

Electroweak precision tests: theory uncertainties

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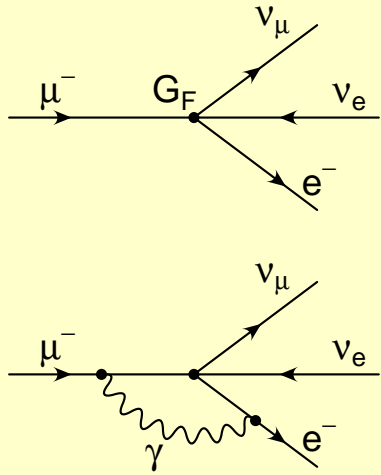
University of Pittsburgh

FCC-ee Physics Workshop, 3-5 Feb 2015

- 1. Electroweak precision observables**
- 2. Current status of SM loop results**
- 3. Future projections**
- 4. Parametrization of new physics**

W mass

μ decay in Fermi Model



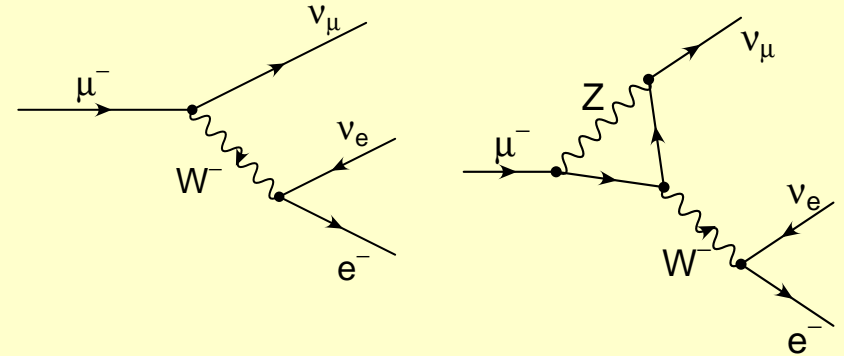
← QED corr.
(2-loop)

$$\Gamma_\mu = \frac{G_F^2 m_\mu^5}{192\pi^3} F\left(\frac{m_e^2}{m_\mu^2}\right) (1 + \Delta q)$$

Ritbergen, Stuart '98

Pak, Czarnecki '08

μ decay in Standard Model



$$\frac{G_F^2}{\sqrt{2}} = \frac{e^2}{8s_w^2 M_W^2} (1 + \Delta r)$$

electroweak corrections

$e^+e^- \rightarrow f\bar{f}$ for $\sqrt{s} \sim m_Z$:

$$\sigma = \mathcal{R}_{\text{ini}} \left[12\pi \frac{\Gamma_{ee}\Gamma_{ff}}{(s - m_Z^2)^2 - m_Z^2\Gamma_Z^2} (1 - \mathcal{P}_e\mathcal{A}_e) + \sigma_{\text{non-res}} \right],$$

$$\Gamma_{ff} = \mathcal{R}_V^f g_{Vf}^2 + \mathcal{R}_A^f g_{Af}^2, \quad \Gamma_Z = \sum_f \Gamma_{ff},$$

$$\mathcal{A}_f = 2 \frac{g_{Vf}/g_{Af}}{1 + (g_{Vf}/g_{Af})^2} = \frac{1 - 4|Q_f| \sin^2 \theta_{\text{eff}}^f}{1 - 4|Q_f| \sin^2 \theta_{\text{eff}}^f + 8(|Q_f| \sin^2 \theta_{\text{eff}}^f)^2}.$$

$$A_{\text{FB}}^f = \frac{3}{4} \mathcal{A}_e \mathcal{A}_f \quad A_{\text{LR}} = \mathcal{A}_e$$

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QED/QCD corrections on ext. fermions

Chetyrkin, Kataev, Tkachov '79

Dine, Saphirstein '79

Celmaster, Gonsalves '80

Gorishnii, Kataev, Larin '88,91

Chetyrkin, Kühn '90

Surguladze, Samuel '91

Kataev '92

Chetyrkin '93

etc...

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$$A_{\text{FB}}^f = \frac{3}{4} \mathcal{A}_e \mathcal{A}_f \quad A_{\text{LR}} = \mathcal{A}_e$$

additional initial-state QED corrections

Kuraev, Fadin '85

Berends, Burgers, v. Neerven '88

Kniehl, Krawczyk, Kühn, Stuart '88

Beenakker, Berends, v. Neerven '89

Bardin et al. '89,91

Montagna, Nicosini, Piccinini '97

etc...

$e^+e^- \rightarrow f\bar{f}$ for $\sqrt{s} \sim m_Z$:

$$\sigma = \mathcal{R}_{\text{ini}} \left[12\pi \frac{\Gamma_{ee}\Gamma_{ff}}{(s - m_Z^2)^2 - m_Z^2\Gamma_Z^2} (1 - \mathcal{P}_e\mathcal{A}_e) + \sigma_{\text{non-res}} \right],$$

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$$A_{\text{FB}}^f = \frac{3}{4} \mathcal{A}_e \mathcal{A}_f \quad A_{\text{LR}} = \mathcal{A}_e$$

electroweak corrections

Known corrections to Δr , $\sin^2 \theta_{\text{eff}}^f$, g_{Vf} , g_{Af} :

- Complete NNLO corrections (Δr , $\sin^2 \theta_{\text{eff}}^l$) Freitas, Hollik, Walter, Weiglein '00
Awramik, Czakon '02; Onishchenko, Veretin '02
Awramik, Czakon, Freitas, Weiglein '04; Awramik, Czakon, Freitas '06
Hollik, Meier, Uccirati '05,07; Degrossi, Gambino, Giardino '14
- “Fermionic” NNLO corrections (g_{Vf} , g_{Af}) Czarnecki, Kühn '96
Harlander, Seidensticker, Steinhauser '98
Freitas '13,14
- Partial 3/4-loop corrections to ρ/T -parameter
 $\mathcal{O}(\alpha_t \alpha_s^2)$, $\mathcal{O}(\alpha_t^2 \alpha_s)$, $\mathcal{O}(\alpha_t \alpha_s^3)$ Chetyrkin, Kühn, Steinhauser '95
Faisst, Kühn, Seidensticker, Veretin '03
Boughezal, Tausk, v. d. Bij '05
Schröder, Steinhauser '05; Chetyrkin et al. '06
Boughezal, Czakon '06

$$(\alpha_t \equiv \frac{y_t^2}{4\pi})$$

	Experiment	Theory error	Main source
M_W	80.385 ± 0.015 MeV	4 MeV	$\alpha^3, \alpha^2\alpha_s$
Γ_Z	2495.2 ± 2.3 MeV	0.5 MeV	$\alpha_{\text{bos}}^2, \alpha^3, \alpha^2\alpha_s, \alpha\alpha_s^2$
σ_{had}^0	41540 ± 37 pb	6 pb	$\alpha_{\text{bos}}^2, \alpha^3, \alpha^2\alpha_s$
$R_b \equiv \Gamma_Z^b / \Gamma_Z^{\text{had}}$	0.21629 ± 0.00066	0.00015	$\alpha_{\text{bos}}^2, \alpha^3, \alpha^2\alpha_s$
$\sin^2 \theta_{\text{eff}}^l$	0.23153 ± 0.00016	4.5×10^{-5}	$\alpha^3, \alpha^2\alpha_s$

Methods for theory error estimates:

- Parametric factors, *i. e.* factors of α, N_c, N_f, \dots
- Geometric progression, *e. g.* $\frac{\mathcal{O}(\alpha^3)}{\mathcal{O}(\alpha^2)} \sim \frac{\mathcal{O}(\alpha^2)}{\mathcal{O}(\alpha)}$
- Renormalization scale dependence (often underestimates error)
- Renormalization scheme dependence (may underestimate error)

■ Geometric perturbative series

$$\alpha_t = \alpha m_t^2$$

$$\mathcal{O}(\alpha^3) - \mathcal{O}(\alpha_t^3) \sim \frac{\mathcal{O}(\alpha^2) - \mathcal{O}(\alpha_t^2)}{\mathcal{O}(\alpha)} \mathcal{O}(\alpha^2) \sim 0.26 \text{ MeV}$$

$$\mathcal{O}(\alpha^2 \alpha_s) - \mathcal{O}(\alpha_t^2 \alpha_s) \sim \frac{\mathcal{O}(\alpha^2) - \mathcal{O}(\alpha_t^2)}{\mathcal{O}(\alpha)} \mathcal{O}(\alpha \alpha_s) \sim 0.30 \text{ MeV}$$

$$\mathcal{O}(\alpha \alpha_s^2) - \mathcal{O}(\alpha_t \alpha_s^2) \sim \frac{\mathcal{O}(\alpha \alpha_s) - \mathcal{O}(\alpha_t \alpha_s)}{\mathcal{O}(\alpha)} \mathcal{O}(\alpha \alpha_s) \sim 0.23 \text{ MeV}$$

$$\mathcal{O}(\alpha \alpha_s^3) - \mathcal{O}(\alpha_t \alpha_s^3) \sim \frac{\mathcal{O}(\alpha \alpha_s) - \mathcal{O}(\alpha_t \alpha_s)}{\mathcal{O}(\alpha)} \mathcal{O}(\alpha \alpha_s^2) \sim 0.035 \text{ MeV}$$

$$\mathcal{O}(\alpha_{\text{bos}}^2) \sim \mathcal{O}(\alpha_{\text{bos}})^2 \sim 0.1 \text{ MeV}$$

■ Parametric prefactors: $\mathcal{O}(\alpha_{\text{bos}}^2) \sim \Gamma_Z \alpha^2 \sim 0.1 \text{ MeV}$

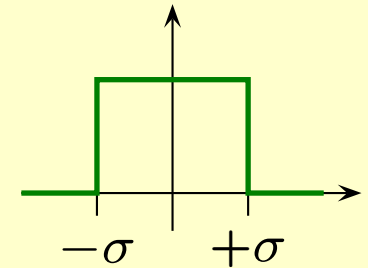
$$\mathcal{O}(\alpha \alpha_s^2) - \mathcal{O}(\alpha_t \alpha_s^2) \sim \frac{\alpha n_{|q}}{\pi} \alpha_s^2 \sim 0.29 \text{ MeV}$$

Total: $\delta \Gamma_Z \approx 0.5 \text{ MeV}$

- Add theory errors from each source **linearly**:

Idea: each value within error range is equally likely

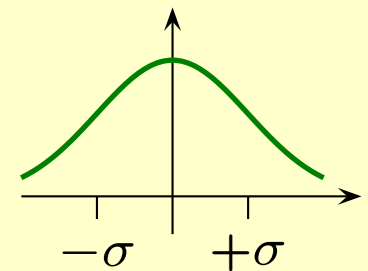
→ Use flat prior in global fits



- Add theory errors from each source **quadratically**:

Idea: different error sources are uncorrelated

→ Use Gaussian prior in global fits (central limit theorem)



	Current exp.	ILC	TLEP	Current perturb.
M_W [MeV]	15	3–5	~ 1	4
Γ_Z [MeV]	2.3	~ 1	~ 0.1	0.5
R_b [10^{-5}]	66	15	$\lesssim 5$	15
$\sin^2 \theta_{\text{eff}}^\ell$ [10^{-5}]	16	1.3	0.3	4.5

→ Existing theoretical calculations adequate for LEP/SLC/LHC,
but not ILC/TLEP!

	ILC	TLEP	perturb. error with 3-loop [†]	Param. error ILC*	Param. error TLEP**
M_W [MeV]	3–5	~ 1	1	2.6	1
Γ_Z [MeV]	~ 1	~ 0.1	$\lesssim 0.2$	0.5	?
R_b [10^{-5}]	15	$\lesssim 5$	5–10	< 1	< 1
$\sin^2 \theta_{\text{eff}}^l$ [10^{-5}]	1.3	0.3	1.5	2	2

[†] **Theory scenario:** $\mathcal{O}(\alpha\alpha_S^2)$, $\mathcal{O}(N_f\alpha^2\alpha_S)$, $\mathcal{O}(N_f^2\alpha^2\alpha_S)$
 (N_f^n = at least n closed fermion loops)

Parametric inputs:

* **ILC:** $\delta m_t = 100$ MeV, $\delta\alpha_S = 0.001$, $\delta M_Z = 2.1$ MeV

****TLEP:** $\delta m_t \lesssim 50$ MeV, $\delta\alpha_S = ?$, $\delta M_Z = 0.1$ MeV

also: $\delta(\Delta\alpha) = 5 \times 10^{-5}$

■ Subtraction of QED radiation contributions

→ Known to $\mathcal{O}(\alpha^2)$, $\mathcal{O}(\alpha^3 L^3)$ for **ISR**,
 $\mathcal{O}(\alpha^2)$ for **FSR** and $\mathcal{O}(\alpha^2 L^2)$ for **A_{FB}**

$$(L = \log \frac{s}{m_e^2})$$

Berends, Burgers, v.Neerven '88

Kniehl, Krawczyk, Kühn, Stuart '88

Beenakker, Berends, v.Neerven '89

Skrzypek '92; Montagna, Nicrosini, Piccinini '97

→ $\mathcal{O}(0.1\%)$ uncertainty on σ_Z , A_{FB}

→ Improvement needed for ILC/TLEP

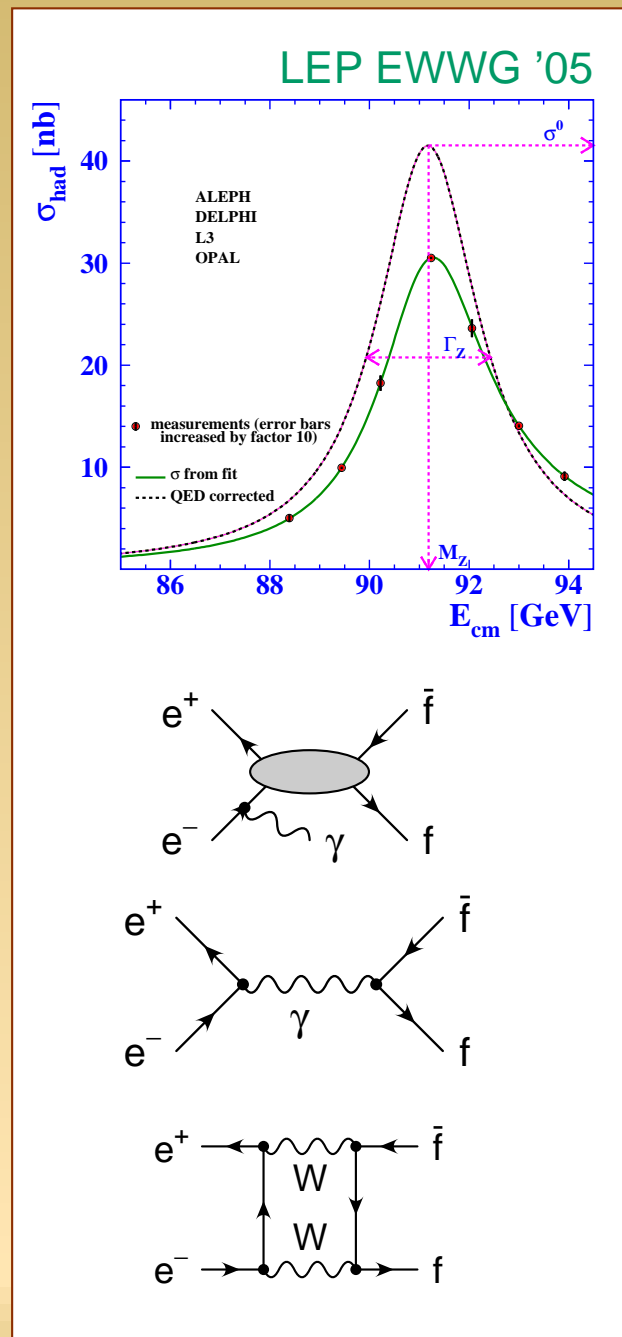
■ Subtraction of non-resonant γ -exchange, γ -Z interf., box contributions, Bhabha scattering

see, e.g., Bardin, Grünewald, Passarino '99

→ $\mathcal{O}(0.01\%)$ uncertainty within SM

(improvements may be needed)

→ Sensitivity to some NP beyond EWPO



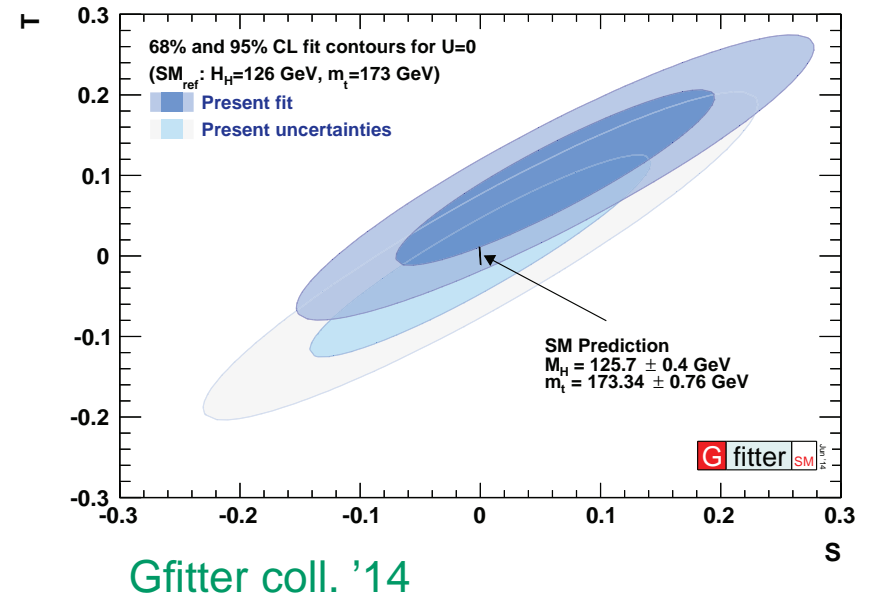
Oblique parameters:

$$\alpha T = \frac{\Sigma_{WW}(0)}{M_W} - \frac{\Sigma_{ZZ}(0)}{M_Z}$$

$$\frac{\alpha}{4s^2c^2} S = \frac{\Sigma_{ZZ}(M_Z^2) - \Sigma_{ZZ}(0)}{M_Z} + \frac{s^2 - c^2}{sc} \frac{\Sigma_{Z\gamma}(M_Z^2)}{M_Z} - \frac{\Sigma_{\gamma\gamma}(M_Z^2)}{M_Z}$$

$$\frac{\alpha}{4s^2} (S + U) = \frac{\Sigma_{WW}(M_W^2) - \Sigma_{WW}(0)}{M_W} - \frac{c}{s} \frac{\Sigma_{Z\gamma}(M_Z^2)}{M_Z} - \frac{\Sigma_{\gamma\gamma}(M_Z^2)}{M_Z}$$

→ Not adequate for new physics that affects flavor ($Z \rightarrow \ell\ell$, $Z \rightarrow bb$, ...)



More general setup: Use pseudo-observables

$$M_W, \Gamma_Z, \sigma_{\text{had}}^0, R_b, R_\ell, A_\ell, A_b, A_{lq} \quad (\ell = e, \mu, \tau) \quad \rightarrow 12 \text{ quantities}$$

Effective field theory: $\mathcal{L} = \sum_i \frac{c_i}{\Lambda^2} \mathcal{O}_i + \mathcal{O}(\Lambda^{-3}) \quad (\Lambda \gg M_Z)$

$$\mathcal{O}_{\phi 1} = (D_\mu \Phi)^\dagger \Phi \Phi^\dagger (D^\mu \Phi) \quad \alpha \Delta T = -\frac{v^2}{2} \frac{c_{\phi 1}}{\Lambda^2}$$

$$\mathcal{O}_{\text{BW}} = \Phi^\dagger B_{\mu\nu} W^{\mu\nu} \Phi \quad \alpha \Delta S = -e^2 v^2 \frac{c_{\text{BW}}}{\Lambda^2}$$

$$\mathcal{O}_{\text{LL}}^{(3)e} = (\bar{L}_L^e \sigma^a \gamma_\mu L_L^e) (\bar{L}_L^e \sigma^a \gamma^\mu L_L^e) \quad \Delta G_F = -\sqrt{2} \frac{c_{\text{LL}}^{(3)e}}{\Lambda^2}$$

$$\mathcal{O}_R^f = i(\Phi^\dagger \overleftrightarrow{D}_\mu \Phi) (\bar{f}_R \gamma^\mu f_R) \quad f = e, \mu, \tau, b, lq$$

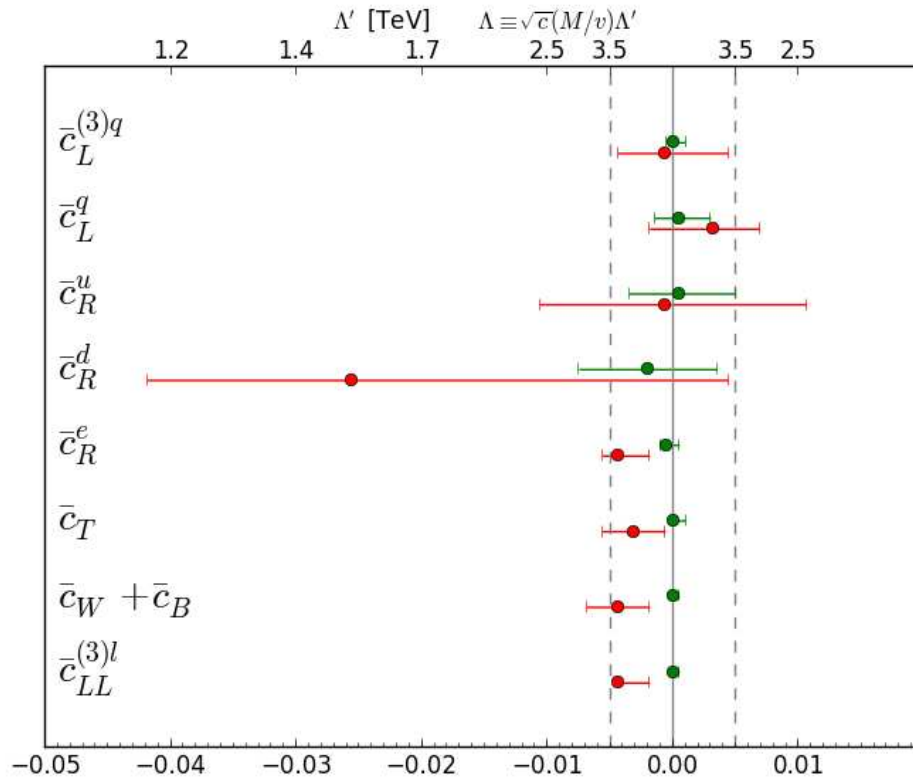
$$\mathcal{O}_L^F = i(\Phi^\dagger \overleftrightarrow{D}_\mu \Phi) (\bar{F}_L \gamma^\mu F_L) \quad F = \begin{pmatrix} \nu_e \\ e \end{pmatrix}, \begin{pmatrix} \nu_\mu \\ \mu \end{pmatrix}, \begin{pmatrix} \nu_\tau \\ \tau \end{pmatrix}, \begin{pmatrix} u, c \\ d, s \end{pmatrix}, \begin{pmatrix} t \\ b \end{pmatrix}$$

$$\mathcal{O}_L^{(3)F} = i(\Phi^\dagger \overleftrightarrow{D}_\mu^a \Phi) (\bar{F}_L \sigma_a \gamma^\mu F_L)$$

More operators than EWPOs

\rightarrow Some can be constrained by $W \rightarrow \ell\nu$, had., $e^+e^- \rightarrow W^+W^-$

Assuming flavor universality:



Significant correlation/
degeneracy between
different operators

Pomaral, Riva '13
Ellis, Sanz, You '14

- Experimental precision from LEP/SLC demands SM prediction with **2-loop corrections** and **partial 3-loop corrections**
- **LHC** will provide independent results for $\sin^2 \theta_{\text{eff}}$ and M_W , but overall precision not improved
- **ILC/TLEP** with $\sqrt{s} \sim M_Z$ will reduce exp. error of some EWPO by $\mathcal{O}(10)$
→ 3-loop (and maybe some 4-loop) corrections needed!
- Procedure for combination of experimental, perturbative and parametric uncertainties needs consensus
- At current and future level of precision, new physics should not be parametrized just by S/T/U parameters
 - For TeV-scale physics: use effective operators (which is the best basis?)
 - For new physics at weak scale and below: directly compute pseudo-obs.

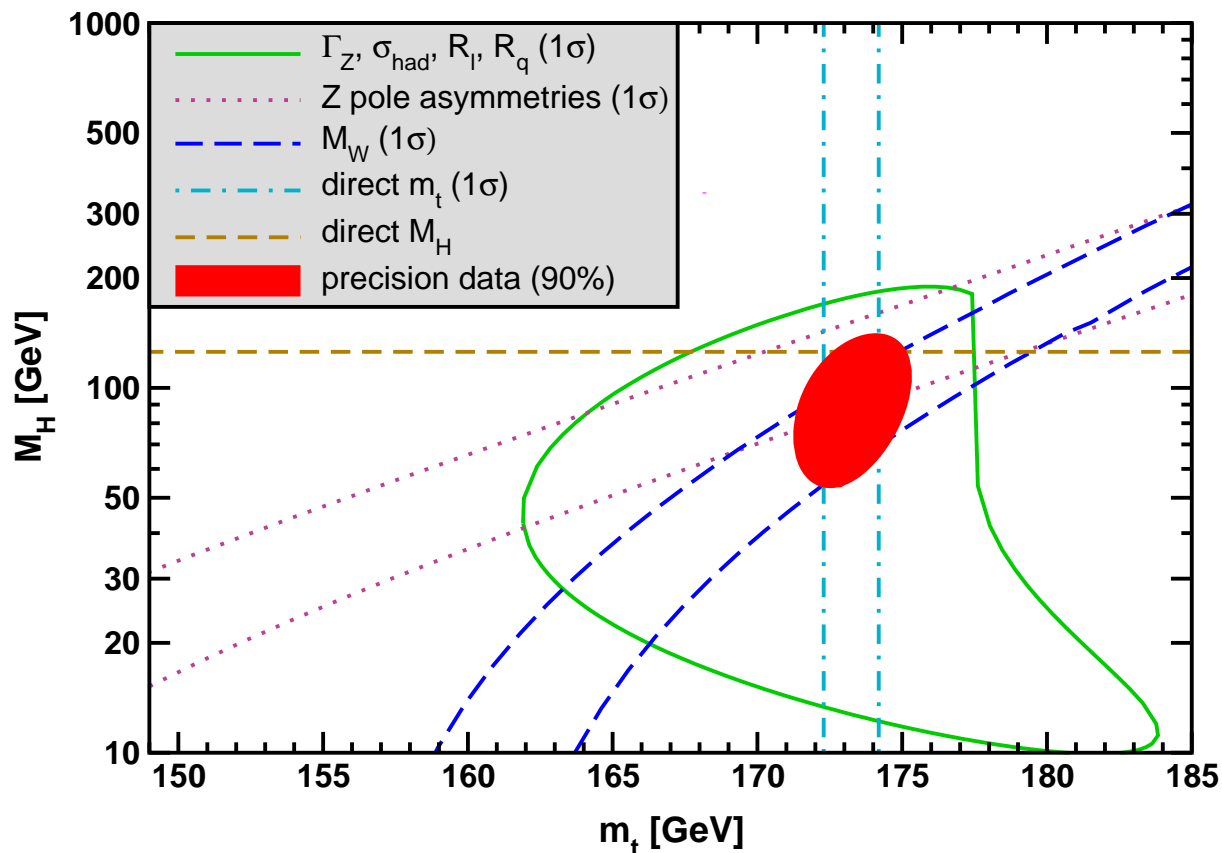
Backup slides

Electroweak precision fit

Standard Model after Higgs discovery:

- Very good agreement over large number of observables
- Sensitivity to TeV-scale new physics

Erlar '13



Direct measurements:

$$M_H = 125.6 \pm 0.4 \text{ GeV}$$
$$m_t = 173.24 \pm 0.95 \text{ GeV}$$

Indirect prediction:

$$M_H = 123.7 \pm 2.3 \text{ GeV}$$

(with LHC BRs)

$$M_H = 89^{+22}_{-18} \text{ GeV}$$

(w/o LHC data)

$$m_t = 177.0 \pm 2.1 \text{ GeV}$$

Z-pole observables

After deconvolution of initial-state QED radiation and subtraction of γ -exchange:

$$\mathcal{A}[e^+e^- \rightarrow f\bar{f}] = \frac{R}{s - s_0} + S + (s - s_0)S'$$

$$s_0 \equiv M_Z^2 - iM_Z\Gamma_Z$$

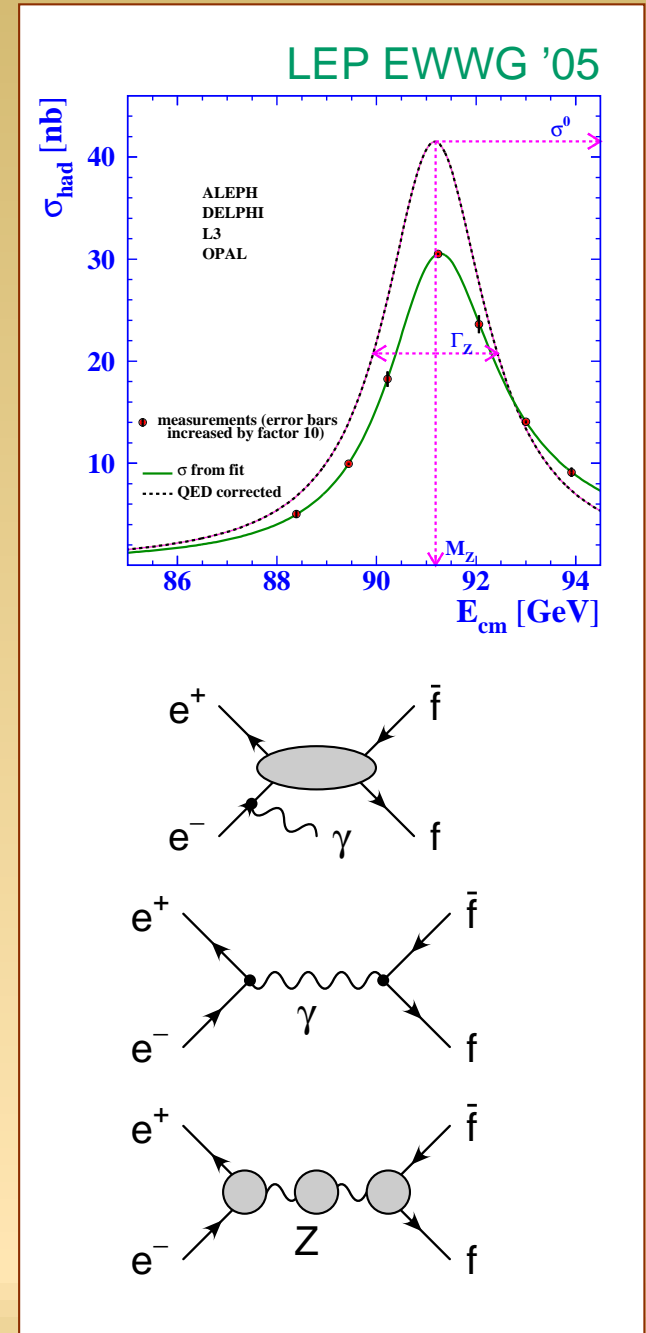
s_0 , R , S , S' are gauge-invariant

Willenbrock, Valencia '91; Sirlin '91; Stuart '91
Gambino, Grassi '00

$$\sigma \sim \frac{1}{(s - M_Z^2)^2 + M_Z^2\Gamma_Z^2} + \text{non.res.}$$

$$M_Z = M_Z^{\text{exp}} - 34 \text{ MeV}$$

$$\Gamma_Z = \Gamma_Z^{\text{exp}} - 0.9 \text{ MeV}$$



Electroweak precision tests: new physics

Constraints on Higgs physics:

- Indirect determination of M_H
- Couplings: need to go beyond SM, e.g. THDM:

$$\left| \frac{g_{hVV}^{\text{THDM}}}{g_{hVV}^{\text{SM}}} \right| = \sin(\beta - \alpha),$$

$$\left| \frac{g_{hff}^{\text{THDM}}}{g_{hff}^{\text{SM}}} \right| = \frac{\cos \alpha}{\sin \alpha} \quad \text{or} \quad \frac{\sin \alpha}{\cos \alpha}$$

- Couplings: effective operators

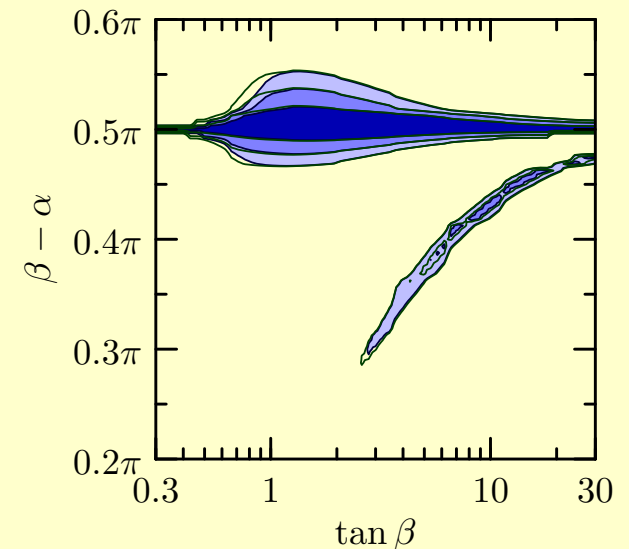
$$\mathcal{O}_W = (D^\mu \Phi)^\dagger W_{\mu\nu} (D^\nu \Phi)$$

$$\mathcal{O}_{WW} = \Phi^\dagger W^{\mu\nu} W_{\mu\nu} \Phi$$

etc.

Mebane, Greiner, Zhang, Willenbrock '13

Chen, Dawson, Zhang '13



Eberhardt, Nierste, Wiebusch '13

