

A menagerie of non-relativistic physics

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Major lesson from the 2000s

Deduce ~universal, non-perturbative features
in ~~SUSY~~ QFT via

1. Spacetime symmetries (AdS/CFT)
2. Anomalies (symmetry protected phases)

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hydrodynamics, low-energy $T > 0$ transport

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Deduce ~universal, non-perturbative features
in ~~SUSY~~ QFT via

1. Spacetime symmetries (AdS/CFT)

2. Anomalies (symmetry protected phases)

existence of gapless edge states,
symmetry-protected response (e.g. QHE, TI)

Anomalies and QHE

Classic example:

- in 3d gapped phase of relativistic QFT, low-energy topological EFT

$$S_{eff} = \frac{k}{4\pi} \int A \wedge dA + c \int \text{tr} \left(\Gamma \wedge d\Gamma + \frac{2}{3} \Gamma^3 \right) + \dots$$

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1. symmetry-protected bulk response
2. characterize anomalies of edge QFT

Anomalies and NR QHE

What one actually finds in (F)QHE is a 3d topological EFT for gapped phases of NR QFT:

$$S_{eff} = \frac{k}{4\pi} \int A \wedge dA + \tilde{\eta} \int A \wedge d\omega + c \int \omega \wedge d\omega + \dots$$

Basic questions:

1. What precisely is ω ?
2. Do $(\tilde{\eta}, c)$ correspond to genuine anomalies?
3. What about QHE in Galilean-invariant theories?

Thermal transport

Lesson from applied AdS/CFT, fluid/gravity, hydro.. :

Spacetime symmetries imply non-perturbative **interrelations** between charge, thermal transport

Manifested in **hydrodynamic** effective description

What about in NR systems? Like **anomalous Hall** effect?

(2+1d systems with ~~parity~~,
theory of transport developed in 1970s)

The plan

- i. Motivation
- ii. Geometry for NR QFT**
- iii. Transport
- iv. Warped CFT

Galilean QFT

Galilean theories possess a boost symmetry and U(1) charge: $[P_i, K_j] = -i\delta_{ij}M$

simplest example:

$$S_{free} = \int d^d x \left\{ \frac{i}{2} \Psi^\dagger \overleftrightarrow{\partial}_0 \Psi - \frac{\delta^{ij}}{2m} \partial_i \Psi^\dagger \partial_j \Psi \right\}$$

Geometry and symmetries

After some work, can deduce that Galilean theories naturally couple to [Son]³, [Geracie, Son, Wu²], [KJ], [others]

$$(n_\mu, h_{\nu\rho}, A_\sigma) \longleftrightarrow (v^\mu, h^{\nu\rho}, A_\sigma)$$

$h_{\mu\nu}$ gives spatial metric $\gamma_{\mu\nu} = n_\mu n_\nu + h_{\mu\nu}$

$$S_{free} = \int d^d x \sqrt{\gamma} \left\{ \frac{iv^\mu}{2} \Psi^\dagger \overleftrightarrow{D}_\mu \Psi - \frac{h^{\mu\nu}}{2m} D_\mu \Psi^\dagger D_\nu \Psi \right\}$$

Geometry and symmetries

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One ought to impose invariance under

1. U(1) gauge transformations

2. Reparameterizations

3. “Milne boosts,” which act as

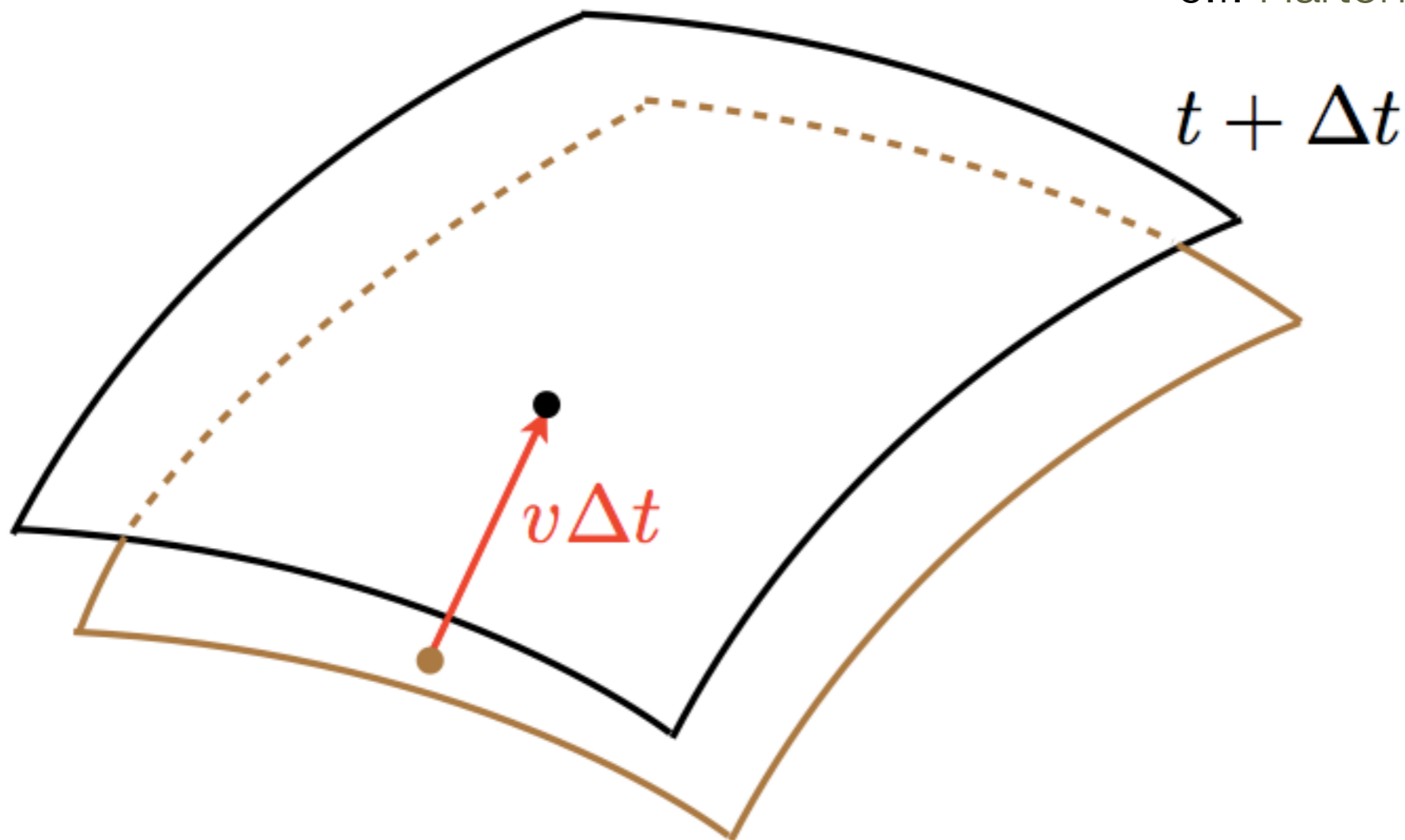
$$\begin{cases} v^\mu \rightarrow v^\mu + \psi^\mu \\ A_\mu \rightarrow A_\mu + \psi_\mu - \frac{1}{2} n_\mu \psi^2 \end{cases}$$

$(v^\mu \psi_\mu = 0)$

Newton-Cartan geometry

We've unwittingly generated (one version of) Newton-Cartan (NC) geometry

c.f. Hartong's, Obers' talks



Cross-checks

1. For NR theories obtained either from large c limits [KJ, Karch] or from DLCQ [Duval, et al] [Christensen, et al], NR QFT couples to NC geometry in a **Milne-invariant** way [KJ]

$$G = 2n_{\mu} dx^{\mu} (dx^{-} + A) + h_{\mu\nu} dx^{\mu} dx^{\nu}$$

2. So this background geometry appears on boundary of asymptotically Schrodinger spacetimes [Christensen, et al]
3. (Others)

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Hydrodynamics - I

- Hydrodynamics is near-universal **effective theory** at $T > 0$
- Describes evolution of $\langle T^{\mu\nu} \rangle, \langle J^\mu \rangle$ near equilibrium
- ~ Current algebra
- Intimately tied to spacetime **symmetries**
- AdS/CFT-inspired mini-revolution (many new **transport coefficients** which do not appear in the textbook treatment)

Hydrodynamics - II

Hydro in one slide:

1. Promote thermodynamic quantities (T, μ, u^μ) to classical dofs (hydro variables)

2. Constitutive relations for $(T^{\mu\nu}, J^\mu)$, e.g.

$$J^\mu = \rho u^\mu + \mathcal{O}(\partial)$$

3. Impose Ward identities as eoms
$$\begin{cases} D_\nu T^{\mu\nu} = F^\mu{}_\nu J^\nu, \\ D_\mu J^\mu = 0, \end{cases}$$

4. Demand a local Second law:

$$\exists S^\mu \text{ s.t. } D_\mu S^\mu \geq 0$$

Hydrodynamics - II

Hydro in one slide:

1. Promote thermodynamic quantities (T, μ, u^μ) to classical **dofs** (hydro variables)

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4. Demand a local Second law:

$$\exists S^\mu \text{ s.t. } D_\mu S^\mu \geq 0 \longrightarrow$$


~understood from first principles!

Hydrodynamics - III

What do we have for NR hydro? [KJ]

1. Hydro variables: (T, μ, u^μ) $(u^\mu n_\mu = 1)$
2. Operators conjugate to background spacetime:

$$\langle J^\mu \rangle, \langle \mathcal{P}_\mu \rangle = h_{\mu\nu} \langle J^\nu \rangle, \langle T_{\mu\nu} \rangle, \langle \mathcal{E}^\mu \rangle$$


$$\langle \mathcal{T}^{\mu\nu} \rangle$$

$$\langle \tilde{\mathcal{E}}^\mu \rangle = \langle \mathcal{E}^\mu \rangle - \left(u_\nu - \frac{1}{2} n_\nu u^2 \right) \langle \mathcal{T}^{\mu\nu} \rangle$$

Hydrodynamics - III

What do we have for NR hydro? [KJ]

3. Ward identities:
[Geracie, Son, Wu²], [KJ]

$$\begin{cases} \left(\tilde{D}_\mu - 2E_\mu^n \right) \tilde{\mathcal{E}}^\mu = \tilde{D}_\mu u_\nu \mathcal{T}^{\mu\nu} \\ \left(\tilde{D}_\nu - E_\nu^n \right) \mathcal{T}^{\mu\nu} = (F^n)^\mu{}_\nu \tilde{\mathcal{E}}^\nu \end{cases}$$

4. Impose local **Second law**

P-violating hydrodynamics in 2+1d

Consider fluid mechanics of $T > 0$ system with ~~parity~~

Relevant for systems which display **anomalous Hall effect**

A careful entropy analysis shows that there are **6** new first-order, **dissipationless** response coefficients (4 independent)

$$\tilde{\eta}, \tilde{\kappa}$$

$$\mathcal{M}, \mathcal{M}_n$$

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Hall viscosity
anomalous thermal Hall conductivity

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magnetic susceptibilities

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Warped CFTs and holography

- Many near-extremal BHs have BTZ near-horizon geometry
- However, most (like Kerr) have a “warped BTZ” near-horizon

$$g = R^2 \left(-\cosh^2 r d\tau^2 + dr^2 + c^2 (d\phi + \sinh \sigma d\tau)^2 \right)$$

Is there a dual CFT? What is it?

Warped CFTs and holography

$$g = R^2 \left(-\cosh^2 r d\tau^2 + dr^2 + c^2 (d\phi + \sinh \sigma d\tau)^2 \right)$$

1. One set of bc (e.g. [Compere, Detournay]) implies
Virasoro \times U(1) **Kac-Moody** symmetries (“WCFT”)
2. Another [Guica] implies dual is **ordinary CFT**

WCFTs from field theory - I

Suppose you have a 2d theory invariant under:

$$x^{\pm} \rightarrow x^{\pm} + c^{\pm}, \quad x^{-} \rightarrow \lambda x^{-} \quad [\text{Hofman, Strominger}]$$

Unitarity, locality enhance these to affine algebras

Two minimal cases:

1. Lorentz-invariance : get ordinary 2d CFT
2. R-moving translation enhances to Virasoro,
L-moving translation to Kac-Moody

→ WCFT!

WCFTs from field theory - II

$$x^- = x^-(X^-),$$

$$x^+ = X^+ + f(X^-)$$

Virasoro:

$$T(x)T(0) \sim \frac{c}{2(x^-)^4}$$

Kac-Moody:

$$P(x)P(0) \sim \frac{k}{(x^-)^2}$$

(R-moving global symmetries also get enhanced to KM)

WCFTs from field theory - III

Can use symmetries to deduce modular properties of torus partition function, “Cardy formula”

[Detournay, Hartman, Hofman]

Correctly reproduces **entropy** of WAdS3 BHs

Q: Should we really care about WCFT?

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One route:

- deduce *anomalies* of WCFT
- can these be matched to anything?
(putative holographic duals, edge theory of QHE, &c)

WCFTs are Galilean CFTs

[Hofman, Rollier]

Show that WCFTs naturally couple to a version of NC geometry which they dub “warped geometry”

Easy to then show that WCFTs are (massless) Galilean CFTs with $z = \infty$ [KJ, WIP]

$$h_{\mu\nu} = h_\mu h_\nu \quad (d = 2)$$

Free scalar WCFT

With this in hand we can easily write free scalar WCFTs:

$$S = \frac{1}{2} \int d^2x \sqrt{\gamma} \left\{ (w^\mu \partial_\mu \varphi)^2 + \left(\frac{N^2}{2} - w^\mu \partial_\mu N \right) \varphi^2 + m^2 \varphi^2 \right\}$$

$$(N = \varepsilon^{\mu\nu} \partial_\mu n_\nu)$$

$$\varphi \rightarrow e^{-\Omega/2} \varphi$$

“conformal mass”

Kac-Moody: $k \sim m$

Anomalies and central charges

WZ consistency \Rightarrow four types of anomalies:

1. Pure boost: $\delta_\psi W \propto k \int d^2x \sqrt{\gamma} \psi$
2. Mixed boost/Weyl: $\delta_\chi W \propto c \int d^2x \sqrt{\gamma} \left\{ -\dot{\Omega} N + \frac{\psi}{2} N^2 \right\}$
3. Gravitational: $\delta_\xi W = \tilde{c} \int \partial_\mu \xi^\nu d\Gamma^\mu{}_\nu + \dots$
4. Flavor: $\delta_\Lambda W \propto \tilde{k} \int \Lambda \cdot dA$

Anomalies and central charges

WZ consistency => four types of anomalies:

1. Pure boost:

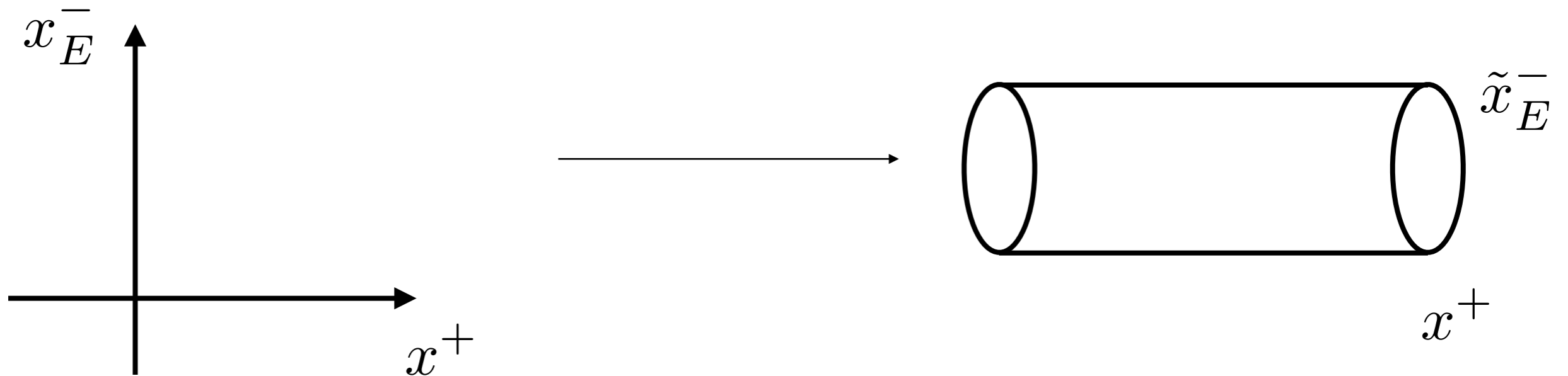
$$\delta_\psi W \propto k \int d^2x \sqrt{\gamma} \psi$$

I don't see how this can be matched to holography, or anything else for that matter

Challenge: Match (?) boost anomaly to WAdS3

Thermodynamics

Just like in 2d CFT, there is a *Weyl/boost* which takes the Euclidean plane to a *thermal cylinder*



Get thermodynamics from “Schwarzian”

(can also get entire hydrostatic partition function)

Conclusions

1. Galilean-invariant theories couple to **NC geometry**
2. Crucial ingredient: **Milne boosts**
3. **NR hydro** modernized; ~ first-principles approach to AHE
4. Warped CFTs are Galilean CFTs with $z = \infty$,
puzzles with anomaly matching to WAdS3
5. [WIP] Anomaly inflow and QHE

Thank you!

Extra slides

Another proposal

[Geracie, Prabhu, Roberts] Give alternative: “Bargmann structure”

Most natural in terms of frame, spin connection

$$g = \begin{pmatrix} 1 & 0 \\ K^i & R^i_j \end{pmatrix}$$

$$\omega = \begin{pmatrix} 0 & 0 \\ \omega^i_0 & \omega^i_j \end{pmatrix}$$

$$\begin{cases} f^A_\mu = (n_\mu, e^i_\mu) \\ F^\mu_A = (v^\mu, E^\mu_i) \end{cases}$$

reduces to above upon restricting

$$\omega^i_0 \wedge e_i \sim F$$

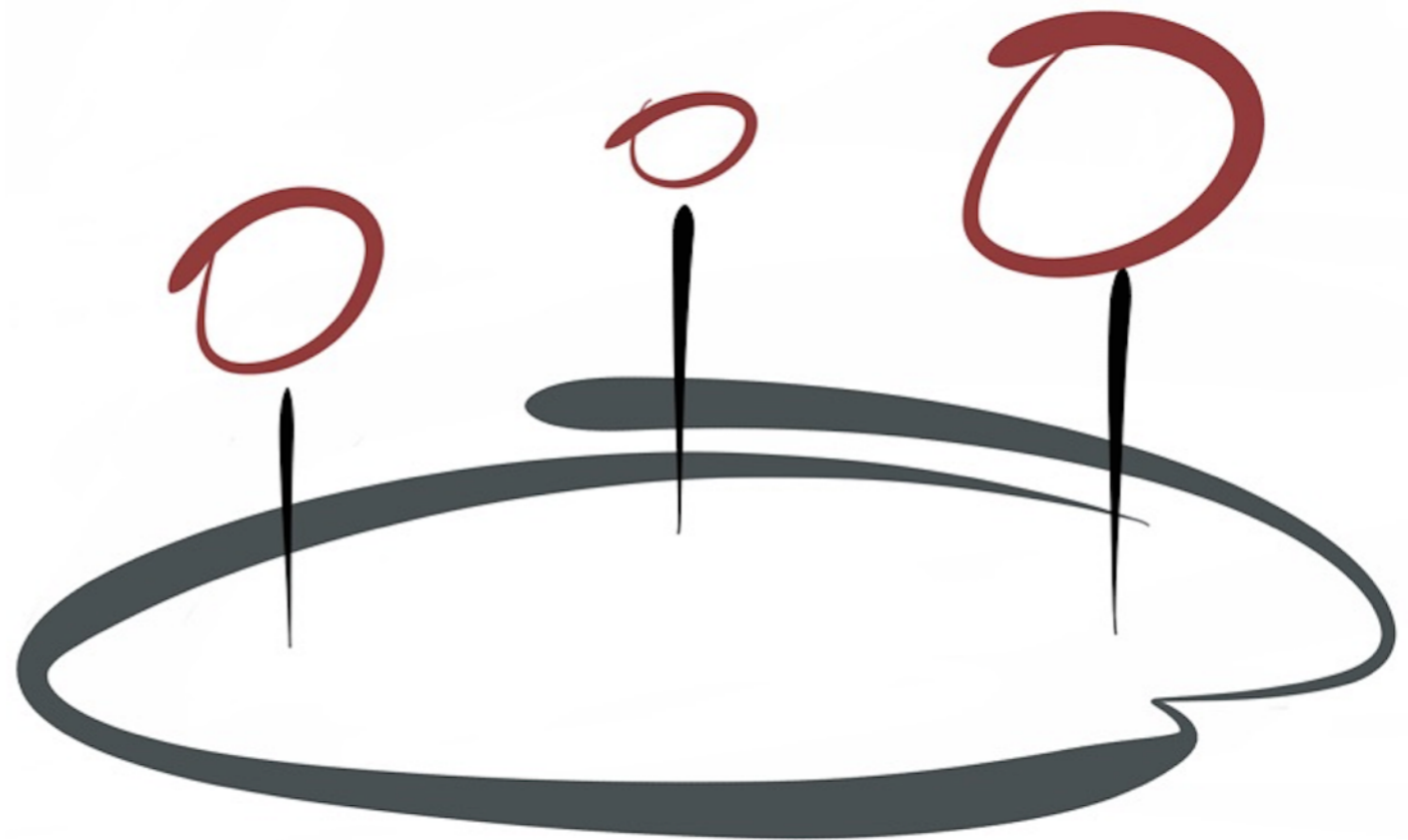
Hydrostatics

Just as in the relativistic case [KJ, et al] [Banerjee, et al]

Intimate relation between local Second law and
hydrostatic equilibrium

Finite correlation length:
effective Euclidean theory
gapped, leads to simple

$$W = \int d^d x \sqrt{\gamma} P(T, \mu) + \mathcal{O}(\partial)$$



WCFTs are Galilean CFTs

Easy to then show that WCFTs are (massless) Galilean CFTs with $z = \infty$

How do we see this?

Consider a theory of $m=0$ Galilean fields, coupled to $(n_\mu, h_{\mu\nu})$

$$\begin{array}{l} h_{\mu\nu} = h_\mu h_\nu \\ h^{\mu\nu} = w^\mu w^\nu \end{array} \longrightarrow \text{Under boosts, } \begin{cases} n_\mu \rightarrow n_\mu \\ h_\mu \rightarrow h_\mu - \psi n_\mu \end{cases}$$

$$z = \infty \longrightarrow \text{Under Weyl, } \begin{cases} n_\mu \rightarrow e^\Omega n_\mu \\ h_\mu \rightarrow h_\mu \end{cases}$$

WCFTs are Galilean CFTs

Global symmetries in flat space fix $n = dx^-$, $h = dx^+$

These are exactly

$$x^- = x^-(X^-),$$

$$x^+ = X^+ + f(X^-)$$