison in progress)



New physics sensitivity Impact of BSM on Top Sector

In composite Higgs models, the **top quark** is often **partially composite**. This results in **form factors in ttZ couplings**, which can be measured at ILC. **Beam polarization is essential** to distinguish **left/right-handed couplings**.





1



$$\begin{split} &\Gamma_{t\bar{t}}^{\mu}(\gamma,Z) = ie \left[\gamma^{\mu} \Big[\widetilde{F}_{1V}^{\gamma,Z} + \widetilde{F}_{1A}^{\gamma,Z} \gamma^{5} \Big] + \frac{(p_{t} - p_{\bar{t}})^{\mu}}{2m_{t}} \Big[\widetilde{F}_{2V}^{\gamma,Z} + \widetilde{F}_{2A}^{\gamma,Z} \gamma^{5} \Big] \right] \\ &\widetilde{F}_{1v}^{x} = -(F_{1v}^{x} + F_{2v}^{x}), F_{1v}^{y} = -2/3, F_{1v}^{z} = \frac{-1}{4 \operatorname{swcw}} (1 - 8/3 \operatorname{sw}^{2}) \\ &\widetilde{F}_{1x}^{x} = -F_{1A}^{x}, F_{1A}^{y} = 0, F_{1A}^{z} = \frac{-1}{4 \operatorname{swcw}} (1 - 8/3 \operatorname{sw}^{2}) \\ &\widetilde{F}_{2v}^{x} = F_{2v}^{x}, F_{2v}^{y} = Q_{t}(g - 2)/2 \propto d_{x}^{y} \\ &\widetilde{F}_{2x}^{x} = -iF_{2x}^{x}, F_{2x}^{x} \propto d_{x}^{x} \\ \end{split}$$ (Q_t = electric charge, g-2 = anom. magn. Morred of the product of the pro

Close to threshold observables depend on $\rm F_{_{1V}}$ + $\rm F_{_{1A.}}$ Full disentangling imprecise for \sqrt{s} < 1 TeV.

Control over beam polarization is vital to distinguish photon and Z form factors!!

Photon-Z interference brings sensitivity to sign of form factors

CP violating form factors F2A are best measured with special CP observables (TESLA TDR)

For a translation to effective operators language, see J.A. Aguilar Saavedra, Nucl. Phys. B812 (2009), arXiv:1308: $L_{eff} = \Sigma \frac{C_x}{\Lambda^2} O_x$ $X_{tt}^R \propto F_{1V}^Z + F_{1A}^Z + C$ $\delta X_{tt}^R = -\Re \left(C_{\varphi u}^{33} \right) \frac{v^2}{\Lambda^2}$

relations with W-t-b and gluon-t-t vertices explicit dependence on new physics scale Roentsch/Schulze (arXiv:1501.05939) Fiolhais/Aguilar-Saavedra (JHEP 1207, 180) Implemented in WHIZARD (F. Bach)





Most general expression for this vertex...

Impact of new physics



Vary anomalous couplings in narrow range around 0 and register changes in cross-section and $A_{_{FB}}$

Repeat at different center-of-mass energies:

380 (black), 420 (red), 500 (green), 1000 (blue) and 3000 GeV (yellow)

Confirm naïve picture for some operator-observable pairs (larger impact at 3 TeV), but not universal...



Impact of new physics on asymmetry





Impact of new physics on x-section





Top quark reconstruction vs. center-of-mass energy



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Top quark reconstruction

Three different final states:

1) Fully hadronic (46.2%) \rightarrow 6 jets

2) Semi leptonic (43.5%) → 4 jets + 1 charged lepton and a neutrino

3) Fully leptonic (10.3%) \rightarrow 2 jets + 4 leptons





Final state reconstruction uses all detector aspects

Top quark selection/reconstruction



Top reconstruction is non-trivial at any center-of-mass energy

Low energy (~500 GeV):

Challenging combinatorics: migrations due to combining wrong W^+/W^- and b/\overline{b} dilute measurements that rely on top quark reconstruction

Distinguishing top from anti-top with lepton in "lepton+jets" and jet charge in "fully hadronic" final state.



Top quark selection/reconstruction



Top reconstruction is non-trivial at any center-of-mass energy

High energy: top jets \rightarrow no combinatorics for 1 TeV and up!

Provided we can deal with the $\gamma\gamma \rightarrow$ hadrons background in fat jets, top reconstruction at high energy may well be more precise than at low energy!



Reconstruction vs. \sqrt{s}

Angle between W-boson and b-quark that are to form the top candidate $t\bar{t}$ production in MG5_aMC@NLO, no ISR, no luminosity spectrum, no polarization, ----- = correct Wb combination ----- = incorrect combination



Migrations known to disappear for boosted top quarks

Too naïve to expect relative syst. uncertainty to be constant vs. \sqrt{s}



Precision on couplings vs. center-of-mass energy



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First extraction of couplings at $\sqrt{s} \neq 500$ GeV

Rerun the extraction of the couplings from measurement of σ , A_{FR} (Roman Poeschl, LAL)

- set fixed integrated luminosity: 2×250 /fb, with P = (+80, -30) and P = (-80, +30), at any center-of-mass energy
- cross-section initially ~constant: σ = 550 pb at 380 GeV, 530 pb at 500 GeV, then rapid drop-off
- the value of A_{FB} drops rapidly as sqrt(s) $\rightarrow 2 m_t$
- assuming stat. dominated uncertainty: $\delta A_{FB} = (1 A_{FB}^2) \times \delta \sigma / \sigma$



For the F1V couplings we find excellent results also at 420 GeV Drop in x-section at center-of-mass energy ≥ 1 TeV only partially recovered by greater instantaneous luminosity \rightarrow sensitivity for F1V degraded by factor 5-10



Nominal beam polarization $(e^{-} 80\%, e^{+} 30\%)$

Electron polarization only

Extraction of axial coupling





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Extraction of dipole moments F2V





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Summary

Lepton colliders can get a very precise top quark mass

- \rightarrow stat. error \sim few tens of MeV
- \rightarrow exp. syst. error \sim few tens of MeV
- \rightarrow theory error \sim 10 MeV
- \rightarrow 1S \rightarrow $\overline{\rm MS}$ conversion \sim 100 MeV

Carefully distinguish targets from prospects

tt Z and tt γ couplings measurement to < 1% are unique opportunity at a lepton collider

Every sqrt(s) regime brings additional potential;

 \rightarrow Coupling measurement (in particular F1AZ) has sweet spot around 420-700 GeV

where $A_{_{FR}}$ and cross-section are large and precise calculations are "easy"

- \rightarrow Dipole moments and 4-fermion contact interactions might show high-scale NP in TeV regime
- \rightarrow Polarization is needed to disentangle photon and Z couplings,

Much more to explore:

top physics potential below threshold



The top quark mass combination, small print

all measurements considered in the present combination, the analyses are calibrated to the Monte Carlo (MC) top-quark mass definition. It is expected that the difference between the MC mass definition and the formal pole mass of the top quark is up to the order of 1 GeV (see Refs. [19,20] and references therein).

to jet calibration and modelling of the $t\bar{t}$ events. Given the current experimental uncertainty on m_{top} , clarifying the relation between the top quark mass implemented in the MC and the formal top quark pole mass demands further theoretical investigations. The dependence of the result on the correlation assumptions between mea-

[19] General-purpose generators for particle physics

Note: it's likely that the GeV is at least partially accounted for in current modelling uncertainty



Scheme dependence – an old debate

Even if it decays (rather than hadronizes) the top is a quark, a coloured object. Mass is not an observable, but must be inferred from measurements.

The scheme makes a difference:

For a top pole mass of 173 GeV, the $\overline{\text{MS}}$ mass at the top mass ~167 GeV

Quantify the difference between pole and any other mass scheme Hoang & Stewart, Nucl.Phys.Proc.Suppl. 185 (2008) 220-226

$$m_{t}^{pole} = m_{t}(R,\mu) + R \Sigma_{k} \Sigma_{n} a_{nk} \left[\frac{\alpha_{s}(\mu)}{4\pi}\right]^{n} \ln^{k}\left(\frac{\mu}{R}\right) \qquad 180 \qquad 180 \qquad \overline{m(m)} \\ 170 \qquad 170 \qquad \overline{m(R)} \\ 160 \qquad 160 \qquad 160 \qquad \overline{m(R)} \\ 160 \qquad 150 \qquad \overline{R} = m(R) \\ 150 \qquad 0 \qquad 150 \qquad \overline{R}$$



Top quark mass: interpretation

Which top quark mass did we implement in our MC?

Matrix Element: ~ pole mass (for NLO)

Parton Shower: ~ pole mass

Hadronization: ?

In practice, the impact on the measurement depends on how sensitive a given observable is to soft physics. Unfortunately, the three-jet invariant mass used to measure the most precise top quark mass is quite sensitive

 $R \sim \Gamma_t \sim PS$ cut-off and:

$$m_t^{pole} = m_t^{MC}(R) + R\alpha_s \frac{(\mu)}{4\pi}$$

Is there a clear and universal relation between the two masses? Can we "discover" this relation?

CMS-FTR-13-017-PAS: The relation between the pole mass and the MC top-quark mass as "not an experimental problem, but a theoretical (or phenomenological) issue."

Theorists may actually manage to do this: see for a serious attempt for bottom: A. Hoang, LCWS14

Perspective for improvement systematics on combination



Break-down of uncertainties on March '14 world average:

Jet energy scale: in situ JES (240 MeV), standardJES (200 MeV), flavourJES (120 MeV) and b-JES (250 MeV)

Statistics:

already < 300 MeV

Modelling:

(strongly correlated even between experiments): Monte Carlo (380 MeV) radiation (210 MeV) colour reconnection (310 MeV)

For a long time we claimed an LHC precision of 1 GeV

Prospect studies for top quark mass precision at Snowmass reported in arXiv:1310.0799, that I sign, concluded: "We estimate that [...] might lead to a top mass extraction with uncertainty as low as 500-600 MeV"



CMS-FTR-13-017-PAS claims the ultimate reach of the "conventional method is 200 MeV, based on "assumptions [that] are optimistic but not unrealistic."

Clearly, the 200 MeV require a lot of work on JES and generators. Time will tell...

Top quark mass - alternatives

Endpoint measurement

CMS, arXiv:1304.5783, currently 2 GeV uncertainty)

CMS estimate 600 MeV precision after the complete LHC programme

Move away from jets

(reduced dependence on shower modelling and JES)

- Extraction from $m_{\mbox{\tiny bl}}$
- Extraction from J/psi spectra t \rightarrow Wb \rightarrow lvb \rightarrow lvJ/ ϕ \rightarrow lvII

How well can we predict top quark





Theory milestone:

full NNLO and NNLL result for top quark pair production at hadron colliders

K-factor (NLO \rightarrow NNLO) ~ 10% Scale stability ~ 5 % Series seems to converge...

Collider	$\sigma_{ m tot}~[m pb]$	scales [pb]	pdf [pb]
Tevatron	7.009	+0.259(3.7%) -0.374(5.3%)	+0.169(2.4%) -0.121(1.7%)
LHC 7 TeV	167.0	+6.7(4.0%) -10.7(6.4%)	+4.6(2.8%) -4.7(2.8%)
LHC 8 TeV	239.1	+9.2(3.9%) -14.8(6.2\%)	$+6.1(2.5\%) \\ -6.2(2.6\%)$
LHC 14 TeV	933.0	$+31.8(3.4\%) \\ -51.0(5.5\%)$	+16.1(1.7%) -17.6(1.9%)

Alternative: top mass from cross-section



Nearly flat, negligible residual MC mass dependence

Top quark mass

Extraction from cross section - revisited

Well-defined mass scheme (pole mass, $\overline{\text{MS}}$ mass) Limited by poor sensitivity: $\Delta m/m \sim 0.2 \Delta \sigma/\sigma$ tt threshold has better sensitivity, but requires theory progress (bound states) currently ~4 GeV uncertainty, PLB728 (2014) 496-517

Now consider the $t\bar{t}g$ cross-section

Alioli, Moch, Uwer, Fuster, Irles, Vos, EPJC73 (2013) 2438, arXiv:1303.6415



Top quark mass

- Measure the normalized differential tt+1jet production cross-section vs. Invariant mass of the tt+1jet system.
- Extract the mass in any (welldefined) scheme. Currently: pole mass
- Theory uncertainty (due to scale and PDF) < 1 GeV
- Experimental uncertainties can be controlled to same level



- $ρ_s μ 1/m(t\bar{t}j)$
 - \rightarrow 1 at threshold
 - \rightarrow 0 for boosted production

Top quark mass from tt + 1 jet events



(difference NLO vs. NLO+PS ~ 300 MeV)

Top quark mass

- Don't you run into a MC mass dependence in the correction of the normalized differential cross-section?
- No, compatible results are obtained for a large range of MC mass values.



Top quark pole mass

M^{pole}=173.7 +2.3 -2.1 GeV

 \rightarrow

 \rightarrow currently the most precise top quark pole mass

1.5 (stat.), 1.4 (syst.), +1.0/-0.5 (theo.)

 \rightarrow room for improvement, even with 2012 data set (ongoing)

