Ideas and Approaches to Predict the CP Phase(s) in the Leptonic Sector

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Outline

- lepton mixing: parametrization and data
- recap: prediction of lepton mixing angles
- ideas for prediction of leptonic CP phase(s)
 - less symmetry
 - corrections
 - involve CP symmetry
- conclusions



Lepton mixing: parametrization

Parametrization

$$U_{PMNS} = \tilde{U} \operatorname{diag}(1, e^{i\alpha/2}, e^{i(\beta/2 + \delta)})$$

with

$$\tilde{U} = \begin{pmatrix} c_{12}c_{13} & s_{12}c_{13} & s_{13}e^{-i\delta} \\ -s_{12}c_{23} - c_{12}s_{23}s_{13}e^{i\delta} & c_{12}c_{23} - s_{12}s_{23}s_{13}e^{i\delta} & s_{23}c_{13} \\ s_{12}s_{23} - c_{12}c_{23}s_{13}e^{i\delta} & -c_{12}s_{23} - s_{12}c_{23}s_{13}e^{i\delta} & c_{23}c_{13} \end{pmatrix}$$

and $s_{ij} = \sin \theta_{ij}$, $c_{ij} = \cos \theta_{ij}$



Lepton mixing: data

Global fits IH [NH] (Gonzalez Garcia et al. ('14) [after NOW 2014]) best fit and 1σ error 3σ range $\sin^2 \theta_{13} = 0.0219[8]^{+0.0011[0]}_{-0.0010}$ $0.0188[6] \le \sin^2 \theta_{13} \le 0.0251[0]$ $\sin^2 \theta_{12} = 0.304^{+0.013}_{-0.012}$ $0.270 < \sin^2 \theta_{12} < 0.344$ $\sin^2 \theta_{23} = \begin{cases} 0.579^{+0.025}_{-0.037} \\ [0.452^{+0.052}_{-0.028}] \end{cases}$ $0.389[2] \le \sin^2 \theta_{23} \le 0.644[3]$ $\delta = \left(1.41[70]^{+0.35[22]}_{-0.34[9]}\right) \pi \qquad 0 \le \delta \le 2\pi$ $lpha \;,\;\; eta$ unconstrained

Well-known mixing pattern: tri-bimaximal (TB) mixing

(Harrison et al. ('02), Xing ('02))

$$U_{TB} = \begin{pmatrix} \sqrt{\frac{2}{3}} & \frac{1}{\sqrt{3}} & 0\\ -\frac{1}{\sqrt{6}} & \frac{1}{\sqrt{3}} & \frac{1}{\sqrt{2}}\\ -\frac{1}{\sqrt{6}} & \frac{1}{\sqrt{3}} & -\frac{1}{\sqrt{2}} \end{pmatrix} \quad J_{CP} = 0$$



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does not fit data well!

And since θ_{13} vanishes, the CP phase δ is unphysical.



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does not fit data well!

And since θ_{13} vanishes, the CP phase δ is unphysical. But nice thing about it: it's derived from symmetries like A_4 , S_4



 \cap

Origin of TB mixing (Lam ('07,'08))

 $U_{PMNS} = U_e^{\dagger} U_{\nu}$

[Masses do not play a role in this approach.]

Use different group for mixing with $\theta_{13} \neq 0$, $\theta_{23} \neq \frac{\pi}{4}$ (de Adelhart Toorop/Feruglio/H ('11))



(de Adelhart Toorop/Feruglio/H ('11))

$$||U_{PMNS}|| = \frac{1}{\sqrt{3}} \begin{pmatrix} \frac{1}{2}\sqrt{4+\sqrt{2}+\sqrt{6}} & 1 & \frac{1}{2}\sqrt{4-\sqrt{2}-\sqrt{6}} \\ \frac{1}{2}\sqrt{4+\sqrt{2}-\sqrt{6}} & 1 & \frac{1}{2}\sqrt{4-\sqrt{2}+\sqrt{6}} \\ \sqrt{1-\frac{1}{\sqrt{2}}} & 1 & \sqrt{1+\frac{1}{\sqrt{2}}} \end{pmatrix}$$

$$\sin^2 \theta_{12} = \frac{4}{8 + \sqrt{2} + \sqrt{6}} \approx 0.337 , \quad \sin^2 \theta_{23} = \frac{4 - \sqrt{2} + \sqrt{6}}{8 + \sqrt{2} + \sqrt{6}} \approx 0.424 ,$$
$$\sin^2 \theta_{13} = \frac{4 - \sqrt{2} - \sqrt{6}}{12} \approx 0.011 , \quad \sin \delta = 0$$

"Accidentally" δ is trivial ...



• surveys of groups G_f , G_e and G_{ν}

(H/Meroni/Vitale ('13), King/Neder/Stuart ('13), Joshipura/Patel ('13), Holthausen et al. ('12), Lam ('12), Fonseca/Grimus ('14)) have shown that all patterns have trimaximal column $(\sin^2 \theta_{12} \gtrsim 1/3)$ and Dirac phase $\delta = 0, \pi$, if mixing angles are accommodated well



Leptonic CP phases: ideas

• surveys of groups G_f , G_e and G_{ν}

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- Way-outs
 - less symmetry
 - corrections
 - involve CP as symmetry





- p. 13/??

(Ge et al. ('11), Hanlon et al. ('13))

- diagonal basis of charged leptons: $U_e = \mathbb{1}$
- neutrino sector is invariant under $Z_2^s(k)$ that is generated by

$$G_1(k) = \frac{1}{2+k^2} \begin{pmatrix} 2-k^2 & 2k & 2k \\ 2k & k^2 & -2 \\ 2k & -2 & k^2 \end{pmatrix}$$

• correlation of Dirac phase δ and mixing angles is derived

$$\cos \delta = -\frac{\left(\sin^2 \theta_{12} - \cos^2 \theta_{12} \sin^2 \theta_{13}\right) \cos 2\theta_{23}}{\sin 2\theta_{12} \sin 2\theta_{23} \sin \theta_{13}}$$



(Ge et al. ('11), Hanlon et al. ('13))

- diagonal basis of charged leptons: $U_e = \mathbb{1}$
- neutrino sector is invariant under $\overline{Z}_2^s(k)$ that is generated by

$$G_2(k) = \frac{1}{2+k^2} \begin{pmatrix} 2-k^2 & 2k & 2k \\ 2k & -2 & k^2 \\ 2k & k^2 & -2 \end{pmatrix}$$

• correlation of Dirac phase δ and mixing angles is derived

$$\cos \delta = -\frac{\left(\sin^2 \theta_{12} \sin^2 \theta_{13} - \cos^2 \theta_{12}\right) \cos 2\theta_{23}}{\sin 2\theta_{12} \sin 2\theta_{23} \sin \theta_{13}}$$



(Hanlon et al. ('13))





Use less symmetry!

(Hernandez/Smirnov ('12))

 $G_f = \mathsf{PSL}(2,7)$



 \checkmark

neutrinos

K

 \checkmark

 $G_e = Z_7$ U_e

 $G_{\nu} = Z_2$ U_{ν}

 $U_{PMNS} = U_e^{\dagger} U_{\nu}$



(Hernandez/Smirnov ('12))

one column is fixed

$$|U_{\mu 1}|^2 = \frac{1}{4\left(1 + \sin\frac{\pi}{14}\right)} \approx 0.204 \quad , \quad |U_{\tau 1}|^2 = \frac{1}{4\left(1 + \cos\frac{\pi}{7}\right)} \approx 0.132$$

so the first column of the PMNS mixing matrix is of the form

$$\begin{pmatrix} |U_{e1}| \\ |U_{\mu 1}| \\ |U_{\tau 1}| \end{pmatrix} \approx \begin{pmatrix} 0.815 \\ 0.452 \\ 0.363 \end{pmatrix}$$

• in particular notice that $|U_{e1}|^2 = \cos^2 \theta_{12} \cos^2 \theta_{13} \approx 0.664$



thick line: $\delta = 0$ dashed line: $\delta = \pi/4$ dotted line: $\delta = \pi/2$

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(Marzocca et al. ('11))

Ansatz

- $U_{PMNS} = U_e^{\dagger} \Psi U_{\nu}$
- U_{ν} is TB or bimaximal (BM)
- $U_e \neq 1$, but

$$U_e = R_{12} \left(\theta_{12}^e \right)$$

• Ψ diagonal matrix with phase



(Marzocca et al. ('11))





(Petcov ('14), Ballett et al. ('14), Girardi et al. ('14))

 $U_{PMNS} = U_e^{\dagger} \Psi U_{\nu}$ with $U_e = R_{23} \left(\theta_{23}^e \right) R_{12} \left(\theta_{12}^e \right)$ and Ψ phases

and U_{ν} is TB, BM, golden ratio (GR) mixings as well as hexagonal mixing (all have in common $\theta_{13}^{\nu} = 0$ and $\theta_{23}^{\nu} = \pi/4$)

Relation of Dirac phase δ , lepton mixing angles θ_{ij} and neutrino mixing angle θ_{12}^{ν}

$$\cos \delta = \frac{\tan \theta_{23} \sin^2 \theta_{12} + \frac{\sin^2 \theta_{13} \cos^2 \theta_{12}}{\tan \theta_{23}} - \sin^2 \theta_{12}^{\nu} \left(\tan \theta_{23} + \frac{\sin^2 \theta_{13}}{\tan \theta_{23}} \right)}{\sin 2\theta_{12} \sin \theta_{13}}$$

(Girardi et al. ('14))



(Ballett et al. ('14))



- p. 24/??

(Ballett et al. ('14))



width of bands is generated by varying θ_{23}



Sum rule

(Ballett et al. ('13))

$$a \approx a_0 + \lambda r \cos \delta$$
 with $\sin \theta_{13} = \frac{r}{\sqrt{2}}$ and $\sin \theta_{23} = \frac{1+a}{\sqrt{2}}$



(Dasgupta/Smirnov ('14))

Ansatz

$$U_{PMNS} = D(\gamma) V_{CKM}^{\dagger} U_X D(\beta)$$
 with U_X is real

and U_{PMNS} fits experimental data Note: this reminds of quark-lepton complementarity

Dirac phase
$$\sin \delta \approx -\sin \delta_q \frac{\sin \theta_{13}^q}{\sin \theta_{13}} \cos \theta_{23}$$

 $\times (1+2 \sin \theta_{13} \tan \theta_{23} \cot 2\theta_{12})$
 $\approx -\sin \delta_q \lambda^2 \text{ and } \lambda \approx 0.22$



(Dasgupta/Smirnov ('14))

Ansatz

$$U_{PMNS} = D(\gamma) V_{CKM}^{\dagger} U_X D(\beta)$$
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Dirac phase $\sin \delta \approx -\sin \delta_q \lambda^2$ and $\lambda \approx 0.22$ Majorana phases $\beta_1 \approx \sin \delta_q \sin \theta_{13}^q \sin \theta_{23} \cot \theta_{12} \approx \sin \delta_q \lambda^3$ $\beta_2 \approx -\sin \delta_q \sin \theta_{13}^q \sin \theta_{23} \tan \theta_{12} \approx -\sin \delta_q \lambda^3$

All leptonic CP phases are suppressed.

Large CP phases can arise from neutrino sector.

Involve CP as symmetry into the game!

(Feruglio/H/Ziegler ('12,'13), Holthausen et al. ('12), Chen et al. ('14), Grimus/Rebelo ('95))



Comments:

- again, masses are not fixed in this approach
- well-known example is " μ - τ reflection symmetry"

(Harrison/Scott ('02,'04), Grimus/Lavoura ('03))

$$\nu_{\mu} \rightarrow \mathsf{CP}(\nu_{\tau}) \text{ and } \nu_{\tau} \rightarrow \mathsf{CP}(\nu_{\mu})$$

consequences

$$|U_{\mu i}| = |U_{\tau i}|$$
 for $i = 1, 2, 3$



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$$\nu_{\mu} \rightarrow \mathsf{CP}(\nu_{\tau}) \text{ and } \nu_{\tau} \rightarrow \mathsf{CP}(\nu_{\mu})$$

consequences

 $\sin \theta_{23} = \cos \theta_{23}$ $\sin 2\theta_{12} \sin \theta_{13} \cos \delta = 0$



Comments:

- again, masses are not fixed in this approach
- well-known example is " μ - τ reflection symmetry" (Harrison/Scott ('02,'04), Grimus/Lavoura ('03))
- all CP phases and mixing angles are strongly correlated, since the PMNS mixing matrix only contains on free real parameter θ



Involve CP as symmetry into the game!

(Feruglio/H/Ziegler ('12,'13), Holthausen et al. ('12), Chen et al. ('14), Grimus/Rebelo ('95))



case I

(Feruglio/H/Ziegler ('12,'13))

$$U_{PMNS} = \frac{1}{\sqrt{6}} \begin{pmatrix} 2\cos\theta & \sqrt{2} & 2\sin\theta \\ -\cos\theta + i\sqrt{3}\sin\theta & \sqrt{2} & -\sin\theta - i\sqrt{3}\cos\theta \\ -\cos\theta - i\sqrt{3}\sin\theta & \sqrt{2} & -\sin\theta + i\sqrt{3}\cos\theta \end{pmatrix} K_{\nu}$$

$$\sin^2 \theta_{13} = \frac{2}{3} \sin^2 \theta , \quad \sin^2 \theta_{12} = \frac{1}{2 + \cos 2\theta} , \quad \sin^2 \theta_{23} = \frac{1}{2}$$

and

$$|\sin\delta| = 1 , \quad |J_{CP}| = \frac{|\sin 2\theta|}{6\sqrt{3}} , \quad \sin\alpha = 0 , \quad \sin\beta = 0$$



case I

(Feruglio/H/Ziegler ('12,'13))

$$U_{PMNS} = \frac{1}{\sqrt{6}} \begin{pmatrix} 2\cos\theta & \sqrt{2} & 2\sin\theta \\ -\cos\theta + i\sqrt{3}\sin\theta & \sqrt{2} & -\sin\theta - i\sqrt{3}\cos\theta \\ -\cos\theta - i\sqrt{3}\sin\theta & \sqrt{2} & -\sin\theta + i\sqrt{3}\cos\theta \end{pmatrix} K_{\nu}$$

$$\sin^2 \theta_{13} \approx 0.023$$
, $\sin^2 \theta_{12} \approx 0.341$, $\sin^2 \theta_{23} = \frac{1}{2}$

and

 $|\sin \delta| = 1$, $|J_{CP}| \approx 0.0348$, $\sin \alpha = 0$, $\sin \beta = 0$ for $\theta \approx 0.185$



(Feruglio/H/Ziegler ('12,'13))




Involve CP as symmetry into the game!

(Di lura/H/Meloni, in preparation)

 $G_f = A_5$ and CP \checkmark neutrinos charged leptons $G_e = Z_5$ $G_{\nu} = Z_2 \times \mathbf{CP}$ U_e U_{ν} \checkmark $U_{PMNS} = U_e^{\dagger} U_{\nu}$

- p. 37/??

(Di lura/H/Meloni, in preparation)

case ||
$$\tan \varphi = 1/\phi, \ \phi = \frac{1}{2} \left(1 + \sqrt{5} \right) \text{ and } \Phi = \frac{2\pi}{5}$$

$$U_{PMNS} = \frac{1}{\sqrt{2}} \begin{pmatrix} -\sqrt{2}\cos\varphi & -\sqrt{2}\sin\varphi & 0\\ -e^{-3i\Phi}\sin\varphi & e^{-3i\Phi}\cos\varphi & -e^{-7i\Phi/4}\\ -e^{-2i\Phi}\sin\varphi & e^{-2i\Phi}\cos\varphi & e^{-3i\Phi/4} \end{pmatrix} R_{13}(\theta) K_{\nu}$$

$$\sin^2 \theta_{13} = \frac{1}{10} \left(5 + \sqrt{5} \right) \, \sin^2 \theta \,, \quad \sin^2 \theta_{12} = \frac{2}{2 + (3 + \sqrt{5}) \, \cos^2 \theta}, \quad \sin^2 \theta_{23} = \frac{1}{2}$$

and

$$|\sin \delta| = 1$$
, $|J_{CP}| = \frac{1}{20\sqrt{2}}\sqrt{5+\sqrt{5}}|\sin 2\theta|$, $\sin \alpha = 0$, $\sin \beta = 0$



(Di lura/H/Meloni, in preparation)

case ||

$$U_{PMNS} = \frac{1}{\sqrt{2}} \begin{pmatrix} -\sqrt{2}\cos\varphi & -\sqrt{2}\sin\varphi & 0\\ -e^{-3i\Phi}\sin\varphi & e^{-3i\Phi}\cos\varphi & -e^{-7i\Phi/4}\\ -e^{-2i\Phi}\sin\varphi & e^{-2i\Phi}\cos\varphi & e^{-3i\Phi/4} \end{pmatrix} R_{13}(\theta) K_{\nu}$$

 $\sin^2 \theta_{13} \approx 0.0219$, $\sin^2 \theta_{12} \approx 0.283$, $\sin^2 \theta_{23} = \frac{1}{2}$

and

 $|\sin \delta| = 1$, $|J_{CP}| \approx 0.0325$, $\sin \alpha = 0$, $\sin \beta = 0$ for $\theta \approx 0.175$





(Di lura/H/Meloni, in preparation)

-n 40/??

(Everett et al. ('15))

Side remark:

- if $G_{\nu} = Z_2 \times Z_2 \times CP$, we recover GR mixing with $\theta_{13} = 0$ and unphysical δ
- if CP is not constrained to correspond to automorphism of flavor group $G_f = A_5$, Majorana phases α and β can be non-trivial (and mixing angles are accommodated well)



Involve CP as symmetry into the game! (H/Meroni/Molinaro ('14); see also Ding et al. ('14))



 $G_e = Z_3$

 U_e

neutrinos



 \checkmark

 $U_{PMNS} = U_e^{\dagger} U_{\nu}$



(H/Meroni/Molinaro ('14))

Case 1)

 $U_{PMNS} = \Omega_1 R_{13}(\theta) K_{\nu}$

with

$$\Omega_{1} = e^{i\phi_{s}} U_{TB} \begin{pmatrix} 1 & 0 & 0 \\ 0 & e^{-3i\phi_{s}} & 0 \\ 0 & 0 & -1 \end{pmatrix} \text{ and } U_{TB} = \begin{pmatrix} \sqrt{\frac{2}{3}} & \frac{1}{\sqrt{3}} & 0 \\ -\frac{1}{\sqrt{6}} & \frac{1}{\sqrt{3}} & \frac{1}{\sqrt{2}} \\ -\frac{1}{\sqrt{6}} & \frac{1}{\sqrt{3}} & -\frac{1}{\sqrt{2}} \end{pmatrix}$$

and

$$\phi_s = \frac{\pi s}{n}$$



(H/Meroni/Molinaro ('14))

Case 1)

$$U_{PMNS} = \Omega_1 R_{13}(\theta) K_{\nu}$$

This mixing matrix has a trimaximal column. Thus, $\sin^2 \theta_{12} \gtrsim 1/3$.

Results for CP phases

 $\sin \delta = 0$ and $\sin \beta = 0$ $\sin \alpha = (-1)^{k_1+1} \sin 6 \phi_s$

Very similar results are obtained, if $G_{\nu} = Z_2 \times Z_2 \times CP$ (King/Neder ('14))



(H/Meroni/Molinaro ('14))

Case 2)

$$U_{PMNS} = \Omega_2 R_{13}(\theta) K_{\nu}$$

with

$$\Omega_2 = e^{i\phi_v/6} U_{TB} R_{13} \left(-\frac{\phi_u}{2}\right) \begin{pmatrix} 1 & 0 & 0\\ 0 & e^{-i\phi_v/2} & 0\\ 0 & 0 & -i \end{pmatrix}$$

and

$$\phi_u = \frac{\pi u}{n}$$
 and $\phi_v = \frac{\pi v}{n}$



(H/Meroni/Molinaro ('14))

Case 2)

$$U_{PMNS} = \Omega_2 R_{13}(\theta) K_{\nu}$$

This mixing matrix has a trimaximal column. Thus, $\sin^2 \theta_{12} \gtrsim 1/3$.

Approximate results for CP phases

$$\sin \delta \approx \pm 1 \mp 3.3 \left(\phi_u - \bar{\phi} \right)^2$$

$$\sin \beta \approx \mp 5.6 \left(\phi_u - \bar{\phi} \right) \pm 23 \left(\phi_u - \bar{\phi} \right)^3 \quad \text{for} \quad \bar{\phi} = 0, \pm \pi$$

$$\sin \alpha \approx -\sin \phi_v \quad \text{for} \quad \bar{\phi} = 0$$



Results for Dirac phase δ and Majorana phase β (with θ expressed in u/n)



- p. 47/??

Results for Majorana phase α (with θ and u fixed with u = 1)



- p. 48/??

Numerical example: n = 8

u	$\sin^2 \theta_{13}$	$\sin^2 \theta_{12}$	$\sin^2 \theta_{23}$	$\sin \delta$	$\sin\beta$
u = 0	0.0218	0.341	0.5	1	0
u = -1	0.0254	0.342	0.387	0	0
u = 1	0.0254	0.342	0.613	0	0

values of $\sin \alpha$ for u = 0

 $\sin \alpha = 0$, $\sin \alpha = 1$ and $\sin \alpha = -1/\sqrt{2}$ values of $\sin \alpha$ for $u = \pm 1$

 $\sin \alpha \approx -0.924$ and $\sin \alpha \approx 0.383$



(H/Meroni/Molinaro ('14))

Case 3a)

$$U_{PMNS} = \Omega_3 R_{12}(\theta) K_{\nu}$$

with

$$\Omega_{3} = \begin{pmatrix} 1 & 0 & 0 \\ 0 & \omega & 0 \\ 0 & 0 & \omega^{2} \end{pmatrix} \Omega_{1} R_{13}(\phi_{m})$$

and

$$\phi_m = \frac{\pi m}{n}$$

We find $\sin^2 \theta_{13}$ and $\sin^2 \theta_{23}$ as functions of ϕ_m .



(H/Meroni/Molinaro ('14))

Case 3a)

 $U_{PMNS} = \Omega_3 R_{12}(\theta) K_{\nu}$

• m/n = 1/16 (m/n = 15/16) leads to good fit of data:

$$\sin^2 \theta_{13} \approx 0.0254$$
 and $\sin^2 \theta_{23} \approx \begin{cases} 0.613 \\ 0.387 \end{cases}$

- solar mixing angle depends on additional parameters ϕ_s and θ



Results for Dirac phase δ for m/n = 1/16



Results for Majorana phases α and β for m/n = 1/16



- p. 53/??

Numerical example: n = 16, m = 1 $\sin^2 \theta_{13} \approx 0.0254$ and $\sin^2 \theta_{23} \approx 0.613$ some viable choices of s

s	$\sin^2 heta_{12}$	$\sin\delta$	$\sin lpha$	$\sin eta$
s = 0	0.304	0	0	0
s = 1	0.304	0.458	0.939	0.662
	0.304	0.0594	-0.939	0.0383
s = 3	0.317	-0.533	0	-0.357



(H/Meroni/Molinaro ('14))

Case 3b.1)

 $U_{PMNS} = \Omega_3 R_{12}(\theta) K_{\nu} P$

- P permutation that changes ordering of columns
- for m = n/2 the first column is like in TB mixing
- thus the solar mixing angle is constrained

$$\sin^2\theta_{12} \lesssim \frac{1}{3}$$



(H/Meroni/Molinaro ('14))

Case 3b.1)

• for m/n = 1/2 also the expressions of the Majorana phases become very simple

 $\sin \alpha \propto \sin 6 \phi_s$ and $\sin \beta \propto \sin 6 \phi_s$

• only the expression for $\sin \delta$ is still complicated; however, lower limit of $|\sin \delta|$ is found

 $|\sin \delta| \gtrsim 0.71$

if mixing angles are accommodated well



Results for Dirac phase δ for m/n = 1/2 and n = 20



– n. 57

Numerical example: n = 8, m = 4 some viable choices of s

s	$\sin^2 heta_{13}$	$\sin^2 heta_{12}$	$\sin^2 heta_{23}$	$\sin\delta$	$\sin\alpha = \sin\beta$
s = 1	0.0220	0.318	0.579	0.936	$-1/\sqrt{2}$
	0.0220	0.318	0.421	-0.936	$-1/\sqrt{2}$
s = 2	0.0216	0.319	0.645	-0.739	1
s = 4	0.0220	0.318	0.5		0



Results for Dirac phase δ for m = 11 and n = 20





Results for Majorana phases α and β for n = 20 and m = 1







(Chen et al. ('14))



- p. 62/??

(Chen et al. ('14))





(Chen et al. ('14))

Remarks:

- if the two CPs are combined, this transformation acts on flavor space
 (certain correspondence to approach with G_ν = Z₂ × CP)
- PMNS mixing matrix contains one real free parameter θ
- if only one CP is preserved in neutrino sector, PMNS mixing matrix contains three real free parameters



Conclusions

- approach with flavor symmetry only leads to trivial Dirac phase
- if we assume less symmetry or corrections, values of δ different from 0 or π can be achieved
- approach with flavor and CP symmetry leads to highly constrained scenario (only one free real parameter):

all mixing parameters are strongly correlated, in particular predictions for all CP phases can be made (variations of approach are possible)



Conclusions

 phenomenology not limited to prediction of Dirac phase measured in neutrino oscillations, but Majorana phases in neutrinoless double beta decay and high energy phases relevant for leptogenesis can be also constrained

Thank you for your attention.



Back up Slides



(Ge et al. ('11))



(a) Dirac CP Phase δ_D



(Ge et al. ('11))





(Hanlon et al. ('13))



- D 70/22

Leptonic CP phases: corrections

(Marzocca et al. ('13))

Ansatz

- $U_{PMNS} = U_e^{\dagger} \Psi U_{\nu}$
- U_{ν} is TB or BM
- U_e now contains two rotations:

 $R_{12}\left(heta_{12}^{e}
ight)$ and $R_{23}\left(heta_{23}^{e}
ight)$

• Ψ diagonal matrix with phase



Leptonic CP phases: corrections

(Marzocca et al. ('13))



- p. 72/??
(Ballett et al. ('14))



width of bands is generated by varying θ_{23}



(Ballett et al. ('14))



width of bands is generated by varying θ_{23}



(Girardi et al. ('14))





(Girardi et al. ('14))





- neutrinos can be their own antiparticles
- if true, a process called $0\nu\beta\beta$ decay is allowed





- neutrinos can be their own antiparticles
- if true, a process called $0\nu\beta\beta$ decay is allowed

$$m_{ee} = \left| \cos^2 \theta_{12} \, \cos^2 \theta_{13} \, m_1 + \sin^2 \theta_{12} \, \cos^2 \theta_{13} \, e^{i\alpha} \, m_2 + \sin^2 \theta_{13} \, e^{i\beta} \, m_3 \right|$$

using the experimentally preferred 3 σ ranges of $\sin^2 \theta_{13}$, $\sin^2 \theta_{12}$ and of the mass splittings and using both possible mass orderings (normal and inverted) and

varying the unknown Majorana phases α and β and the lightest neutrino mass m_0 we get ...



- neutrinos can be their own antiparticles
- if true, a process called $0\nu\beta\beta$ decay is allowed



(H/Molinaro, in preparation)

Case 2) with n = 8, u = 0 and normal ordering



– p. 80/??

(H/Molinaro, in preparation)

Case 2) with n = 8, u = 0 and inverted ordering



- p. 81/??

(H/Molinaro, in preparation)

Case 2) with n = 8, u = 1 and normal ordering



(H/Molinaro, in preparation)

Case 2) with n = 8, u = 1 and inverted ordering



(H/Molinaro, in preparation)

Case 3a) with n = 16, m = 1, s = 0 and normal ordering



- p. 84/??

(H/Molinaro, in preparation)

Case 3a) with n = 16, m = 1, s = 0 and inverted ordering



(H/Molinaro, in preparation)

Case 3a) with n = 16, m = 1, s = 1 and normal ordering



- p. 86/??

(H/Molinaro, in preparation)

Case 3a) with n = 16, m = 1, s = 1 and inverted ordering



- p. 87/??

(H/Molinaro, in preparation)

Case 3a) with n = 16, m = 1, s = 3 and normal ordering



– p. 88/??

(H/Molinaro, in preparation)

Case 3a) with n = 16, m = 1, s = 3 and inverted ordering

