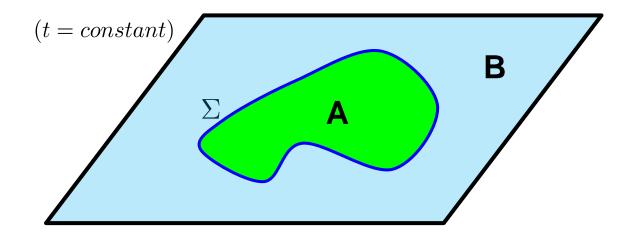
Gauge/Gravity Duality 2015

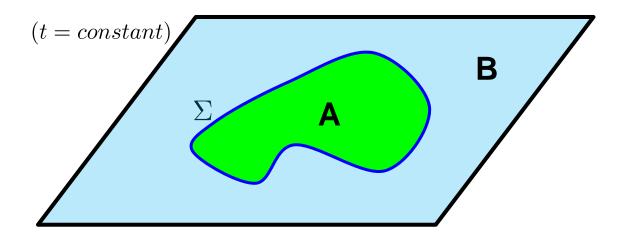
Mutual Information & the F-Theorem

with Casini, Huerta & Yale

- general tool; divide quantum system into two parts and use entropy as measure of correlations between subsystems
- in QFT, typically introduce a (smooth) boundary or entangling surface $\Sigma\,$ which divides the space into two separate regions
- integrate out degrees of freedom in "outside" region
- remaining dof are described by a density matrix ρ_A
 - \longrightarrow calculate von Neumann entropy: $S_{EE} = -Tr \left[\rho_A \log \rho_A \right]$

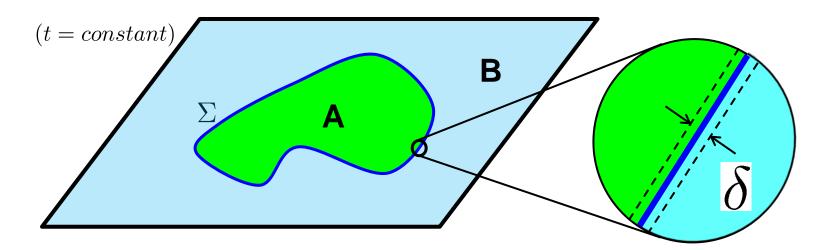


- remaining dof are described by a density matrix ρ_A
 - \longrightarrow calculate von Neumann entropy: $S_{EE} = -Tr \left[\rho_A \log \rho_A \right]$



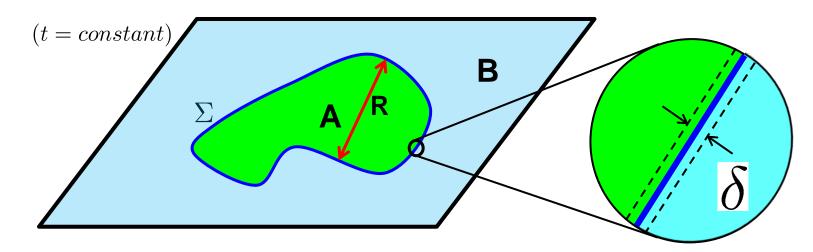
• result is UV divergent!

- remaining dof are described by a density matrix ρ_A
 - \longrightarrow calculate von Neumann entropy: $S_{EE} = -Tr \left[\rho_A \log \rho_A \right]$



• result is UV divergent! dominated by short-distance correlations

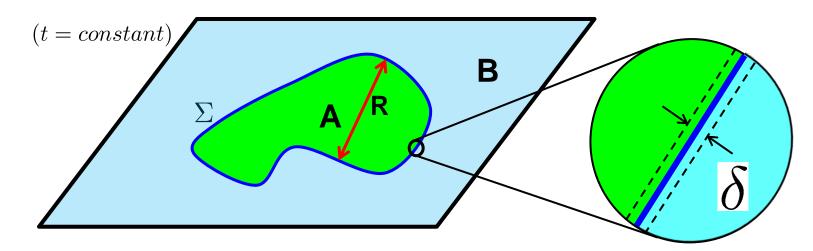
- remaining dof are described by a density matrix ρ_A
 - \longrightarrow calculate von Neumann entropy: $S_{EE} = -Tr \left[\rho_A \log \rho_A \right]$



- result is UV divergent! dominated by short-distance correlations
- must regulate calculation: $\delta = \text{short-distance cut-off}$

$$S = c_0 \frac{R^{d-2}}{\delta^{d-2}} + c_2 \frac{R^{d-4}}{\delta^{d-4}} + \cdots \qquad d = \text{spacetime dimension}$$

- remaining dof are described by a density matrix ρ_A
 - \longrightarrow calculate von Neumann entropy: $S_{EE} = -Tr \left[\rho_A \log \rho_A \right]$



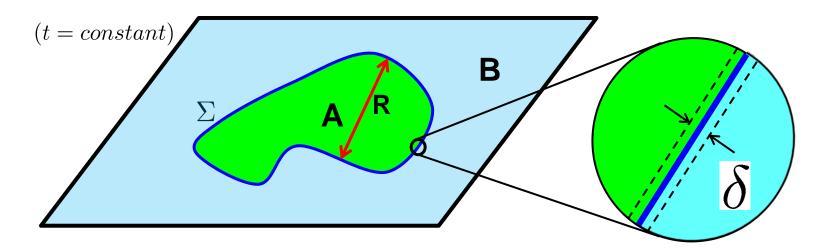
- result is UV divergent! dominated by short-distance correlations
- must regulate calculation: $\delta =$ short-distance cut-off

$$S = c_0 \frac{R^{d-2}}{\delta^{d-2}} + c_2 \frac{R^{d-4}}{\delta^{d-4}} + \cdots d$$
 = spacetime dimension

• careful analysis reveals geometric structure, eg, $S = \tilde{c}_0 \frac{A_{\Sigma}}{\delta^{d-2}} + \cdots$

- remaining dof are described by a density matrix ρ_A

 \longrightarrow calculate von Neumann entropy: $S_{EE} = -Tr \left[\rho_A \log \rho_A \right]$



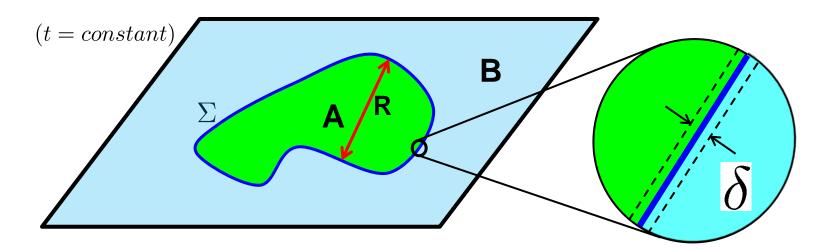
• must regulate calculation: $\delta = \text{short-distance cut-off}$

 $S = c_0 \frac{R^{d-2}}{\delta^{d-2}} + c_2 \frac{R^{d-4}}{\delta^{d-4}} + \cdots \qquad d = \text{spacetime dimension}$

- leading coefficients sensitive to details of regulator, eg, $\delta \to 2\delta$

- remaining dof are described by a density matrix ρ_A

 \longrightarrow calculate von Neumann entropy: $S_{EE} = -Tr \left[\rho_A \log \rho_A \right]$



• must regulate calculation: $\delta = \text{short-distance cut-off}$

 $S = c_0 \frac{R^{d-2}}{\delta^{d-2}} + c_2 \frac{R^{d-4}}{\delta^{d-4}} + \cdots \qquad d = \text{spacetime dimension}$

- leading coefficients sensitive to details of regulator, eg, $\delta \to 2\delta$
- find universal information characterizing underlying QFT in subleading terms, eg, $S = \cdots + c_d \log (R/\delta) + \cdots$

Entanglement C-theorem conjecture:

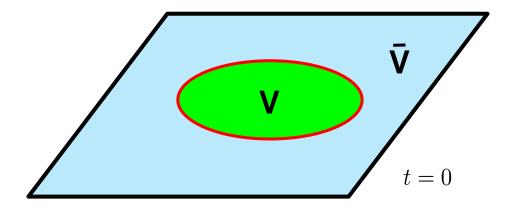
(RM & Sinha '10)

 identify central charge with universal contribution in entanglement entropy of ground state of CFT across sphere S^{d-2} of radius R:

$$S_{univ} = \begin{cases} (-)^{\frac{d}{2}-1} 4 a_d \log(2R/\delta) & \text{for even } d \\ (-)^{\frac{d-1}{2}} 2\pi a_d & \text{for odd } d \end{cases}$$

for RG flows connecting two fixed points

$$(a_d)_{UV} \geq (a_d)_{IR}$$



Entanglement C-theorem conjecture:

(RM & Sinha '10)

 identify central charge with universal contribution in entanglement entropy of ground state of CFT across sphere S^{d-2} of radius R:

$$S_{univ} = \begin{cases} (-)^{\frac{d}{2}-1} 4 a_d \log(2R/\delta) & \text{for even } d \\ (-)^{\frac{d-1}{2}} 2\pi a_d & \text{for odd } d \end{cases}$$

for RG flows connecting two fixed points

$$(a_d)_{UV} \geq (a_d)_{IR}$$

- unified framework to consider c-theorem for odd or even d
- \longrightarrow connect to Cardy's conjecture: $a_d = A$ for any CFT in even d

Entanglement C-theorem conjecture:

(RM & Sinha '10)

 identify central charge with universal contribution in entanglement entropy of ground state of CFT across sphere S^{d-2} of radius R:

$$S_{univ} = \begin{cases} (-)^{\frac{d}{2}-1} 4 a_d \log(2R/\delta) & \text{for even } d \\ (-)^{\frac{d-1}{2}} 2\pi a_d & \text{for odd } d \end{cases}$$

for RG flows connecting two fixed points

$$(a_d)_{UV} \geq (a_d)_{IR}$$

- unified framework to consider c-theorem for odd or even d
- → connect to Cardy's conjecture: $a_d = A$ for any CFT in even d eg, d=2 → $a_2 = \frac{c}{12}$, Zamolodchikov's c-theorem (entanglement proof: Casini & Huerta, '04)

d=4 $\longrightarrow a_4 = a$, a-theorem (Komargodski & Schwimmer, '11)

F-theorem: d=3

(Jafferis, Klebanov, Pufu & Safdi '11)

- examine partition function for broad classes of 3-dimensional quantum field theories on three-sphere (SUSY gauge theories, perturbed CFT's & O(N) models)
- in all examples, $F = -\log Z(S^3) > 0$ and decreases along RG flows

 \longrightarrow conjecture: $F_{UV} > F_{IR}$

• also naturally generalizes to higher odd d

F-theorem: d=3

(Jafferis, Klebanov, Pufu & Safdi '11)

- examine partition function for broad classes of 3-dimensional quantum field theories on three-sphere (SUSY gauge theories, perturbed CFT's & O(N) models)
- in all examples, $F = -\log Z(S^3) > 0$ and decreases along RG flows

 \longrightarrow conjecture: $F_{UV} > F_{IR}$

- also naturally generalizes to higher odd *d*
- coincides with entanglement c-theorem (Casini, Huerta & RM '11)

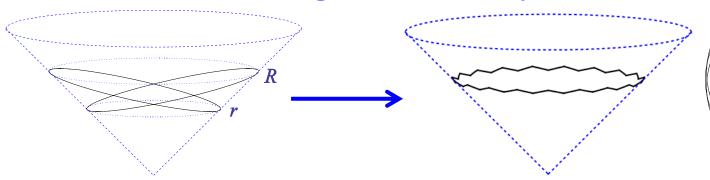
F-theorem: d=3

(Jafferis, Klebanov, Pufu & Safdi '11)

- examine partition function for broad classes of 3-dimensional quantum field theories on three-sphere (SUSY gauge theories, perturbed CFT's & O(N) models)
- in all examples, $F = -\log Z(S^3) > 0$ and decreases along RG flows

→ conjecture: $F_{UV} > F_{IR}$

- also naturally generalizes to higher odd d
- coincides with entanglement c-theorem (Casini, Huerta & RM '11)
- entanglement proof of F-theorem: (Casini & Huerta '12)
 F-theorem for d=3 RG flows established using unitarity, Lorentz invariance and strong subadditivity



Why is universal term in S_{EE} universal?

$$S_{univ} = \begin{cases} (-)^{\frac{d}{2}-1} 4 a_d^* \log(2R/\delta) & \text{for even } d \\ (-)^{\frac{d-1}{2}} 2\pi a_d^* & \text{for odd } d \end{cases}$$

(Schwimmer & Theisen)

• QFT intution: log divergences define physical cuts but finite *p* polynomials subject to renormalization ambiguities

Why is universal term in S_{EE} universal?

$$S_{univ} = \begin{cases} (-)^{\frac{d}{2}-1} 4 a_d^* \log(2R/\delta) & \text{for even } d \\ (-)^{\frac{d-1}{2}} 2\pi a_d^* & \text{for odd } d \end{cases}$$

(Schwimmer & Theisen)

- QFT intution: log divergences define physical cuts but finite *p* polynomials subject to renormalization ambiguities
- ----> even d seems okay but odd d might be problematic?

Why is universal term in S_{EE} universal?

$$S_{univ} = \begin{cases} (-)^{\frac{d}{2}-1} 4 a_d^* \log(2R/\delta) & \text{for even } d \\ (-)^{\frac{d-1}{2}} 2\pi a_d^* & \text{for odd } d \end{cases}$$

(Schwimmer & Theisen)

- QFT intution: log divergences define physical cuts but finite *p* polynomials subject to renormalization ambiguities
- ----> even d seems okay but odd d might be problematic?
- shifting $\delta \to \delta' = \delta + \alpha \, m \delta^2$, constant term polluted by UV data

d=3:
$$S(R) = \frac{2\pi R}{\delta} c_0 - 2\pi a_3$$

Why is universal term in S_{EE} universal?

$$S_{univ} = \begin{cases} (-)^{\frac{d}{2}-1} 4 a_d^* \log(2R/\delta) & \text{for even } d \\ (-)^{\frac{d-1}{2}} 2\pi a_d^* & \text{for odd } d \end{cases}$$

(Schwimmer & Theisen)

- QFT intution: log divergences define physical cuts but finite *p* polynomials subject to renormalization ambiguities
- ----> even d seems okay but odd d might be problematic?
- shifting $\delta \to \delta' = \delta + \alpha \, m \delta^2$, constant term polluted by UV data

 \longrightarrow sure but no scales in CFT, so no scale m!!

d=3:
$$S(R) = \frac{2\pi R}{\delta} c_0 - 2\pi a_3$$

Why is universal term in S_{EE} universal?

$$S_{univ} = \begin{cases} (-)^{\frac{d}{2}-1} 4 a_d^* \log(2R/\delta) & \text{for even } d \\ (-)^{\frac{d-1}{2}} 2\pi a_d^* & \text{for odd } d \end{cases}$$

(Schwimmer & Theisen)

- QFT intution: log divergences define physical cuts but finite *p* polynomials subject to renormalization ambiguities
- ----> even d seems okay but odd d might be problematic?
- shifting $\delta \to \delta' = \delta + \alpha \, m \delta^2$, constant term polluted by UV data

 \longrightarrow sure but no scales in CFT, so no scale m!!

 \longrightarrow scales from RG flow can appear in final S_{EE}!!

(eg, Hertzberg & Wilczek; Banerjee)

d=3:
$$S(R) = 2\pi R \left(\frac{c_0}{\delta} + m c_1\right) - 2\pi a_3$$

Why is universal term in S_{EE} universal?

$$S_{univ} = \begin{cases} (-)^{\frac{d}{2}-1} 4 a_d^* \log(2R/\delta) & \text{for even } d \\ (-)^{\frac{d-1}{2}} 2\pi a_d^* & \text{for odd } d \end{cases}$$

(Schwimmer & Theisen)

- QFT intution: log divergences define physical cuts but finite *p* polynomials subject to renormalization ambiguities
- ----> even d seems okay but odd d might be problematic?
- shifting $\delta \to \delta' = \delta + \alpha \, m \delta^2$, constant term polluted by UV data

 \longrightarrow sure but no scales in CFT, so no scale m!!

 \longrightarrow scales from RG flow can appear in final S_{EE}!!

(eg, Hertzberg & Wilczek; Banerjee)

• in regulators, tension between Lorentz inv. and unitarity

 \longrightarrow latter emerge in $\delta \rightarrow 0$ limit, but regulator exposed in S_{EE}

(Liu & Mezei)

 divergences determined by local geometry of entangling surface with covariant regulator, eg,

$$S_{CFT} = c_0 \frac{R^{d-2}}{\delta^{d-2}} + c_2 \frac{R^{d-4}}{\delta^{d-4}} + \dots + (-)^{\frac{d-1}{2}} 2\pi a_d + O(\delta/R)$$

• can isolate finite term with appropriate manipulations, eg,

d=3:
$$S_3(R) = RS'(R) - S(R)$$

d=4:
$$S_4(R) = R^2 S''(R) - RS'(R)$$

(Liu & Mezei)

• divergences determined by local geometry of entangling surface with covariant regulator, eg,

$$S = c_0(\mu_i \delta) \frac{R^{d-2}}{\delta^{d-2}} + c_2(\mu_i \delta) \frac{R^{d-4}}{\delta^{d-4}} + \dots + (-)^{\frac{d-1}{2}} 2\pi a_d(\mu_i \delta) + O(\delta/R)$$

• can isolate finite term with appropriate manipulations, eg,

d=3:
$$S_3(R) = RS'(R) - S(R)$$

d=4:
$$S_4(R) = R^2 S''(R) - RS'(R)$$

(Liu & Mezei)

 divergences determined by local geometry of entangling surface with covariant regulator, eg,

$$S = c_0(\mu_i \delta) \frac{R^{d-2}}{\delta^{d-2}} + c_2(\mu_i \delta) \frac{R^{d-4}}{\delta^{d-4}} + \dots + (-)^{\frac{d-1}{2}} 2\pi a_d(\mu_i \delta) + O(\delta/R)$$

- can isolate finite term with appropriate manipulations, eg,
 - d=3: $S_3(R) = RS'(R) S(R)$ d=4: $S_4(R) = R^2 S''(R) - RS'(R)$ c-function of Casini & Huerta
 - (unfortunately, holographic experiments indicate $S_d(R)$ are **not** good C-functions for d>3 not monotonic)

(Liu & Mezei)

 divergences determined by local geometry of entangling surface with covariant regulator, eg,

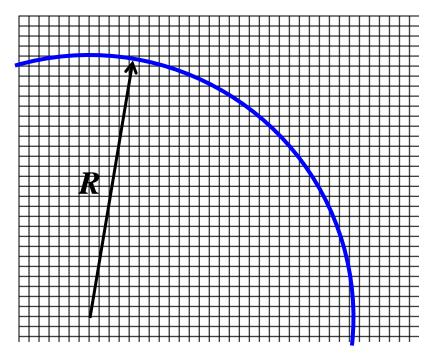
$$S = c_0(\mu_i \delta) \frac{R^{d-2}}{\delta^{d-2}} + c_2(\mu_i \delta) \frac{R^{d-4}}{\delta^{d-4}} + \dots + (-)^{\frac{d-1}{2}} 2\pi a_d(\mu_i \delta) + O(\delta/R)$$

- can isolate finite term with appropriate manipulations, eg,
 - d=3: $S_3(R) = RS'(R) S(R)$ d=4: $S_4(R) = R^2 S''(R) - RS'(R)$ c-function of Casini & Huerta

(unfortunately, holographic experiments indicate $S_d(R)$ are **not** good C-functions for d>3 – not monotonic)

- approach demands special class of regulators: "covariant"
 is result artifact of choosing "nice" regulator??
- if a_d is physical, we should be able to use any regularization which defines the continuum QFT

$$d = 3$$
: $S(R) = \frac{2\pi R}{\delta} c_0 - 2\pi a_3$



$$d = 3 : S(R) = \frac{2\pi R}{\delta} c_0 - 2\pi a_3$$
• circumference always uncertain to $O(\delta)$

$$R \rightarrow R' = R + \alpha \delta,$$

$$d = 3$$
: $S(R) = \frac{2\pi R}{\delta} c_0 - 2\pi a_3$

- circumference always uncertain to $O(\delta)$

$$R \to R' = R + \alpha \, \delta,$$

 $\longrightarrow a_3$ always polluted by UV

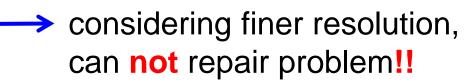
$$d = 3: S(R) = \frac{2\pi R}{\delta} c_0 - 2\pi a_3$$

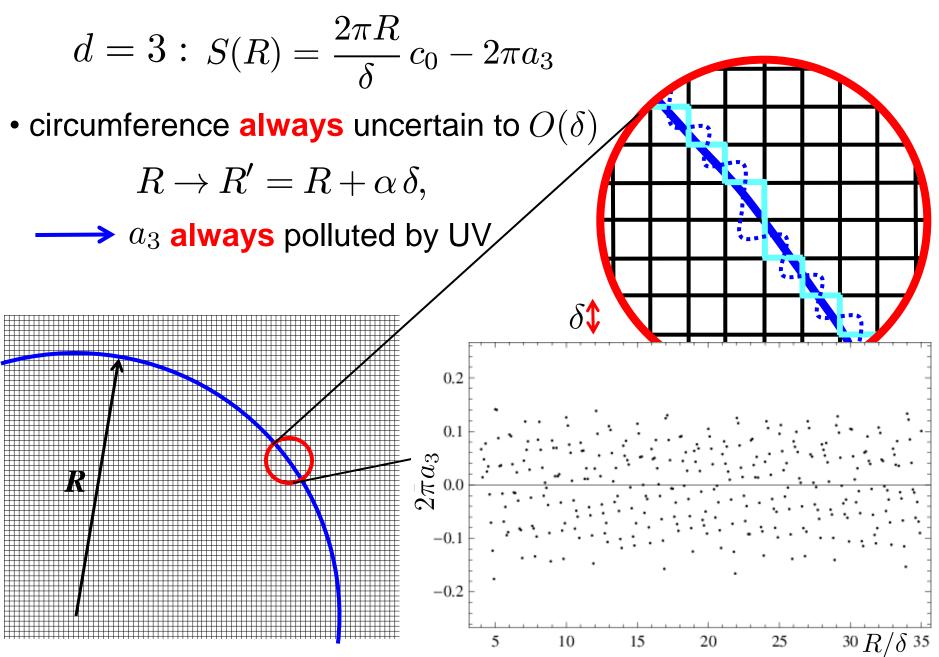
• circumference always uncertain to $O(\delta)$

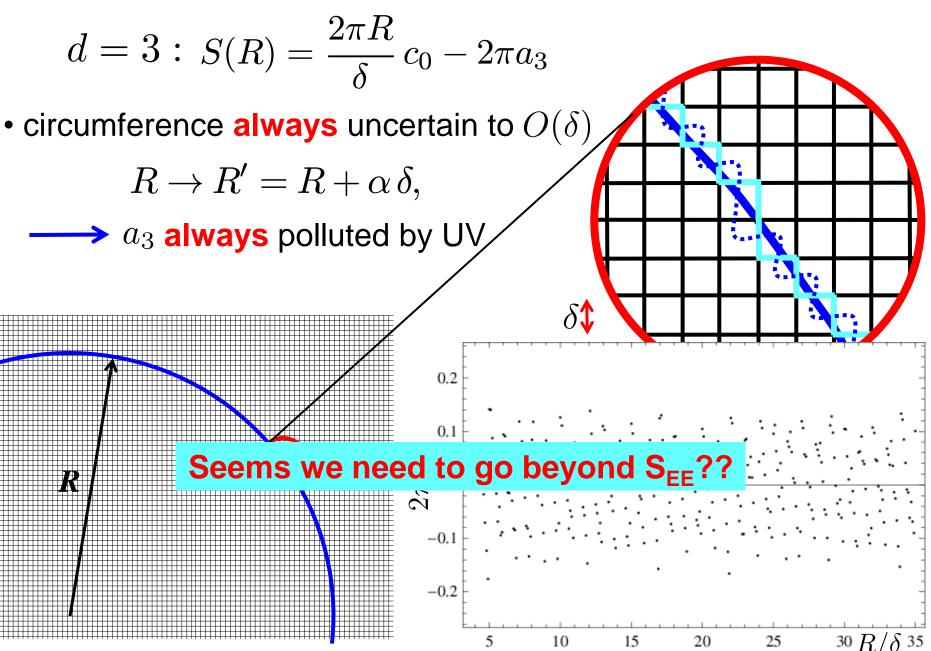
$$R \to R' = R + \alpha \, \delta,$$

 $\rightarrow a_3$ always polluted by UV/

R







1. C-function must be dimensionless, well-defined quantity, which is independent of the regularization scheme

----> computable with any regulator

1. C-function must be dimensionless, well-defined quantity, which is independent of the regularization scheme

computable with any regulator

2. C-function must be intrinsic to fixed point of interest

independent of details of RG flows

1. C-function must be dimensionless, well-defined quantity, which is independent of the regularization scheme

computable with any regulator

- C-function must be intrinsic to fixed point of interest
 independent of details of RG flows
- 3. C-function must decrease monotonically along any RG flows connecting a UV fixed point to an IR fixed point

1. C-function must be dimensionless, well-defined quantity, which is independent of the regularization scheme

computable with any regulator

- C-function must be intrinsic to fixed point of interest
 → independent of details of RG flows
- 3. C-function must decrease monotonically along any RG flows connecting a UV fixed point to an IR fixed point

- S_{EE} seems to fail to satisfy criteria 1 & 2
- alternate choice? alternate measure of entanglement?

Mutual Information:

- another measure of entanglement between two systems
- for non-intersecting regions A and B:

 $I(A,B) = S(A) + S(B) - S(A \cup B)$

- can be defined without reference to S_{EE}
 (Araki; Narnhofer)
- bounds correlators between A and B (Wolf, Verstraete, Hastings & Cirac)

$$I(A,B) \ge \frac{|\langle \mathcal{O}_A \, \mathcal{O}_B \rangle_c|^2}{2 \|\mathcal{O}_A\|^2 \, \|\mathcal{O}_B\|^2}$$

finite! UV divergences in S(A) and S(B) canceled by S(A U B)

Mutual Information:

- another measure of entanglement between two systems
- for non-intersecting regions A and B:

 $I(A,B) = S(A) + S(B) - S(A \cup B)$

- can be defined without reference to S_{EE}
 (Araki; Narnhofer)
- bounds correlators between A and B (Wolf, Verstraete, Hastings & Cirac)

$$I(A,B) \ge \frac{|\langle \mathcal{O}_A \, \mathcal{O}_B \rangle_c|^2}{2 \|\mathcal{O}_A\|^2 \, \|\mathcal{O}_B\|^2}$$

- finite! UV divergences in S(A) and S(B) canceled by S(A U B)
- if c-function defined with mutual information
 criterion 1 will automatically be satisfied

Mutual Information:

- another measure of entanglement between two systems
- for non-intersecting regions A and B:

 $I(A,B) = S(A) + S(B) - S(A \cup B)$

- can be defined without reference to S_{EE}
 (Araki; Narnhofer)
- bounds correlators between A and B (Wolf, Verstraete, Hastings & Cirac)

$$I(A,B) \ge \frac{|\langle \mathcal{O}_A \, \mathcal{O}_B \rangle_c|^2}{2 \|\mathcal{O}_A\|^2 \, \|\mathcal{O}_B\|^2}$$

- finite! UV divergences in S(A) and S(B) canceled by S(A U B)
- if c-function defined with mutual information

criterion 1 will automatically be satisfied

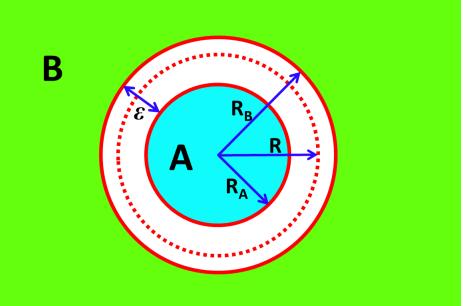
criterion 2 & 3 will be satisfied with further care

$$I(A,B) = S(A) + S(B) - S(A \cup B)$$

consider following geometry:

$$R_A = R - \left(\frac{1}{2} - \alpha\right)\varepsilon$$
$$R_B = R + \left(\frac{1}{2} + \alpha\right)\varepsilon$$

or
$$R = \frac{R_A + R_B}{2} - \alpha \varepsilon$$



• using $S(A) = S(\overline{A})$ for pure state: $I(A, B) = S(A) + S(\overline{B}) - S(\overline{A \cup B})$ two disks ~ R ______ narrow annulus

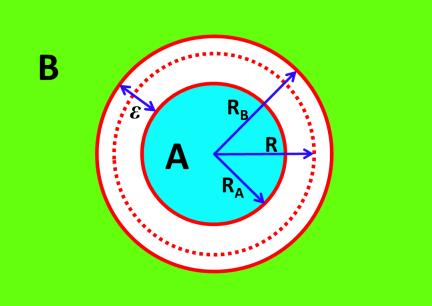
work in continuum limit: $R \gg \varepsilon \gg \delta$ (R and ε are macroscopic scales)

$$I(A,B) = S(A) + S(B) - S(A \cup B)$$

• consider following geometry:

$$R_A = R - \left(\frac{1}{2} - \alpha\right)\varepsilon$$
$$R_B = R + \left(\frac{1}{2} + \alpha\right)\varepsilon$$

or
$$R = \frac{R_A + R_B}{2} - \alpha \varepsilon$$

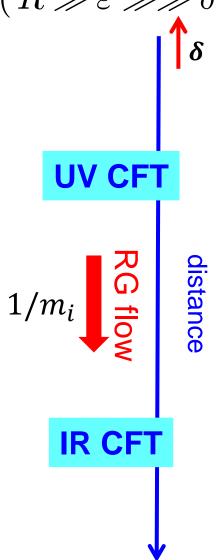


- work in continuum limit: $R\gg \varepsilon \gg \delta$
- mutual information takes form:

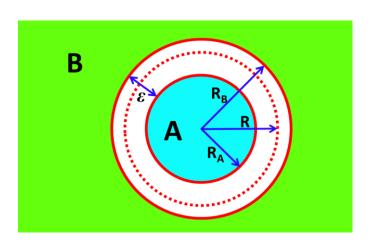
$$I(A,B) = 2\pi R \left(\frac{\tilde{c}_0}{\varepsilon} + \tilde{c}_1\right) - 4\pi \,\tilde{a}_3 + O(\varepsilon/R)$$

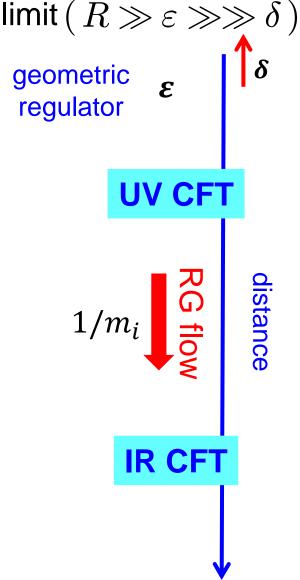
- mutual information "regulates" entanglement entropy of disk
- work with renormalized QFT in continuum limit ($R \gg \varepsilon \gg \delta$)

- mutual information "regulates" entanglement entropy of disk
- work with renormalized QFT in continuum limit ($R \gg \varepsilon \gg \delta$)
- Strategy:

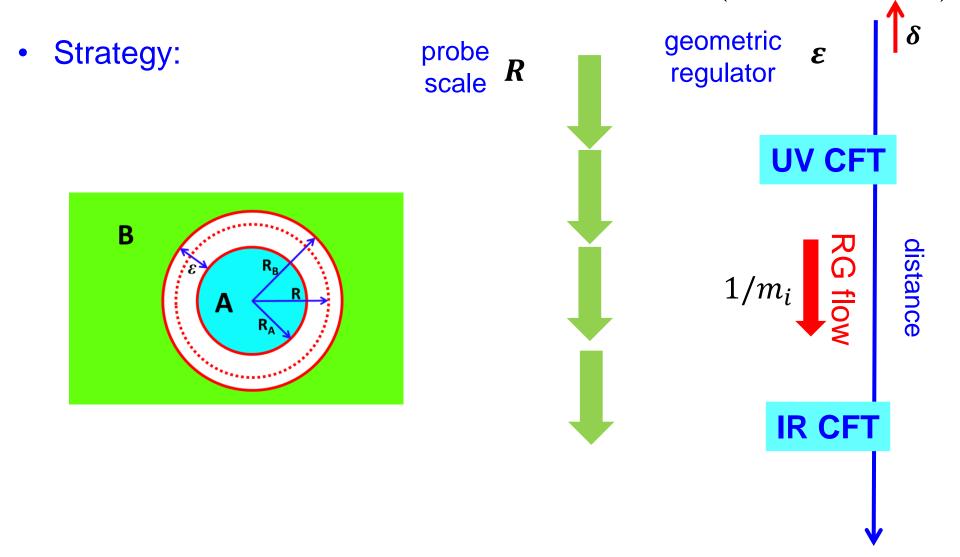


- mutual information "regulates" entanglement entropy of disk
- work with renormalized QFT in continuum limit ($R \gg \varepsilon \gg \delta$)
- Strategy:

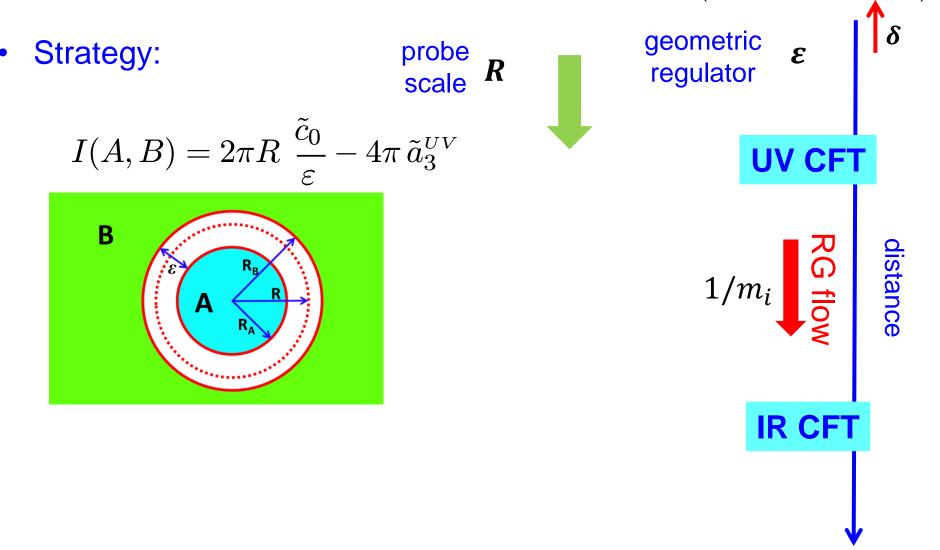




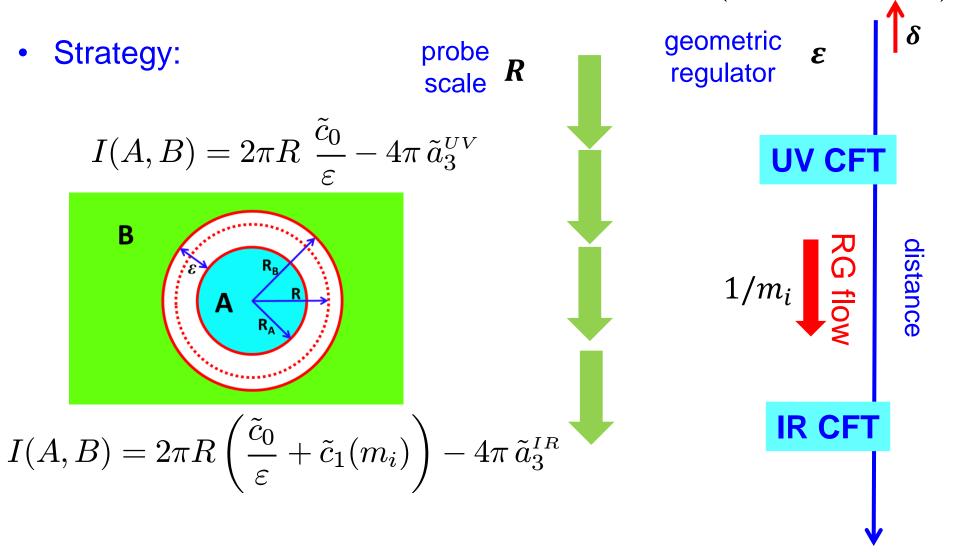
- mutual information "regulates" entanglement entropy of disk
- work with renormalized QFT in continuum limit ($R \gg \varepsilon \gg \delta$)



- mutual information "regulates" entanglement entropy of disk
- work with renormalized QFT in continuum limit ($R \gg \varepsilon \gg \delta$)



- mutual information "regulates" entanglement entropy of disk
- work with renormalized QFT in continuum limit ($R \gg \varepsilon \gg \delta$)

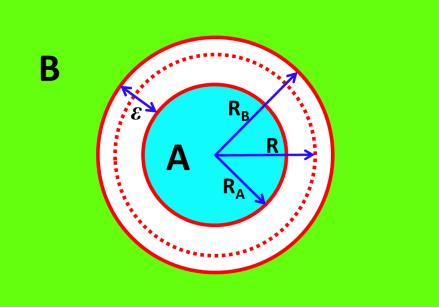


$$I(A,B) = S(A) + S(B) - S(A \cup B)$$

• consider following geometry:

$$R_A = R - \left(\frac{1}{2} - \alpha\right)\varepsilon$$
$$R_B = R + \left(\frac{1}{2} + \alpha\right)\varepsilon$$

or
$$R = \frac{R_A + R_B}{2} - \alpha \varepsilon$$



- work in continuum limit: $R\gg \varepsilon \gg \gg \delta$
- mutual information takes form:

$$I(A,B) = 2\pi R \left(\frac{\tilde{c}_0}{\varepsilon} + \tilde{c}_1\right) - 4\pi \,\tilde{a}_3 + O(\varepsilon/R)$$

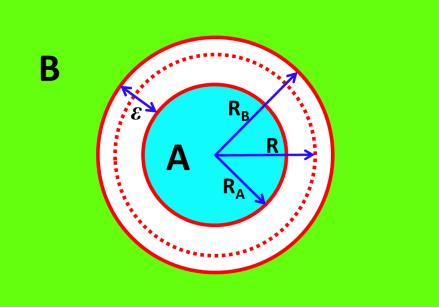
• criterion 2? is \tilde{a}_3 intrinsic to fixed point??

$$I(A,B) = S(A) + S(B) - S(A \cup B)$$

consider following geometry:

$$R_A = R - \left(\frac{1}{2} - \alpha\right)\varepsilon$$
$$R_B = R + \left(\frac{1}{2} + \alpha\right)\varepsilon$$

or
$$R = \frac{R_A + R_B}{2} - \alpha \varepsilon$$



- work in continuum limit: $R\gg \varepsilon \gg \gg \delta$
- mutual information takes form:

$$I(A,B) = 2\pi R \left(\frac{\tilde{c}_0}{\varepsilon} + \tilde{c}_1\right) - 4\pi \,\tilde{a}_3 + O(\varepsilon/R)$$

• ambiguity: $\alpha \to \alpha' = \alpha + \delta \alpha$, $\tilde{a}_3 \to \tilde{a}'_3 = \tilde{a}_3 + \tilde{c}_0 \, \delta \alpha$

• can we choose α such that \tilde{a}_3 is independent of higher scales?

- can we choose α such that \tilde{a}_3 is independent of higher scales?
- consider probing at IR critical point where *m*, lowest mass scale in RG flow:: $R \gg 1/m \gg \varepsilon$

2/m

- correlations near boundary nonconformal
- high energy contribution to I(A,B): local and extensive

$$I(A,B)_{HE} = 2\pi R \left(\sigma_0 + \frac{\sigma_1}{R} + \frac{\sigma_2}{R^2} + \cdots\right)$$

- can we choose α such that \tilde{a}_3 is independent of higher scales?
- consider probing at IR critical point where *m*, lowest mass scale in RG flow:: $R \gg 1/m \gg \varepsilon$

2/m

- correlations near boundary nonconformal
- high energy contribution to I(A,B): local and extensive

$$I(A,B)_{HE} = 2\pi R \left(\sigma_0 \left(\frac{\sigma_1}{R} \right) \frac{\sigma_2}{R^2} + \cdots \right)$$

• can we choose α to eliminate σ_1 ??

- can we choose α such that \tilde{a}_3 is independent of higher scales?
- consider probing at IR critical point where *m*, lowest mass scale in RG flow:: $R \gg 1/m \gg \varepsilon$
- correlations near boundary nonconformal
- high energy contribution to I(A,B): local and extensive

$$I(A,B)_{HE} = 2\pi R \left(\sigma_0 \left(\frac{\sigma_1}{R} \right) \frac{\sigma_2}{R^2} + \cdots \right)$$

- can we choose α to eliminate σ_1 ??
- for general strip (with small curvatures):

$$I(A,B)_{HE} = \int ds \, \left(\sigma_0 - \sigma_1 \,\mathbf{n} \cdot \partial_s \mathbf{t} - \sigma_2 \,\mathbf{t} \cdot \partial_s^2 \mathbf{t} + \cdots\right)$$

2/m

- can we choose α such that \tilde{a}_3 is independent of higher scales?
- consider probing at IR critical point where *m*, lowest mass scale in RG flow:: $R \gg 1/m \gg \varepsilon$
- correlations near boundary nonconformal
- high energy contribution to I(A,B): local and extensive

$$I(A,B)_{HE} = 2\pi R \left(\sigma_0 \left(\frac{\sigma_1}{R} \right) \frac{\sigma_2}{R^2} + \cdots \right)$$

- can we choose α to eliminate σ_1 ??
- for general strip (with small curvatures):

$$I(A,B)_{HE} = \int ds \, \left(\sigma_0 - \sigma_1 \,\mathbf{n} \cdot \partial_s \mathbf{t} - \sigma_2 \,\mathbf{t} \cdot \partial_s^2 \mathbf{t} + \cdots\right)$$

2/m

• σ_1 must vanish if reflection symmetry $\longrightarrow \alpha = 0$

$$I(A,B) = S(A) + S(B) - S(A \cup B)$$

• consider following geometry:

$$R_A = R - \varepsilon/2$$
$$R_B = R + \varepsilon/2$$

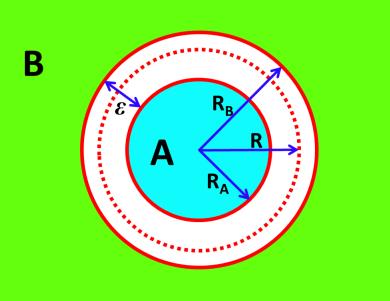
$$R = \frac{R_A + R_B}{2}$$

- work in continuum limit: $R \gg \varepsilon \gg \delta$
- mutual information takes form:

$$I(A,B) = 2\pi R \left(\frac{\tilde{c}_0}{\varepsilon} + \tilde{c}_1\right) - 4\pi \,\tilde{a}_3 + O(\varepsilon/R)$$

• fixing $\alpha = 0$ ensures \tilde{a}_3 is intrinsic to fixed point

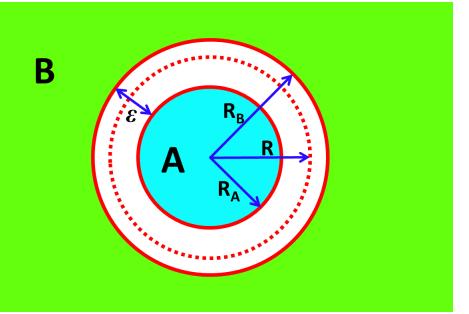
criteria 1 and 2 are satisfied!!



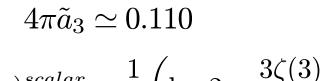
consider following geometry:

$$R = \frac{R_A + R_B}{2}$$

- work in continuum limit: $R \gg \varepsilon \gg \delta$



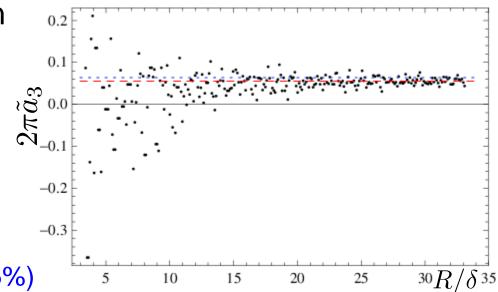
 calculate for a free scalar on a square lattice:



$$(4\pi a_3)^{scalar} = \frac{1}{4} \left(\log 2 - \frac{\delta \zeta(0)}{2\pi^2} \right)$$

\$\approx 0.127\$

($R: \varepsilon: \delta = 33: 6: 1$, result good to 15%)



1. C-function must be dimensionless, well-defined quantity, which is independent of the regularization scheme

computable with any regulator

- C-function must be intrinsic to fixed point of interest
 Independent of details of RG flows
- 3. C-function must decrease monotonically along any RG flows connecting a UV fixed point to an IR fixed point
- defining \tilde{a}_3 with mutual information & fixing $\alpha = 0$ ensures criteria 1 and 2 are satisfied; must still consider criterion 3

1. C-function must be dimensionless, well-defined quantity, which is independent of the regularization scheme

computable with any regulator

- C-function must be intrinsic to fixed point of interest
 Independent of details of RG flows
- 3. C-function must decrease monotonically along any RG flows connecting a UV fixed point to an IR fixed point
- defining \tilde{a}_3 with mutual information & fixing $\alpha = 0$ ensures criteria 1 and 2 are satisfied; must still consider criterion 3
- all of above easily generalizes to higher odd dimensions

1. C-function must be dimensionless, well-defined quantity, which is independent of the regularization scheme

-----> computable with any regulator

- C-function must be intrinsic to fixed point of interest
 Independent of details of RG flows
- 3. C-function must decrease monotonically along any RG flows connecting a UV fixed point to an IR fixed point
- defining \tilde{a}_3 with mutual information & fixing $\alpha = 0$ ensures criteria 1 and 2 are satisfied; must still consider criterion 3
- all of above easily generalizes to higher odd dimensions
- monotonic flow follows as in entropic proof of F-theorem

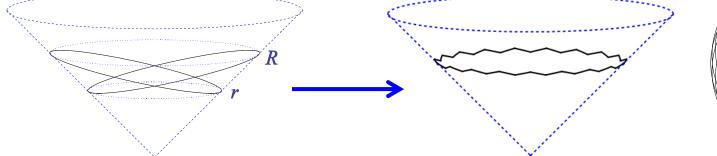
(Casini & Huerta '12)

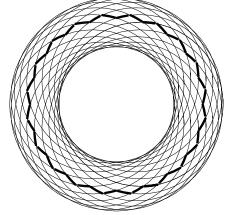
Entanglement proof of F-theorem:

 F-theorem for d=3 RG flows established using unitarity, Lorentz invariance and strong subadditivity

 $\sum S(X_i) \ge S(\cup_i X_i) + S(\cup_{\{ij\}} (X_i \cap X_j)) + S(\cup_{\{ijk\}} (X_i \cap X_j \cap X_k)) + \dots + S(\cap_i X_i)$

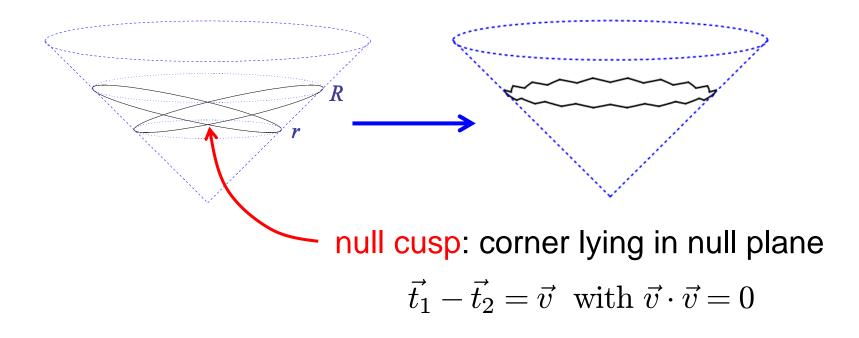
 geometry more complex than d=2: consider many circles intersecting on null cone





- no corner contribution from intersection in null plane
- define: C(R) = RS'(R) S(R)
- for d=3 CFT: $S(R) = \frac{2\pi R}{\delta}c_0 2\pi a_3 \longrightarrow C_{\rm CFT}(R) = 2\pi a_3$
- with SSA and "continuum" limit $\longrightarrow \partial_R C(R) \leq 0$
- hence C(R) decreases monotonically and $[a_3]_{\rm UV} > [a_3]_{\rm IR}$

- key ingredients:
- a) unitary & Lorentz invariant regularization of EE defined on regions with smooth boundaries except for "null cusps"
- b) regulated EE satisfies strong subaddivity for sets whose union and intersection only generates more "null cusps"
- c) wiggly circles have EE which approaches that of circle with same perimeter as the number of null cusps goes to ∞

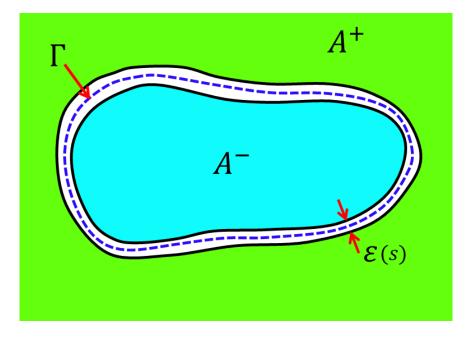


• mutual information approach satisfy these key ingredients?

• mutual information approach satisfy these key ingredients? yes

- mutual information approach satisfy these key ingredients? yes
- consider region A with smooth boundary Γ
- expand boundary: $\Gamma_{\pm} = \Gamma \pm \frac{1}{2} \epsilon(s) \hat{n}(s)$

 $I(A^+, A^-) = \tilde{c}_0 \oint_{\Gamma} \frac{ds}{\varepsilon(s)} + \frac{I_0(A)}{\varepsilon(s)} + O(\varepsilon)$

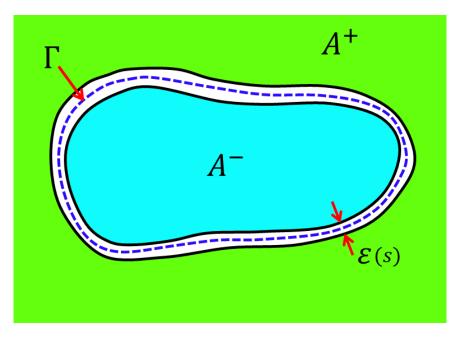


- mutual information approach satisfy these key ingredients? yes
- consider region A with smooth boundary Γ
- expand boundary: $\Gamma_{\pm} = \Gamma \pm \frac{1}{2} \epsilon(s) \hat{n}(s)$

 $I(A^+, A^-) = \tilde{c}_0 \quad \oint_{\Gamma} \frac{ds}{\varepsilon(s)} + \underbrace{I_0(A)}_{\varepsilon(s)} + O(\varepsilon)$

 regulated EE: property of A; independent of framing

eg, for circle $I_0(A) = 2\pi R \, \tilde{c}_1(m_i) - 4\pi \, \tilde{a}_3$

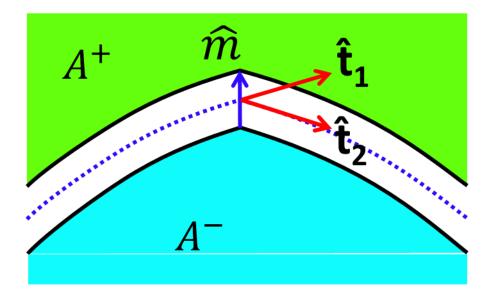


- mutual information approach satisfy these key ingredients? yes
- consider region A with smooth boundary Γ with null cusps
- expand boundary: $\Gamma_{\pm} = \Gamma \pm \frac{1}{2} \epsilon(s) \hat{n}(s)$

 $I(A^+, A^-) = \tilde{c}_0 \ \oint_{\Gamma} \frac{ds}{\varepsilon(s)} + \ I_0(A) + \sum f(q_{1i}, q_{2i}) + O(\varepsilon)$

 additional contributions for null cusps characterized by two local invariants:

 $q_1 = \widehat{m} \cdot \widehat{t}_1 \quad q_2 = \widehat{m} \cdot \widehat{t}_2$



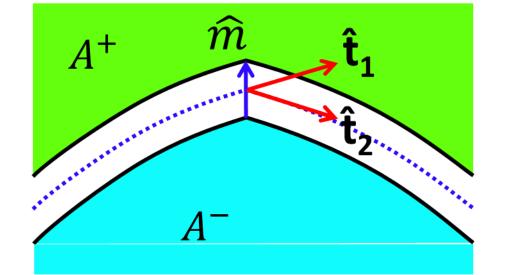
- mutual information approach satisfy these key ingredients? yes
- consider region A with smooth boundary Γ with null cusps
- expand boundary: $\Gamma_{\pm} = \Gamma \pm \frac{1}{2} \epsilon(s) \hat{n}(s)$

 $I(A^+, A^-) = \tilde{c}_0 \ \oint_{\Gamma} \frac{ds}{\varepsilon(s)} + \ I_0(A) + \sum f(q_{1i}, q_{2i}) + O(\varepsilon)$

 additional contributions for null cusps characterized by two local invariants:

 $q_1 = \widehat{m} \cdot \widehat{t}_1 \quad q_2 = \widehat{m} \cdot \widehat{t}_2$

• *I*₀(*A*) still satisfies SSA:



 $I_0(A) + I_0(B) \ge I_0(A \cup B) + I_0(A \cap B)$

1. C-function must be dimensionless, well-defined quantity, which is independent of the regularization scheme

computable with any regulator

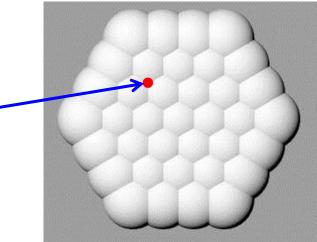
- C-function must be intrinsic to fixed point of interest
 Independent of details of RG flows
- 3. C-function must decrease monotonically along any RG flows connecting a UV fixed point to an IR fixed point
- defining \tilde{a}_3 with mutual information & fixing $\alpha=0\,$ ensures criteria 1 and 2 are satisfied
- monotonic flow follows as in entropic proof of F-theorem

have properly established F-theorem in d=3

• is there entropic proof of c-theorem in higher dimensions?

→ need a new idea?

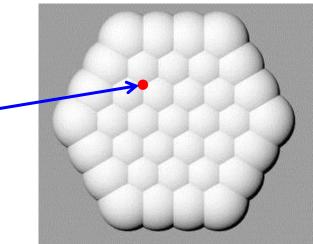
higher dim. intersections lead to subleading divergences which trivialize SSA inequality



• is there entropic proof of c-theorem in higher dimensions?

→ need a new idea?

higher dim. intersections lead to subleading divergences which trivialize SSA inequality

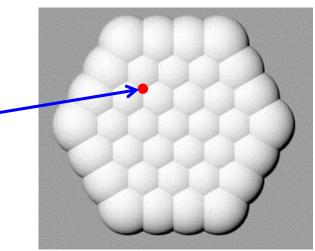


 d=4 a-theorem proved with more "standard" QFT techniques (Komargodski & Schwimmer)

• is there entropic proof of c-theorem in higher dimensions?

→ need a new idea?

higher dim. intersections lead to subleading divergences which trivialize SSA inequality

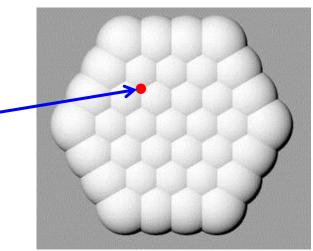


- d=4 a-theorem proved with more "standard" QFT techniques (Komargodski & Schwimmer)
- hybrid approach proposed (Solodukhin): still needs development

• is there entropic proof of c-theorem in higher dimensions?

→ need a new idea?

higher dim. intersections lead to subleading divergences which trivialize SSA inequality



- d=4 a-theorem proved with more "standard" QFT techniques (Komargodski & Schwimmer)
- hybrid approach proposed (Solodukhin): still needs development
- can c-theorems be proved for higher dimensions? eg, d=5 or 6
 - → again, entropic approach needs a new idea
 - dilaton-effective-action approach requires refinement for d=6 (Elvang, Freedman, Hung, Kiermaier, RM & Theisen; Elvang & Olson)

Conclusions and Questions:

- entanglement lends new insights into c-theorems
- using mutual information, properly established d=3 F-theorem
- how much of Zamolodchikov's structure for d=2 RG flows extends higher dimensions?
- d=3 entropic C-function not always stationary at fixed points (Klebanov, Nishioka, Pufu & Safdi)
- same already observed for d=2; special case or generic? need a better C-function?
- does scale invariance imply conformal invariance beyond d=2?

"more or less" in d=4 (Luty, Polchinski & Rattazzi; Dymarsky, Komargodski, Schwimmer & Theisen)

SSA ---- NEC (Bhattacharya etal; Lashkari et al; Lin etal)

 what more entanglement/quantum information have to teach us about RG flows, holography or nonperturbative QFT?