## Is the composite fermion a Dirac particle?

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• Fractional quantum Hall effect

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- Composite fermion orthodoxy

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- Fractional quantum Hall effect
- Composite fermion orthodoxy
- The old puzzle of particle-hole symmetry
- The solution to the puzzle

## Hall conductivity/resistivity



$$j_i = \sigma_{ij} E_j$$

$$E_i = \rho_{ij} j_j$$

i, j = x, y

#### Fractional QH effect



Landau levels of 2D electron in B field



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Filling factor  $\nu = \frac{n}{B/2\pi}$ 

Landau levels of 2D electron in B field



#### Energy scales



Interesting limit:  $eB/mc >> \Delta (m \rightarrow 0)$ only lowest Landau level (LLL) states survives

No small parameter

(Wilczek 1982)



(Wilczek 1982)



(Wilczek 1982)



(Wilczek 1982)



(Wilczek 1982)

Attaching flux changes statistics



 $(-1)\exp(i\pi) = (+1)$ 

(Wilczek 1982)



(Wilczek 1982)



(Wilczek 1982)

Attaching flux changes statistics



 $(-1)\exp(2i\pi) = (-1)$ 

(Wilczek 1982)





per F



per F



average per F



FQHE for = IQHE for original fermions composite fermions



original fermions

composite fermions (n=2)

### Mathematically

Lopez, Fradkin Halperin, Lee, Read

$$\mathcal{L} = i\psi^{\dagger}(\partial_{0} - iA_{0} + ia_{0})\psi - \frac{1}{2m}|(\partial_{i} - iA_{i} + ia_{i})\psi|^{2} + \frac{1}{4\pi p}\epsilon^{\mu\nu\lambda}a_{\mu}\partial_{\nu}a_{\lambda} + \cdots$$

$$\uparrow$$

$$\# \text{ of attached}$$
flux quanta

 $\mathbf{\nabla} \times \mathbf{a} = 2\pi p \, \psi^{\dagger} \psi$ 

## Comments on flux attachment

- No small expansion parameter: p~1
- Difficulty with energy scales, especially in the limit  $m \rightarrow 0$
- Nevertheless, explains a number of facts
  - Jain sequences
  - Gapless v=1/2 state

## Jain's sequences



### v=1/2 state

- After flux attachment: average magnetic field=0
- Ground state: a gapless Fermi liquid
  - described by Halperin-Lee-Read (HLR) field theory
- Substantial experimental evidence

#### Particle-hole symmetry Girvin 1984



Formalized as an anti-unitary transformation

exact symmetry on the LLL, when mixing of higher LLs negligible

• Comparing states on two Jain sequences

$$v = 1/3$$
  $v = 2/3$ 

• Comparing states on two Jain sequences

$$v = 2/5$$
  $v = 3/5$ 

• Comparing states on two Jain sequences

$$v = 3/7$$
  $v = 4/7$ 

• Comparing states on two Jain sequences

CF picture does not respect PH symmetry

## PH symmetry of v=1/2 Fermi liquid

How can the inside of a circle be equivalent to the outside?





### PH symmetric CFs?

$$\nu = \frac{n}{2n+1} \longrightarrow \nu_{\rm CF} = n$$

$$\nu = \frac{n+1}{2n+1} \longrightarrow \nu_{\rm CF} = n+1$$

#### PH symmetric CFs?







Can the filling factor of an IQH state be half-integer?

## Relativistic model with FQHE



$$S = \int d^3x \, i\bar{\psi}\gamma^{\mu}(\partial_{\mu} - iA_{\mu})\psi - \frac{1}{4e^2} \int d^4x \, F_{\mu\nu}^2$$

(Graphene: 4 types of fermions, electrons and photons have different speeds)

## Ground state in finite magnetic field



Ground state is not determined without interaction When e<sup>2</sup><<I: exactly the same FQH problem

 $\mu$ =0: half-filled Landau level

#### The offset of 1/2

$$\Delta \nu = 1 \bigwedge^{\nu} \frac{\nu = 1/2}{\nu = -1/2}$$

$$\nu = \nu_{\rm NR} - \frac{1}{2}$$

## Offset of 1/2 in graphene





**Figure 4** | **QHE for massless Dirac fermions.** Hall conductivity  $\sigma_{xy}$  and longitudinal resistivity  $\rho_{xx}$  of graphene as a function of their concentration at B = 14 T and T = 4 K.  $\sigma_{xy} \equiv (4e^2/h)\nu$  is calculated from the measured

# Discrete symmetries of (2+1)D Dirac fermion



CP and PT are symmetries at  $\mu$ =0

PH symmetry = CT = (CP)(PT) is a symmetry before projection to LLL

## PH symmetry for Dirac fermion

Jain sequences

$$\nu_{\rm NR} = \frac{n}{2n+1} \longrightarrow \nu = -\frac{1}{2(2n+1)}$$
$$\nu_{\rm NR} = \frac{n+1}{2n+1} \longrightarrow \nu = \frac{1}{2(2n+1)}$$

$$\nu_{\rm CF} = n + \frac{1}{2}$$

$$2\nu = \frac{1}{2\nu_{\rm CF}}~?$$

#### Particle-vortex duality?

original fermion

magnetic field

density

composite fermion

density

magnetic field

#### Particle-vortex duality?



This suggests the following effective action for CFs:

$$S = \int d^3x \left[ i\bar{\psi}\gamma^{\mu}(\partial_{\mu} + 2ia_{\mu})\psi + \frac{1}{2\pi}\epsilon^{\mu\nu\lambda}A_{\mu}\partial_{\nu}a_{\lambda} + \cdots \right]$$
$$-\frac{1}{4e^2}\int d^4x F^2_{\mu\nu}$$

$$S = \int d^3x \left[ i\bar{\psi}\gamma^{\mu}(\partial_{\mu} + 2ia_{\mu})\psi + \frac{1}{2\pi}\epsilon^{\mu\nu\lambda}A_{\mu}\partial_{\nu}a_{\lambda} + \cdots \right]$$

$$j^{\mu} = \frac{\delta S}{\delta A_{\mu}} = \frac{1}{2\pi} \epsilon^{\mu\nu\lambda} \partial_{\nu} a_{\lambda}$$

$$\frac{\delta S}{\delta a_0} = 0 \longrightarrow \langle \psi \bar{\gamma}^0 \psi \rangle = \frac{B}{4\pi}$$

### Dirac composite fermions

- No Chern-Simons interaction ada
  - ada would break CP and CT
- conflict with flux attachment idea?
- composite fermions have Berry phase π around Fermi surface

#### A few related works

- "Composite Dirac liquid" on surface of topological insulators Mross, Essin, Alicea 2014
- Emergent Fermi surface from mirror symmetry Hook, Kachru, Torroba, Wang 2014

## Consequences

• Exact particle hole symmetry in linear respose

• at 
$$\nu = \frac{1}{2}$$
,  $\sigma_{xy} = \frac{1}{2}$  exactly

• New particle-hole symmetric gapped nonabelian state at v=1/2:

 $\left\langle \epsilon^{\alpha\beta}\psi_{\alpha}\psi_{\beta}\right\rangle \neq 0$ 

## CS theory as the NR limit



When CP is broken, CF has mass

In the NR limit: NR action for CF

Integrating out Dirac sea: Chern-Simons interaction between CF

## Conclusion and open questions

- PH symmetry: a challenge for CF picture
- Proposal: Dirac CF with gauge, non-CS interaction
- a CP-invariant Pfaffian-like state
- Alternative: PH symmetry spontaneously broken Barkeshli Mulligan Fisher 2015
- Open questions:
  - derivation of the effective theory
  - experimental measurement of the Berry phase