

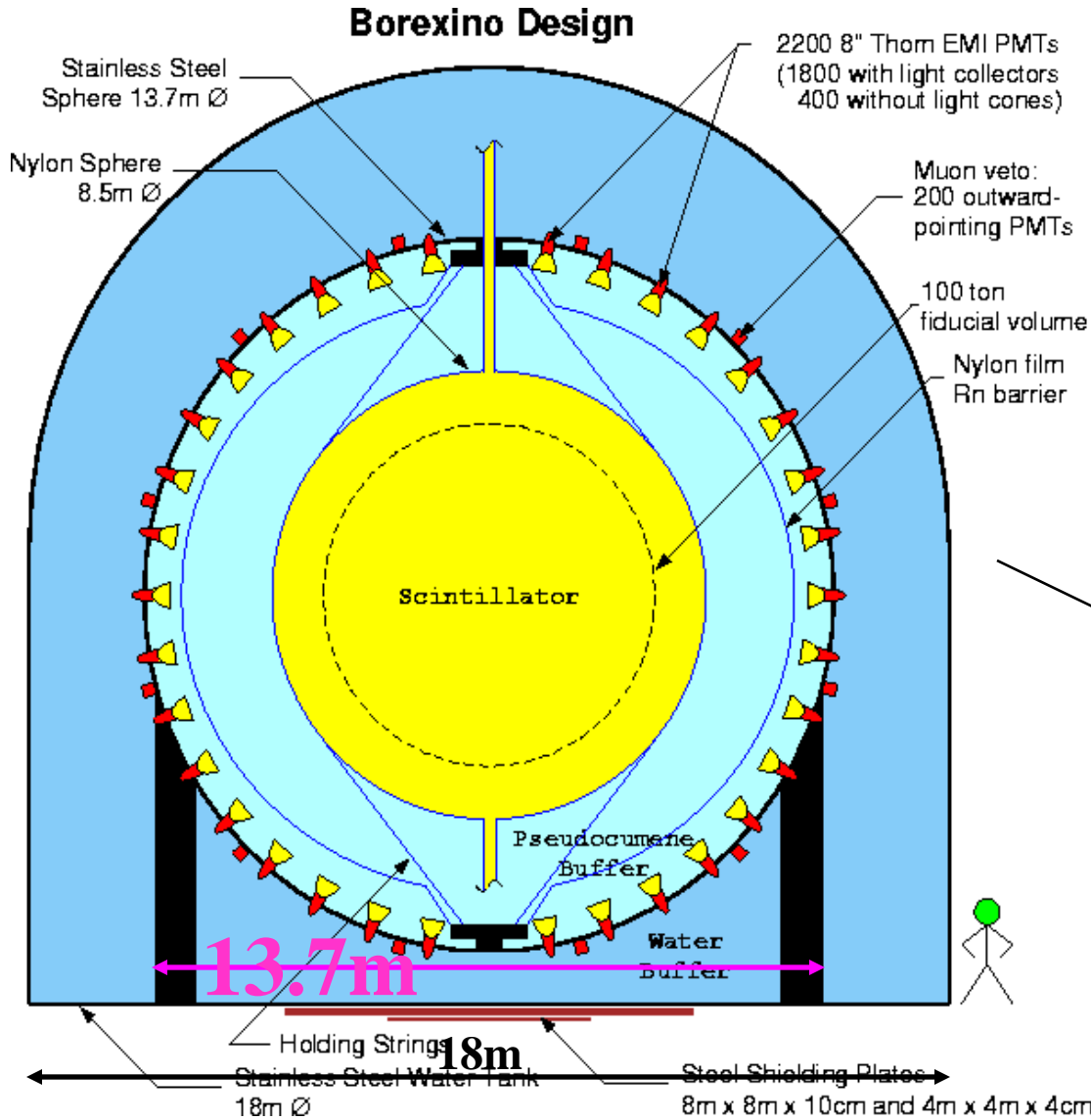
The measurement of the solar neutrinos from the pp nuclear reaction

Oleg Smirnov, JINR, Dubna

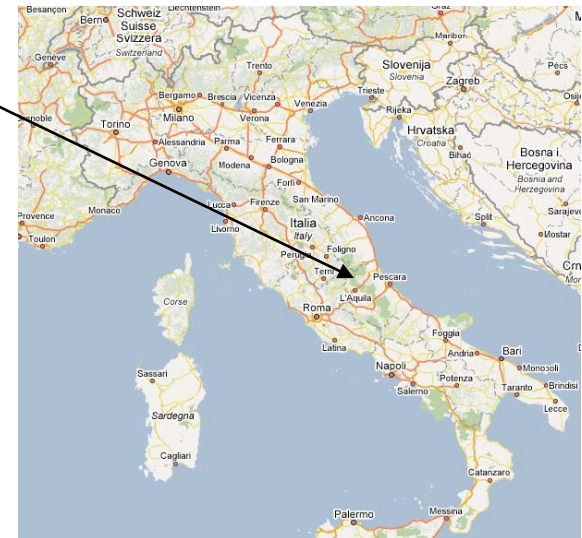
On behalf of Borexino collaboration

Workshop on the Borexino Physics, LNGS, September 5, 2014

BOREXINO (in operation from May, 2007)

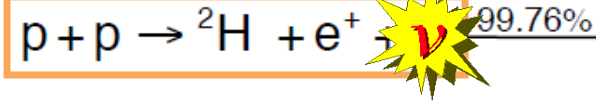


- 278 t of liquid organic scintillator PC + PPO (1.5 g/l)
- (v,e)-scattering with low threshold (~200 keV)
- Outer muon detector

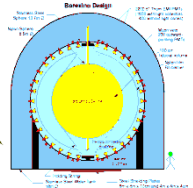
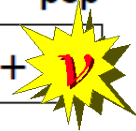
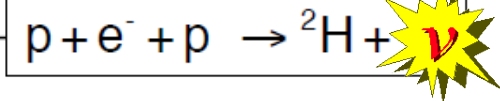


pp-chain

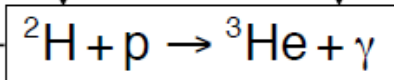
pp



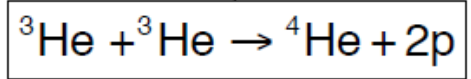
pep



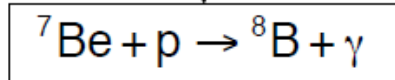
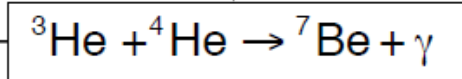
83.30%



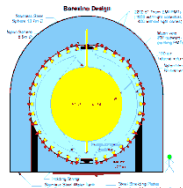
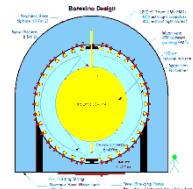
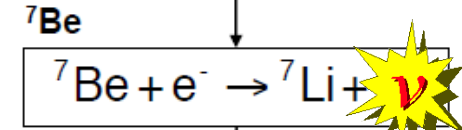
16.70%



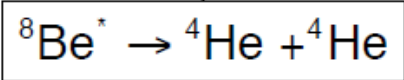
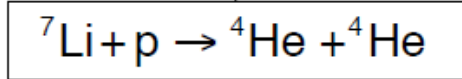
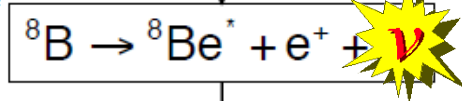
0.12%

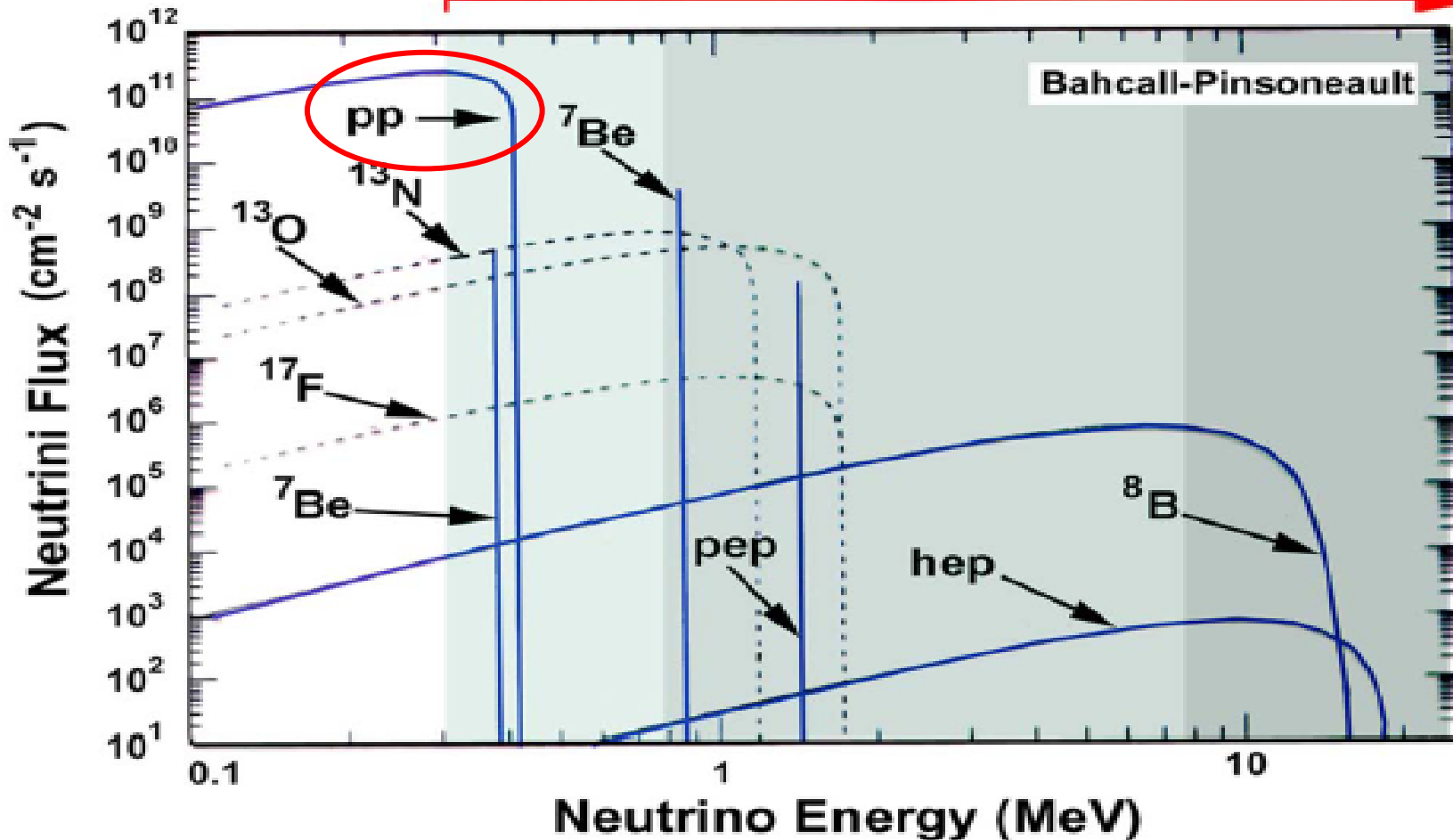


99.88%



${}^8\text{B}$



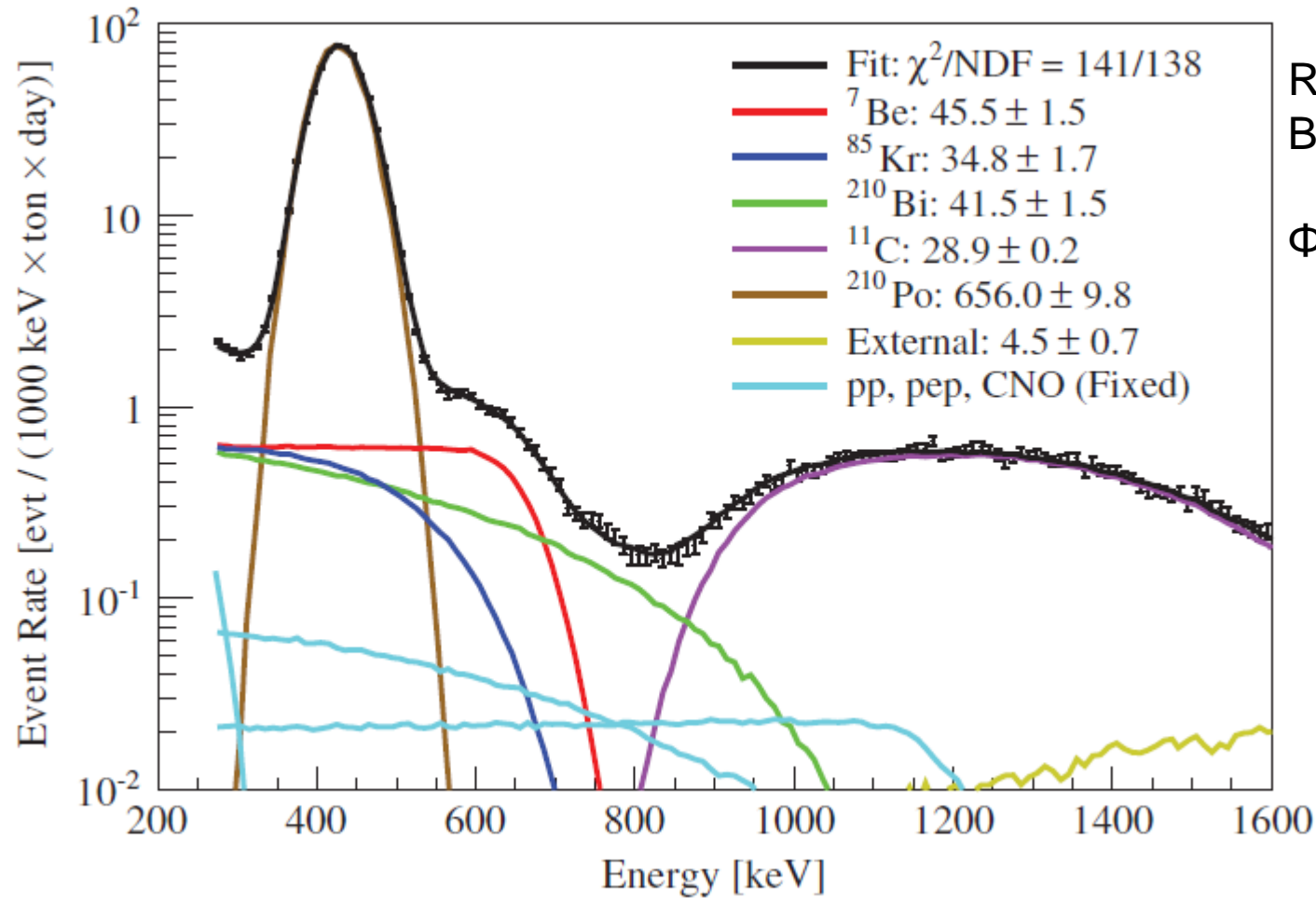


SSM pp neutrino flux $\sim 6 \cdot 10^{10} \text{ cm}^{-2} \text{ s}^{-1}$ (>90% of all solar neutrinos)

${}^7\text{Be}$ $5 \cdot 10^9 \text{ cm}^{-2} \text{ s}^{-1}$ ($\sim 10\%$)

Other s– much less

^7Be neutrino study : above ^{14}C end-point

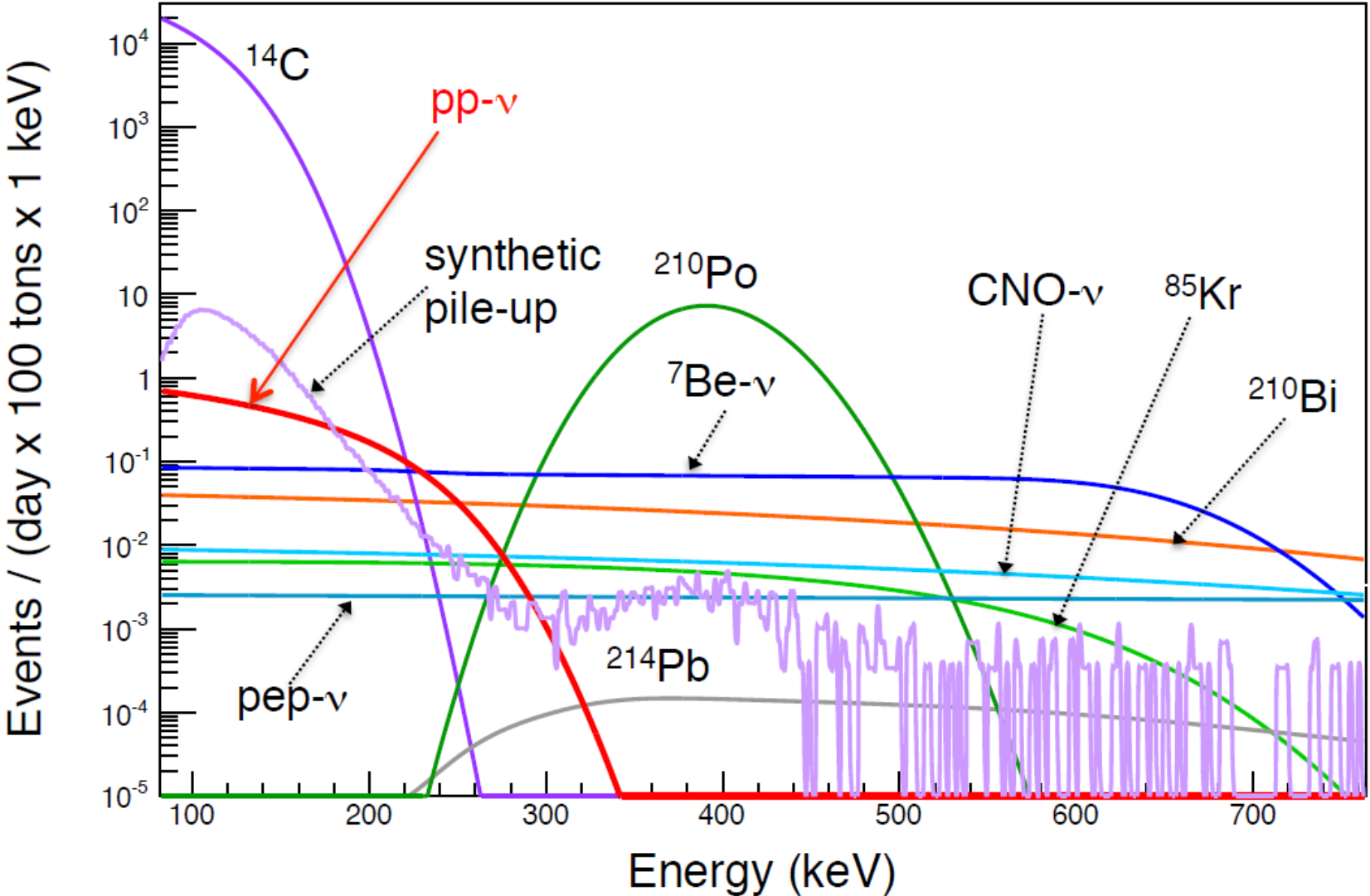


Radiochemical+
Borexino ^7Be :

$$\Phi(\text{pp}) = (6.14 \pm 0.61) 10^{10} \text{ cm}^{-2} \text{ s}^{-1}$$

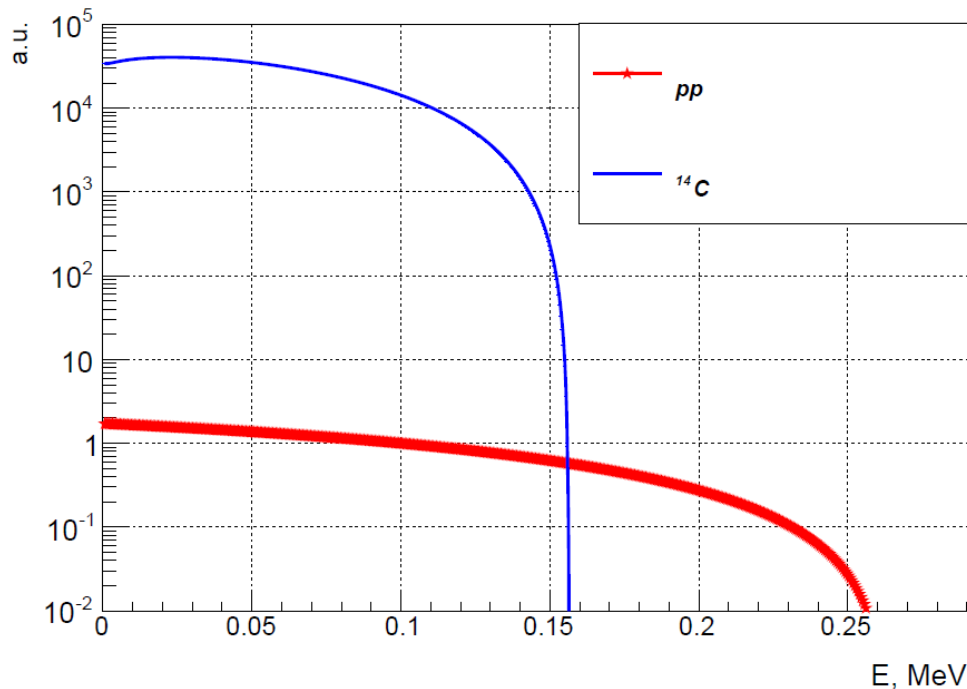
“Precision Measurement of the ^7Be Solar Neutrino Interaction Rate in Borexino”
PRL 107, 141302 (2011)

Expected Borexino Spectrum in energy scale



Why spectral measurement of pp-neutrino is possible?

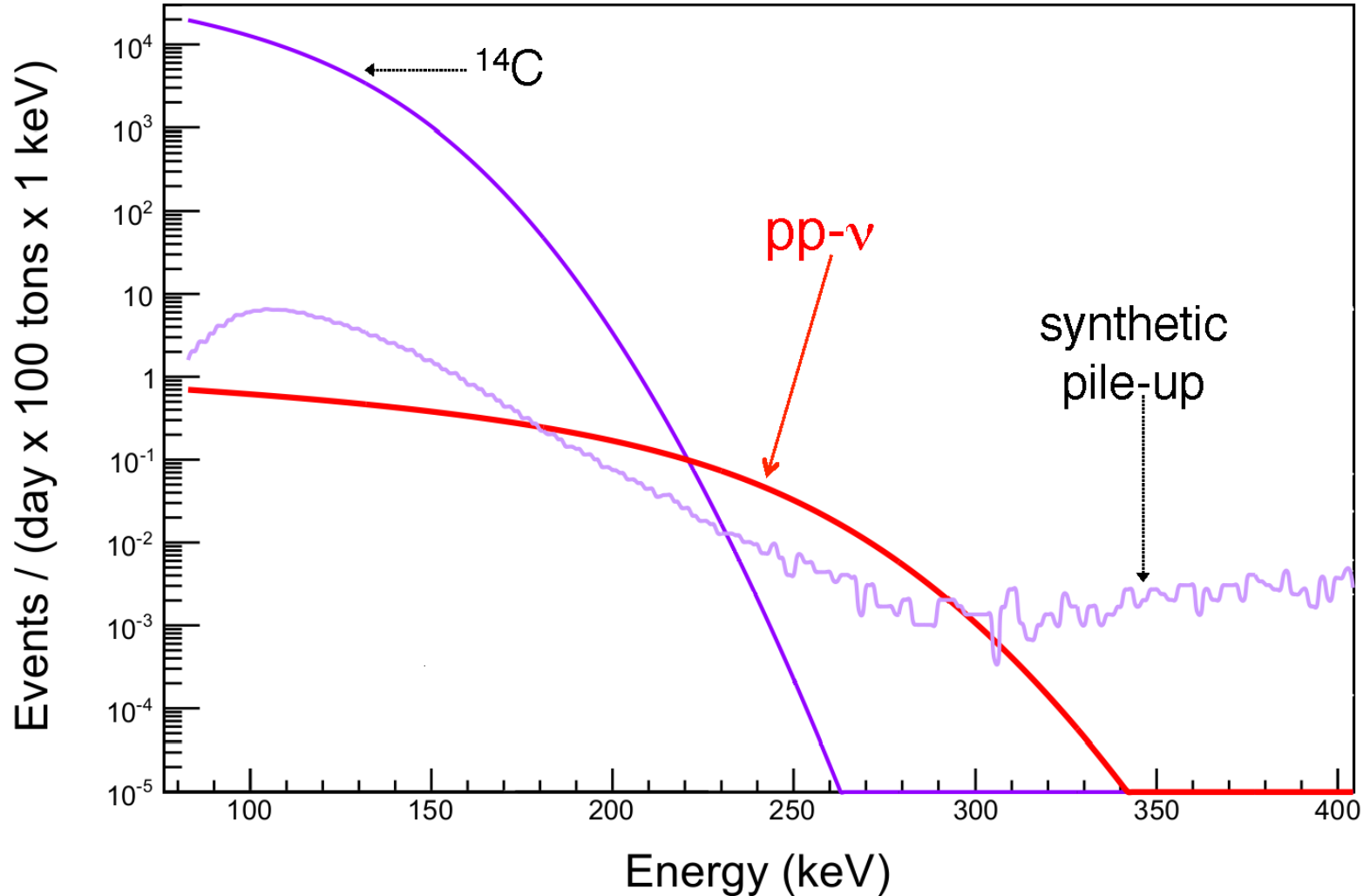
- Electron recoils spectrum for pp-neutrino has 264 keV end-point
- Main background for the measurement in liquid organic scintillator is ^{14}C , 156 keV end-point
- Spectra of other identified backgrounds are almost flat in the energy region of interest



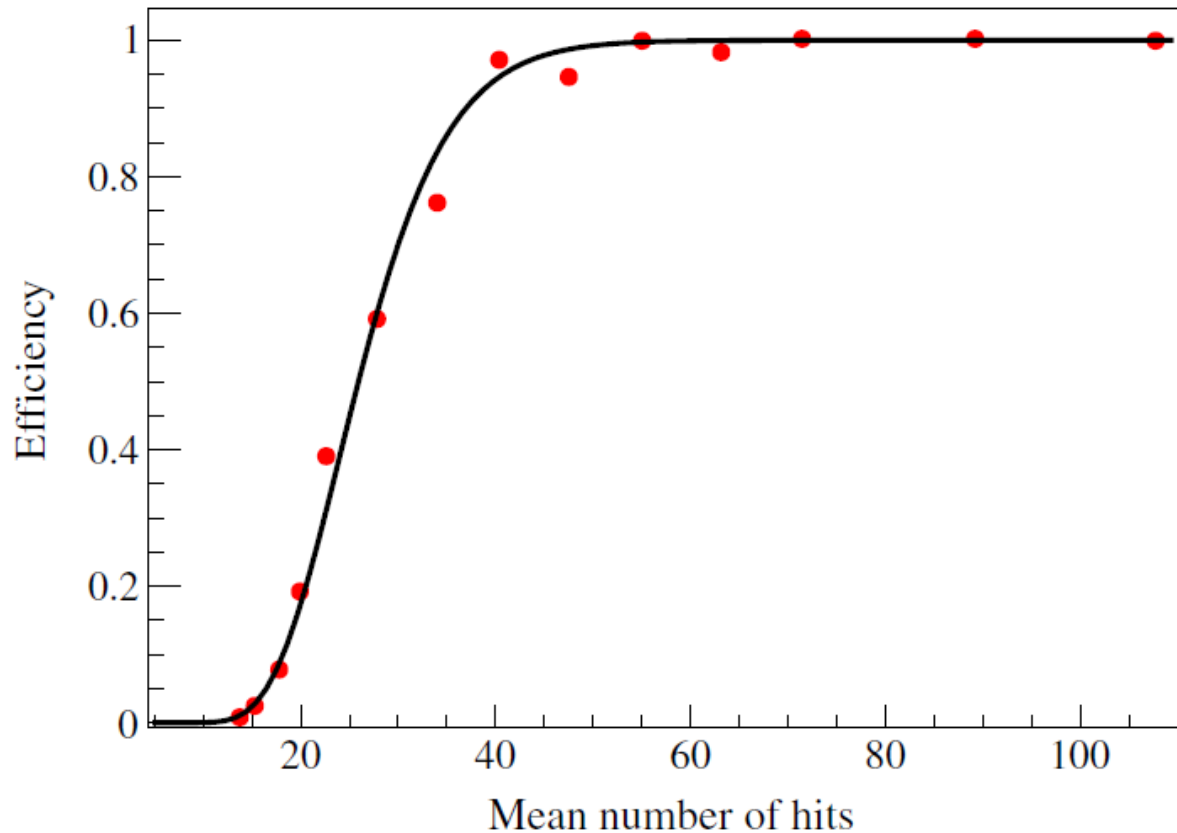
The measurement is feasible with:

- 1) low energy threshold
- 2) good energy resolution (~10% at 200 keV)
- 3) low radioactive backgrounds.
- 4) low ^{14}C count in view of the tail and pile-up

Effect of energy resolution and pile-up



The threshold on e-recoil



60 hits (of 2000 PMTs) \sim 120 keV \rightarrow

we can go to much lower threshold

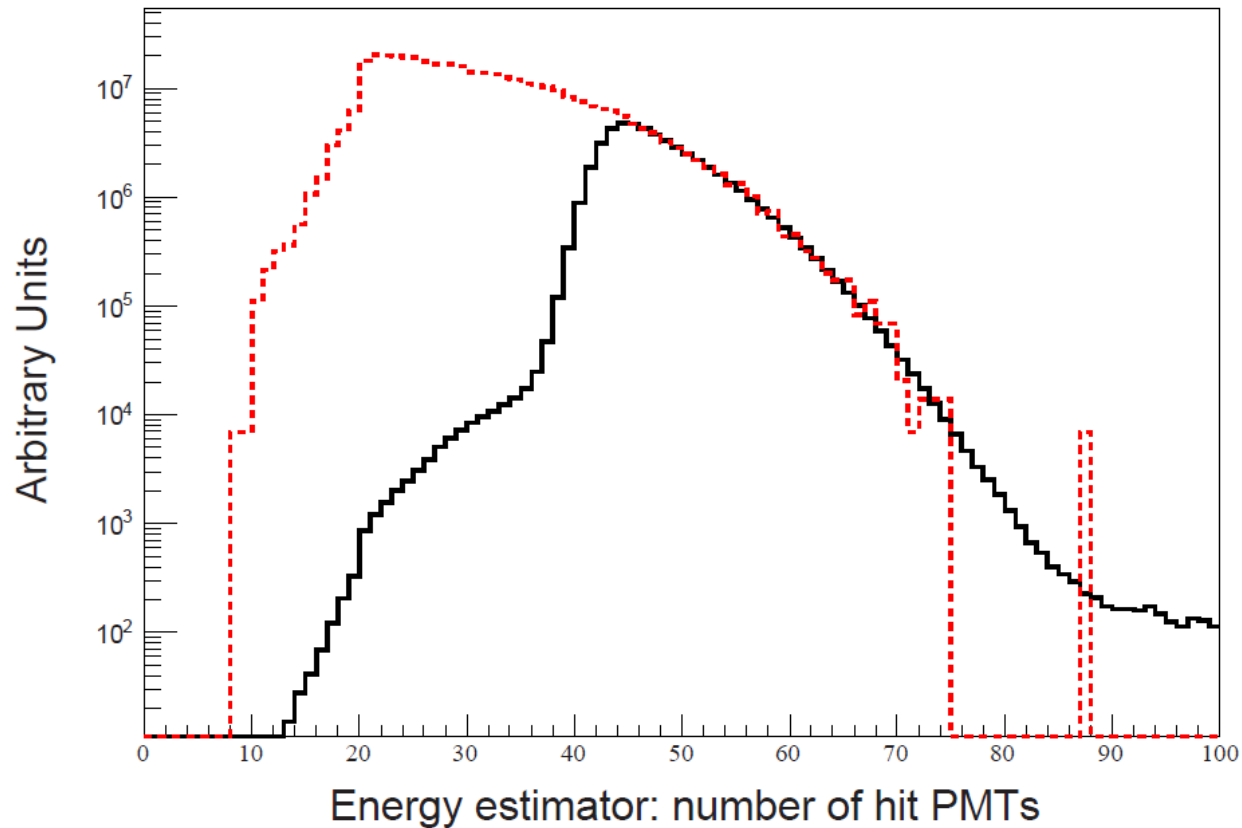
Data selection

- The data used for this analysis were acquired from January 2012 to May 2013 (408 days of data; Borexino Phase II).
- Data cuts – the same of ${}^7\text{Be}$ analysis excluding energy-dependent cuts (i.e. soft α/β and some others) to preserve shapes
- no coincidence with muon events : a 300 ms veto is applied following muons crossing the scintillator and buffer volumes, and a 2 ms veto following muons crossing only the water tank
- FV cut – the same of ${}^7\text{Be}$ ($R < 3.021$ && $|Z| < 1.67$ corresponding to 75.5 tonnes of LS). Independence on energy checked.

pp-neutrino signal extraction

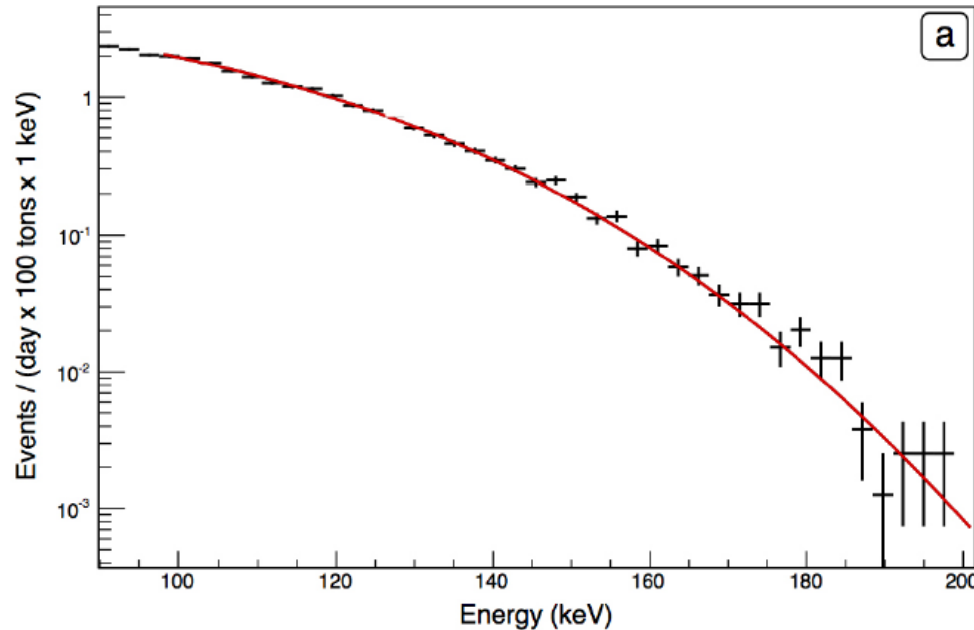
- **The pp neutrino rate is extracted by fitting experimental spectrum with the expected spectra of signal and backgrounds.**
- **We need to study carefully:**
 - **Spectral shapes involved in the fit**
 - **Rate of ^{14}C and pile-up events**
- **Critical issue concerning spectral shape is the analytical description of the energy scale, the analytical description of the energy resolution and the shape of the detector response to the monoenergetic event (the scintillation line shape).**

Independent ^{14}C rate estimation using second cluster events



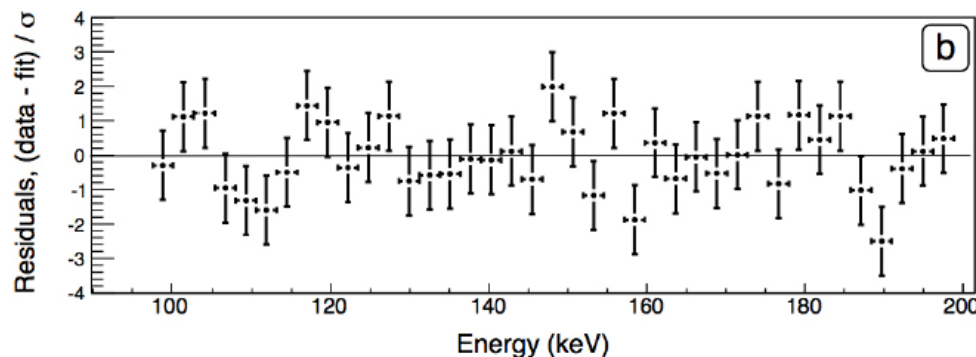
The ^{14}C rate was determined independently from the main analysis, by looking at a sample of data in which the event causing the trigger is followed by a second event (in red) within the acquisition time window of 16 ms. The main events spectrum are shown in the same scale in black.

^{14}C rate estimation using second cluster events



Time acquisition window is 16 ms. Only the data with $\Delta T > 8$ ms were used (to exclude PMTs afterpulses)
Fitting the data with theoretical β -spectrum shape we obtained:
 40 ± 1 Bq/100 t rate.

The corresponding $^{14}\text{C}/^{12}\text{C}$ isotopical abundance is
 $(2.7 \pm 0.1) \cdot 10^{-18}$ g/g

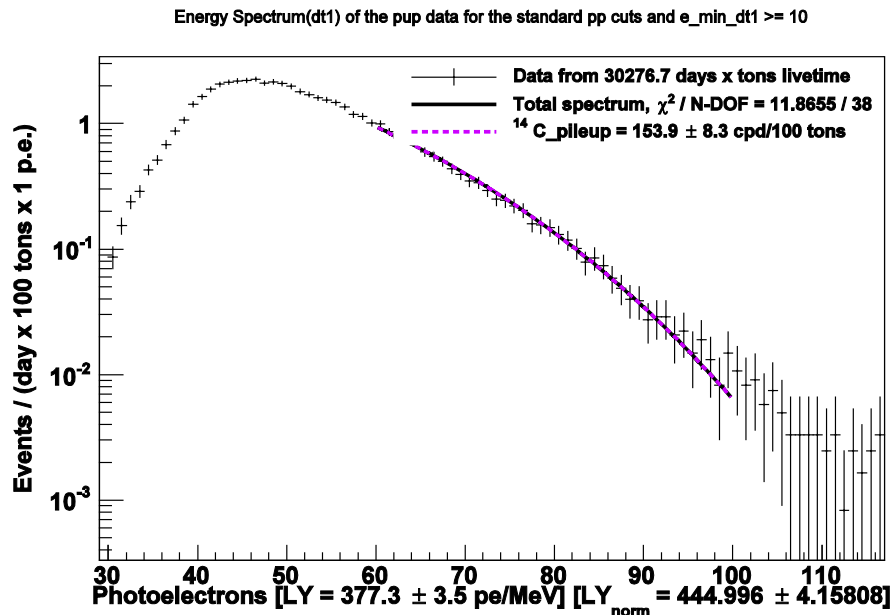


^{14}C β -spectrum has some deviations from the allowed shape parameterized as $S(E) = 1 + \alpha E$

Methods to construct pile-up

- Synthetic pile-up Data+Data

- Real triggered events without any selection cuts are artificially overlapped with random data samples and processed with reconstruction code in a standard way



Fit of synth.pile-up with analytical function:
 $154 \pm 10 \text{ cpd}/100 \text{ tones}$

Compatible with naïve rate estimation using ^{14}C counting rate only (40 Bq/100 t).

Energy scale

Energy estimator: number of triggered pmts (Npmts)

Conversion of energy to Npmts depends on:

- Light yield (LY) – **free**
- Birks' parameter (k_B)- **fixed** at the value found from the calibration data ($k_B=0.0109$ cm/MeV)
- Average number of live PMTs (N_{Live}) – calculated for data set
- Geometric correction parameter (gc) **fixed** at value found with MC (no sensitivity in the low energy part)

$$npmts = N_{Live} \cdot \frac{1 - e^{-\frac{Q(E)}{N_{Live}}}}{1 + gc \cdot \frac{Q(E)}{N_{Live}}}$$

$$Q(E) = LY \cdot E \cdot f(k_B, E)$$

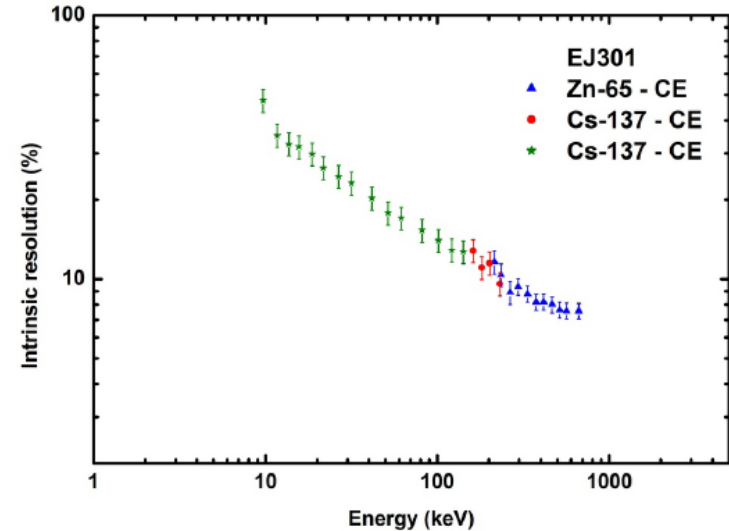
E → Q → Npmts
conversion

The same as in
⁷Be analysis

Energy resolution

energy resolution for npmts variable depends on:
(2 free):

- v_T - **free**. Takes into account spatial non-uniformity of the light collection (basically ^{210}Po parameter because it has negligible contribution in the low energy part of the spectrum). The same for Po and ^{14}C
- σ_{int} - **free**. Intrinsic line width (extra width compared to $\sqrt{N_{\text{photons}}}$) for β -particles (absent or negligible for α).
- v_f - **fixed** at calculated value (variance of the number of live Pmts).
- $v_1 = 0.17$ - **fixed** at value found with MC (no sensitivity in the low energy part).



Intrinsic resolution measured for an EJ301 liquid scintillator. Various colours represent the data obtained with different radiation sources placed at various angles. (figure from 2012 JINST 7 P06011)

$$\sigma_N^2 = npmts \cdot \left(1 - \frac{npmts}{N_{\text{Live}}}(1 + v_1)\right) + npmts^2 \cdot (v_T \cdot npmts \cdot (1 + v_f) + v_f) + \sigma_{\text{int}}^2$$

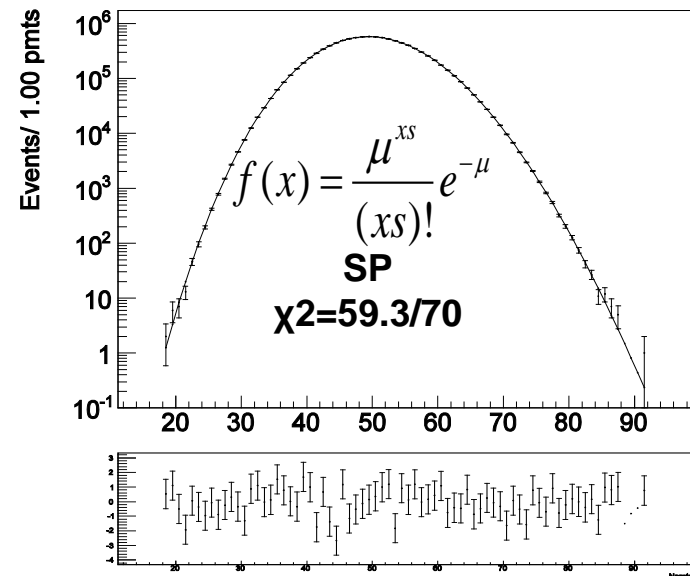
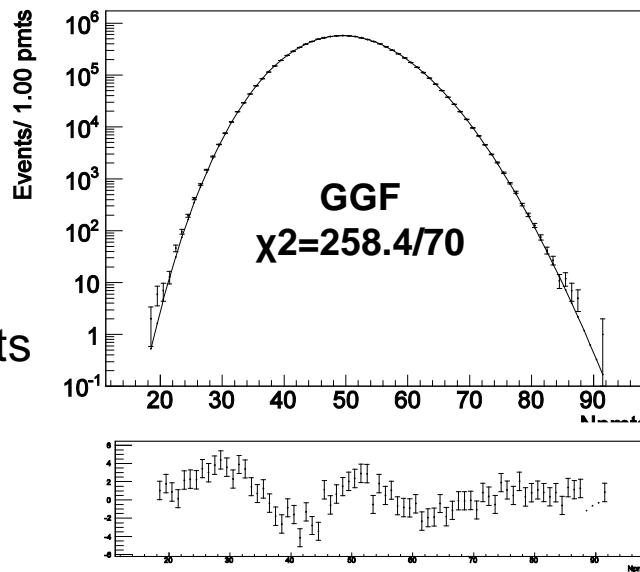
The same as in ^7Be analysis + Intrinsic line width

Analytical description of Scintillation Line shape

Line shape – the shape of the detector's response for uniformly distributed monoenergetic particle.

Generalized Gamma Function (${}^7\text{Be}$ analysis) substituted with Scaled Poisson

Reason: very high statistics of ${}^{14}\text{C}$ (compared to ${}^{210}\text{Po}$ in ${}^7\text{Be}$ analysis)



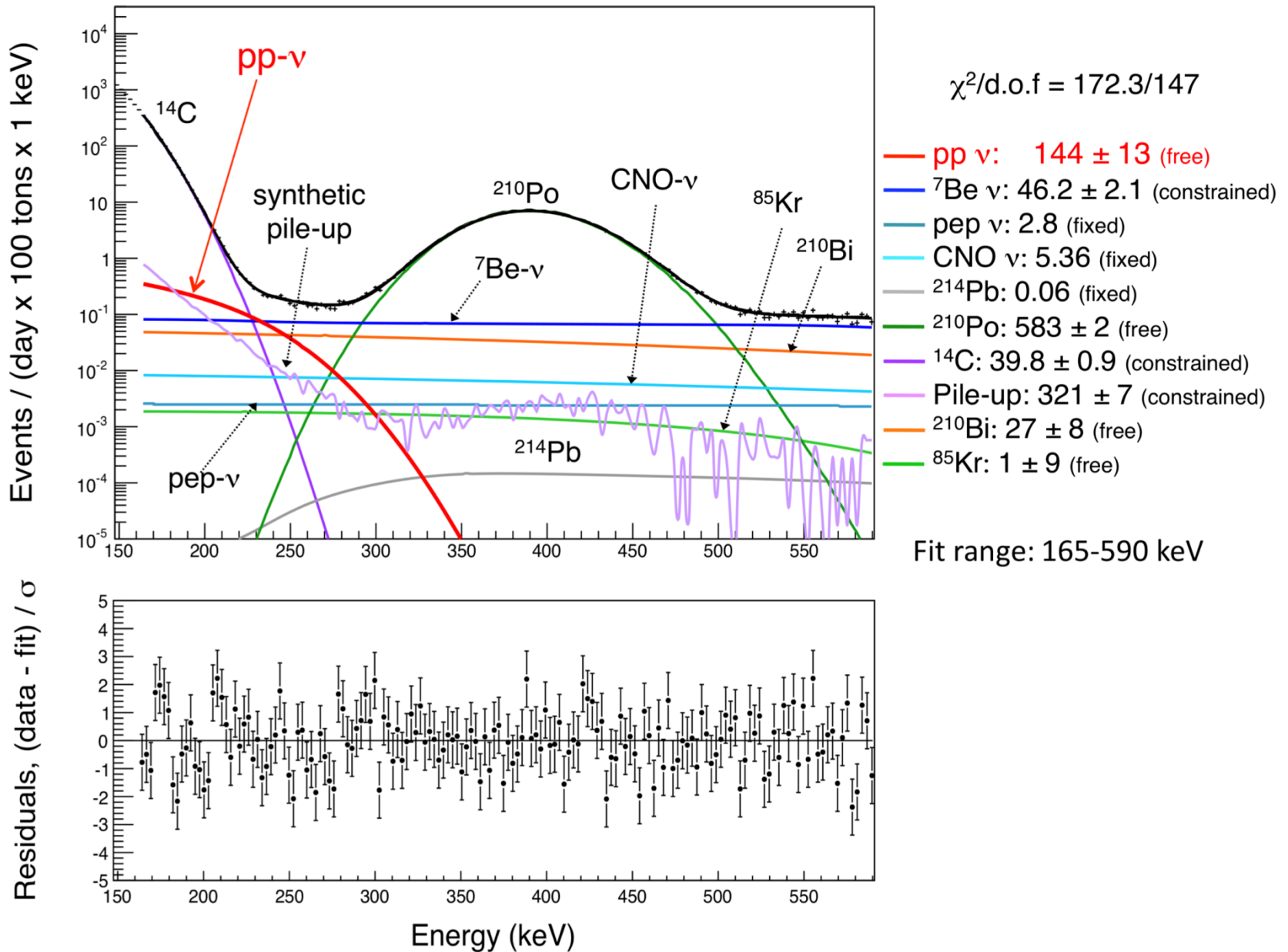
10⁷ MC events
in FV
<N>=60

The mean values and variances of both shapes are the same

“Standard” fit

- Variable: npmts
- Range: [60-220 Npmts] (corresponds to 165-590 keV)
- Response function shape: Scaled Poisson
- **Free** spectral components: ^{14}C , ^{210}Bi , ^{210}Po , pp, ^{85}Kr
- **Constrained** spectral components: ^7Be (paper central value)
- **Fixed** spectral components: pep+CNO+ ^8B (SSM/LMA(HM))
- Energy scale variables: LY **free**, k_B **fixed**, f_{eq} **calculated**
- Energy resolution vars: σ_{int} and v_T **free**; v_f **calculated**
- Synthetic pile-up **constrained**
- Free: LY+2 resolution vars+Po position+5 spectral components + constrained pup and ^7Be
- **9 free parameters and 2 (pileup and ^7Be) constrained**

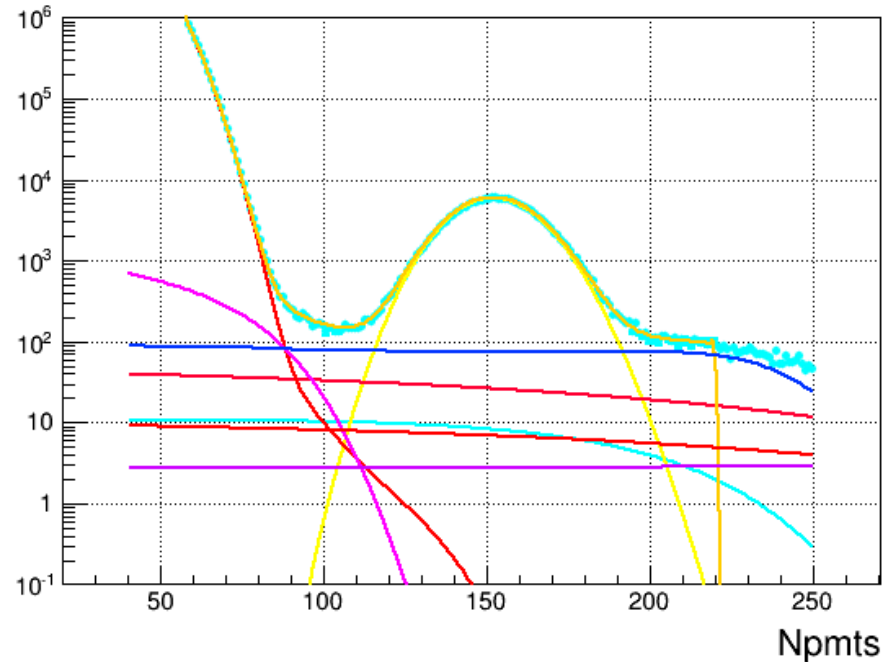
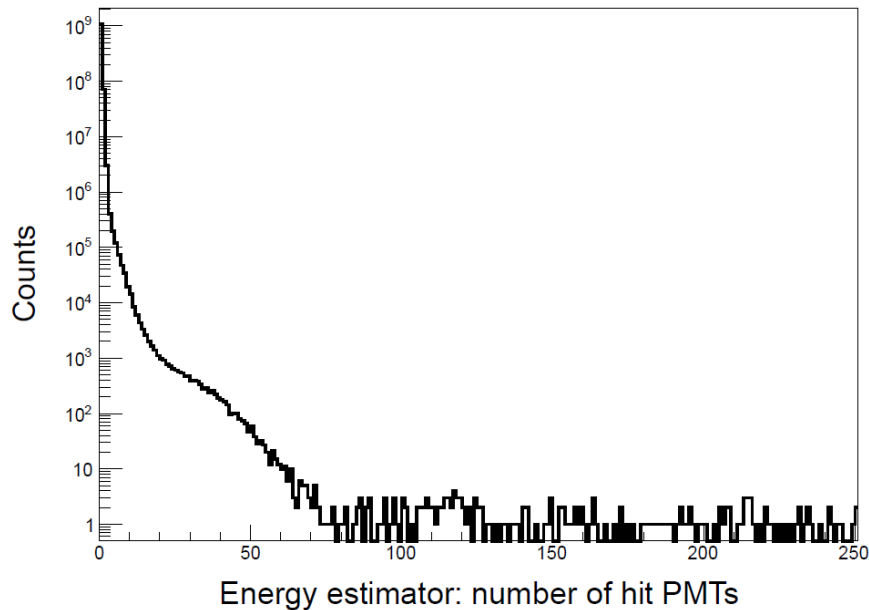
The results of the standard spectral fit



Cross-check of pile-up

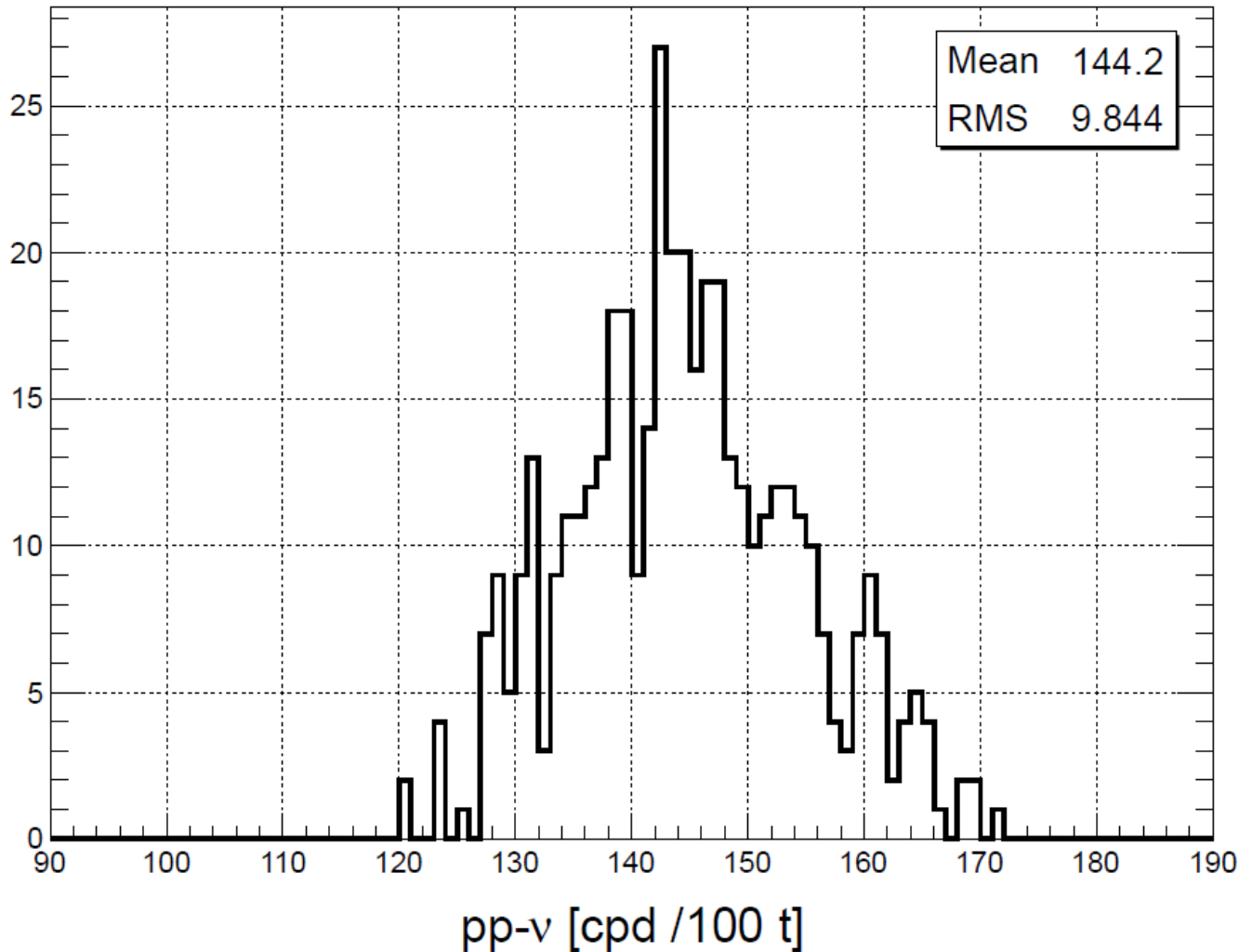
Convolution with randomly sampled data (trivial once randomly sampled data are provided)

Used for cross checks (“convolution” method), the effect can be seen in the tail of ^{14}C in the plot.



Systematics estimation

Fit Occurrence



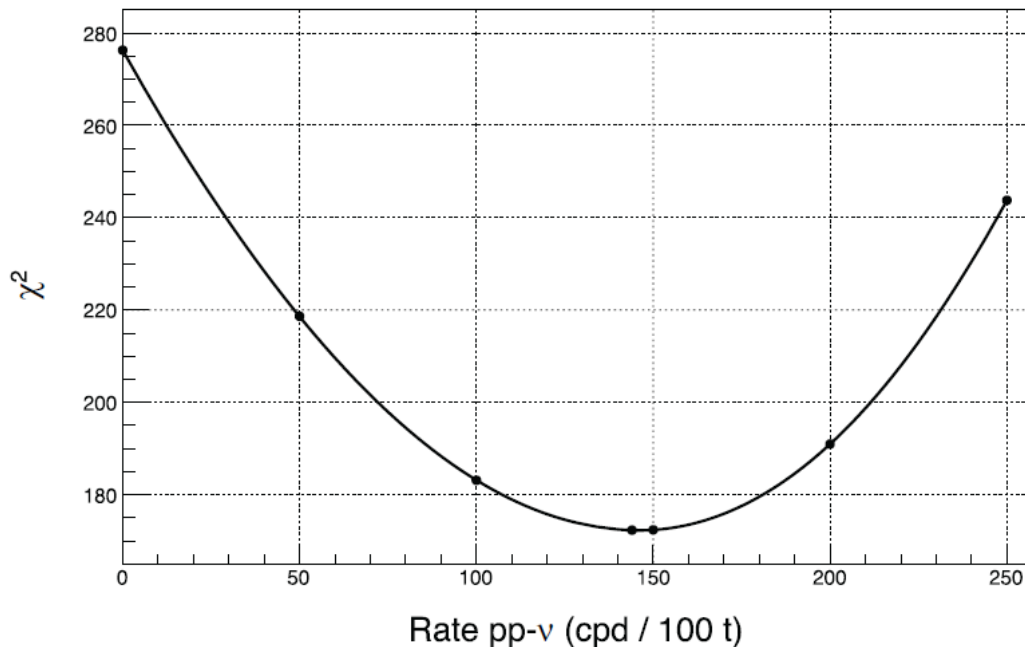
Values are obtained by varying the fit conditions, including the fit energy range, synthetic-versus-analytic pile-up spectral shape, and energy estimator. The distribution shown is peaked around reported value of 144 cpd/100 t.

First real-time measurement of pp-neutrino flux ($\sim 11\%$ precision)

$$pp = 144 \pm 13 \text{ (stat)} \pm 10 \text{ (syst) cpd/100 t}$$

compared to expected (MSW/LMA, HM)

$$131 \pm 2 \text{ cpd/100 t}$$



pp neutrino flux:

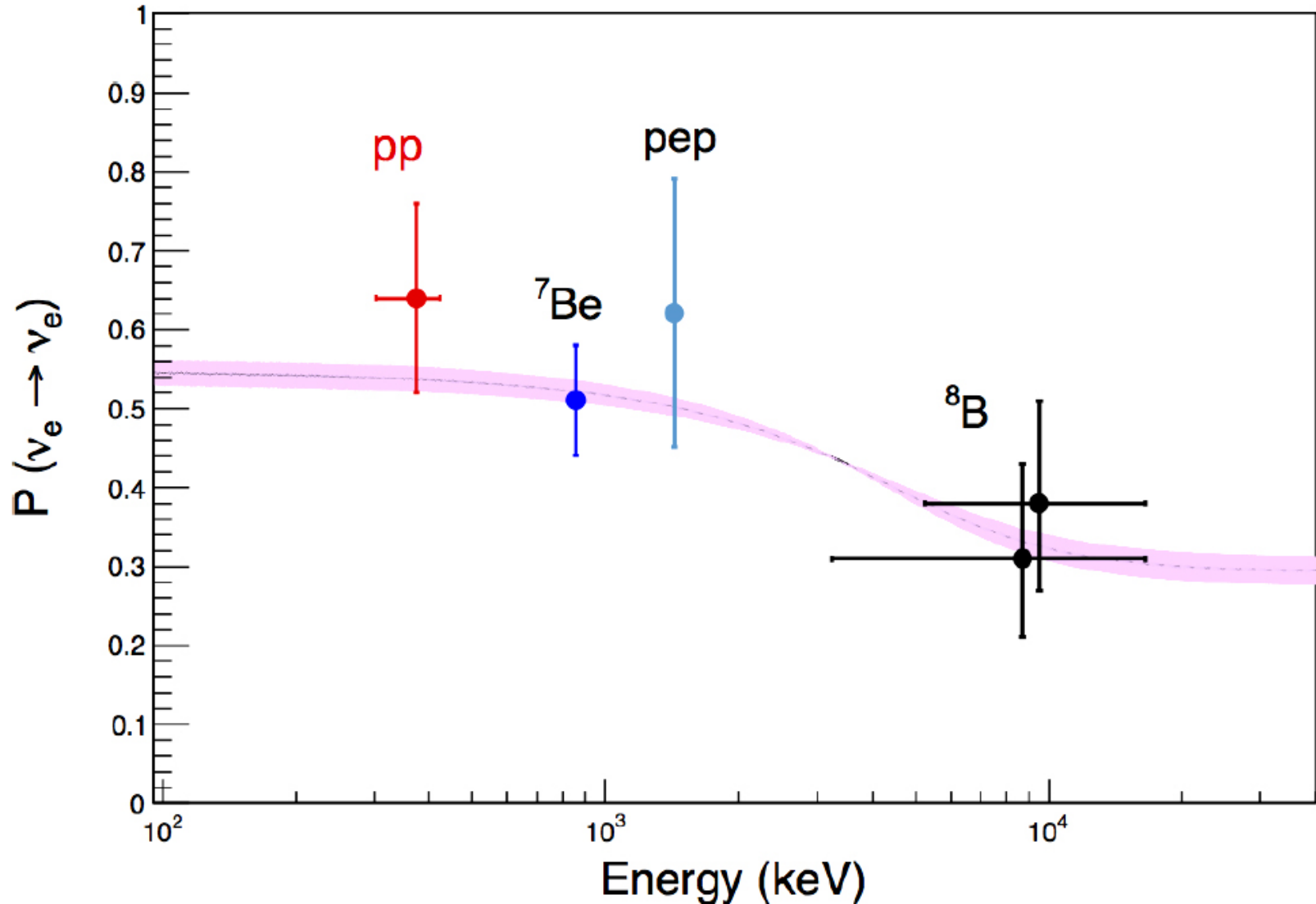
$$(6.6 \pm 0.7) \cdot 10^{10} \text{ cm}^{-2} \text{ s}^{-1}$$

VS

$$(5.98 \pm 0.04) \cdot 10^{10} \text{ cm}^{-2} \text{ s}^{-1}$$

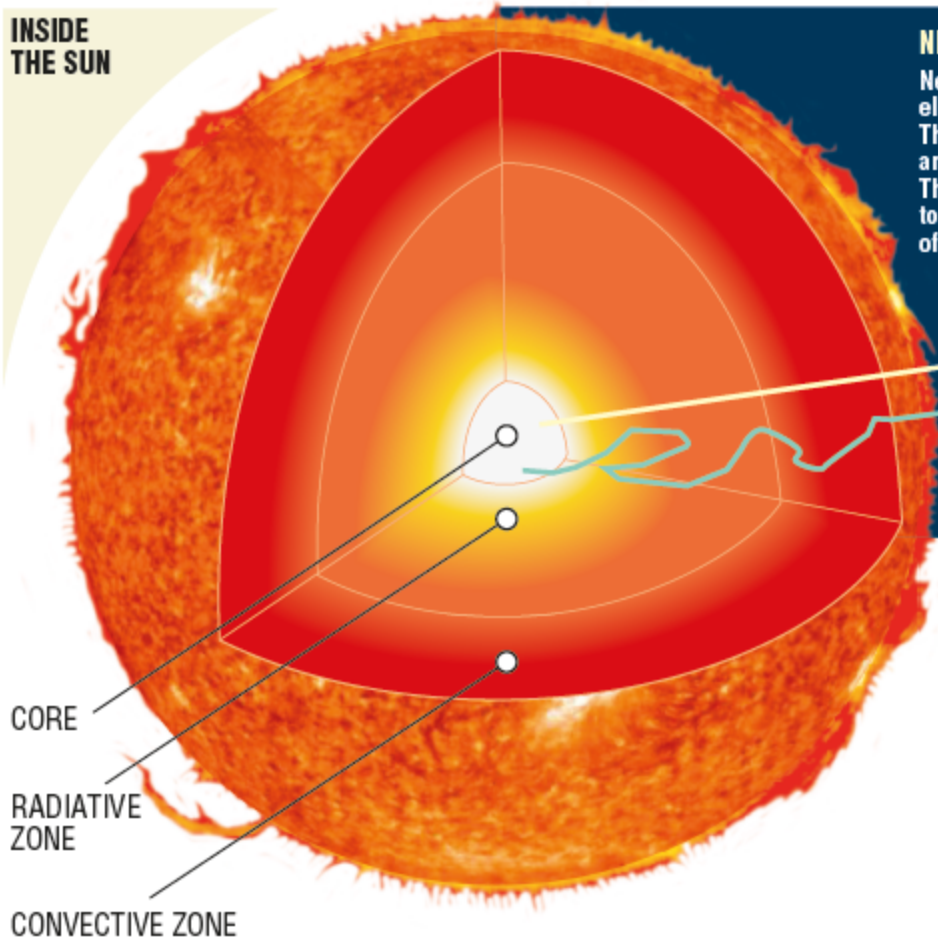
Zero pp count is excluded at 10σ level

Borexino measured electron neutrino survival probability for 4 different nuclear reactions



THE SUN AS BOREXINO SEES IT IN REAL TIME

INSIDE THE SUN



NEUTRINOS

Neutrinos are particles with no electric charge and a tiny mass. They rarely interact with matter and may cross it undisturbed. That's why they take 8 minutes to get there from the core of the Sun to the Earth.

PHOTONS

The radiation studied so far is made up of photons, which interact with solar matter. It takes about 100,000 years for it to reach the Sun's surface and reach Earth.

8 minutes

100,000 years



Gran Sasso mountain

Laboratory

By analyzing P-P neutrino emission, Borexino has shown that the energy produced today in the Sun's core is the same as that produced 100,000 years ago.

CORE

RADIATIVE ZONE

CONVECTIVE ZONE

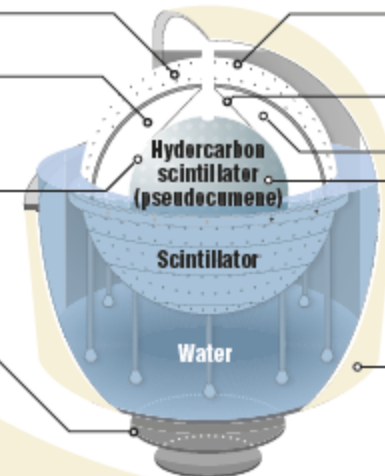
THE BOREXINO DETECTOR: HOW IT WORKS

Stainless steel sphere
13,7 m diameter

Thin nylon film
(radon gas barrier)

Nylon sphere
8,5 m diameter

Shielding
steel dishes



Multiplication
200 photomultiplier tubes (facets)

Vessel retention

2.200 photomultiplier tubes (facets)

organic liquid

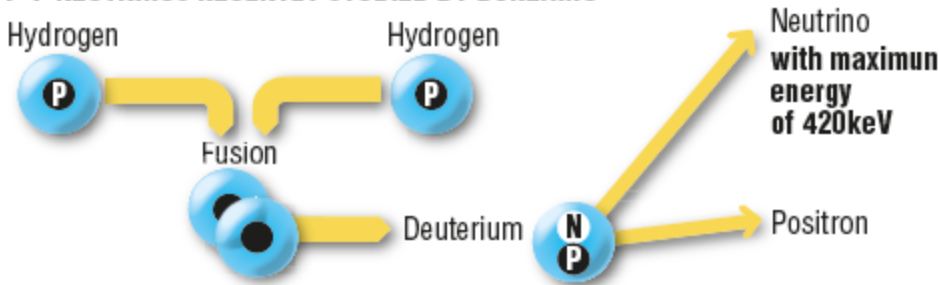
Stainless steel

Water

10



THE THERMONUCLEAR FUSION REACTION THAT PRODUCES THE P-P NEUTRINOS RECENTLY STUDIED BY BOREXINO



Borexino displays a russian doll structure. Surrounded by 2,400 tons of highly purified water, a stainless steel sphere contains 1,000 tons of a liquid hydrocarbon (pseudocumene). At its center, within a smaller nylon sphere, are 300 tons of scintillating liquid.

Within this innermost sphere neutrinos interact with the liquid scintillator producing small flashes of light.

The photomultiplier tubes, acting as ultra-sensitive artificial eyes, detect and record the light flashes produced by the neutrinos.

Borexino observes dozens of these signals every day.



Milano



Perugia



Genova



Napoli



TU Dresden



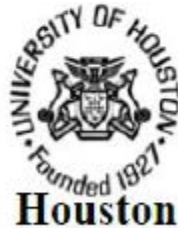
Jagiellonian
Kraków



the Borexino Collaboration



JINR
Dubna



Houston



Paris



MOSCOW



Los Angeles



Princeton



UMass
Amherst



St. Petersburg



Kurchatov
Moscow



Spares

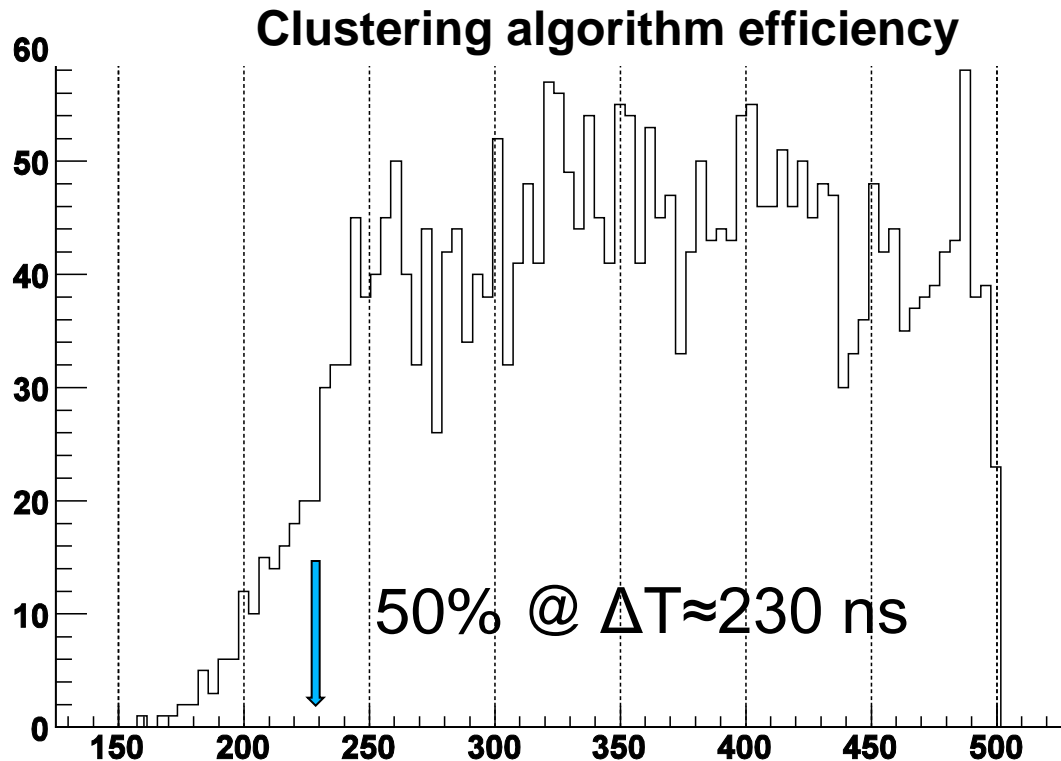
List of proposals on pp-neutrino search

Project (reference)	Method	Threshold, keV	Resolution	Mass, t	Reaction	pp events/d
LENS [3]	^{176}Yb , LS	301(ν)	7% (1 MeV)	20 (8% nat ^{176}Yb)	$^{176}\text{Yb} + \nu_e \rightarrow$ $^{176}\text{Lu} + e^-$	0.5
INDIUM [4]	^{115}In , LS	118(ν)	5–10% (1 MeV)	4	$^{115}\text{In} + \nu_e \rightarrow$ $^{115}\text{Sn}^*(613) + e^-$	1.0
GENIUS [5]	^{76}Ge , Scattering	11(e^-) 59(ν)	0.3% (300 keV)	1 10	$\nu + e^- \rightarrow \nu + e^-$	1.8 18
HERON [6]	Superfluid ^4He , rotons/phonons + UV	50(e^-) 141(ν)	8.3% (364 keV)	10	$\nu + e^- \rightarrow \nu + e^-$	14
XMASS [7]	Xe, LS	50(e^-) 141(ν)	17.5% (100 keV)	10	$\nu + e^- \rightarrow \nu + e^-$	14
HELLAZ [8]	He (5 atm), TPC	100(e^-) 217(ν)	6% (800 keV)	2000 m^3	$\nu + e^- \rightarrow \nu + e^-$	7
MOON [9]	Drift chambers	168(ν)	12.4% FWHH (1 MeV)	3.3	$\nu_e + ^{100}\text{Mo} \rightarrow$ $^{100}\text{Tc} + e^-$	1.1
MUNU [10]	TPC, CF_4 , direction	100(e^-) 217(ν)	16% FWHH (1 MeV)	0.74 (200 m^3)	$\nu + e^- \rightarrow \nu + e^-$	0.5
NEON [11]	He, Ne, scintillator	20(e^-) 82(ν)	16% FWHH (100 keV)	10	$\nu + e^- \rightarrow \nu + e^-$	18
Present work	LS	170(e^-) 310(ν)	10.5% (200 keV)	10	$\nu + e^- \rightarrow \nu + e^-$	1.8

Energy variables

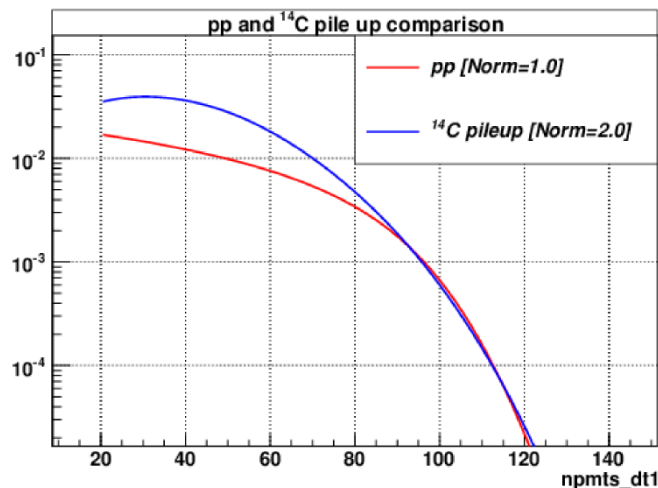
Two of them: **npmts_dt1** and **npmts_d2** (**non-normalized** to 2000 PMTs in contrast to what have been done in ^7Be analysis)

Both are based on the number of triggered PMTs in the fixed-length energy window (230 ns for dt1 and 400 ns for dt2), easier to model the pile-up. Pile-up is scaled as dT while other components remains unchanged - possibility for pile-up cross-checks.



Naïve ^{14}C pile-up rate estimation

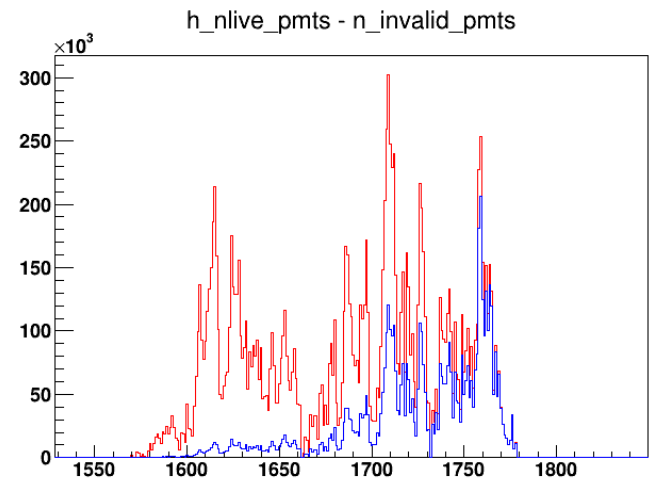
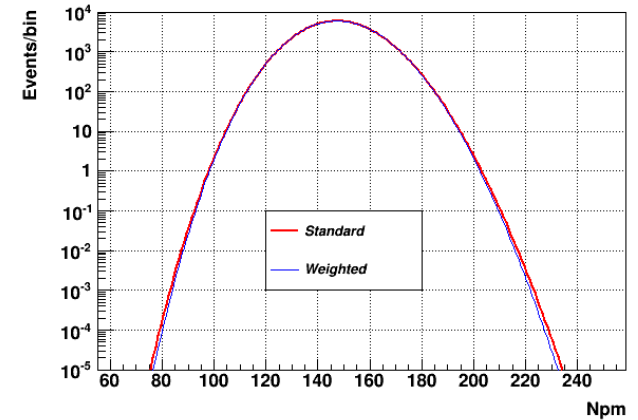
- Estimated count rate of ^{14}C in IV is 108 Bq for 270 t
- In one day we got $(108)^2 \cdot 8.64 \cdot 10^4 \cdot 4 \cdot 10^{-7} = 403$ cpd of ^{14}C pile-up in $\Delta T = 400$ ns gates in IV or 150 cpd in 100 tones;
- For $\Delta T = 230$ ns the pile-up rate reduces to 232 ± 12 cpd in IV or:
 86 ± 5 cpd/100 tones
- **We need the shape and the counting rate/**both the shape and the counting rate could be influenced by the reconstruction algorithm




Comparison of pp and analytical pileup shape


Other factors influencing the scintillation line shape

- **Varying number of live PMTs** : could be taken into account precisely as we know NLivePmt for every event. The response function can be weighted over NLivePmts . The improvement of χ^2 is of 2-3 with all the best parameters of the fit unchanged, and with a drawback of heavy calculations.
- **Additional spread because of the varying light collection efficiency over the detector** : small compared to the “statistical” line width, the “basic” shape remains unchanged
- The decay of Po and decrease of NLivePmt in time lead to the higher “effective” LY for Po and lower “effective” variance of the Po line : could be calculated analytically ($\text{Fill}(1,1) \rightarrow \text{Fill}(1, e^{-(t-t_0)/T})$)



Scaled Poisson vs GGF

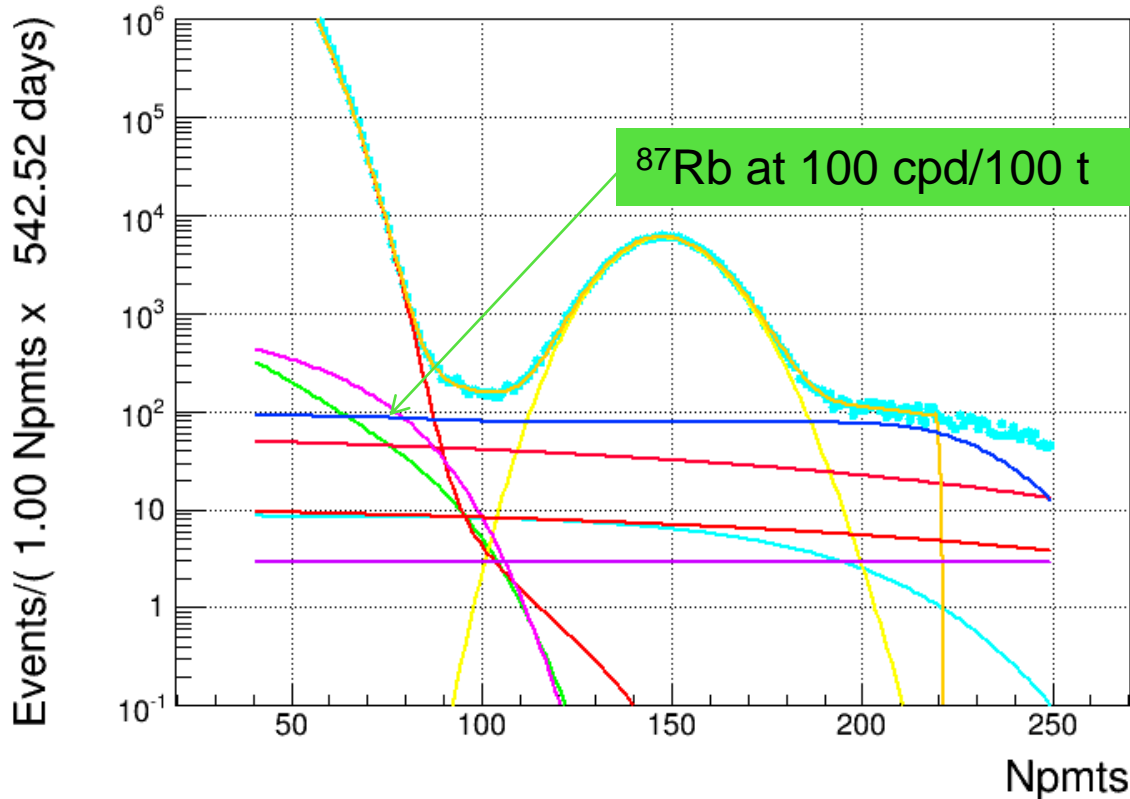

$$f(x) = \frac{\mu^{xs}}{(xs)!} e^{-\mu}$$


$$f(x) = 2\beta^\alpha \frac{x^{2\alpha-1}}{\Gamma(\alpha)} e^{-\beta x^2}$$

	1σ	2σ	3σ	4σ
Gauss	$1.5 \cdot 10^{-1}$	$2.10 \cdot 10^{-2}$	$1.2 \cdot 10^{-3}$	$2.5 \cdot 10^{-5}$
GGF	$1.5 \cdot 10^{-1}$	$2.45 \cdot 10^{-2}$	$1.9 \cdot 10^{-3}$	$7.8 \cdot 10^{-5}$
SP	$1.5 \cdot 10^{-1}$	$2.45 \cdot 10^{-2}$	$2.0 \cdot 10^{-3}$	$8.4 \cdot 10^{-5}$

Fraction of events in tail for 3 base functions ($\langle N \rangle = 49.9$, source at center)

^{87}Rb – could it change the result?



Second forbidden β -decay
with $E_0=283.3$ keV end-point

$$\Delta\chi^2=+1$$

@

100 cpd of $^{87}\text{Rb}/100\text{t}$

with corresponding pp
count decrease of 40
cpd/100 t

Correlation: 1 cpd/100t of ^{87}Rb simulates 0.4 cpd/100t of pp

^{87}Rb vs ^{40}K

- though the Rb is quite abundant in nature (the twenty-third most abundant element in the Earth's crust) the typical abundance of Rb is factor $(2-4) \cdot 10^3$ lower than that of K (which is one of most abundant and contributes 1.5% of the crust weight).
- Taking into account the abundance of radioactive isotopes: 0.278 of ^{87}Rb against $1.17 \cdot 10^{-4}$ of ^{40}K , and the ratio of live-times 47.2 vs 1.28 billions yr, we get the ratio for the typical activity of ^{87}Rb :

$$\frac{\text{Activity}(\text{Rb})}{\text{Activity}(\text{K})} = \frac{\text{abund}({}^{87}\text{Rb})}{\text{abund}({}^{40}\text{K})} \cdot \frac{T_{1/2}^{40\text{K}}}{T_{1/2}^{87\text{Rb}}} \cdot \frac{\text{At}(\text{Rb})}{\text{At}(\text{K})} = 64.5 \cdot \frac{\text{At}(\text{Rb})}{\text{At}(\text{K})}$$

Location	K (ppb by Atoms)	Rb (ppb by Atoms)	R(Rb/K)	Activity ratio
Universe	100	0.1	10 ⁻³	0.065
Sun	100	0.4	4·10 ⁻³	0.26
Meteorite (carbonaceous)	370000	770	2.1·10 ⁻³	0.14
Crustal rocks	7800000	14000	1.8·10 ⁻³	0.12
Sea water	65800	8.7	1.3·10 ⁻⁴	0.08
Human	32000	340	1.1·10 ⁻²	0.7

- The chemistry of the elements is the same, Rb substitutes K, so we can expect the same purification factor if started from the natural ratio of Rb/K.
 - **^{40}K limit from the pep-analysis is <0.11 cpd/100 t (68% C.L.)*** **$<4.6 \cdot 10^{-16}$ Knat [g/g Sc]**
- *published value is <0.4 cpd/100 t corresponding to 95% C.L. ($<1.7 \cdot 10^{-15}$ Knat [g/g Sc])