

Heavy quark physics

LaThuile 2015: XXIX Rencontres de Physique de la Vallee d'Aoste
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CMS collaboration

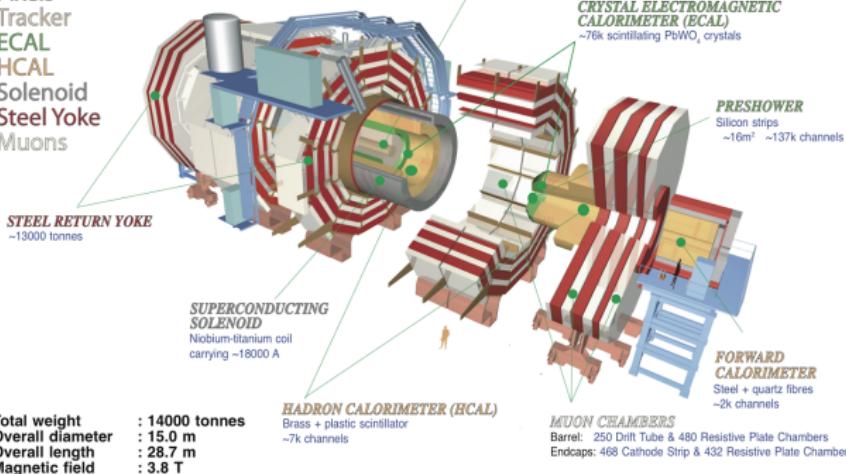
March 3, 2015



The Compact Muon Solenoid detector

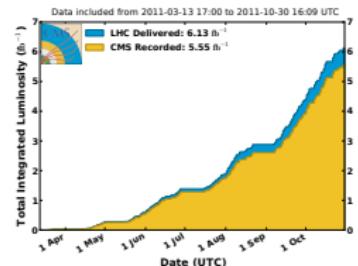
CMS Detector

Pixels
Tracker
ECAL
HCAL
Solenoid
Steel Yoke
Muons

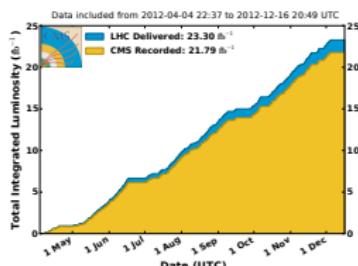


- Dimuon mass resolution: 32 MeV (barrel) and 75 MeV (endcap)
- Impact parameter resolution: 10 μm @ 100 GeV and 20 μm @ 10 GeV

CMS Integrated Luminosity, pp, 2011, $\sqrt{s} = 7 \text{ TeV}$



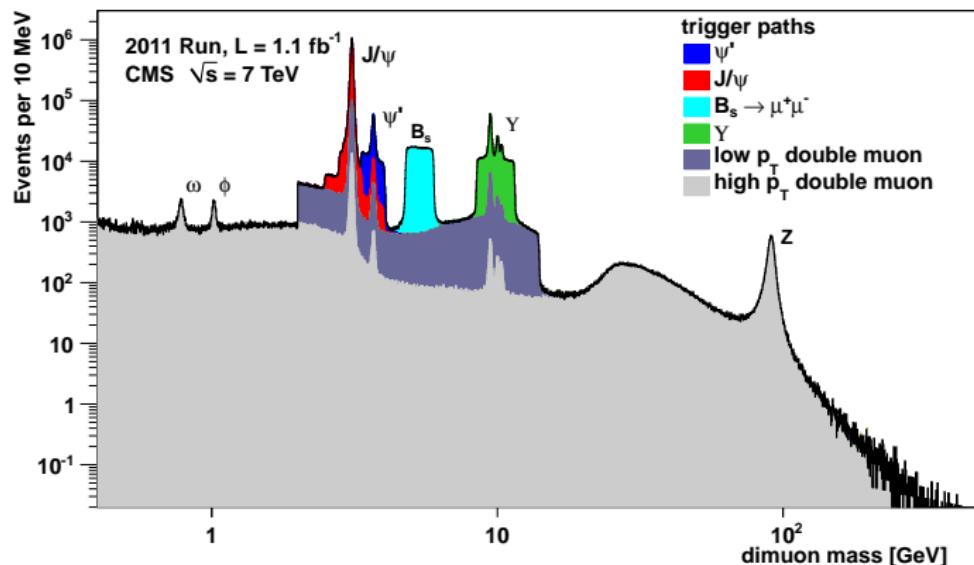
CMS Integrated Luminosity, pp, 2012, $\sqrt{s} = 8 \text{ TeV}$



- 2010: $L=10^{32} \text{ cm}^{-2}\text{s}^{-1}$
- 2011: $L=4 \cdot 10^{33} \text{ cm}^{-2}\text{s}^{-1}$
- 2012: $L=7.5 \cdot 10^{33} \text{ cm}^{-2}\text{s}^{-1}$



Heavy flavour physics triggers



Trigger selection

- Muon kinematics
- Dimuon kinematics and vertex probability
- Flight length significance and pointing angle

Outline

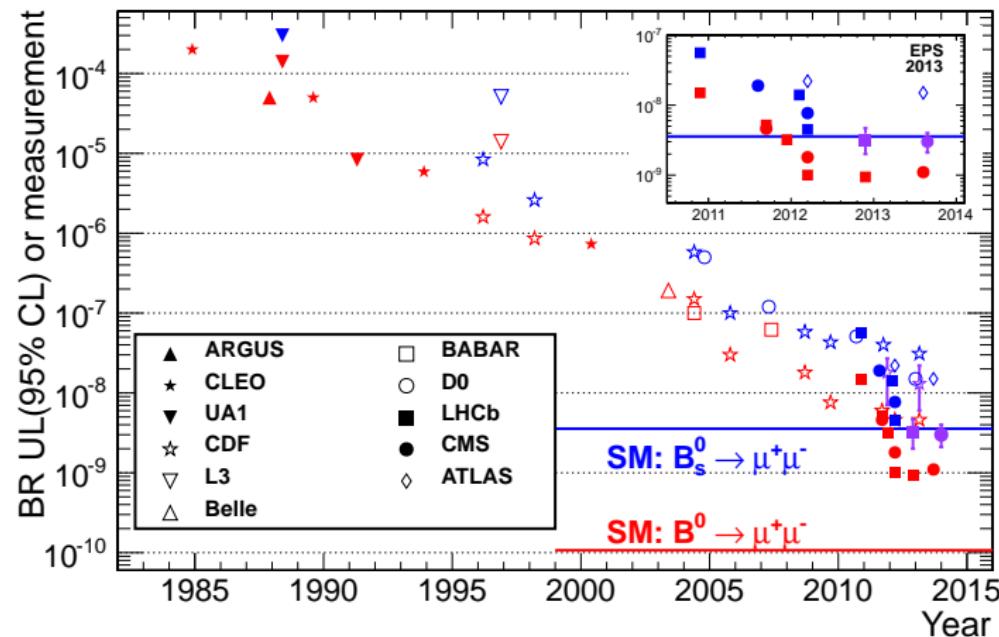
Recent results on physics of beauty and charm quarks in proton-proton collisions at CMS:

- Search for physics beyond the standard model (by the mean of indirect evidence)
 - * Search for $B_{s,d} \rightarrow \mu\mu$ decays at CMS
 - * CP-violating phase ϕ_s using $B_s \rightarrow J/\psi\phi(1020)$ channel
- Study of standard model of particle physics, test of QCD and particle properties
 - * $B^+ \rightarrow \psi(2S)\phi K^+$ rare decay
 - * B_c decays
 - * $\mathcal{B}(B_s \rightarrow J/\psi f_0(980)) / \mathcal{B}(B_s \rightarrow J/\psi\phi(1020))$ ratio
 - * $J/\psi, \psi(2S), \Upsilon(nS), \text{and } \chi_{bn}$ production

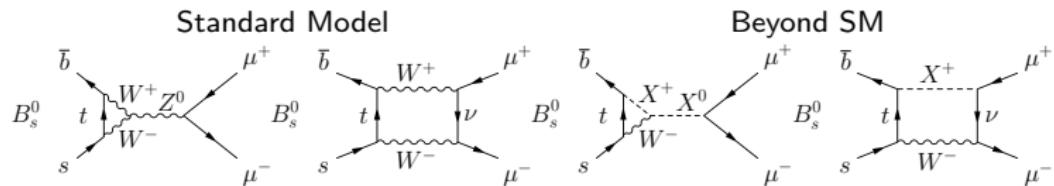


Search for $B_{s,d} \rightarrow \mu\mu$ decays at CMS

$B_{s,d} \rightarrow \mu\mu$ search history



$B_{s,d} \rightarrow \mu\mu$ motivation



- Highly suppressed decay in SM
 - Forbidden at tree level \rightarrow FCNC transitions only possible through penguin or box diagram
 - Cabibbo $|V_{td}| < |V_{ts}|$ and helicity suppressed
- SM predictions
 - $\mathcal{B}(B_s \rightarrow \mu\mu) = (3.66 \pm 0.23) \cdot 10^{-9}$ (PRL112, 101801, 2014)
 - $\mathcal{B}(B_d \rightarrow \mu\mu) = (1.06 \pm 0.10) \cdot 10^{-10}$ (PRL112, 101801, 2014)
- Sensitivity to NP
 - 2HDM and $m(H^\pm)$
 - MSSM $\tan \beta$
 - Leptoquarks
 - 4th generation top

$B_s \rightarrow \mu\mu$ CMS measurement [PRL 111 (2013) 101804]

- Selections:

- BDT muon selection → low muon fake rate
- Selection done using a BDT with 12 input variables for B candidates
- Trained with MC signals sample and with data sidebands
- Same BDT for $B^\pm \rightarrow J/\psi K^\pm$ normalisation and $B_s \rightarrow J/\psi \phi(1020)$ control channels
- No pile-dependence, no B_s mass dependence

- BDT categories:

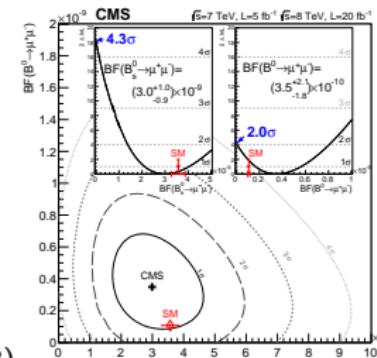
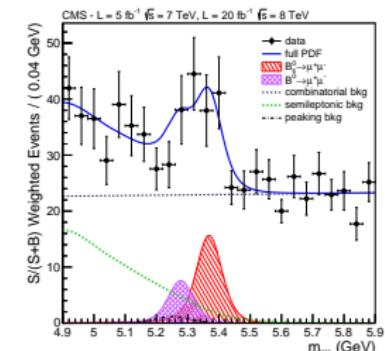
- BDT response divided in multiple bins depending on data period
2012(4) or 2011(2)
- Two further categories: barrel and endcap regions (different mass resolution)

- Unbinned ML fit with per-event resolution:

- Signal: fit of both B_s and B^0 shapes
- Peaking backgrounds ($B^0 \rightarrow K\pi$, $B_s \rightarrow KK$)
- Combinatorial and semileptonic BG

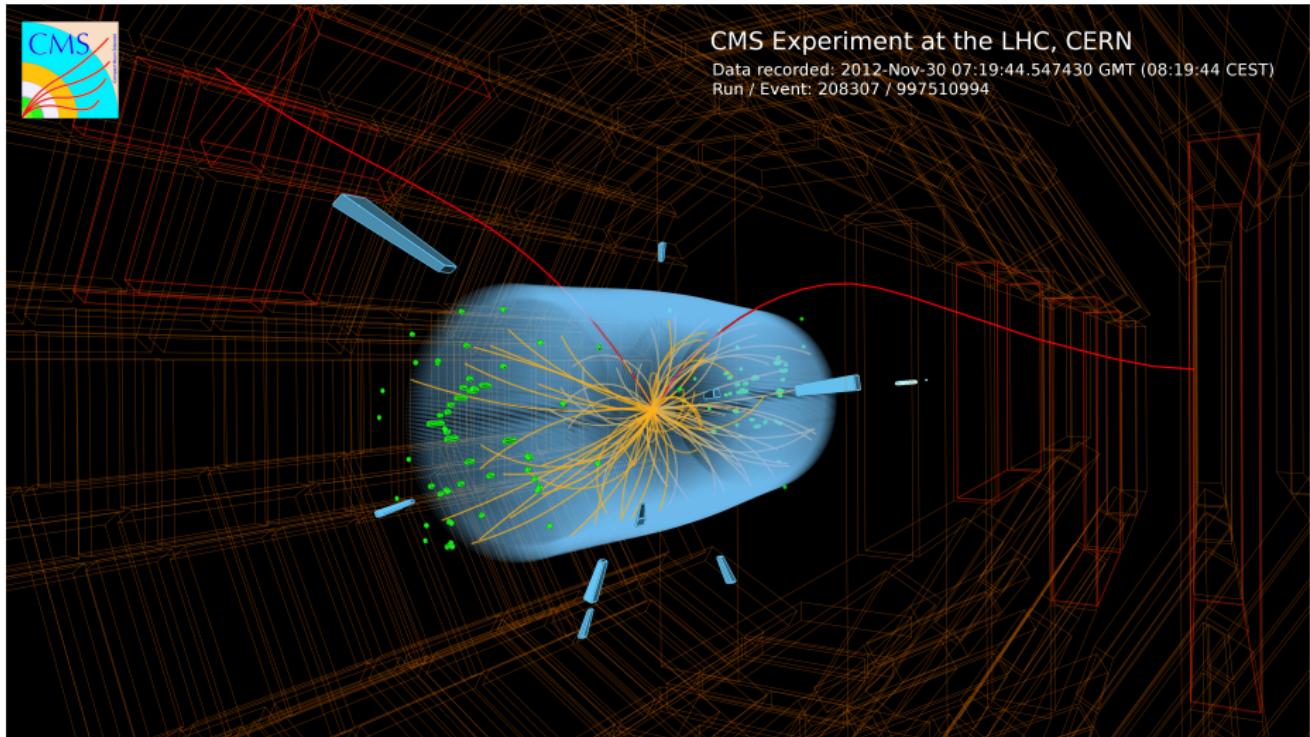
$$\mathcal{B}(B_s \rightarrow \mu\mu) = (3.0^{+0.9}_{-0.8}(\text{stat.})^{+0.6}_{-0.4}(\text{syst.})) \cdot 10^{-9} \quad S = 4.3\sigma \text{ (Exp. 4.8)}$$

$$\mathcal{B}(B_d \rightarrow \mu\mu) < 1.1 \cdot 10^{-9} \text{ (95% C.L.)}$$





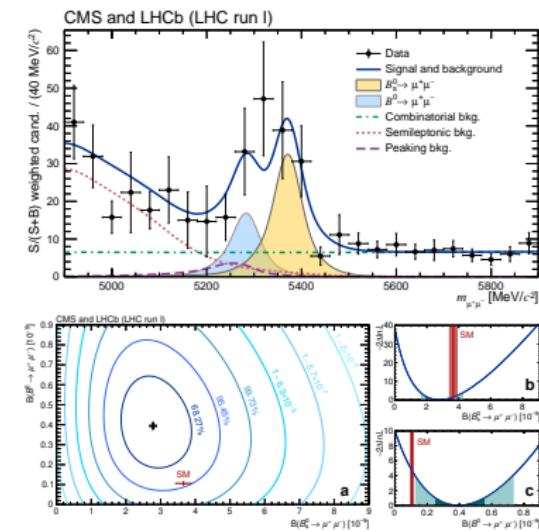
$B_{s,d} \rightarrow \mu\mu$ Event display



- Data: 25 fb^{-1} (CMS) and 3 fb^{-1} (LHCb)
- Selection: BDT with 20 categories, 12 CMS categories which depend on \sqrt{s} , detector region, and BDT ranges + 8 LHCb categories which depend on BDT ranges
- Common parameters: hadronisation fraction f_d/f_s , $B(B^\pm \rightarrow J/\psi K^\pm)$
- Improvement of the CMS analysis: $\Lambda_b \rightarrow p\mu\nu$ branching ratio updated to the last prediction

$$\mathcal{B}(B_s \rightarrow \mu\mu) = (2.8^{+0.7}_{-0.6}) \cdot 10^{-9} \quad S=6.2\sigma \quad (\text{Exp: } 7.4\sigma)$$

$$\mathcal{B}(B_d \rightarrow \mu\mu) = (3.9^{+1.6}_{-1.4}) \cdot 10^{-10} \quad S=3.0\sigma \quad (\text{Exp: } 0.8\sigma)$$

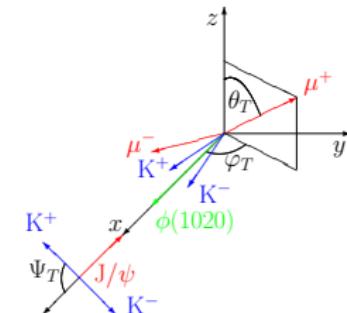
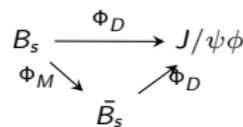




CP-violating phase ϕ_s using $B_s \rightarrow J/\psi\phi(1020)$ channel

$B_s \rightarrow J/\psi\phi(1020)$ ϕ_s motivation

- The B_s meson decays into a final state $J/\psi\phi(1020)$, which is a mixture of two CP eigenstates (odd/even);
- Interference between direct $B_s \rightarrow J/\psi\phi(1020)$ decay and decay via mixing gives rise to a CP violation phase $\phi_s \approx -2\beta_s = -0.0363^{+0.0016}_{-0.0015}$ rad (SM), where $\beta_s = \arg(-V_{ts} V_{tb}^* / V_{cs} V_{cb}^*)$;
- If the ϕ_s differs w.r.t. the SM prediction, then New Physics might contribute in the mixing box diagram.



- We measure the weak phase ϕ_s and the decay width difference $\Delta\Gamma_s$ by disentangling the two CP final states of the B_s with a tagged angular analysis;
- Opposite side lepton tagging is used to define the flavour at production time.



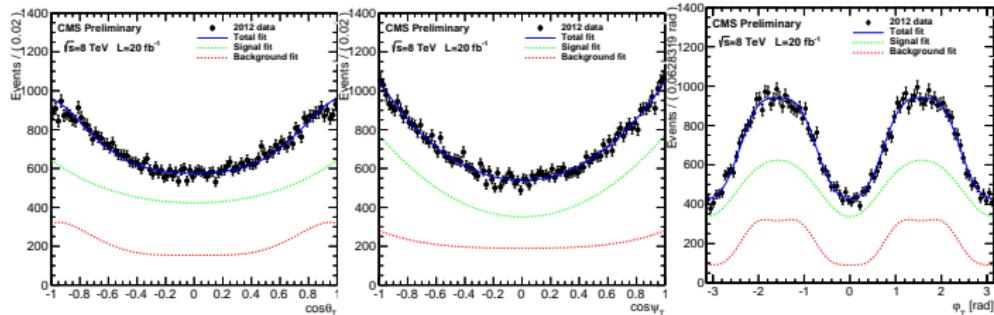
$B_s \rightarrow J/\psi \phi(1020)$ ϕ_s measurement

$$\frac{d^4\Gamma(B_s(t))}{d\Theta dt} = X(\Theta, \alpha, t) = \sum_{i=1}^{10} O_i(\alpha, t) \cdot g_i(\Theta),$$

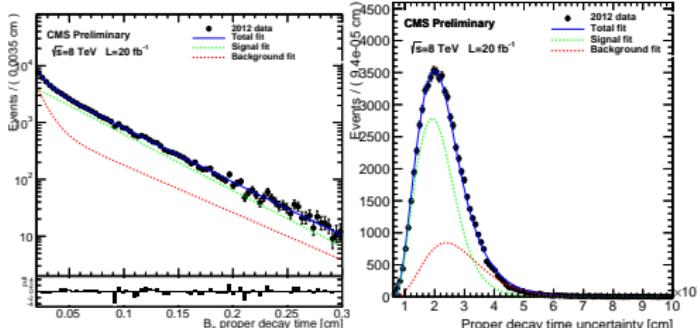
$$O_i(\alpha, t) = N_i e^{-\Gamma_s t} \left[a_i \cosh\left(\frac{1}{2}\Delta\Gamma_s t\right) + b_i \sinh\left(\frac{1}{2}\Delta\Gamma_s t\right) + c_i \cos(\Delta m_s t) + d_i \sin(\Delta m_s t) \right]$$

- b_i and d_i depend on $\sin \phi_s$ and $\cos \phi_s$
- $\Delta\Gamma_s > 0$: we use previous LHCb results
- α physics parameters ($\Delta\Gamma_s$, ϕ_s , $c\tau$, $|A_0|^2$, $|A_S|^2$, $|A_\perp^2|$, $\delta_{||}$, $\delta_{S\perp}$, δ_\perp)
- Averaged performance of the combined (μ and e) tagger: $\omega = 32.2 \pm 0.3\%$,
 $\epsilon_{tag} = 7.67 \pm 0.04\%$, $P_{tag} = 0.97 \pm 0.04\%$;

$B_s \rightarrow J/\psi\phi(1020) \phi_s$ results [CMS-PAS-BPH-13-012]



- Fit of the tagged model on the 2012 data, with a gaussian constraint of Δm_s to the PDG value $17.69 \pm 0.08 \text{ } \hbar/\text{ps}$
- Extended maximum-likelihood fit



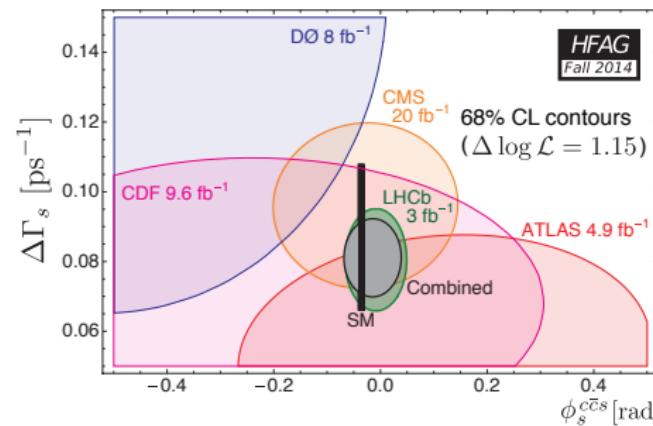
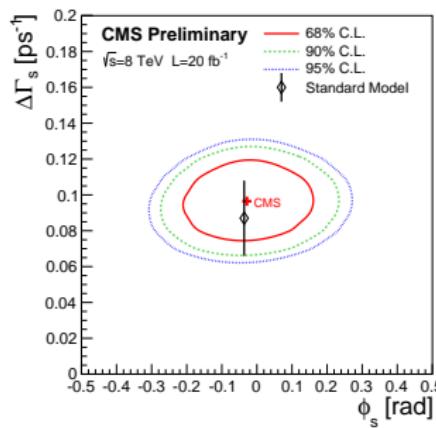
| Parameter | Fit results |
|-----------------------------------|-------------------|
| $ A_0 ^2$ | 0.511 ± 0.006 |
| $ A_S ^2$ | 0.015 ± 0.016 |
| $ A_{\perp} ^2$ | 0.242 ± 0.008 |
| $\Delta\Gamma_s [\text{ps}^{-1}]$ | 0.096 ± 0.014 |
| $\delta_{\parallel} [\text{rad}]$ | 3.48 ± 0.09 |
| $\delta_{S\perp} [\text{rad}]$ | 0.34 ± 0.24 |
| $\delta_{\perp} [\text{rad}]$ | 2.73 ± 0.36 |
| $\phi_s [\text{rad}]$ | -0.03 ± 0.11 |
| $c\tau [\mu\text{m}]$ | 447 ± 3 |

$B_s \rightarrow J/\psi\phi(1020)$ ϕ_s summary

- Analysing the 2012 CMS data (20.0 fb^{-1}), we selected 49k B_s signal events. We obtain:

$$\begin{aligned}\phi_s &= -0.03 \pm 0.11 \text{ (stat.)} \pm 0.03 \text{ (syst.) rad} \\ \Delta\Gamma_s &= 0.096 \pm 0.014 \text{ (stat.)} \pm 0.007 \text{ (syst.) ps}^{-1}\end{aligned}$$

- Contour plot (stat. only), constraining $\Delta\Gamma_s > 0$:

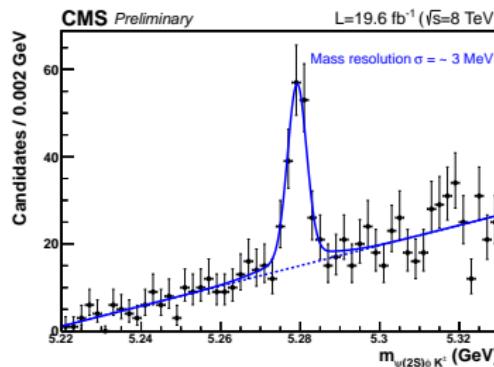




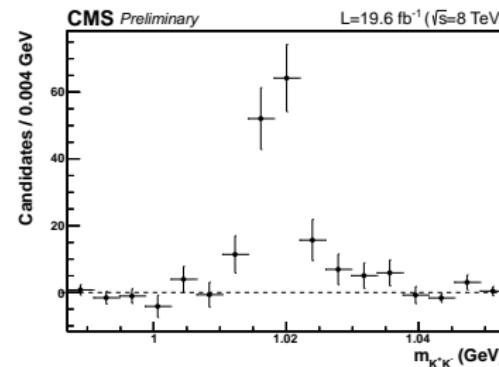
$B^+ \rightarrow \psi(2S)\phi K^+$ rare decay

$B^+ \rightarrow \psi(2S)\phi K^+$ rare decay [BPH-13-009 twiki]

Resonant peak



Background subtracted plot KK mass

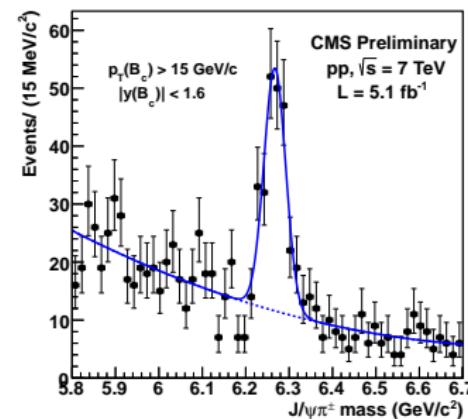
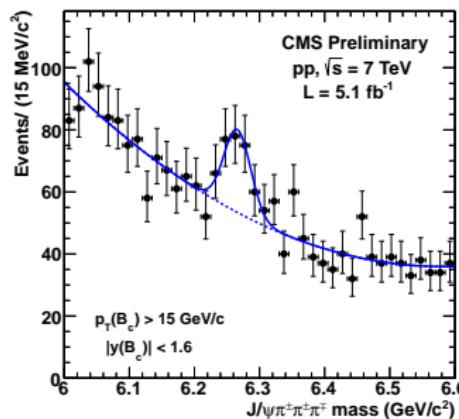


- Resonant structure due to $B^+ \rightarrow \psi(2S)\phi K^+$ decay has been found
- About 144 events identified as signal
- Next step: branching fraction measurement



B_c decays

The $B_c^\pm \rightarrow J/\psi \pi^\pm$ and $B_c^\pm \rightarrow J/\psi \pi^\pm \pi^\pm \pi^\mp$ decay modes are studied in CMS in the kinematic region where the transverse momentum of the B_c^\pm meson is greater than 15 GeV/c and within the central rapidity region $|y| < 1.6$.



$$R_{B_c} = \frac{\mathcal{B}(B_c^\pm \rightarrow J/\psi \pi^\pm \pi^\pm \pi^\mp)}{\mathcal{B}(B_c^\pm \rightarrow J/\psi \pi^\pm)} = 2.55 \pm 0.80(\text{stat}) \pm 0.33(\text{syst})^{+0.04}_{-0.01}(\tau_{B_c})$$

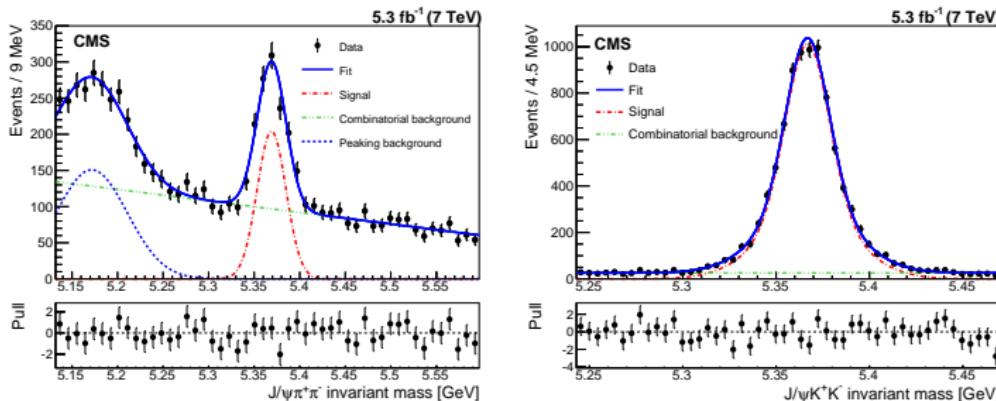
$$R_{c/u} = \frac{\sigma(B_c^\pm) \times \mathcal{B}(B_c^\pm \rightarrow J/\psi \pi^\pm)}{\sigma(B^\pm) \times \mathcal{B}(B^\pm \rightarrow J/\psi K^\pm)} = (0.48 \pm 0.05(\text{stat}) \pm 0.03(\text{syst}) \pm 0.05(\tau_{B_c})) \times 10^{-2}$$



$$\frac{\mathcal{B}(B_s \rightarrow J/\psi f_0(980))}{\mathcal{B}(B_s \rightarrow J/\psi \phi(1020))}$$

ratio

$\frac{\mathcal{B}(B_s \rightarrow J/\psi f_0(980))}{\mathcal{B}(B_s \rightarrow J/\psi \phi(1020))}$ ratio [arXiv:1501.06089 Sub. to PLB]



$$R_{f_0/\phi} = \frac{\mathcal{B}(B_s \rightarrow J/\psi f_0) \mathcal{B}(f_0 \rightarrow \pi\pi)}{\mathcal{B}(B_s \rightarrow J/\psi \phi) \mathcal{B}(\phi \rightarrow KK)} = \frac{N_{\text{obs}}^{f_0}}{N_{\text{obs}}^\phi} \times \epsilon_{\text{reco}}^{\phi/f_0} = 0.140 \pm 0.013(\text{stat}) \pm 0.018(\text{syst})$$

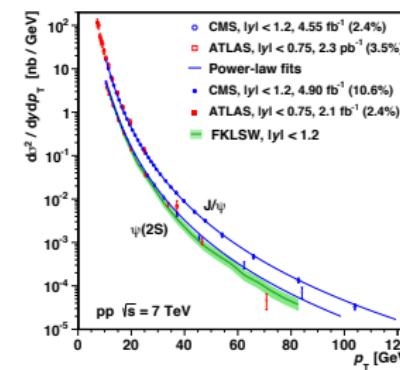
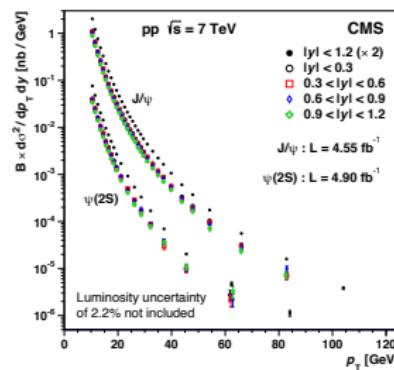
- $N_{\text{obs}}^{f_0} = 873$ and $N_{\text{obs}}^\phi = 8377$
- Consistent with predictions (PRD 79 (2009) 074024)
- Consistent with previous measurements

| Systematics | Uncert. |
|--------------------------------------|---------|
| Fit model | 2.1% |
| f_0 mass window width | 6.4% |
| MC simulation (f_0 natural width) | 8.6% |
| Decay model in MC generation | 6.2% |

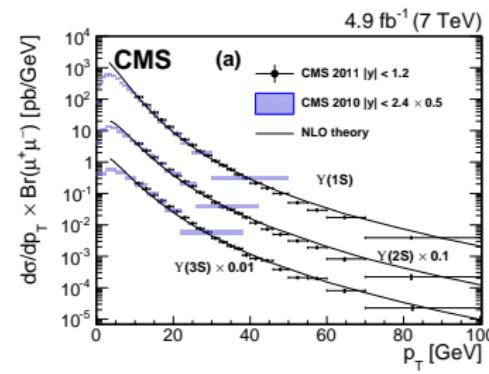
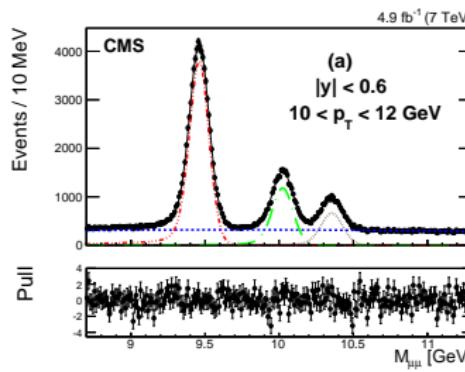


J/ψ , $\psi(2S)$, $\Upsilon(nS)$, and χ_{bn} production

- Double differential prompt J/ψ and $\psi(2S)$ cross section and rapidity integrated ($|y| < 1.2$) differential cross section (assuming unpolarised scenario)
- 2011 data with extended p_T range up to 120 GeV (J/ψ) and 100 GeV ($\psi(2S)$)
- Acceptance measured for different polarisation scenario

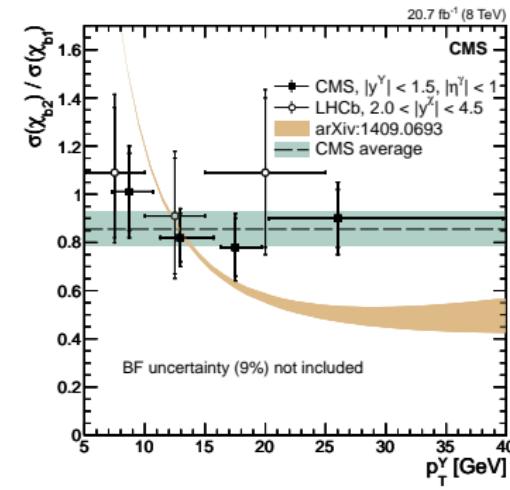
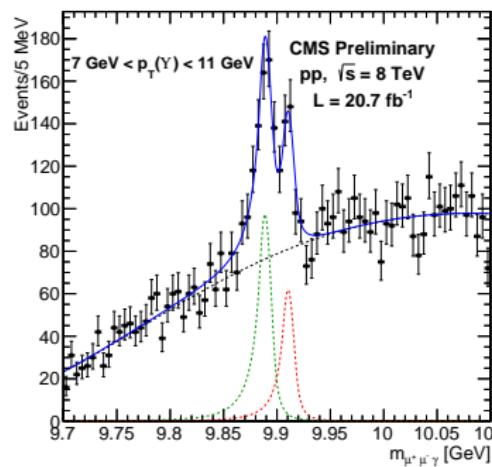


- Differential prompt $\Upsilon(1S,2S,3S) \rightarrow \mu\mu$ production as function of p_T in the rapidity range $|y| < 1.2$
- NLO calculations from [Gong, et al.] extended by the authors to cover the range $p_T < 100$ GeV
- The measurements show a transition from an exponential to a power-law behavior at $p_T \approx 20$ GeV for the three Υ states



$\chi_{b1}(1P)$ and $\chi_{b2}(1P)$ production (P-wave) [arXiv:1409.5761 Sub. to PLB]

- Use of converted photons to reconstruct $\Upsilon(1S)+\gamma$ final states in four bins
- Cross section ratio $\sigma(\chi_{b2})/\sigma(\chi_{b1})$ shows no evident p_T dependence
- Not completely in agreement with the NRQCD model predictions





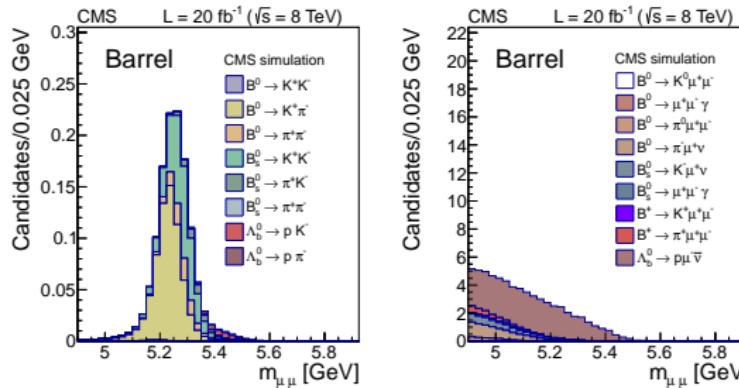
Summary

- CMS take advantage of its performance in the reconstruction of high p_T muon to measure the cross section of charmonion and bottomium states at unprecedented p_T range
- Good muon reconstruction permits CMS to be competitive in the search for $B_s \rightarrow \mu\mu$ first observation
- Rare decays of B_c and B^+ mesons are reconstructed and can be used in next analyses (e.g. cross section and lifetime measurement)
- CP-violating phase ϕ_s measurement can take advantage of $B_s \rightarrow J/\psi f_0(980)$ reconstruction



Additional slides

$B_{s,d} \rightarrow \mu\mu$ rare backgrounds



$$\mathcal{B}(B_s \rightarrow \mu\mu) = \frac{N_{B_s}}{N_{\text{norm}}} \frac{f_d}{f_u} \frac{\epsilon_{\text{norm}}}{\epsilon_{B_s}} \mathcal{B}_{\text{norm}} \quad (1)$$



$\frac{\mathcal{B}(B_s \rightarrow J/\psi f_0(980))}{\mathcal{B}(B_s \rightarrow J/\psi \phi(1020))}$ ratio non-resonant BG

- $f_0 \rightarrow \pi\pi$ invariant mass restricted to 50 MeV
- Assumption that $\pi\pi$ non-resonant component negligible component
- Assumption confirmed by PRD 89 029006
- Dedicated systematics:
 - varying the assumed width of the f_0 (50 MeV) in the MC simulation by $\pm 20\%$, which is within the 90% CL of our measured f_0 mass model
 - varying the $f_0(890)$ mass window up to ± 100 MeV to account for possible backgrounds (basically $f_0(1500)$ and $f_2(1270)$)

$B_s \rightarrow J/\psi \phi(1020) \phi_s$ signal model

We use the same notations as LHCb [arXiv:1304.2600]:

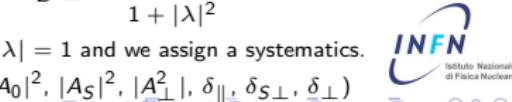
$$\frac{d^4\Gamma(B_s(t))}{d\Theta dt} = X(\Theta, \alpha, t) = \sum_{i=1}^{10} O_i(\alpha, t).g_i(\Theta),$$

$$O_i(\alpha, t) = N_i e^{-\Gamma_s t} \left[a_i \cosh\left(\frac{1}{2}\Delta\Gamma_s t\right) + b_i \sinh\left(\frac{1}{2}\Delta\Gamma_s t\right) + c_i \cos(\Delta m_s t) + d_i \sin(\Delta m_s t) \right]$$

| i | $g_i(\theta_T, \psi_T, \phi_T)$ | N_i | a_i | b_i | c_i | d_i |
|-----|---------------------------------------------------------------------------------------------------------------------------------------------------------------------|---------------------------|----------------------------------------|----------------------------------------|--------------------------------------|----------------------------------------|
| 1 | $2 \cos^2 \psi_T (1 - \sin^2 \theta_T \cos^2 \phi_T)$ | $ A_0(0) ^2$ | 1 | D | C | $-S$ |
| 2 | $\sin^2 \psi_T (1 - \sin^2 \theta_T \sin^2 \phi_T)$ | $ A_{ }(0) ^2$ | 1 | D | C | $-S$ |
| 3 | $\sin^2 \psi_T \sin^2 \theta_T$ | $ A_{\perp}(0) ^2$ | 1 | $-D$ | C | S |
| 4 | $-\sin^2 \psi_T \sin 2\theta_T \sin \phi_T$ | $ A_{ }(0)A_{\perp}(0) $ | $C \sin(\delta_{\perp} - \delta_{ })$ | $S \cos(\delta_{\perp} - \delta_{ })$ | $\sin(\delta_{\perp} - \delta_{ })$ | $D \cos(\delta_{\perp} - \delta_{ })$ |
| 5 | $\frac{1}{\sqrt{2}} \sin 2\psi_T \sin^2 \theta_T \sin 2\phi_T$ | $ A_0(0)A_{ }(0) $ | $\cos(\delta_{ } - \delta_0)$ | $D \cos(\delta_{ } - \delta_0)$ | $C \cos(\delta_{ } - \delta_0)$ | $-S \cos(\delta_{ } - \delta_0)$ |
| 6 | $\frac{1}{\sqrt{2}} \sin 2\psi_T \sin 2\theta_T \sin \phi_T$ | $ A_0(0)A_{\perp}(0) $ | $C \sin(\delta_{\perp} - \delta_0)$ | $S \cos(\delta_{\perp} - \delta_0)$ | $\sin(\delta_{\perp} - \delta_0)$ | $D \cos(\delta_{\perp} - \delta_0)$ |
| 7 | $\frac{2}{3}(1 - \sin^2 \theta_T \cos^2 \phi_T)$ | $ A_S(0) ^2$ | 1 | $-D$ | C | S |
| 8 | $\frac{1}{3}\sqrt{6} \sin \psi_T \sin^2 \theta_T \sin 2\phi_T$ | $ A_S(0)A_{ }(0) $ | $C \cos(\delta_{ } - \delta_S)$ | $S \sin(\delta_{ } - \delta_S)$ | $\cos(\delta_{ } - \delta_S)$ | $D \sin(\delta_{ } - \delta_S)$ |
| 9 | $\frac{1}{3}\sqrt{6} \sin \psi_T \sin 2\theta_T \cos \phi_T$ | $ A_S(0)A_{\perp}(0) $ | $\sin(\delta_{\perp} - \delta_S)$ | $-D \sin(\delta_{\perp} - \delta_S)$ | $C \sin(\delta_{\perp} - \delta_S)$ | $S \sin(\delta_{\perp} - \delta_S)$ |
| 10 | $\frac{4}{3}\sqrt{3} \cos \psi_T (1 - \sin^2 \theta_T \cos^2 \phi_T)$ | $ A_S(0)A_0(0) $ | $C \cos(\delta_0 - \delta_S)$ | $S \sin(\delta_0 - \delta_S)$ | $\cos(\delta_0 - \delta_S)$ | $D \sin(\delta_0 - \delta_S)$ |
| | $C = \frac{1 - \lambda ^2}{1 + \lambda ^2}, \quad S = -\frac{2 \lambda \sin \phi_S}{1 + \lambda ^2}, \quad D = -\frac{2 \lambda \cos \phi_S}{1 + \lambda ^2}$ | | | | | |

$|\lambda|$ includes possible contribution from CP violation in direct decay, we assume $|\lambda| = 1$ and we assign a systematics.

$\Delta\Gamma_s > 0$: we use previous LHCb results. α physics parameters ($\Delta\Gamma_s, \phi_s, c\tau, |A_0|^2, |A_S|^2, |A_{||}|^2, \delta_{||}, \delta_{S\perp}, \delta_{\perp}$)



$B_s \rightarrow J/\psi \phi(1020)$ ϕ_s systematics table

| Source of uncertainty | $ A_0 ^2$ | $ A_S ^2$ | $ A_{\perp} ^2$ | $\Delta\Gamma_s \text{ ps}^{-1}$ | $\delta_{\parallel} \text{ rad}$ | $\delta_{S\perp} \text{ rad}$ | $\delta_{\perp} \text{ rad}$ | $\phi_s \text{ rad}$ | $c\tau \mu\text{m}$ |
|---------------------------------|-----------|-----------|-----------------|----------------------------------|----------------------------------|-------------------------------|------------------------------|----------------------|---------------------|
| Statistical uncertainty | 0.0058 | 0.016 | 0.0077 | 0.0138 | 0.092 | 0.24 | 0.36 | 0.109 | 3.0 |
| Proper time efficiency | 0.0015 | - | 0.0023 | 0.0057 | - | - | - | 0.002 | 1.0 |
| Angular efficiency | 0.0060 | 0.008 | 0.0104 | 0.0021 | 0.674 | 0.14 | 0.66 | 0.016 | 0.8 |
| Model bias | 0.0008 | - | - | 0.0012 | 0.025 | 0.03 | - | 0.015 | 0.4 |
| Proper decay resolution | 0.0009 | - | 0.0008 | 0.0021 | 0.004 | - | 0.02 | 0.006 | 2.9 |
| BG mistag modelling | 0.0021 | - | 0.0013 | 0.0018 | 0.074 | 1.10 | 0.02 | 0.002 | 0.7 |
| Flavour tagging | - | - | - | - | - | - | 0.02 | 0.005 | - |
| PDF modelling assumpt. | 0.0016 | 0.002 | 0.0021 | 0.0021 | 0.010 | 0.03 | 0.04 | 0.006 | 0.2 |
| $ \lambda $ as a free parameter | 0.0001 | 0.005 | 0.0001 | 0.0003 | 0.002 | 0.01 | 0.03 | 0.015 | - |
| Kaon p_T re-weighting | 0.0094 | 0.020 | 0.0041 | 0.0015 | 0.085 | 0.11 | 0.02 | 0.014 | 1.1 |
| Total systematics | 0.0116 | 0.022 | 0.0117 | 0.0073 | 0.685 | 1.12 | 0.66 | 0.032 | 3.5 |

- The systematic uncertainties from the proper decay time efficiency and the kaon p_T distribution originate from residual differences between data and simulation;
 - For the proper decay time efficiency, the statistical uncertainty of the efficiency model is propagated to the results
 - For the kaon p_T , the discrepancy between the data and the simulation is taken into account as a systematic uncertainty
- The model bias systematic comes from the biases of the observables obtained with toy pseudo-experiments of the fitted model;
- We left $|\lambda|$ as a free parameter to evaluate the effect of possible other contributions in the weak phase.